

# Toroidal and Simple Orbifold Compactification

PHY 250

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# Analytic and Anti-Analytic

Analytic function is complex differentiable:

$$f(x, y) = u(x, y) + iv(x, y)$$

$$z = x + iy$$

$$dz = dx + idy$$

Cauchy-Reimann conditions are:

$$\frac{du}{dx} = \frac{dv}{dy}$$

$$\frac{dv}{dx} = -\frac{du}{dy}$$

Anti-analytic function has  $\frac{df}{dz} = 0$

Using

$$\frac{df}{dz} = \frac{df}{dx} \frac{dx}{dz} + \frac{df}{dy} \frac{dy}{dz}$$

The Anti-Analytic version of the Cauchy-Reimann conditions are:

$$\frac{du}{dx} = -\frac{dv}{dy}$$

$$\frac{dv}{dx} = \frac{du}{dy}$$

Hence the name Anti-Analytic!

# Generalized Toroidal Compactification

On 5/26 we did simple compactification to  $S^1 \times M^{25}$

Generalized compactification:  $d = 26 - k =$  number of non-compact dimensions

$$X^m \cong X^m + 2\pi R, d \leq m \leq 25$$

Space-time is  $M^d \times T^k$ , a k-Torus

Geometry of Torus depends on internal metric  $G_{mn}$

$k > 1$  so antisymmetric tensor has scalar components  $B_{mn}$ , Kaluza-Klein gauge bosons  $A_\mu^m$  and antisymmetric gauge bosons  $B_{m\mu}$

The Graviton-Dilaton action is:

$$S_{G-D} = \frac{(2\pi R)^k}{2k_0^2} \int d^d x \sqrt{-G_d} e^{-2\Phi_d} \left[ R_d + 4\partial_\mu \Phi_d \partial^\mu \Phi_d - \frac{1}{4} G^{mn} G^{pq} \{ \partial_\mu G_{mp} \partial^\mu G_{nq} + \partial_\mu B_{mp} \partial^\mu B_{nq} \} - \frac{1}{4} G_{mn} F_{\mu\nu}^m F^{\mu\nu n} - \frac{1}{4} G^{mn} H_{m\mu\nu} H_n^{\mu\nu} - \frac{1}{12} H_{\mu\nu\lambda} H^{\mu\nu\lambda} \right]$$

Where  $k_0$  is a normalized field constant rescaled by redefinition of  $\Phi$ ,  $\Phi_d = \Phi - \frac{1}{4} \ln(\det(G_{mn}))$

$H_{m\mu\nu}$ =generalized field strength,  $F$ =Faraday tensor,  $R_d$  comes from  $G_{\mu\nu}$

# Toroidal Compactification, cont.

Now we have a contribution from the antisymmetric tensor background  $B_{mn}$  which contributes to the world-sheet Lagrangian:

$$B_{mn} \partial_a (\epsilon^{ab} X^m \partial_b X^n)$$

Since this is a total derivative for  $B_{mn}$  it has no effect locally, so the world-sheet is still a Conformal Field Theory. Focusing on the zero-mode contribution, and inserting into the world-sheet action:

$$X^m(\sigma^1, \sigma^2) = x^m(\sigma^2) + w^m R \sigma^1$$

Yields (a dot indicates derivative with respect to world-sheet time  $\sigma^2$ )

$$L = \frac{1}{2\alpha'} G_{mn} (\dot{x}^m \dot{x}^n + w^m w^n R^2) + \frac{i}{\alpha'} B_{mn} \dot{x}^m w^n R$$

$$p_m = -\frac{\partial L}{\partial v^m} = \frac{1}{\alpha'} (G_{mn} v^n - B_{mn} w^n R)$$

Canonical momenta

Where  $v^m = i\dot{x}^m$  because we are using Euclidean signature metric

# Results of Toroidal Compactification

Using Minkowski time,

$$\begin{aligned}v^m &= \partial_0 x^m \\ p_m &= \frac{n_m}{R} \\ v_m &= \alpha' \frac{n_m}{R} + B_{mn} w^n R\end{aligned}$$

The zero-mode contribution to the world-sheet Hamiltonian is:

$$\frac{1}{2\alpha'} G_{mn} (v^m v^n + w^m w^n R^2)$$

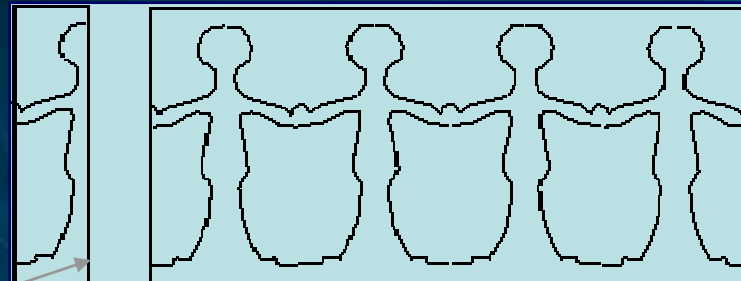
And the closed string mass (sum condition) is:

$$\begin{aligned}m^2 &= \frac{1}{2\alpha'^2} G_{mn} (v_L^m v_L^n + v_R^m v_R^n) + \frac{2}{\alpha'} (N + \tilde{N} - 2) \\ v_{L,R}^m &= v^m \pm w^m R\end{aligned}$$

