

Problemen Set 1

1) a) Show $\Gamma_{10}^1 = \Gamma_{01}^1 = \frac{\dot{a}}{a}$ and $\Gamma_{11}^0 = \frac{a\ddot{a}}{c(1-kr^2)}$ for FLRW-metric

$$ds^2 = -dt^2 + a(t)^2 \left(\frac{dr^2}{1-kr^2} + r^2(d\theta^2 + \sin^2\theta d\phi^2) \right)$$

Proof: $\Gamma_{ik}^j = \frac{1}{2} g^{jl} \cdot \left(\frac{\partial g_{il}}{\partial x^k} + \frac{\partial g_{ik}}{\partial x^l} - \frac{\partial g_{kl}}{\partial x^i} \right)$

With $i, j, k, l = t, r, \theta, \phi$

$$\Gamma_{10}^1 = \Gamma_{01}^1 \text{ because of symmetry (1.13)}$$

$$\Gamma_{01}^1 = \frac{1}{2} g^{11} \cdot \left(\frac{\partial g_{01}}{\partial x^1} + \frac{\partial g_{11}}{\partial x^0} - \frac{\partial g_{01}}{\partial x^1} \right)$$

with $g^{ab} = \begin{pmatrix} -c^2 & 0 & 0 & 0 \\ 0 & \frac{1}{1-kr^2} & 0 & 0 \\ 0 & 0 & r^2 & 0 \\ 0 & 0 & 0 & r^2 \sin^2\theta \end{pmatrix} \cdot \begin{pmatrix} dt \\ dr \\ ds \\ d\phi \end{pmatrix} \cdot \begin{pmatrix} dt \\ dr \\ ds \\ d\phi \end{pmatrix}$

$$= \underbrace{\frac{1}{2} g^{10} \cdot \left(\frac{\partial g_{00}}{\partial t} \dots \right)}_0 + \underbrace{\frac{1}{2} g^{11} \cdot \left(\frac{\partial g_{01}}{\partial t} \dots \right)}_0 + \underbrace{\frac{1}{2} g^{12} \cdot \left(\frac{\partial g_{02}}{\partial t} \dots \right)}_0 + \underbrace{\frac{1}{2} g^{13} \cdot \left(\dots \right)}_0$$

$$= \frac{1}{2} \cdot g^{11} \cdot \left(\underbrace{\frac{\partial g_{01}}{\partial r}}_0 + \underbrace{\frac{\partial g_{11}}{\partial t}}_0 - \underbrace{\frac{\partial g_{01}}{\partial r}}_0 \right)$$

$$= \frac{1}{2} \cdot \underbrace{g^{11}}_0 \cdot \left(\frac{\partial}{\partial t} \frac{a(t)}{1-kr^2} \right)$$

$g^{ab} = \frac{1}{g^{ab}}$ weil Diagonale = 0 $\Rightarrow \frac{1}{g^{ab}}$ ist inverses

$$= \frac{1}{2} \cdot \frac{(1 - kr^2)}{c^2} \cdot \frac{2\alpha \cdot \ddot{\alpha}}{(1 - kr^2)}$$

$$= \frac{\ddot{\alpha}}{\alpha} \quad S_{\alpha y}!$$

$$\Gamma_{11}^0 = \frac{1}{2} g^{00} \cdot \left(\frac{\partial g_{10}}{\partial x^1} + \frac{\partial g_{11}}{\partial x^1} - \frac{\partial g_{01}}{\partial x^0} \right)$$

$$= \frac{1}{2} \cdot (g^{00} \dots + g^{01} \dots + g^{02} \dots + g^{03})$$

0

$$= \frac{1}{2} g^{00} \cdot \left(\underbrace{\frac{\partial g_{10}}{\partial x^1} + \frac{\partial g_{11}}{\partial x^1}}_0 - \frac{\partial g_{01}}{\partial x^0} \right)$$

$$= \frac{1}{2} g^{00} \cdot \frac{\partial g_{11}}{\partial t}$$

$$= \frac{1}{2} \underbrace{g^{00}}_{-\frac{1}{c^2}} \cdot \frac{-2\alpha \cdot \ddot{\alpha}}{1 - kr^2}$$

