

The Determination of Vehicle Drag Contributions from Coast-Down Tests

R. A. White and H. H. Korst
Dept. of Mechanical and Industrial Engineering,
University of Illinois

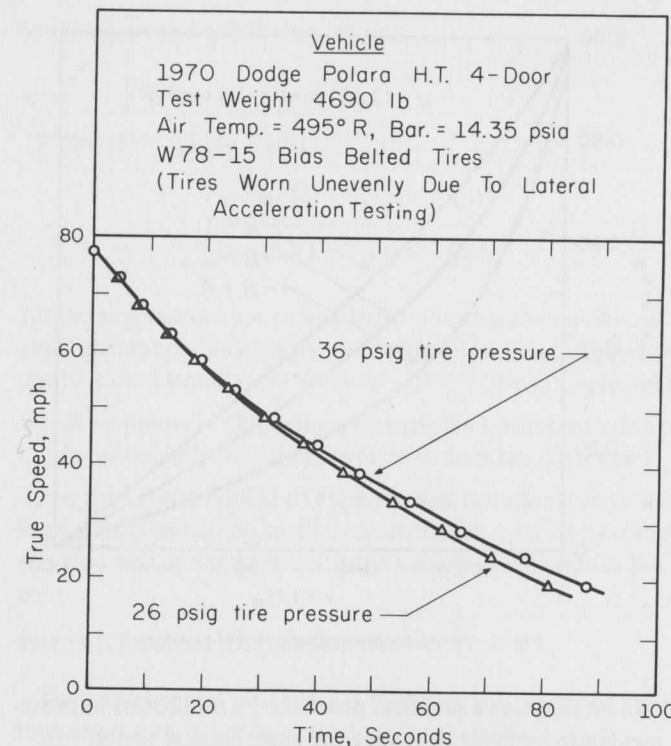


Fig. 1 - Velocity-time coast-down test

$$-\frac{d}{dt}(\text{kinetic energy}) = \frac{-v(m + \Delta m)}{g_c} \frac{dv}{dt} = \frac{1}{2g_c} \rho v^3 AC_d + vR \quad (1)$$

It is important to note that the mass term contains here an additive (Δm) term to account for the effective inertia mass of the rotating components (primarily the wheels).

Without loss of generality, one may separate the variables, obtaining

$$-\frac{dv}{\frac{\rho AC_d v^3}{2m_{eff}} + \frac{Rg_c}{m_{eff}}} = dt \quad (2)$$

where $m_{eff} = m + \Delta m$, and Δm is constant when the kinetic energy of the rotating masses can be directly related to the translational speed of the vehicle.

One notes now that recently the automotive industry has adopted the bias-belted tires as standard equipment while a few makes appear with radials. These types of tires have a markedly constant rolling resistance up to approximately 70 mph, limiting the variation to approximately 18% (Fig. 2). The standard bias-construction low-profile tire which is becoming popular has an even smaller change in this range, namely about 15%. Thus, if the coast-down tests are made for speeds below that which the major rise in rolling resistance occurs, (approximately 60-80 mph depending on the type and style of tire) (10, 11), the rolling resistance can be taken as a constant to a good approximation as will be borne out by the

