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https://github.com/julianmak/academic-notes

The repository principally contains the compiled products rather than the source for size reasons.

- Associated Python code (as Jupyter notebooks mostly) will be held on the same repository. The source data however might be big, so I am going to be naughty and possibly just refer you to where you might get the data if that is the case (e.g. JRA-55 data). I know I should make properly reproducible binders etc., but I didn't...
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OCES 2003 : Descriptive Physical Oceanography

(a.k.a. physical oceanography by drawing pictures)

Lecture 16: Dynamics 2 (waves and dynamic mechanisms)

Outlook of the next few lectures

Dynamics important, next few lectures on

- waves (this Lec. + 16, 18) and instabilities (Lec. 17)
 - → because waves are easier to talk about without maths...

Highlight gross features (i.e. those that can be drawn...)

- ▶ how to describe waves (Lec. 15)
- types of waves (Lec. 16)
 - \rightarrow consequence + leading to instabilities
- ▶ instabilities (Lec. 17)
 - → parcel-type (mechanistic) arguments for instability
- ▶ tides (particularly as internal gravity waves) (Lec. 18)

Outline

- gravity waves
 - \rightarrow gravity/buoyancy as restoring mechanism (\sqrt{gH})
- ▶ inertial waves
 - \rightarrow Coriolis as restoring mechanism (f)
 - \rightarrow e.g. Rossby waves, Kelvin waves
- ▶ inertial-gravity + internal waves $(\sqrt{gH} \text{ or } N, \text{ and } f)$
 - \rightarrow extra depth dimension to deal with
 - \rightarrow Brunt–Väisälä or buoyancy frequency N
- propagation mechanism (Rossby wave example)
 - → kinematic argument with **vorticity**

Key terms: buoyancy frequency, gravity waves, inertial waves, Rossby waves, Kelvin waves, vorticity inversion



Recap: what goes down must come up

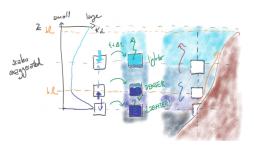


Figure: Schematic of the diffusive upwelling.

▶ diapycnal mixing contribute upwelling, strongest in boundary layers
 → broad diffusive boundary intensified upwelling

what causes the boundary intensification of κ_d ? dynamics!

- at the surface, lots of things... (convection, waves, Langmuir turbulence etc.)
- ▶ at the bottom, probably tidal conversion (Lec. 18) \rightarrow internal gravity waves (Lec. 16) \rightarrow shear instabilities (Lec. 17)



Recap: waves and dispersion relation

- waves are ubiquitous physical features
 - \rightarrow depends on physics
- wave described by the dispersion relation $\omega = \mathcal{F}(k)$
 - \rightarrow physics dictate the form of \mathcal{F}









Figure: Examples of systems supporting waves. All figures from Wikipedia except the cello one



Figure: Gravity waves with signal at the sea surface (as darker and lighter bands). Taken at HKUST.

- ► (linear) waves can interfere with each other
 - → constructive or destructive
 - → interference can lead to steepening and breaking ("becoming"

Recap: wave propagation

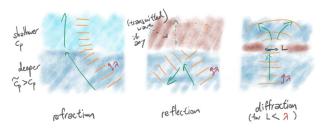


Figure: Schematic of refraction, reflection (and transmission), and diffraction. See Lec. 15.

b phase speed (in a direction) and group velocity as (note $\omega = \mathcal{F}(k)$)

$$c_{p,x} = \frac{\omega}{k}, \qquad c_{g,x} = \frac{\partial \omega}{\partial k}$$

- → individual wave vs. wavepacket behaviour
- \rightarrow contribute to wave phenomenon (e.g. refraction from

$$c_p = c_p(x)$$



Wave steepening and breaking



Figure: Schematic of mixing by (irreversible) wave breaking, with contours reconnecting leading to e.g. diapycnal mixing.

- growing waves by instability
 - \rightarrow convective and/or shear (see Lec. 17)
 - → mixing of material **across** isopycnals after reconnection, leading to diapycnal mixing
- feedback onto MOC (see Lec. 14)

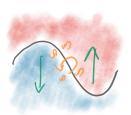


Figure: Velocity shear from waves can lead to mixing.