$$= \frac{\alpha \cdot \ddot{\alpha}}{c^2 \cdot (1 - kr^2)} \quad S_{\alpha y}!$$

b) Show $R_{00} = -3 \frac{\ddot{\alpha}}{\alpha}$

$R_{ij} = R^k_{ikj} \rightarrow$ Trace of Riemann-Tensor R^i_{ijk}

$$R^i_{ijk} = \frac{\partial \Gamma^i_{jk}}{\partial x^k} + \Gamma^i_{mk} \Gamma^m_{jk} - \frac{\partial \Gamma^i_{jk}}{\partial x^m} - \Gamma^i_{ml} \Gamma^m_{jk}$$

$$\Gamma_{11}^0 = \frac{a\dot{a}}{c^2(1-\kappa r^2)} \quad \Gamma_{22}^0 = \frac{1}{c^2} a\ddot{a}r^2 \quad \Gamma_{33}^0 = \frac{1}{c^2} a\ddot{a}r^2 \sin^2 \theta$$

$$\Gamma_{10}^1 = \Gamma_{01}^1 = \frac{\dot{a}}{a} \quad \Gamma_{11}^1 = \frac{kr}{1-\kappa r^2} \quad \Gamma_{22}^1 = -r(1-\kappa r^2) \quad \Gamma_{33}^1 = -r(1-\kappa r^2) \sin^2 \theta$$

$$\Gamma_{20}^2 = \Gamma_{02}^2 = \frac{\dot{a}}{a} \quad \Gamma_{21}^2 = \Gamma_{12}^2 = \frac{1}{r} \quad \Gamma_{33}^2 = -\sin \theta \cos \theta$$

$$\Gamma_{30}^3 = \Gamma_{03}^3 = \frac{\dot{a}}{a} \quad \Gamma_{31}^3 = \Gamma_{13}^3 = \frac{1}{r} \quad \Gamma_{32}^3 = \Gamma_{23}^3 = \cot \theta.$$

I hope there is no easier way to do this,

which I am overlooking :C

$$R_{jkl}^i = \frac{\partial \Gamma_{jl}^i}{\partial x^k} + \Gamma_{mk}^i \Gamma_{jl}^m - \frac{\partial \Gamma_{jk}^i}{\partial x^l} - \Gamma_{ml}^i \Gamma_{jk}^m$$

$j (= 0)$ for $i=k = \text{range}(0, 4)$ for $m = \text{range}(0, 4)$

$$= R_{000}^0 + R_{010}^1 + R_{020}^2 + R_{030}^3$$

$$= \frac{\partial \Gamma_{00}^i}{\partial x^k} + \Gamma_{mk}^i \Gamma_{00}^m - \frac{\partial \Gamma_{0k}^i}{\partial x^0} - \Gamma_{m0}^i \Gamma_{0k}^m$$

$$= \underbrace{\frac{\partial \Gamma_{00}^0}{\partial x^0} + \frac{\partial \Gamma_{00}^1}{\partial x^1} + \frac{\partial \Gamma_{00}^2}{\partial x^2} + \frac{\partial \Gamma_{00}^3}{\partial x^3}}_0 + \dots$$

$$= 0 + \Gamma_{m0}^0 \Gamma_{00}^m - \dots // i=k=0$$

$$= 0 + \underbrace{\Gamma_{00}^0 \Gamma_{00}^0 + \Gamma_{10}^0 \Gamma_{00}^1 + \Gamma_{20}^0 \Gamma_{00}^2 + \Gamma_{30}^0 \Gamma_{00}^3}_0 - \dots$$

$$= 0 + \Gamma_{n1}^1 \Gamma_{00}^m - \dots // i=1$$

$$= 0 + \Gamma_{01}^1 \Gamma_{00}^0 + \dots // m=0$$

$\frac{\ddot{a}}{a} \cdot 0 \rightarrow$ since $\Gamma_{00}^m = 0$, this term cancels out!

