

Structure preserving methods via Global Flux quadrature: divergence-free preservation with continuous Finite Element

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SUPG

- Streamline **upwind** Petrov-Galerkin
- Stabilization for Continuous Galerkin Finite Element methods
- Advection dominated problems

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SUPG for a **system** of linear hyperbolic conservation laws

$$\partial_t q + J^x \partial_x q + J^y \partial_y q = 0, \quad q : \Omega_h \times \mathbb{R}^+ \rightarrow \mathbb{R}^S.$$

Take $V_h^K := \{\varphi \in \mathcal{C}(\Omega_h) : \varphi|_E \in \mathbb{P}^K(E) \forall E \in \Omega_h\}$. SUPG is $\forall \varphi \in (V_{h,0}^K)^S$ find $q \in (V_h^K)^S$ such that

$$0 = \int (\varphi) (\partial_t q + J^x \partial_x q + J^y \partial_y q) dx$$

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$$0 = \int (\varphi + \alpha \Delta x \partial_x \varphi \operatorname{sign}(J^x) + \alpha \Delta y \partial_y \varphi \operatorname{sign}(J^y)) (\partial_t q + J^x \partial_x q + J^y \partial_y q) \, dx$$

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$$\begin{aligned} 0 &= \int (\varphi + \alpha \Delta x \partial_x \varphi \operatorname{sign}(J^x) + \alpha \Delta y \partial_y \varphi \operatorname{sign}(J^y)) (\partial_t q + J^x \partial_x q + J^y \partial_y q) dx \\ &= \int \varphi (\partial_t q + J^x \partial_x q + J^y \partial_y q) dx + \alpha \int (\Delta x \partial_x \varphi \operatorname{sign}(J^x) + \Delta y \partial_y \varphi \operatorname{sign}(J^y)) \partial_t q dx \\ &\quad + \alpha \int \Delta x \partial_x \varphi \operatorname{sign}(J^x) (J^x \partial_x q + J^y \partial_y q) dx + \alpha \int \Delta y \partial_y \varphi \operatorname{sign}(J^y) (J^x \partial_x q + J^y \partial_y q) dx. \end{aligned}$$

Acoustics equations and involutions

Acoustics equation

$$\begin{cases} \partial_t u + \partial_x p = 0, \\ \partial_t v + \partial_y p = 0, \\ \partial_t p + \partial_x u + \partial_y v = 0, \end{cases}$$

$$\begin{cases} \partial_t \underline{v} + \nabla p = 0, \\ \partial_t p + \nabla \cdot \underline{v} = 0, \end{cases}$$

$$\partial_t q + J^x \partial_x q + J^y \partial_y q = 0$$

$$q = \begin{pmatrix} u \\ v \\ p \end{pmatrix}, \quad J^x = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix}, \quad J^y = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}.$$

Involution

The system of linear acoustics possesses an involution:

$$\partial_t (\nabla \times \underline{v}) = \nabla \times \partial_t \underline{v} = -\nabla \times \nabla p = 0,$$

Equilibria

$$q : \partial_t q = 0, \quad \begin{cases} \nabla \cdot \underline{v} = 0 \\ p \equiv C \in \mathbb{R} \end{cases}$$

Other equations sharing div-free equilibria

SW, Euler, Maxwell, low Mach

Typical problems

4.1 Low Mach number limit

81

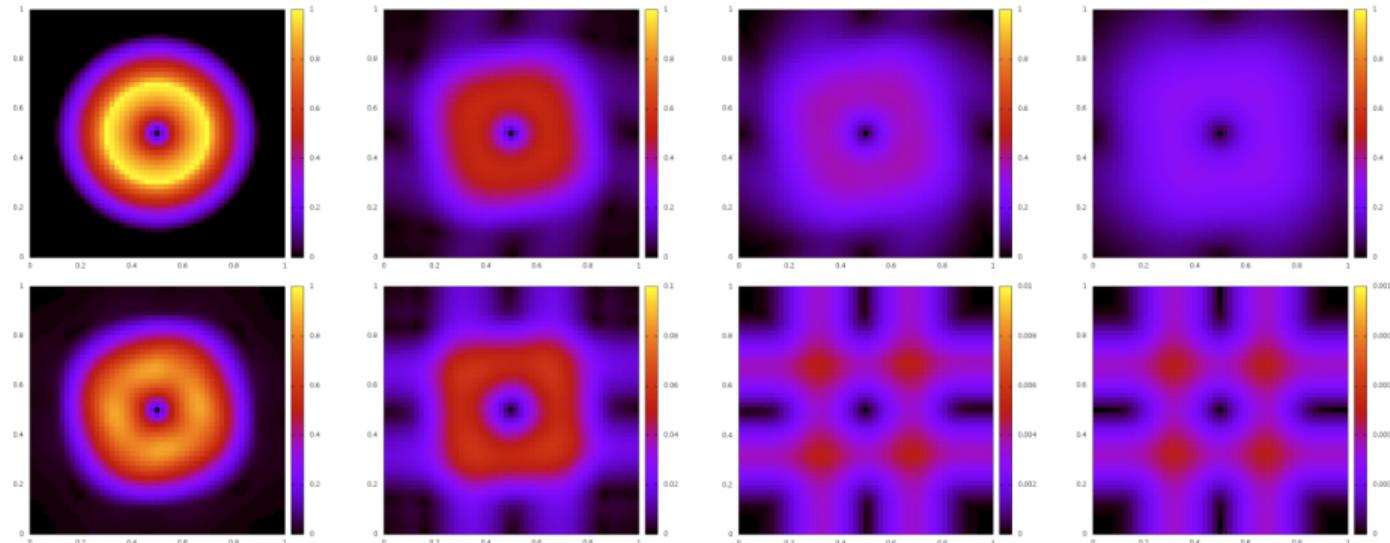


Figure 4.1: Simulation results for a vortex setup for $t = 0, 1, 2, 3$ (from left to right). Colour coded is $\sqrt{u^2 + v^2}$. Top row: Euler equations. Bottom row: Acoustic equations.

12

¹Barsukow, W. Low Mach number finite volume methods for the acoustic and Euler equations, Ph.D. thesis, 2018.

²Finite Volume Upwind numerical flux simulations.

Typical problems: SUPG

Let's try with SUPG.

Hope

$$\int (\varphi + \alpha \Delta x \partial_x \varphi \operatorname{sign}(J^x) + \alpha \Delta y \partial_y \varphi \operatorname{sign}(J^y)) \begin{pmatrix} \partial_t u + \partial_x p \\ \partial_t v + \partial_y p \\ \partial_t p + \partial_x u + \partial_y v \end{pmatrix} dx = 0$$

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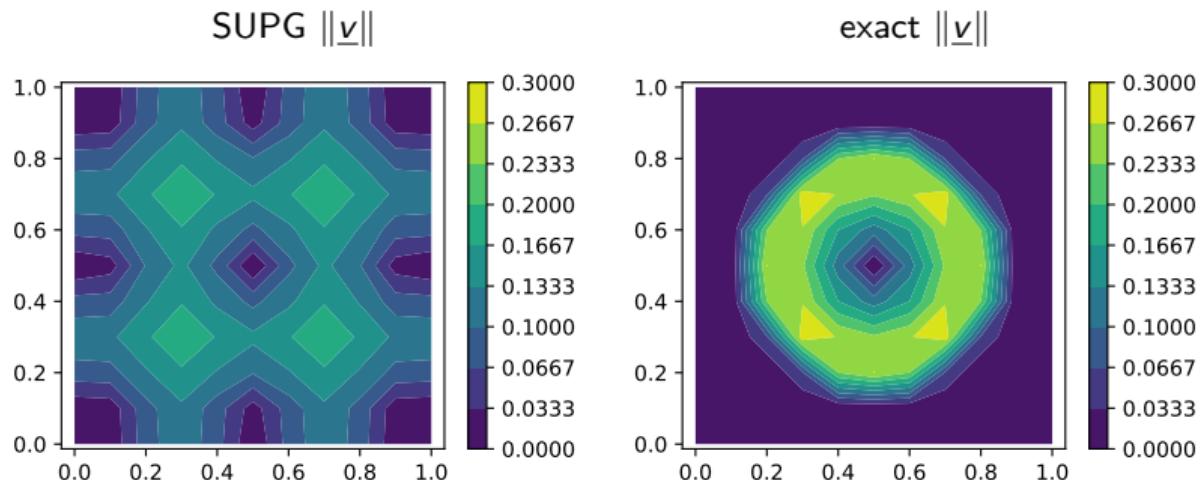
Typical problems: SUPG

Discretization

- Cartesian grid
- CG-FEM
- SUPG
- \mathbb{Q}^1
- $N_x = 10$
- $N_y = 10$

Test

- Vortex \underline{v}
- $p \equiv 1$
- Long time simulation
 $T = 100$



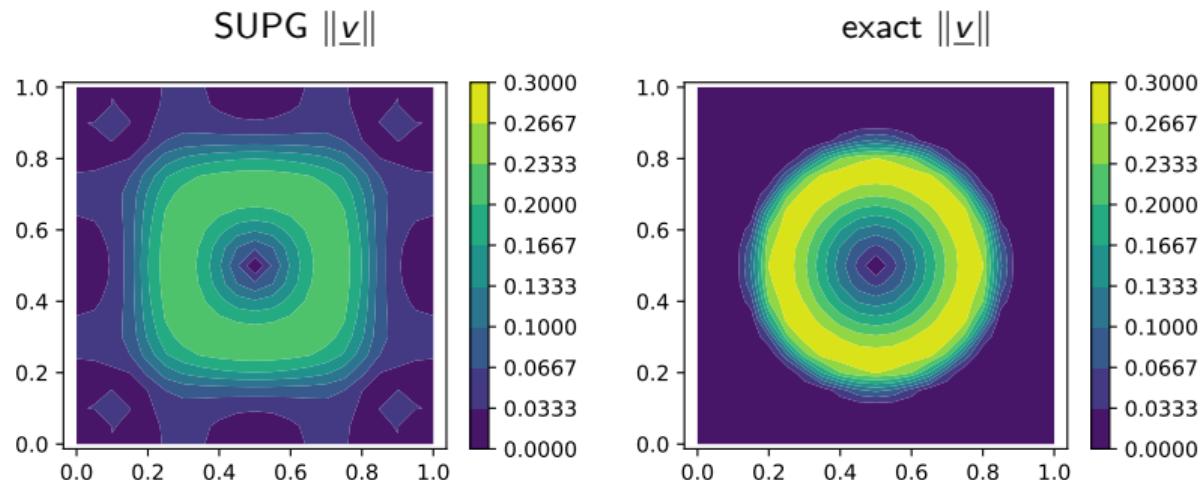
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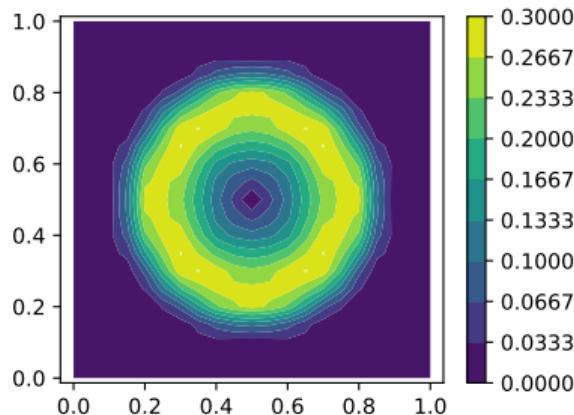
Discretization

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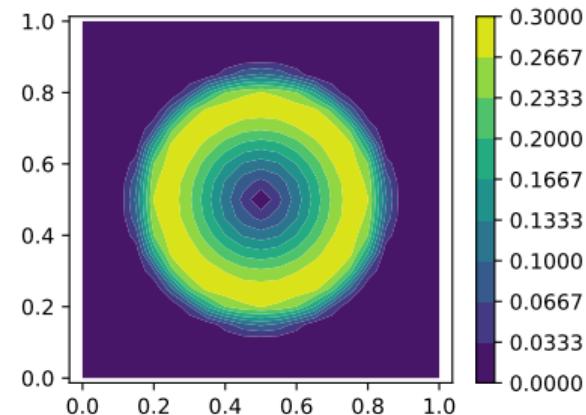
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SUPG $\|\underline{v}\|$



exact $\|\underline{v}\|$



Typical problems: SUPG

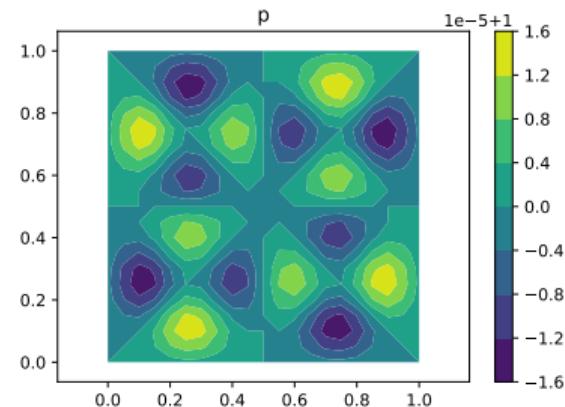
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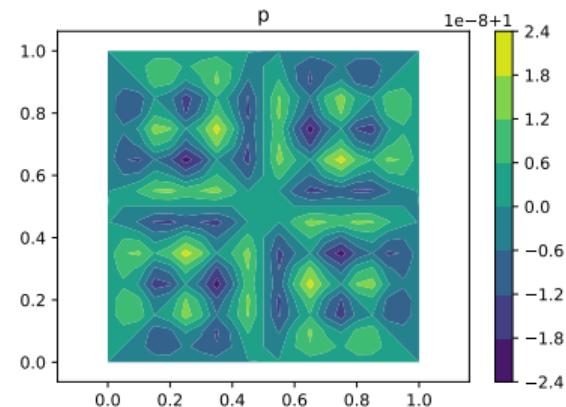
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SUPG $p, \mathbb{Q}^1, N_x = N_y = 20$



exact $p, \mathbb{Q}^2, N_x = N_y = 10$



Why also SUPG?

Ideal

- At equilibrium $\partial_t q = 0$

SUPG formulation

$$\int \varphi(\partial_t u + \partial_x p) + \alpha \Delta x \partial_x \varphi (\partial_t p + \partial_x u + \partial_y v) = 0$$
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- At equilibrium $\partial_t q = 0$
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$$\begin{aligned} & \int \varphi (\partial_t u + \partial_x p) + \alpha \Delta x \partial_x \varphi (\partial_t p + \partial_x u + \partial_y v) = 0 \\ & \int \varphi (\partial_t v + \partial_y p) + \alpha \Delta y \partial_y \varphi (\partial_t p + \partial_x u + \partial_y v) = 0 \\ & \int \varphi (\partial_t p + \partial_x u + \partial_y v) + \\ & \alpha \Delta x \partial_x \varphi (\partial_t u + \partial_x p) + \alpha \Delta y \partial_y \varphi (\partial_t v + \partial_y p) = 0 \end{aligned}$$

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Practice

We just know that the combination of the operators is equal to 0.

