

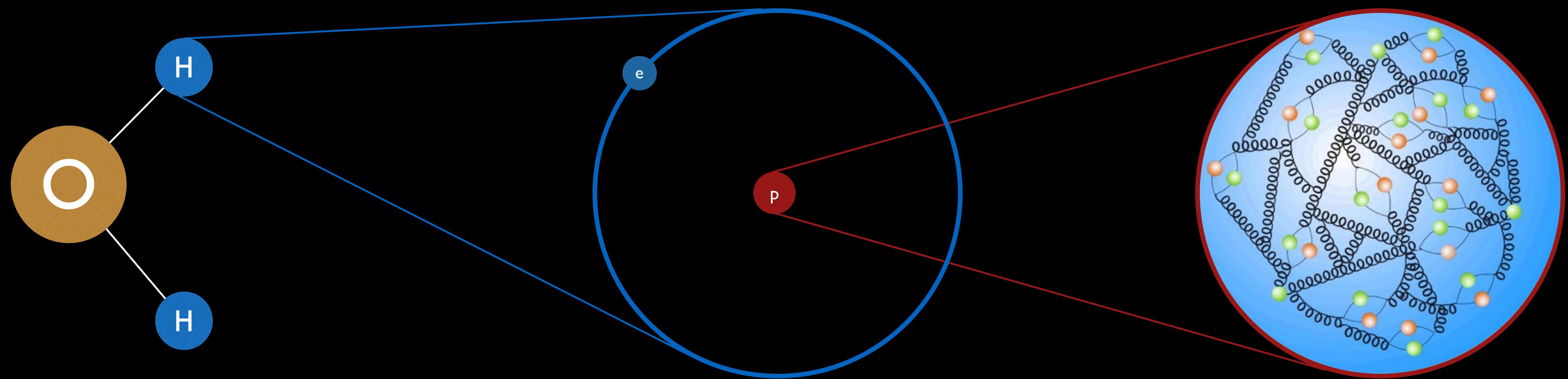
Hidden Higgses and Dark Matter *at* Current and Future Colliders

Adarsh Pyarelal



THE UNIVERSITY
OF ARIZONA

What are we made of?

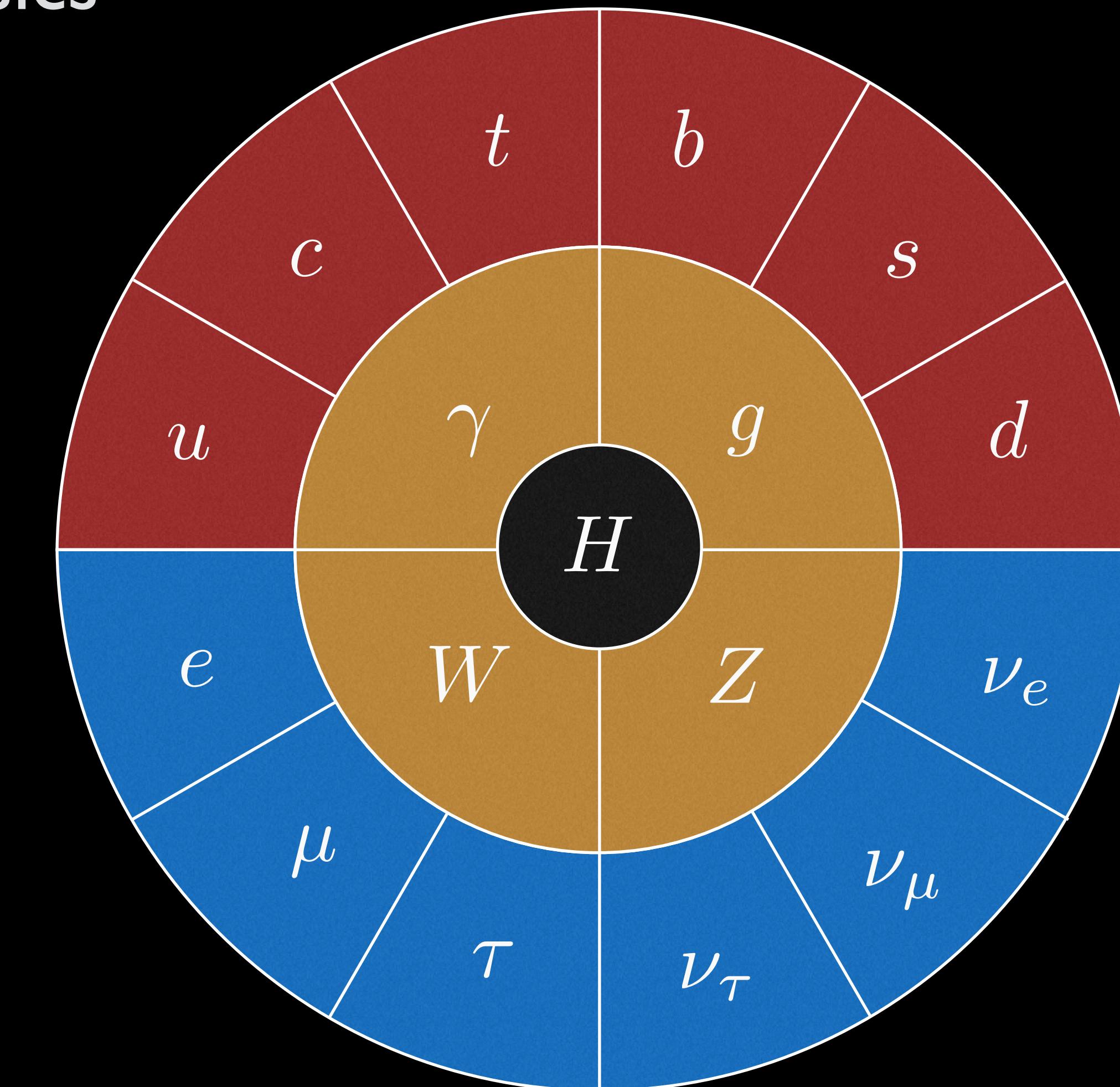
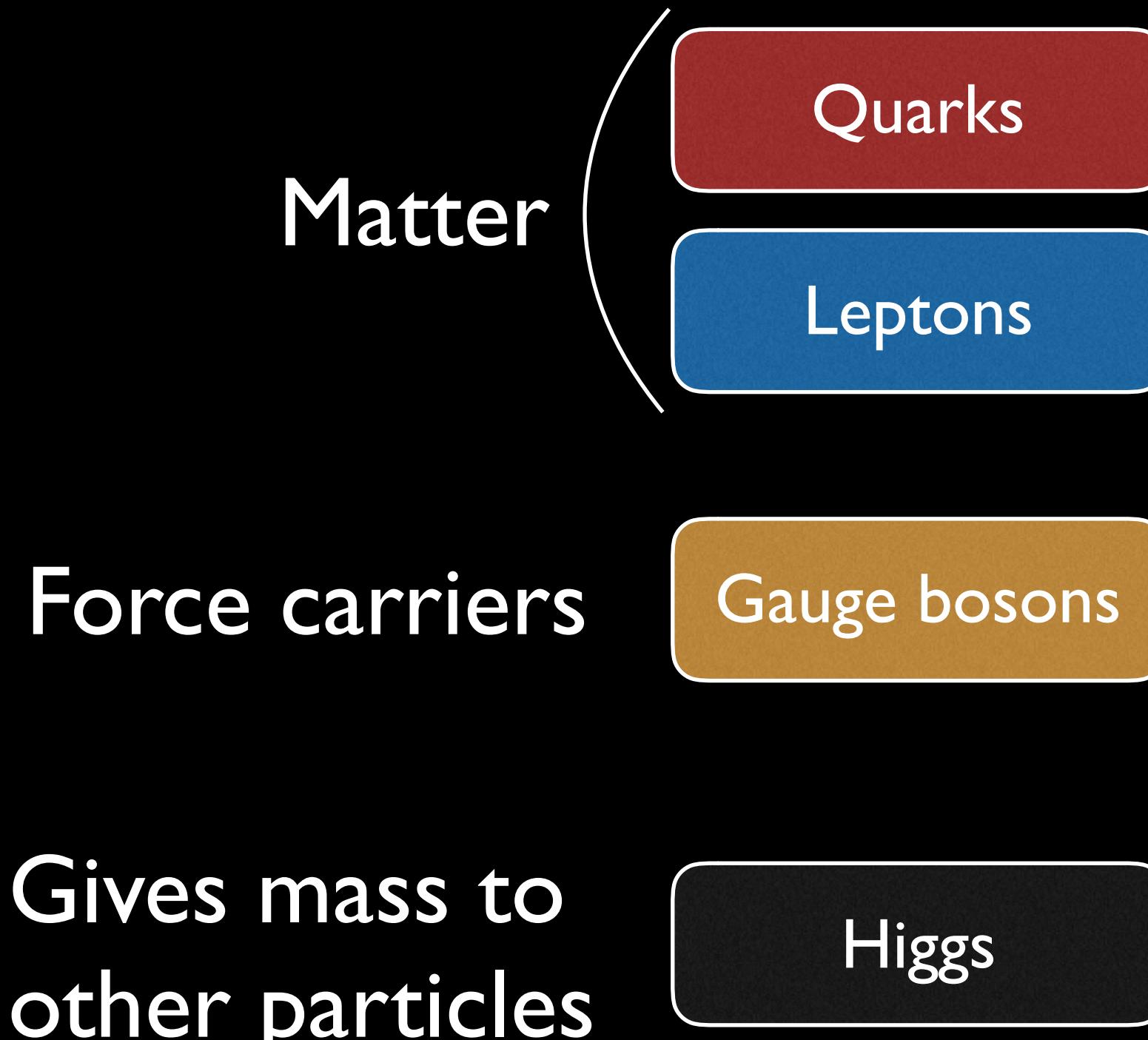


Water molecule
(H_2O)

Hydrogen atom

Proton

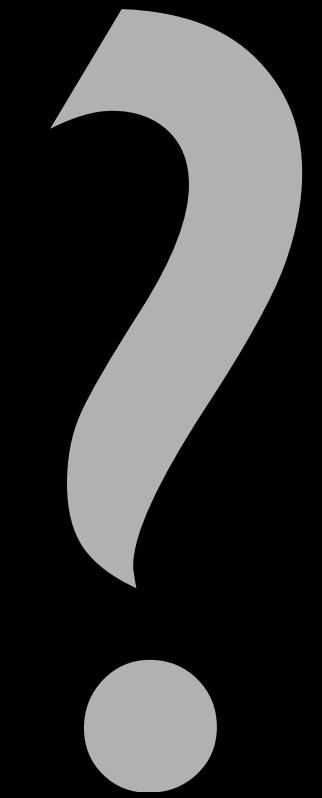
The Standard Model of Particle Physics



Unanswered Questions

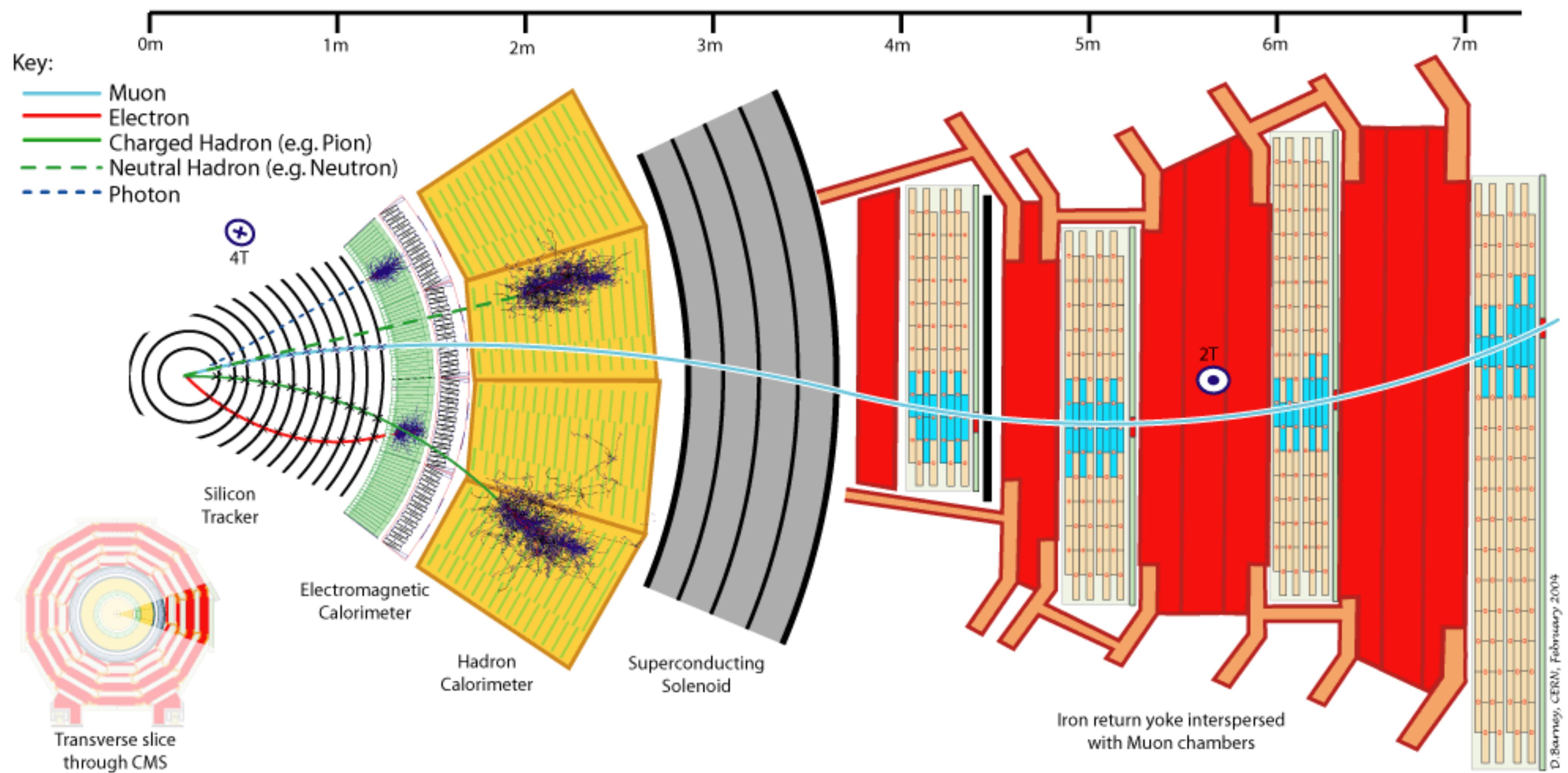
Higgs mass?

Neutrino oscillations?



Dark Matter?

And others...



