

STERILE NEUTRINOS

"The existence of hundreds of new particles may seem far fetched,
but is in principle not impossible."

(Boyarski et al. 2018)



Antoniadis, Carreira, & Karchev. GRAPPA Student Seminar (2019)

MOTIVATION

Neutrinos

- Electrically neutral
- Only interact weakly
- Long-living
- Oscillate \Rightarrow Massive \Rightarrow DM candidate?

Neutrinos

- $\rho_\nu = n_\alpha \sum m_i \quad , \quad n_\alpha = \frac{4}{6} \frac{\zeta(3)}{\pi^2} T_\nu^3$
- $T_\nu \simeq \left(\frac{4}{11}\right)^{\frac{1}{3}} T_\gamma \simeq 10^{-4} eV$
- DM candidates $\Rightarrow \sum m_i \simeq 11.5 eV$
- But $\sum m_i < 0.12 eV$

Neutrinos

- $m_\nu \ll T_\nu$
 - ⇒ Neutrinos decouple relativistic
 - ⇒ No time for large structure formation
- Hot DM

Neutrinos

- $\sum m_i < 0.12\text{eV}$
- Hot DM
 - ⇒ Not a viable DM candidate

The SM does not contain a viable DM candidate

*“particles that are, from the point of view of ordinary weak interactions, **sterile**, i.e. practically unobservable, since they have the “incorrect” helicity”*

(Bruno Pontecorvo, 1967)

Sterile neutrinos

- With a mass of a few keV, sterile neutrinos could make up all DM
- Could also solve the neutrino mass problem

Sterile neutrinos

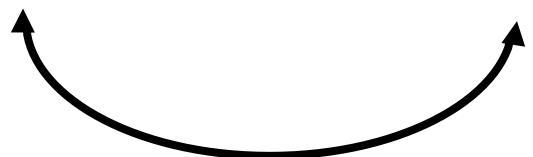
$$\blacksquare \quad \mathcal{L} = \mathcal{L}_{SM} + i\overline{\nu_{Ri}}\partial\nu_{Ri} - \left[\underbrace{\frac{1}{2}\overline{\nu_{Ri}^c}(M)_{ij}\nu_{Rj}}_{\text{Kinetic}} - \underbrace{Y_{ai}\overline{l_{La}}\epsilon\phi^*\nu_{Ri}}_{\text{Majorana}} + h.c. \right]_{\text{Interactions}}$$

- At low energies,

$$\Rightarrow \mathcal{L}_{eff} = \mathcal{L}_{SM} + \frac{1}{2}\overline{l_L}\tilde{\phi}YM_M^{-1}Y^T\tilde{\phi}^Tl_L^c$$

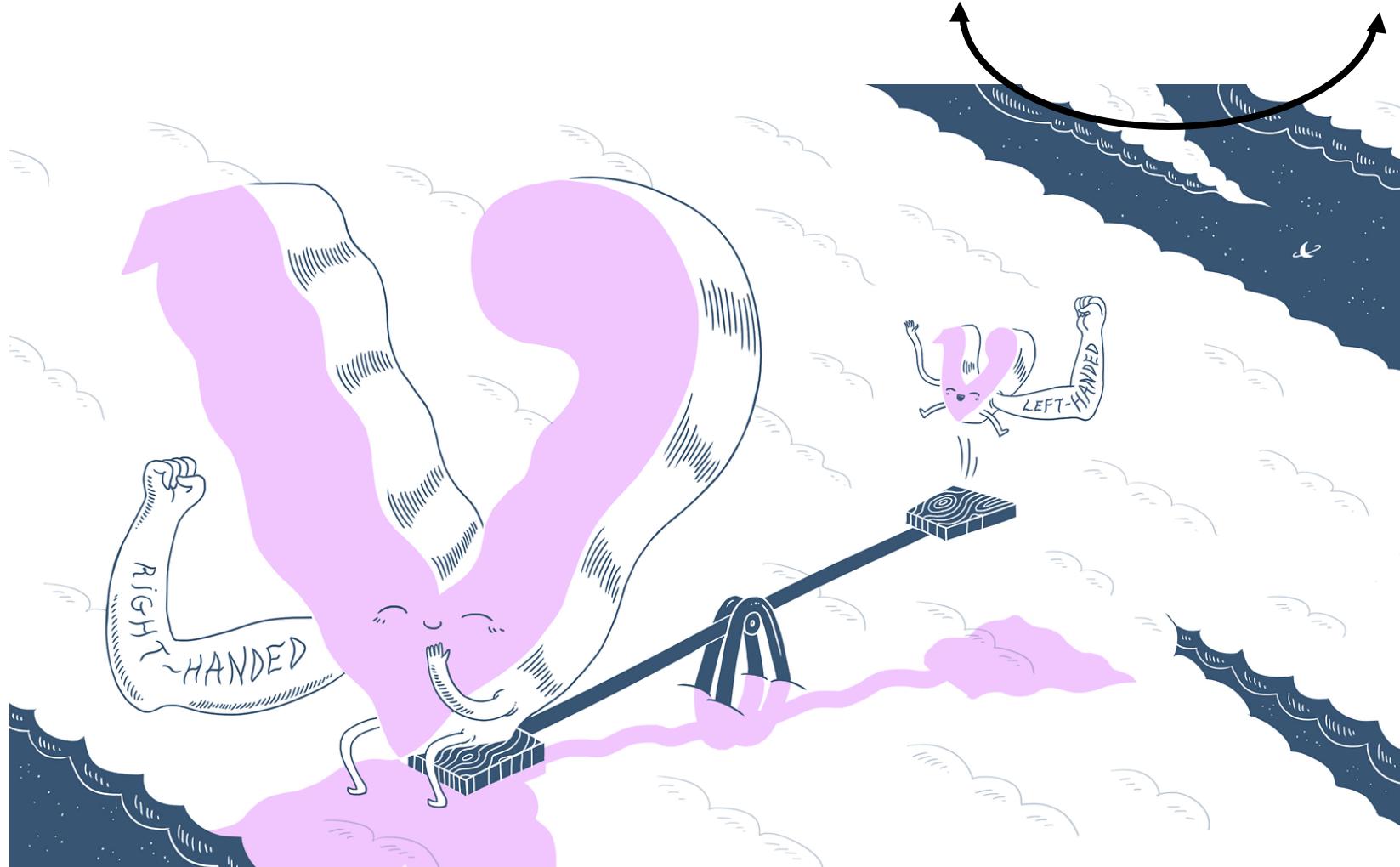
Sterile neutrinos

- $\mathcal{L}_{eff} = \mathcal{L}_{SM} + \frac{1}{2} \overline{l}_L \tilde{\phi} Y M_M^{-1} Y^T \tilde{\phi}^T l_L^c$
- The Higgs mechanism then generates a Majorana mass term of the form $\frac{1}{2} \overline{\nu}_L m_\nu \nu_L^c$
- With $m_\nu = -v^2 Y M_M^{-1} Y^T \Rightarrow$ Seesaw mechanism



