



Superconducting Triangular Islands as a Platform for Manipulating Majorana Zero Modes

Aidan Winblad
Hua Chen

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Motivation

Kitaev
Triangle

Hollow
Triangle

Braiding

Summary



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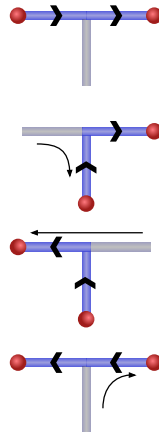
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- P-wave superconductors contain half-quantum vortices.
 - Majorana fermions located at core of a vortex.
 - Braiding vortices exhibits Non-Abelian statistics.
- 1D p-wave superconductors host Majorana fermions on end points.
 - Measurements in real systems:
 - V. Mourik, *Science* **336**, 1003 (2012)
 - S. Nadj-Perge, *Science* **346**, 602 (2014)
 - L. Schneider, *Nat. Nanotechnol.* **17**, 384 (2022)
- Quasi-1D T-junction
 - Braiding of Majorana fermions is defined for 2D.
 - In practice challenging to make, but still feasible and seriously pursued.



Alicea, *Nature Phys.* **7**, 412 (2011)

arXiv:
2309.11607



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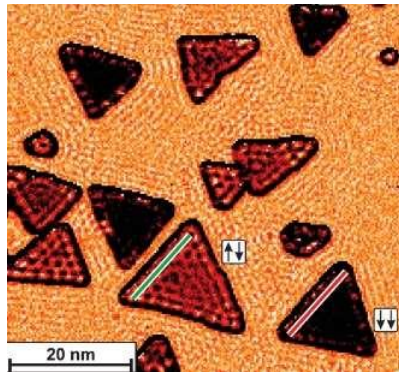
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- Consider triangular islands, topologically similar to T-junctions.
- Islands of three-fold rotational symmetry occur naturally in epitaxial growth on close-packed metal surfaces.
- Good platform for transition from 2D to 1D topological superconductor.



Triangular Co islands on Cu(111).
Pietzsch et al., *PRL* **96**, 237203 (2006)

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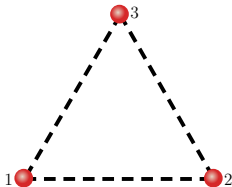


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$$\mathcal{H} = \sum_{\langle j,l \rangle} \left[-te^{i\phi_{jl}} c_j^\dagger c_l + \Delta e^{i\theta_{jl}} c_j c_l + h.c. \right] - \sum_j \mu c_j^\dagger c_j$$



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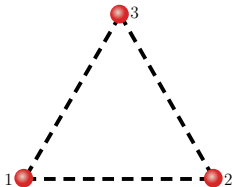
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$$\begin{aligned} (\phi_{12}, \phi_{23}, \phi_{31}) &= \left(0, -\frac{\pi}{3}, -\frac{\pi}{3}\right) = \phi_1 \\ &\rightarrow \left(-\frac{\pi}{3}, -\frac{\pi}{3}, 0\right) = \phi_2 \\ &\rightarrow \left(-\frac{\pi}{3}, 0, -\frac{\pi}{3}\right) = \phi_3 \\ &\rightarrow \phi_1 \end{aligned}$$



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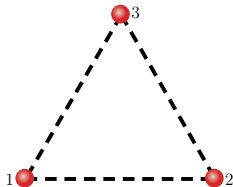
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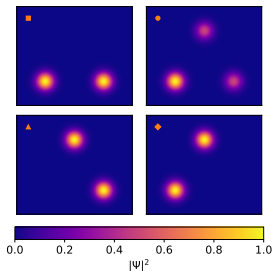
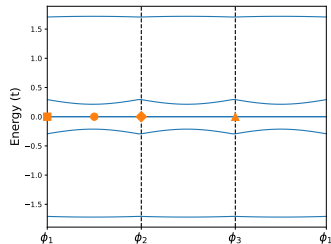
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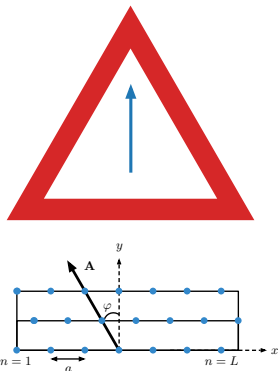
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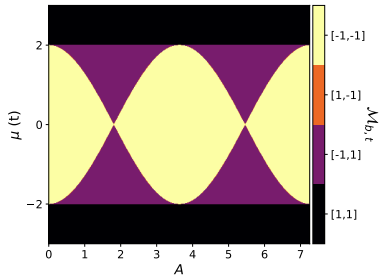
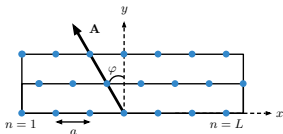
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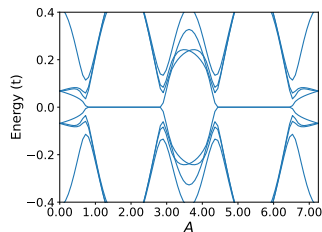
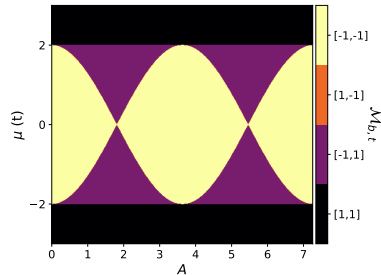
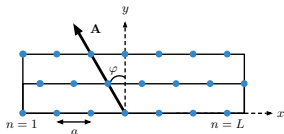
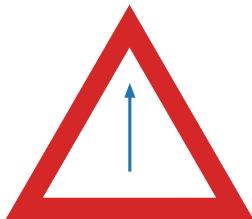