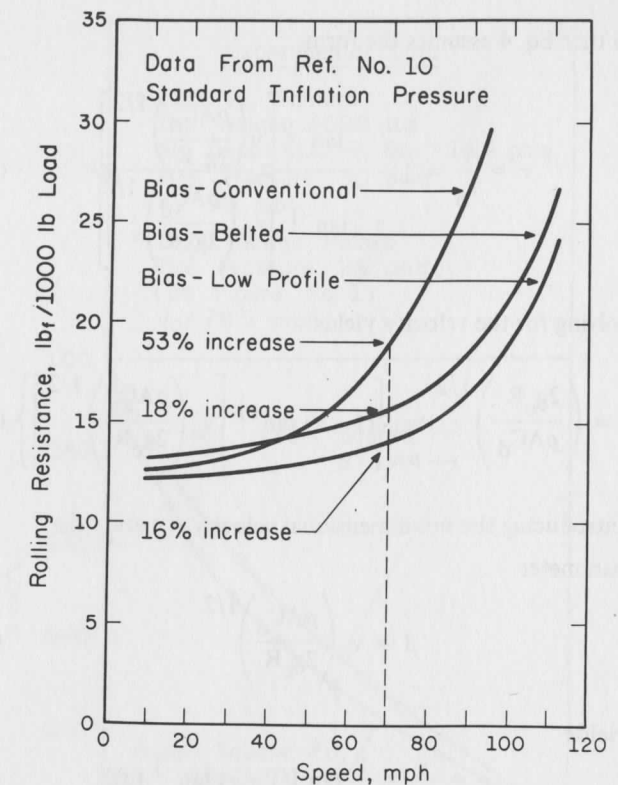


Fig. 2 - Rolling resistance of time versus speed

analysis. In addition, the assumption of a constant aerodynamic drag coefficient is reasonable due to the very large Reynolds numbers encountered. Consequently, Eq. 2 can now be integrated to yield

$$-\left[\frac{2m_{eff}^2}{\rho AC_d g_c R}\right]^{1/2} \tan^{-1} \left[v \left(\frac{\rho AC_d}{2g_c R} \right)^{1/2} \right] = t + \text{constant} \quad (3)$$

The constant of integration can be evaluated as $v = v_0$ for $t = 0$ (v_0 is the initial speed at start of coast down). Thus, one obtains

$$t = \left[\frac{2m_{eff}^2}{\rho AC_d g_c R} \right]^{1/2} \left\{ \tan^{-1} \left[v_0 \left(\frac{\rho AC_d}{2g_c R} \right)^{1/2} \right] - \tan^{-1} \left[v \left(\frac{\rho AC_d}{2g_c R} \right)^{1/2} \right] \right\} \quad (4)$$

One nondimensionalizes the time variable t by the total coast-down time t_0 , where $v = 0$ (which shall be subjected to further scrutiny)

$$t_0 = \left[\frac{2m_{eff}^2}{\rho AC_d g_c R} \right]^{1/2} \tan^{-1} \left[v_0 \left(\frac{\rho AC_d}{2g_c R} \right)^{1/2} \right] \quad (5)$$

THE PROBLEM OF AERODYNAMIC and rolling resistance characteristics of cars and trucks is of considerable importance to vehicle engineers as the two major contributions to external vehicle drag. Numerous techniques have been developed for their experimental determination each with particular advantages and disadvantages. For aerodynamic purposes, the methods include wind-tunnel testing of scale models (1-3)*, the testing of full-size production cars (4, 5), and coast-down testing (6-8). Wind-tunnel testing, while well established in the aircraft industry, when applied to determining the aerodynamic drag characteristics of both models and full-size cars has raised many questions of interpretation of data due to ground plane simulation problems (2, 9) and model scale effects (Reynolds number, boundary layer transition, and separation). These limitations imposed on accuracy of results, together with the inherently high costs of wind-tunnel testing are certainly severe drawbacks.

On the other hand, the coast-down technique, which avoids these problems, is affected by the limited accuracy germane to the finding of derivatives to experimentally established curves (7) and lumps together the effects of aerodynamic and other (mostly rolling) resistances. The methods for the determination of rolling resistance are primarily those of the

*Numbers in parentheses designate References at end of paper.

ABSTRACT

The problem of aerodynamic and rolling resistance characteristics of cars and trucks is of considerable importance to vehicle engineers as the two major contributions to external

testing of individual tires on drums or tow rigs, or the actual towing of full-size cars in a protective box (9, 10).

If these shortcomings of the coast-down technique can be overcome, its simplicity and inexpensiveness make it most attractive. The presently proposed method is based on the mathematical analysis of a simplified dynamic model which not only allows the separation of aerodynamic and rolling resistance forces, but also utilizes the closed mathematical form of the solution to eliminate the need for differentiating an experimentally determined data curve.