Seesaw mechanism

$$m_\nu = -v^2 Y M_M^{-1} Y^T$$



Seesaw mechanism



$$m_\nu = -v^2 Y M_M^{-1} Y^T$$

$$m_\nu \simeq 0.1\text{eV}$$

$$Y \simeq 1$$



$$M_N \simeq \text{GUT scale}$$

Seesaw mechanism



$$m_\nu = -v^2 Y M_M^{-1} Y^T$$

$$m_\nu \simeq 0.1\text{eV}$$

$$Y \ll 1$$



$M_N \simeq \text{keV scale}$
(Dark Matter)

PRODUCTION MECHANISMS



Thermal
Production via
Mixing ("freeze in")

Non-
resonant
Production

Resonant
Production



Thermal Production via New
Gauge Interactions ("freeze
out")



Non-thermal Production in the
decay of heavier particles

Thermal production via mixing ("freeze in")

- Weak interaction
- Active-sterile mixing (seesaw model)

$$|\nu_a\rangle = \cos \theta |\nu_1\rangle + \sin \theta |\nu_2\rangle$$

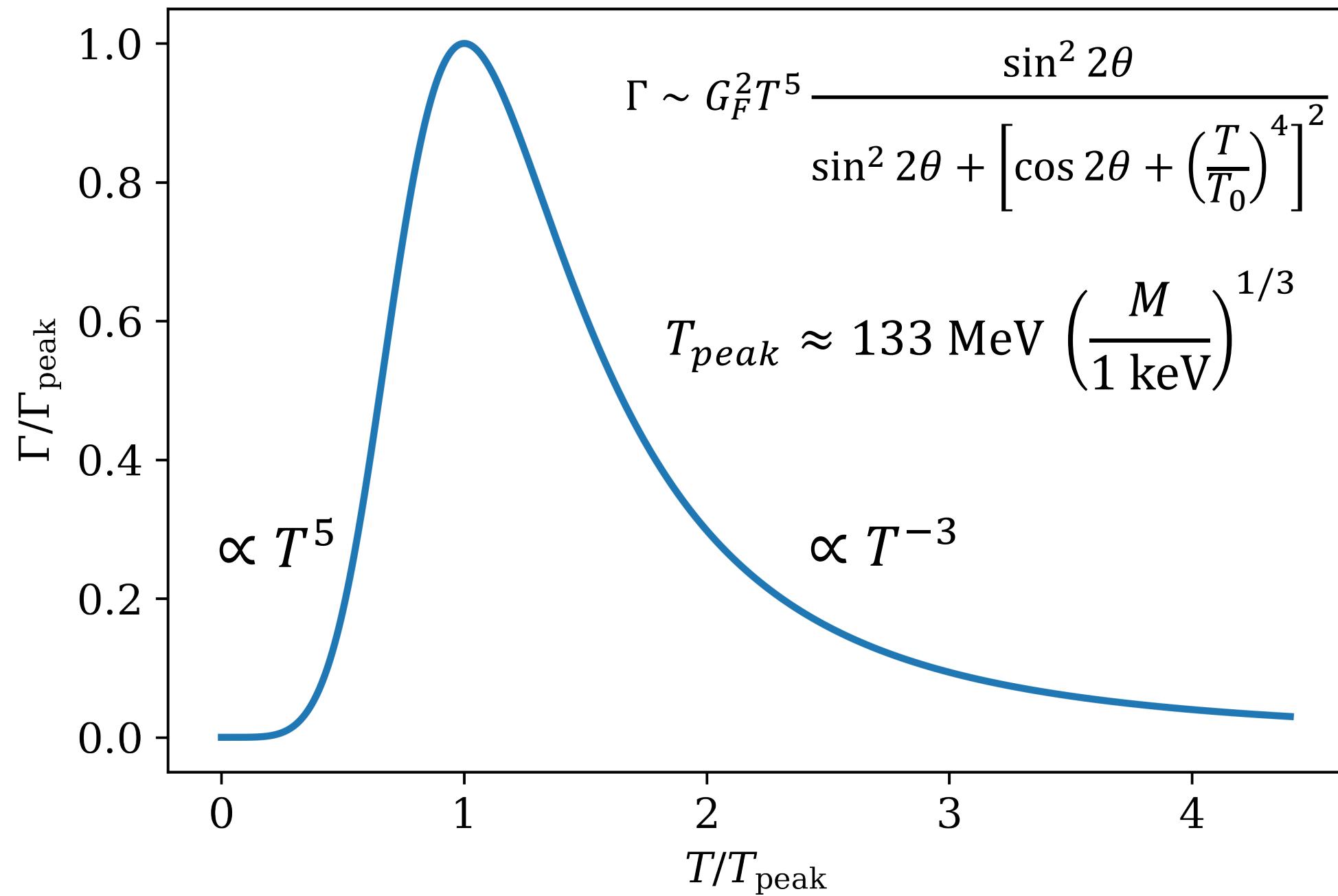
$$|\nu_s\rangle = -\sin \theta |\nu_1\rangle + \cos \theta |\nu_2\rangle$$

Non-resonant production

- Coherent oscillations between active and sterile states
- Decoherent scatterings
 - *Destroy the coherence of the quantum state*
 - *Force it into weak interaction*

$$|\nu_a\rangle = \cos \theta_m(t) |\nu_1(t)\rangle + \sin \theta_m(t) |\nu_2(t)\rangle$$
$$|\nu_s\rangle = -\sin \theta_m(t) |\nu_1(t)\rangle + \cos \theta_m(t) |\nu_2(t)\rangle$$

$$\Gamma \sim G_F^2 T^5 \frac{\sin^2 2\theta}{\sin^2 2\theta + \left[\cos 2\theta + \left(\frac{T}{T_0} \right)^4 \right]^2}$$



