

## **BASELIUS COLLEGE KOTTAYAM**

#### **DEPARTMENT OF PHYSICS**

## **PROJECT REPORT**

# CALCULATING HUBBLE CONSTANT USING PANTHEON TYPE Ia SUPERNOVAE DATASET

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## Calculating Hubble Constant Using Pantheon Type Ia Supernovae Dataset

A Project Report Submitted for the Partial Fulfillment of the requirements for the degree of Master of Science

in

#### **Physics**

by

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Under the guidance of

Dr. Moncy V John



DEPARTMENT OF PHYSICS
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#### **CERTIFICATE**

This is to certify that we have examined the project entitled "Calculating Hubble Constant Using Pantheon Type Ia Supernovae Dataset", submitted by Adarsh T Saji (Roll Number: 200011012171), a postgraduate student of Department of Physics in partial fulfillment for the award of degree of Masters of Science with specialization in Physics. We hereby accord our approval of it as a study carried out and presented in a manner required for its acceptance in partial fulfillment for the post graduate degree for which it has been submitted. The project has fulfilled all the requirements as per the regulations of the institute and has reached the standard needed for submission.

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I, Adarsh T Saji, hereby declare that the project entitled "Calculating Hubble Constant Using Pantheon Type Ia Supernovae Dataset" submitted in partial fulfillment of the requirements for the award of the degree of Master of Science from Mahatma Gandhi University Kottayam, is a bonafide record of the work carried out by myself under the guidance and supervision of Dr. Moncy V John Visiting Professor, School of Pure and Applied Physics, Mahatma Gandhi University, Kottayam

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Place: Kottayam

Date: 22nd September 2022

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#### **ABSTRACT**

The Pantheon Type Ia supernovae catalogue Scolnic et.al(2018) data set which contains the data of 1048 supernovae acts as the base for this dissertation. Calculation of the Hubble constant using this data is our primary aim. We also look into the existing discrepancy in the values of  $H_0$  which is called as Hubble tension. We also examine more closely at the many models that have been developed in cosmology. In this analysis we classify the models using the deceleration parameter  $(q_0)$  [measure of the cosmic acceleration of the expansion of space in a Friedman-Lemaitre-Robertson-Walker universe]. Using he distance modulus equation for each model we calculate the  $H_0$  value. The likelihood of occurrence of an obtained  $H_0$  value is then measured using the Bayesian theorem. Error estimation of the values are also done.

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## Chapter 1

## Introduction

Cosmology is a branch of astronomy that deals with the origin, structure, and space-time relationships of the universe as a mathematically driven science, cosmological physics is usually thought to be extremely precise. But it has its downsides. One being the difficulty to evaluate Hubble's constant precisely. The Hubble constant plays a crucial part in cosmology as it describes how fast the universe is expanding, which in turn helps in calculating the universe's age and history. Modern scientific cosmology is usually considered to have begun in 1917 with Albert Einstein's publication of his final modification of general relativity in the paper "Cosmological Considerations of the General Theory of Relativity". General theory of relativity prompted all the major cosmogonists to explore its possible astronomical ramifications, thereby enhancing the ability to explore distant bodies. On 1929, Edwin Hubble calculated this constant from star measurements. Hubble used his telescope experience to make measurements of Cepheid variables. He used the work of fellow astronomer Henrietta Leavitt to predict the brightness of these stars, which enabled him to calculate their distance from earth. These measurements confirmed that the universe extends far beyond the milky way. This is regarded as the most significant cosmological milestone. But the problem is, even after 93 years of its introduction, scientists are still puzzled about its exact value. The problem is that having two different figures for the Hubble constant measured from different perspectives would simply invalidate the cosmological model. Cosmologists using the plank satellite to study the cosmic microwave background and have arrives at a high precision value of of the expansion rate which is called as the Hubble constant. Another set of astronomers observing stars and galaxies have also obtained a value of  $H_0$  however the two values disagree. This discrepancy is called as Hubble tension.

#### 1.1 Hubble Tension

Cosmology is a branch of astronomy that studies the origin, structure, and spacetime relationships of the universe. Because it is a mathematically driven science, cosmological physics is widely regarded as extremely precise. However, astronomers have reached a fundamental stumbling block in their understanding of the universe: they cannot agree on how fast it is expanding. This is now the hottest and most mysterious research topic in cosmology.

Scientists first realized the universe was expanding in the 1920s, when the US astronomer Edwin Hubble found that the greater the distance between two galaxies, the faster they are moving apart. Hubble constant, a constant of proportionality in the relation between the velocities of remote galaxies and their distances. It expresses the rate of expansion of the universe. It remains one of the most important scientific discoveries ever made. But, even if the universe was expanding at an increasing rate, one key question remained: what is the precise rate of this expansion? Just how quickly is the cosmos flying apart? To be more specific, what exactly is the value of the Hubble constant? It is a highly valuable and highly sought-after value because it will reveal much about the origin, age, evolution, and, ultimately, fate of the universe. But even after 93 years of its introduction, scientists are still puzzled about its exact value.

The problem is that having two different figures for the Hubble constant measured from different perspectives. The first method is to use astronomical measurements to examine nearby objects and determine how fast they are moving. This is a local approach.

The other method for determining the Hubble constant has involved astronomers studying the rippling pattern of light known as the cosmic microwave background, which were formed just after the universe's big bang 13.8 billion years ago. These findings demonstrate how the early universe's expansion would have most likely resulted in an expansion that astronomers can now measure.

Until recently, these two approaches produced estimates that appeared to be consistent with each other, despite the fact that both measurements had significant uncertainties. Everyone believes that as the two values are tested with increasing precision, the differences between them will disappear. Unfortunately for astronomers hoping for a quick fix, this has not occurred. In reality, the inverse has occurred. The disparity has grown more pronounced. This disparity, known as the Hubble tension, has been growing for years, as study after study of both the early and late universe yields ever more precise results, leaving scientists on both sides worried and perplexed. After all, either faction could be measuring the universe incorrectly. However, the tension could be a true reflection of reality, necessitating exotic new physics and a dramatic revision of our understanding of cosmic evolution. The lower estimate of the Hubble constant has gotten a little lower over the years,

while the higher estimate has gotten even higher. Today, those who use cosmic background data to calculate the Hubble constant get a value of 67.4 plus or minus 0.5. By contrast the local approach gives a figure of 73.5 plus or minus 1.4.

The dissimilarity may not sound great but it is significant. This isn't just a case of two experiments disagreeing. We're measuring something entirely different. One is a measurement of how quickly the universe is expanding as we see it today. The other is a prediction based on early universe physics and measurements of how fast it should be expanding. If these values do not agree, it is very likely that we are missing a factor in the cosmological model that connects the two eras. In short, something appears to be missing from our understanding of the universe, and the Hubble constant has become the focus of a heated debate over the nature of this intangible influence.

Changing the Hubble constant from 67.4 to 73.5 implies that it must have been flying apart faster than previously assumed, implying that it is younger than the currently accepted age of 13.8 billion years. In fact, it would shorten the time to 12.7 billion years. And this does cause issues. The universe contains some very old stars with estimated ages of around 12 billion years. After all, stars take a long time to form.

There have been over 300 proposals for solutions to the cosmology crisis to date. Some argue for more physics in the CMB era. Some claim that dark energy did something strange in the recent past. Some fundamentally alter physics, interfering with our observations of supernovae. However, no single proposal can account for the wealth of cosmological evidence, and there is no consensus on a solution. The fact of concern has been that discrepancy has been increasing with the number of studies being conducted on both the early and the late universe. So we wouldn't be able to say what the age of the universe was until we had put our physics right.