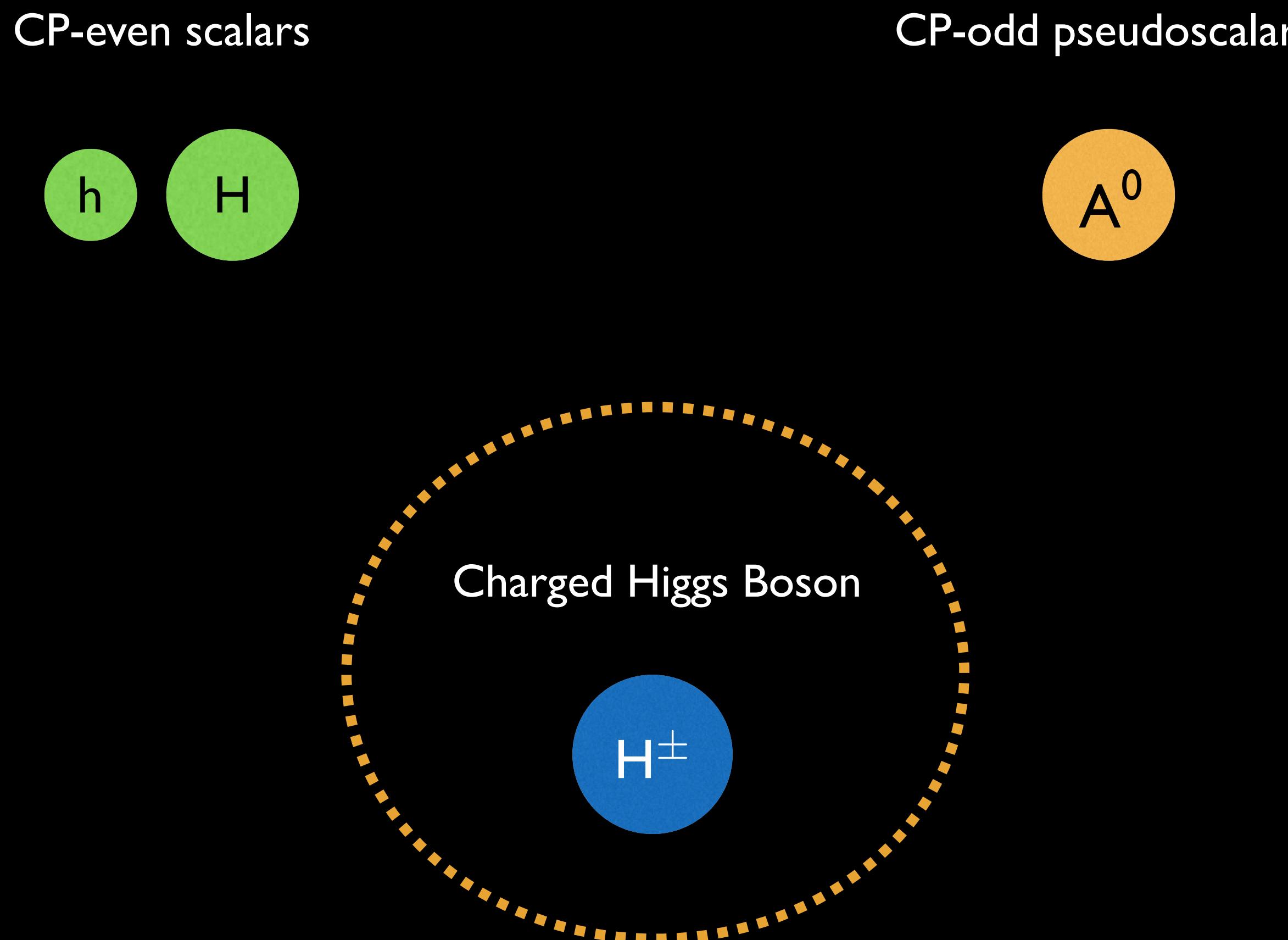


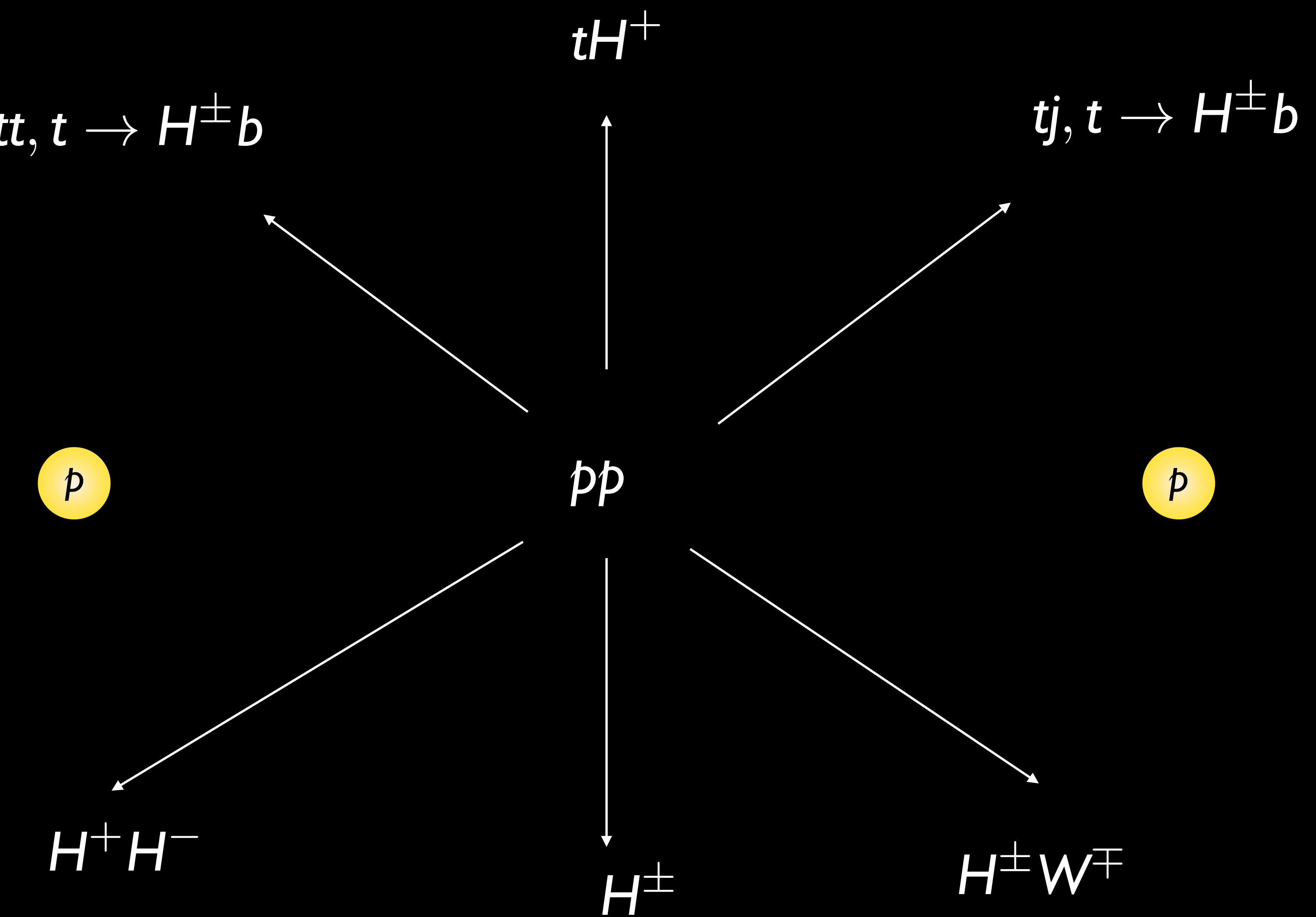
Two Higgs Doublet Models

Minimal Supersymmetric Standard Model

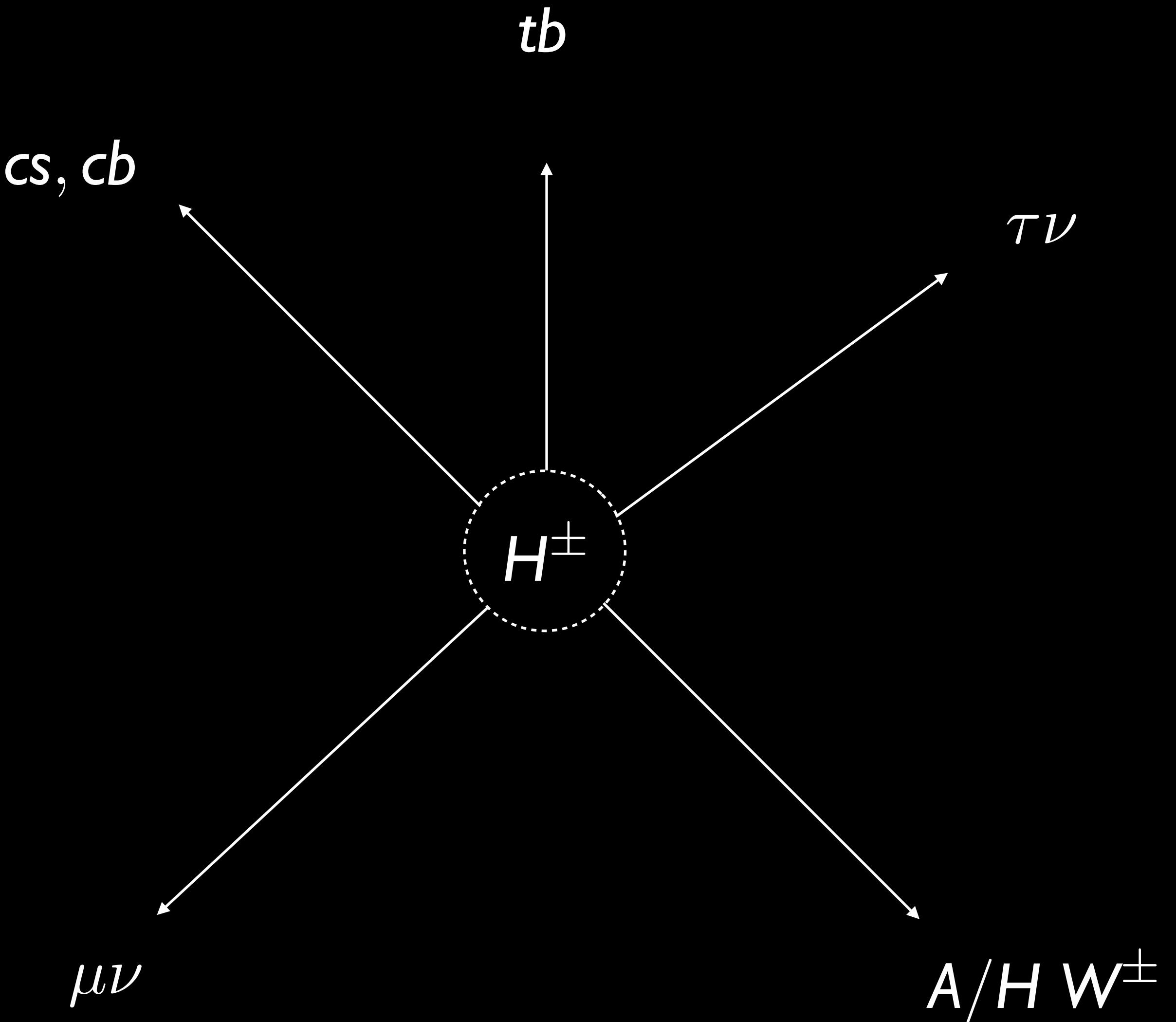
Focus of this talk

Higgs bosons in Two-Higgs Doublet Model

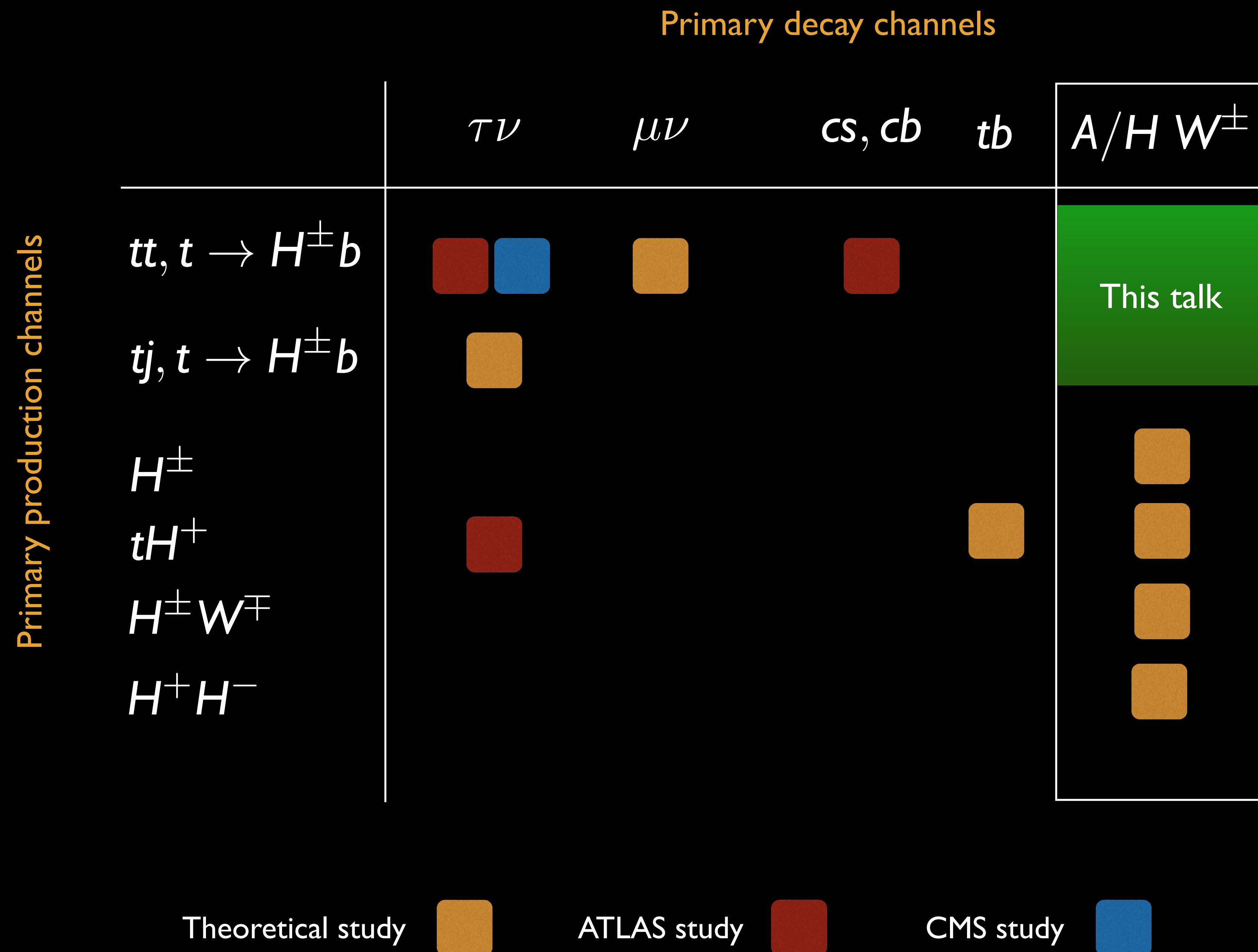




Primary production modes for the charged Higgs boson



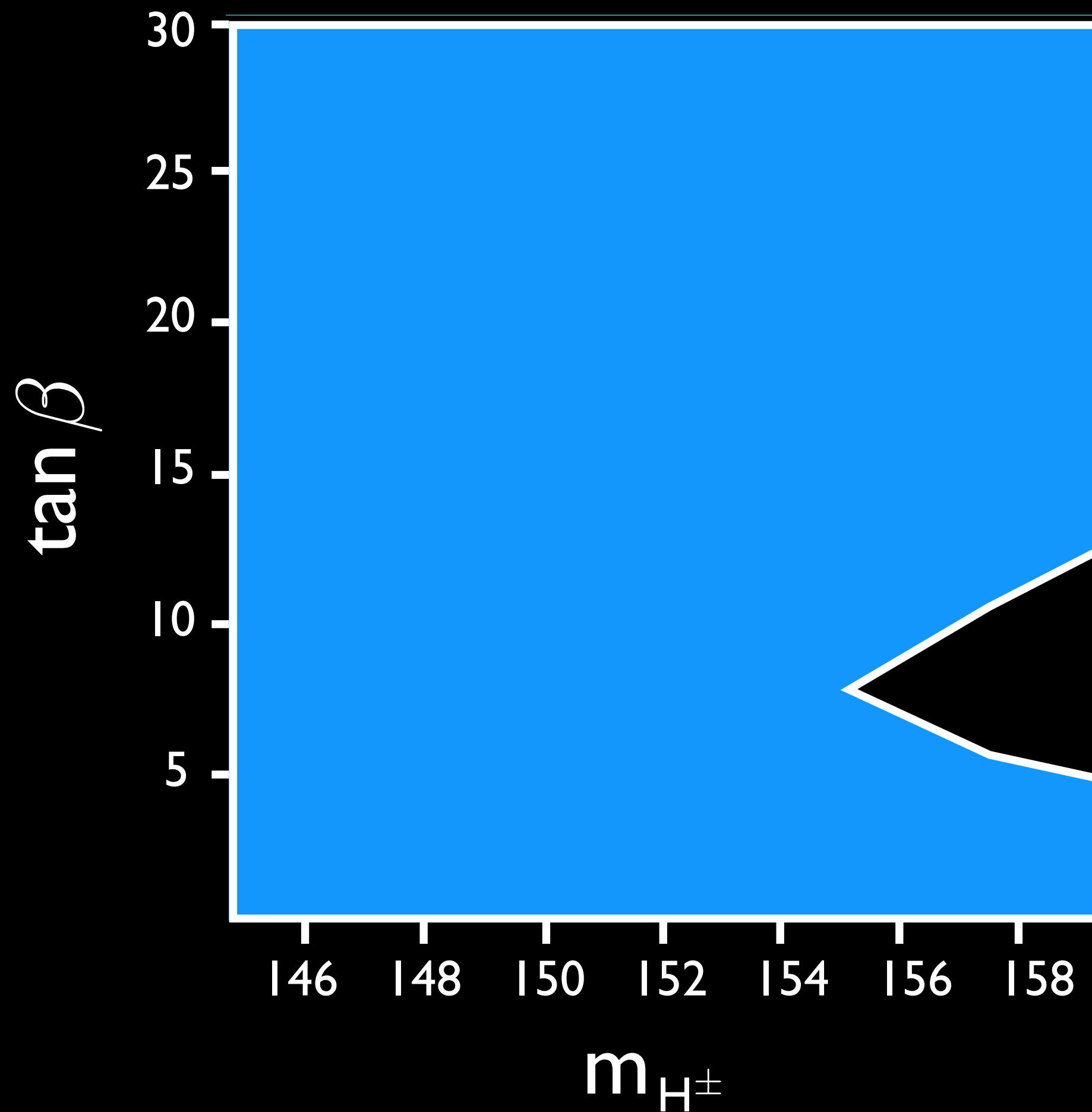
Primary decay channels for the charged Higgs



Charged Higgs Challenges

	H^\pm	top	H^\pm
Light ($< 174 \text{ GeV}$)	✓	Soft final state particles	Heavy ($> 174 \text{ GeV}$)
	✓	Off-shell production - rare	✓
	✓	Cleaner leptonic decays suppressed	✓

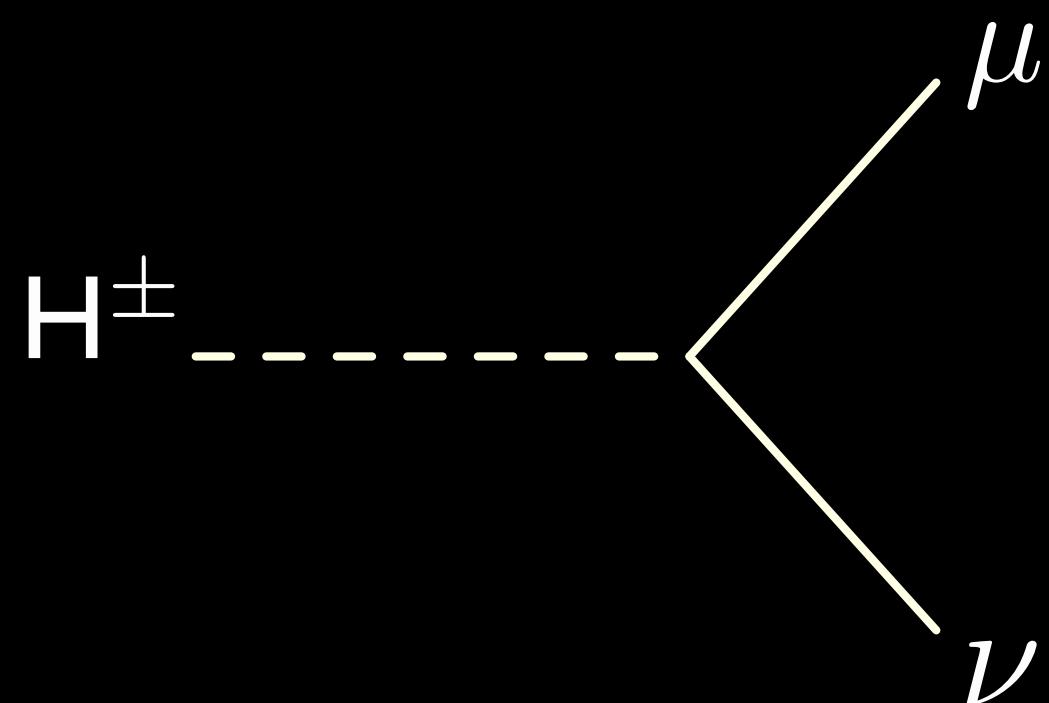
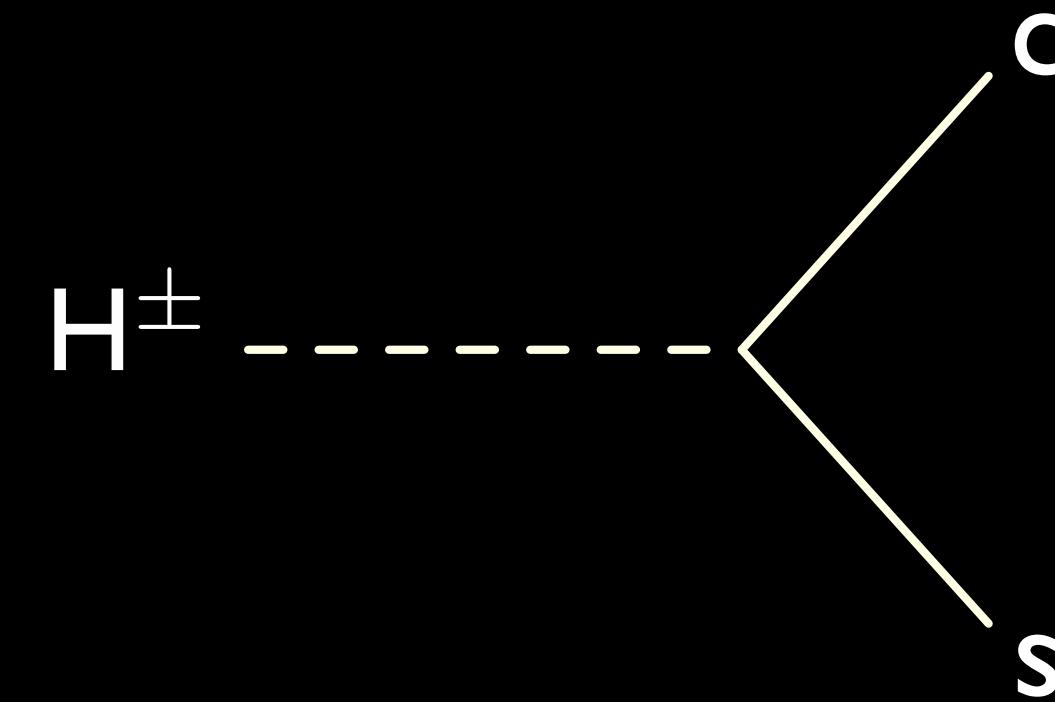
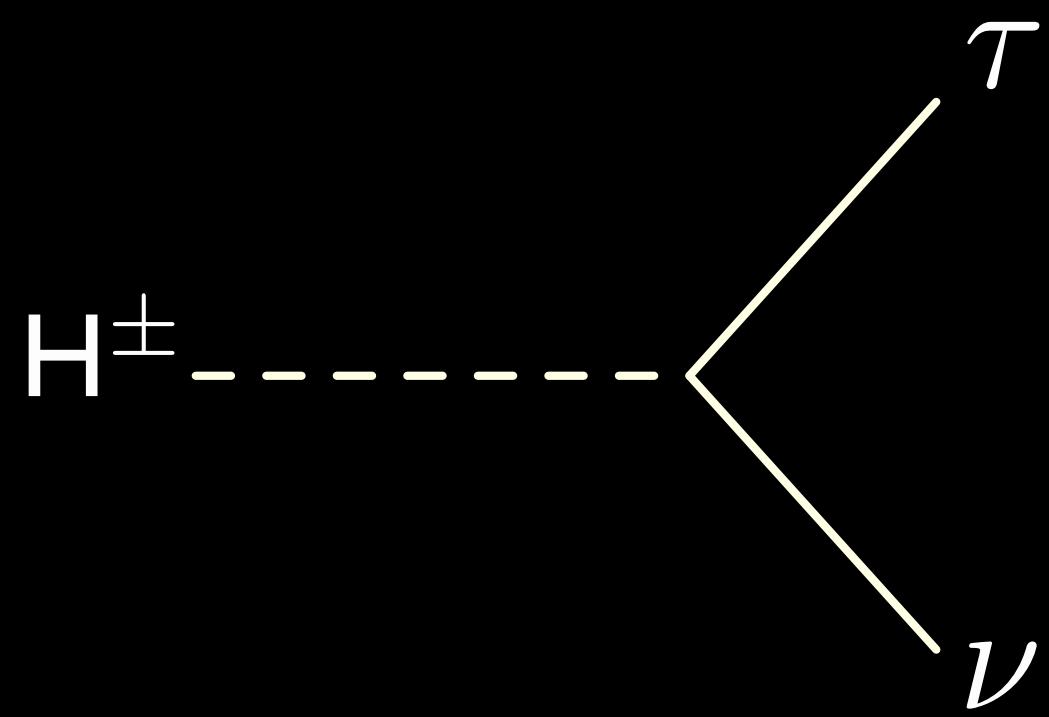
Current Limits on Light Charged Higgs