Resonant production

- Lepton asymmetries

$$l_{\nu_a} \equiv (n_{\nu_a} - n_{\bar{\nu}_a})/n_\gamma$$

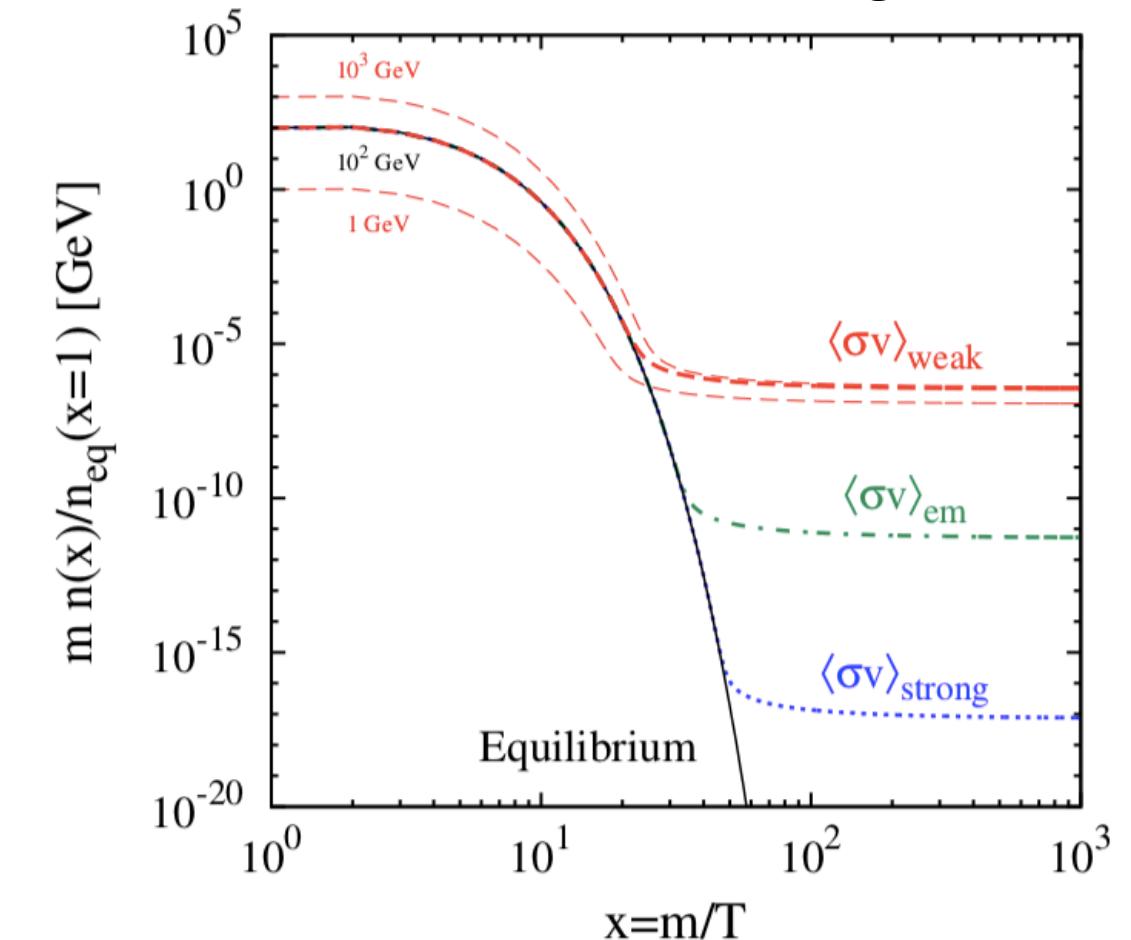
- Mikheyev–Smirnov–Wolfenstein (MSW) effect
 - *Enhances the transition between two kinds of neutrinos in a changing lepton density*
- Disfavoured by constraints from Lyman- α forest, dwarf galaxy number counts and X-ray data
- Uncertain how large lepton asymmetries could be generated

New gauge interactions (“freeze-out”)