A representation of the coast-down process in terms of dimensionless variables also establishes a dimensionless parameter which characterizes the resistance behavior of any given vehicle. Furthermore, vehicles of considerably different design can be correlated by a single normalized coast-down function.

ANALYSIS

In typical coast-down tests, data can be presented in the form of velocity-time curves similar to Fig. 1; it is from such information that our analysis shall now proceed.

In the absence of a propulsive force, the time rate of decrease of kinetic energy (both translational and rotating) is attributed to the power absorbed by aerodynamic drag and rolling resistance R , respectively.

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so that Eq. 4 assumes the form

$$\tau = \frac{t}{t_0} = 1 - \frac{\tan^{-1} \left[v \left(\frac{\rho A C_d}{2g_c R} \right)^{1/2} \right]}{\tan^{-1} \left[v_0 \left(\frac{\rho A C_d}{2g_c R} \right)^{1/2} \right]} \quad (6)$$

Solving for the velocity yields

$$v = \left(\frac{2g_c R}{\rho A C_d} \right)^{1/2} \tan \left\{ (1 - \tau) \tan^{-1} \left[v_0 \left(\frac{\rho A C_d}{2g_c R} \right)^{1/2} \right] \right\} \quad (7)$$

Introducing the nondimensional velocity $v = v/v_0$, the

parameter

$$\beta = v_0 \left(\frac{\rho A C_d}{2g_c R} \right)^{1/2} \quad (8)$$

yields

$$v = \frac{v}{v_0} = \frac{\beta}{1} \tan \left[(1 - \tau) \tan^{-1}(\beta) \right] \quad (9)$$

Eq. 9 represents a one-parameter family, in (β) , of curves as

shown in Fig. 3, in the dimensionless v - τ plane.

By comparing test data for any vehicle with the curves in Fig. 3, the appropriate β value can be obtained. With knowl-

edge of the total coast-down time, it is now possible to write simple algebraic expressions for the rolling resistance and drag

coefficient. These are, as obtained from Eqs. 5 and 8,

$$C_d = \frac{2m_{eff} \beta \tan^{-1}(\beta)}{v_0 \rho A} \quad (10)$$

and

$$R = \text{rolling resistance} = \frac{v_0 m_{eff} \tan^{-1}(\beta)}{\beta t_{0gc}} \quad (11)$$

Eqs. 9-11 allow determination of the important characteristics

of the vehicle in a straightforward manner from a simple coast-

down test. It is also obvious that a comparison of vehicles on

the basis of their drag characteristics can be made by examini-

ng β values. Furthermore, the coast-down behavior of all cars

(for that matter, all vehicles) can be represented by a single

curve in the nondimensional plane by the selection of a

normalized value; for example, $\beta = 1$, * through the proper

choice of v_0 .

*One notes that a higher value of β should be selected if it is

desirable to correlate cars with relatively high aerodynamic

drag contribution, or for covering coast-down tests from high

initial speeds.

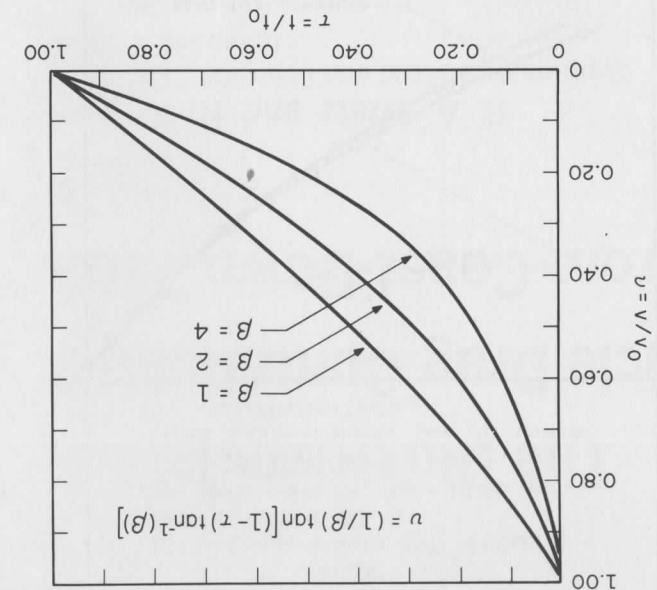


Fig. 3 - Nondimensional velocity-time plane

In establishing a practical procedure, a number of improve-
ments can be made to reduce the experimental scatter for
obtaining t_0 without direct measurement, and for more

accurately identifying the parameter β .