#### 1.2 Cosmological Models

Let's first discuss some of the major cosmological models. A cosmological model describes the universe's largest-scale structures and dynamics and allows researchers to investigate fundamental questions about the universe's origin, structure, evolution, and ultimate fate.

#### 1.2.1 Einstein's Static Model

Einstein used three assumptions that were outside the scope of his equations to derive his 1917 cosmological model. The first assumption was that the universe is homogeneous and isotropic in general (i.e., the same everywhere on average at any instant in time). The second assumption was that the universe was homogeneous and isotropic, with a closed spatial geometry. Einstein's third assumption was that

the universe as a whole is static, meaning that its large-scale properties do not change over time. This assumption was made prior to Hubble's observational discovery of the universe's expansion. As a result, it is known as the Einstein Static Model. It should be noted that this model is unrealistic in light of Hubble's law.

#### 1.2.2 de-Sitter's Model

In 1917, the Dutch astronomer Willem de-Sitter realised that by removing all matter, he could obtain a static cosmological model that differed from Einstein's. Since there is no matter to move, the solution remains stationary. If some test particles are reintroduced into the model, the cosmological term would push them apart. Astronomers began to wonder if this effect could not explain the spiral galaxies' recession.

#### 1.2.3 Friedmann-Lemaître Model

Friedmann and Georges Lemaître independently discovered realistic solutions to Einstein's equations in 1922 and 1927, respectively. These evolutionary models correspond to cosmologies based on the Big Bang. Friedmann and Lemaître adopted Einstein's spatial homogeneity and isotropy assumption. They rejected his assumption of time independence, however, and considered both positively curved ("closed") and negatively curved ("open") universes. The difference between Friedmann's and Lemaître's approaches is that the former set the cosmological constant to zero, whereas the latter allowed for it to have a nonzero value.

The geometry of space in Friedmann's closed models is similar to Einstein's original model; however, there is a curvature to both time and space. Unlike Einstein's model, in which time runs eternally at each spatial point on an uninterrupted horizontal line that extends infinitely into the past and future, time in Friedmann's version of a closed universe has a beginning and an end when material expands from or is recompressed to infinite densities. These are known as the "big bang" and "big squeeze" instants, respectively.

Final expression of luminosity distance D in closed model is given by,

$$D = \frac{c}{H_0} \frac{1}{q_0^2} [q_0 z + (q_0 - 1)(\sqrt{1 + 2zq_0} - 1)]$$
(1.1)

where  $q_0$  is the deceleration parameter. It is a dimensionless measure of the cosmic expansion of space in Friedmann-Lemaître-Robertson-Walker universe. The formula was first derived by Mattig in 1958.

The spatial and temporal behaviour of open Friedmann models differs from that of closed models. The total volume of space and the number of galaxies contained in an open universe are infinite. The three-dimensional spatial geometry is one of uniform negative curvature, which means that if circles are drawn with very long

lengths of string, the circumference to length ratio is greater than  $2\pi$ . The universe's temporal history begins again with infinite density expansion from a big bang, but this time the expansion continues indefinitely, and the average density of matter and radiation in the universe would eventually become negligibly small. In such a model, time has a beginning but no end.

The calculation in this case is similar to that for the closed model, with the difference that the trigonometric functions are replaced by hyperbolic ones.

Final expression for luminosity distance are same for closed and open model.

From eq.(1.1) for  $q_0 = 1$  we get,

$$D = \frac{c}{H_0} z \tag{1.2}$$

And for  $q_0 = 0$  we get,

$$D = \frac{c}{H_0} z (1 + \frac{z}{2}) \tag{1.3}$$

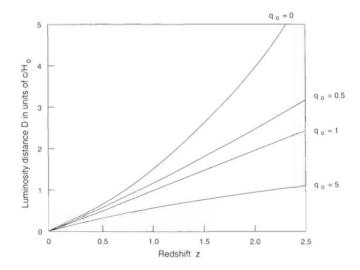


Figure 1.1: The luminosity distance expressed as a function of the redshift for various  $q_0$  values.

Source: Introduction to Cosmology by Jayant V.Narlikar

Plot  $D(q_0, z)$  as a function of z for various parametric value of  $q_0$ . All curves start off with linear Hubble law for small redshift, but then spread out, with only the curve for  $q_0 = 1$  staying linear all the way.[1]

#### 1.2.4 Einstein-de Sitter Model

In 1932, Einstein and de Sitter proposed that the cosmological constant be set to zero, and they developed a homogeneous and isotropic model that provides the separating case between the closed and open Friedmann models; that is, Einstein and de Sitter assumed that the universe's spatial curvature is neither positive nor

negative, but rather zero. The Einstein-de Sitter universe's spatial geometry is Euclidean (infinite total volume), but space-time is not globally flat (i.e., not exactly the space-time of special relativity). Time begins with a big bang again, and the galaxies recede indefinitely, but the recession rate (Hubble's "constant") asymptotically approaches zero as time advances to infinity. The final equation for luminosity distance of Einstein-de Sitter Model is

$$D = \frac{2c}{H_0} [(1+z) - (1+z)^{\frac{1}{2}}]$$
 (1.4)

If we put  $q_0 = 0.5$  in eq.(1.1) it is easy to see that the result eq.(1.4) for the Einstein-de Sitter Model also follows from the same formula.[1]

#### 1.2.5 Standard Model / ACDM Model

The ΛCDM model is a depiction of the big bang cosmological model, which states that the universe has three primary components. A cosmological constant symbolised by Lambda, which is related with dark energy, cold dark matter, and finally ordinary matter. It is now considered to be the standard model of big bang cosmology. This is because it provides reasonably good account for the properties of the cosmos like existence and the structure of the cosmic microwave background, large scale distribution of galaxies and also the accelerated expansion of the universe. The  $\Lambda \text{CDM}$  model incorporates metric space expansion which is extensively established through red shift of spectral absorption or emission lines in light from distant galaxies. The letter  $\Lambda$  represents the cosmological constant that is associated with dark energy in empty space which is used to explain the accelerating expansion of the space against the influence of gravity. Dark matter is postulated in order to account for gravitational effects observed in very large scale structures that cannot be accounted for by the amount of observed matter. Cold dark matter is currently hypothesized to be non-baryonic, cold, dissipation less and collision less. The model uses the Friedmann–Lemaître–Robertson–Walker metric, the Friedmann equations and the cosmological equations of state to describe the observable universe from right after the inflationary epoch to present and future. Currently, the concordance model (currently accepted or a model of the universe that assumes a minimum number of parameters) is the ΛCDM model (which includes cold dark matter and a cosmological constant). According to this model, the Universe is 13.7 billion years old and composed of 4% baryonic matter, 23% dark matter, and 73% dark energy. The Hubble constant for this model is 71 km/s/Mpc.

#### 1.2.6 Eternal Coasting Model

The eternal coasting model[2] is a model that is very similar to the Friedman Robertson-Walker model which assumes the presence of dark energy along with with matter and radiation, the only difference being is that eternal coasting models constrained by the equation of state  $\rho + 3p = 0$ , where  $\rho$  is the energy density. This model predicted that the value  $H_0t_0$  was unity and the data that we are currently obtaining is in agreement with this. The luminosity distance in this model is given by

$$D = \frac{c}{H_0}(1+z)\ln(1+z)$$
 (1.5)

#### 1.3 Motive of the project

The largest dataset of SN 1a samples, that is available to the public is known as the pantheon data[3]. It contains 1048 supernovae on the redshift range 0 < z < 2.3. Pantheon+ data, the largest dataset of SN 1a samples, has not yet been published. Our motivation was to calculate the value of the Hubble constant from these Type Ia supernovae data for various available cosmological models. In addition, we will use this data to try to understand the concept of the Hubble tension. After obtaining the values we try to find the best suitable model to explain the big bang theory with help of the Bayesian theorem. Each model has a prescribed equation for obtaining the value for distance modulus rearranging this equation we obtain the value for calculating the value of  $H_0$ . With values of zCMB from the pantheon data we calculate the Hubble constant. This calculation is done by selecting certain ranges for z for example, z < 0.05, z < 0.15 etc. Now to find the most suitable values in each of these models with the help of Bayesian theorem. This is done by acquiring the value for  $exp(-chi^2/2)$ . In order to do this we have to find the calculated value for zCMB, With which  $\chi$  can be known. We then plot a graph with  $H_0$  and  $exp(-\chi^2/2)$  from which we can obtain the likelihood of the parameter which in this case is the likelihood of  $H_0$ .