$$\Gamma_N = G_F^2 T^5 \left(\frac{m_W}{m_{W_R}} \right)^4$$

$m_{W_R} > 1 \text{ TeV}$

Steigman et al. 2012

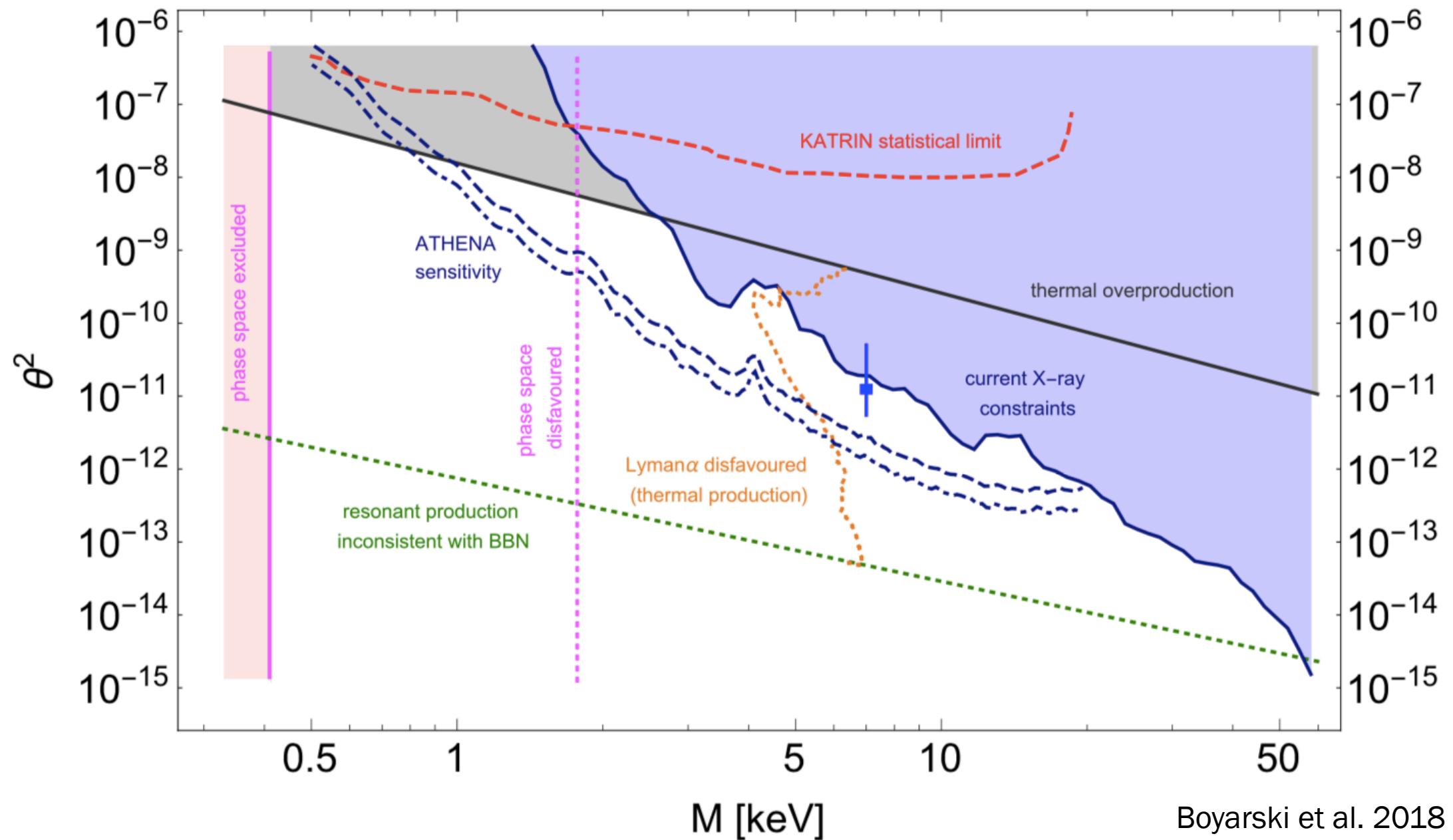


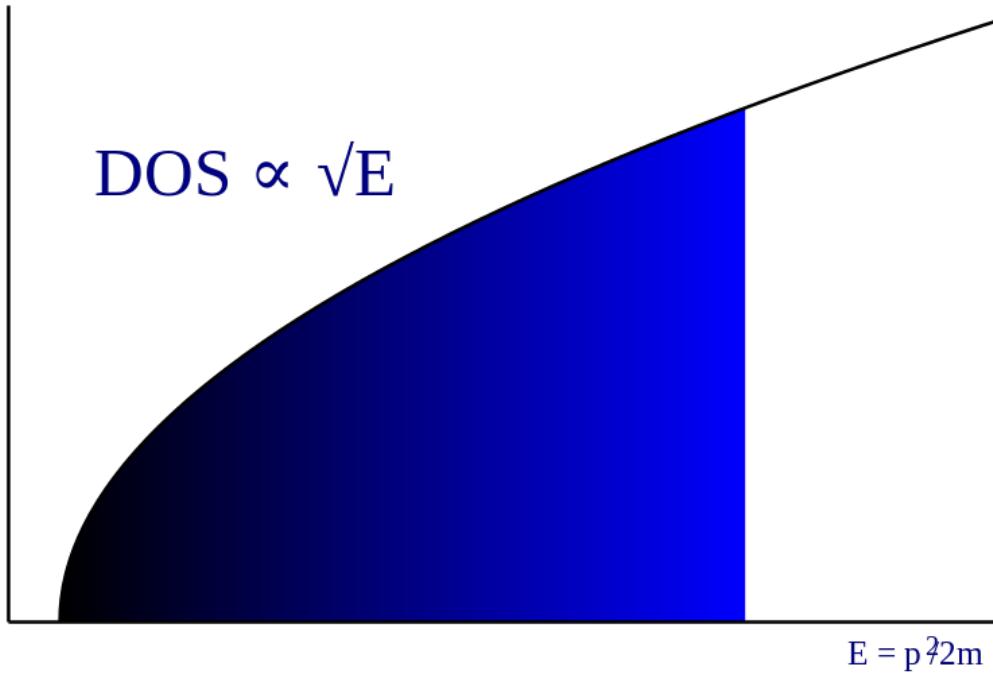
$$\Omega_N/\Omega_{DM} = \frac{1}{S} \left(\frac{10.75}{g_{*S}(t_f)} \right) \left(\frac{M}{\text{keV}} \right) \times 100$$

Non-thermal Production in the decay of heavier particles

- Decay far from thermal equilibrium
- Decay of a heavier parent particle
- Sterile neutrino in keV scale
- highly relativistic
- Possible decays of pions, the Higgs boson, W bosons