And the level matching condition becomes

$$0 = G_{mn} (v_L^m v_L^n - v_R^m v_R^n) + 4\alpha' (N - \tilde{N}) = 4\alpha' (n_m w^m + N - \tilde{N})$$

# What is an Orbifold?



Fixed Line

## Mathematically

- A generalization of a manifold that allows the presence of the points whose neighborhood is diffeomorphic to a coset of  $R^n$  such as  $R^n/\Gamma$  (where  $\Gamma$  is a finite group). J.H. Conway defines 17 types

## Physically

- A coset  $M/G$  where  $M$  is a manifold (or theory) and  $G$  is a group of its isometries (or symmetries) - not necessarily all of them. In string theory, these symmetries do not have to have a geometric interpretation. [1]



# Example Orbifolds

Instead of periodic identification of  $X^m$ , consider points under reflection:

$$X^{25} \cong -X^{25}$$

Or more generally:

$$X^m \cong -X^m, 26-k \leq m \leq 25$$

There is a space of fixed points at  $X^{26-k} = \dots = X^{25} = 0$ .

The noncompact space identified in this manner is the quotient space  $\mathbb{R}^k/\mathbb{Z}_2$

We can also use the same toroidal compactification as previously to form  $T^k/\mathbb{Z}_2$

For the noncompact case, there is 1 fixed point identified by:

$$X^{25} \cong X^{25} + 2\pi R$$

For the compact case, there are  $2^k$  fixed points with each  $X^m = 0$  or  $\pi R$

These singular spaces are orbifolds.

An orbifold is thus a generic manifold defined as a quotient space  $M/G$  where  $M$  is a manifold (or theory) and  $G$  is the (not necessarily complete) group of isometries/symmetries of the manifold/theory. Intuitively, it can be thought of as a mapping rotated about a point, line, or surface (i.e. manifold), where these points become singular points defining a "twisted sector":

$$X^{25}(\sigma^1 + 2\pi) = -X^{25}(\sigma^1)$$

# Simple Orbifold Compactification

For the simple, compact 1-d orbifold  $S^1/Z_2$ , the untwisted sector produces the spectrum:

$$|N, \tilde{N}; k^\mu, n, w\rangle \rightarrow (-1)^{N^{25} + \tilde{N}^{25}} |N, \tilde{N}; k^\mu, -n, -w\rangle$$

So this reverses compact winding and momentum. To form linear combinations invariant under this, the states that are massless at generic  $R$  have  $n=w=0$ , so the number of 25-excitations is even. The spacetime graviton, antisymmetric tensor, dilaton, and tachyon survive. However, the Kaluza-Klein gauge bosons are no longer in the spectrum.

In the twisted sector:

$$X^{25}(z, \bar{z}) = i \left( \frac{\alpha'}{2} \right)^{\frac{1}{2}} \sum_{m=-\infty}^{\infty} \frac{1}{m + \frac{1}{2}} \left( \frac{\alpha_{m+\frac{1}{2}}^{25}}{z^{m+\frac{1}{2}}} + \frac{\tilde{\alpha}_{m+\frac{1}{2}}^{25}}{\bar{z}^{m+\frac{1}{2}}} \right)$$

This antiperiodicity forbids any center-of-mass coordinate or momentum, so the string cannot move away from the  $X^{25}=0$  fixed point. A similar mode expansion occurs for the fixed point:

$$X^{25}(\sigma^1 + 2\pi) = 2\pi R - X^{25}(\sigma^1)$$

The expression is the same, with an additional constant term of  $\pi R$ . The mass-shell and level-matching conditions for the twisted sector is:

$$m^2 = \frac{4}{\alpha} \left( N - \frac{15}{16} \right), N = \tilde{N}$$



# Conclusion

In general, this twisting can be thought of as gauging the discrete group  $H$ . Under T-duality, at specific radii, a toroidal theory at  $R=2\alpha'^{1/2}$  can be shown to be equivalent to an orbifold theory  $R=\alpha'^{1/2}$

The Euler characteristic for an orbifold can be defined as:

$$X = V - E + F$$

Where  $V$ =vertices,  $E$  = edges, and  $F$  = faces

It turns out the the number of generations in a particle theory is equal to  $X/2$

In general, we want to compactify to  $M^4 \times V^6$  where  $V$  is the Calabi-Yau space, which is defined as a manifold with nonvanishing harmonic spinors.

# References

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<http://www.physicsdaily.com/physics/Orbifold>
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