$$= 0 + 0 - \underbrace{\frac{\partial \Gamma_{0k}^i}{\partial x^0}}_{\partial t} - \dots$$

$$= 0 + 0 - \left(\frac{\partial \Gamma_{01}^1}{\partial t} + \frac{\partial \Gamma_{02}^2}{\partial t} + \frac{\partial \Gamma_{03}^3}{\partial t} \right) - \dots // \Gamma_{00}^0 = 0$$

$$= 0 + 0 - \underbrace{\frac{\partial}{\partial t} \left(3 \frac{\dot{a}}{a} \right)}_{3 \frac{\ddot{a} \cdot a - \dot{a}^2}{a^2}} - \Gamma_{m0}^i \Gamma_{0k}^m // \Gamma_{m0}^0 \Gamma_{00}^m = \Gamma_{00}^1 \Gamma_{01}^0 = 0$$

$$= 0 + 0 - 3 \cdot \frac{\ddot{a} \cdot a - \dot{a}^2}{a^2} - \left(\underbrace{\Gamma_{10}^1 \Gamma_{01}^1}_{\frac{\dot{a}}{a} \cdot \frac{a}{\dot{a}}} + \underbrace{\Gamma_{20}^2 \Gamma_{02}^2}_{\frac{\dot{a}}{a} \cdot \frac{a}{\dot{a}}} + \dots \right)$$

$$= 0 + 0 - 3 \cdot \left(\frac{\ddot{a} \cdot a - \dot{a}^2}{a^2} + \frac{\dot{a}^2}{a^2} \right)$$

$$= -3 \cdot \left(\frac{\ddot{a} \cdot a - \dot{a}^2 + \dot{a}^2}{a^2} \right)$$

$$= -3 \cdot \frac{\ddot{a}}{a} \quad \text{Slay!}$$

c) Show that $G_0^0 = -3 \left(\frac{1}{c^2} \left(\frac{\ddot{a}}{a} \right)^2 + \frac{k}{a^2} \right)$

$$G_{UV} = R_{UV} - \frac{1}{2} R g_{UV}$$

$$R_{00} = -3 \frac{\ddot{a}}{a}, \quad R = 6 \left(\frac{1}{c^2} \left(\frac{\ddot{a}}{a} + \frac{\dot{a}^2}{a^2} \right) + \frac{k}{a^2} \right)$$

$$\text{with } g^{ab} = \begin{pmatrix} -c^2 & 0 & 0 & 0 \\ 0 & \frac{1}{1+kr^2} & 0 & 0 \\ 0 & 0 & r^2 & 0 \\ 0 & 0 & 0 & r^2 \sin^2\theta \end{pmatrix} \cdot \begin{pmatrix} dt \\ dr \cdot \alpha^2 \\ ds \cdot \alpha^2 \\ dp \cdot \alpha^2 \end{pmatrix}$$

$$\hookrightarrow g_{00} = -\frac{1}{c^2}$$

\rightarrow Index heben mit Metrik

$$G^j_v = g^{ij} G_{iv}$$

$$G_{vv} = R_{vv} - \frac{1}{2} R g_{vv}$$

$$\begin{aligned} G^j_v &= g^{00} R_{00} - \frac{1}{2} R g^{00}_{00} \\ &= -\frac{1}{c^2} \left(-\frac{3\ddot{\alpha}}{\alpha} \right) - \frac{1}{2} \left(6 \cdot \left(\frac{1}{c^2} \cdot \left(\frac{\ddot{\alpha}}{\alpha} + \frac{\dot{\alpha}^2}{\alpha^2} \right) + \frac{k}{\alpha^2} \right) \cdot -\frac{1}{c^2} \right) \cdot -c^2 \\ &= 3 \frac{\ddot{\alpha}}{c^2 \alpha} - 3 \cdot \left(\frac{\ddot{\alpha}}{c^2 \alpha} + \frac{\dot{\alpha}^2}{c^2 \alpha^2} + \frac{k}{\alpha^2} \right) \\ &= -3 \cdot \left(\frac{1}{c^2} \cdot \left(\frac{\ddot{\alpha}}{\alpha} + \frac{k}{\alpha^2} \right) \right) \quad \text{Slay!} \end{aligned}$$

d) Show H^2 follows from G_0^0, T_v^v

$$G_{vv} + \lambda g_{vv} = \frac{8\pi G}{c^4} T_{vv} \quad // 144, \text{ only equation I found to contain everything}$$