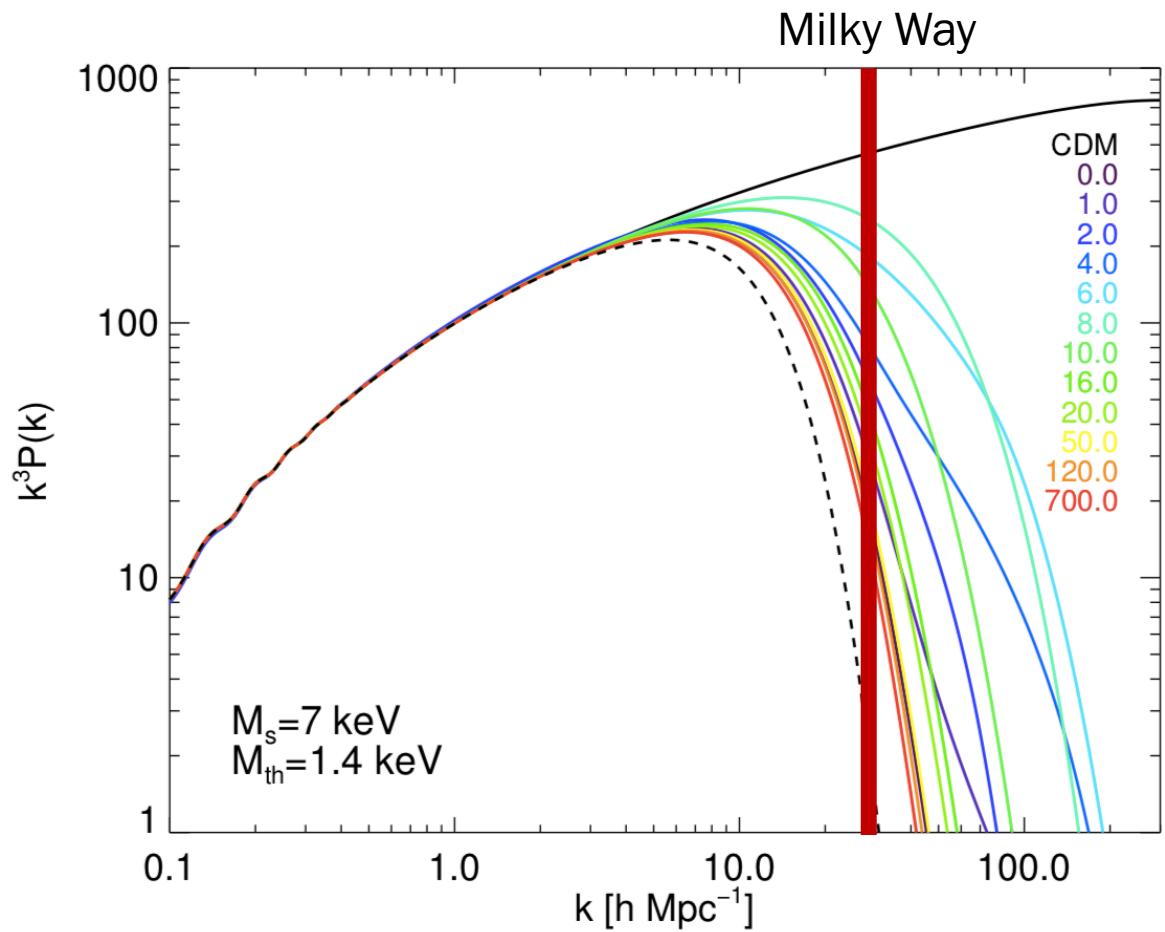
OBSERVATIONAL EVIDENCE





$$M_{DM} > \frac{9\pi\hbar^3}{4\sqrt{2}G_N^{3/2}} \frac{1}{\sqrt{M_{halo}R^3}}$$
$$\approx 100 \text{ eV}$$

Phase space saturation



Warm
dark
matter

Lensing

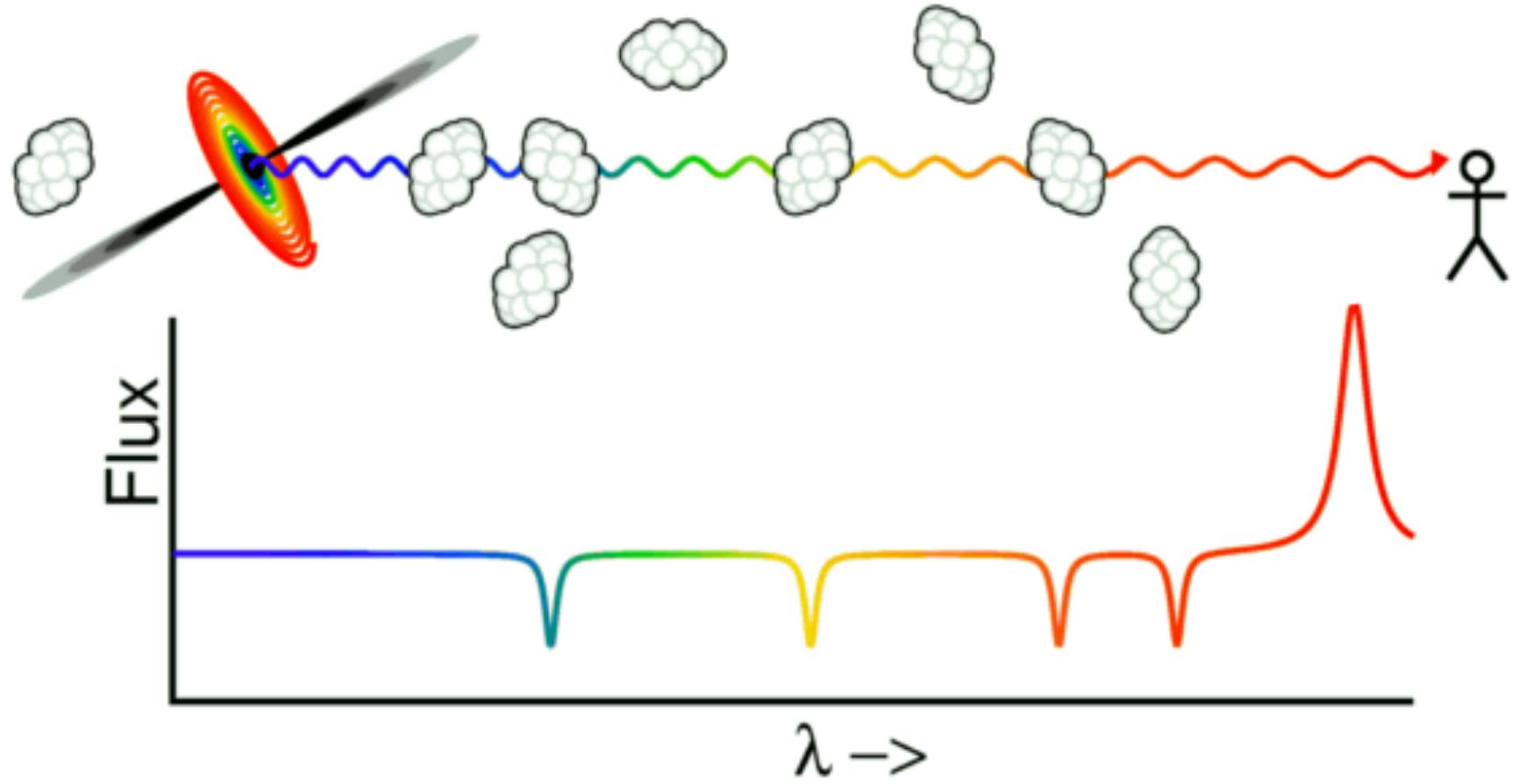
- Lensing applicable at low z
- CDM/WDM difference most pronounced at high-z