Applying the method of least squares to Eq. 9 for determin-

ing optimal values of β and t_0 requires that

$$\frac{\partial}{\partial \beta} \sum [v_i - v(\beta, \tau_i)]^2 = 0 \quad (12a)$$

$$\frac{\partial}{\partial t_0} \sum [v_i - v(\beta, \tau_i)]^2 = 0 \quad (12b)$$

Carrying out the operations of Eq. 12 yields

$$\beta^2 - \beta \left\{ \frac{\sum v_i + \cos^2 [\tan^{-1}(\beta)] \sum v_i^2}{\sum v_i \cos^2 [\tan^{-1}(\beta)] \sum v_i^2} \right\} \frac{\sum v_i^2}{\sum v_i^2} + \frac{\cos^2 [\tan^{-1}(\beta)] \sum v_i^2}{\sum v_i^2} = 0 \quad (13)$$

where:

$$\sum v_i = \sum v_i \tan [(1 - \tau_i) \tan^{-1}(\beta)]$$

$$\sum v_i^2 = \sum v_i^2 \tan^2 [(1 - \tau_i) \tan^{-1}(\beta)]$$

$$\sum v_i^3 = \sum \frac{v_i^3 (1 - \tau_i)}{\cos^2 [(1 - \tau_i) \tan^{-1}(\beta)]}$$

$$\sum v_i^4 = \sum \frac{v_i^4 (1 - \tau_i)}{\cos^2 [(1 - \tau_i) \tan^{-1}(\beta)]}$$

$\tau_i = t_i/t_0$ ratios.

*Note that the unknown t_0 appears implicitly in the

where m is the actual mass of the car.

$$F \left| 1000 = \frac{R_{gc}}{mg} 1000 \quad (15)$$

by

The rolling resistance as calculated from Eq. 11 can be
brought into the conventional form of force/1000 lbf weight

density ρ .
from Eq. 10, it is necessary to determine the effective mass

(m_{eff}) and the frontal area of the vehicle A , as well as the air
seen from Eqs. 10 and 11. To obtain the drag coefficient C_d

of Fig. 4. It does, however, directly effect the calculation of
both aerodynamic drag coefficient and the rolling resistance as

t_0 value is not explicitly seen in the nondimensional (v - τ) plot
The apparent usefulness of obtaining an accurate (calculated)

test in favor of the remainder.
to avoid overemphasizing the importance of any portion of the

a relatively uniform spacing of the experimental data is used,
fit approach used in Eqs. 13 and 14, care should be taken that

"weighting" scheme has been introduced into the least square
square fit" theoretical curve is most gratifying. Since no

agreement between experimental data points and the "least
solution corresponding to Eq. 9. As can be seen in Fig. 4, the

analytically determined β value then establishes the theoretical
in Fig. 1 by a dimensionless $v = v/v_0$ versus $\tau = t/t_0$ plot. The

one can now replace the original velocity-time diagram shown
simultaneously for β and t_0 . With t_0 being thus determined,

of velocity-time value pairs, Eqs. 13 and 14 can be solved
approximately 75-20 mph. After obtaining a representative set

time relationship during coast down of the vehicle from
The experimental procedure is to determine the speed versus

PRESENTATION OF RESULTS
zero.

large effects caused by small irregularities in road surface and
down time is susceptible to experimental variation due to the

t_0 is a particularly useful improvement since the total coast-
digital computer). Obtaining a statistically optimized value of

can be solved simultaneously for β and t_0 * (for example, on a
experimental points in every coast-down test. Eqs. 13 and 14

$$\sum \frac{v_i^2 \tau_i}{\cos^2 [(1 - \tau_i) \tan^{-1}(\beta)]} - \frac{\sum v_i \tau_i \tan [(1 - \tau_i) \tan^{-1}(\beta)]}{\sum \frac{v_i^2 \tau_i}{\cos^2 [(1 - \tau_i) \tan^{-1}(\beta)]}} = 0 \quad (14)$$

Similarly, from Eq. 12b one obtains

VEHICLE DRAG CONTRIBUTIONS

CONCLUSIONS

cannot be fully appraised.