Moreover, we would like to know the relation between the following

$$\ker \left[\int \varphi (\partial_x u + \partial_y v) dx \quad \forall \varphi \in V_{h,0}^K \right] \quad \not\subset \quad \ker \left[\int \partial_x \varphi (\partial_x u + \partial_y v) dx \quad \forall \varphi \in V_{h,0}^K \right].$$

How to solve the problem?

Recipe?

- 2D operators with more recognizable kernels
- Recast 2D operators to 1D operators to easily study their kernels
- Divergence operator that should be a Kronecker product of operators

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Global Flux

Reminder of what is Global Flux 1D (for balance laws)

$$\partial_t U + \partial_x F(U) = S(U)$$

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Global Flux SUPG for acoustics

Define $\sigma_x(x, y) := \int_{y_0}^y u(x, s) ds$ and $\sigma_y(x, y) := \int_{x_0}^x v(s, y) ds$, with $\sigma_x, \sigma_y \in V_h^K(\Omega_h)$, $\Phi := \sigma_x + \sigma_y$.

$$\int \varphi (\partial_t u + \partial_x p) + \alpha \Delta x \partial_x \varphi (\partial_t p + \partial_x \partial_y (\sigma_x + \sigma_y)) = 0$$

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$$\int \varphi(\partial_t u + \partial_x p) + \alpha \Delta x \partial_x \varphi (\partial_t p + \partial_x \partial_y \Phi) = 0$$

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$$\int \varphi(\partial_t u + \partial_x p) + \alpha \Delta x \partial_x \varphi (\partial_t p + \partial_x \partial_y \Phi(u, v)) = 0$$

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Changes in equilibrium

$$\nabla \cdot \underline{v} = 0$$

$$\implies \partial_x \partial_y (\sigma_x + \sigma_y) = 0$$

$$\iff \sigma_x + \sigma_y = f(x) + g(y)$$

Discrete operators on Cartesian grid

- $\int \partial_x \varphi \partial_x \partial_y \Phi = D_x^x \otimes D_y \Phi$
- $\int \partial_y \varphi \partial_x \partial_y \Phi = D_x \otimes D_y^y \Phi$
- $\int \varphi \partial_x \partial_y \Phi = D_x \otimes D_y \Phi$
- $\Phi_{i,j} := \int_{y_0}^{y_j} u dy + \int_{x_0}^{x_i} v dx$
- $(D_x)_{ij} = \int \phi_i(x) \partial_x \phi_j(x) dx$
- $(D_x^x)_{ij} = \int \partial_x \phi_i(x) \partial_x \phi_j(x) dx$
- $(I_x)_{ij} = \int_{x_0}^{x_j} \phi_j(x) dx$
- $\Phi = \text{Id}_x \otimes I_y u + I_x \otimes \text{Id}_y v$

Discretization

- $u, v, p, \sigma_x, \sigma_y, \Phi \in V_h^K$

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- $p \equiv C$ and u, v :

$$\Phi_{ij} = \int_{y_0}^{y_j} u(x_i, y) dy + \int_{x_0}^{x_i} v(x, y_j) dx$$

- $\Phi_{ij} = f_i + g_j$
- $\Phi_{ij} - \Phi_{i0} - \Phi_{0j} + \Phi_{00} = 0$ for all i, j

Properties summary

Discretization

- $u, v, p, \sigma_x, \sigma_y, \Phi \in V_h^K$

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Dissipation of spurious modes

- Divergence operator $D_x \otimes D_y$ has spurious equilibria
- $D_x^x \otimes D_y$ or $D_x \otimes D_y^y$ dissipate essentially all spurious equilibria (we have a proof)

Involution

- It is “possible” to compute the discrete involution, but not so nice

FEM details

- Lagrangian basis functions
- Gauss–Lobatto nodes for quadrature
- Gauss–Lobatto nodes for basis function
- Tensor product/Kronecher product to 2D structures

FEM and quadrature

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SUPG-GF FEM discretization

$$\Phi := \text{Id}_x \otimes I_y u + I_x \otimes \text{Id}_y v$$

$$0 = M_x \otimes M_y \partial_t u + D_x \otimes M_y p + \alpha \Delta x (D^x \otimes M_y \partial_t p + D_x^x \otimes D_y I_y u + D_x^x I_x \otimes D_y v),$$

$$0 = M_x \otimes M_y \partial_t v + M_x \otimes D_y p + \alpha \Delta y (M_x \otimes D^y \partial_t p + D_x \otimes D_y^y I_y u + D_x I_x \otimes D_y^y v),$$

$$0 = M_x \otimes M_y \partial_t p + D_x \otimes D_y I_y u + D_x I_x \otimes D_y v +$$

$$\alpha (\Delta x D^x \otimes M_y \partial_t u + \Delta y M_x \otimes D^y \partial_t v + (\Delta x D_x^x \otimes M_y + \Delta y M_x \otimes D_y^y) p).$$

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Safety check!

Convergence of method on nonstationary problem with exact solution

$$\begin{cases} u(x, y, t) = -\frac{1}{2c} (\cos(\alpha\xi(x, y) + ct) - \cos(\alpha\xi(x, y) - ct)) \cos(\theta), \\ v(x, y, t) = -\frac{1}{2c} (\cos(\alpha\xi(x, y) + ct) - \cos(\alpha\xi(x, y) - ct)) \sin(\theta), \\ p(x, y, t) = \frac{1}{2} (\cos(\alpha\xi(x, y) + ct) + \cos(\alpha\xi(x, y) - ct)), \end{cases}$$

Smooth nonstationary test: oblique flow

Safety check!

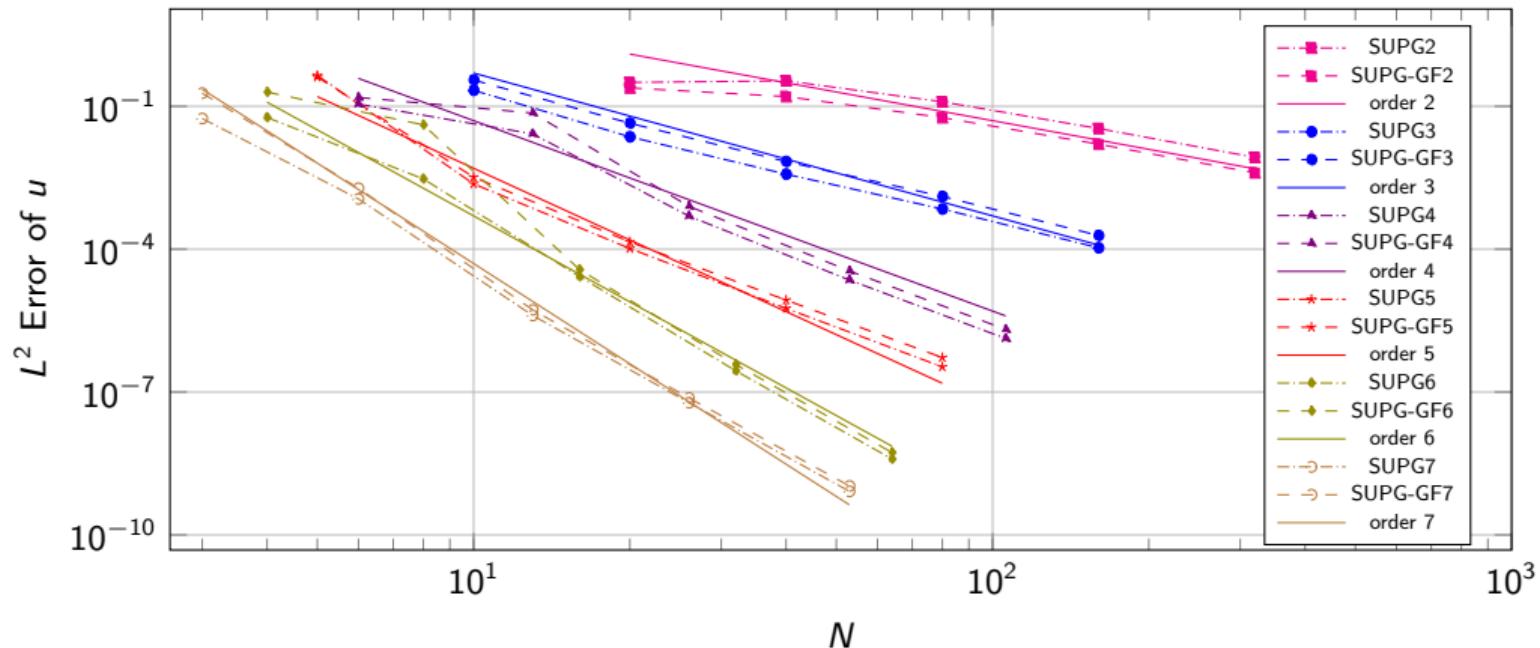


Figure: Oblique flow: convergence of L^2 error of u with respect to the number of elements in x

Simulation of vortex

$$\begin{cases} u(x, y) = f(\rho(x, y)) \cdot (y - y_0) \\ v(x, y) = -f(\rho(x, y)) \cdot (x - x_0) \\ p(x, y) = 1 \end{cases}$$

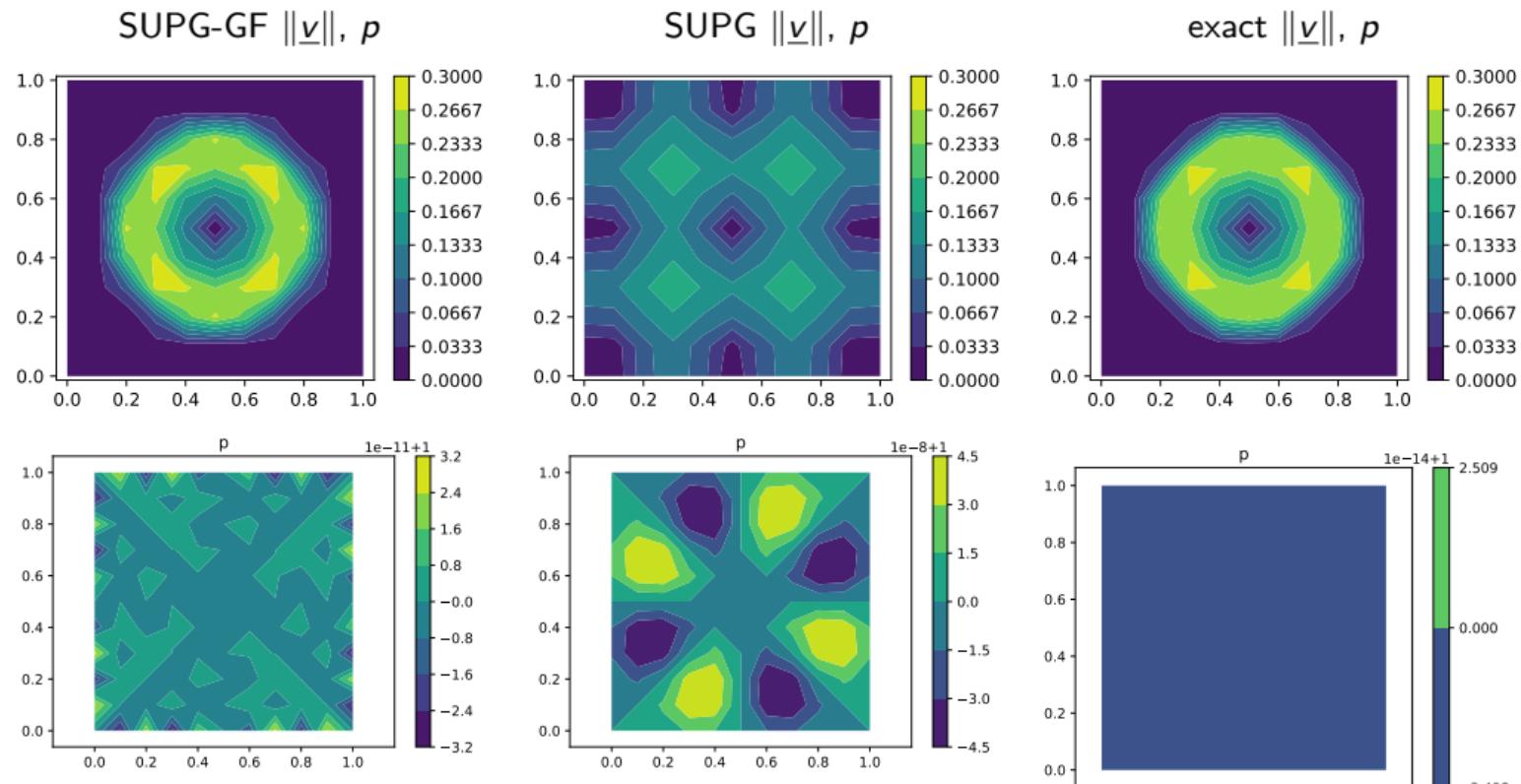
with $\rho(x, y) = \frac{\sqrt{(x-x_0)^2+(y-y_0)^2}}{r_0}$ with $r_0 = 0.45$ the radius of the support.