$$\hookrightarrow G_{00} + \lambda g_{00} = \frac{8\pi G}{c^4} T_{00}$$

$$\begin{aligned} &= R_{00} - \frac{1}{2} R g_{00} + \lambda g_{00} = \frac{8\pi G}{c^4} \cdot \rho \\ &= -3 \frac{\ddot{\alpha}}{\alpha} - \frac{1}{2} \cdot \left(6 \left(\frac{1}{c^2} \left(\frac{\ddot{\alpha}}{\alpha} + \frac{\dot{\alpha}^2}{\alpha^2} \right) + \frac{k}{\alpha^2} \right) \right) \cdot -\frac{1}{c^2} - \frac{\lambda}{c^2} = \frac{8\pi G}{c^4} \rho \end{aligned}$$

$$= -3 \frac{\ddot{\alpha} \cdot c^2}{\alpha} - \left(3 \frac{\ddot{\alpha}}{c^2 \alpha} + 3 \frac{\dot{\alpha}^2}{c^2 \alpha^2} + 3 \frac{k}{\alpha^2} \right) - \lambda = \frac{8\pi G}{c^4} \rho \cdot c^2$$

$$= -3 \left[\frac{\dot{a}c^2}{a} + \frac{\dot{a}}{c^2a} + \frac{\ddot{a}^2}{c^2a^2} + \frac{kc}{a^2} \right] = \frac{8\pi G}{c^2} p + \lambda \quad // . c^2, : 3, - \frac{k}{a^2}$$

$$-\underbrace{\frac{\dot{a}^4}{a}}_{\text{Cancels out for } a=1?} + \frac{\dot{a}}{a} + \frac{\ddot{a}^2}{a^2} = \frac{8\pi G p}{3} + \frac{\lambda c^2}{3} - \frac{kc^2}{a^2}$$

$$\rightarrow \left(\frac{\dot{a}}{a} \right)^2 = \frac{8\pi G p}{3} + \frac{\lambda c^2}{3} - \frac{kc^2}{a^2}$$

2) starts on the next page :)



$$2) H^2 := \left(\frac{a}{\dot{a}}\right)^2 = \Lambda_0 a^{-4} + \Lambda_m a^{-3} + \Lambda_k a^{-2} + \Lambda_\Lambda$$

where $\Lambda_k = 1 - \Lambda_r - \Lambda_m - \Lambda_\Lambda$

a) Find equalities

$$H_0 = 67,74$$

$$\Lambda_m = 0.3089 \pm 0.0062$$

$$\Lambda_\Lambda = 0.6911 \pm 0.062$$

$$\Lambda_k = 0 \quad ;)$$

$$\Lambda_r = 8.4 \cdot 10^{-5}$$

$$(1.72)$$

$$\text{Equality, } \Lambda_{m,n}: \frac{\Lambda_m}{a^3 \Lambda_n} = 1$$

$$\hookrightarrow a^3 = \frac{\Lambda_m}{\Lambda_n}$$

$$a = \sqrt[3]{\frac{0.3089}{0.6911}}$$

Unsicherheiten sind für
Experimentalphysiker \approx

$$a(t) = 0.7645$$

$$\rightarrow t \approx 73.8 \cdot 0.7645 = 10.5501 \text{ GJ}$$

$$\text{Equality, } n,r: \quad a = \frac{\Lambda_r}{\Lambda_m} = \frac{8.4 \cdot 10^{-5}}{0.3089} = 0.00027$$

$$\rightarrow t \approx 0.0037 \text{ GJ}$$

Discussion: For a rather short time, the universe was dominated by radiation. For the next ~ 10.5 billion years, the universe is in the matter dominated era..

? Are the constants dependent on each other?

b) Find the solutions for $\Lambda_i = \Lambda_i$, $\Lambda_j = 0$

$$\text{Radiation: } H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \Lambda_k a^{-2}$$

$$\hookrightarrow \dot{a}^2 = \Lambda_k a^{-2}$$

Trennen über Variablen:

$$\left(\frac{da}{dt}\right)^2 = \Lambda_k a^{-2}$$

$$\frac{da}{dt} \cdot a = \sqrt{\Lambda_k}$$

$$a da = \sqrt{\Lambda_k} dt$$

$$\frac{a^2}{2} = \sqrt{\Lambda_k} t + C$$

$$a(t) = (2t\sqrt{\Lambda_k} + C)^{\frac{1}{2}}$$

(for whatever value Λ_k
would have for $\Lambda_j = 0$)

$$\text{Matter: } \left(\frac{da}{dt}\right)^2 = \Lambda_m a^{-1}$$

$$\frac{da}{dt} \sqrt{a} = \sqrt{\Lambda_m}$$

$$\frac{2}{3} a^{3/2} = \sqrt{\Lambda_m} t + C$$

$$a(t) = \left(\frac{3}{2} \cdot (\sqrt{\Lambda_m} t + C)^{\frac{2}{3}}\right)^{\frac{1}{2}}$$

