

SN 2020jgb

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ABSTRACT

Keywords: keywords

1. INTRODUCTION

2. OBSERVATIONS

2.1. Detection and Classification

SN 2020jgb was first discovered by the Zwicky Transient Facility (ZTF; Bellm et al. 2019; Graham et al. 2019) on 2020 May 03.463 UT (MJD 58972.463) with the 48-inch Samuel Oschin Telescope (P48) at Palomar Observatory. The internal designation is ZTF20aayhacx. It was detected at a magnitude of 19.86 in ZTF *g*-band, and J2000 coordinates $\alpha = 17^{\text{h}}53^{\text{m}}12^{\text{s}}.651$, $\delta = -00^{\circ}51'21''.81$. The last non-detection was on 2020 April 27.477 (MJD 58966.477; 5.99 days before the first detection) up to a limiting magnitude of 20.7 in ZTF *r*-band.

Classification, ...

2.2. Optical Photometry

We obtained *gr*-band photometry of SN 2020jgb with the ZTF camera. A Galactic extinction of $E(B - V) = 0.404$ is reported by the maps of Schlafly & Finkbeiner (2011), for which we correct all our photometry using the extinction model proposed by Fitzpatrick (1999). We do not account for any additional host extinction due to the lack of any Na I D absorption in our spectra (**Is it in the outskirts?**).

2.3. Optical Spectroscopy

2.4. Near-infrared (NIR) Spectroscopy

We obtained one NIR spectrum of the transient using the Gemini near-infrared spectrometer (GNIRS; Elias et al. 1998) on the Gemini North telescope on 2020 June 9 (≈ 22 days after *r*-band peak), for an integration time of 2400 s. The spectra were reduced with the `PyPeIt` Python package (Prochaska et al. 2020; Prochaska et al. 2020).

3. ANALYSIS

Table 1. Spectroscopic Observations of SN 2020jgb

t_{obs}	Phase	Telescope/	R	Range	Air
(MJD)	(d)	Instrument	($\Delta\lambda/\lambda$)	(Å)	Mass
58,976.42	−9.7	P60/SED	100	3770–9220	1.23
58,982.12	−4.2	NOT/ALFOSC	360	4000–9620	1.17
58,990.43	+3.9	P60/SED	100	3770–9220	1.23
58,997.44	+10.7	P60/SED	100	3770–9220	1.29
58,998.00?	+11.2?	Shane/Kast	500?	3620–10720	
59,008.41	+21.3	P60/SED	100	3770–9220	1.28
59,010.00?	+22.9?	P200/DBSP	700	3200–9500	
59,023.58	+36.1	Keck I/LRIS	1100	3200–10250	2.04
59,107.29	+117.3	Keck I/LRIS	1100	3200–10250	1.31
59,143.26	+152.2	Keck I/LRIS	1100	3200–10250	2.16

NOTE—Phase is measured relative to $t_{r,\text{peak}}$ in the host galaxy rest frame. The resolution R is reported for the central region of the spectrum.

3.1. Photometric Properties

- sub-luminous
- first light time, peak time
- color evolution

3.2. Spectroscopic Properties

- infrared Ca II triplet (Ca II IRT)
- tentative He I absorption at $\approx 1 \mu\text{m}$

3.3. NIR spectra

In Figure 3, the NIR spectrum is presented along with three spectra in the sample of Marion et al. (2009) at similar phases. SN 2020jgb shows a strong absorption feature at $\approx 0.99 \mu\text{m}$, which is not seen in normal SNe Ia. This feature was still prominent two weeks later, as detected by LRIS on Keck (see Figure 5), though it was only partially covered due to the limitation of bandwidth. In general, SN 2020jgb highly resembles