$$f(\rho) = 2\gamma e^{-\frac{1}{2(1-\rho)^2}} \sqrt{\frac{g}{r_0(1-\rho)^3}}$$

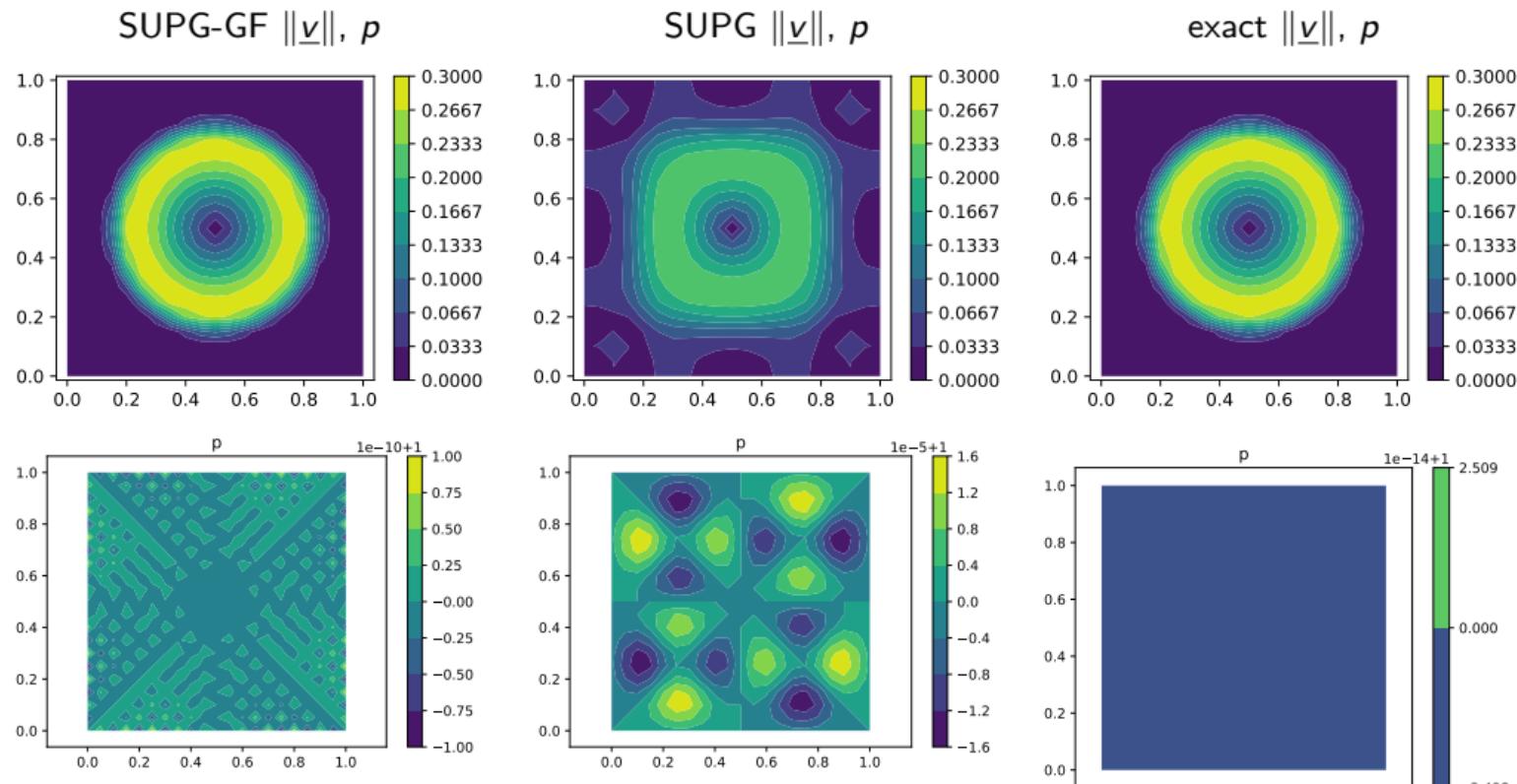
with $g = 9.81$, $\gamma = 0.2$ if $\rho < 1$, else 0.

$$T = 100$$

Simulation of vortex: \mathbb{Q}^1 , $N_x = N_y = 10$

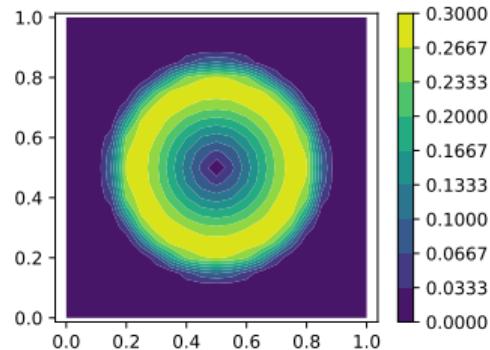


Simulation of vortex: \mathbb{Q}^1 , $N_x = N_y = 20$

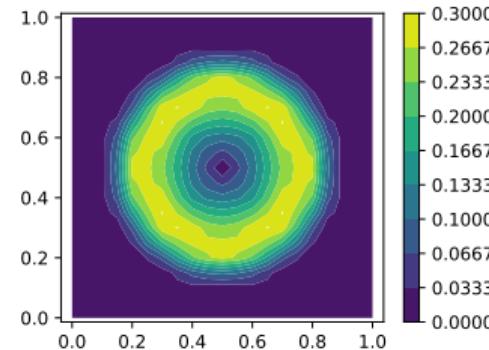


Simulation of vortex: \mathbb{Q}^2 , $N_x = N_y = 10$

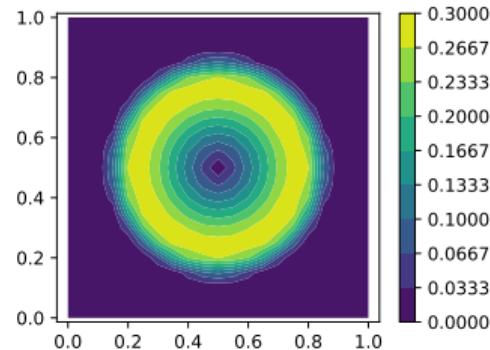
SUPG-GF $\|\underline{v}\|, p$



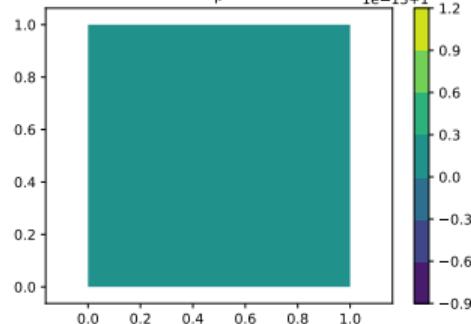
SUPG $\|\underline{v}\|, p$



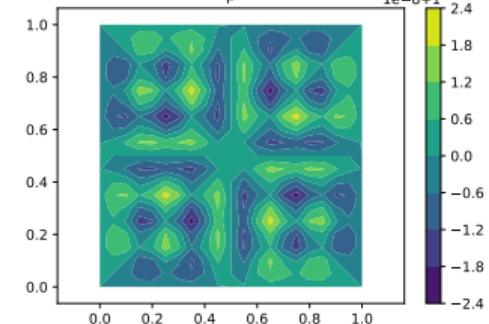
exact $\|\underline{v}\|, p$



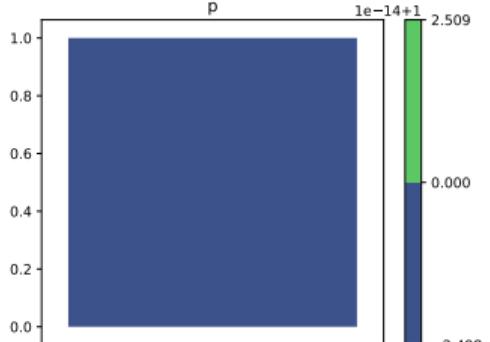
p



p



p



Simulation of vortex: errors

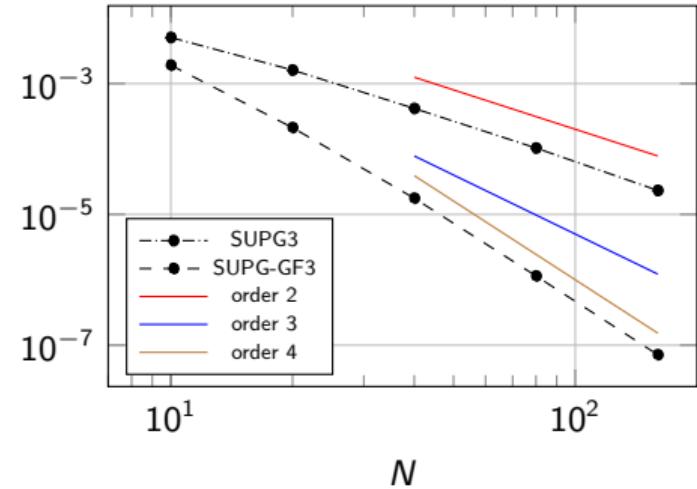
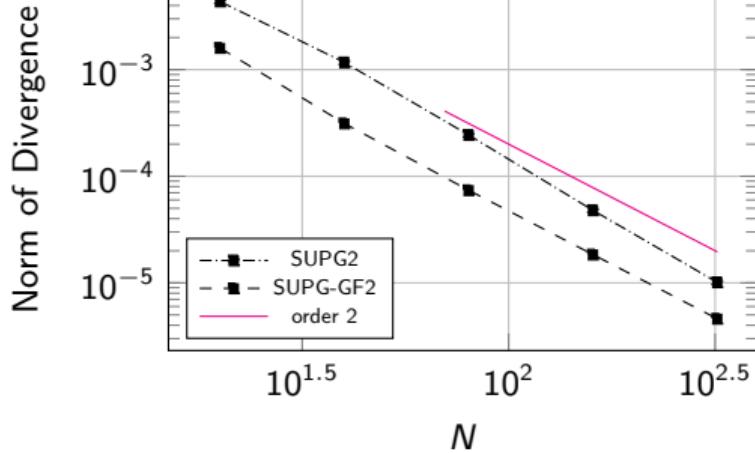


Figure: Smooth vortex: convergence of L^2 error of u with respect to the number of elements in x

Simulation of vortex: errors

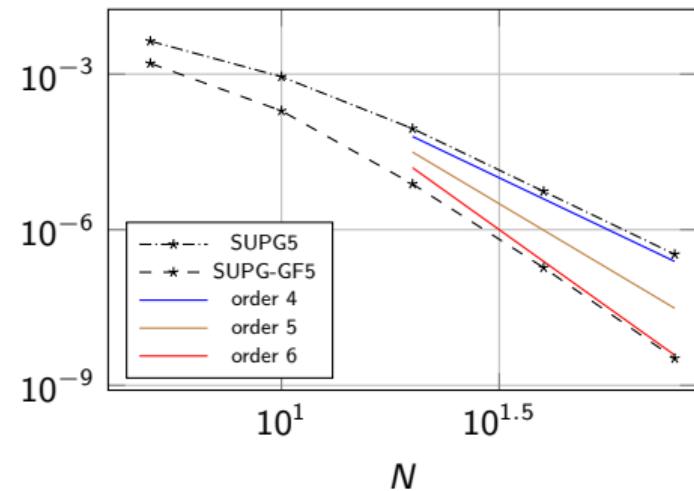
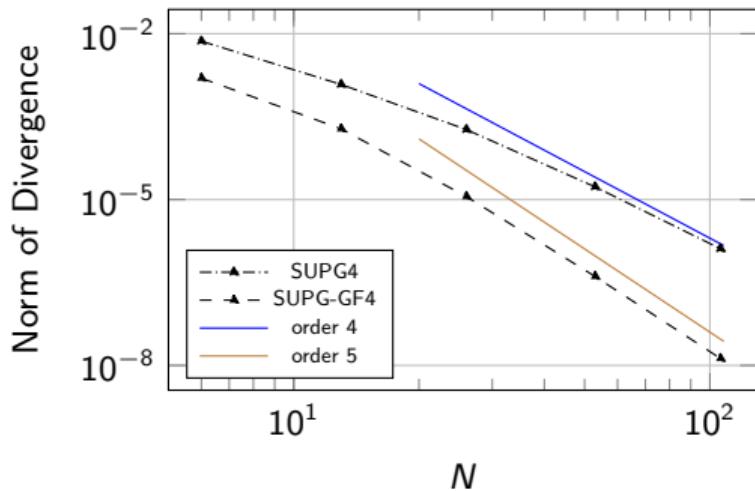


Figure: Smooth vortex: convergence of L^2 error of u with respect to the number of elements in x

Vortex simulation: divergence error

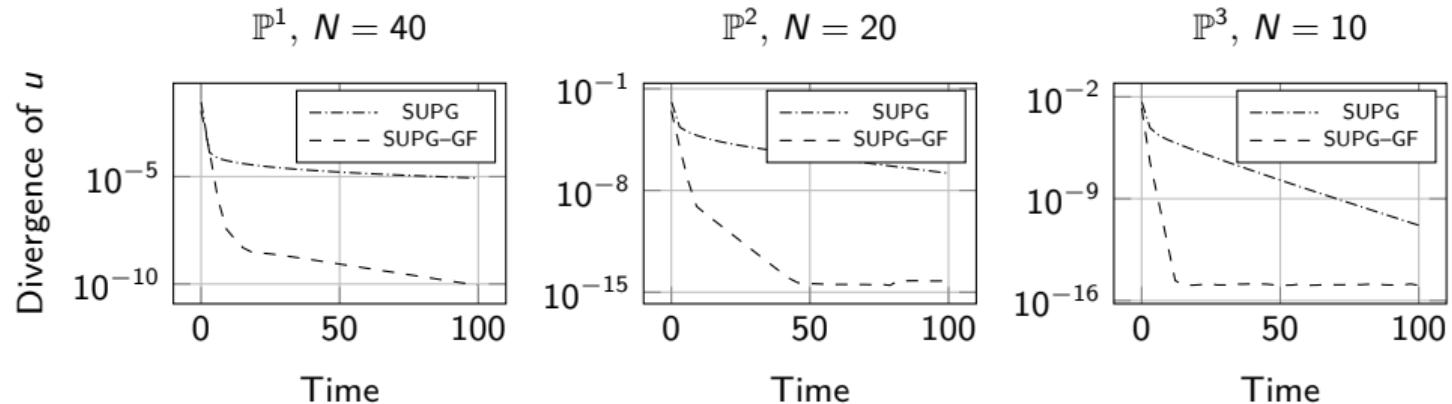


Figure: Norm of discrete divergence of u for SUPG ($\partial_x u + \partial_y v$) and SUPG-GF ($\partial_x \partial_y (\sigma_x + \sigma_y)$) simulations with respect to time for different orders

Pressure perturbation

- Gaussian centered in $\underline{x}_p = (0.4, 0.43)$
- scaling coefficient $r_0 = 0.1$
- radius $\rho(\underline{x}) = \sqrt{\|\underline{x} - \underline{x}_p\|}/r_0$
- final time $T = 0.35$

$$\delta_p(\underline{x}) = \varepsilon e^{-\frac{1}{2(1-\rho(\underline{x}))^2} + \frac{1}{2}},$$

Vortex perturbation

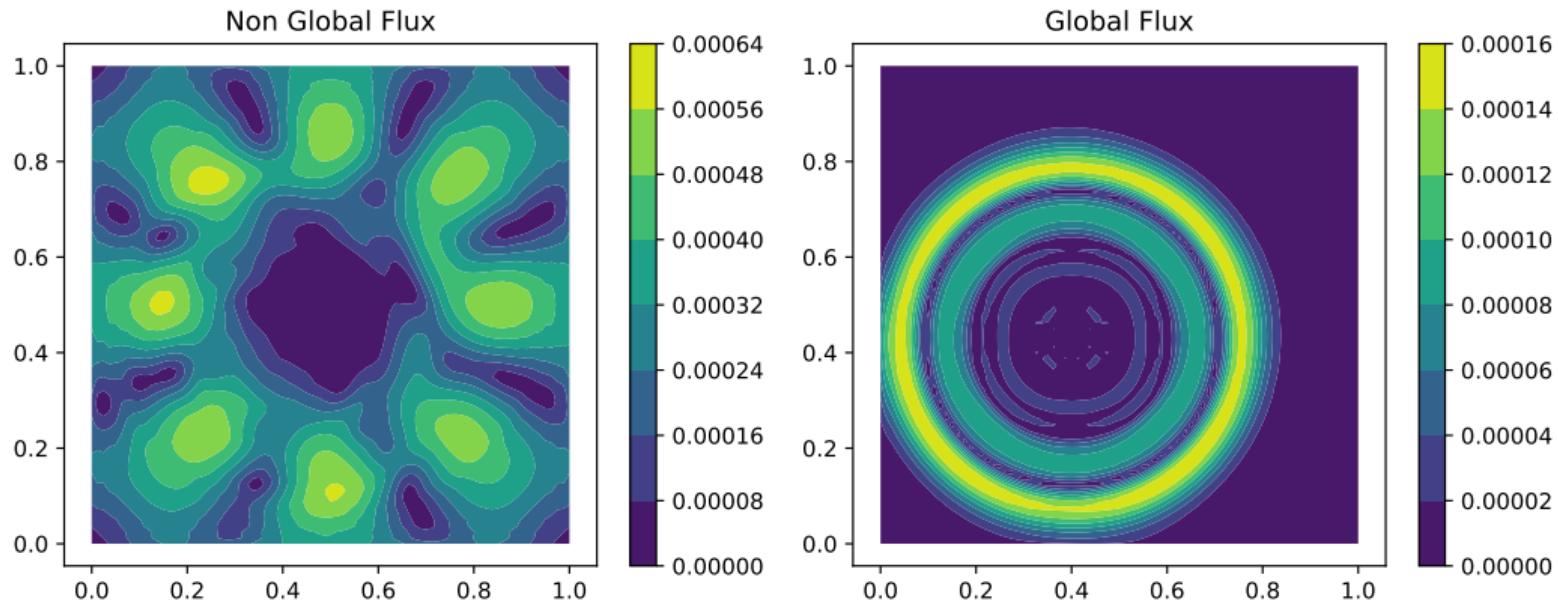


Figure: Perturbation($\varepsilon = 10^{-3}$) test. Plot of $\|\underline{u}_{eq} - \underline{u}_p\|$, with \underline{u}_{eq} the equilibrium obtained with a cheap optimization process. \mathbb{P}^1 with 80×80 cells and 6561 dofs.