Lensing

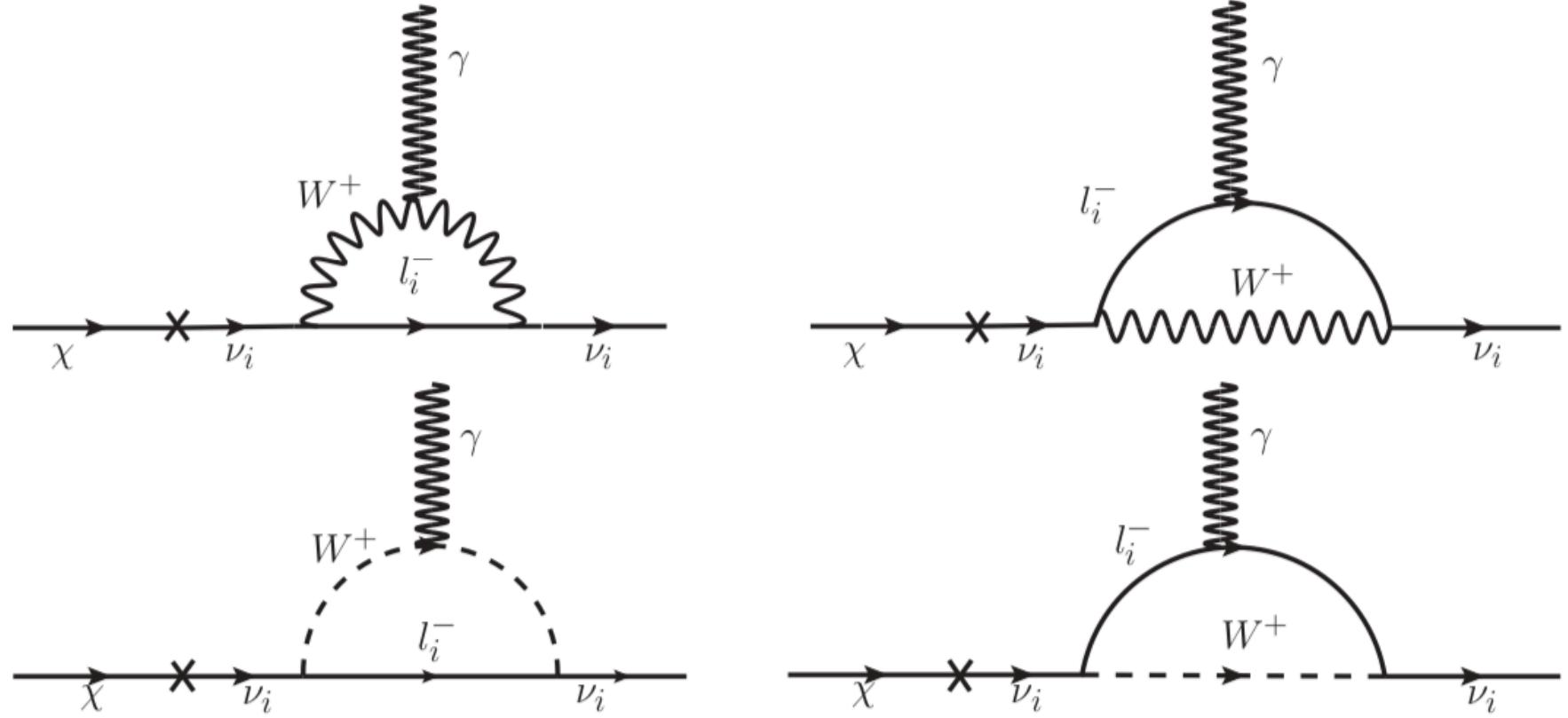
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Counting haloes

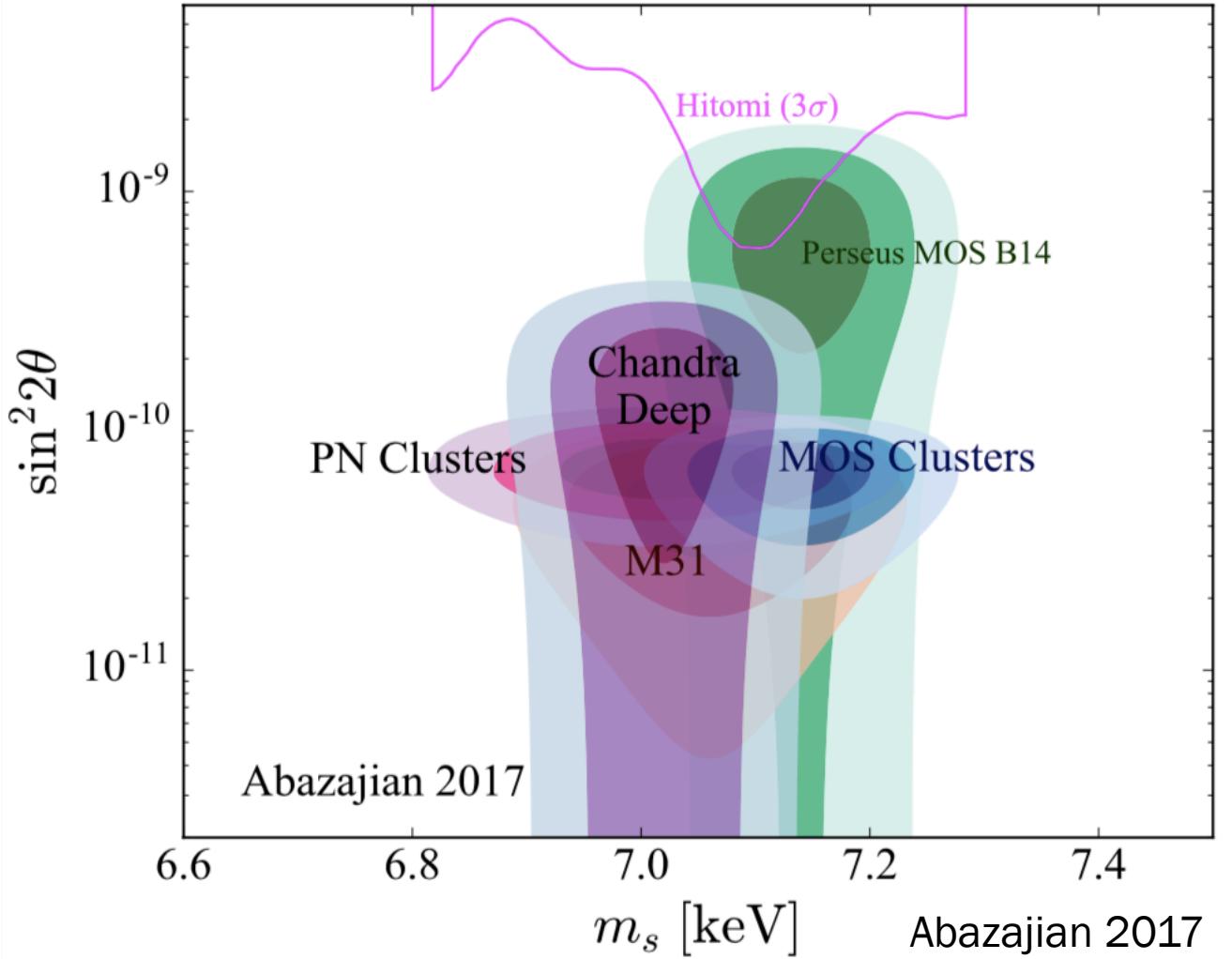
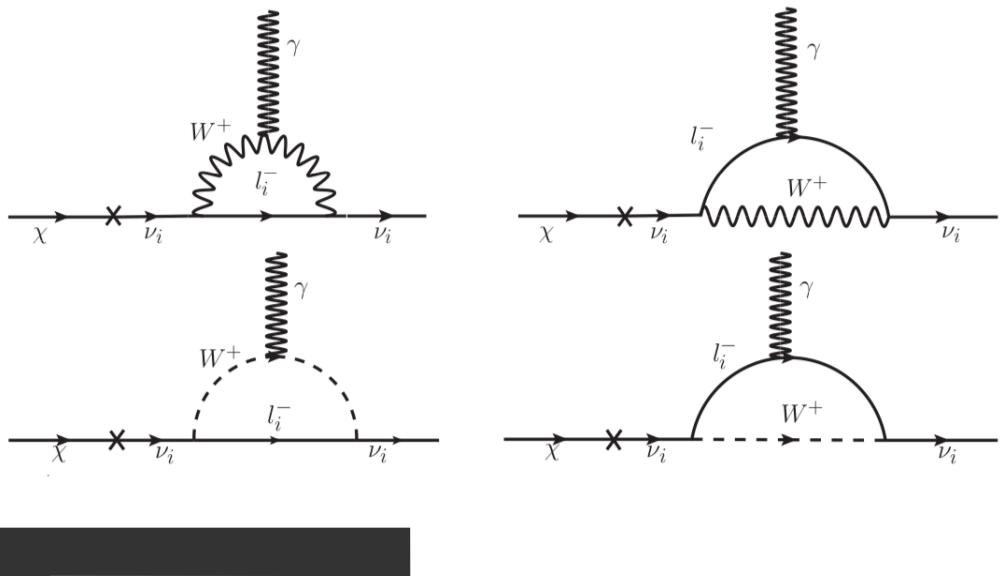
- Nonlinear structure formation
 - *realization variance (possibly larger than physical effects)*
- Small haloes countable only for very local universe (MW, Andromeda)
 - *Milky Way has fewer satellites than predicted*
- Baryonic structures may not trace DM
- Effects of pressure