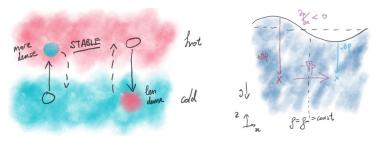


Figure: Gravity as restoring force. Pictures adapted from ones used in Lec. 7.

- deviation from resting isopycnal experiences restoring force from gravity (buoyancy)
 - → left case: internal isopycnal (as an **isotherm**)
 - \rightarrow right case: sea surface is the isopycnal
- ▶ weak damping, restoring force, overshoots ⇒ oscillatory motion (up and down in this case)

For simplicity, consider a **homogeneous** (i.e. $\rho = \text{const}$) layer of fluid (cf. Lec. 11 + 12, right case in previous Figure) until we get to internal waves

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In this instance, dispersion relation for gravity waves is given by (without derivation)

$$\omega^2 = gk \tanh(kH)$$
 \Rightarrow $\omega = \pm \sqrt{gk \tanh(kH)}$

- ► H is water depth, and tanh = hyperbolic tangent, goes from -1 to 1
 - \rightarrow note symmetry in both directions (the \pm sign)
 - ($\omega \leq 0$ cases are just **shifts** in the **phase**)

$$\omega = \pm \sqrt{gk \tanh(kH)}$$

▶ for **deep** water waves ($kH \gg 1$) and **shallow** water waves ($kH \ll 1$),

$$\omega_{\mathrm{deep}} = \pm \sqrt{gk}, \qquad \omega_{\mathrm{shallow}} = \pm k\sqrt{gH}$$

- $\rightarrow kH \gg 1$ so $tanh(kH) \rightarrow 1$
- $ightarrow kH \ll 1$ with $anh(kH) pprox kH + O((kH)^3)$ (do a Taylor expansion)
- deep water waves are depth-independent and dispersive
- ▶ shallow water waves are slower in shallow waters $(c_p \sim \sqrt{H})$ and non-dispersive $(c_p = c_g)$



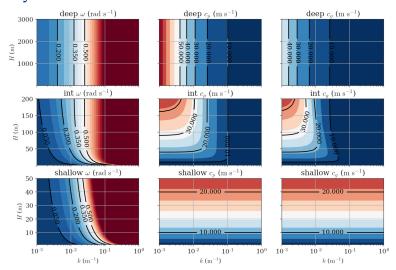
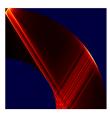


Figure: Water wave ω , c_p and c_g over (k,H) space, with k shown on a log axis. k chosen so wavelengths are roughly between 50 m to 5 km (recall $k=2\pi/\lambda$). Also note the transitions from shallow to intermediate to deep are really to do with $kH\sim H/\lambda$. See waves.ipynb.

Inertial waves

Inertial waves has the Coriolis "force" act as the restoring force

- generic for rotating systems
 - \rightarrow planetary interiors, stars, galactic disks
- mostly arise in context of internal waves
 - → at surface buoyancy effects can dominate
- ► limited in frequency by inertial frequency f_0 (cf. Coriolis parameter)
 - → revisit later when talking about inertia-gravity waves



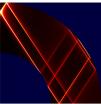


Figure: Inertial wave attractors in a homogeneous planetary interior at different tidal forcings. From Gordon Ogilvie (2009, Mon. Not. Royal. Astro. Soc).

Inertial-Gravity waves

In reality Coriolis and buoyancy effects both contribute

- ▶ large-scale and/or slow \Rightarrow Coriolis important (because Ro \ll 1), classify as inertial waves
 - \rightarrow e.g. Rossby waves
- ightharpoonup small-scale and/or fast \Rightarrow Coriolis unimportant
 - \rightarrow e.g. internal gravity waves
- somewhere in between? Poincaré or inertia-gravity waves

$$\omega = \pm \sqrt{f_0^2 + gH(k_x^2 + k_y^2)}$$

- \rightarrow for $gH(k_x^2 + k_y^2) \gg f_0$, recover gravity waves
- \rightarrow for $gH(k_x^2 + k_y^2) \ll f_0$, recover inertial waves