Vortex perturbation

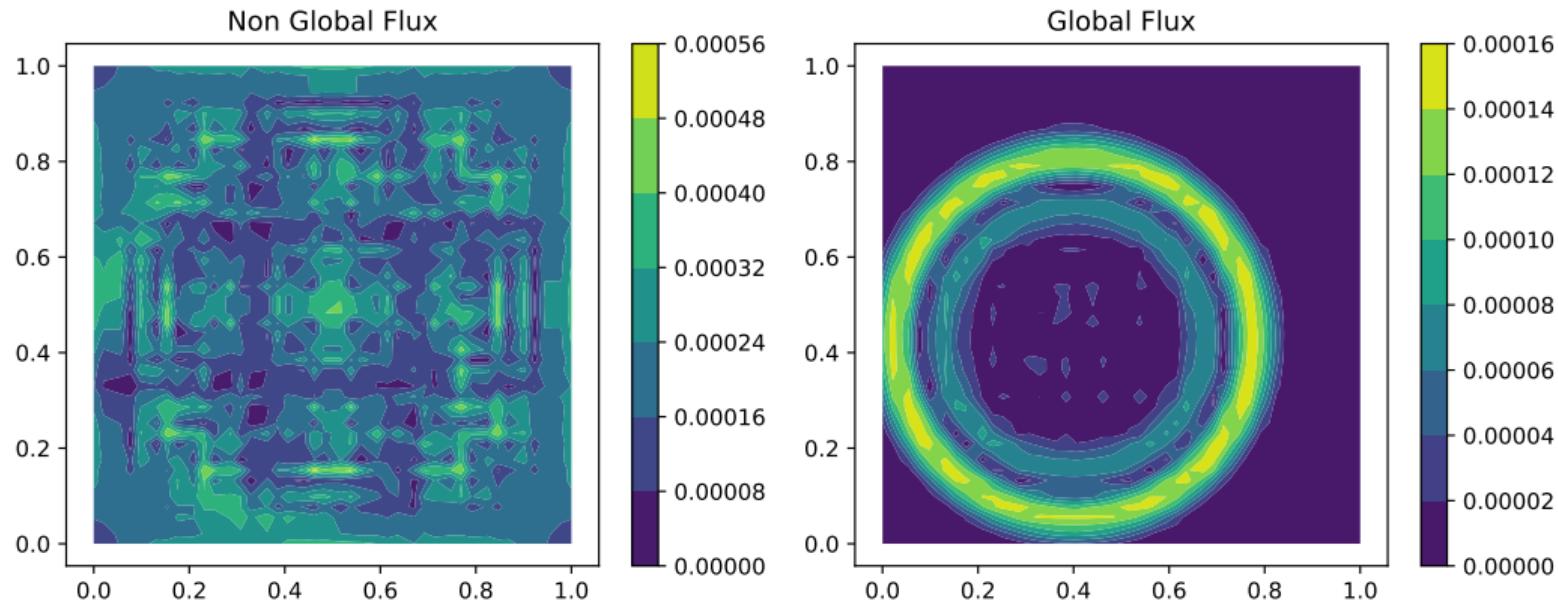


Figure: Perturbation($\varepsilon = 10^{-3}$) test. Plot of $\|\underline{u}_{eq} - \underline{u}_p\|$, with \underline{u}_{eq} the equilibrium obtained with a cheap optimization process. \mathbb{P}^3 with 13×13 cells and 1600 dofs.

Vortex perturbation

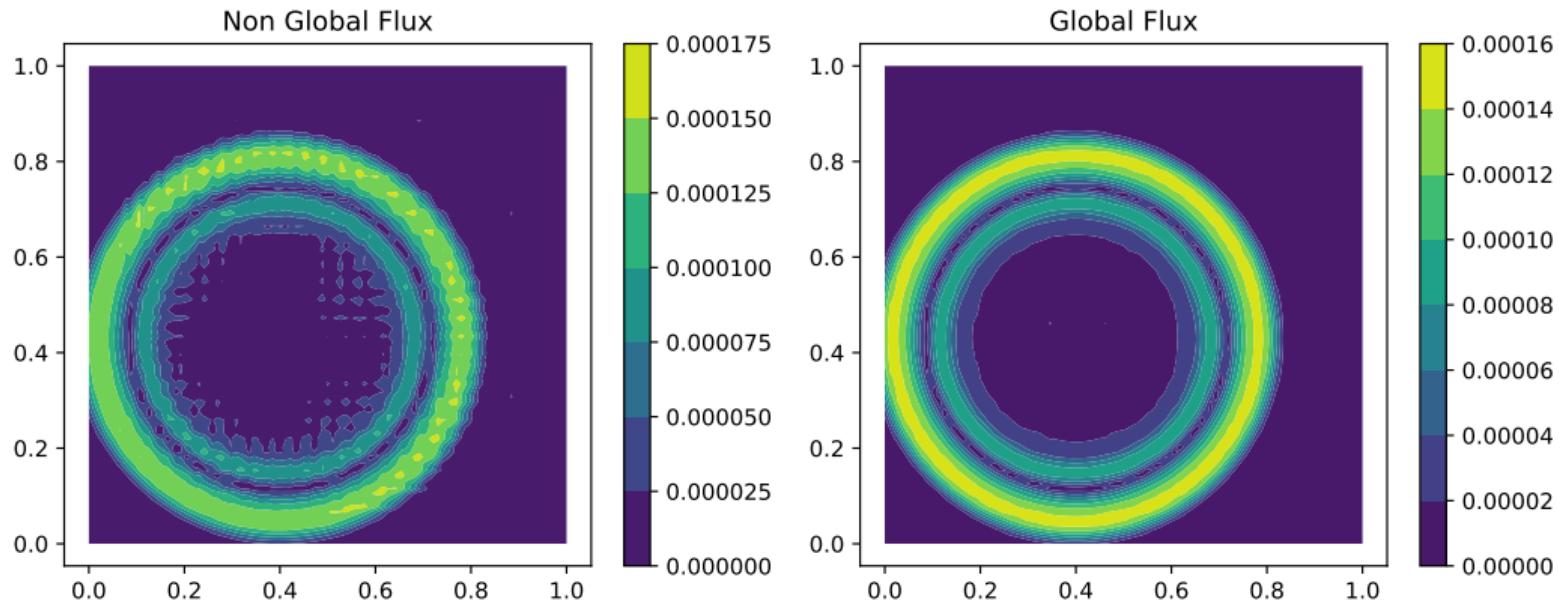


Figure: Perturbation($\varepsilon = 10^{-3}$) test. Plot of $\|\underline{u}_{eq} - \underline{u}_p\|$, with \underline{u}_{eq} the equilibrium obtained with a cheap optimization process. \mathbb{P}^3 with 26×26 cells and 6241 dofs.

Other models

Acoustic with Coriolis

$$\partial_t \begin{pmatrix} u \\ v \\ p \end{pmatrix} + \partial_x \begin{pmatrix} p \\ 0 \\ u \end{pmatrix} + \partial_y \begin{pmatrix} 0 \\ p \\ v \end{pmatrix} + c_f \begin{pmatrix} -v \\ u \\ 0 \end{pmatrix} = 0.$$

Acoustic with source term

$$\partial_t \begin{pmatrix} u \\ v \\ p \end{pmatrix} + \partial_x \begin{pmatrix} p \\ 0 \\ u \end{pmatrix} + \partial_y \begin{pmatrix} 0 \\ p \\ v \end{pmatrix} + \begin{pmatrix} 0 \\ 0 \\ s \end{pmatrix} = 0.$$

Stommel Gyre

$$\partial_t p = - \operatorname{div} \mathbf{u}$$

$$\partial_t \mathbf{u} = - \operatorname{grad} p + \phi \mathbf{u}^\perp - R \mathbf{u} + \boldsymbol{\tau}$$

Model

- Other stabilizations (OSS, CIP)
- Other equations

Triangular meshes

- Haven't tried yet
- In principle, we can still define $\Phi : \int^y u + \int^x v$ in each element
- Question: will it be that effective?
- Kernels? Maybe difficult to write, still working

THANKS!!



References

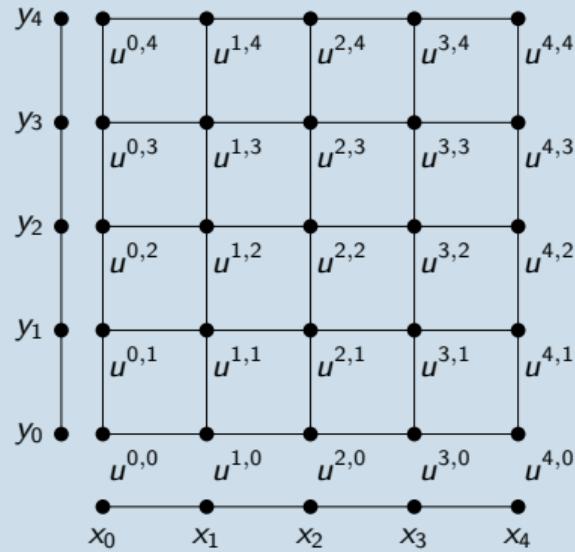
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- M. Ciallella, D. Torlo and M. Ricchiuto. *Arbitrary high order WENO finite volume scheme with flux globalization for moving equilibria preservation*. Journal of Scientific Computing, 96(2):53, 2023.
- W. Barsukow, M. Ricchiuto and D. Torlo. *Structure preserving methods via Global Flux quadrature: divergence-free preservation with continuous Finite Element*. In preparation, 2024.
- davidetorlo.it

Detailed definition of Global Flux SUPG

Definition of σ_x , σ_y

Cartesian grid, Lagrangian basis functions in Lobatto points (x_i, y_j) in each direction.

So, $\phi_i(x_k) = \delta_{ik}$ and $\psi_j(y_\ell) = \delta_{j\ell}$ and



$$u(x, y) = \sum_{i,j} \varphi_{ij}(x, y) u^{i,j} = \sum_{i,j} \phi_i(x) \psi_j(y) u^{ij}$$

$$u(x_i, y_j) = u^{i,j}$$

$$\sigma_x(x, y) = \sum_{i,j} \phi_i(x) \psi_j(y) \sigma_x^{i,j}$$

$$\sigma_x(x_i, y_j) = \sigma_x^{i,j}$$

$$\sigma_x(x, y) = \int_{y_0}^y u(x, s) ds$$

$$\sigma_x^{i,j} = U(x_i, y_j) = \int_{y_0}^{y_j} u(x_i, s) ds = \sum_{k,\ell} \phi_k(x_i) \int_{y_0}^{y_j} \psi_\ell(s) ds u^{k,\ell}$$

So, even if both $\sigma_x, u \in V_h^K$, in quadrature points, we have that exactly $u(x_i, y_j) = \int_{y_0}^{y_j} \sigma_x(x_i, y) dy$.

Myth buster

Global Flux is not global!

- In principle $\sigma_x(x, y) = \int_{y_B}^y u(x, s)ds$ should be integrated from the beginning (bottom) of the domain y_B !
- In practice we always use $\partial_x \partial_y \sigma_x(x, y)$ integrated in one cell!!!!
- So,

$$\sigma_x(x, y) = \int_{y_B}^y u(x, s)ds = \underbrace{\int_{y_B}^{y_0} u(x, s)ds}_{\text{constant in one cell!}} + \int_{y_0}^y u(x, s)ds$$

whatever constant we bring from outside the cell, is canceled out

$$\partial_y \sigma_x(x, y) = \partial_y \int_{y_B}^y u(x, s)ds = \partial_y \int_{y_B}^{y_0} u(x, s)ds + \partial_y \int_{y_0}^y u(x, s)ds = \partial_y \int_{y_0}^y u(x, s)ds$$

Why SUPG-GF works so better?