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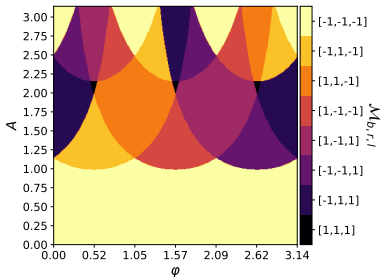
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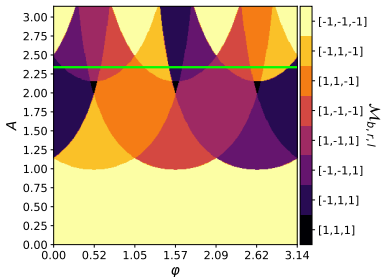
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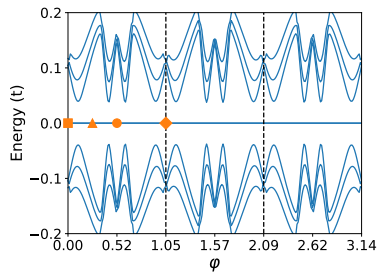
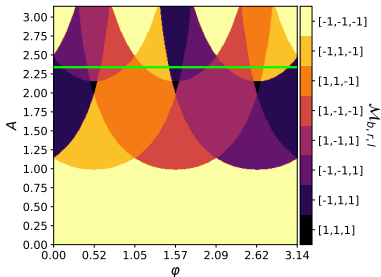
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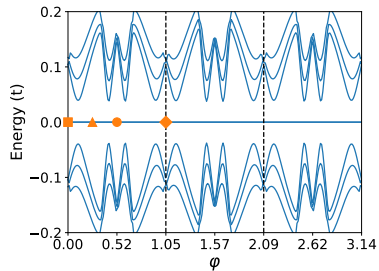
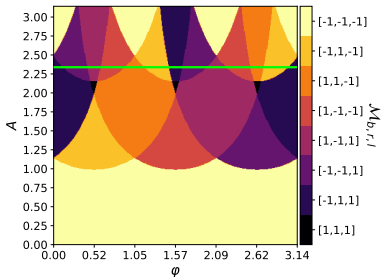
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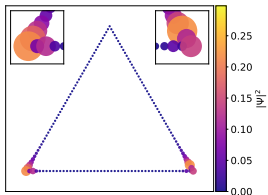
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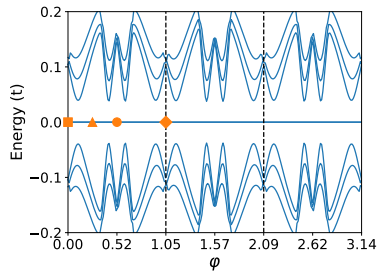
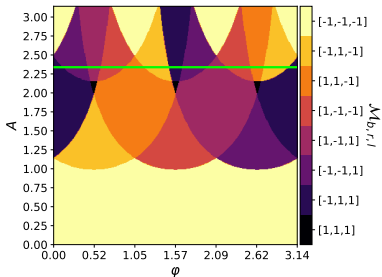
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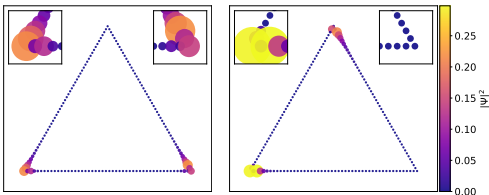
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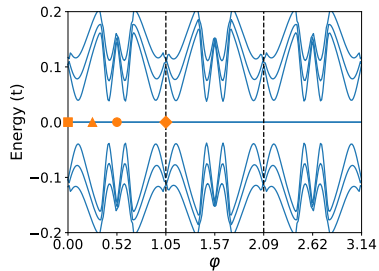
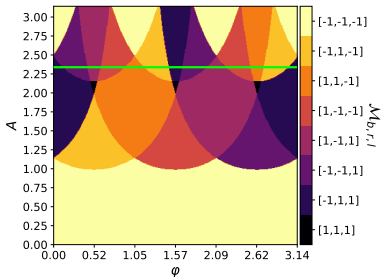
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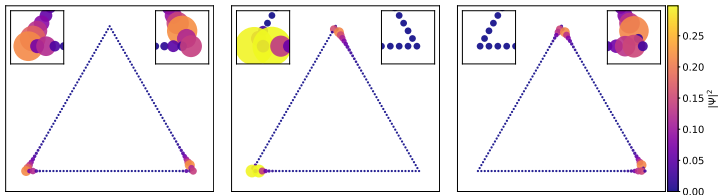