(12) can be fully resolved by our comments is questionable.
Yet, the accuracy of the experimental method reported in (12)

between our C_d values for runs 6 and 7 and that reported in
of runs 6 and 7). Whether the somewhat larger discrepancy

accurate in isolating aerodynamic drag and rolling resistance
most data from other sources (runs 2-5), and sufficiently

method of analysis appears to be in reasonable agreement with
Altogether, the information resulting from the present

tions affecting our experimental results.
Comments have been included to expose particular condi-

Results obtained from coast-down tests with a variety of
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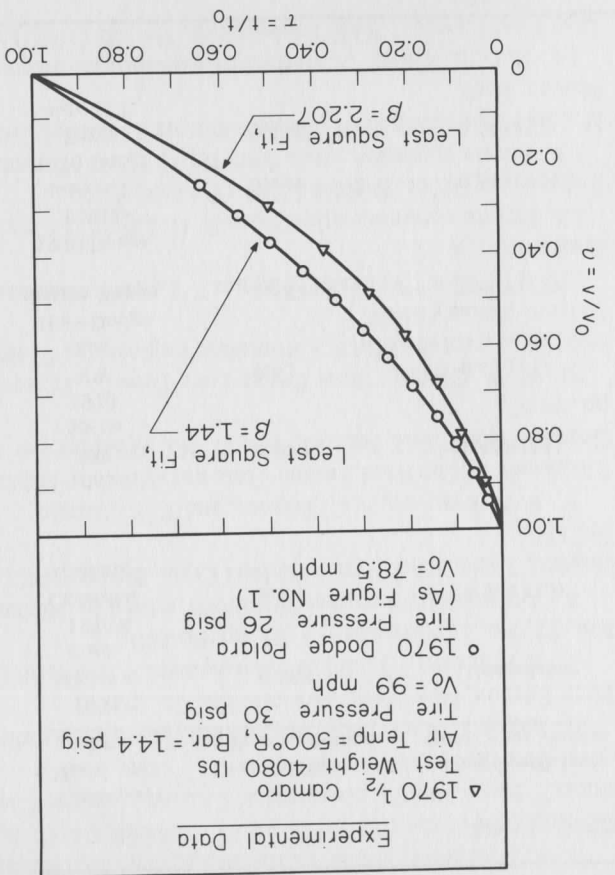
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Fig. 4 - Comparison of experimental data and least square fit in dimensionless velocity-time plane



Furthermore, it follows from the analysis, that coast-down
optimization procedure.

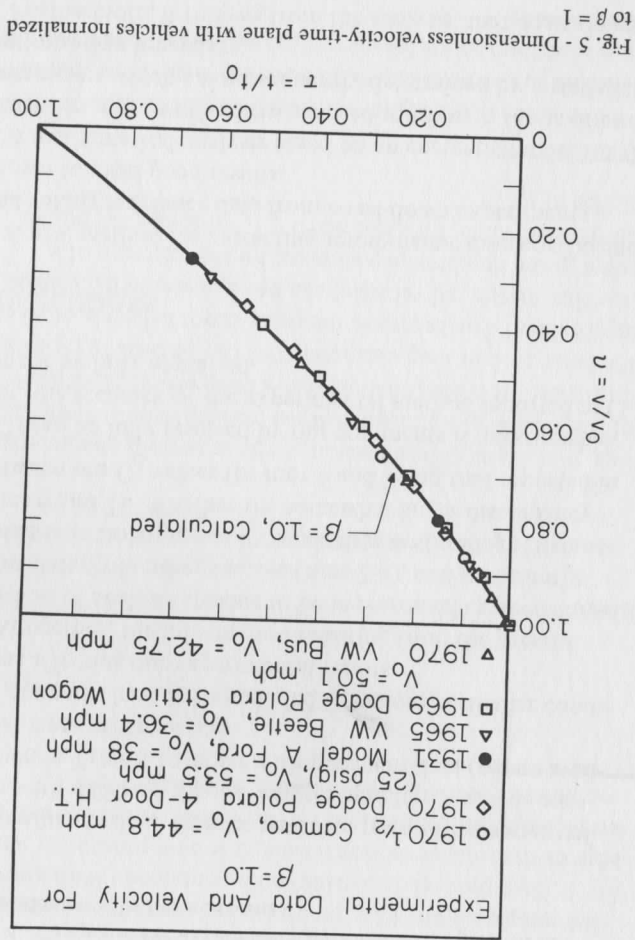
A mathematical analysis based on an energetic model for the
coast-down process leads to a closed solution in terms of two