Clearly divergence-free preserving

- Which divergence? $\partial_x \partial_y (\sigma_x + \sigma_y) \approx \partial_x \partial_y \left(\int^y u(x, s) ds + \int^x v(s, y) ds \right) = \partial_x u + \partial_y v$

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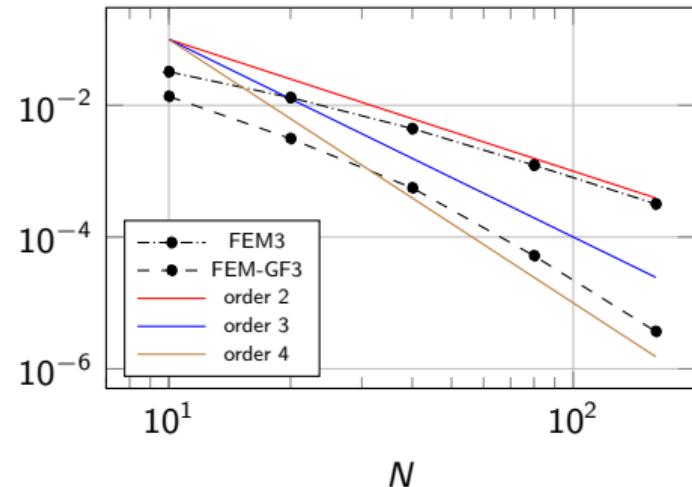
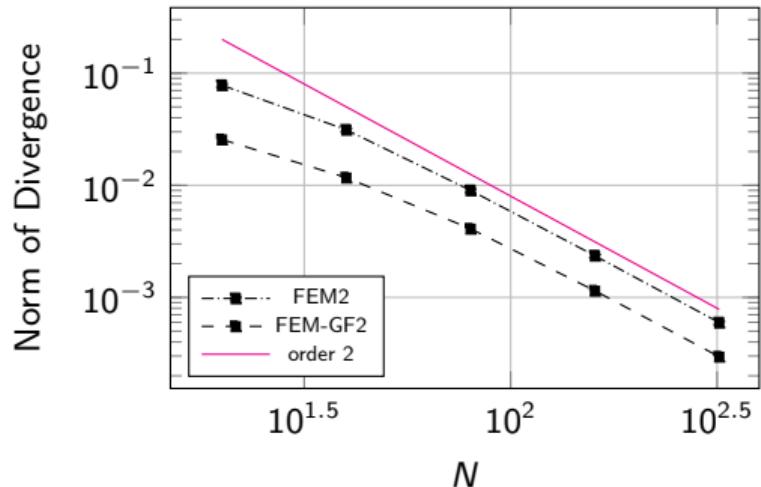


Figure: Smooth vortex: convergence of divergence operator on exact IC with respect to the number of elements in x

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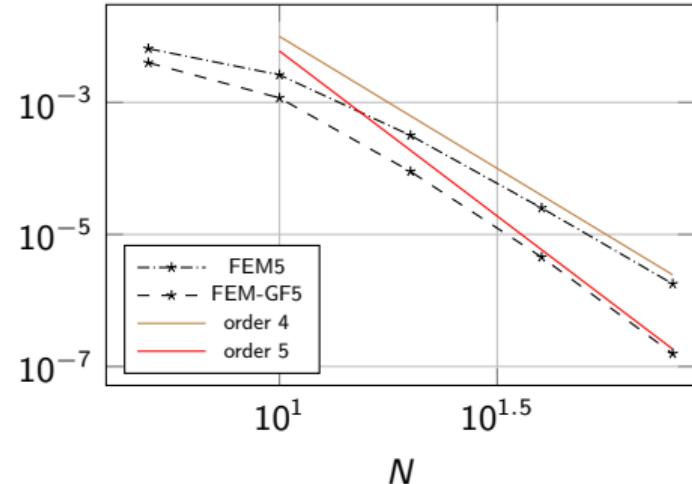
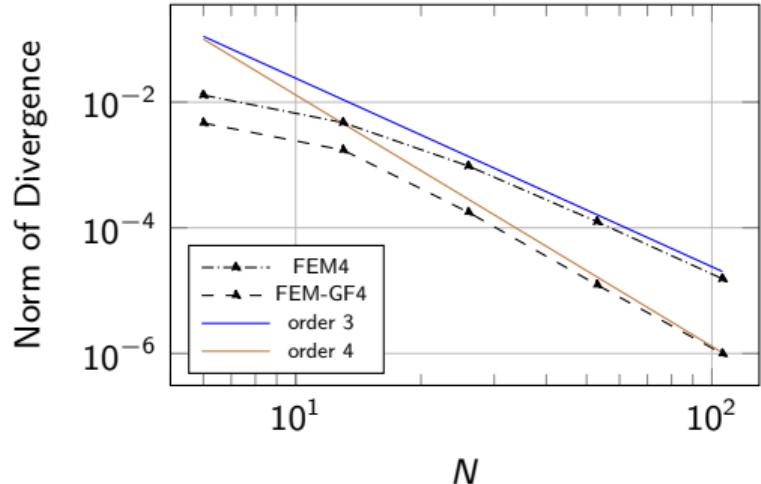


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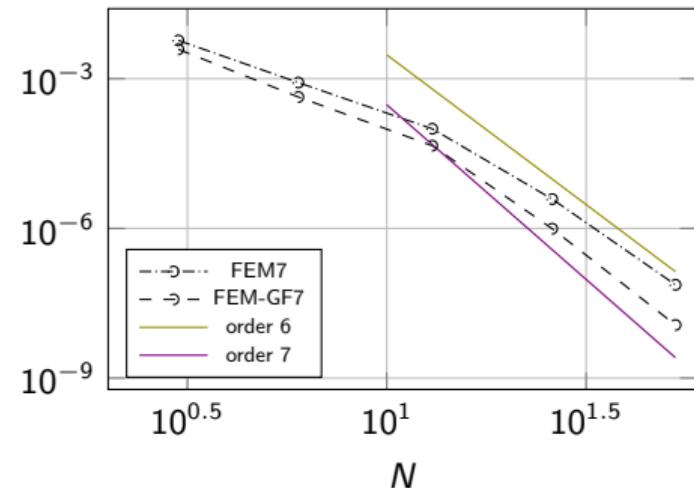
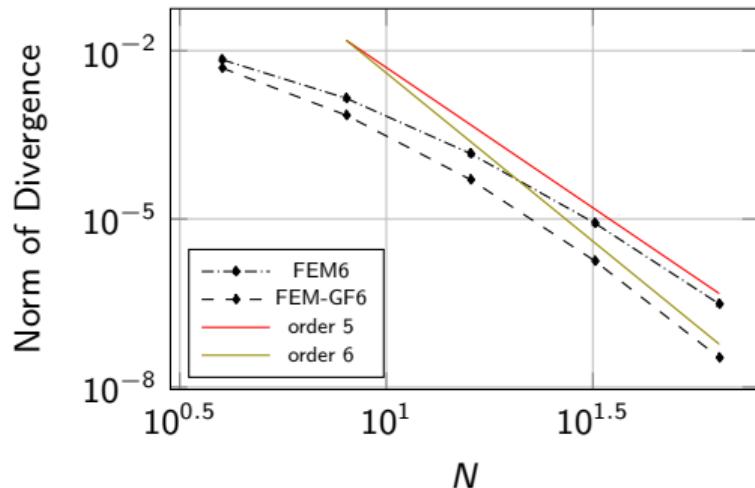


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Why SUPG-GF works so better?

Clearly divergence-free preserving

- Which divergence? $\partial_x \partial_y (\sigma_x + \sigma_y) \approx \partial_x \partial_y \left(\int^y u(x, s) ds + \int^x v(s, y) ds \right) = \partial_x u + \partial_y v$
- If we know that $\partial_x \partial_y (\sigma_x + \sigma_y) = 0$ and $p \equiv c$ then equilibrium

Why SUPG-GF works so better?

New operators kernels

$$\Phi = \sigma_x + \sigma_y$$

$$\int_{\Omega_h} \varphi(x, y) \partial_x \partial_y (\Phi) dx dy = 0 \quad \forall \varphi \in V_{h,0}^K$$

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$$\int_{\Omega_h} \partial_y \varphi(x, y) \partial_x \partial_y (\Phi) dx dy = 0 \quad \forall \varphi \in V_{h,0}^K$$

Matrix formulation

$$(D_x)_{ij} := \int \phi_i(x) \partial_x \phi_j(x) dx \quad (D_x^x)_{ij} := \int \partial_x \phi_i(x) \partial_x \phi_j(x) dx$$

$$\Phi = \sigma_x + \sigma_y \quad \Phi \in \mathbb{R}^{(N_x K+1) \times (N_y K+1)}$$

$$(D_x \otimes D_y) \Phi = 0 \quad D_x, D_x^x \in \mathbb{R}^{(N_x K-1) \times (N_x K+1)}$$

$$(D_x^x \otimes D_y) \Phi = 0 \quad D_y, D_y^y \in \mathbb{R}^{(N_y K-1) \times (N_y K+1)}$$

$$(D_x \otimes D_y^y) \Phi = 0$$

Kernels of Kronecker products

$$M_x \otimes M_y \Phi = 0 \iff$$

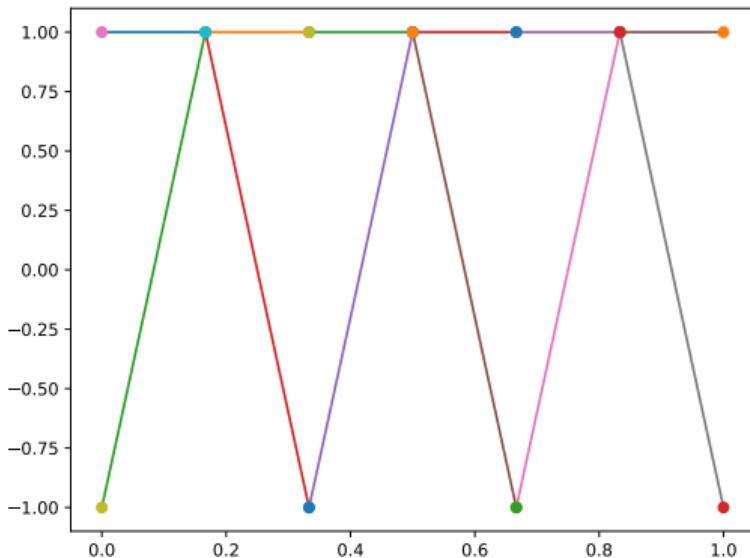
$$M_x \Phi^{\cdot, j} = 0 \forall j \text{ or } M_y \Phi^{i, \cdot} = 0 \forall i$$

We can pass from the study of the 2D operators to the 1D operators!

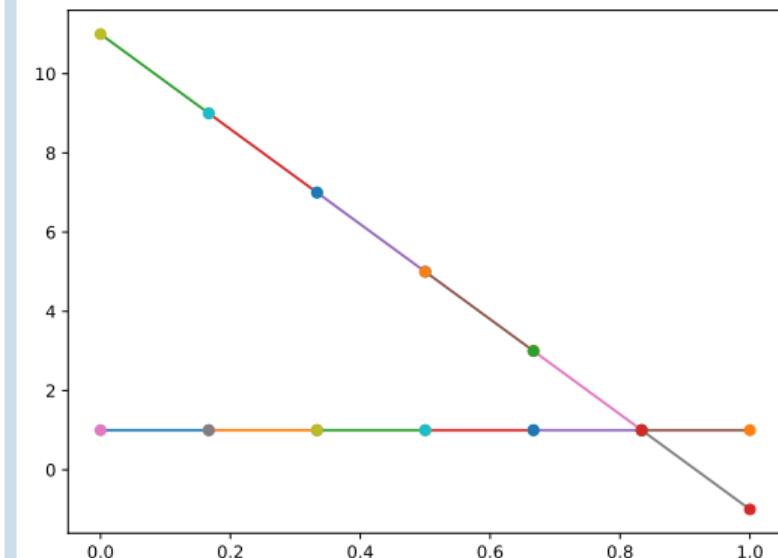
Reminder: before it was not possible because we had a combination of operators $D_x u + D_y v = 0$.

One dimensional kernels of D_x and D_x^x

Kernel of D_x



Kernel of D_x^x

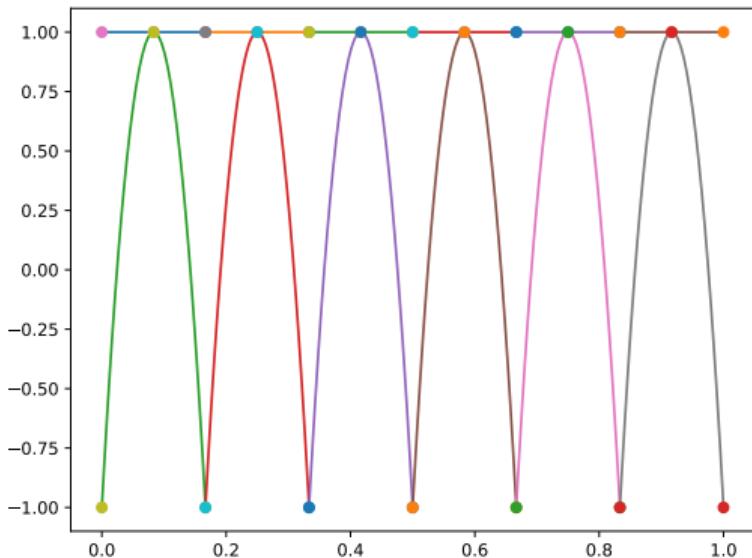


Operators

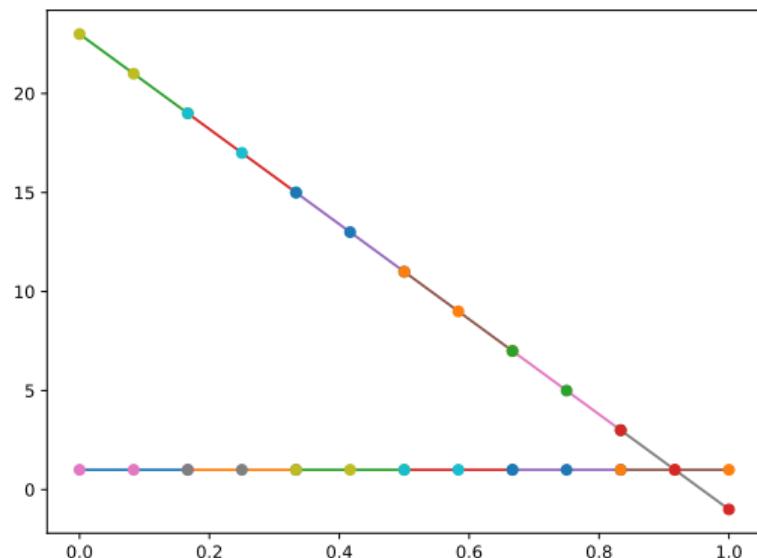
- Divergence $D_x \otimes D_y$
- Stabilization $D_x^x \otimes D_y$, $D_x \otimes D_y^y$