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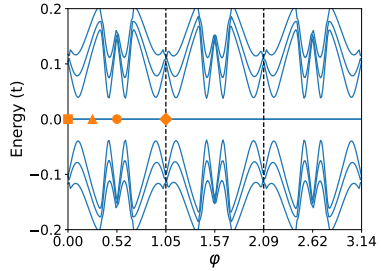
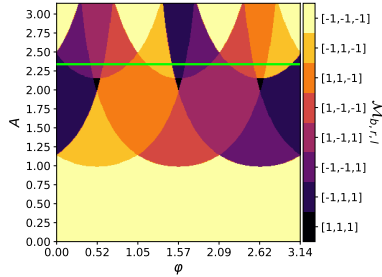
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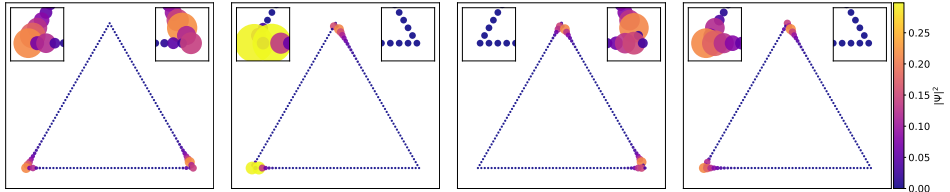
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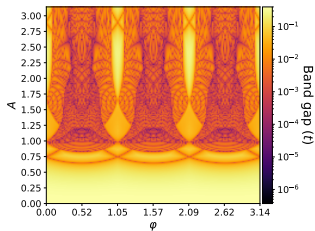
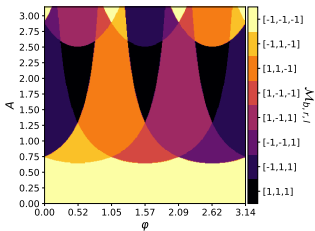
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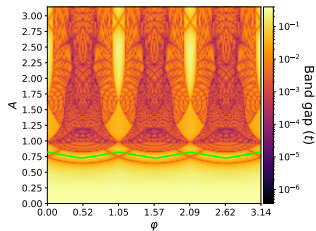
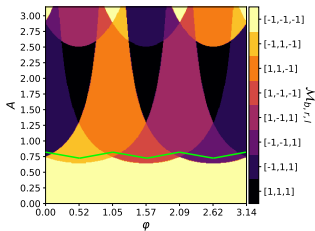
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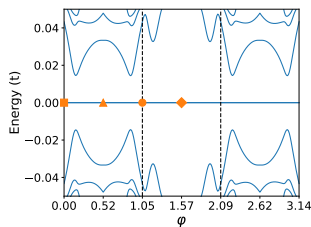
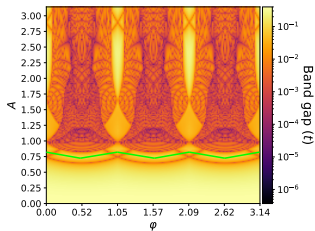
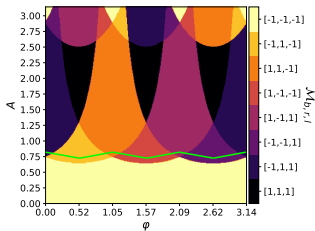
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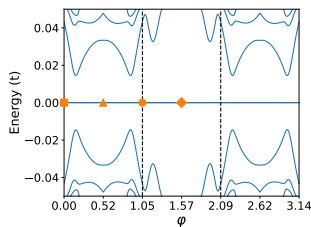
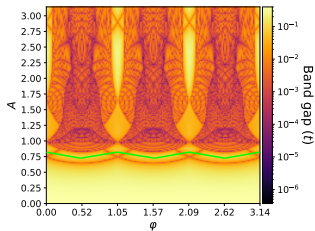
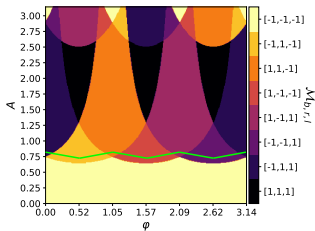
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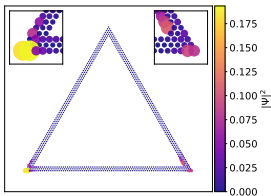
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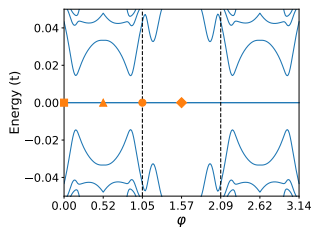
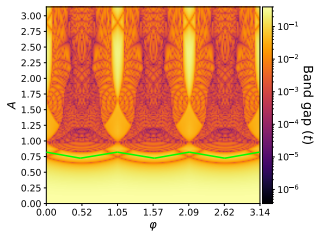