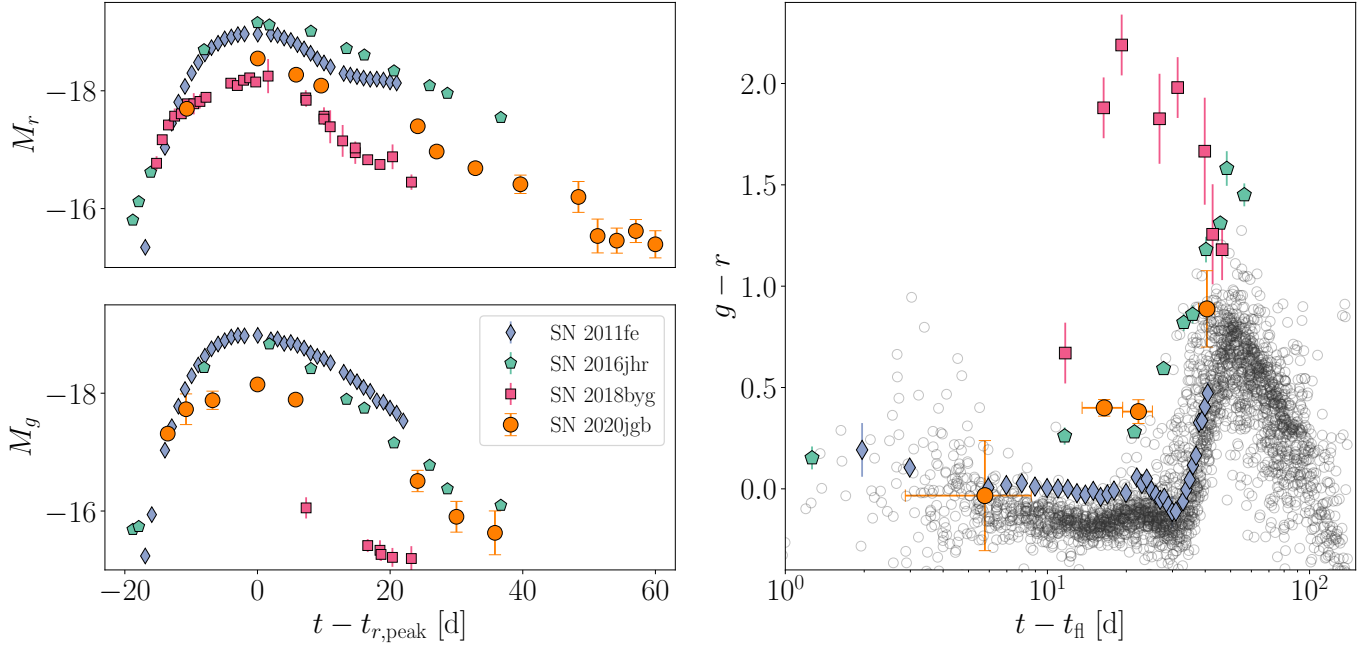


Figure 1. *Left:* comparison of the multi-color (g and r bands) light curves of SN 2020jgb to the normal SN Ia SN 2011fe and the He double detonation candidate SN 2018byg. *Right:* comparison of $g - r$ color evolution to SN 2011fe and SN 2018byg, as well as 62 normal SNe Ia (open circles) with prompt observations within 5 days of first light by ZTF (Bulla et al. 2020). The shaded region denotes the 1- σ credible interval of the color of SN 2020jgb until ≈ 40 days after the peak, estimated using Gaussian process.

normal SNe Ia in NIR band. The shape of the continuum redwards to $\approx 1.2 \mu\text{m}$ is significantly altered by line-blanketing of Fe-group elements synthesized in the SN interior, as opposed to the Fe-group elements in the outermost region as ashes of shell helium burning. Just like normal SNe Ia, SN 2020jgb also show enhancement of flux at about 1.3, 1.55, 2.0, 2.1, and $2.25 \mu\text{m}$, accompanied by several Co II absorption lines. It is especially similar to SN 2004da at +25 days after maximum as the increase in flux at $\approx 1.55 \mu\text{m}$, known as the H -band break, has become less prominent.

What is not seen in usual SNe Ia is the wide, deep absorption at $\approx 0.99 \mu\text{m}$, indicating its peculiarity. According to Marion et al. (2009), normal SNe Ia are nearly featureless in spectra around $1 \mu\text{m}$ a few weeks past the week. After ~ 3 weeks, some Fe II features (0.9998 & $1.0500 \mu\text{m}$) start to develop, but are not even nearly as prominent as the $0.99 \mu\text{m}$ feature. However, there are other He-shell double detonation candidates showing complexity in this region. In the currently small sample of five candidates, two objects (SN 2016jhr and SN 2019ofm) do not have any available NIR spectra, while the other three all exhibit strong absorption features near $1 \mu\text{m}$, though the spectra were obtained at different phases, as shown in Figure 4. The $1 \mu\text{m}$ feature for SN 2016hmk lies slightly redder than the other

two, which corresponds to an expansion velocity lower by $\sim 3,000 \text{ km/s}$, assuming they all have the same origin.