parameters which are subsequently determined by a statistical
shown to yield good results.

A new method for extracting aerodynamic drag coefficients
and rolling resistance data from coast-down experiments is

Table 1 - Comparison of Drag Contribution

Car Make, Model, and Year	Drag Coefficient by Present Method	Drag Coefficient from Other Sources	Rolling resistance, 1000 lb load, and Tire Pressure	V_0 mph	β	Comments
1931 Ford Model A	0.838	Unavailable	29.2 lb/1000 at 25 psig	50.0	1.314	Drugging mechanical brakes account for the high rolling resistance
1970 1/2 Chevrolet Camaro	0.491	0.521 (13)	12.1 lb/1000 at 30 psig	99.0	2.207	Mechanical transmission, F-70 belted bias construction tires. Ref. 13 C_D value is from full-scale wind tunnel tests
1965 VW Beetle 1970 VW Bus	0.450	0.458 (14)	12.9 lb/1000 at 26 psig	68.7	1.886	Bias construction tires on this vehicle
1969 Dodge Station Wagon	0.587	0.607 (12)	18.6 lb/1000 at 25 psig	78.5	1.564	Car equipped with studded snow tires and automatic transmission
1970 Dodge 4-dr. H.T. Polara	0.499	0.543 (12)	16.35 lb/1000 at 35 psig	77.2	1.541	Car equipped with tires worn unevenly from lateral acceleration tests. Only variable is tire pressure, showing separation of rolling resistance and aerodynamic drag
1970 Dodge 4-dr. H.T. Polara	0.498	0.543 (12)	18.7 lb/1000 at 25 psig	77.2	1.441	

Fig. 5 - Dimensionless velocity-time plane with vehicles normalized to $\beta = 1$

behavior can be correlated by a single, nondimensionalized speed-time curve. Experiments conducted with a variety of different vehicles (but also with fans and compressors) have verified this, as shown in Fig. 5.

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NOMENCLATURE

A	=	frontal area of vehicle, ft^2
D	=	aerodynamic drag force, lbf
g	=	gravitational acceleration, ft-s^{-2}
g_c	=	gravitational constant, $\text{lbm-lbf}^{-1}\text{-ft-s}^{-2}$
m	=	mass of vehicle, lbm
m_{eff}	=	effective mass of vehicle, lbm
v	=	velocity ft-s^{-1}
V_0	=	initial velocity, ft-s^{-1}
V_0	=	initial velocity, rolling distance, mph
R	=	rolling resistance, lbf

VEHICLE DRAG CONTRIBUTIONS

t	=	time, s
t_0	=	total coast-down time from V_0 to $V = 0$, s
Δm	=	additive mass, lbm
ρ	=	air density, lbm-ft^{-3}
<i>Dimensionless Groups</i>		
v	=	dimensionless velocity, V/V_0
τ	=	dimensionless time, t/t_0
C_D	=	aerodynamic drag coefficient, $2Dg_c/AV^2\rho$
$F/1000$	=	dimensionless rolling resistance (Eq. 15)
β	=	coast-down parameter, Eq. 8
<i>Subscripts</i>		
i	=	dimensionless data point pairs of v_i and τ_i

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