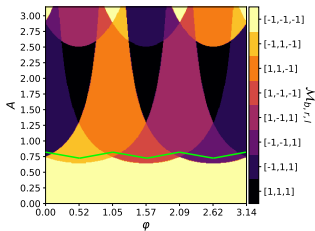
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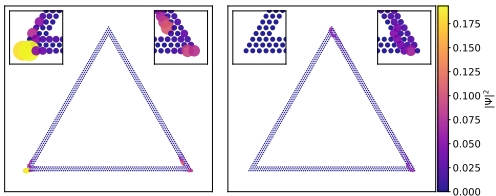
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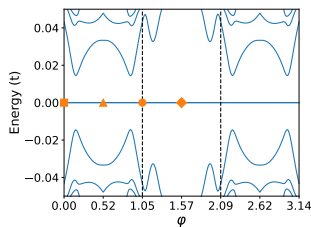
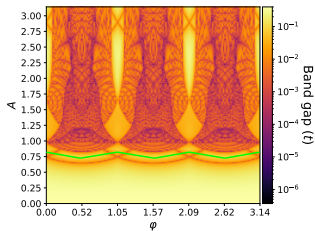
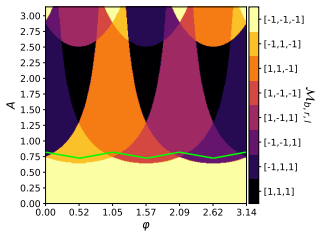
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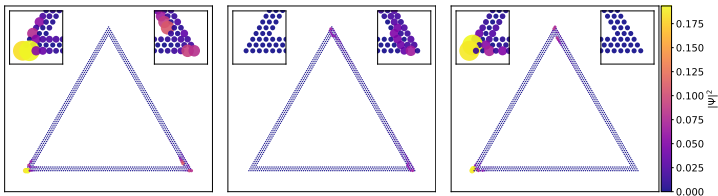
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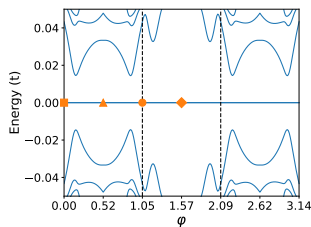
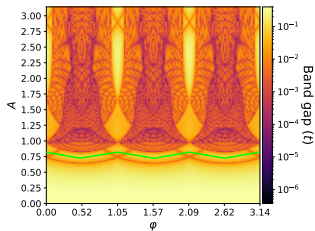
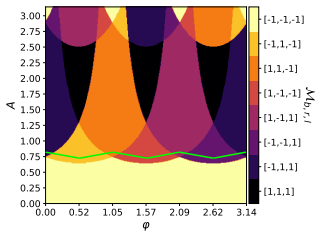
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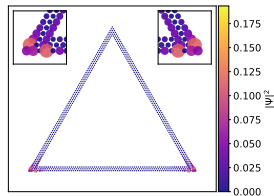
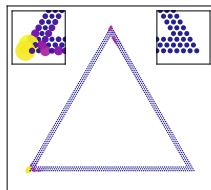
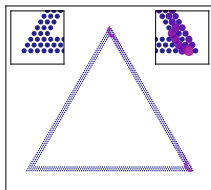
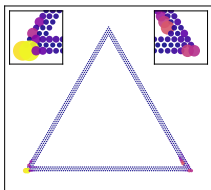
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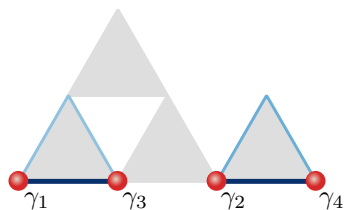
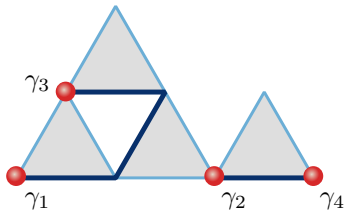
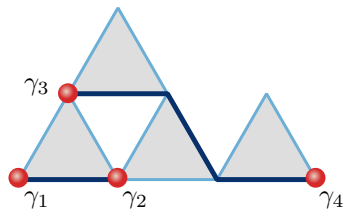
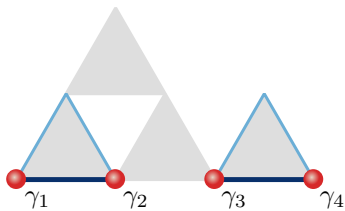
Kitaev
Triangle