Rossby deformation radius

$$\omega = \pm \sqrt{f_0^2 + gH(k_x^2 + k_y^2)}$$

boundary between gravity and inertial regimes is roughly

$$L_d = \frac{\sqrt{gH}}{f_0}$$

- the Rossby deformation radius (for shallow water system)
 - \rightarrow roughly also the boundary where geostrophic approximation should hold (see Lec. 8 + 13)
 - \rightarrow estimates in a few slides

introduce a useful quantity

$$N = \sqrt{-\frac{g}{\rho_0} \frac{\partial \rho}{\partial z}}$$

- ► Brunt–Väisälä or buoyancy frequency (units: s⁻¹)
 - $\rightarrow N^2$ normally used
 - \rightarrow note $\partial \rho / \partial z < 0$ for stable stratification, i.e.

$$N^2 > 0$$

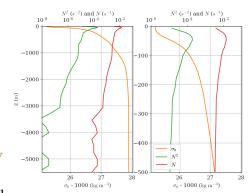


Figure: σ_0 (see Lec. 6) and the associated N^2 and N. $N^2 \ll 1$ means weakly stratified (weak density gradients), whilst $N^2 < 0$ shows unstable stratification (none in this case, but see Lec. 17). See plot-eos.ipynb.

simplistic view (!): $\sqrt{gH} \rightarrow N$

Generally then, internal inertia-gravity waves described by

$$\omega = \pm \sqrt{\frac{f_0^2 k_z^2 + N^2 k_x^2}{k_x^2 + k_z^2}} \approx \pm \sqrt{f_0^2 + \frac{N^2 k_x^2}{k_z^2}} \quad \text{(for } |k_z| \gg |k_x|\text{)}$$

- ► atmosphere and ocean has $N/f_0 = O(10^1 \text{ to } 10^2)$
 - \rightarrow so really we have gravity waves influenced by rotation
 - → refer to them here as internal waves

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- ▶ note that $|f_0| \le |\omega| \le |N|$
 - \rightarrow since $0 \le k_{x,z}^2/(k_x^2 + k_z^2) \le 1$
 - \rightarrow frequency is **much lower** than gravity waves

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internal tides to be seen as internal waves (Lec. 18)



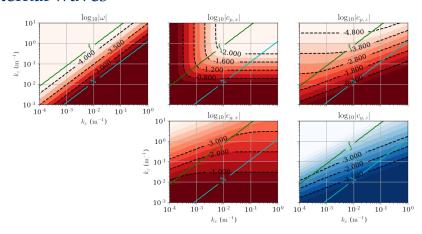


Figure: Inertial-gravity waves (with the $k_z\gg k_x$ approximation) ω , $c_{p,x}$, $c_{p,y}$, $c_{g,x}$ and $c_{p,y}$ as a log-log plot in (k_x,k_z) space, with $f=5\times 10^{-5}$ and $N=3\times 10^{-3}$ (oceanic relevant values). The contours denote the exponent x of $|10^x|$ and the colour shading denotes the sign (more blue = more negative actual values rather than exponents, more red = more positive actual values rather than exponents, since k_x and k_z is chosen to be positive, everything except $c_{g,z}$ is positive. Contours of f and N plotted with an offset plotted to show the boundary beyond which everything is either gravity waves or inertial oscillations. See waves . ipynb.

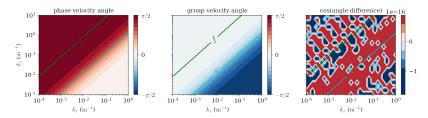


Figure: Inertial-gravity waves (with the $k_z \gg k_x$ approximation) phase velocity c_p angles and group velocity c_g angles (in radians, relative to the horizontal, and note $\pi/2 = 90^{\circ}$). The final panel shows $c_p \cdot c_g = |c_p||c_g|\cos\theta$ (which is zero up to rounding errors). Contours of f and N plotted with an offset plotted as in the previous diagram. See waves . i.pynb.

note that, for inertial-gravity waves (left as a bonus exercise),

$$c_p \cdot c_g = 0$$

ightarrow i.e. phase and group velocities are **perpendicular** to each other (see Lec. 4)



Deformation radius

 boundary given again by the Rossby deformation radius (for the continuously stratified case)

$$L_d = \frac{NH}{f}$$

- \rightarrow $L_{d,atmos} = O(1000 \text{ km})$, scale of cyclones and anti-cyclones, i.e. weather systems form (synoptic structures)
- \rightarrow $L_{d,ocean} = O(50 \text{ km})$, scale of ocean eddies
- ▶ latitude (through *f*) and *H* dependent
 - \rightarrow smaller L_d for **high** latitudes and **shallow** regions
 - \rightarrow consequence for **geostrophic approximation?** (e.g. shelves and coasts)