## Chapter 2

### **Hubble Constant**

The Hubble Constant is one of the most significant numbers in cosmology because it is needed to estimate the size and age of the universe. This number indicates the rate at which the universe is expanding, since "The Big Bang". In this section we discuss about Hubble's law, Method of calculating Hubble Constant and Hubble tension.

#### 2.0.1 Hubble's Law

In 1929, one of the most astonishing discovery about the motion of galaxies was published by Edwin Hubble. That discovery would demolish many of the previously held beliefs. Using the observations of distant galaxies, he showed that the universe is expanding. Not only that, the light from distant galaxies is systematically increased in wavelength, the fractional increase being proportional to the distance D of the galaxy from us. Thus if  $\lambda$  is the wavelength of light sent out by the galaxy, and  $\lambda + \Delta \lambda = \lambda_0$  the light received, then

$$z = \frac{\Delta \lambda}{\lambda} \propto D \tag{2.1}$$

The quantity z is called the redshift. Hubble interpreted this as a Doppler effect (As an object moves away from us, the light waves emitted by the objects are stretched out, which makes them have lower pitch and moves towards end of the electromagnetic spectrum, where light has a longer wavelength. This is called redshift) and attributed a velocity of recession v = cz to the source galaxy.

Hubble observation (1.1) can be written as

$$cz = H_0 D (2.2)$$

And also be written as

$$v = H_0 D (2.3)$$

Where  $H_0$  is called Hubble constant. The Hubble constant is most frequently quoted

in km/s/Mpc. Though the Hubble constant  $H_0$  is roughly constant in the velocity-distance space at any given moment in time. Hubble found that  $H_0 \approx 530 km/s/Mpc$ 

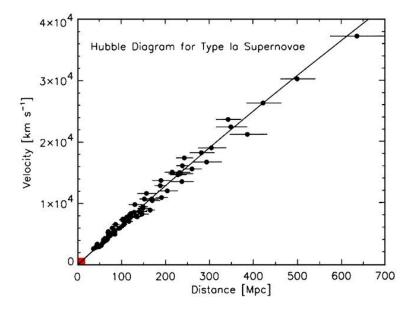


Figure 2.1: The Hubble diagram for type Ia supernovae. Source: PNAS

The Hubble constant  $H_0$  has units of inverse time; the Hubble time  $t_H$  is simply defined as the inverse of Hubble constant.

i.e., 
$$t_H = \frac{1}{H_0} = \frac{1}{530 km/s/Mpc} = 1.8 \times 10^9 \text{ years}$$

In these units  $H_0^{-1} \approx 1.8 \times 10^9$  years as found by Hubble.

The Hubble time is the age it would have had if the expansion had been linear, and it is different from the real age of the universe because the expansion is not linear.

Although Hubble originally obtained  $H_0 \approx 530 km/s/Mpc$ , the present estimate of  $H_0$  is much lower. We believed to lie in the range of  $50 \le H_0 \le 100$ . Hubble Space Telescope and some ground based telescopes have narrowed this range down to around 55 to 75.

#### 2.0.2 Hubble Constant

The Hubble constant describes how fast the universe is expanding in different distances from a particular point in space. It was first calculated in the 1920s by Edwin Hubble, he made the observation that cloud-like celestial objects were actually distant galaxies that were present outside our own galaxy. Using the research done by Henrietta Leavitt on Cepheid, Hubble derived the Cepheid's distance which lead to the formulation of the Hubble's law that was explained in the previous section.

Hubble constant acts as a constant of proportionality in the Hubble law. It is denoted as  $H_0$ . Hubble's original value for  $H_0$  was 500 km per Mpc in cosmological units. But modern techniques have helped in refining the initial measurement. How much was Hubble's value off by, still is a topic of debate. After the discovery that the universe is not only expanding but also accelerating in its expansion it became necessary to modify the existing model with new data, including the addition of "dark energy"- a force that pushes everything apart in the universe. After this discovery they tried to pin down the value of the Hubble constant using two methods.

- 1. Using the data of Cepheid variables and other astrophysical sources, which gave the value of  $H_0$  to be 73 km/s/Mpc
- 2. Using the information of European Space Agency's plank satellite where  $H_0$  was found to be 67 km/s/Mpc.

The two values obtained even though are different they are extremely precise without any overlapping between their error bars. If any one of these methods are considered to be wrong it can lead to a domino effect which can affect all the advancements that we have made in the field of cosmology. This continuous to be the biggest topic of debate in cosmology. New methods such as using the LIGO to obtain new set of data or gravitational lensing- which occurs when extremely massive object wraps and bends the space time like magnifying glass- could clear the discrepancy. one measurement is about how fast the universe is expanding as of today, while the other measurement is a prediction that is based on the physics of the early universe, this discrepancy shows us that we are missing a factor that connects these two eras in our cosmological model. This difference also affects the calculation of age of the universe, cutting of more than a billion years of existence in one case.

The introduction of the concept of dark energy can help us to validate the discrepancy in the value of  $H_0$ . One idea contains the use of a sub atomic particle that travels close to the velocity of light, these entities are called as dark radiations, these particles can affect the speed of expansion of the universe. Another idea is the presence of a special intense dark- energy produced after the big bang which expanded the universe faster than astronomers had previously anticipated. But at

the moment no one knows when and where the final answer for the Hubble constant lies.

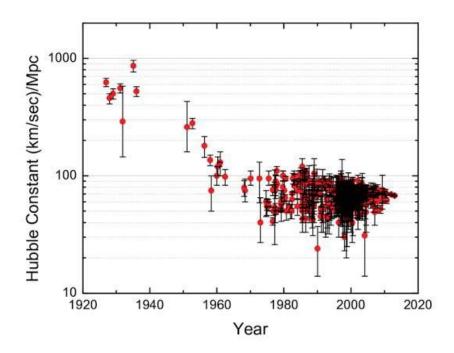


Figure 2.2: Historical evolution of Hubble constant measurements. Source: uploaded by Boris Pritychenko (https://www.researchgate.net/profile/Boris-Pritychenko)

#### 2.0.3 Method of calculating Hubble Constant

Considerable progress has been made in determining the Hubble constant over the past two decades. Even though the accuracy of these methods is a topic up for debate. In this section we briefly discuss the different methods present, the major focus being high precision distance determination methods like Cepheid, tip of the red giant, maser galaxies, surface brightness fluctuations, the Tully-Fisher relation and Type 1a supernovae which is the method that is being used in this dissertation to determine the Hubble constant.

A Cepheid is a variable star having a regular cycle of brightness with a frequency related to its luminosity, so allowing estimation of its distance from the earth. This method provided the initial step on the cosmic distance ladder. Hubble used the observational data for Cepheid in galaxies to deduce a law, which stated that the more remote the galaxy is from the Milky Way, the faster its recession is. The period luminosity(pl) relation for Cepheid is a relation linking the luminosity of pulsating variable stars with their pulsation period. This pl relation provided a powerful tool that astronomers use to measure the distance to nearby galaxies and hence calculate

the Hubble's constant.

