

Lecture 8 outline: Linear phase, Group delay, and simple filter design

Reading: P&M section 5.4

This lecture assumes we have already:

- Covered magnitude and phase response $H(z)$ of LTI systems
- Reviewed pole/zero descriptions of $H(z)$

After learning this material, students should be able to:

- Calculate group delay
- Understand why linear phase systems are preferable
- Use pole/zero placement to design simple filter types

Skills that will be used in later work include:

- Understanding of group delay / linear phase concepts

Warm –up question:

Assume we have a real-valued input signal $x[n]$, and want a real-valued output signal $y[n]$. What does this imply for pole-zero placement of $H(z)$?

Outline

- 1) Brief review of general types of ideal filters and their magnitude response
- 2) Ideally, filters would also have *linear phase*
 - a. Systems with pure delay
 - b. Symmetric system responses
 - c. Concept of generalized linear phase
- 3) Group delay: the time delay of different frequencies moving through system
 - a. Definition
 - b. Some physical insight
- 4) Filter design using pole/zero placement
 - a. Basic idea: put poles near frequencies where large response is desired, zeros where low response is desired
 - b. Constraints include: real response, stable response, specified magnitude at certain frequencies
- 5) Example: notch filter design

Lecture: linear phase, group delay, simple filter design

1) Brief review - ideal filters



For LTI system $y[n] = h[n] * x[n]$

$$Y(\omega) = \underbrace{H(\omega)} X(\omega)$$

↑ passes some freq, attenuates others

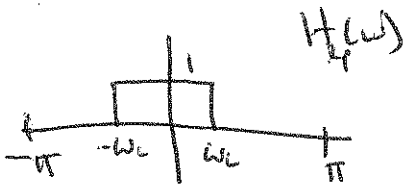
thus

$$|Y(\omega)| = |H(\omega)| |X(\omega)|$$

$$\angle Y(\omega) = \angle H(\omega) + \angle X(\omega)$$

Example ideal filters:

① low pass



$$\rightarrow h_{lp}(n) = \frac{\sin \omega_c n}{\pi n}$$



* note this is not physically realizable - ∞ long $h[n]$

② high pass



$$|H_{hp}(\omega)| = 1 - |H_{lp}(\omega)|$$

$$\rightarrow h_{hp}(n) = \delta(n) - \frac{\sin \omega_c n}{\pi n}$$

Q)

Ideally, what would phase be? ^{often} zero or ^{easily} explain / compensated

remember F.T. pair: $\delta(n-nd) \Leftrightarrow e^{-j\omega nd}$

→ ~~the~~ Simple example: delay system

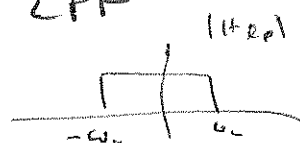
$$h = \delta(n-nd) \rightarrow y(n) = h * x = x(n-nd)$$

$$\text{so } H(\omega) = e^{-j\omega nd}$$

$$|H| = 1, \quad \angle H = -\omega nd \in \text{Linear phase} \quad (\sim \omega)$$

next example: linear phase, ~~the~~ ideal LPF

$$H_{lp}(\omega) = \begin{cases} e^{-j\omega nd} & |\omega| < \omega_c \\ 0 & \text{else} \end{cases}$$



$$\angle H_{lp} = \begin{cases} -\omega nd, & |\omega| < \omega_c \\ 0 & \text{else} \end{cases}$$

again linear

$$h_{lp}(n) = \frac{\sin \omega_c (n-nd)}{\pi (n-nd)} = \text{delayed version of first lpf}$$

Idea: linear phase \rightarrow delay in time.

Otherwise the signal phase is not changed.

[Often, we want to just attenuate magnitude in noise regions; in these cases, linear phase / time delay leaves desired signal untouched]

- Examples of distortion later

③ Group delay.

→ ~~now~~ if we plot $\angle H$ vs. ω , we can see if phase is linear.