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 - \rightarrow smaller L_d for **high** latitudes and **shallow** regions
 - \rightarrow consequence for **geostrophic approximation?** (e.g. shelves and coasts)
- internal L_d defined analogously (normally smaller than above)



Kelvin waves (more on this in Lec. 18)

A type of boundary wave

- need f and a boundary
 - → could be land (coastal Kelvin waves) (see Lec. 18)
 - \rightarrow could be a wave guide (e.g. equator where f changes sign, equatorial Kelvin waves) (see OCES 4001, El-Niño, QBO etc.)
- needs f but propagates at the gravity wave speed, with

$$\omega = k\sqrt{gH}$$

- → non-dispersive
- → fairly fast (gravity wave speed)

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- \rightarrow fairly fast (gravity wave speed)
- NOTE the lack of ±!

Kelvin waves (more on this in Lec. 18)

 boundary introduces asymmetry in this case: general solution like

$$\eta \sim \mathrm{e}^{\pm f_0 y / \sqrt{gH}} \cos(kx - \omega t)$$

- \rightarrow take $y \le 0$ to be **boundary**, if $f_0 > 0$ (NH), need minus sign, and vice-versa
- → wave propagates cyclonically (same sign as f)
- ▶ taking $f_0 > 0$ (NH),

$$\eta \sim e^{-y/L_d} \cos(kx - \omega t),$$

so decay over the $L_d = \sqrt{gH}/f_0$

Rossby waves (more on this later)

A (particularly important) type of **inertial** wave

- requires a gradient in background vorticity
 - $\rightarrow \partial f/\partial y = \beta$ (planetary case)
 - ightarrow background flow $-\partial U/\partial y \sim \nabla \times u$ (see later and Lec. 17)
- dispersion relation given by (on β -plane)

$$\omega = -\frac{\beta k_x}{k_x^2 + k_y^2}$$

 \rightarrow note that Rossby waves propagate to the **west** (more generally, **retrograde** or against the mean flow) Since

$$c_{p,x} = \frac{\omega}{k_x} = -\frac{\beta}{k_x^2 + k_y^2} < 0,$$

and **long** waves ($k_x \ll 1$) are **fast**(er)



Rossby waves (more on this later)

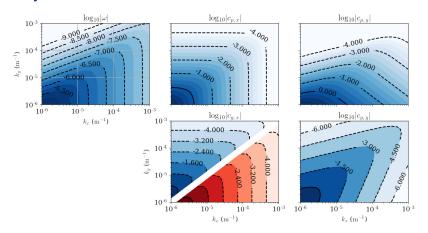


Figure: Rossby waves ω , $c_{p,x}$, $c_{p,y}$, $c_{g,x}$ and $c_{p,y}$ as a log-log plot in (k_x, k_y) space, with magnitude also as logs. The contours denote the exponent x of $|10^x|$ and the colour shading denotes the sign (more blue = more negative actual values, more red = more positive actual values); since k_x and k_y is chosen to be positive, everything except $c_{g,x}$ is negative. Choice of k_x and k_y correspond to wavelengths roughly between 6 km to 6000 km (Rossby waves are usually seen as planetary-scale waves). See waves.ipynb.

Propagation mechanism: Rossby waves

Rossby waves propagate west-ward (or, more generally, retrograde)

$$c_{p,x} = \frac{\omega}{k_x} = -\frac{\beta}{k_x^2 + k_y^2} < 0$$
why?

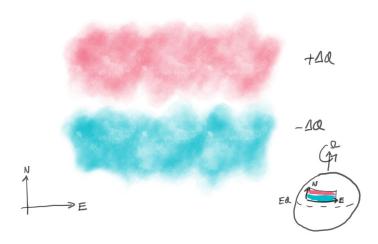
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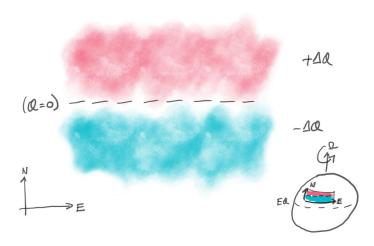
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 why?