Tip of the Red Giant Branch (TRGB) is a distance indicator used in astronomy. it uses the luminosity of the brightest red giant branch stars in a galaxy as reference point to measure the distance to that galaxy. TRGB offers an alternative to Cepheid as it provides an accurate ( $\sim 5\%$ ) distance to galaxies within 10 Mpc with a single orbit of Hubble space telescope time.

The next method is by using maser galaxies. An astrophysical maser is a naturally occurring source of stimulated spectral line emission, typically in the microwave portion of the electromagnetic spectrum. In this technique we utilize the mapping of 22.2 GHz water maser sources in the accretion disks of massive black holes located in spiral galaxies with active galactic nuclei. In the simplest version of this technique a rotation curve is measured along the major axis of the disk; proper motions are measured on the near side of the disk minor axis, and a comparison of the angular velocities in the latter measurement with the absolute velocities in km s-1 in the former measurements yields the distance.

For distance to elliptical galaxies and early-type spiral bulge populations we use the Surface Brightness Fluctuation method(SBF). Both TRGB and SFB use the properties of the red giant branch luminosity function to estimate distances. here pixel to pixel varience is measured, the variance in a pixel is taken to be a function of distance simply because the total number of discrete sources contributing to any given pixel increases with the square of the distance. The major difference between TRGB and SBF method is that TRGB completely relies on the brightest red giant stars while SBF method uses a luminosity weighed integral.

The Tully-Fisher relation at present is one of the most widely applied methods for distance measurements, providing distances to thousands of galaxies both in the general field and in groups and clusters. The relation can be understood in terms of of the virial relation applied to rotationally supported disk galaxies, under the assumption of a constant mass to light ratio. However, a detailed self-consistent physical picture that reproduces the Tully-Fisher relation and the role of dark matter in producing almost universal spiral galaxy rotation curves still remain a challenge.

Type Ia supernovae, unlike Cepheids, are bright enough to be seen from relatively greater distances. Astronomers compare the luminosity and apparent brightness of distant supernovae to determine the distance to which the universe's expansion can be seen. They compare those distance measurements to how light from supernovae is stretched to longer wavelengths due to space expansion. They use these two values to calculate the Hubble constant.

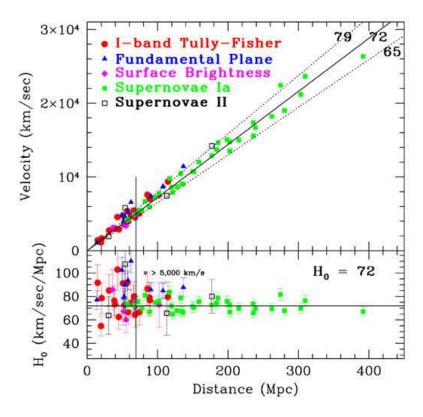


Figure 2.3: Hubble diagram using different methods. Source : Measuring and Understanding the Universe - W. Freedman

## Chapter 3

## Calculating Hubble Constant using SN Ia data

Using the data obtained from the Pan-STARRSI (PS1) Medium Deep Survey, scolnic et.el.2018 [3] deduced the optical light curves, redshifts and classification for 365 spectroscopically confirmed Type Ia supernovae. Improvements were done on PS1 SN photometry, astrometry, and calibrations which help in reducing the systematic uncertainties in the PS1 SN 1a distances. By combining the subset of 279 PS1 SNe Ia (0.03 < z < 0.68) with useful distance estimates of SNe 1a from the sloan Digital Sky ,SNLS and Hubble Space Telescope samples thus forming the largest combined sample of SNe Ia, with a total of 1048 SNe Ia within the range of 0.01 < z < 2.3 hence forming the what we call as the "Pantheon Sample" (fig.3.1)

In this section value of  $H_0$  is calculated by varying the deceleration parameter  $(q_0 = 0, q_0 = \frac{1}{2}, q_0 = 1)$ . Then we exclusively select low redshift region and calculate  $H_0$ . and finally  $H_0$  for the entire dataset (1048 supernovae).

The sample set contains Target ID, zCMB, zHEL (redshifts in two different reference frames, DZ(Error in redshifts), MB (Apparent Magnitude), DMB (Error in MB). Let's discuss these terms in detail.

#### 3.0.1 Redshift

Redshift is an extremely important phenomenon in cosmology and astronomy. When electromagnetic radiation are emitted or reflected from an object, it is shifted toward the less energetic (higher wavelength) end of the spectrum. Using this phenomenon allows for the first steps toward understanding features of our galaxy and even the universe as a whole. Our data set contains two types of redshifts: Zhel and ZCMB. ZHEL is the redshift in the Sun's reference frame. ZCMB is the redshift in CMB(Cosmic Microwave Background) frame.

#### 3.0.2 Apparent Magnitude

Apparent magnitude can be defined as the measure of the brightness of a star or any other astronomical entity observed from Earth. An object's apparent magnitude depends upon various factors like intrinsic luminosity, its distance from the earth, and scattering of the radiation due to the presence of interstellar dust along the line of sight of the observer. The distinction between luminosity and apparent brightness is that luminosity is an intrinsic property of the star, with a fixed value, while apparent magnitude varies with the distance of the observer. In this analysis, we denote apparent brightness with 'm'. We obtained the apparent brightness of about 1048 supernovas from Scolnic et al. (2018), using which we calculated the values of the Hubble constant. In the data sheet it is denoted as MB.

#### 3.0.3 Absolute Magnitude

Absolute magnitude(M) is the measure of the luminosity of a celestial object on an inverse logarithmic astronomical magnitude scale. If an object is viewed from a distance of exactly 10 parsecs, its absolute magnitude equals its apparent magnitude, without extinction of its light due to the presence of interstellar dust. Absolute magnitude can be specified for different wavelength regions. As the luminosity of an object increases, the numeric value of its apparent brightness decreases. During this analysis we have taken the value of M to be -19.3. This assumption is based on the fact that we only consider type 1a supernovae. The value of M changes with the type entity we measure. Type 1a supernovae occur when a white dwarf accumulates too much mass to resist the force of gravity. This always occurs when the mass of the star reaches the Chandrasekhar limit. All type Ia supernovae have approximately the same absolute magnitude, for this very reason they are commonly used as standard candles to determine the distance to a galaxy once the stretch factor is accounted for.