→ otherwise, hard to learn much from $\angle H$

→ a more useful way to capture distortion is group delay

$$\tau(\omega) = \text{grad} [H(\omega)] = -\frac{d}{d\omega} \arg[H(\omega)]$$

→ has units of ^{time / delay} ($\frac{1}{\text{Vs}} = \text{s}$)

→ so $1/\tau$ is like velocity (how fast different frequencies move through system)

Some insight (hopefully)

→ For pure delay, we saw $H(\omega) = e^{-j\omega nd}$

$$\angle H = -\omega nd$$

$$-\frac{\partial \angle H}{\partial \omega} = nd \leftarrow \text{everything has delay } nd \forall \omega$$

→ group delay / group velocity is a concept from physics - the rate at which an envelope of the signal (information) propagates

→ example of mechanical dispersive system:

→ ~~like~~ on frozen pond, hit the ice. ~~For~~ Higher-freq sands travel quicker, lower-freq slower so original signal is smeared out.

more general:

$$\angle H(\omega) = \tan^{-1} \left(\frac{\text{Im}[H(\omega)]}{\text{Re}[H(\omega)]} \right)$$

$$\equiv \tan^{-1}(A(\omega))$$

remember (or look up) $\frac{d}{dx} \tan^{-1} x = \frac{1}{1+x^2}$

then use chain rule

$$\tau(\omega) = \frac{1}{1+A^2(\omega)} \frac{dA(\omega)}{d\omega}$$

Other examples of linear phase:

a) pure delay - above

b) symmetric system:

version 1 $h_1[n] = h_1[-n]$

as h_1 is odd, $H_1(\omega)$ is real

$$\angle = 0 \rightarrow \tau = 0$$

version 2: symmetric, about a shifted point

$$h_2[n] = h_1[n - n_0]$$

$$H_2(\omega) = e^{-j\omega n_0} H_1(e^{j\omega})$$

$$\angle H_2 = \angle H_1 - \omega n_0$$

c) Generalized linear phase:

these systems are of form

$$H(\omega) = e^{j\beta} e^{-j\omega\alpha} I(\omega)$$

τ_{real}

here,

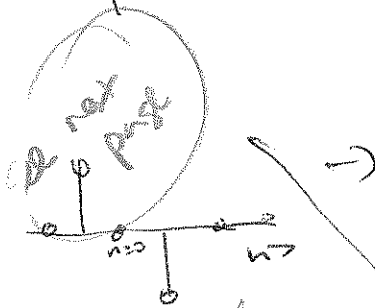
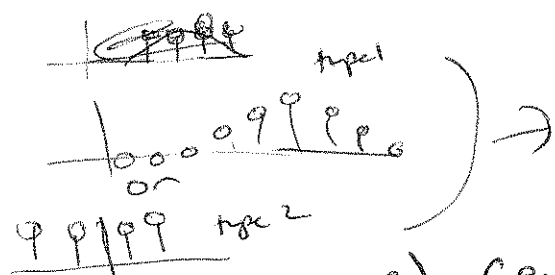
$$\angle H = \beta - \omega\alpha$$

α constant

$$-\frac{\partial \angle H}{\partial \omega} = \alpha \leftarrow \text{constant delay}$$

Some ~~nonlinear~~ filters can achieve generalized linear phase ~~that is not~~ in their pass band - good enough

⊛ Linear phase is easy to do w/ FIR, hard w/ IIR



$$H(\omega) = j e^{-j\omega} \sin(\omega)$$

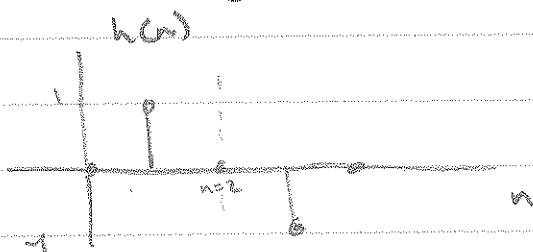
$$= e^{j\pi/2} e^{-j\omega} \sin(\omega)$$

leave 7:

$$H(\omega) = 2j e^{j\omega\pi/2} \sin(\omega\pi/2)$$

$$= 2 e^{j\pi/2} e^{j\omega\pi/2} \sin(\omega\pi/2)$$

example generalized linear phase system



anti-symmetric
about $n=2$

$$H(z) = (1)z^{-1} - (1)z^{-3}$$

$$= z^{-2}(z + z^{-1})$$

$$H(\omega) = H(z) \Big|_{z=e^{j\omega}} = e^{j2\omega} (e^{j\omega} - e^{-j\omega})$$