Excluded, assuming
 $BR(H^\pm \rightarrow \tau\nu) = 100\%$

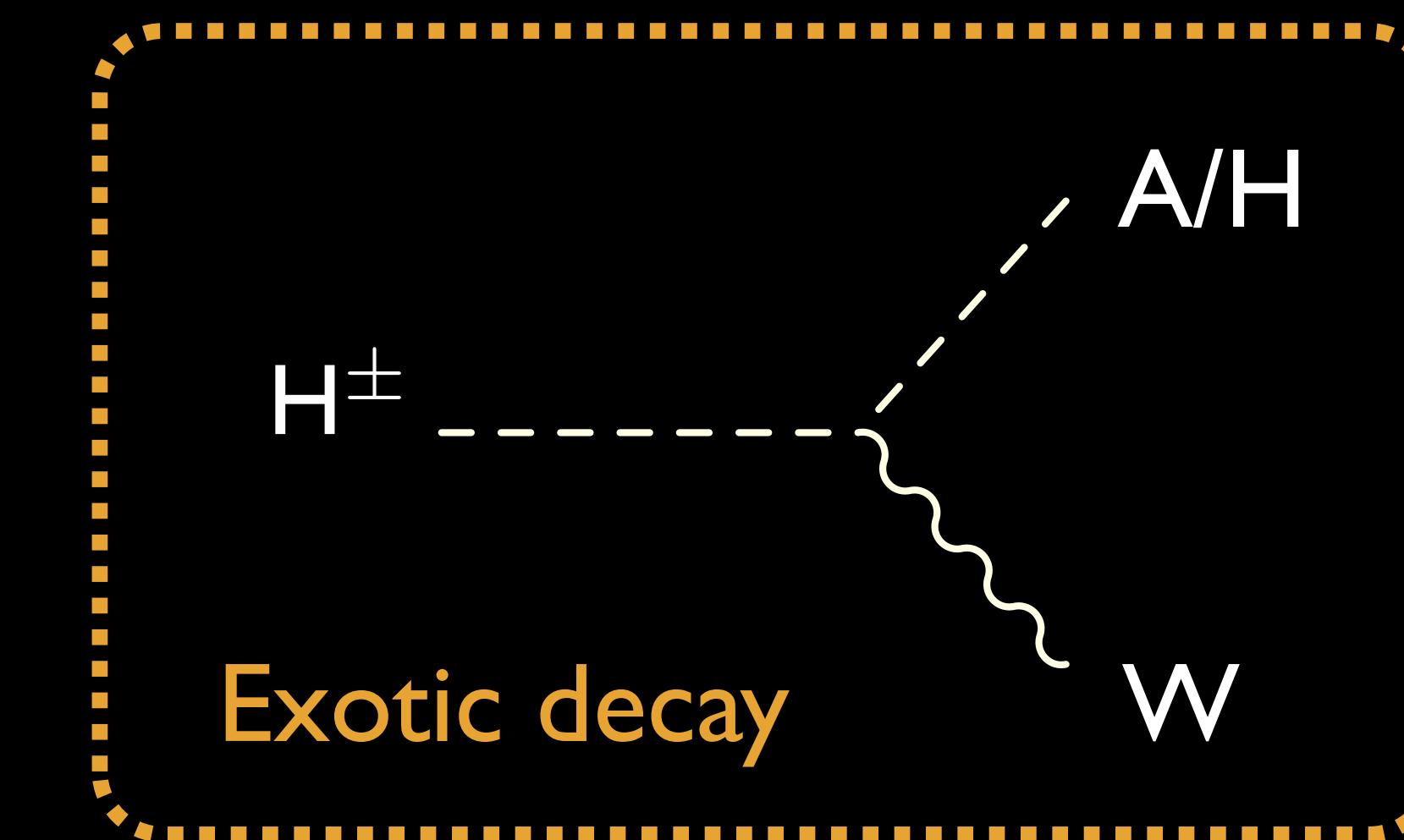
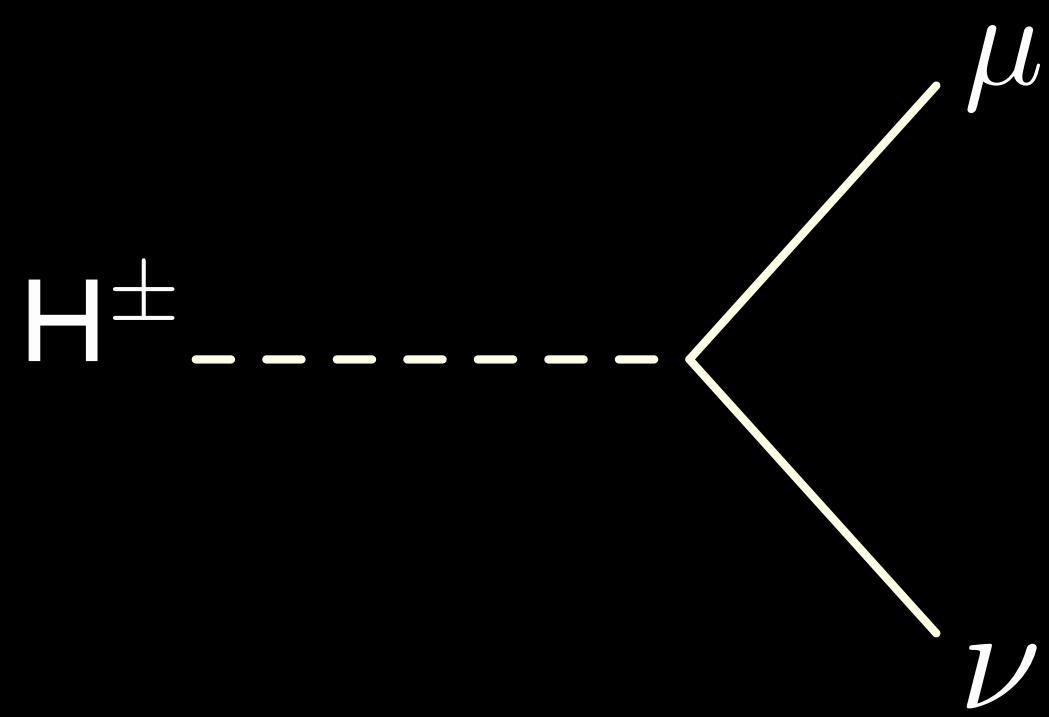
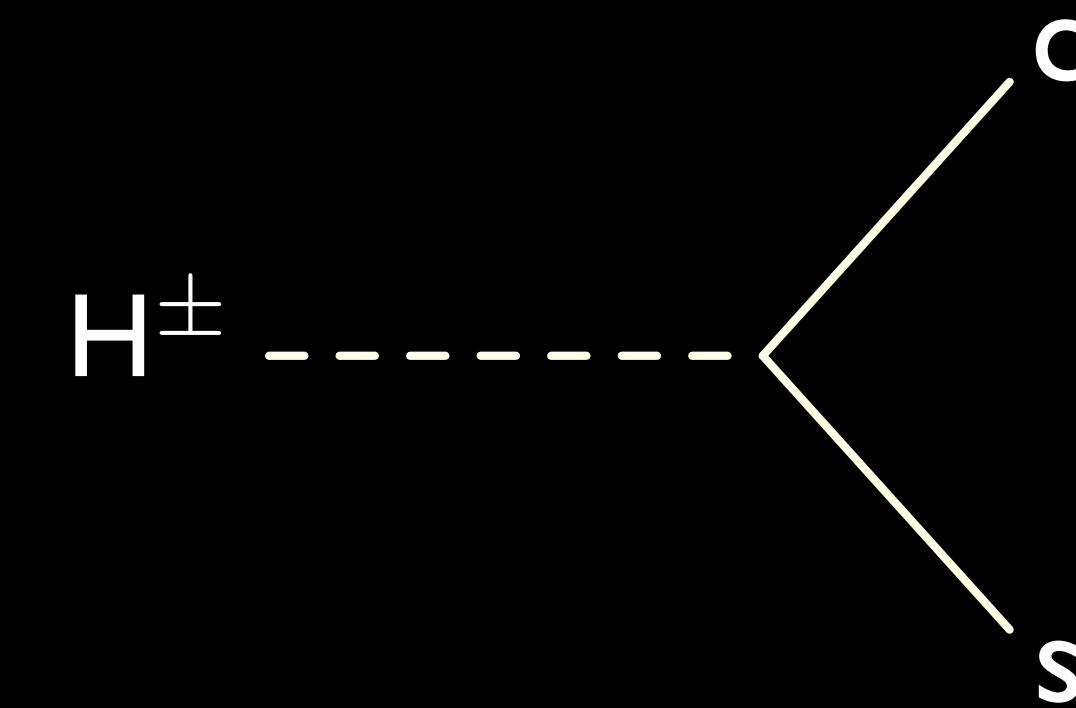
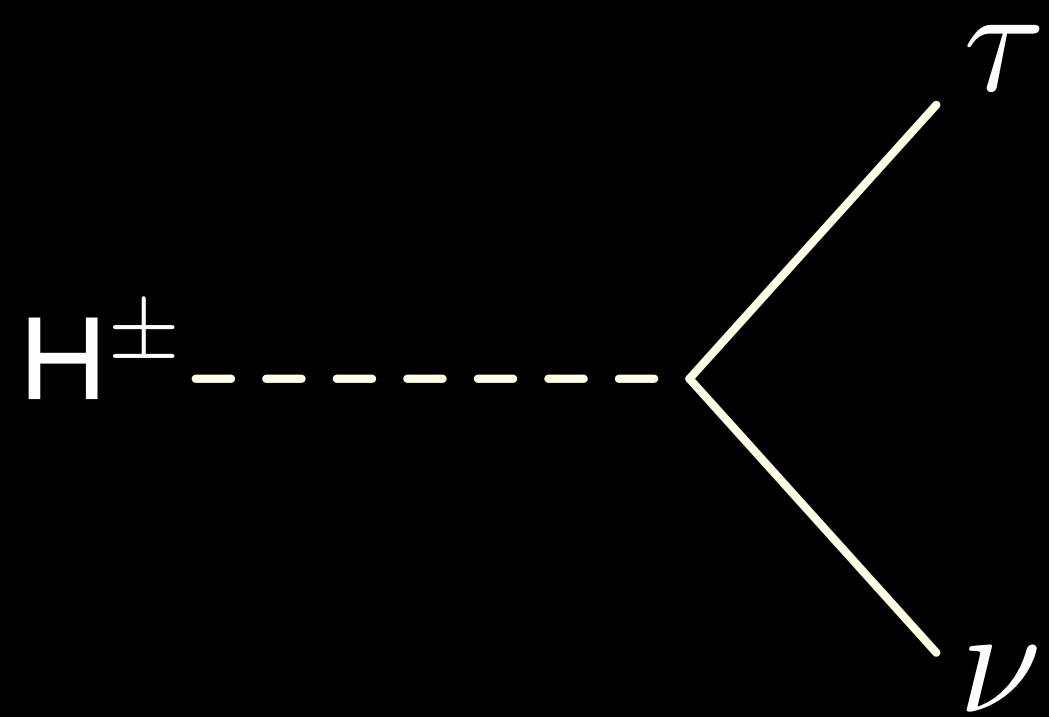
Source:
CMS Analysis, September 2014

Primary decay modes of the Light Charged Higgs



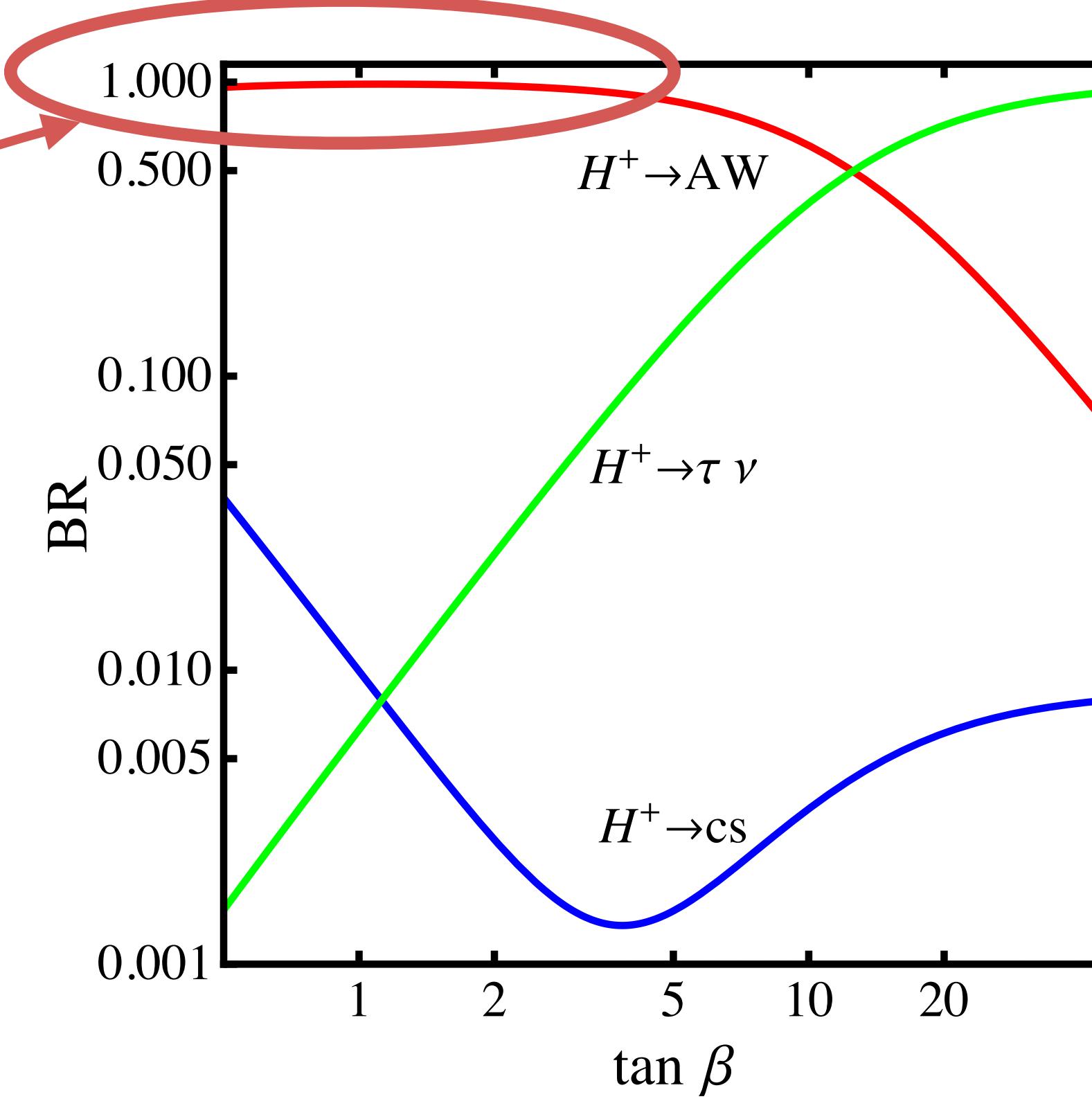
Decays to Standard
Model particles

Primary decay modes of the Light Charged Higgs

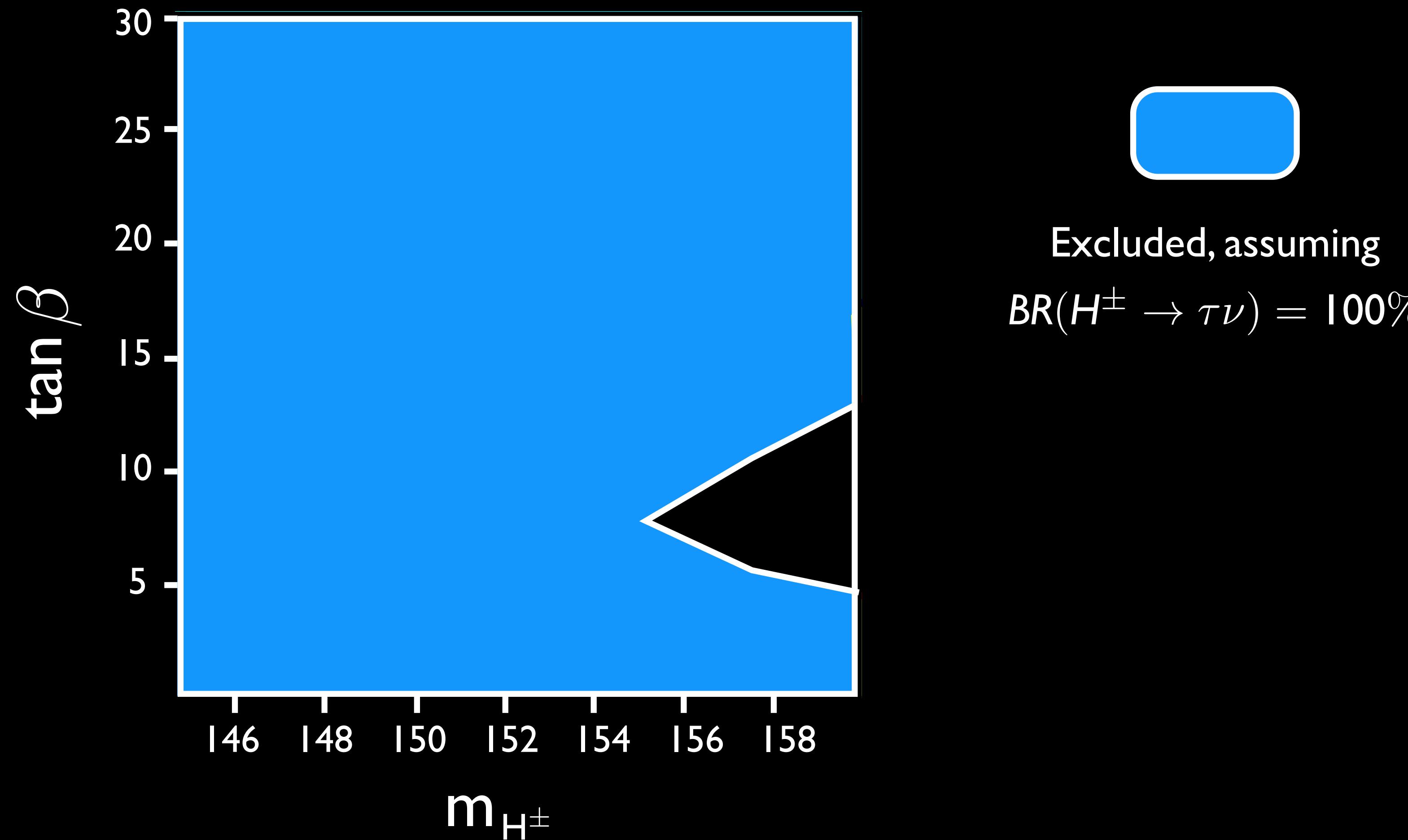


Branching Ratios

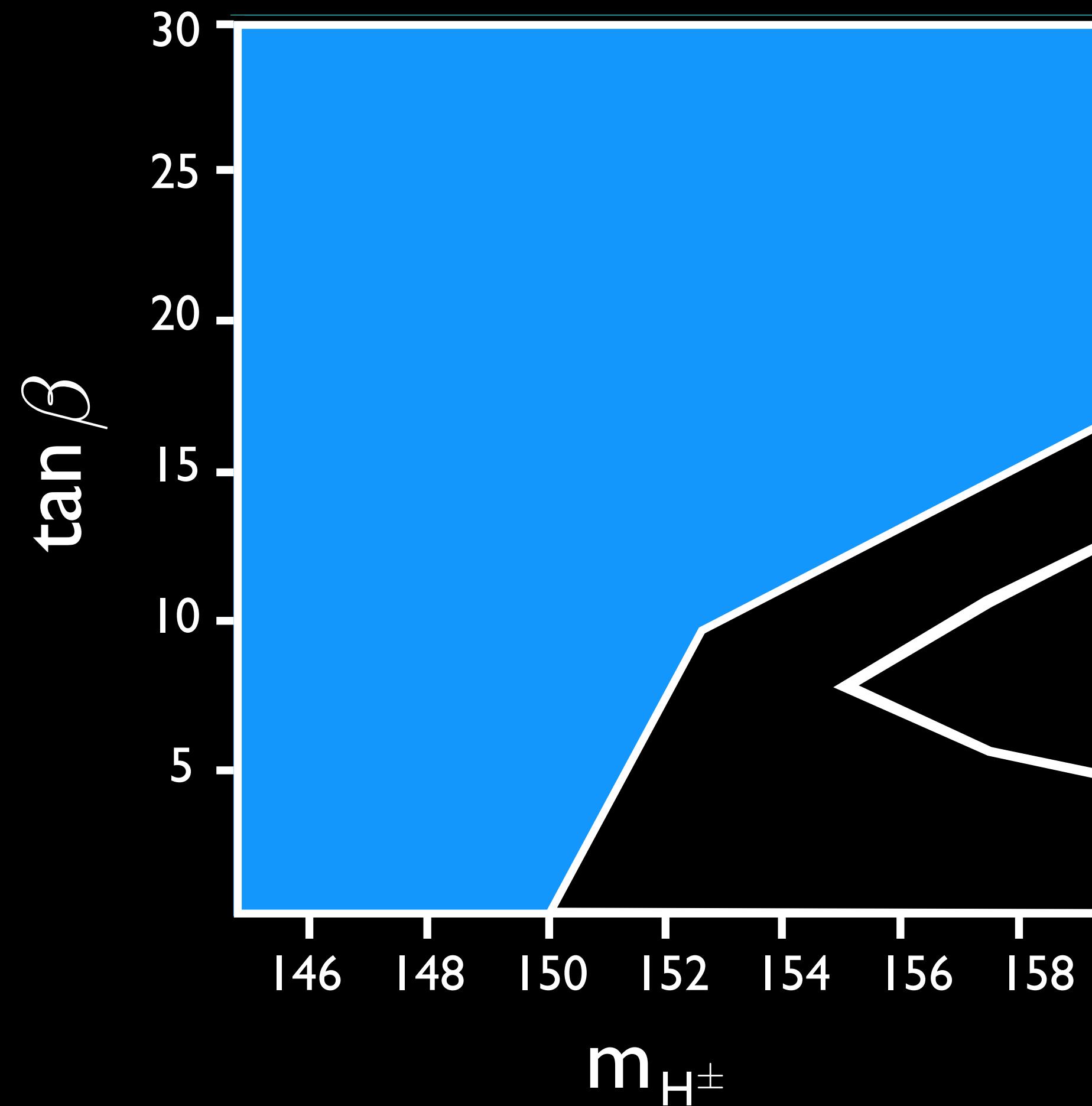
Exotic decay
dominates for
low $\tan \beta$



