## 1 Wavy heliospheric current sheet

Both components of the drift velocity are modified by the waviness:

• There is a Heaviside step function in the Parker field:

$$\vec{B} = A_c B_0 \left(\frac{r_0}{r}\right)^2 \left(\hat{r} - \frac{\Omega r \sin \theta}{V_{sw}} \hat{\phi}\right) H(\theta' - \theta),$$

where  $\theta'$  is the angular extent of the HCS. In Strauss et al 2012 this is given by eq 11:

$$\theta' = \frac{\pi}{2} + \sin^{-1} \left[ \sin \alpha \sin \left( \phi - \phi_0(t) + \frac{\Omega r}{V_{sw}} \right) \right],$$

where  $\phi_0(t) = \phi_0 + \Omega t$  and  $\phi_0$  is an arbitrary phase. In the code, the step function appears when calculating the gradient/curvature part of the drift velocity.

## TODO:

- How is  $\phi_0$  determined?
- $-\alpha$  should also be time dependent. How should I implement this? Parameter file should probably point to table of  $\alpha$  vs t, which the parameter object can then parse. The date at which the particle was observed at Earth can be specified in the run file.
- The HCS drift velocity is also affected. It takes the form (eq 17)

$$\vec{v}_{hcs} = v_{hcs} \left[ \cos A_{\xi} \xi \sin \Psi \hat{r} + \sin A_{\xi} \xi \hat{\theta} + \cos A_{\xi} \xi \cos \Psi \hat{\phi} \right] A_{c} \operatorname{sgn} q.$$

The angle  $\xi$  ( $\beta$  in the reference) between the tangent to the HCS and the radial line passing through the point on the HCS is given by

$$\tan^2 \xi = \left[ -r \frac{\partial \theta'}{\partial r} \right]^2 + \left[ \frac{1}{\sin \theta} \frac{\partial \theta'}{\partial \phi} \right]^2$$

$$\implies \tan \xi = \frac{\Omega r}{V_{sw}} \frac{1}{\sin \Psi} \frac{\sqrt{\sin^2 \alpha - \cos^2 \theta'}}{\sin \theta'} \text{ (for Parker field)}.$$

The sign of  $\xi$ ,  $A_{\xi} = \pm 1$ , is equal to the sign of  $\partial \theta' / \partial r$ :

$$A_{\xi} = \operatorname{sgn} \cos \left( \phi - \phi_0(t) + \frac{\Omega r}{V_{syn}} \right).$$

 $v_{hcs}$  is given by the usual approximation:

$$v_{hcs} = \left[0.457 - 0.412 \frac{L}{r_L} + 0.915 \frac{L^2}{r_L^2}\right] v,$$

where L is the smallest distance from the particle's position to the HCS. If  $L > 2r_L$ ,  $v_{hcs}$  is zero. Finally,  $\xi$  is evaluated at the HCS point which minimizes L.

The hard part of these modifications is calculating L, which cannot be done analytically and is expensive. If

• 
$$|\theta - \frac{\pi}{2}| \le \alpha$$
,

- $\theta < \frac{\pi}{2} \alpha$  and  $L_{+}^{th} \leq 2r_{L}$ , where  $L_{+}^{th} = r\cos(\alpha + \theta)$  is the distance from the particle to the surface bounding the HCS above,
- Or  $\theta > \frac{\pi}{2} + \alpha$  and  $L_{-}^{th} = -r\cos(\theta \alpha) \leq 2r_L$ ,

we compute L using the Nelder-Mead simplex algorithm. Here are the steps for minimizing  $L(r_s, \phi_s)$ , where  $r_s$  and  $\phi_s$  are the coordinates of a point in the sheet:

- 0. Choose an initial set of points  $\vec{x}_1$ ,  $\vec{x}_2$ ,  $\vec{x}_3$ .
- 1. Order vertices:  $L(\vec{x}_1) \leq L(\vec{x}_2) \leq L(\vec{x}_3)$ .
- 2. Compute  $\vec{x}_O = \frac{1}{2}(\vec{x}_1 + \vec{x}_2)$ .
- 3. Compute reflected point  $\vec{x}_R = \vec{x}_O + \alpha(\vec{x}_O \vec{x}_3)$ . If  $L(\vec{x}_1) \leq L(\vec{x}_R) < L(\vec{x}_2)$ , replace  $\vec{x}_3$  with  $\vec{x}_R$  and go to step 1.
- 4. If  $L(\vec{x}_R) < L(\vec{x}_1)$ , compute expanded point  $\vec{x}_E = \vec{x}_O + \gamma(\vec{x}_R \vec{x}_O)$ . If  $L(\vec{x}_E) < L(\vec{x}_R)$ , replace  $\vec{x}_3$  with  $\vec{x}_E$  and go to step 1. Else, replace  $\vec{x}_3$  with  $\vec{x}_R$  and go to step 1.
  - Else, go to step 5.
- 5. Compute contracted point  $\vec{x}_C = \vec{x}_O + \rho(\vec{x}_3 \vec{x}_O)$ . If  $L(\vec{x}_C) < L(\vec{x}_3)$ , replace  $\vec{x}_3$  with  $\vec{x}_C$  and go to step 1. Else, go to step 6.
- 6. For all but the  $\vec{x}_1$ , replace  $\vec{x}_i$  with  $\vec{x}_1 + \sigma(\vec{x}_i \vec{x}_1)$ . Go to step 1.

Standard values for the constants are  $\alpha = 1$ ,  $\gamma = 2$ ,  $\rho = 1/2$ ,  $\sigma = 1/2$ .

The only tough part of this is choosing the initial set of points. Since  $\theta'$  is "periodic" in  $r \to r + 4\frac{\Omega}{V_{sw}}$ ,

$$\left\{ \left(r - 2\frac{\Omega}{V_{sw}}, \phi\right), \left(r + 2\frac{\Omega}{V_{sw}}, \phi\right), \left(r, \phi + \frac{\pi}{6}\right) \right\}$$

seem like good guesses. This will need to be tested, though.