There are several elements that may be associated with this feature, none of which is fully satisfying. The most attractive possibility is the strong He I line at $1.0830 \mu\text{m}$, as has predicted in sub-Chandrasekhar-mass He-shell double detonation models when considerable amount of helium in the shell is left unburnt (Boyle et al. 2017). Figure 5 shows that the $1 \mu\text{m}$ feature, if associated with He I $\lambda 1.0830 \mu\text{m}$, has a high velocity ($\approx 26,000 \text{ km/s}$), yet similar as the HVF of Ca II IRT ($\approx 24,000 \text{ km/s}$). The expansion velocity in the ejecta is roughly linearly proportional to the radius, so such a high velocity indicates that both the Ca II IRT and the tentative He I absorption line form far outside the normal photosphere, which has a velocity of only $\approx 10,000 \text{ km/s}$. In the sense, the He-shell double detonation scenario, in which the unburnt helium locates at the outermost ejecta, is indeed supported.

Still, this helium detection remains skeptical, since other He I are not unambiguously detected, such as the He I $\lambda 2.0581 \mu\text{m}$. Considering a line velocity of $\approx 26,000 \text{ km/s}$ and a host galaxy redshift of 0.0307, this line will be blueshifted to $\approx 1.95 \mu\text{m}$ in the observer frame, so will be strongly blended by the strong telluric lines within $1.8\text{--}2.0 \mu\text{m}$. After telluric correction, the signal to noise

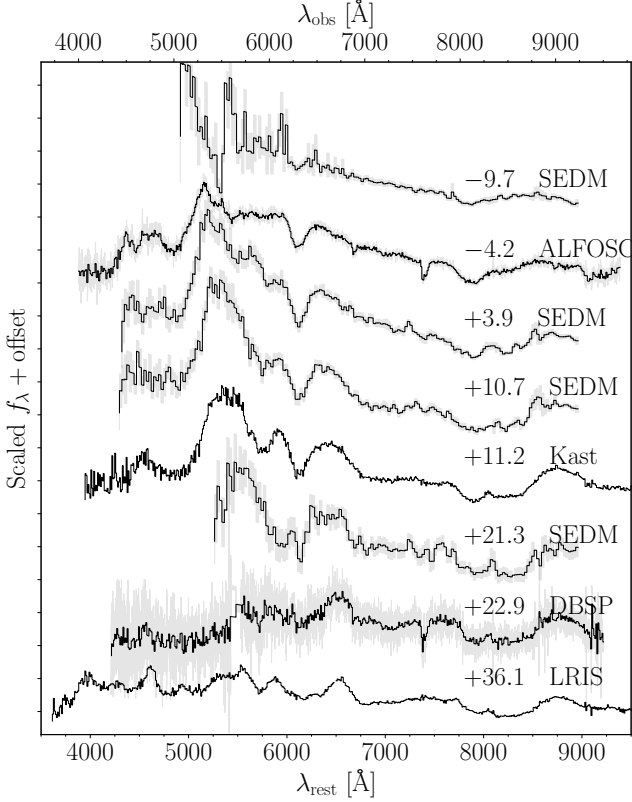


Figure 2. Optical spectroscopic sequence of SN 2020jgb. Rest frame phases (days) relative to the r -band peak and instruments used are posted next to each spectrum. The black curves are binned spectra with a bin size of 10 Å, except for the SEDm spectra, whose resolution is lower. The 1- σ uncertainties of raw spectra are shown in grey. Only regions with SNR > 3 after binning are plotted.

ratio reaches ~ 5 , with which we still cannot see any significant absorption feature. An upper limit of the equivalent width is determined to be $< 2\%$ of the $1.0830 \mu\text{m}$ line, while theoretically, the $2.0581 \mu\text{m}$ line is supposed to be only a factor of 6-12 weaker, depending on temperature (Marion et al. 2009). Another fact is that the $1 \mu\text{m}$ feature is as strong as the He I $\lambda 1.0830 \mu\text{m}$ in many helium-rich core-collapse supernovae, say, Type Ib supernovae, in which the He I $\lambda 2.0581 \mu\text{m}$ is weaker than the $1.0830 \mu\text{m}$ line yet still prominent (Shahbandeh et al. 2022). If the $1 \mu\text{m}$ feature is associated with He I, it would be very unusual if the $2 \mu\text{m}$ feature is not seen at all, even if somehow blended by the telluric lines.