Current Limits



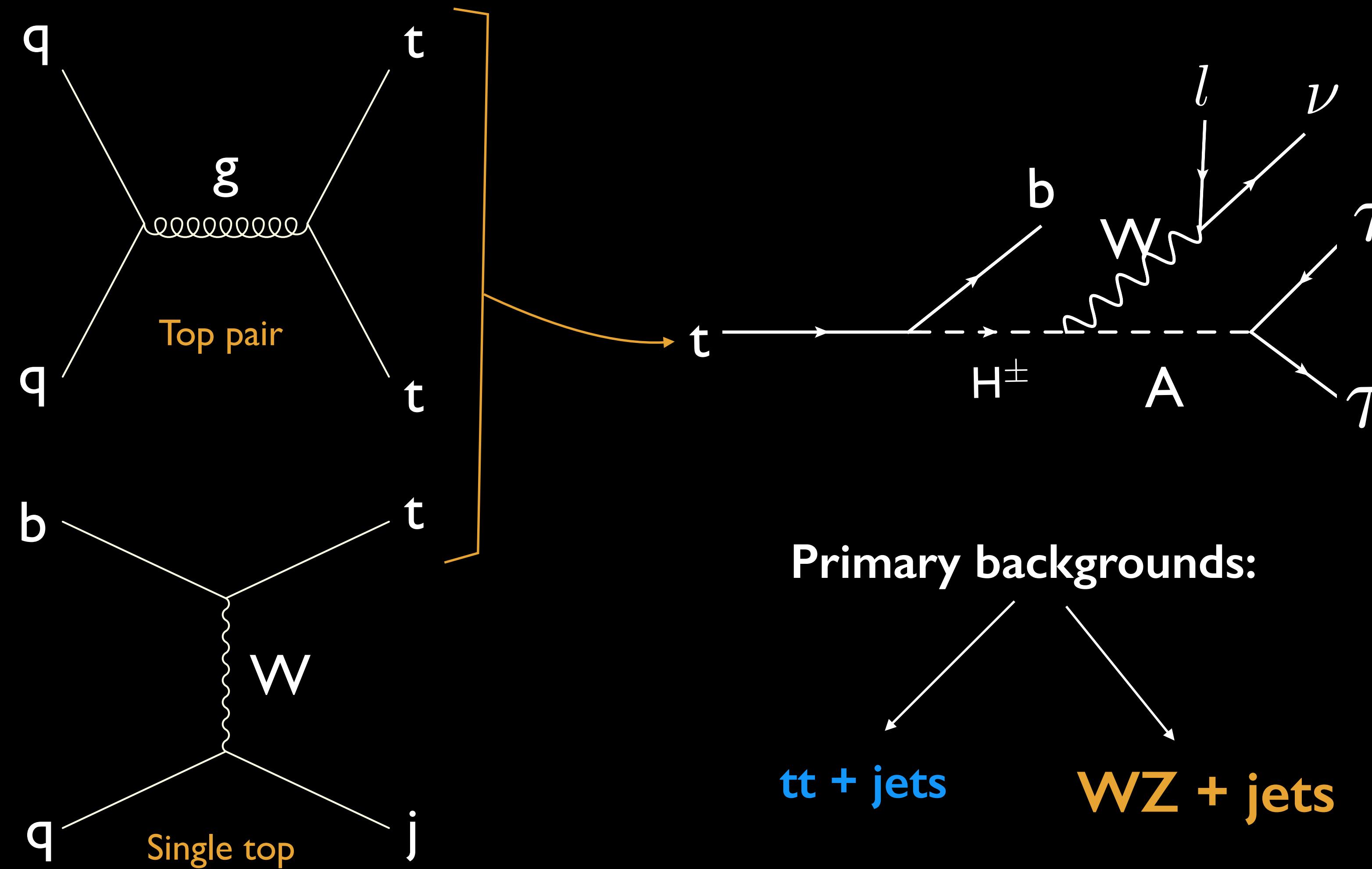
Weakened Limits



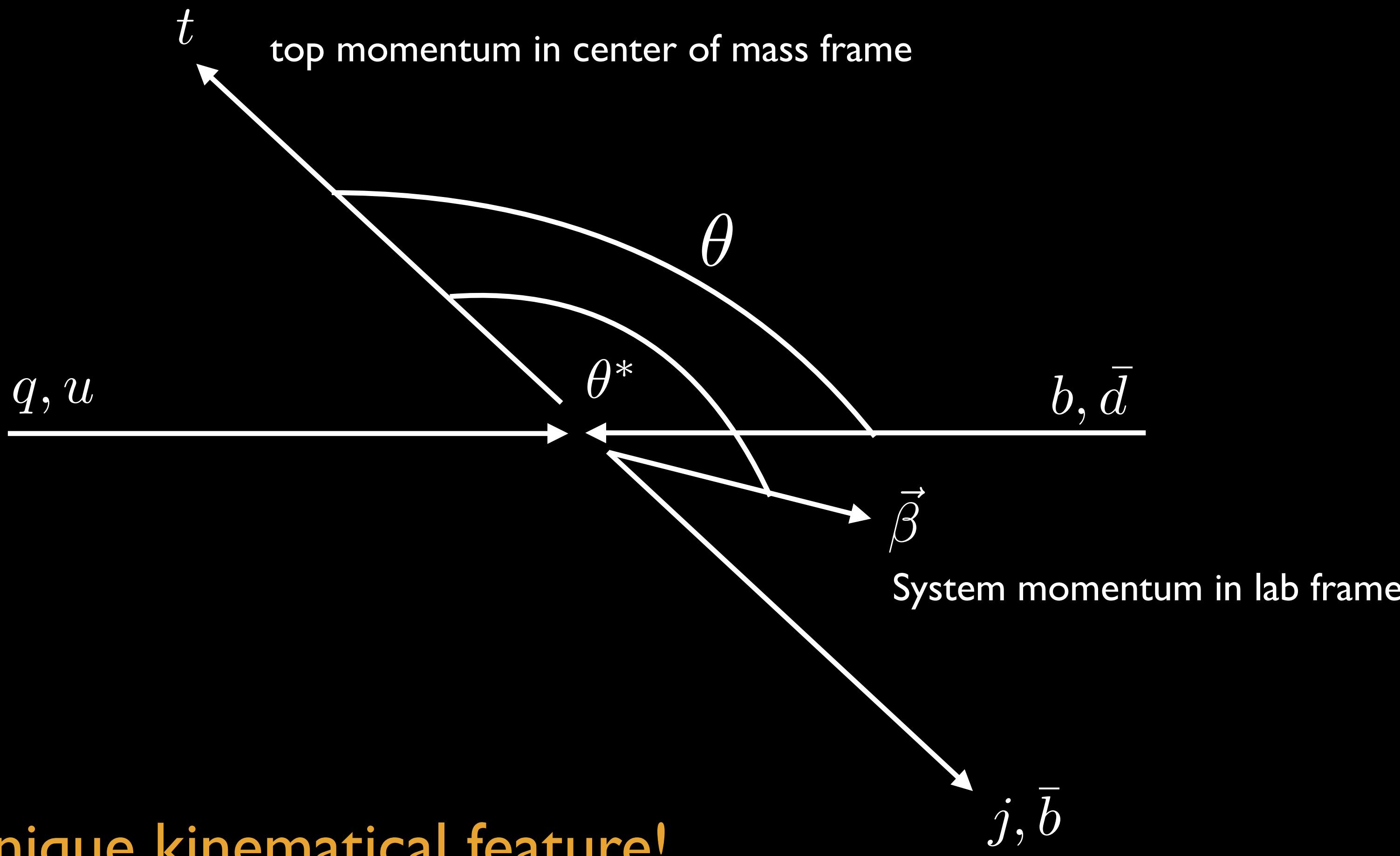
Excluded, if
 $H^\pm \rightarrow AW$ is possible

(Assuming existence of
A with mass 70 GeV)

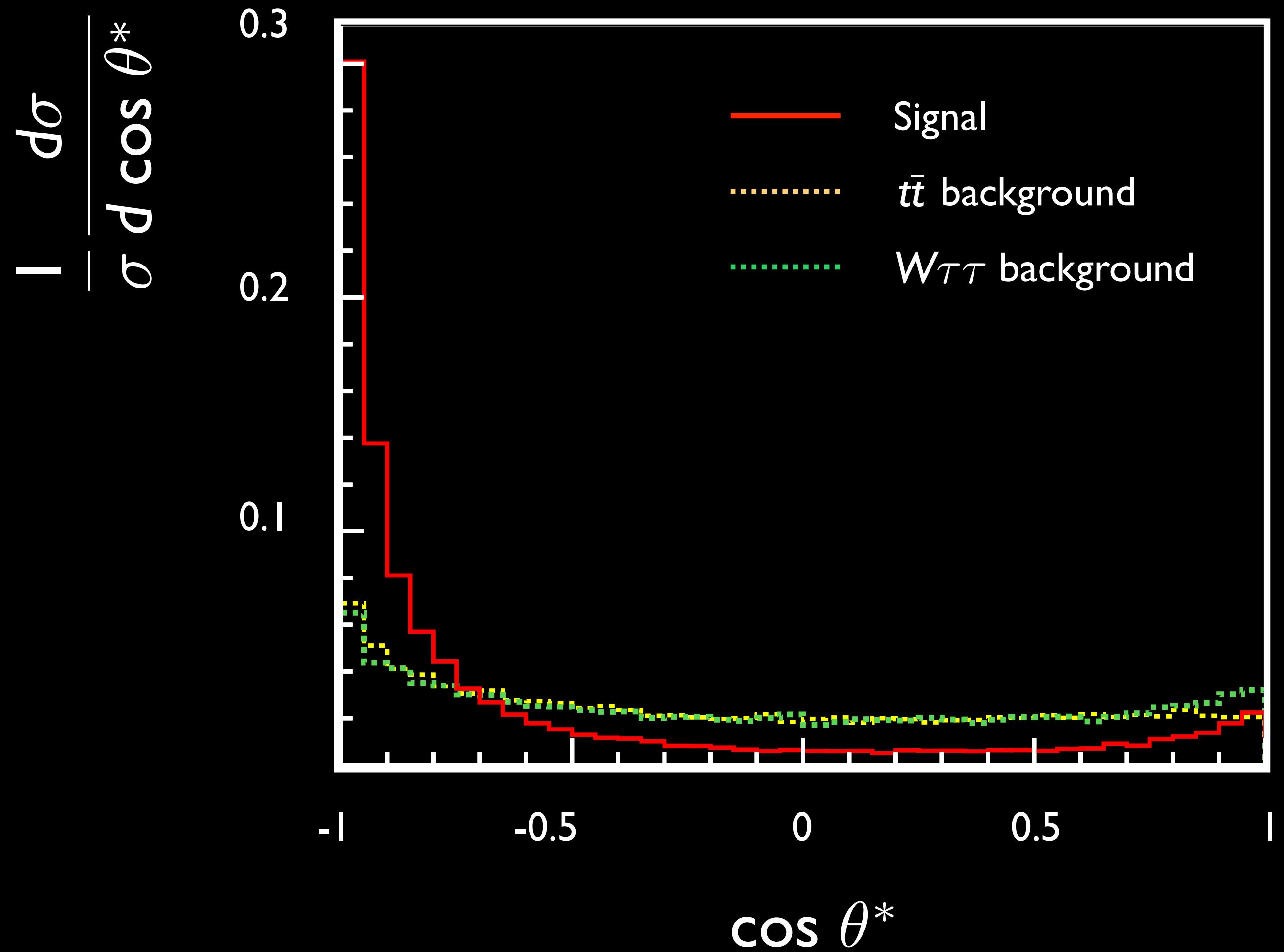
Our search channel



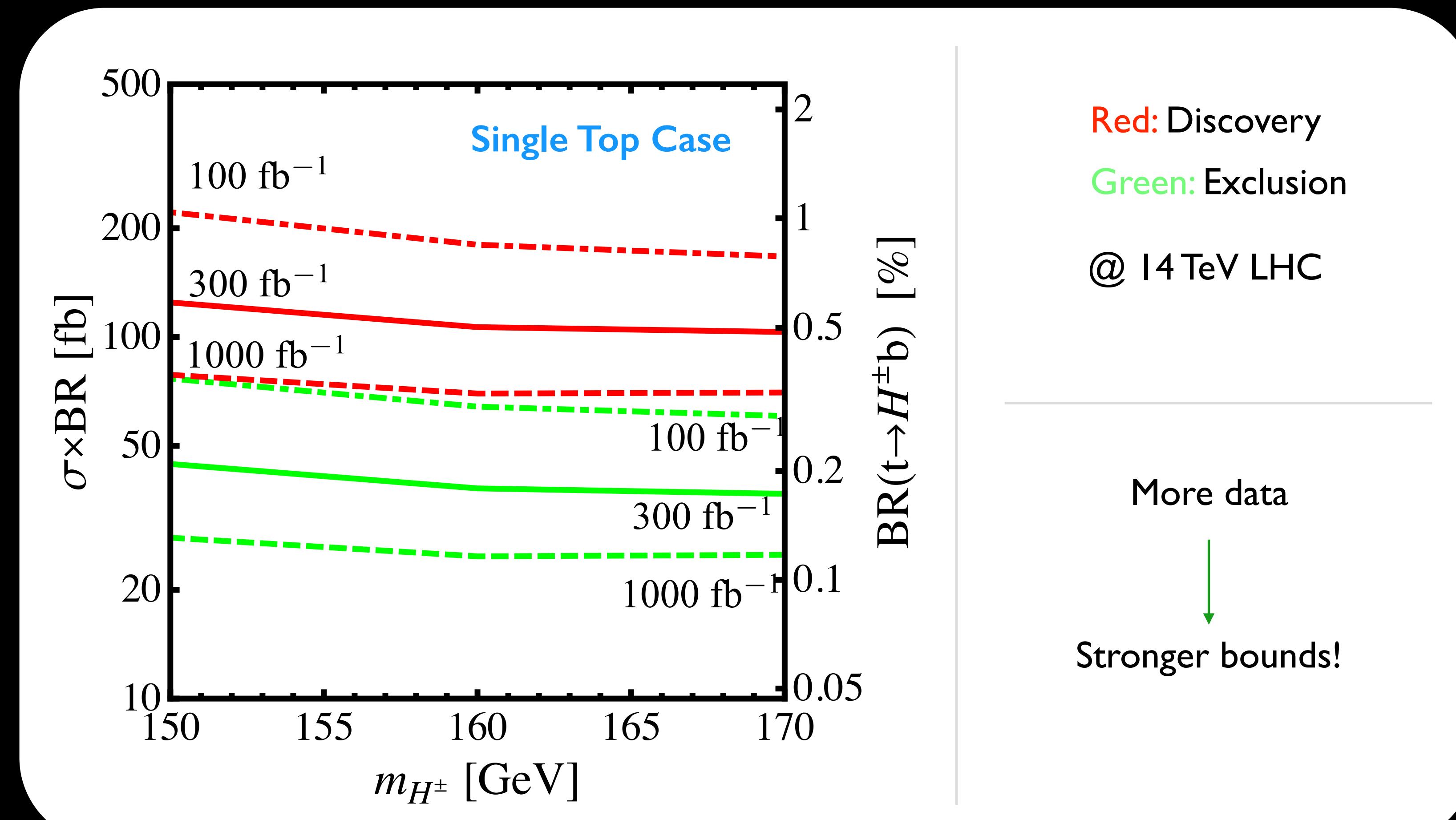
Angular Correlations in Single Top Production



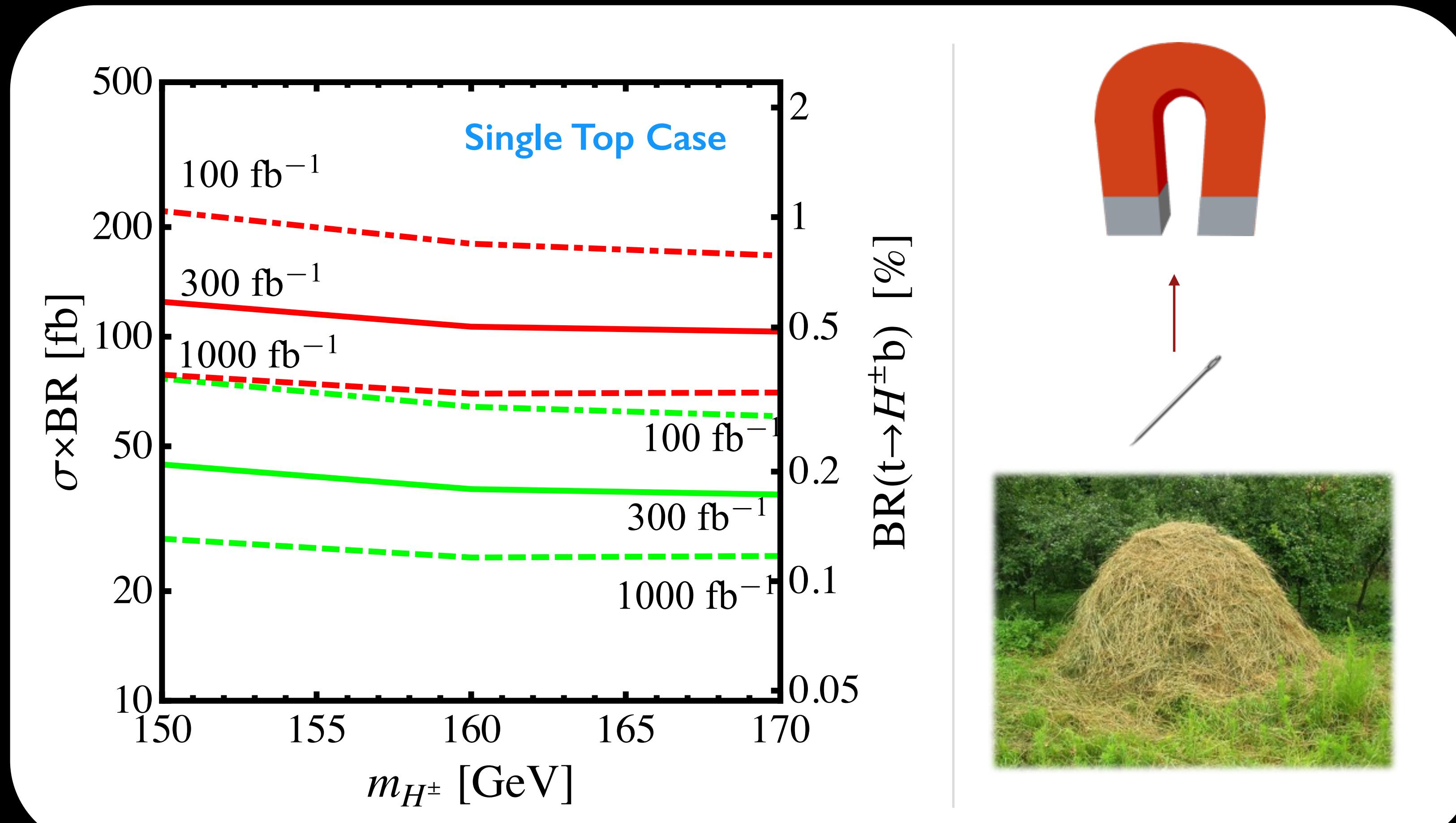
Angular distribution - Signal vs. Background