Key bits to the pictorial/parcel (cf. Lec 5 for temperature) argument:

- ► the initial wave conserves and carries vorticity (spini-ness, recall Lec. 4, 11, 12) into the external environments
 - → these are now vorticity anomalies
- vorticity anomalies induces a velocity/flow (because spini-ness)
- ▶ induced flow seen to self-advect the wave and move it to the West (retrograde in the general case)





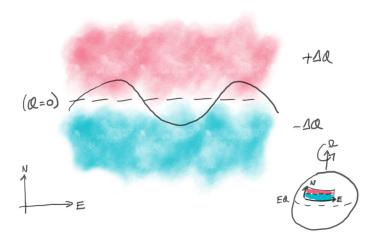
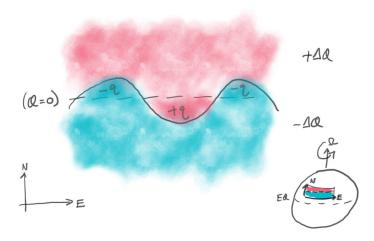
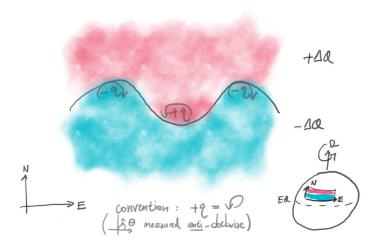
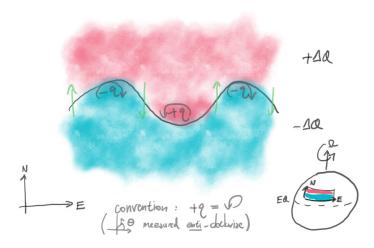


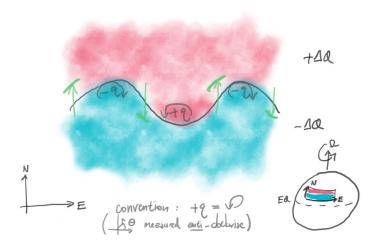
Figure: Rossby wave propagation schematic.

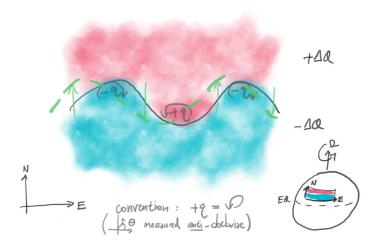




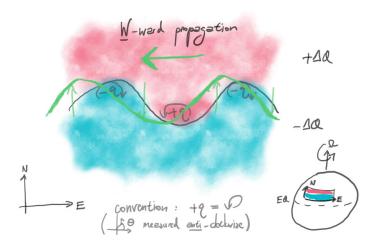














Summary

- gravity waves (gravity/buoyancy)
- inertial waves (Coriolis)
- inertial-gravity waves (general)
 - \rightarrow internal waves have

$$|f| \le |\omega| \le |N|$$

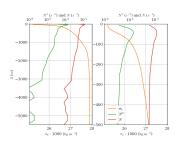


Figure: σ_0 (see Lec. 6) and the associated N^2 and N. See plot_eos.ipvnb.

► Brunt–Väisälä or buoyancy frequency *N*

$$N = \sqrt{-\frac{g}{\rho_0} \frac{\partial \rho}{\partial z}}$$

 \rightarrow measure of stratification strength (see also Lec. 17)



Summary

- parcel argument for west-ward
 Rossby wave propagation
 - → conservation of vorticity
 - → vorticity anomalies induces flow
 - $\rightarrow self\text{-}advecting$

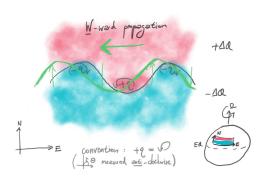


Figure: Rossby wave propagation schematic.

▶ generalisations exist (e.g. internal gravity waves in Harnik et al., 2008, J. Atmos. Sci)

Summary

- parcel argument for west-ward Rossby wave propagation
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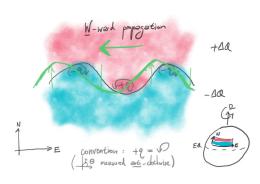


Figure: Rossby wave propagation schematic.

- peneralisations exist (e.g. internal gravity waves in Harnik et al., 2008, J. Atmos. Sci)
- ▶ two such waves interacting? (see Lec. 17)
 - → potential for instabilities