One dimensional kernels of D_x and D_x^x

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Kernel of D_x^x

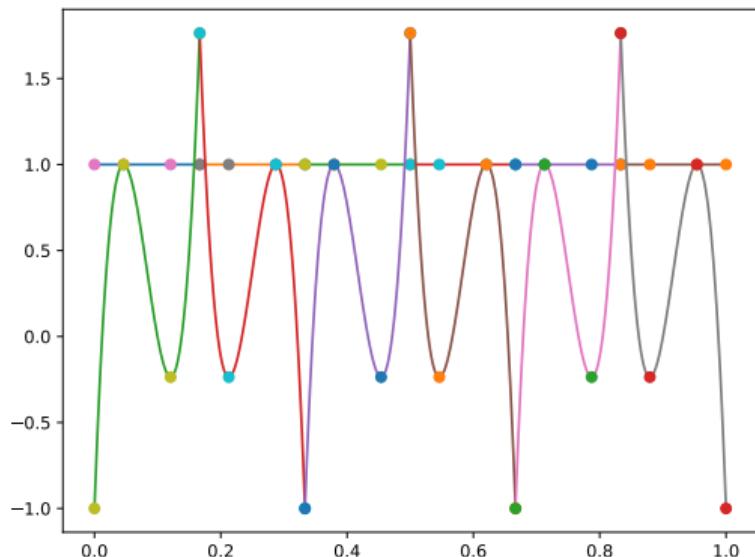


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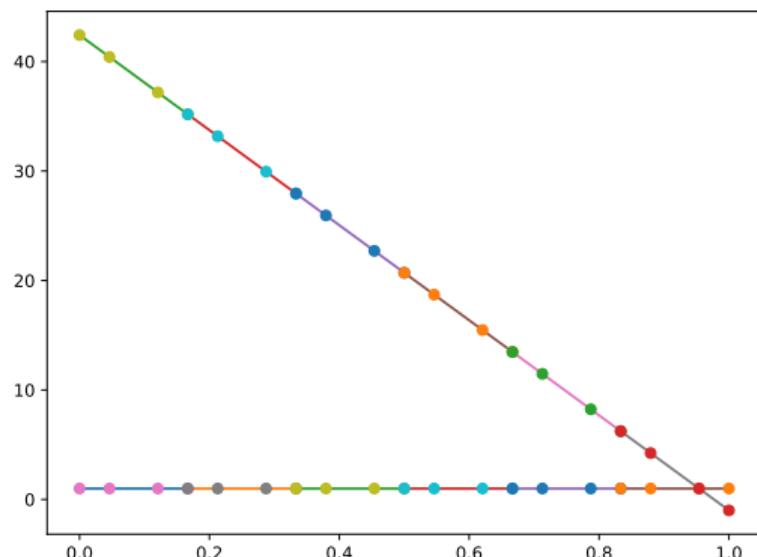
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Deferred Correction Iterative procedure

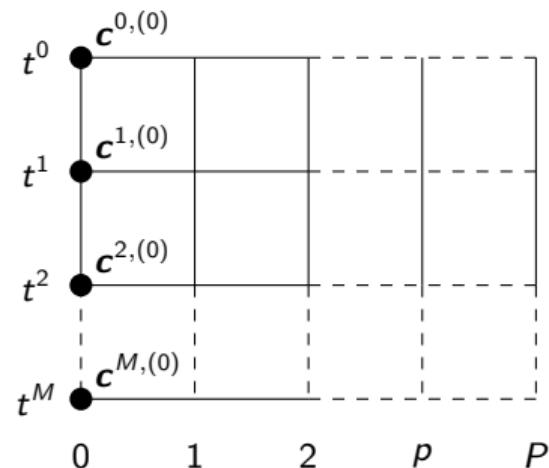
How to combine two methods keeping the accuracy of the second and the stability and simplicity of the first one?

$$\underline{\mathbf{c}}^{0,(p)} := \mathbf{c}(t_n), \quad p = 0, \dots, P,$$

$$\underline{\mathbf{c}}^{m,(0)} := \mathbf{c}(t_n), \quad m = 1, \dots, M$$

$$T^1(\underline{\mathbf{c}}^{(p)}) = T^1(\underline{\mathbf{c}}^{(p-1)}) - T^2(\underline{\mathbf{c}}^{(p-1)}) \text{ with } p = 1, \dots, P.$$

- $T^1(\underline{\mathbf{c}}) = 0$, first order accuracy, easily invertible.
- $T^2(\underline{\mathbf{c}}) = 0$, high order Q .



DeC Theorem

- T^1 coercive with constant $\mathcal{O}(1)$
- $T^1 - T^2$ Lipschitz with constant $\mathcal{O}(\Delta t)$

DeC converges and $\min(P, Q)$ is the order of accuracy.

Deferred Correction Iterative procedure

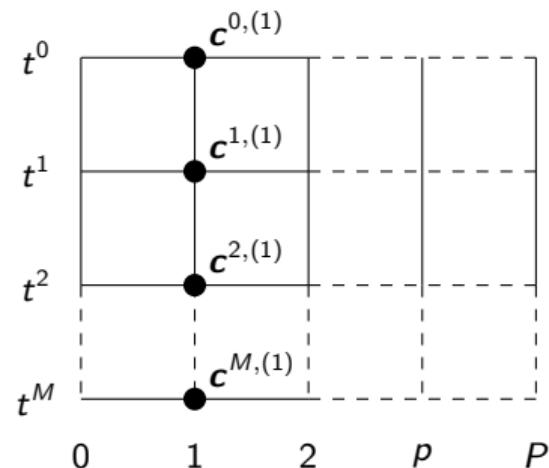
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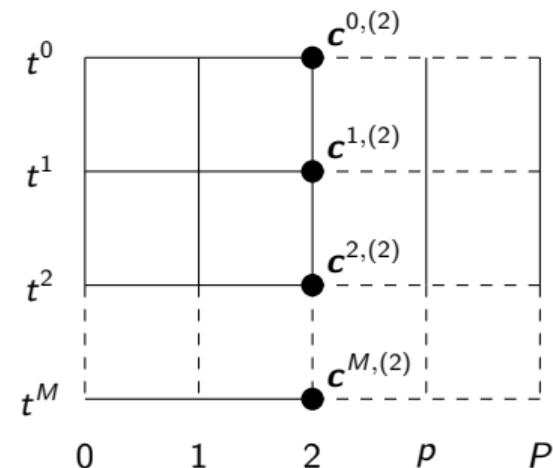
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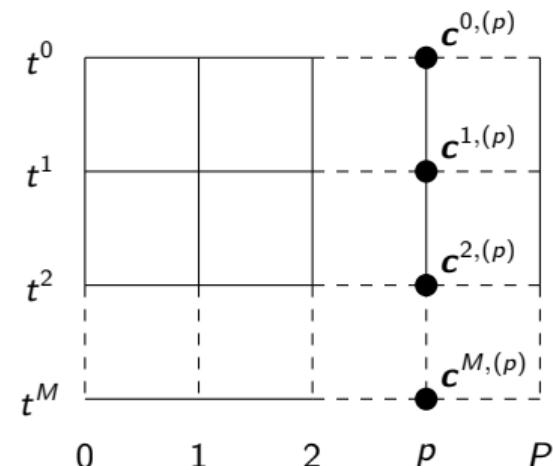
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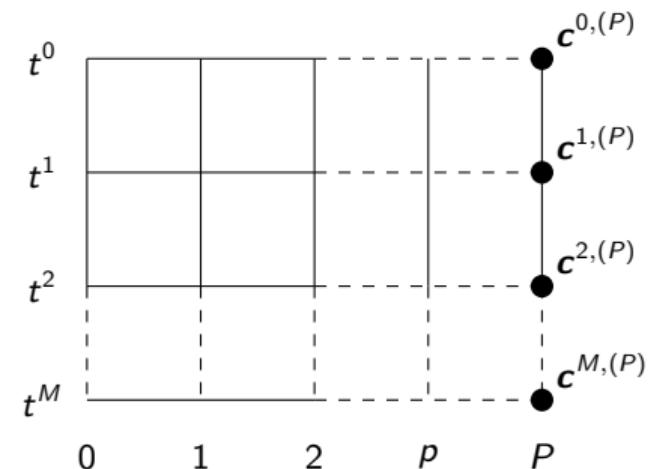
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$$T_u^{2,m}(\underline{\underline{q}}) = M_x \otimes M_y \frac{u^m - u^0}{\Delta t} + D_x \otimes M_y \sum_r \theta_r^m p^r + \alpha h (D^x \otimes M_y \frac{p^m - p^0}{\Delta t} + D_x^x \otimes D_y I_y \sum_r \theta_r^m u^r + D_x^x I_x \otimes D_y \sum_r \theta_r^m v^r), \quad (3a)$$

$$T_v^{2,m}(\underline{\underline{q}}) = M_x \otimes M_y \frac{v^m - v^0}{\Delta t} + M_x \otimes D_y \sum_r \theta_r^m p^r + \alpha h (M_x \otimes D^y \frac{p^m - p^0}{\Delta t} + D_x \otimes D_y^y I_y \sum_r \theta_r^m u^r + D_x I_x \otimes D_y^y \sum_r \theta_r^m v^r), \quad (3b)$$

$$T_p^{2,m}(\underline{\underline{q}}) = M_x \otimes M_y \frac{p^m - p^0}{\Delta t} + D_x \otimes D_y I_y \sum_r \theta_r^m u^r + D_x I_x \otimes D_y \sum_r \theta_r^m v^r + \alpha h (D^x \otimes M_y \frac{u^m - u^0}{\Delta t} + M_x \otimes D^y \frac{v^m - v^0}{\Delta t} + (D_x^x \otimes M_y + M_x \otimes D_y^y) \sum_r \theta_r^m p^r). \quad (3c)$$

$$T_u^{1,m}(\underline{\underline{q}}) = M_x \otimes M_y \frac{u^m - u^0}{\Delta t} + \beta^m D_x \otimes M_y p^0, \quad (2a)$$

$$T_v^{1,m}(\underline{\underline{q}}) = M_x \otimes M_y \frac{v^m - v^0}{\Delta t} + \beta^m M_x \otimes D_y p^0, \quad (2b)$$

$$T_p^{1,m}(\underline{\underline{q}}) = M_x \otimes M_y \frac{p^m - p^0}{\Delta t} + \beta^m (D_x \otimes D_y I_y u^0 + D_x I_x \otimes D_y v^0). \quad (2c)$$

$$\begin{aligned} T_u^{2,m}(\underline{\underline{q}}) = & M_x \otimes M_y \frac{u^m - u^0}{\Delta t} + D_x \otimes M_y \sum_r \theta_r^m p^r + \\ & \alpha h (D^x \otimes M_y \frac{p^m - p^0}{\Delta t} + \textcolor{red}{D_x^x \otimes D_y I_y} \sum_r \theta_r^m u^r + \textcolor{red}{D_x^x I_x \otimes D_y} \sum_r \theta_r^m v^r), \end{aligned} \quad (3a)$$

$$\begin{aligned} T_v^{2,m}(\underline{\underline{q}}) = & M_x \otimes M_y \frac{v^m - v^0}{\Delta t} + M_x \otimes D_y \sum_r \theta_r^m p^r + \\ & \alpha h (M_x \otimes D^y \frac{p^m - p^0}{\Delta t} + \textcolor{red}{D_x \otimes D_y^y I_y} \sum_r \theta_r^m u^r + \textcolor{red}{D_x I_x \otimes D_y^y} \sum_r \theta_r^m v^r), \end{aligned} \quad (3b)$$

$$\begin{aligned} T_p^{2,m}(\underline{\underline{q}}) = & M_x \otimes M_y \frac{p^m - p^0}{\Delta t} + \textcolor{red}{D_x \otimes D_y I_y} \sum_r \theta_r^m u^r + \textcolor{red}{D_x I_x \otimes D_y} \sum_r \theta_r^m v^r + \\ & \alpha h (D^x \otimes M_y \frac{u^m - u^0}{\Delta t} + M_x \otimes D^y \frac{v^m - v^0}{\Delta t} + (D_x^x \otimes M_y + M_x \otimes D_y^y) \sum_r \theta_r^m p^r). \end{aligned} \quad (3c)$$

Vortex with Coriolis

Acoustic with Coriolis

$$\partial_t \begin{pmatrix} u \\ v \\ p \end{pmatrix} + \partial_x \begin{pmatrix} p \\ 0 \\ u \end{pmatrix} + \partial_y \begin{pmatrix} 0 \\ p \\ v \end{pmatrix} + c_f \begin{pmatrix} -v \\ u \\ 0 \end{pmatrix} = 0.$$

GF for Acoustic with Coriolis

$$\partial_t \begin{pmatrix} u \\ v \\ p \end{pmatrix} + \begin{pmatrix} \partial_x(p + \sigma_y) \\ \partial_y(p - \sigma_x) \\ \partial_x \partial_y(\sigma_x + \sigma_y) \end{pmatrix} = 0$$

FEM change

$$T_u^{2,m}(\underline{\underline{q}}) += -c_f M_x \otimes M_y v^m$$

$$T_v^{2,m}(\underline{\underline{q}}) += c_f M_x \otimes M_y u^m$$

$$T_p^{2,m}(\underline{\underline{q}}) += c_f \alpha \sum_r \theta_r^m (M_x \otimes D^y u^r - D^x \otimes M_y v^r)$$

GF-FEM change

$$T_u^{2,m}(\underline{\underline{q}}) += -c_f \textcolor{red}{D_x I_x} \otimes M_x v^m$$

$$T_v^{2,m}(\underline{\underline{q}}) += c_f M_x \otimes \textcolor{red}{D_y I_y} u^m$$

$$T_p^{2,m}(\underline{\underline{q}}) += c_f \alpha \sum_r \theta_r^m (M_x \otimes \textcolor{red}{D_y^y I_y} u^r - \textcolor{red}{D_x^x I_x} \otimes M_y v^r)$$

Test

- $$\begin{cases} u(x, y) = -f(\rho(x, y)) \cdot (y - y_0), \\ v(x, y) = f(\rho(x, y)) \cdot (x - x_0), \\ p(x, y) = 1 - c_f \cdot g(\rho(x, y)), \end{cases}$$
- $\rho(x, y) = \sqrt{x^2 + y^2}$
- $f(\rho) := 20e^{-100\rho^2}$
- $g(\rho) := \frac{1}{10}e^{-100\rho^2}$
- Domain $\Omega = [0, 1]^2$