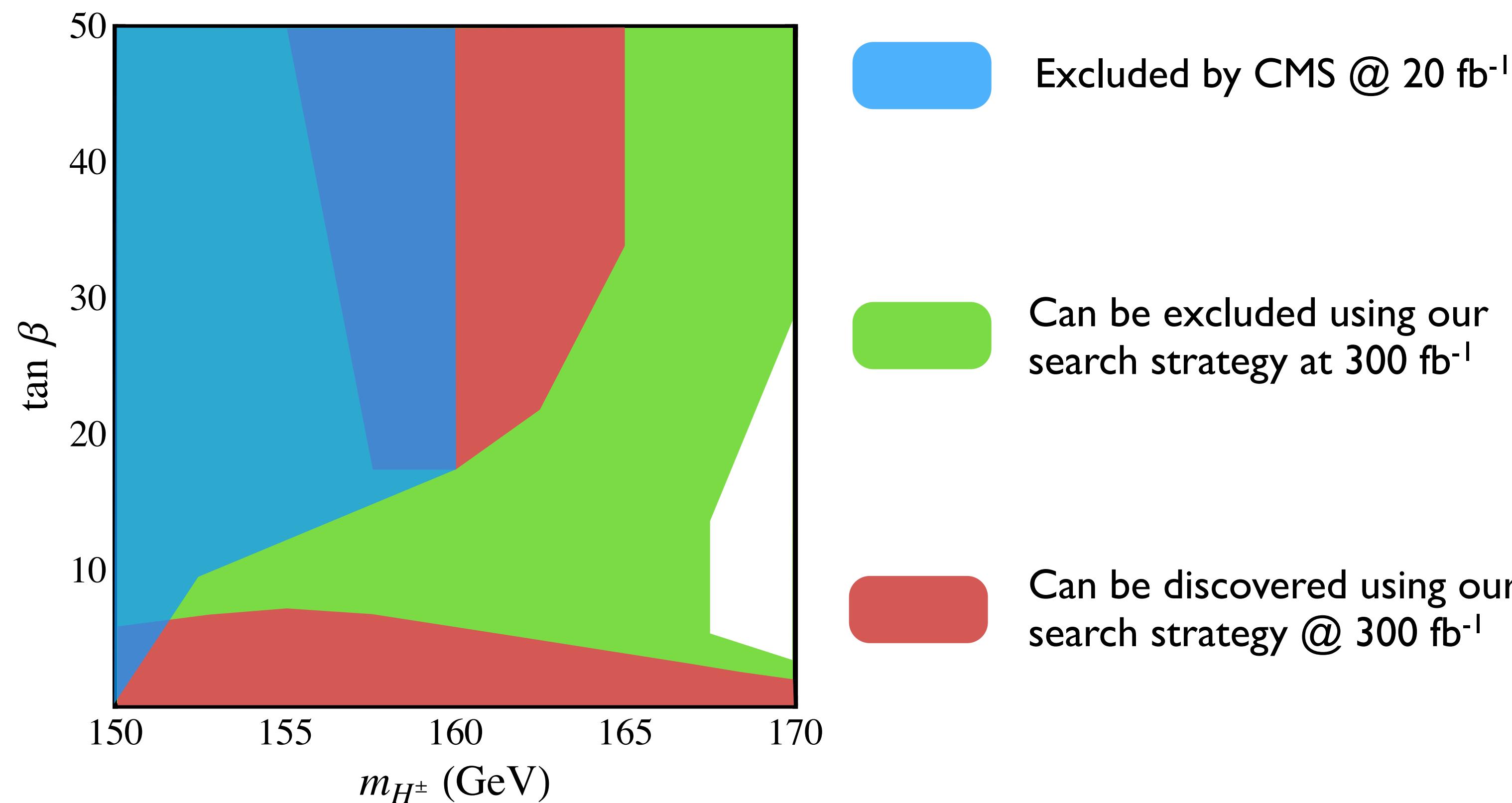
Model-independent Limits



Model-independent Limits

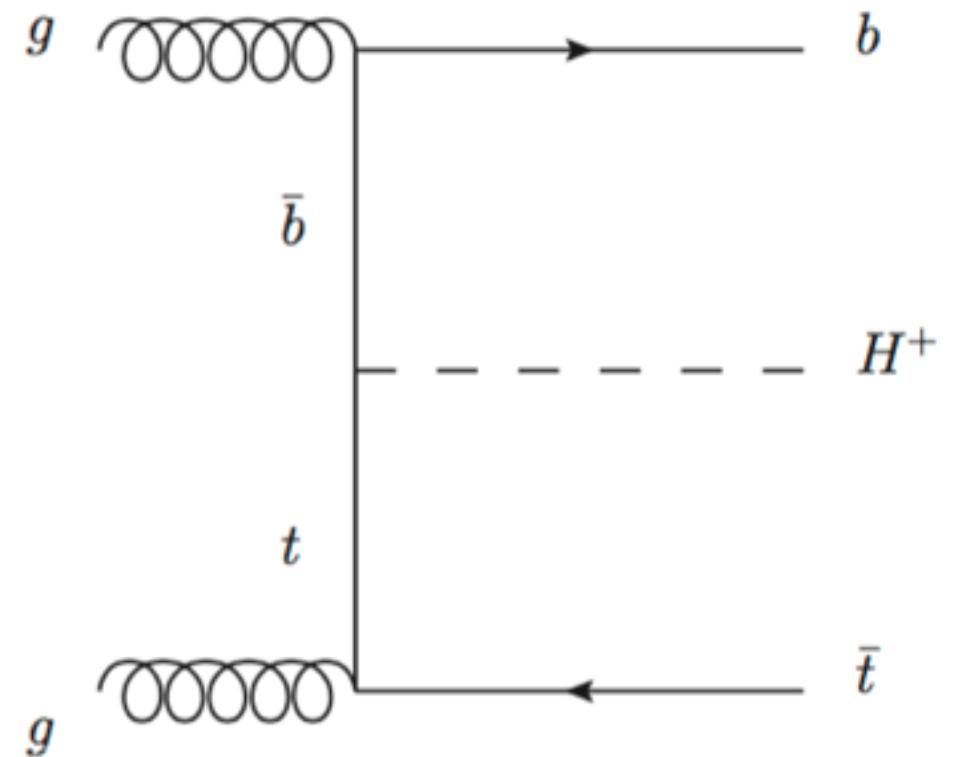


2HDM Implications



Heavy Charged Higgses

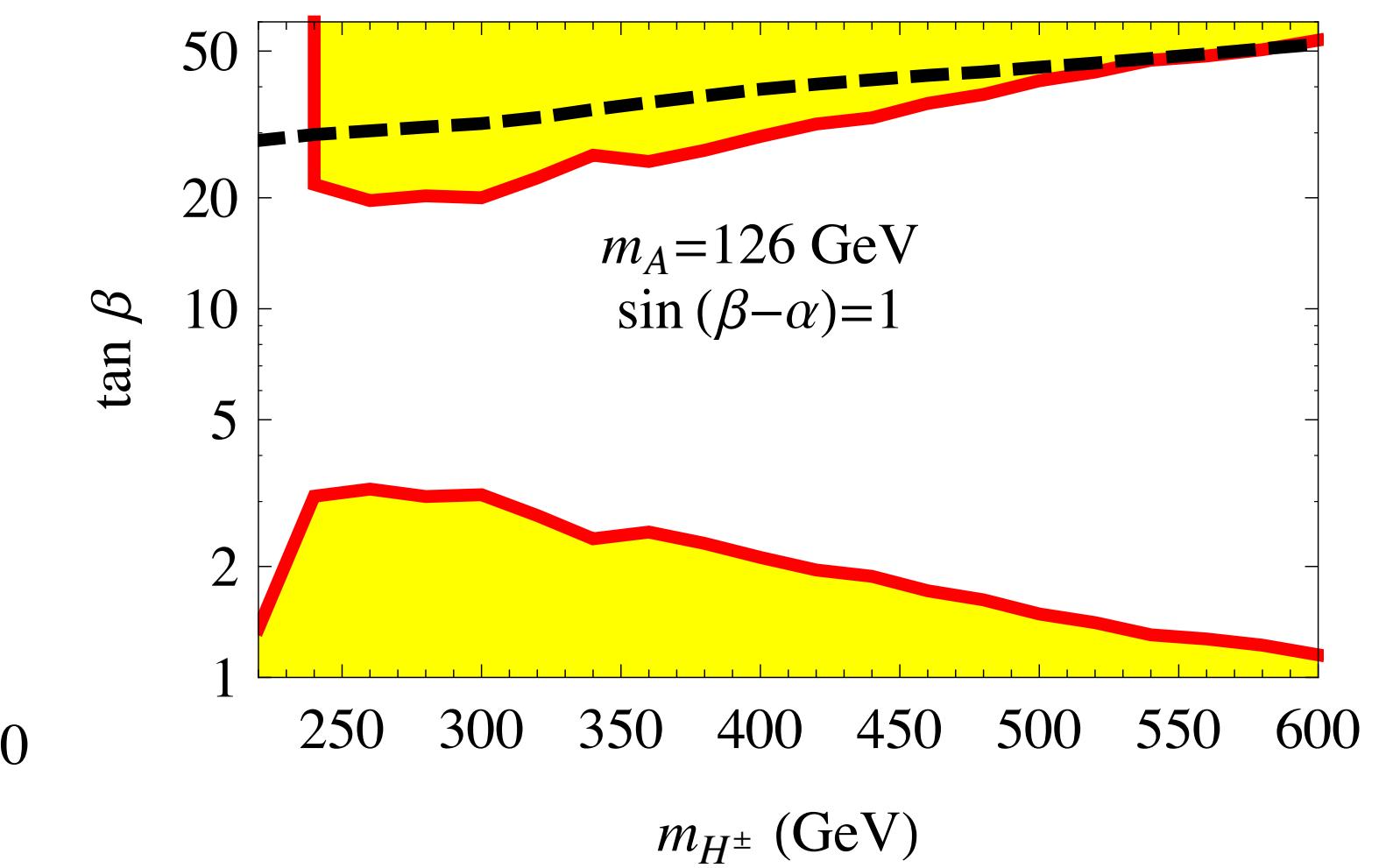
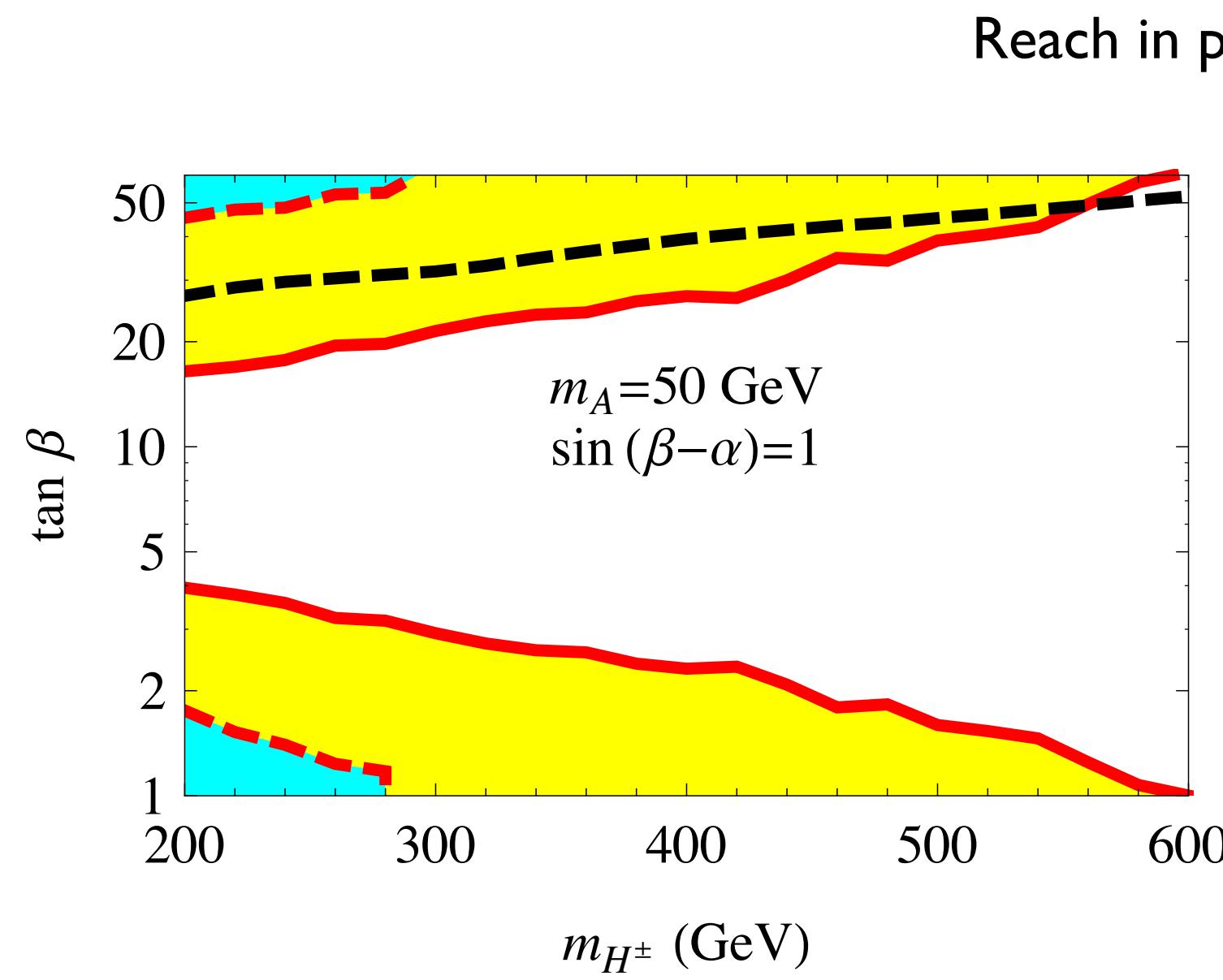
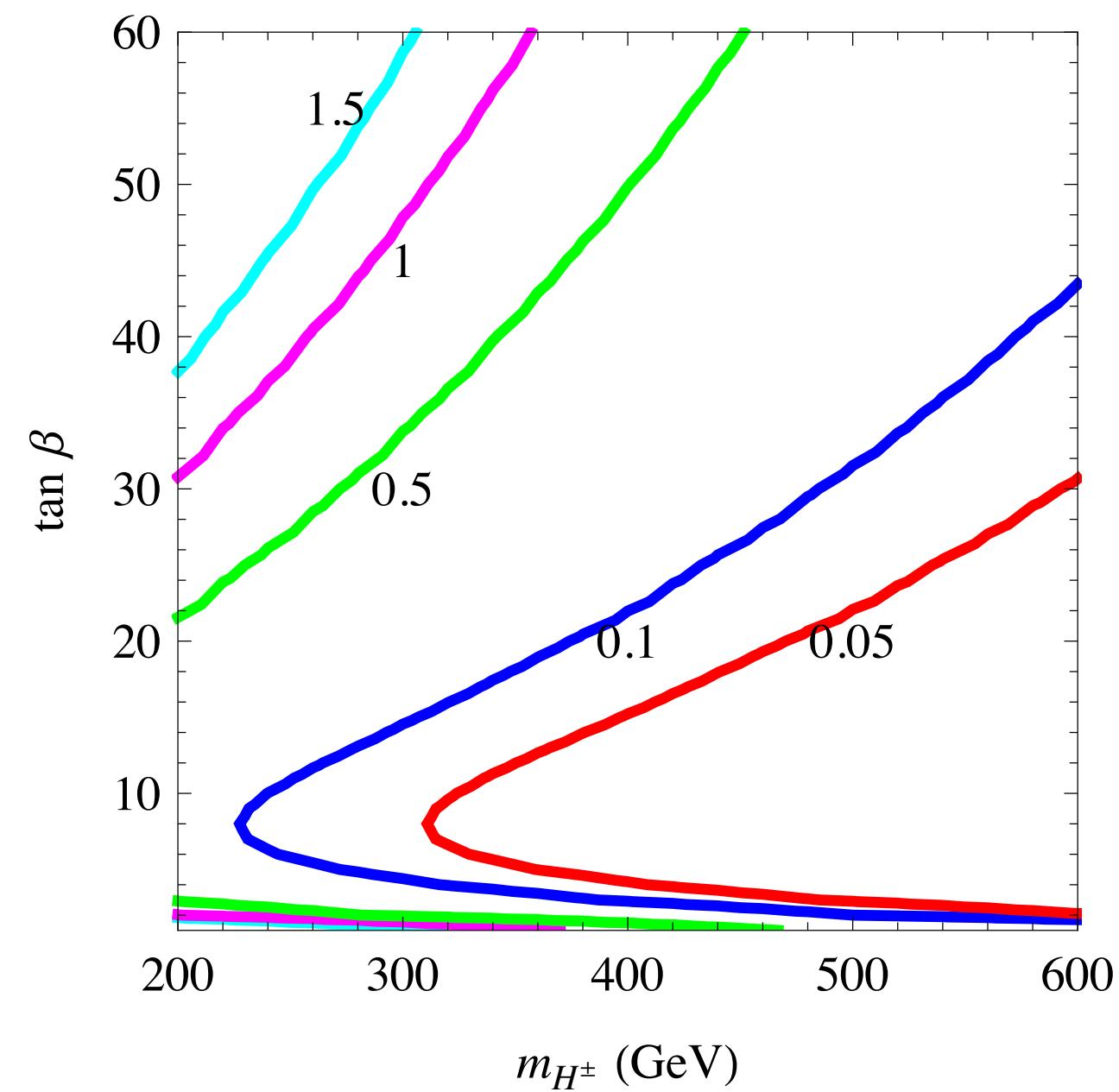
Production in association with a top quark



$H^\pm \rightarrow tb$, large BG
 $H^\pm \rightarrow \tau\nu$ or cs ,
 suppressed Br.

B. Coleppa, F. Kling, and S. Su,
 Charged Higgs search via AW/HW channel,
 JHEP 12 (2014) 148, [arXiv:1408.4119].

Production cross section (pb) @ 14 TeV



Exotic Higgs Decays - Big picture

- Consistent with theories that explain matter-antimatter imbalance in the universe.
- Can dominate over SM final states in regions of parameter space.

In the near future...