(for whatever value Λ_m
would have for $\Lambda_j = 0$)

$$\Lambda: \left(\frac{\dot{a}}{a}\right)^2 = \Lambda_1$$

$$\left(\frac{da}{dt}\right)^2 = \Lambda_1 a^2$$

$$\frac{da}{dt} = \sqrt{\Lambda_1} \cdot a$$

$$\frac{da}{a} = \sqrt{\Lambda_1} dt$$

A certified Head-classic:

$$a = e^{\sqrt{\Lambda_1} t + C}$$

$$\text{Curvature: } \dot{\alpha}^2 = \sqrt{\kappa}$$

$$d\alpha = \sqrt{\kappa} dt$$

$$\alpha = \sqrt{\kappa} t + C$$

c) Provide these expressions in the limiting cases when $\kappa_m = 0, 0.3, 1, 1.3$ and plot them, using e.g. Python
 So I can use something "different" 😊

Verify that $\alpha(n)$ for $t(n)$ works for $(H)^2$

$$\alpha(n) = \frac{\sqrt{\kappa_m}}{2(1-\kappa_m)} \cdot \cosh(\sqrt{1-\kappa_m}n) - 1$$

$$t(n) = \frac{\sqrt{\kappa_m}}{2(1-\kappa_m)} \cdot \frac{\sinh(\sqrt{1-\kappa_m}n)}{\sqrt{1-\kappa_m}} - n$$

This took me way too long:

$$\left(\frac{\dot{\alpha}}{\alpha}\right)^2 = \frac{\kappa_m}{\alpha^2} + \frac{\kappa_k}{\alpha^2}$$

$$\rightarrow \dot{\alpha}^2 = \frac{\kappa_m}{\alpha} + \kappa_k$$

$$= \left(\frac{d\alpha(n)}{dt(n)}\right)^2 = \left(\frac{da}{dt} \frac{dt}{dn}\right)^2 = \left(\frac{da}{dt} \cdot \left(\frac{dn}{dt}\right)^{-1}\right)^2$$

$$= \frac{\frac{\sqrt{\kappa_m}}{2(1-\kappa_m)}}{\frac{\sqrt{\kappa_m}}{2(1-\kappa_m)}} \cdot \frac{\left(\sqrt{1-\kappa_m} \cdot \sinh(\sqrt{1-\kappa_m}n) \cdot 1\right)^2}{\left((\cosh(\sqrt{1-\kappa_m}n) \cdot 1) - 1\right)^2}$$

$$= \frac{(1-\lambda_m) \cdot \sinh^2(\sqrt{1-\lambda_m} n)}{\cosh^2(\sqrt{1-\lambda_m} n) - 2 \cdot \cosh(\sqrt{1-\lambda_m} n) + 1} \quad // (\cosh^2(x) - \sinh^2(x) = 1)$$

$$= \frac{(1-\lambda_m) \cdot \sinh^2(\sqrt{1-\lambda_m} n)}{\sinh^2(\sqrt{1-\lambda_m} n) - 2 \cdot \cosh(\sqrt{1-\lambda_m} n)}$$

From equation (1), we can tell that

$$\alpha = \frac{\lambda_m}{2 \cdot (1-\lambda_m)} \cdot (\cosh(\sqrt{1-\lambda_m} n) - 1)$$