| Target ID<br>(sortable) | ZCMB<br>(sortable)                      | ZHEL<br>(sortable) | DZ<br>(sortable) | MB (µ - M <sub>B</sub> )<br>see Eqn. 3 in <u>Scolnic</u><br>et al. 2018<br>(sortable)   | DMB (error in MB)<br>(sortable)         |
|-------------------------|---|--------------------|------------------|---|---|
| Type filter             | Type filter                             | Type filter        | Type filter      | Type filter   |   |
| 03D1au                  | 0.50309                                 | 0.50309            | 0.0              | 22.93445  | 0.12605                                 |
| 03D1aw                  | 0.58073                                 | 0.58073            | 0.0              | 23.52355  | 0.1372                                  |
| 3D1ax                   | 0.4948                                  | 0.4948             | 0.0              | 22.8802   | 0.11765                                 |
| 3D1bp                   | 0.34593                                 | 0.34593            | 0.0              | 22.11525  | 0.111                                   |
| 3D1co                   | 0.67767                                 | 0.67767            | 0.0              | 24.0377   | 0.2056                                  |
| 3D1ew                   | 0.8665                                  | 0.8665             | 0.0              | 24.34685  | 0.17385                                 |
| 3D1fc                   | 0.33094                                 | 0.33094            | 0.0              | 21.7829   | 0.10685                                 |
| 3D1fq                   | 0.79857                                 | 0.79857            | 0.0              | 24.3605   | 0.17435                                 |
| 3D3aw                   | 0.44956                                 | 0.44956            | 0.0              | 22.78895  | 0.14135                                 |
| 3D3ay                   | 0.37144                                 | 0.37144            | 0.0              | 22.28785  | 0.1245                                  |
| 3D3ba                   | 0.29172                                 | 0.29172            | 0.0              | 21.47215  | 0.12535                                 |
| 3D3bl                   | 0.35582                                 | 0.35582            | 0.0              | 22.05915  | 0.12645                                 |
| 3D3cd                   | 0.46127                                 | 0.46127            | 0.0              | 22.62945  | 0.13775                                 |
| 3D4ag                   | 0.2836                                  | 0.2836             | 0.0              | 21.40915  | 0.1028                                  |
| 3D4at                   | 0.63222                                 | 0.63222            | 0.0              | 23.66065  | 0.20445                                 |
| 200 200                 | (2-2-2-2-2-2-2-2-2-2-2-2-2-2-2-2-2-2-2- | 0-0110-1000        | 0.0              |   | 100000000000000000000000000000000000000 |
| 3D4au                   | 0.4664                                  | 0.4664             | 0.0              | 23.21635  | 0.1587                                  |
| 3D4cj                   | 0.94688                                 |                    | 124              | someoffice of   | L-1-20001-102                           |
| 3D4cx                   |   | 0.94688            | 0.0              | 24.41155  | 0.1642                                  |
| 3D4cy                   | 0.92491                                 | 0.92491            | 0.0              | 24.7416   | 0.1923                                  |
| 3D4cz                   | 0.69315                                 | 0.69315            | 0.0              | 23.79645  | 0.24445                                 |
| 3D4dh                   | 0.62522                                 | 0.62522            | 0.0              | 23.46015  | 0.1303                                  |
| 3D4di                   | 0.89693                                 | 0.89693            | 0.0              | 24.40275  | 0.16645                                 |
| 3D4dy                   | 0.60825                                 | 0.60825            | 0.0              | 23.3767   | 0.1266                                  |
| 3D4fd                   | 0.78904                                 | 0.78904            | 0.0              | 24.2188   | 0.11585                                 |
| 3D4gf                   | 0.57828                                 | 0.57828            | 0.0              | 23.2373   | 0.13895                                 |
| 3D4gg                   | 0.59027                                 | 0.59027            | 0.0              | 23.3354   | 0.1666                                  |
| 4D1aj                   | 0.71963                                 | 0.71963            | 0.0              | 23.82525  | 0.2057                                  |
| 4D1dc                   | 0.21075                                 | 0.21075            | 0.0              | 20.64495  | 0.10525                                 |
| 4D1de                   | 0.7666                                  | 0.7666             | 0.0              | 24.24705  | 0.1593                                  |
| 94D1ff                  | 0.85852                                 | 0.85852            | 0.0              | 24.1237   | 0.1478                                  |
| 4D1hd                   | 0.36791                                 | 0.36791            | 0.0              | 22.24455  | 0.1019                                  |
| 4D1hx                   | 0.55875                                 | 0.55875            | 0.0              | 23.1648   | 0.1335                                  |
| 4D1hy                   | 0.84852                                 | 0.84852            | 0.0              | 24.33175  | 0.14855                                 |
| 4D1iv                   | 0.9964                                  | 0.9964             | 0.0              | 24.69965  | 0.1408                                  |
| 4D1jd                   | 0.7766                                  | 0.7766             | 0.0              | 24.0942   | 0.17835                                 |
| 4D1jg                   | 0.58294                                 | 0.58294            | 0.0              | 23.34385  | 0.1258                                  |
| 4D1kj                   | 0.58375                                 | 0.58375            | 0.0              | 23.2686   | 0.09915                                 |
| 04D1ks                  | 0.79656                                 | 0.79656            | 0.0              | 23.85445  | 0.1588                                  |
| 4D1oh                   | 0.58873                                 | 0.58873            | 0.0              | 23.30085  | 0.13005                                 |
| 4D1ow                   | 0.91348                                 | 0.91348            | 0.0              | 24.37895  | 0.13735                                 |
| 4D1pc                   | 0.76859                                 | 0.76859            | 0.0              | 24.2078   | 0.16885                                 |
| 4D1pd                   | 0.94846                                 | 0.94846            | 0.0              | 24,66675  | 0.15105                                 |
| 4D1pg                   | 0.5138                                  | 0.5138             | 0.0              | 23.12745  | 0.1329                                  |
| 4D1pp                   | 0.73362                                 | 0.73362            | 0.0              | 23.8481   | 0.15575                                 |
| 4D1pu                   | 0.63771                                 | 0.63771            | 0.0              | 23.37   | 0.2329                                  |
| 4D1qd                   | 0.7656                                  | 0.7656             | 0.0              | 24.10805  | 0.1549                                  |
| 4D1rh                   | 0.43487                                 | 0.43487            | 0.0              | 22.5623   | 0.1145                                  |
| 4D1rx                   | 0.98341                                 | 0.98341            | 0.0              | 24.828  | 0.153                                   |
| 4D1sa                   | 0.58375                                 | 0.58375            | 0.0              | 23.46355  | 0.14025                                 |
| 4D1si                   | 0.70064                                 | 0.70064            | 0.0              | 23.7492   | 0.1738                                  |
| 2007H531                |   | CONT. CTC. L.C.    | 52552            | CONTRACTOR |   |

Figure 3.1: Scolnic et el. 2018 Supernovae Catalog

#### 3.0.4 Distance Modulus

The distance modulus can be easily defined as the difference between the apparent magnitude(m) and the absolute magnitude(M). The equation that we use in this analysis to obtain distance modulus is given as

$$\mu = m - M = 5log(D) - 5 \tag{3.1}$$

where D is the distance in pc.

Another equation that we can use is the modified form of the Tripp equation which is given as

$$\mu = m - M + \alpha x \mathbf{1} - \beta c + \Delta m + \Delta n \tag{3.2}$$

where  $\alpha$  is the coefficient of relation between luminosity and stretch,  $\beta$  is the coefficient of relation between luminosity and colour,  $\Delta m$  is the mass correction based on the host galaxy of the SN and  $\Delta n$  is a distance correction based on predicted biases from stimulation.

#### 3.1 Calculation of $H_0$

#### 3.1.1 $q_0 = 1$ (Hubble's Law)

From Hubble's law eq.(2.2) we get,

$$H_0 = zc/D (3.3)$$

The value of D can be obtained from (3.1),

$$\mu = m - M = 5log(D) - 5 \tag{3.1}$$

Rearranging,

$$D = \frac{1}{\log \frac{\mu + 5}{5}} \tag{3.4}$$

Since low redshift will give more accurate value of  $H_0$  due to reduced extinction of radiation we initially choose the range for z to be less than 0.05. There are 157 supernovae between this range.

The average value of  $H_0$  with in this range is 70.7156 km/s/Mpc similarly calculating the average value of  $H_0$  by changing the range from z less than 0.05(157 supernovae), less than 0.1(211 supernovae), less than 0.15(299 supernovae) we obtained 70.5494, 68.7576 respectively.