Other possibilities include the Mg II $\lambda 1.0927 \mu\text{m}$, the C I $\lambda 1.0693 \mu\text{m}$, and the Fe II $\lambda 1.0500 \mu\text{m}$ & $\lambda 1.0863 \mu\text{m}$. The Mg II $\lambda 1.0927 \mu\text{m}$ is prevalent in the NIR spectra of SNe Ia, but usually disappears within a week

after the peak luminosity (Marion et al. 2009), while our spectra were obtained three weeks/a month after the peak. A stronger Mg II line at $0.9227 \mu\text{m}$ is not detected either. Also, the problematically high radial velocity of $\approx 30,000 \text{ km/s}$ is not seen in normal SNe Ia, and is over 20% faster than the HVF of Ca II IRT at the same phase.

The C I $\lambda 1.0693 \mu\text{m}$ line from the unburnt carbon is much less frequently seen than the Mg II $\lambda 1.0927 \mu\text{m}$. Hsiao et al. (2019) presented a sample of five SNe Ia with C I detections, showing the C I feature is strongest for those fainter, fast-declining objects. However, in their sample, the C I feature is always accompanied by the stronger Mg II $1.0927 \mu\text{m}$ line, and is a pre-maximum feature which fades away as the luminosity peaks. The required expansion velocity for C I is $\approx 22,000 \text{ km/s}$, which is consistent with the HVF of Ca II IRT, but still overwhelmingly faster than the estimated velocity for the sample in Hsiao et al. (2019) ($\sim 10,000$ - $12,000 \text{ km/s}$).

The Fe II features in SNe Ia usually start to develop three weeks after the peak, about the same phase when we got our GNIRS spectrum. Two Fe II lines, $\lambda 0.9998 \mu\text{m}$ and $\lambda 1.0500 \mu\text{m}$, are actually visible on each of the wing of the $1 \mu\text{m}$ feature. The Fe II $\lambda 1.0863 \mu\text{m}$ line is not yet seen. They correspond to an expansion velocity of $\approx 8,000 \text{ km/s}$, which is consistent with the PVF of the Ca II IRT at the same epoch, and most matches the same two lines for normal SNe Ia (Marion et al. 2009), making the identification more reliable. Obviously, these two Fe II features are wider and shallower than the strong feature between them. We fit the whole feature with three Gaussian profiles, two of which are set to be the blueshifted Fe II $\lambda 0.9998 \mu\text{m}$ and $\lambda 1.0500 \mu\text{m}$. We find that the shallower and wider Fe II lines contributes to only $\sim 40\%$ of the total equivalent width, and the rest $\sim 60\%$ comes from the central feature, which cannot be accounted for by any Fe II feature at the same velocity. Given the similarity of the Fe-group line blanketing between the GNIRS spectrum with the spectrum of SN 2004da at +25 days, the distribution of Fe-group elements inside each supernova ejecta should be somehow similar, so the central region of the $1 \mu\text{m}$ feature is not likely to be associated with Fe II.