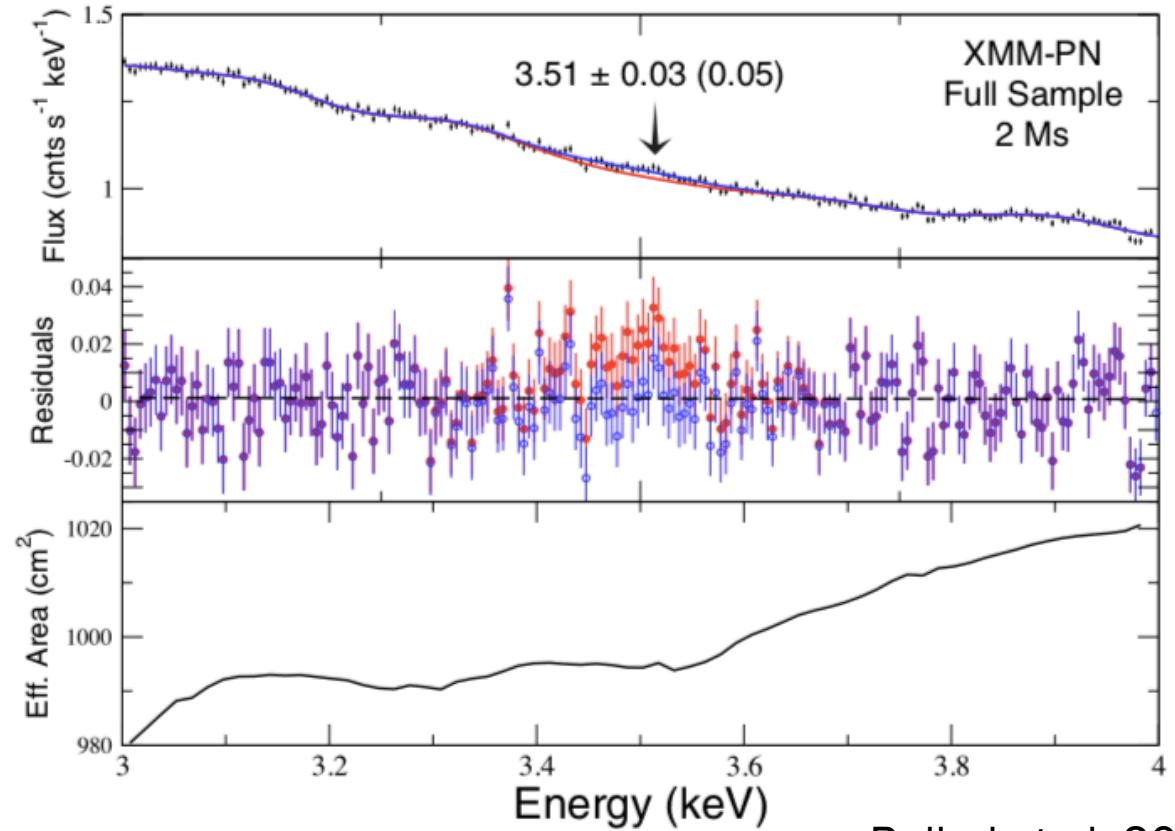
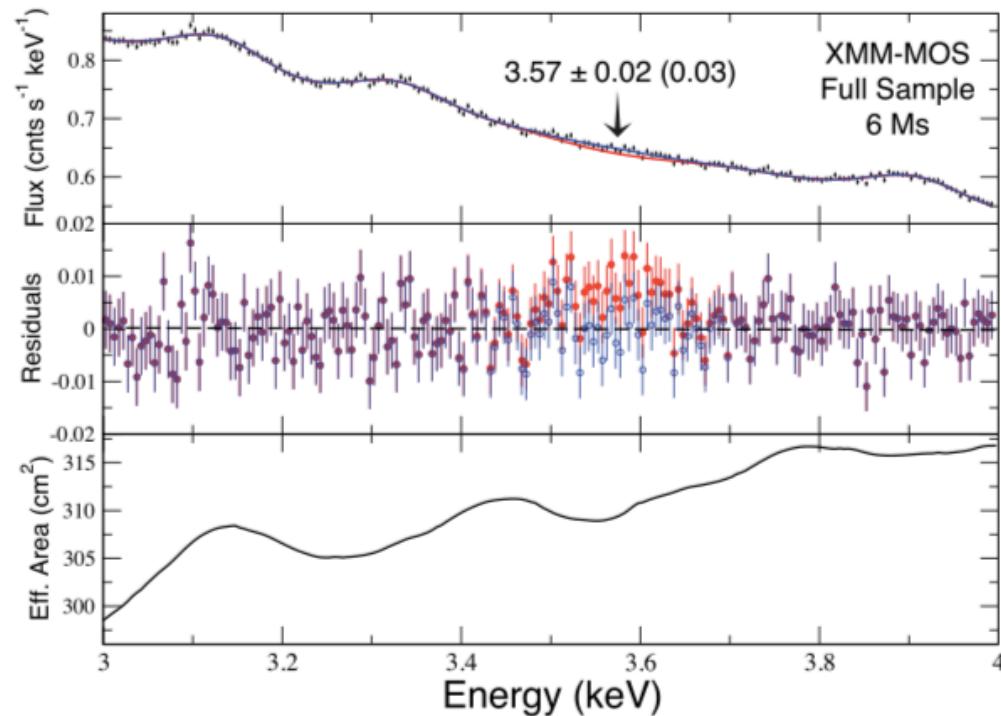
Lyman- α forest