Table 5.1: Summary of all the possible exotic decay modes and the dominant final states for non-SM Higgs bosons

Parent Higgs	Decay type	Channels in 2HDM	Main Final States
Neutral Higgs H, A	HH	$H \rightarrow AA, hh$	$(bb/\tau\tau/WW/ZZ/\gamma\gamma)(bb/\tau\tau/WW/ZZ/\gamma\gamma)$
	HZ	$H \rightarrow AZ, A \rightarrow HZ, hZ$	$(ll/qq/\nu\nu)(bb/\tau\tau/WW/ZZ/\gamma\gamma)$
	$H^+ H^-$	$H \rightarrow H^+ H^-$	$(tb/\tau\nu/cs)(tb/\tau\nu/cs)$
	$H^\pm W^\mp$	$H/A \rightarrow H^\pm W^\mp$	$(l\nu/qq')(tb/\tau\nu/cs)$
Charged Higgs	HW^\pm	$H^\pm \rightarrow hW^\pm/HW^\pm/AW^\pm$	$(l\nu/qq')(bb/\tau\tau/WW/ZZ/\gamma\gamma)$

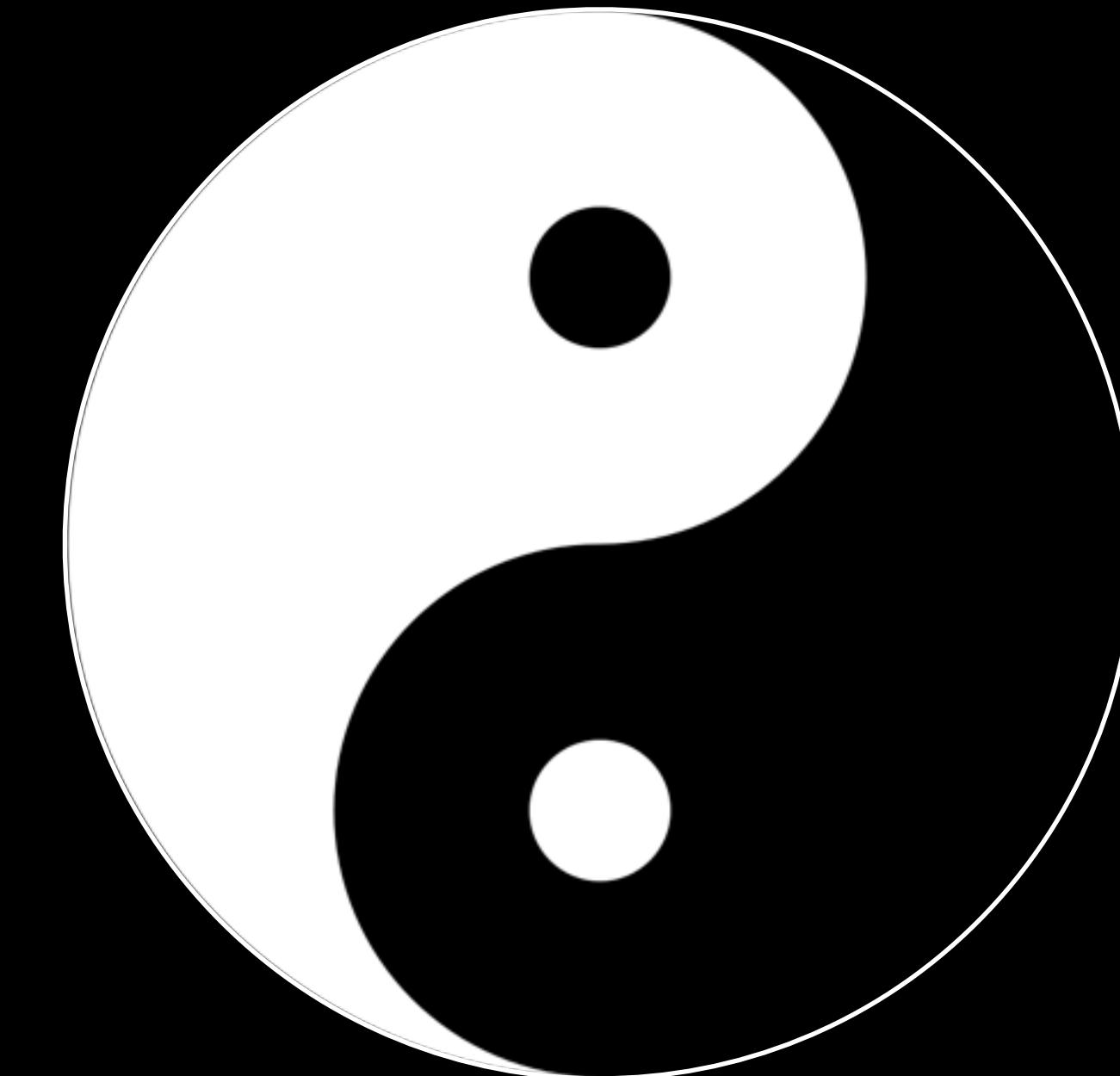
Name	Plane	Decays	m_{12}^2
IA	$m_A > m_H = m_{H^\pm}$	$A \rightarrow H^\pm W^\mp$	$m_H^2 s_\beta c_\beta$
		$A \rightarrow HZ$	0
IB	$m_A < m_H = m_{H^\pm}$	$H \rightarrow AZ/AA$	0
		$H^\pm \rightarrow AW^\pm$	
IIA	$m_H > m_A = m_{H^\pm}$	$H \rightarrow AZ/AA$	0
		$H \rightarrow H^+ H^- / H^\pm W^\pm$	
IIB	$m_H < m_A = m_{H^\pm}$	$A \rightarrow HZ$	$m_H^2 s_\beta c_\beta$
		$H^\pm \rightarrow HW^\pm$	0
III	$m_A = m_H = m_{H^\pm}$ vs $c_{\beta-\alpha}$	$A \rightarrow hZ, H^\pm \rightarrow hW^\pm$	$m_H^2 s_\beta c_\beta$
		$H \rightarrow hh$	0

Table 5.2: Summary of benchmark planes for exploring exotic Higgs decays in 2HDMs that survive theoretical and experimental constraints. Source: [4].

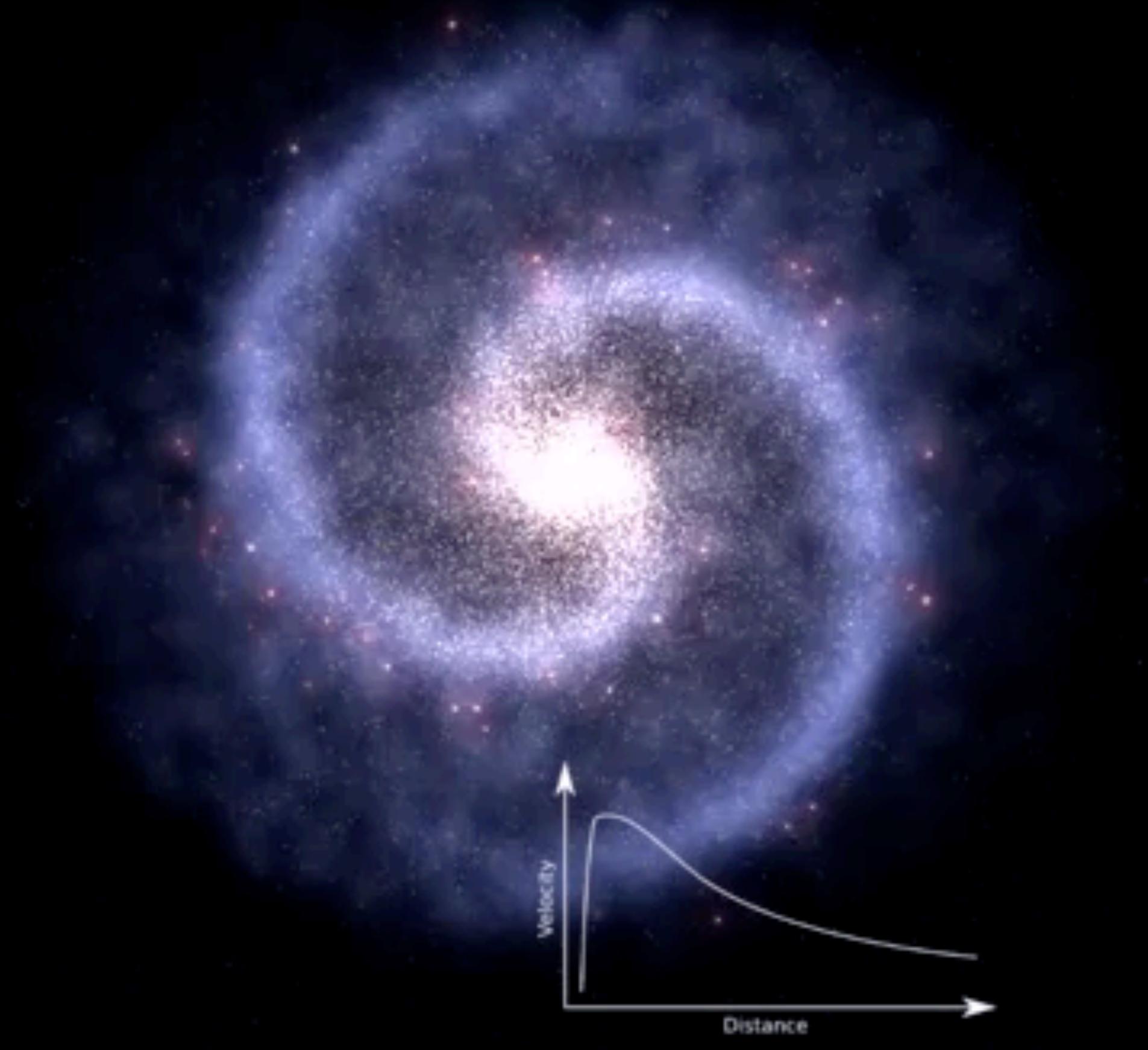
F Kling, J M No, and S Su.
“Anatomy of Exotic Higgs Decays in 2HDM”.
arXiv: 1604.01406

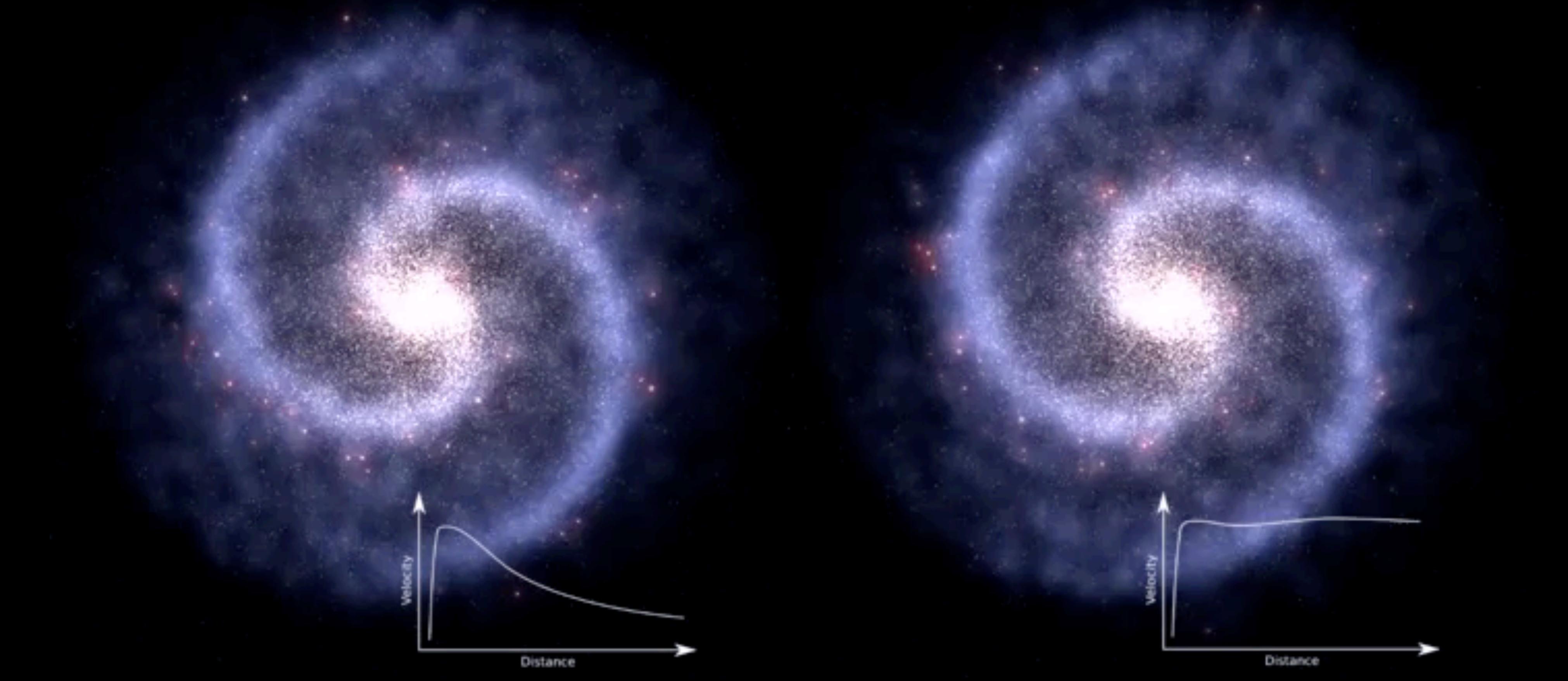
Conclusion

Conventional
decay channels



Exotic decay
channels





Without dark matter

With dark matter
(edges rotate faster)

Dark Matter = Weakly Interacting
Massive Particles (WIMPs)?

WIMPs from Supersymmetry (SUSY)

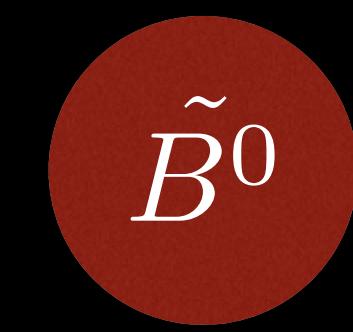


Supersymmetry transformation

A horizontal double-headed arrow pointing from left to right, positioned below the text "Supersymmetry transformation".