Therefore:

$$H^2 = \frac{(1-\lambda_m) \sinh^2(\sqrt{1-\lambda_m} n)}{\sinh^2(\sqrt{1-\lambda_m} n) - 2 \cdot \cosh(\sqrt{1-\lambda_m} n)} = \frac{\frac{\lambda_m}{1-\lambda_m} \cdot (\cosh(\sqrt{1-\lambda_m} n) - 1)}{\frac{\lambda_m}{2(1-\lambda_m)} \cdot (\cosh(\sqrt{1-\lambda_m} n) - 1)} + \lambda_m$$

$$= \frac{2(1-\lambda_m)}{\cosh(\sqrt{1-\lambda_m} n) - 1} + (1-\lambda_m)$$

Divide by $(1-\lambda_m)$:

$$\rightarrow \frac{\sinh^2(\sqrt{1-\lambda_m} n)}{\sinh^2(\sqrt{1-\lambda_m} n) - 2 \cdot \cosh(\sqrt{1-\lambda_m} n)} = \frac{2}{\cosh(\sqrt{1-\lambda_m} n) - 1} + 1$$

$$= \frac{\sinh^2(\sqrt{1-\lambda_m} n)}{(\cosh(\sqrt{1-\lambda_m} n) - 1)^2} = \frac{2}{\cosh(\sqrt{1-\lambda_m} n) - 1} + 1$$

$$\rightarrow \sinh^2(\sqrt{1-\eta_m} n) = 2 \cdot \cosh(\sqrt{1-\eta_m} n) - 1 + (\cosh(\sqrt{1-\eta_m} n) - 1)^2$$

$$\rightarrow \sinh^2(\sqrt{1-\eta_m} n) = \cosh^2(\sqrt{1-\eta_m} n) - 1$$

$$\therefore \sinh^2(x) = \cosh^2(x) - 1$$