Hollow
Triangle

Braiding

Summary

arXiv:
2309.11607





Summary

Superconducting
triangular islands
as a platform for
manipulating
Majorana zero
modes

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Motivation

Kitaev
Triangle

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Braiding

Summary

- Introduction of Peierls phase allows for a minimal Kitaev triangle, reducing fermionic sites down to 3.
- Vector potential field and its rotation allows additional tunability of topology.
- MZM can be hosted and braided on a network of triangular islands.

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Majorana fermion notation and coupling isolations

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The complex fermion operator can be written as a superposition of two Majorana fermions $c_j = \frac{1}{2}(a_j + ib_j)$. Due to the nature of Majorana fermions, $a_j^\dagger = a_j$, the creation operator is $c_j^\dagger = \frac{1}{2}(a_j - ib_j)$.

$$H = -\frac{i\mu}{2} \sum_j a_j b_j - \frac{i}{2} \sum_{\langle jl \rangle} [(t \sin \phi_{jl} - \Delta \sin \theta_{jl}) a_l a_j + (t \sin \phi_{jl} + \Delta \sin \theta_{jl}) b_l b_j + (t \cos \phi_{jl} - \Delta \cos \theta_{jl}) a_l b_j - (t \cos \phi_{jl} + \Delta \cos \theta_{jl}) b_l a_j].$$

$$(t \sin \phi_{jl} - \Delta \sin \theta_{jl}) a_l a_j, \quad (1)$$

$$(t \sin \phi_{jl} + \Delta \sin \theta_{jl}) b_l b_j, \quad (2)$$

$$(t \cos \phi_{jl} + \Delta \cos \theta_{jl}) a_l b_j, \quad (3)$$

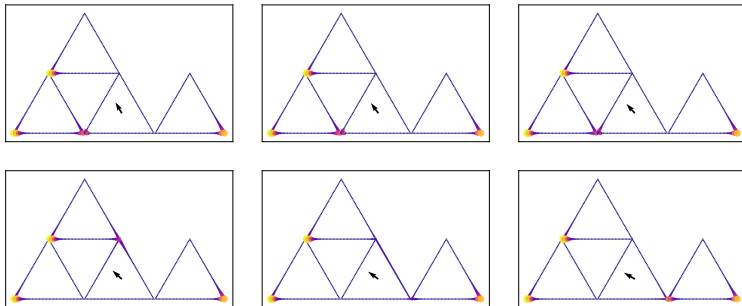
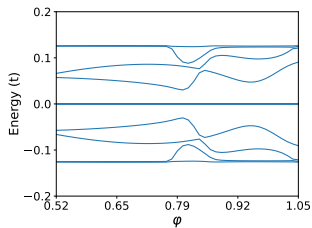
$$(t \cos \phi_{jl} - \Delta \cos \theta_{jl}) b_l a_j \quad (4)$$



Braiding MZM in a Small Network of Triangles

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