4. HOST GALAXY

5. MODEL COMPARISONS

6. DISCUSSION AND CONCLUSION

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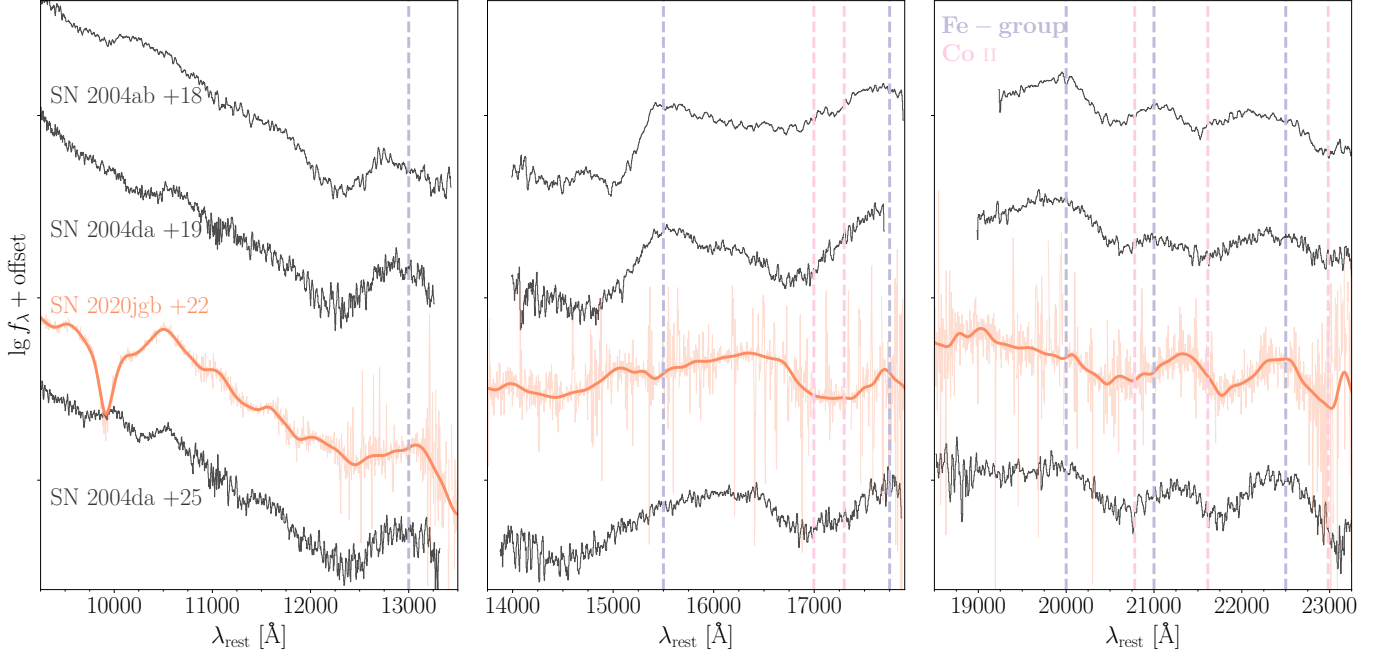


Figure 3. The NIR spectra of SN2020jgb and two SNe Ia with normal maximum luminosity (SN 2004ab and SN 2004da, Marion et al. 2009), taken about three weeks after the peak. For each spectrum, the continuum at $\gtrsim 1.2 \mu\text{m}$ is significantly reshaped by the Fe-group blanketing (emission features, blue vertical lines) and Co II absorption (pink vertical lines).

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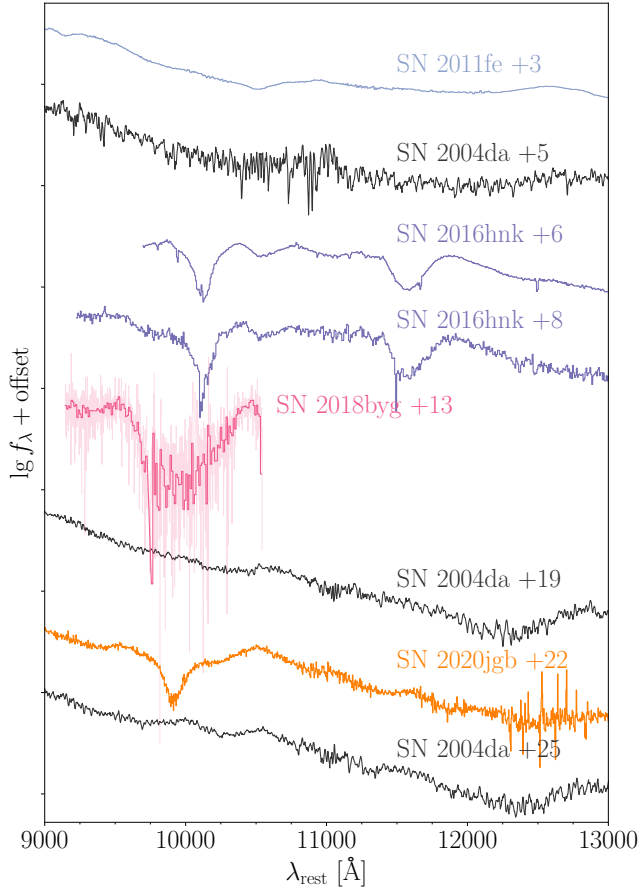


Figure 4. The NIR spectra (9,000 to 13,000 Å) of a few normal SNe Ia (SN 2011fe and SN 2004da) and three He-shell double detonation candidates, which are all subluminous SNe Ia (SN 2016hmk, SN 2018byg, and this source, SN 2020jgb).