Now we calculate  $H_0$  for complete dataset (high redshift region) thereby obtaining  $H_0$  to be 60.3170 km/s/Mpc.

| A             | B  ZCMB (sortable) WMB (u - MB) | C<br>see Eqn. 3 in Scolnic et al. 2018 (sortable) ▼DMB (erro | D<br>or in MB) (sortable) ▼m-M | E        | H K D in Mpc  | ZCMB*c H0 |
|---------------|---------------------------------|--|--------------------------------|----------|---------------|-----------|
| Full Screen T | 0.08859                         | 18.64345   | 0.10605                        | 37,94345 | 387.873402584 | 26577 68  |
| _             | 0.094                           | 18.889   | 0.10285                        | 38.189   | 434.310170886 | 28200 64  |
| Full Screen   | 0.06641                         | 17.9881  | 0.10283                        | 37.2881  | 286.826979833 | 19923 69  |
| 1 6558        | 0.05715                         | 17.9881  | 0.1134                         | 36.95695 | 246.257802674 | 17145 69  |
| 2 12779       | 0.0713                          | 18.5397  | 0.1072                         | 37.8397  | 369.777089658 | 23649 63  |
| 3 12781       | 0.08328                         | 18.7438  | 0.11605                        | 38.0438  | 406.218782331 | 24984 6   |
| 7 12898       | 0.08286                         | 18.5179  | 0.11005                        | 37.8179  | 366.083369396 | 24858 67  |
| 12050         | 0.08162                         | 18.5175  | 0.1035                         | 37.811   | 364.921961009 | 24486 67  |
| 0 17240       | 0.0718                          | 18.0121  | 0.1254                         | 37.3121  | 290.014692268 | 21540 7   |
| 6 17784       | 0.03743                         | 16.7708  | 0.1166                         | 36.0708  | 163.741965801 | 11229 68  |
| 4 18241       | 0.0939                          | 18.94955   | 0.1100                         | 38.24955 | 446.591034314 | 28170 63  |
| 5 21502       | 0.08798                         | 18.6796  | 0.14945                        | 37.9796  |               | 26394 66  |
|               | 0.08798                         | 18.62285   | 0.1057                         | 37.92285 | 394.384647047 | 25602 66  |
| 722           |                                 |  |                                |          | 384.211181714 | 27834 6   |
| 3 774         | 0.09278                         | 18.8576  | 0.109                          | 38.1576  | 428.075132721 |           |
| 3 3592        | 0.08575                         | 18.608   | 0.10035                        | 37.908   | 381.592649119 | 25725 67  |
| 7 7876        | 0.07523                         | 18.2988  | 0.10185                        | 37.5988  | 330.948182215 | 22569 68  |
| 1 10028       | 0.06481                         | 17.87655   | 0.1109                         | 37.17655 | 272.464546982 | 19443 71  |
| 3 10805       | 0.04437                         | 16.92415   | 0.11995                        | 36.22415 | 175.723562794 | 13311 75  |
| 17186         | 0.07877                         | 18.7576  | 0.11205                        | 38.0576  | 408.808578008 | 23631 57  |
| 17258         | 0.08833                         | 18.4809  | 0.11615                        | 37.7809  | 359.898469681 | 26499 7   |
| 19968         | 0.05573                         | 17.56165   | 0.10915                        | 36.86165 | 235.68394554  | 16719 70  |
| 2001ah        | 0.05948                         | 17.71775   | 0.12625                        | 37.01775 | 253.25031837  | 17844 70  |
| 4 2001az      | 0.04093                         | 17.0528  | 0.11665                        | 36.3528  | 186.448974924 | 12279 65  |
| 5 2001da      | 0.01705                         | 14.85645   | 0.1832                         | 34.15645 | 67.8094153619 | 5115      |
| 5 2001en      | 0.01531                         | 14.59755   | 0.1613                         | 33.89755 | 60.1880121542 | 4593 76   |
| 2001fe        | 0.01472                         | 14.6158  | 0.14905                        | 33.9158  | 60.6959900694 | 4416 72   |
| 2001gb        | 0.02673                         | 15.97105   | 0.16165                        | 35.27105 | 113.294805976 | 8019 70   |
| 9 2001G       | 0.01732                         | 14.8718  | 0.1587                         | 34.1718  | 68.2904540079 | 5196 76   |
| 2001ic        | 0.04437                         | 16.72085   | 0.19805                        | 36.02085 | 160.018428132 | 13311 83  |
| 2001V         | 0.01567                         | 14.4307  | 0.15405                        | 33.7307  | 55.736539335  | 4701 84   |
| 2002bf        | 0.02477                         | 15.6593  | 0.1718                         | 34.9593  | 98.143151588  | 7431 75   |
| 2002bz        | 0.038                           | 16.7664  | 0.13185                        | 36.0664  | 163.410515487 | 11400 69  |
| 1 2002ck      | 0.0299                          | 16.2413  | 0.123                          | 35.5413  | 128.309850807 | 8970 6    |
| 2002cr        | 0.01012                         | 13.90745   | 0.19825                        | 33.20745 | 43.8016025363 | 3036 69   |
| 2002de        | 0.0258                          | 16.1648  | 0.1376                         | 35.4648  | 123.86824944  | 7740 62   |
| 7 2002dp      | 0.01038                         | 14.0496  | 0.2048                         | 33.3496  | 46.7648989225 | 3114 6    |
| 3 2002eu      | 0.03791                         | 16.5851  | 0.2001                         | 35.8851  | 150.321118993 | 11373 75  |
| 9 2002G       | 0.03507                         | 16.45115   | 0.16785                        | 35.75115 | 141.328581496 | 10521 74  |
| 2002ha        | 0.01268                         | 14.4486  | 0.171                          | 33.7486  | 56.1978887188 | 3804 67   |
| 2002he        | 0.02525                         | 15.82815   | 0.13295                        | 35.12815 | 106.079142398 | 7575 71   |
| 2 2002kf      | 0.01908                         | 15.3534  | 0.14115                        | 34.6534  | 85.2471758147 | 5724 67   |
| 3 2003ae      | 0.03489                         | 16.36835   | 0.1667                         | 35.66835 | 136.041057746 | 10467 76  |
| 4 2003ch      | 0.0295                          | 16.3438  | 0.12735                        | 35.6438  | 134.511680961 | 8850 65   |
| 5 2003cq      | 0.03416                         | 16.57565   | 0.1563                         | 35.87565 | 149.668360042 | 10248 68  |
| 6 2003cq      | 0.03410                         | 16.7928  | 0.1303                         | 36.0928  | 165 409329042 | 12201 73  |

Figure 3.2: Calculating  $H_0$  from scolnic et.al 2018 supernovae catalog in the range of z<0.05

#### 3.1.2 $q_0 = 0.5$ (Einstein-de Sitter Model)

From Einstein de-Sitter model we have the equation for luminosity distance eq.(1.4) Rearranging,

$$H_0 = \frac{2c}{D}[(1+z) - (1+z)^{\frac{1}{2}}]$$
(3.5)

where D can be obtained from eq.(3.4)

Initially taking the value of z to be less than 0.05, the value of  $H_0$  can be obtained as 71.8407. Next we calculate  $H_0$  similar to the previous section, by considering the entire dataset, we obtain  $H_0$  as 63.8739 km/s/Mpc.

#### **3.1.3** $q_0 = 0$

Rearranging eq.(1.3) we get,

$$H_0 = \frac{c}{D}z(1+\frac{z}{2}) \tag{3.6}$$

For range z < 0.05, we can obtain  $H_0$  as 72.3320. similarly for complete dataset the value of  $H_0$  is found to be 69.067 km/s/Mpc.