Say!

Originally, I wanted to use my own modified python-version:

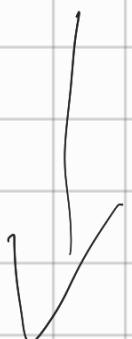
<https://github.com/actopozipc/German-Python-Interpreter>

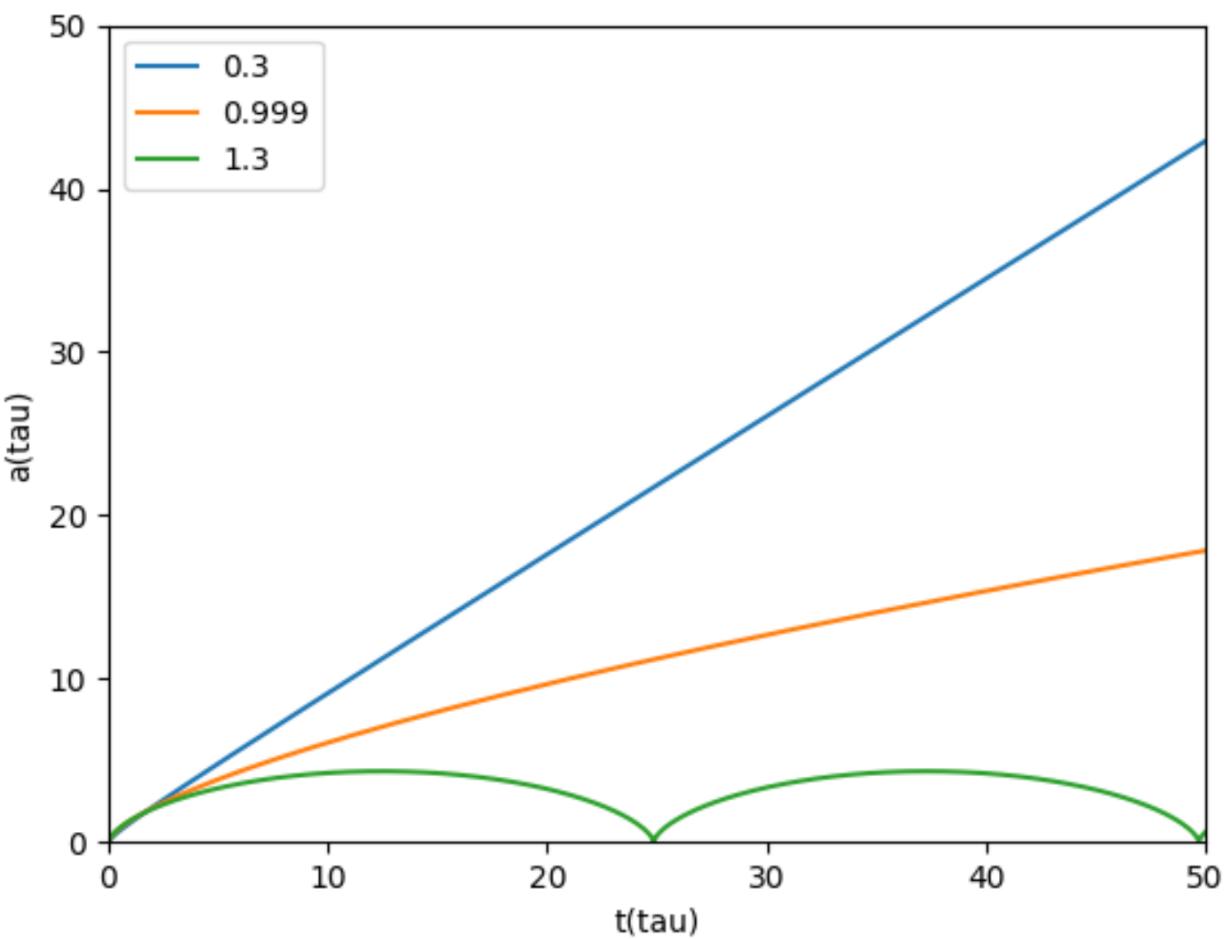
```
a = lambda omega_m, eta: (omega_m/(2*(1-omega_m)))*((np.cosh(np.emath.sqrt(1-omega_m)*eta)-1)) #emath needed for omega_m=1.3 >:(  
t = lambda omega_m, eta: (omega_m/(2*(1-omega_m)))*((np.sinh(np.emath.sqrt(1-omega_m)*eta)/(np.emath.sqrt(1-omega_m))-eta))  
a_values = []  
t_values = []  
für omega_m in [0.3,0.999,1.3]:  
    liste(karte(lambda t_value: t_values.erweiterte(karte(lambda a_value: np.real(t(a_value, t_value)), [omega_m])), np.arange(0,10,0.1)))  
    liste(karte(lambda t_value: a_values.erweiterte(karte(lambda a_value: np.real(a(a_value, t_value)), [omega_m])), np.arange(0,10,0.1)))  
    plt.plot(t_values, a_values, label=f"\{omega_m\}")  
a_values.aufräumen()  
t_values.aufräumen()
```