Neutralino Dark Matter



Bino

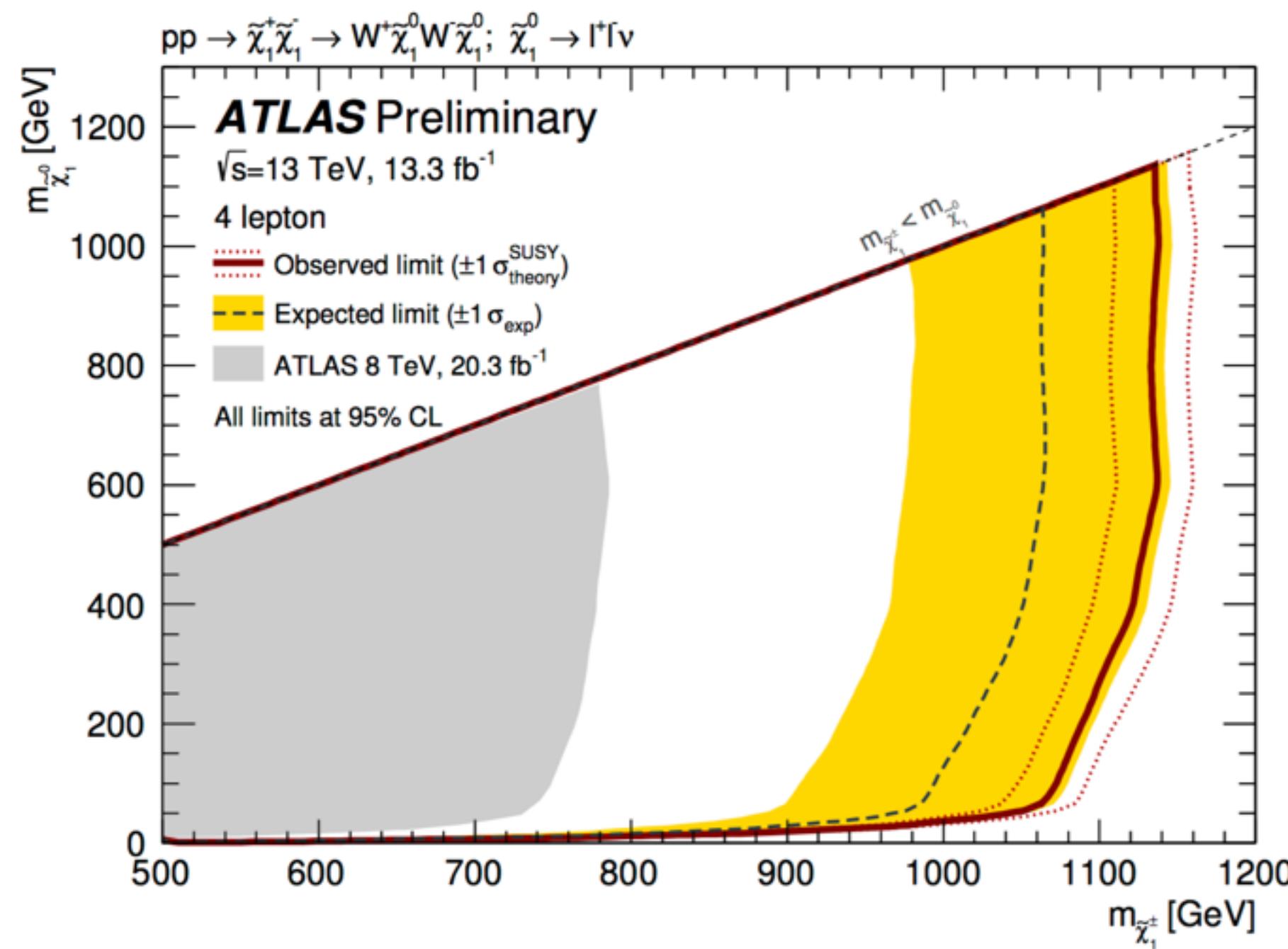


Neutral Higgsinos



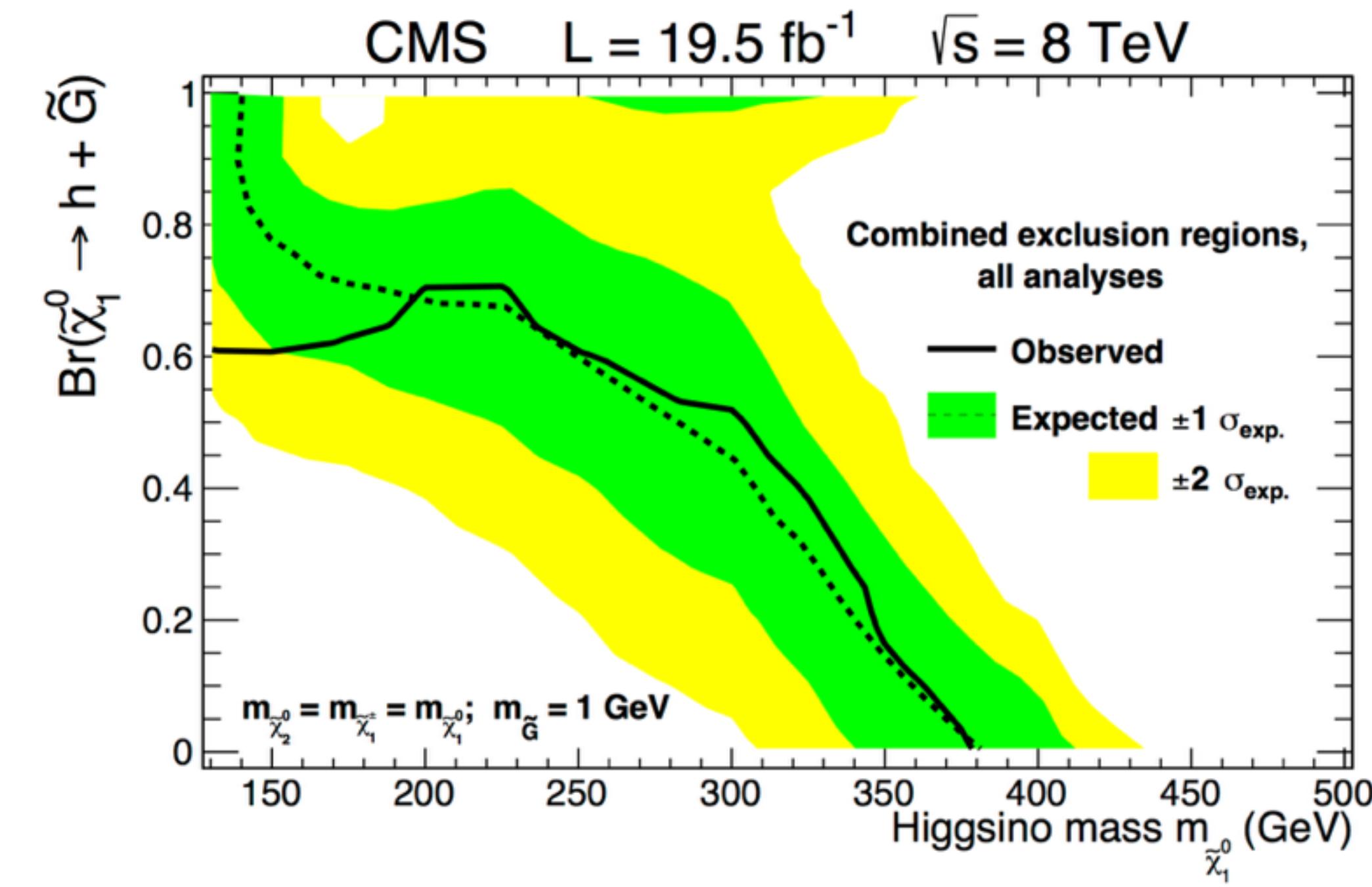
Neutral Wino

Limits on winos



ATLAS-CONF-2016-075

Limits on higgsinos

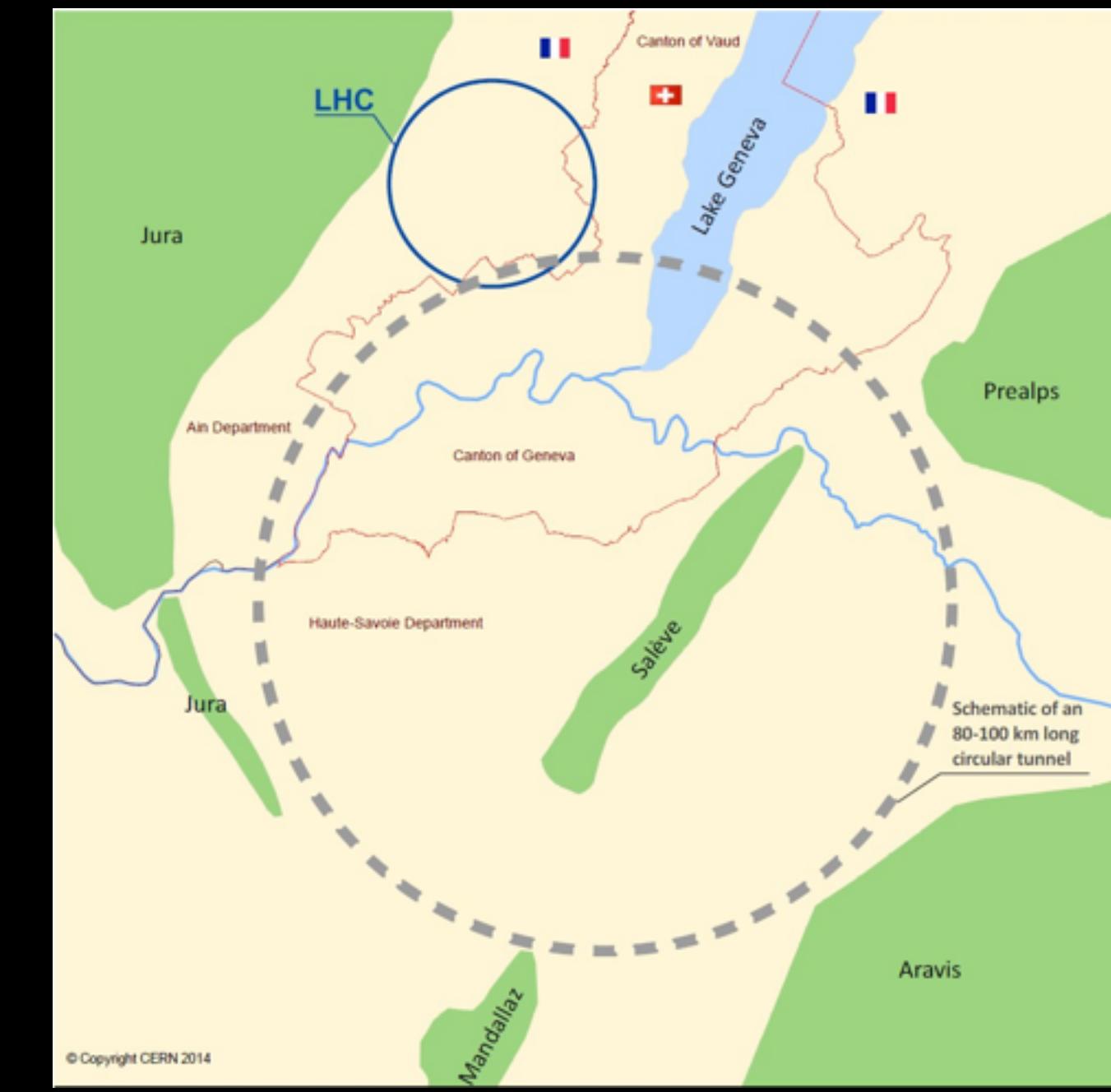


CMS-SUS-14-002

Future Colliders

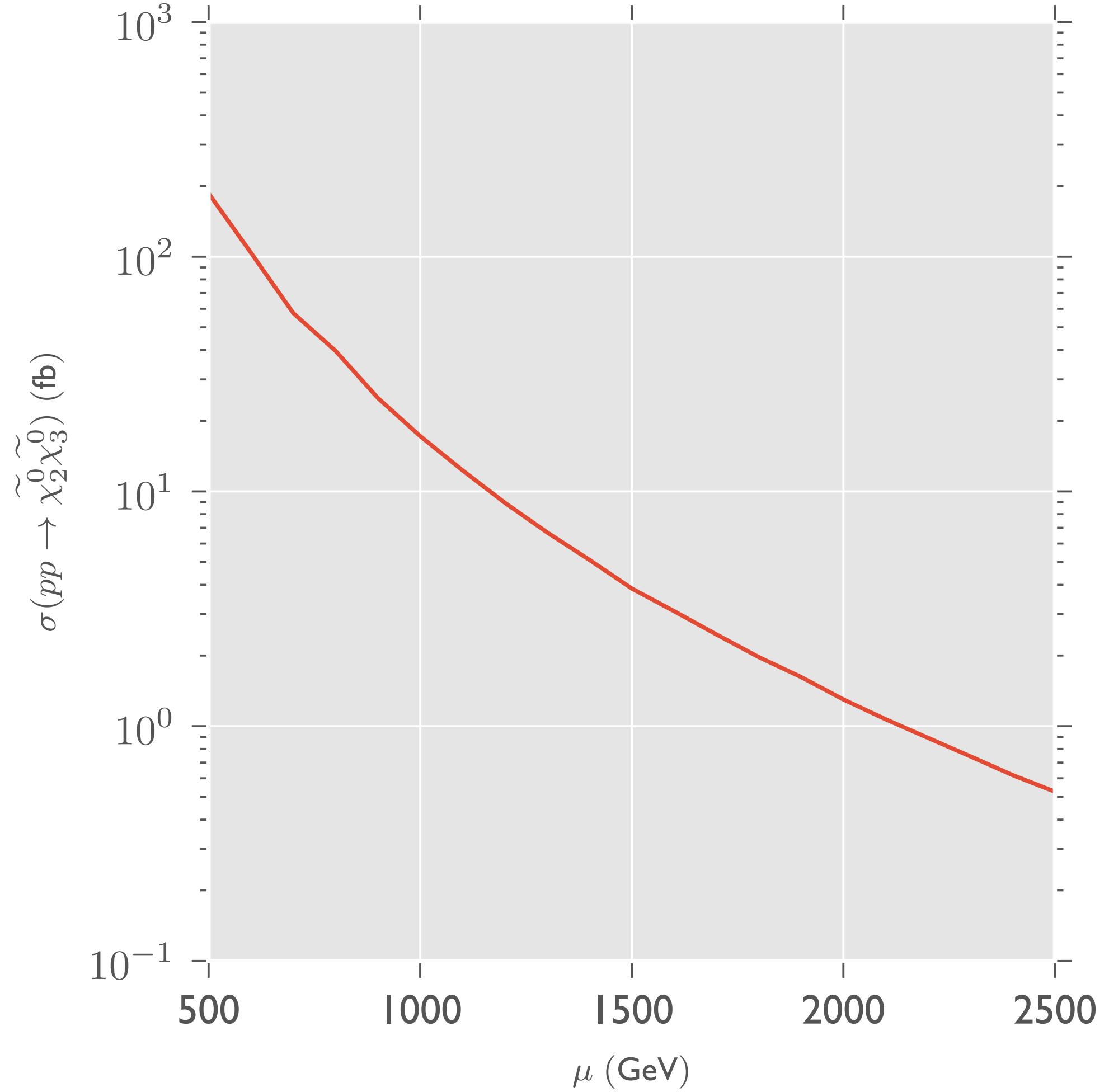


CEPC-SppC (China)

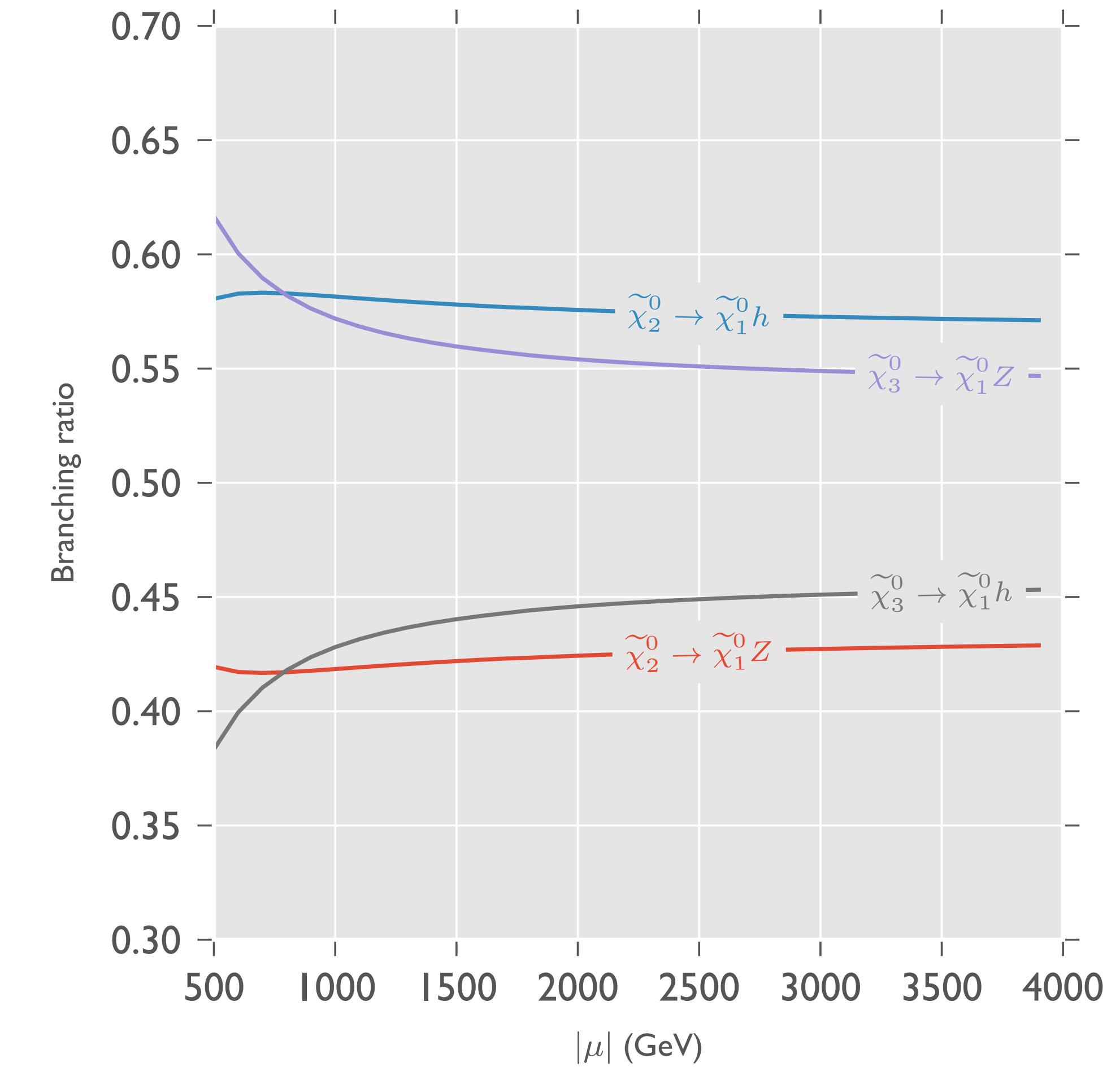


FCC-hh (CERN)

Higgsino pair production cross-section @ 100 TeV

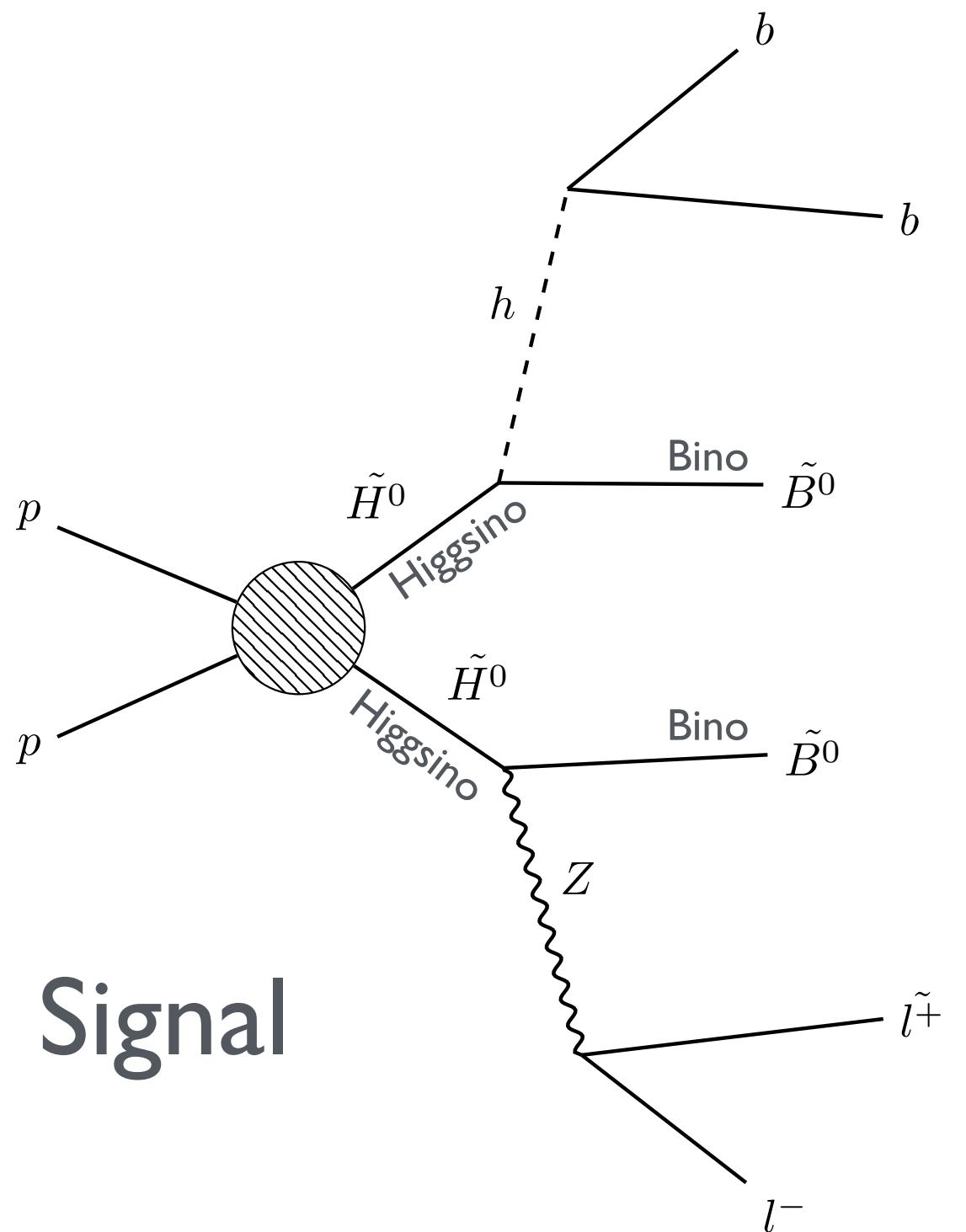


Higgsino branching ratios



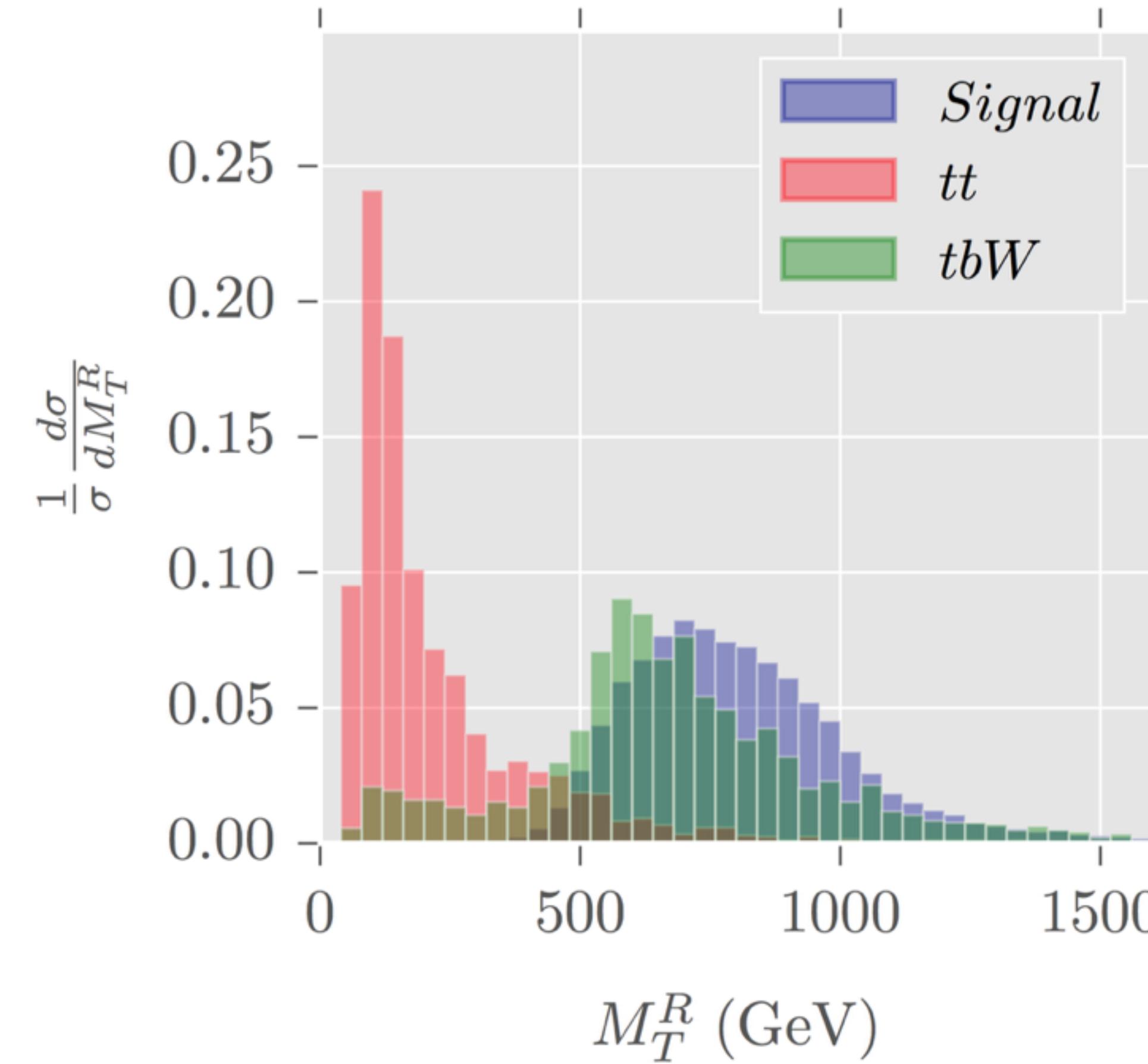
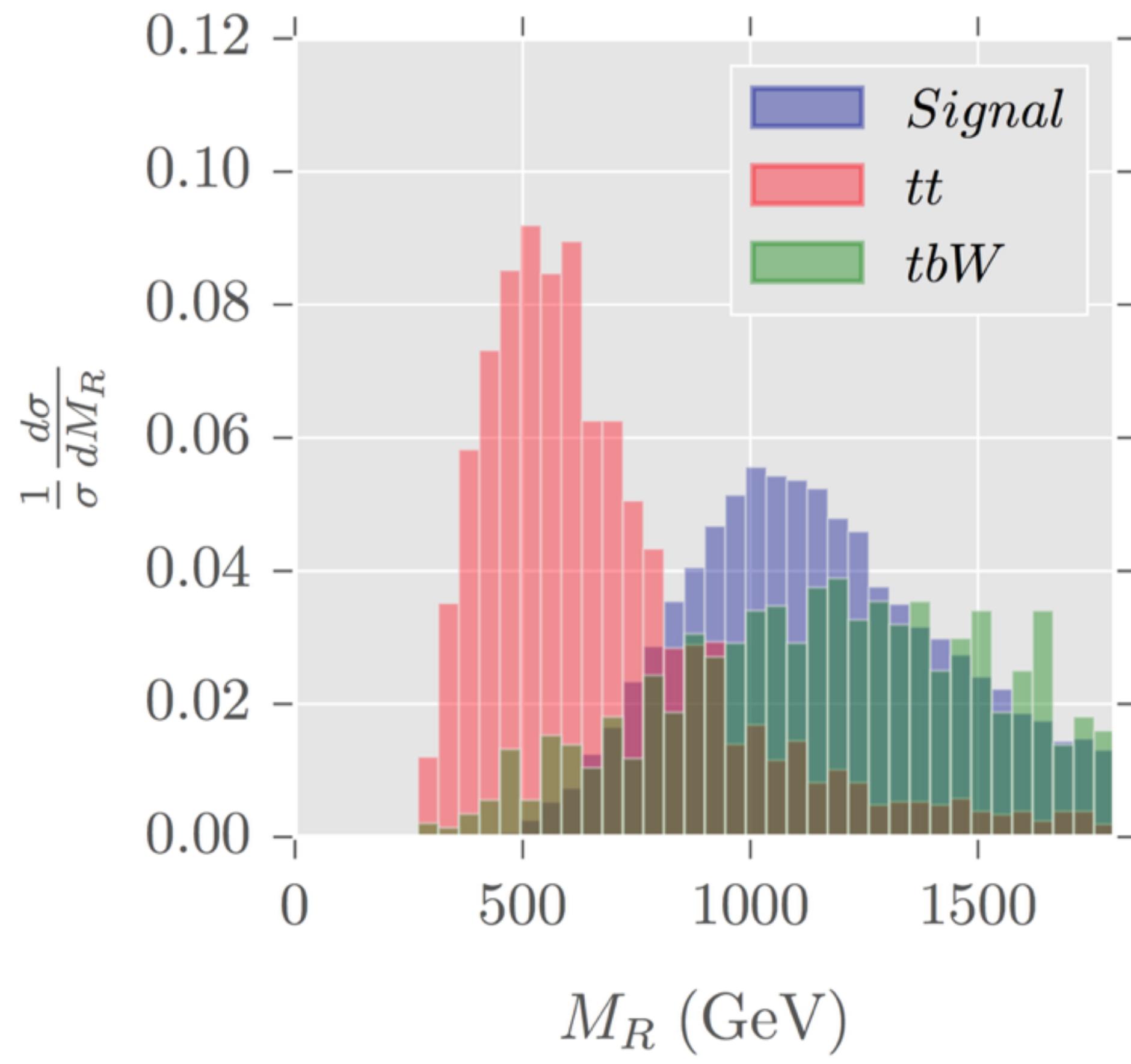
Bino mass = 25 GeV