$$= 2j e^{j2\omega} \left(\frac{e^{j\omega} - e^{-j\omega}}{2j} \right)$$

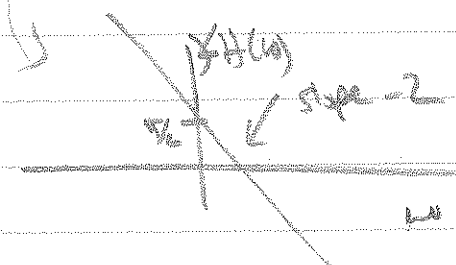
$$= 2j e^{j2\omega} \sin(\omega)$$

but $j = e^{j\pi/2}$, so $H(\omega) = 2 e^{j(\pi/2 - 2\omega)} \sin(\omega)$

$\angle H(\omega) = \pi/2 - 2\omega$

\nearrow
real
Ruler

$\angle(\omega) = -d/d\omega \angle H(\omega) = 2 //$ non-dispersive



3.30: show if $z=re^{j\theta}$ is zero
so. is $z = \frac{1}{r} e^{j\theta}$

1.4 3.30

ZERO LOCATIONS OF LINEAR-PHASE FILTERS

Sec 4.4.4

in Mitra

The zeros of the transfer function $H(z)$ of a linear-phase filter lie in specific configurations.

We can write the symmetry condition

$$h(n) = h(N-1-n)$$

in the Z domain. Taking the Z -transform of both sides gives

$$H(z) = z^{-(N-1)} H(1/z). \quad (25)$$

Recall that we are assuming that $h(n)$ is real-valued. If z_0 is a zero of $H(z)$,

$$H(z_0) = 0,$$

then

$$H(z_0^*) = 0.$$

(Because the roots of a polynomial with real coefficients exist in complex-conjugate pairs.)

Using the symmetry condition (25), it follows that

$$H(z_0) = z_0^{-(N-1)} H(1/z_0) = 0$$

and

$$H(z_0^*) = (z_0^*)^{-(N-1)} H(1/z_0^*) = 0$$

or

$$H(1/z_0) = H(1/z_0^*) = 0.$$

If z_0 is a zero of a (real-valued) linear-phase filter, then so are

$$z_0^*, 1/z_0, 1/z_0^*.$$

$h(n) = a\delta(n) + b\delta(n-1) + c\delta(n-2)$
 ex $H(z) = az + bz^{-1} + cz^{-2}$
 $H(z_0) = az_0 + bz_0^{-1} + cz_0^{-2}$
 $H(1/z_0) = az_0^{-1} + bz_0 + cz_0^2$
 $H(z_0^*) = az_0^* + bz_0^{*-1} + cz_0^{*-2}$
 $H(1/z_0^*) = az_0 + bz_0^{-1} + cz_0^{-2}$
 $N=3$

more fact

TYPE I: ODD-LENGTH SYMMETRIC

The frequency response of a length $N = 5$ FIR Type I filter can be written as follows.

$$H^f(\omega) = h_0 + h_1 e^{-j\omega} + h_2 e^{-2j\omega} + h_1 e^{-3j\omega} + h_0 e^{-4j\omega} \quad (1)$$

$$= e^{-2j\omega} (h_0 e^{2j\omega} + h_1 e^{j\omega} + h_2 + h_1 e^{-j\omega} + h_0 e^{-2j\omega}) \quad (2)$$

$$= e^{-2j\omega} (h_0 (e^{2j\omega} + e^{-2j\omega}) + h_1 (e^{j\omega} + e^{-j\omega}) + h_2) \quad (3)$$

$$= e^{-2j\omega} (2h_0 \cos(2\omega) + 2h_1 \cos(\omega) + h_2) \quad (4)$$

$$= A(\omega) e^{j\theta(\omega)} \quad (5)$$

where

$$\theta(\omega) = -2\omega, \quad A(\omega) = 2h_0 \cos(2\omega) + 2h_1 \cos(\omega) + h_2.$$

Note that $A(\omega)$ is real-valued and can be both positive and negative.

In general, for a Type I FIR filters of length N :

$$H^f(\omega) = A(\omega) e^{j\theta(\omega)}$$

$$A(\omega) = h(M) + 2 \sum_{n=0}^{M-1} h(n) \cos((M-n)\omega).$$

$$\theta(\omega) = -M\omega$$

$$M = \frac{N-1}{2}.$$