Vortex with Coriolis

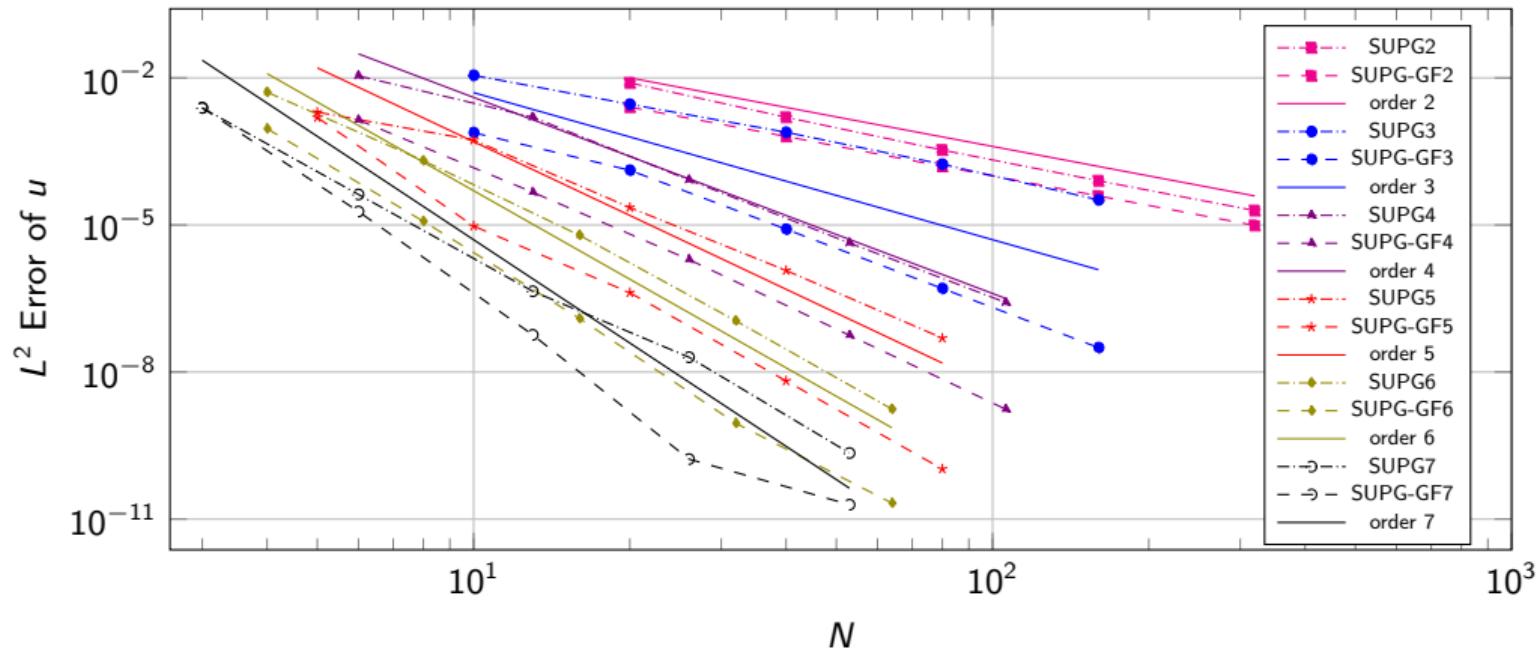
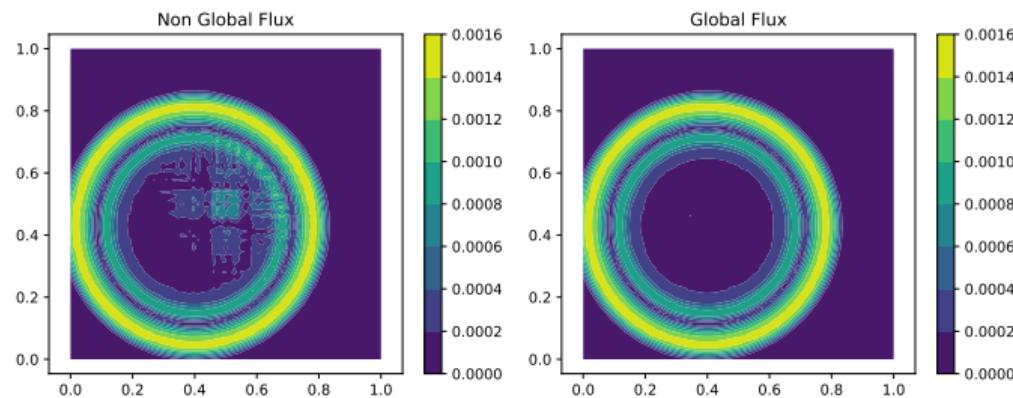
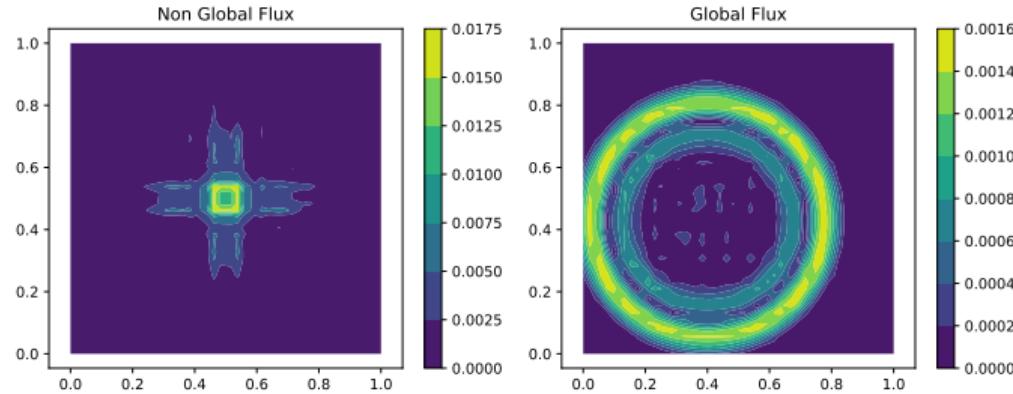


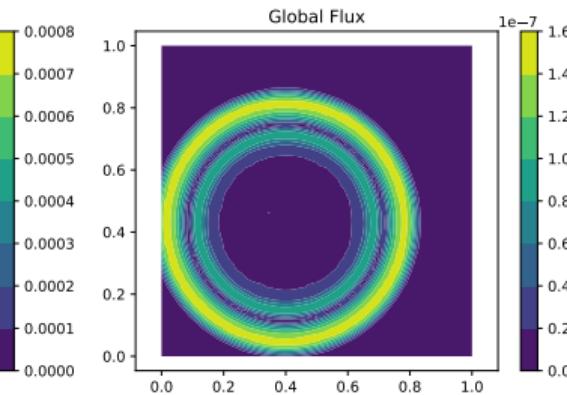
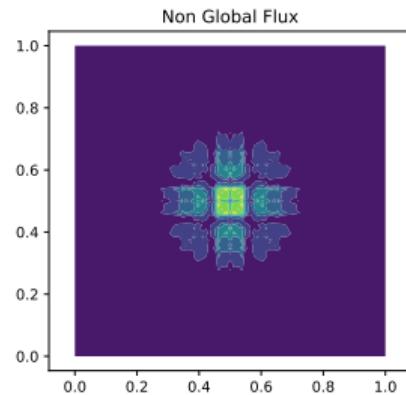
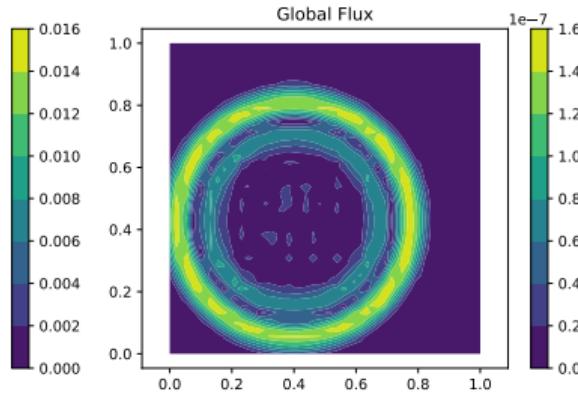
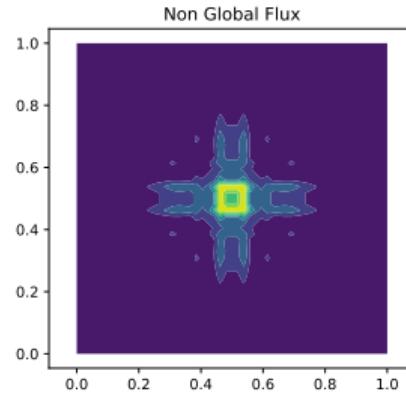
Figure: Coriolis vortex: convergence of L^2 error of u with respect to the number of elements in x

Vortex with Coriolis



Perturbation($\varepsilon = 10^{-2}$) test.
Plot of $\|\underline{u}_{eq} - \underline{u}_p\|$, with \underline{u}_{eq} the analytical equilibrium.
Top \mathbb{P}^3 with 13 cells, bottom \mathbb{P}^3 with 26 cells.

Vortex with Coriolis



Perturbation($\varepsilon = 10^{-6}$) test.
Plot of $\|\underline{u}_{eq} - \underline{u}_p\|$, with \underline{u}_{eq} the analytical equilibrium.
Top \mathbb{P}^3 with 13 cells, bottom \mathbb{P}^3 with 26 cells.

Source term

Consider the source equations

$$\begin{cases} \partial_t \underline{u} + \nabla p = 0, \\ \partial_t p + \nabla \cdot \underline{u} = s, \end{cases} \quad (4)$$

where an equilibrium solution can be found as

$$\begin{cases} p(x, y) \equiv p_0 \in \mathbb{R}, \\ \underline{u}(x, y) = \nabla^\perp \phi_1(x, y) + \nabla \phi_2(x, y), \\ s(x, y) = \Delta \phi_2(x, y), \end{cases} \quad (5)$$

for ϕ_1, ϕ_2 smooth enough. The first term of the velocity, i.e., $\nabla^\perp \phi_1(x, y)$ is analogous to the vortexes defined in (12) and it is divergence-free, while the second term and the source terms balance each other. We will consider the smooth steady vortex (12) for the first part of \underline{u} , while we will use $\phi_2(x, y) := \frac{1}{100} e^{-100||\underline{x} - \underline{x}_0||_2^2}$, with $\underline{x}_0 = (0.65, 0.39)^T$.

Source term

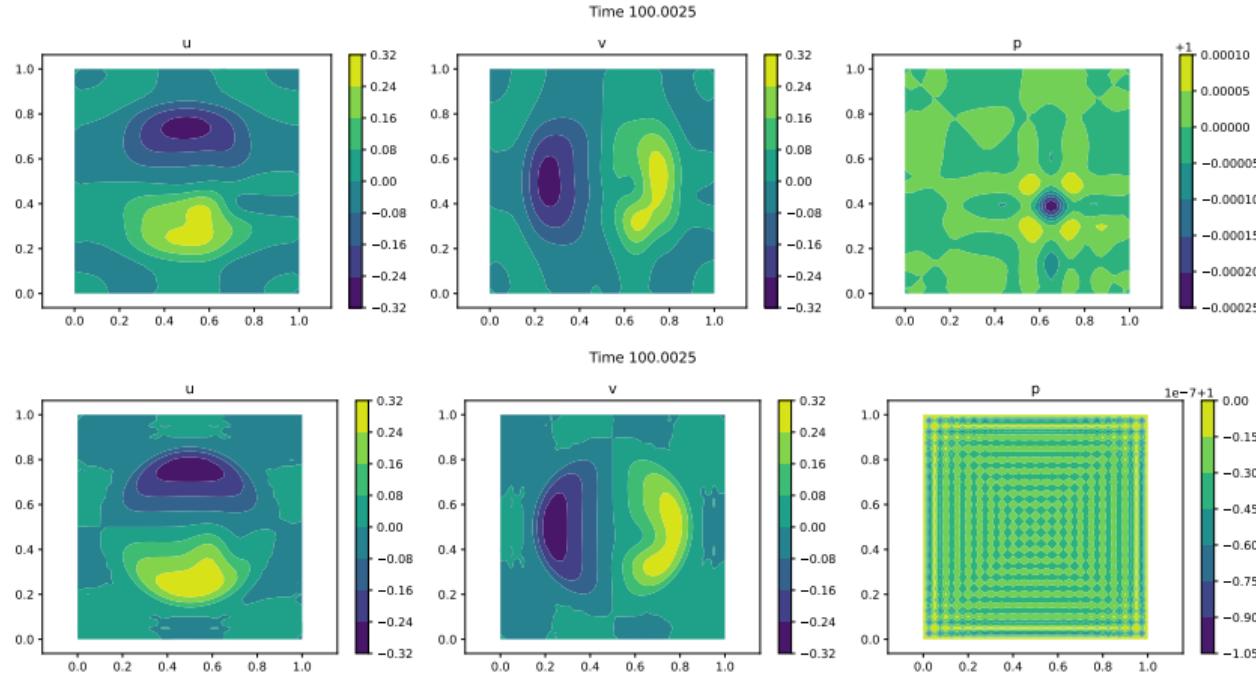


Figure: Simulation of vortex with source term at time $T = 100$ with \mathbb{P}^1 elements and 40×40 cells. SUPG scheme (top) and SUPG-GF scheme (bottom)