#### 3.1.4 Eternal coasting model

By rearranging the equation of luminosity distance of eternal coasting model eq.(1.5) we get,

$$H_0 = \frac{c}{D}(1+z)\ln(1+z)$$
 (3.7)

For low redshift range (z < 0.05),  $H_0$  is 72.3222 km/s/MPc . For complete dataset the value of  $H_0$  is 67.9062 km/s/Mpc.

## Chapter 4

## **Bayesian Parameter Estimation**

## 4.1 $\chi^2$ (Chi square) Test

A chi-square  $(\chi^2)$  statistic is a test that compares a model to actual observed data. Given the size of the sample and the number of variables in the relationship, the chi-square statistic compares the size of any discrepancies between the expected and actual results.  $\chi^2$  can be used to determine whether two variables are related or independent.

$$\chi^2 = \left(\frac{Calculated - Observed}{Error}\right)^2 \tag{4.1}$$

The calculated and observed values of apparent magnitude(m) has been taken during the evaluation of  $\chi^2$  value. Eq.(4.1) becomes,

$$\chi^2 = \left(\frac{m_c - m_o}{DMB}\right)^2 \tag{4.2}$$

 $m_c$  can be deduced from the formula,

$$m_c = (5loqD - 5) + M$$
 (4.3)

where absolute magnitude(M) is taken to be -19.3. Also for the evaluation of  $\chi^2$  here we required  $D_{calculated}$  which in turn is obtained by using an expected value of  $H_0$  within the equations for D which varies depending upon the models used (Ref.eqns.(1.2),(1.3),(1.4)).

Using the above equations value of  $\chi^2$  is procured. [4]

#### 4.2 Bayesian theory

The Bayesian theory can be used to assess the probability of a hypothesis. The Bayes theorem can be derived with the help of two fundamental axioms. Let A, B, C represent prepositions and  $\bar{A}$  represents the denial of A. AB is used to denote some

prepositions that is true only if 'A and B' are true. Similarly the symbol A + B denotes a preposition that is true if 'either A or B; is true.

Axiom 1(sum rule) - 
$$P(A \mid C) + P(\bar{A} \mid C) = 1$$

Axiom 2(product rule) - 
$$P(AB \mid C) = P(A \mid BC)P(B \mid C)$$

Here, the vertical bar '|' is the conditional symbol, indicating what propositions are assumed for the assignment of the probability. As a result, while estimating probabilities, we will aim to explicitly include any relevant background information.

We know that AB and BA are identical from the definition of joint occurrence of A and B, and simplifying we get

$$P(A \mid BC) = P(A \mid C) \frac{P(B \mid AC)}{P(B \mid C)}$$

This is known as the Bayes theorem. Given a data D with some background information I, we can calculate the probability of the hypothesis H by substituting A = H, B = D and C = I, hence modifying the theorem as

$$(H \mid DI) = P(H \mid I)P(D \mid HI)P(D \mid I)$$

the RHS of the above equation is called the posterior probability for H.

The first term  $P(H \mid I)$  is what that gives the probability of the hypothesis. This is assigned before the analysis of the hypothesis hence the name prior probability. It needs to be acknowledged that there is some unpredictability in giving this value, as prior knowledge may vary or even be lacking.

The second term in the numerator  $P(D \mid HI)$ , is the probability for the data with the assumption that both the hypothesis and the information are true.

The quantity on the denominator  $P(D \mid I)$ , is the probability to get the data irrespective of the hypothesis.

The Bayesian theorem helps in calculating the probability of hypothesis, this method differs from the frequentest approach where probability is measured for random variable. However, we should observe that the axioms that lead to Bayes' theorem are relatively arbitrary. Another flaw in utilising Bayes' theory is the issue of allocating prior probability. An important feature of the Bayes theorem is that we can adjust the assessment of a hypothesis when new data regarding the hypothesis is acquired. The Bayesian approach is a more general one as a variety of hypothesis are assessed with their posterior probability. This evaluation requires the prior possibility assignment of the hypothesis and the the probability of the data.

#### 4.2.1 Bayesian parameter estimation

The theory or parameterized model is made up of a set of exclusive hypotheses that are labelled by certain parameters. For some value of each of these parameters, the model is considered to be true. Because the truth of this model is known, it is included in information I in the Bayes' theorem. Continuous or discrete parameters might be used. In this scenario, the hypothesis we wish to evaluate is the set of values for those parameters in the theory. The probability evaluation is done using the theorem that is given as

$$P(\theta \mid DI) = P(\theta \mid I) \frac{P(D \mid \theta I)}{P(D \mid I)}$$

Where the terms in the numerator are the prior probability and the likelihood function. The probability  $P(D \mid \theta I)$  is the probability for the data, given the truth of the value of the parameter  $\theta$ . It is called the the likelihood for the value of the parameter  $\theta$  and can be calculated as  $exp(-\chi^2/2)$ .

Likelihood of parameter = 
$$exp(-\chi^2/2)$$
 (4.4)

The first term on the RHS of the above equation  $P(\theta \mid I)$  is the probability for the value of the parameter  $\theta$ . Before analysing the data. Hence it is called as the prior probability for  $\theta$ . We note that the term in the denominator,  $P(D \mid I)$  is independent of  $\theta$  and that it serves only as a normalisation constant.[5]

#### 4.3 Estimation of Likelihood

Here the likelihood is calculated using the equation (4.4), with the values obtained for  $exp(-\chi^2/2)$  we plot a graph by taking the Hubble constant value between the range 40 to 100 km/s/Mpc.

#### 4.3.1 $q_0 = 1$ (Linear Hubble Law)

Initially we set the limit for z to be less than 0.05. With the procedures prescribed in the previous section we deduce the values for  $\chi^2$  and  $exp(-\chi^2/2)$ . With the obtained values for  $exp(-\chi^2/2)$  we plot a graph with the previously defined range for  $H_0$ .

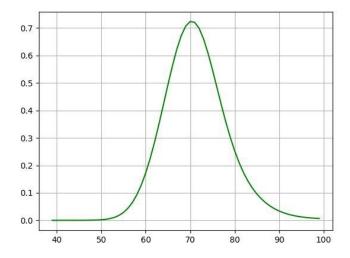


Figure 4.1: Plot of  $H_0$  versus  $\exp(-\chi^2/2)$  in the range z <0.05

From the graph it is visible that the curve we have obtained is nearly a gaussian distribution, whose peak gives us the likelihood. Therefore, from the above curve likelihood is 71.

Next we select the range to be z < 0.1

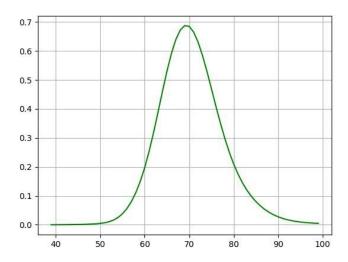


Figure 4.2: Plot of  $H_0$  versus  $exp(-\chi^2/2)$  in the range z <0.1

For this range likelihood is 70 km/s/Mpc.

Similarly we change the range to z < 0.15 and plotting,

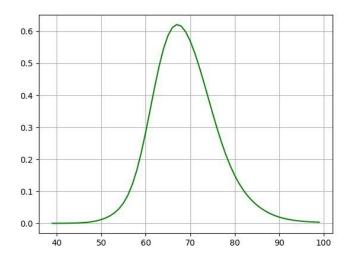


Figure 4.3: Plot of  $H_0$  versus  $exp(-\chi^2/2)$  in the range z<0.15

And for z < 0.15, likelihood is 68 km/s/Mpc.