Sadly, there seems to be a bug in my lambda implementation, so I will do it

with default python:

```
a = lambda omega_m, eta: (omega_m/(2*(1-omega_m)))*((np.cosh(np.emath.sqrt(1-omega_m)*eta)-1)) #emath needed for omega_m=1.3 >:(  
t = lambda omega_m, eta: (omega_m/(2*(1-omega_m)))*((np.sinh(np.emath.sqrt(1-omega_m)*eta)/(np.emath.sqrt(1-omega_m))-eta))  
a_values = []  
t_values = []  
for omega_m in [0.3,0.999,1.3]:  
    list(map(lambda t_value: a_values.extend(map(lambda a_value: np.real(a(a_value, t_value)), [omega_m])), np.arange(0,10,0.1)))  
    list(map(lambda t_value: t_values.extend(map(lambda a_value: np.real(t(a_value, t_value)), [omega_m])), np.arange(0,10,0.1)))  
    plt.plot(t_values, a_values, label=f"\{omega_m\}")  
    a_values.clear()  
    t_values.clear()  
plt.xlabel("t(tau)")  
plt.ylabel("a(tau)")  
plt.show()
```





So, when does the big crunch happen?

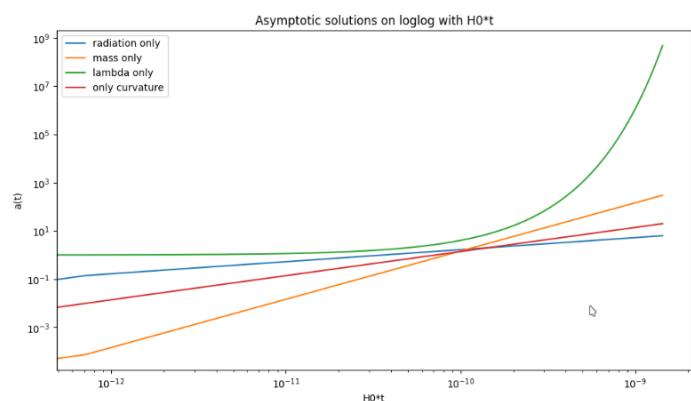
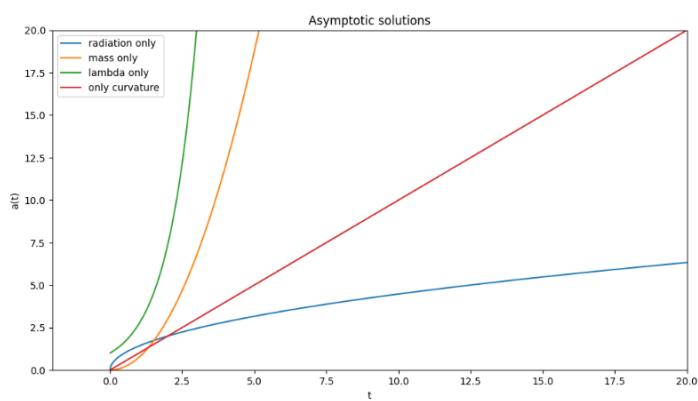
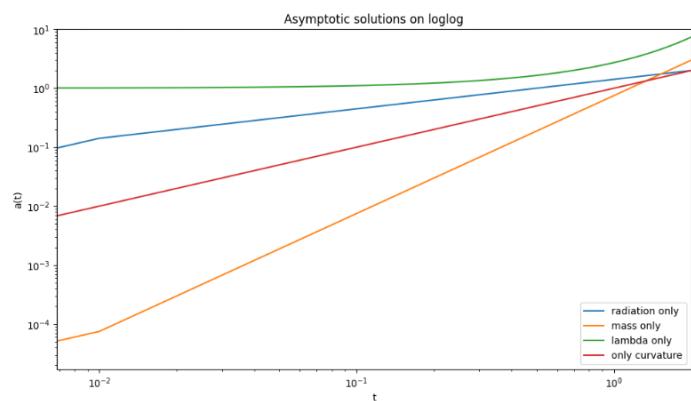
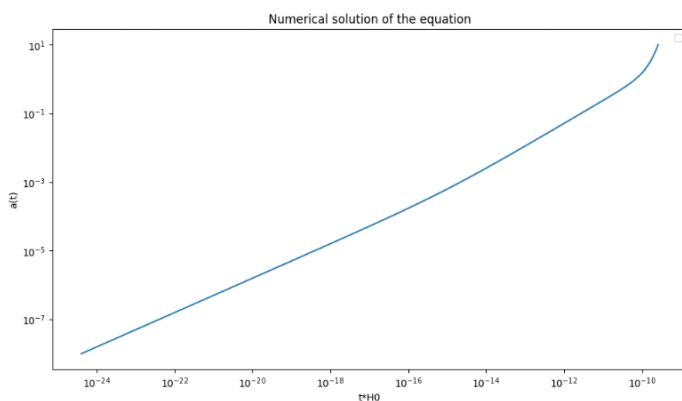
We can find out numerically:

```
big_crunch = lambda l1,l2: [l2[i] for i in range(len(l1)) if l1[i]<0.001]
```

And, as expected from the plot, it doesn't happen for $L_m = 0.3, 1$
but is a returning event for $L_m = 1.3$:

```
[0.0, 2.5008751458474787e-05] 0.3
[0.0] 1
[-0.0, 24.854863755913854, 49.709713900414435] 1.3
```

o) Numerical solutions for given parameters



The first window shows the numerical solution.

I wasn't sure about the x-axis for the asymptotic

solutions, so I made 3: t, log(t) and log(H_0*t)

```
names = ["radiation only", "mass only", "lambda only", "only curvature"]
for name_index, f in enumerate([lambda t: np.sqrt(2*t), lambda t: (3/2 * t)**2/3, lambda t: np.exp(t), lambda t: t]):
    axs[1,0].plot(a,f(a), label=names[name_index])
    axs[0,1].loglog(a,f(a), label=names[name_index])
    axs[1,1].loglog(H0*a,f(a), label=names[name_index])
```

(I will upload the full code)