Source term

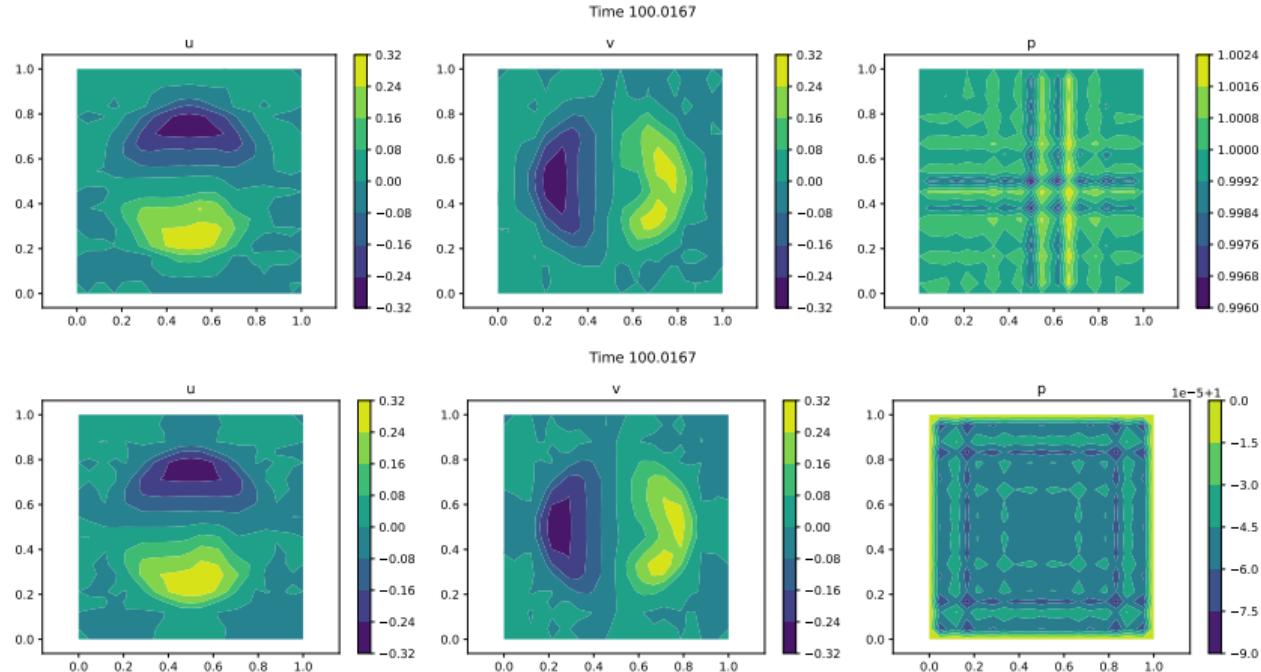
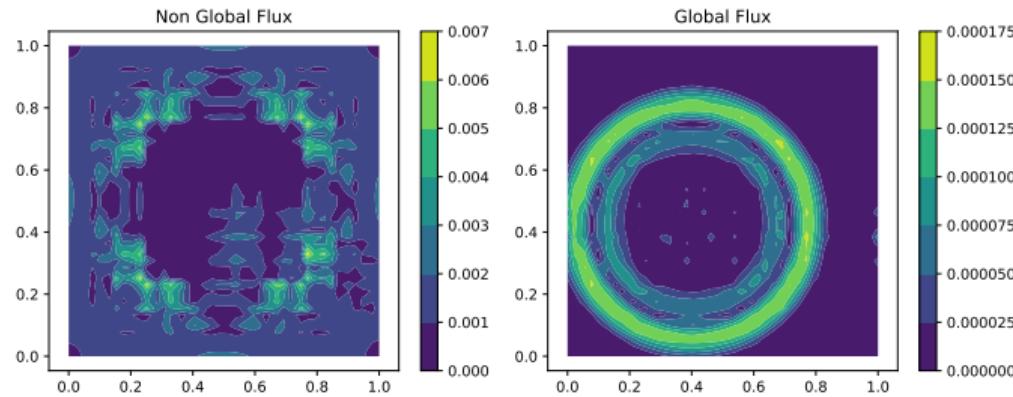
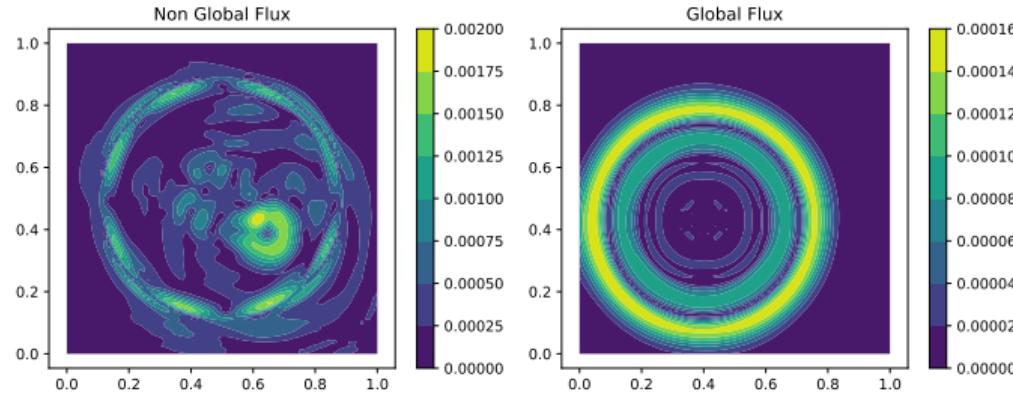


Figure: Simulation of vortex with source term at time $T = 100$ with P^3 elements and 6×6 cells. SUPG scheme (top) and SUPG-GF scheme (bottom)

Source term



Vortex with Source

- Perturbation($\varepsilon = 10^{-3}$) test with source term.
- Plot of $\|\underline{u}_{eq} - \underline{u}_p\|$
- Top \mathbb{P}^1 with 80×80 cells
- Bottom \mathbb{P}^3 with 13 cells

Stommel Gyre

$$\partial_t p = - \operatorname{div} \mathbf{u}$$

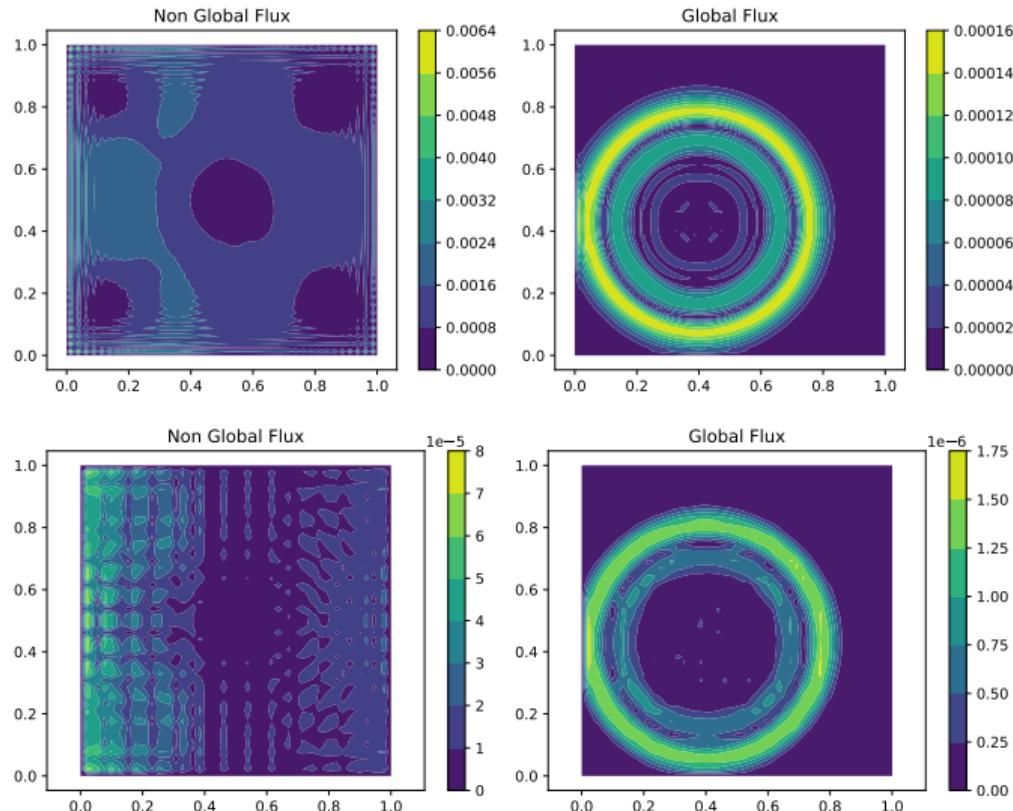
$$\partial_t \mathbf{u} = - \operatorname{grad} p + \phi \mathbf{u}^\perp - R \mathbf{u} + \boldsymbol{\tau}$$

Parameters

- $-R\mathbf{u}$ friction
- $\boldsymbol{\tau}$ wind forcing
- linearized shallow water equations, with a reference depth $h_0 = 1$, and with the gravity acceleration $g = 1$
- R is constant
- $\boldsymbol{\tau} = (-F \cos(\pi y/b), 0)$
- F constant
- well known steady solution due to Henry Stommel^a

^aH. Stommel, The westwards intensification of wind-driven ocean currents, Trans.Amer.Geophys.Union 29(2), 1948

Stommel Gyre



Perturbation test for SG

- Plot of $\|\underline{u}_{eq} - \underline{u}_p\|$
- Top \mathbb{P}^1 with 80×80 cells
 $\varepsilon = 10^{-3}$
- Bottom \mathbb{P}^3 with 13 cells
 $\varepsilon = 10^{-5}$

2D Riemann Problem

- Center $\underline{x}_0 = (0.5, 0.5)$
- Domain $\Omega = [0, 1]^2$
- ICs

$$u(\underline{x}) = \begin{cases} 1, & \text{if } x > 0.5 \text{ and } y > 0.5, \\ 0, & \text{else,} \end{cases} \quad v(\underline{x}) = 0, \quad p(\underline{x}) = 0.$$

- The perpendicular component v has a logarithmic singularity in the center of the RP for all $t > 0$:

$$v(x, y, t) = \frac{1}{2\pi} \mathcal{L} \left(\frac{\sqrt{(x - x_0)^2 + (y - y_0)^2}}{ct} \right),$$

$$\mathcal{L}(s) := \log \left(\frac{1 + \sqrt{1 - s^2}}{s} \right) = -\log \left(\frac{s}{2} \right) - \frac{s^2}{4} + \mathcal{O}(s^4).$$

2D Riemann Problem

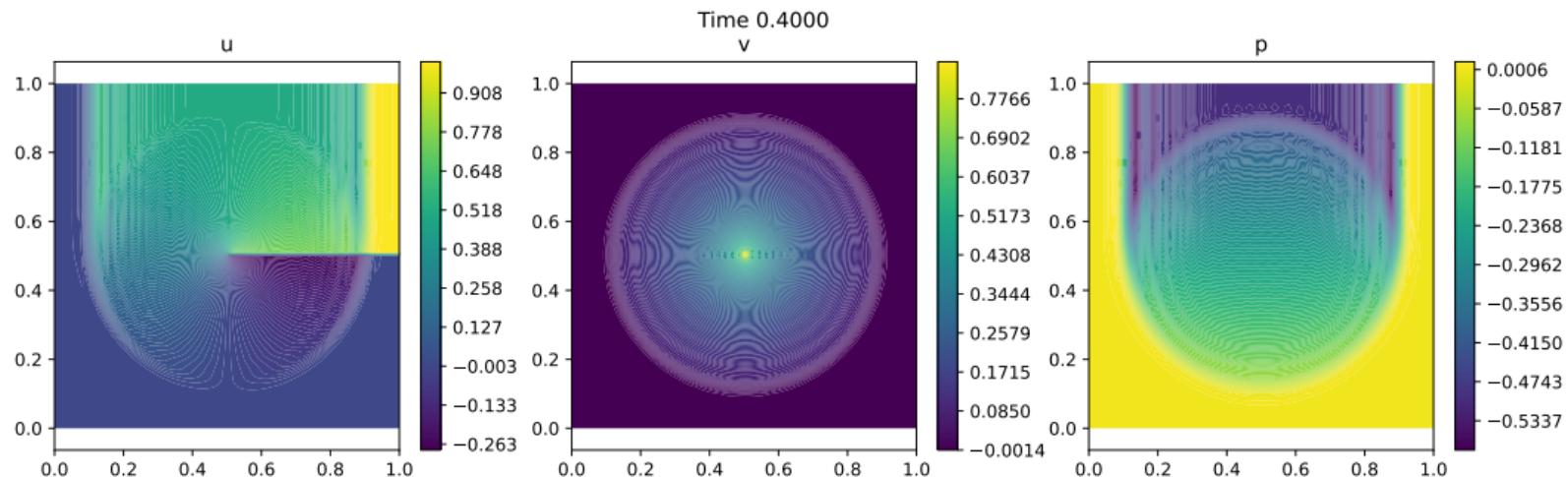


Figure: Riemann Problem. Simulation at time $T = 0.4$ with \mathbb{P}^2 elements and 50×50 cells with SUPG–GF scheme

2D Riemann Problem

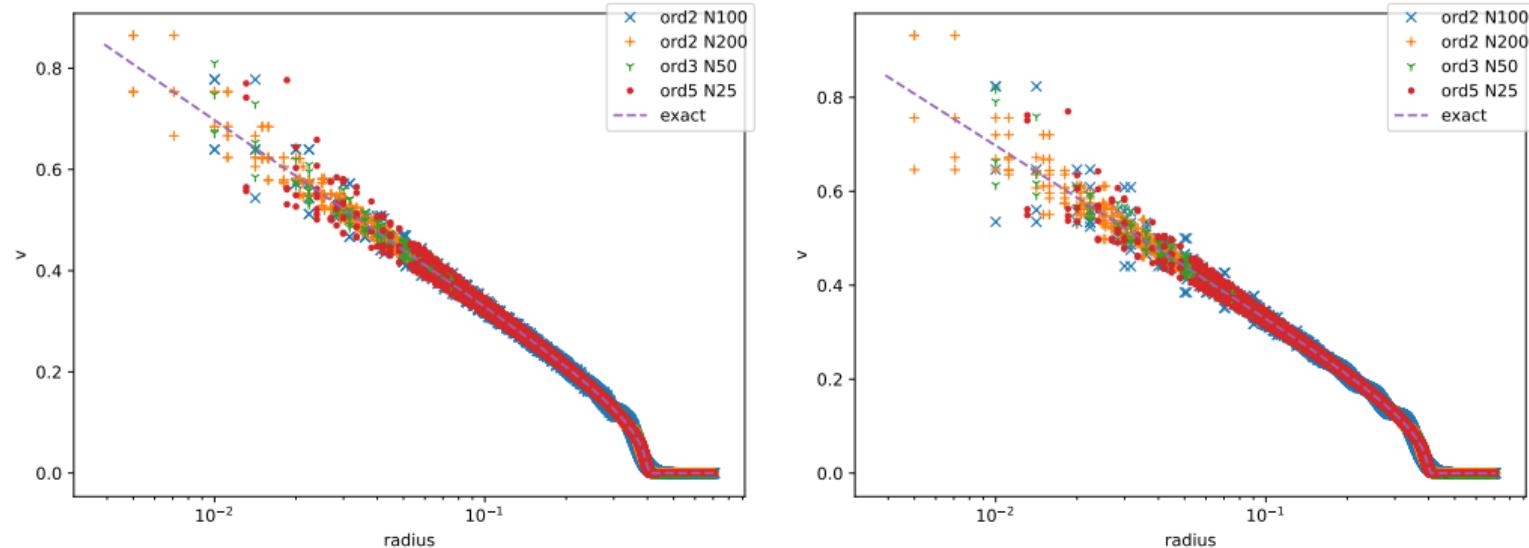


Figure: Riemann Problem. Distribution of the solution v for different elements and meshes. Left SUPG scheme, right SUPG-GF scheme

Analytical involution

$$\nabla \times \partial_t \underline{v} = \nabla \times (\nabla p) = 0$$

Analytical involution

$$\nabla \times \partial_t \underline{v} = \nabla \times (\nabla p) = 0$$

Discrete involution

- If linear method of line
 $\partial_t Q = F(Q) \implies Q^{n+1} = M(F, Q^n)$
- Operator E such that $E \cdot (M(F, Q)) = M(E \cdot F, Q) = 0$ for all Q

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2D SUPG \mathbb{Q}^1 involution operator

$$E := \begin{pmatrix} D_x \left((D_y)^2 M_x - \alpha^2 \Delta y D_y^y \left(\Delta y D_y^y M_x + \Delta x D_x^x M_y \right) \right) \\ D_y \left(- (D_x)^2 M_y + \alpha^2 \Delta x D_x^x \left(\Delta y D_y^y M_x + \Delta x D_x^x M_y \right) \right) \\ \alpha \left(-\Delta x D_x^x (D_y)^2 M_x + (D_x)^2 \Delta y D_y^y M_y \right) \end{pmatrix}.$$

Analytical involution

$$\nabla \times \partial_t \underline{v} = \nabla \times (\nabla p) = 0$$

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2D SUPG \mathbb{Q}^p involution operator

Not feasible