Finally we take the entire dataset

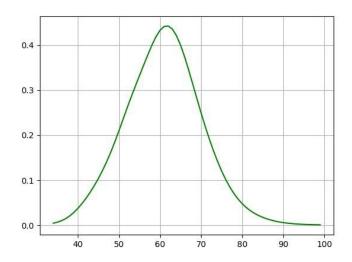


Figure 4.4: Plot of  $H_0$  versus  $exp(-\chi^2/2)$  for the complete dataset  $(q_0 = 1)$ 

Here likelihood is 63 km/s/Mpc.

**4.3.2** 
$$\mathbf{q}_0 = \frac{1}{2}(Einstein - deSitterModel)$$

Here we calculate the value of  $exp(-\chi^2/2)$  by taking the entire dataset (1048 supernovae)

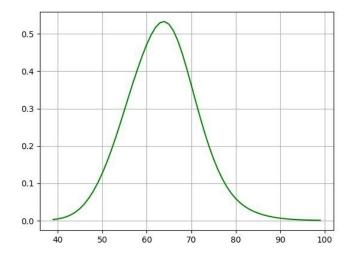


Figure 4.5: Plot of  $H_0$  versus  $exp(-\chi^2/2)$  for the complete dataset  $(q_0 = \frac{1}{2})$ 

From the above graph the peak is at 65 km/s/ Mpc which is the likelihood.

#### **4.3.3** $q_0 = 0$

Similar to the previous subsection here also we take the entire dataset to attain the values of  $\exp(-\chi^2/2)$ 

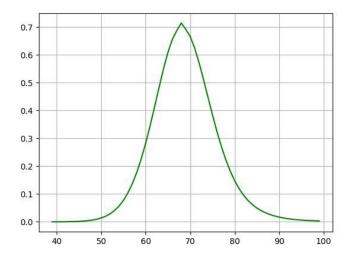


Figure 4.6: Plot of  $H_0$  versus  $exp(-\chi^2/2)$  for the complete dataset  $(q_0=0)$ 

Similar to the previous graphs here the likelihood can be obtained from the peak which is 69 km/s/Mpc.

#### 4.4 Error Estimation

If we observe all the graphs that we have obtained, it can be seen that they all are almost a gaussian distribution. The standard deviation of a gaussian distribution gives the error in our estimation. In this dissertation we calculate the sigma value directly the graph. if we observe the curve it can be seen that The distribution has equal amount of data on either side of the peak, the standard deviation quantifies the variability of the curve. The assumption that we can make about a data that follows a gaussian distribution is that the area under the curve is relative to how many standard deviation we are away from the mean. The area between plus and minus one standard deviation from the mean contains 68% of the data .Two standard deviation contains 95% of data and three standard deviation contains 99.8% of data.



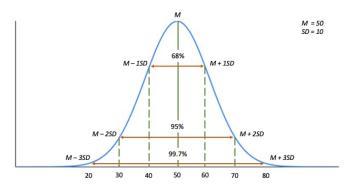


Figure 4.7: Standard deviations in a gaussian distribution. Source: Scribbr.

Using the method provided in the above graph we can acquired the standard deviation( $\sigma$ ) of all the  $H_0$  versus  $\exp(-\chi^2/2)$  graphs.

- For  $q_0 = 1$  within the ranges z < 0.05, z < 0.1, z < 0.15 and for the entire dataset the values of  $\sigma$  are 7.83,7.67,7.67 and 7.33 respectively.
- Similarly for  $q_0 = 0$  the value of sigma for entire dataset is 7.83
- Likewise for  $q_0 = 0.5$  for the entire dataset the value of sigma is 8.17

## Chapter 5

## Conclusion

Throughout this dissertation we have tried to acquire the value of  $H_0$  using the supernovae data from scolnic et.el. 2018

Initially we calculate the value of  $H_0$  in four ways using four different models. We had obtained the value of  $H_0$  for  $q_0 = 1$ ,  $q_0 = 0.5$ ,  $q_0 = 0$  and eternal coasting model in the low redshift region (z < 0.05) as 70.7156 km/s/Mpc, 71.8407 km/s/Mpc, 72.3320 km/s/Mpc and 72.3222 km/s/Mpc respectively. If we analyze these  $H_0$  values they can found to be compatible. Hence the low redshift region can be considered to model independent, as visible in fig.(1.1).

As we increase the range of z, the values of  $H_0$  become less compatible with each other. Due to which they are considered to be model dependent. This notion can be better understood if we look into the values of  $H_0$  for  $q_0 = 1$ ,  $q_0 = 0.5$ ,  $q_0 = 0$  and eternal coasting model when the entire dataset is taken into consideration. The values are  $H_0 = 60.3170 km/s/Mpc$ ,  $H_0 = 63.8739 km/s/Mpc$ ,  $H_0 = 69.067 km/s/Mpc$  and  $H_0 = 67.9062 \text{ km/s/Mpc}$  respectively.

We had used Bayesian method to find the most suitable value of  $H_0$ . This suitable value was acquired using the peak of  $H_0$  versus  $exp(-\chi^2/2)$  curve. This peak is called as the likelihood of parameter  $(H_0)$ . The curve that we had plotted using  $H_0$  and  $exp(-\chi^2/2)$  is a normal curve (gaussian distribution). Any error that might occur can be calculated using the standard deviation, since it is a gaussian distribution.

Combining all the informations that we had stated previously the most suitable  $H_0$  value along with it error for different models was found to be,

for 
$$q_0 = 1$$
,  $H_0 = 63 \pm 7.33 km/s/Mpc$   
for  $q_0 = 0.5$ ,  $H_0 = 65 \pm 8.17 km/s/Mpc$   
for  $q_0 = 0$ ,  $H_0 = 69 \pm 7.83 km/s/Mpc$ 

Likewise the value of  $H_0$  model independent situation  $H_0 = 71 \pm 7.83 km/s/Mpc$ . The Hubble's constant is the current rate of expansion of the universe, but the rate predicted by the standard model is far slower than the rate discovered by the most accurate local observations. There are a lot of ways to evaluate the value of Hubble's constant, which creates the primary problem of its evaluation difficult. The existing dilemma is that we haven't been able to assign a single value for  $H_0$ . In this paper, we used the local approach method to calculate the value for  $H_0$  using the data from the scolnic et. el. 2018 data set.

A new approach of cosmologists is to resolve the difference between the predicted and observed values without jeopardising the consistency of the standard model. This can be done with the help of other cosmological phenomenons, But weather such a phenomenon exist is another topic of discussion. Certain cosmologists have now identified a certain aspect of all cosmological models that has been overlooked, which is that most dimensionless cosmic observables are substantially invariant when gravitational free-fall rates and photon-electron scattering rates are scaled uniformly. This finding made us realise that the phenomenon of cosmic background radiation can be used to bridge the gap between the predicted and observed values for the Hubble constant. This discovery aided us in reconciling cosmic background measurements, resulting in the development of a cosmological model in which a scaling transformation for the  $H_0$  value can be performed without violating any measurements of values not protected by symmetry. The "mirror world" dark sector would result in effective scaling of gravitational free-fall speeds while keeping the precisely determined mean photon density a constant. Further modifications to this model may help us to solve the two remaining constraints: the inferred primordial deuterium and helium abundances. This finding has opened the doors to a new possibility: the existence of a parallel universe that is remarkably similar to our universe but can only be seen through the gravitational influence on our world.

Research is also being carried out to identify whether this gap is partly due to the mistakes in measurements while looking for any new missing elements to complete this model. This has grown in importance as the disparity between the values has become more evident as the quality of the data incorporated in the analysis has also increased.

At the end of the day, it is astonishing to see that the discrepancy which we had hoped in reducing has only seen growth with the number of studies being conducted on both early and late universe. This tension might be a true mirror of reality, necessitating exotic new physics and a radical reworking of our knowledge of the cosmos.

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