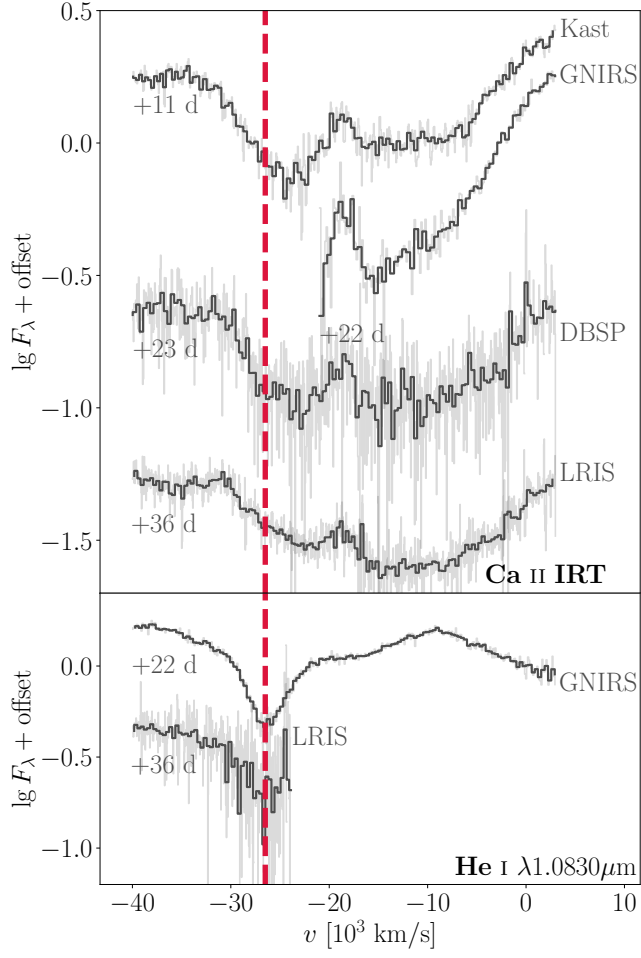


Figure 5. Spectra in the velocity space, comparing the high-velocity component of Ca II IRT and the absorption feature at $\approx 0.99 \mu\text{m}$ assuming it is associated with He I at $1.0830 \mu\text{m}$.

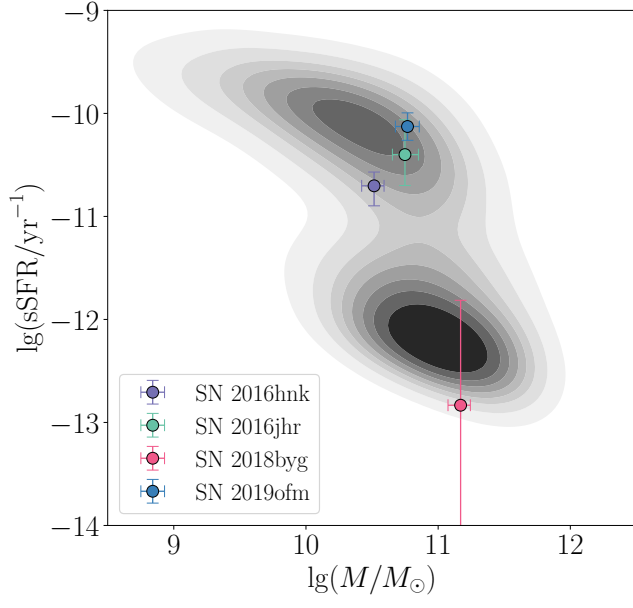


Figure 6. The specific star formation rate (sSFR) and the galactic mass for the host galaxies of He-shell double detonation candidates.