Photon emission & the 3.5 keV line



Photon emission & the 3.5 keV line



Bulbul et al. 2014

Significant?

A novel scenario for the possible X-ray line feature at ~3.5 keV: Charge exchange with bare sulfur ions

Liyi Gu¹, Jelle Kaastra^{1, 2}, A. J. J. Raassen^{1, 3}, P. D. Mullen⁴, R. S. Cumbee⁴, D. Lyons⁴, and P. C. Stancil⁴

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Dark matter searches going bananas: the contribution of Potassium (and Chlorine) to the 3.5 keV line

Tesla Jeltema^{1*} and Stefano Profumo^{1†}

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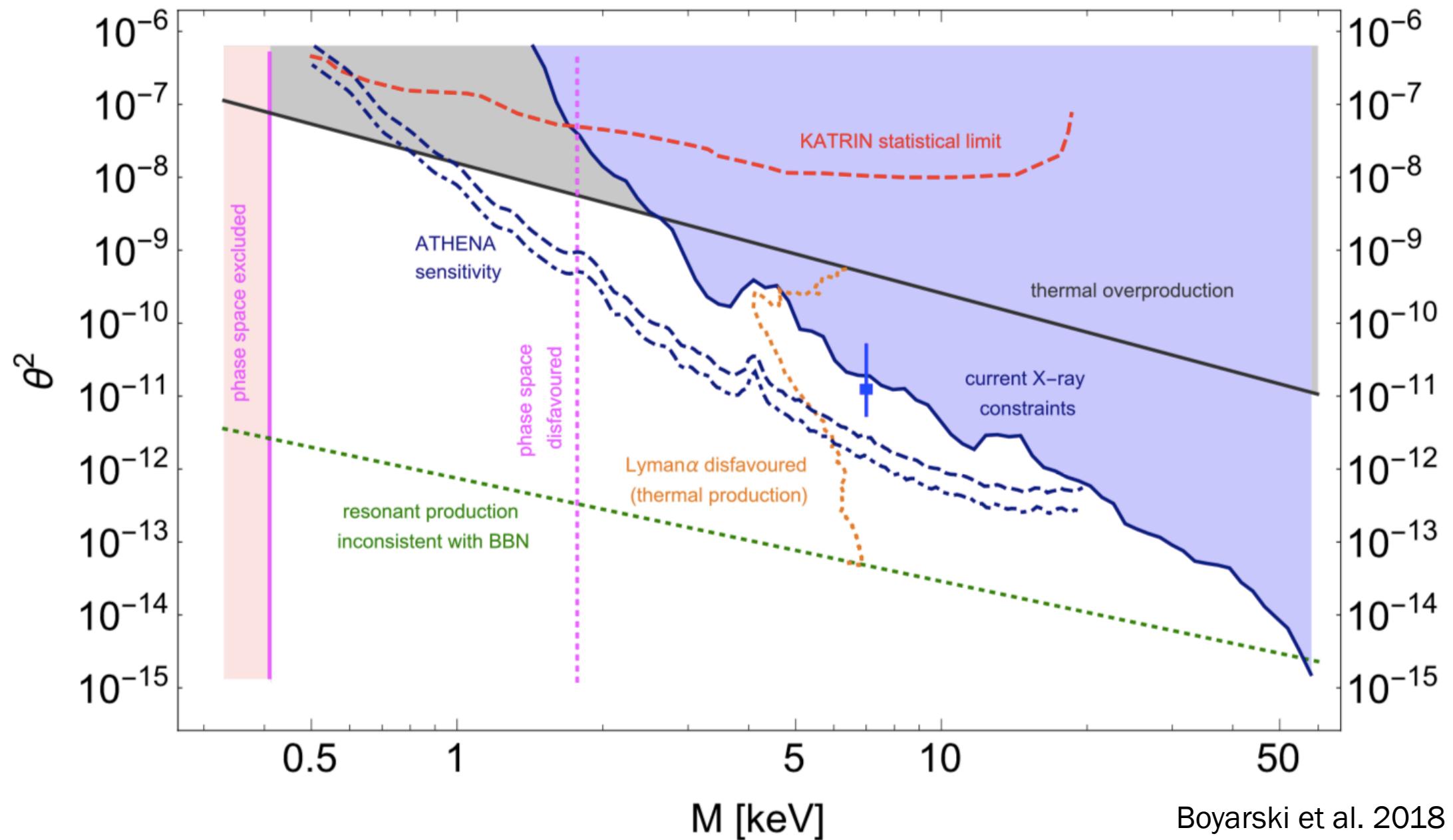
**COMMENT ON ‘DARK MATTER SEARCHES GOING BANANAS: THE CONTRIBUTION OF POTASSIUM
(AND CHLORINE) TO THE 3.5 KEV LINE’**

ESRA BULBUL (1), MAXIM MARKEVITCH (2), ADAM R. FOSTER (1), RANDALL K. SMITH (1), MICHAEL LOEWENSTEIN (2),
SCOTT W. RANDALL (1)

(1) Harvard-Smithsonian Center for Astrophysics, (2) NASA/GSFC

Draft version September 16, 2014





 ν_e  ν_μ  ν_τ ?