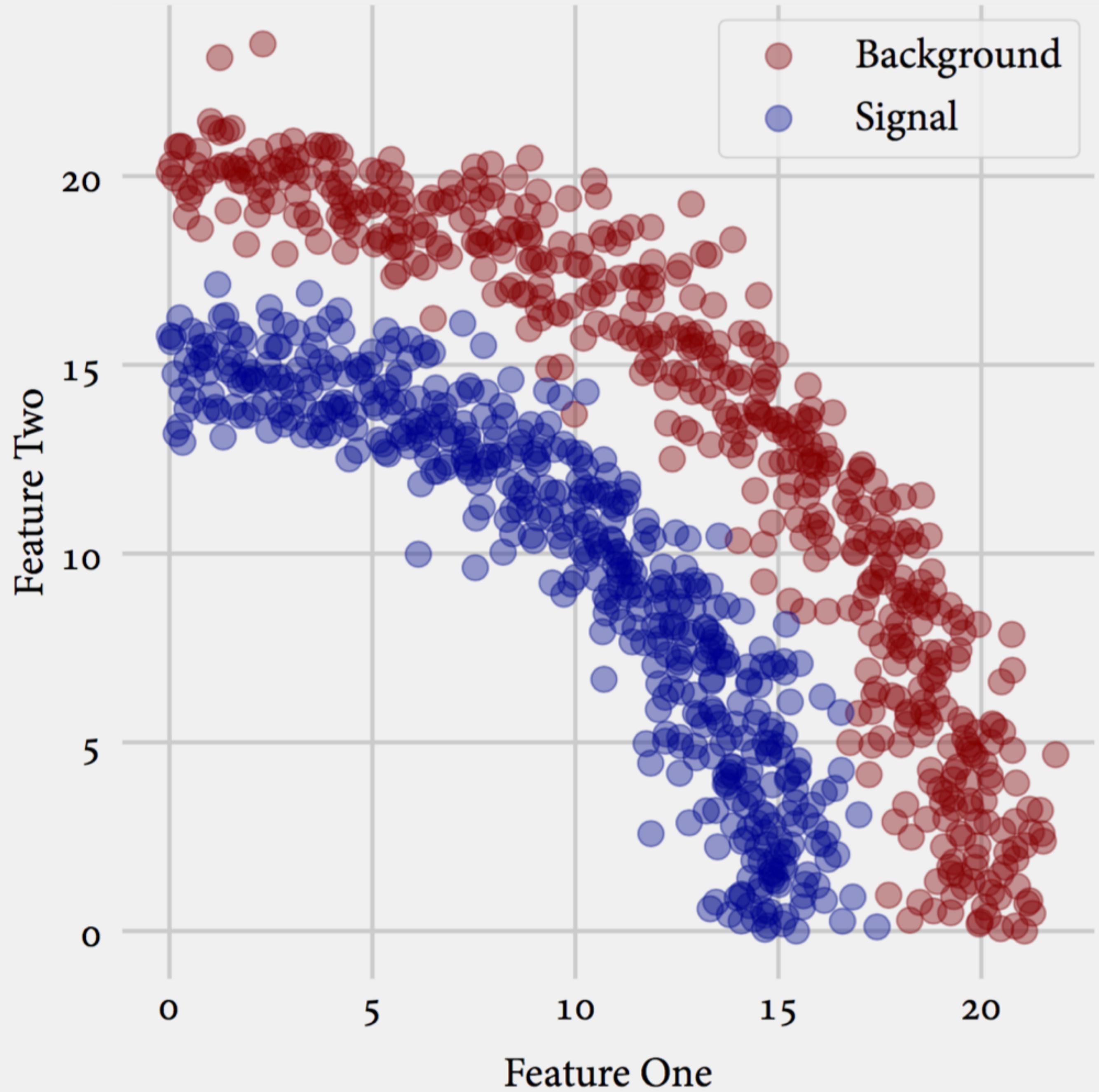
Our search channel



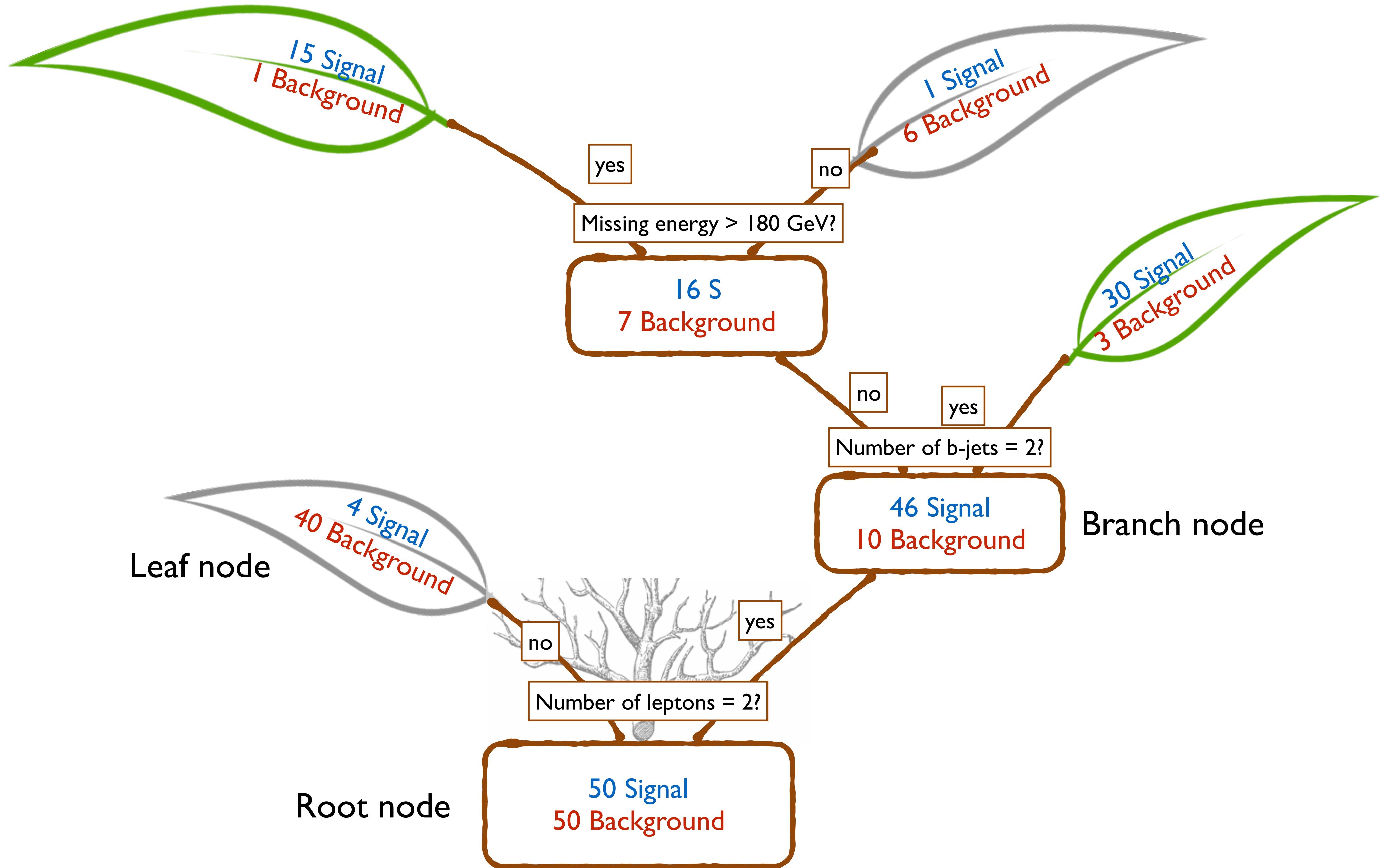
- Final state:
 - ◆ 2 b jets,
 - ◆ two same flavor, opposite charge leptons
 - ◆ missing energy
- Backgrounds:
 - * $t\bar{t}$
 - * $t b W$
 - * $b b W W$

Razor variable kinematic distributions

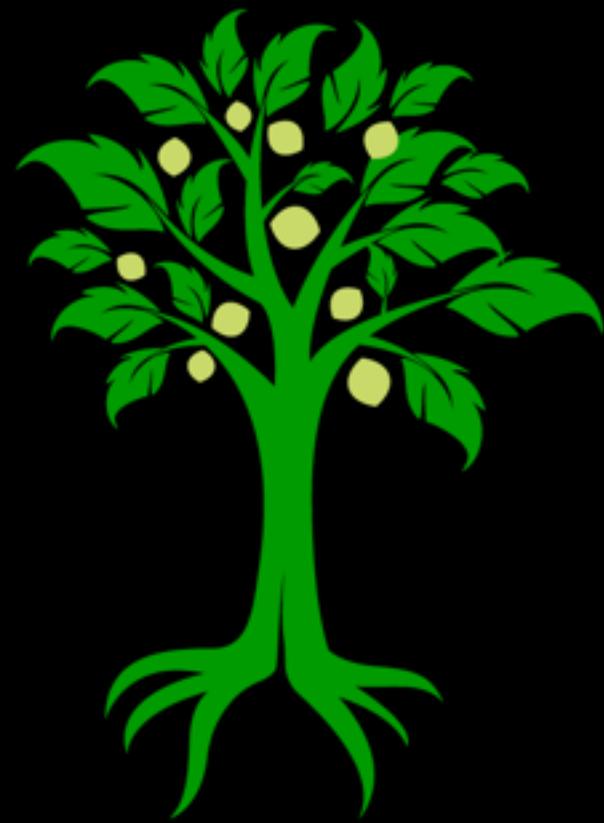


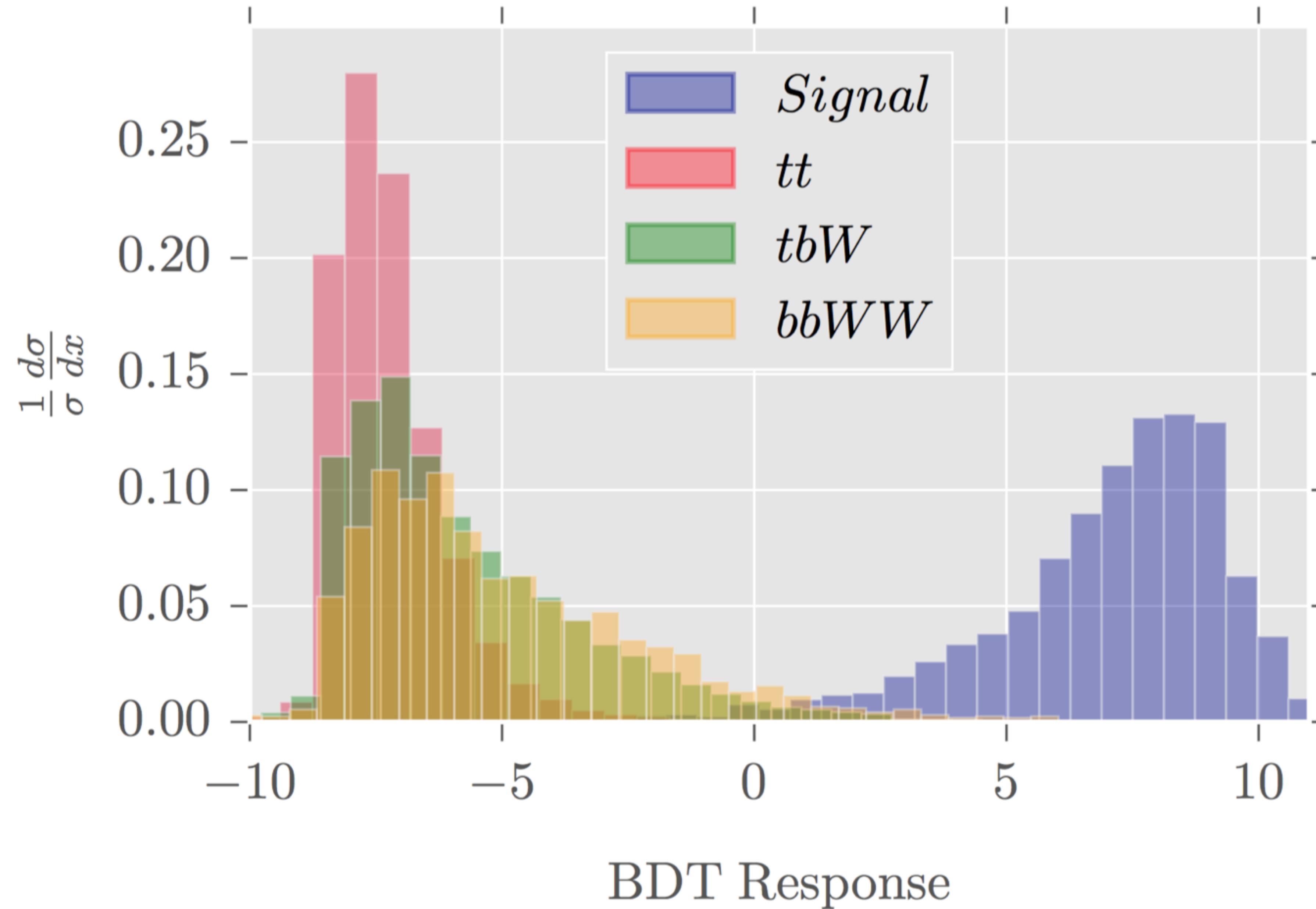


Example of nonlinear decision boundary in 2D feature space

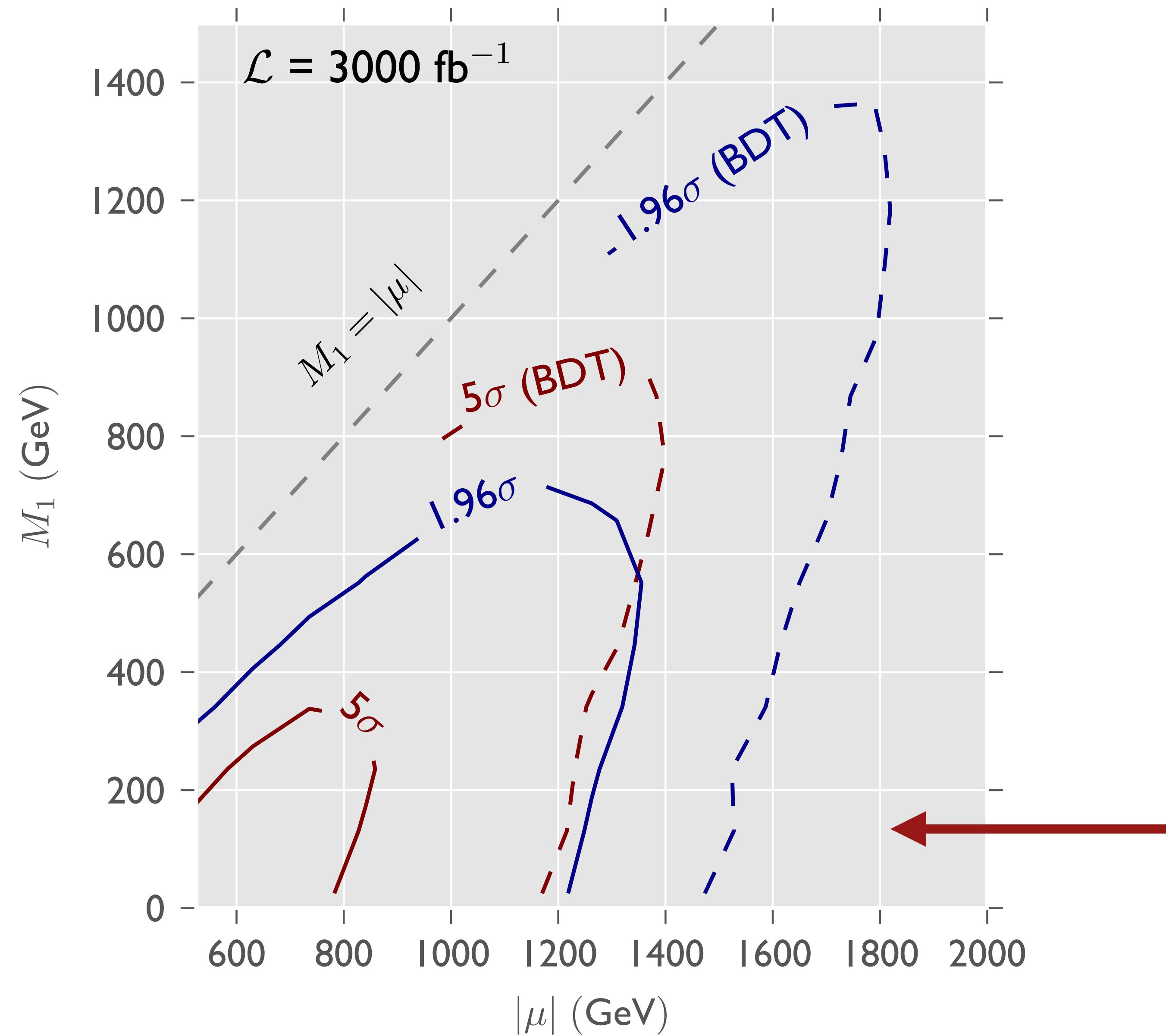


Final Score





Reach in parameter space



Collimated b-jets reduce
reach in this area

Conclusions

- Exotic Higgs decays are a good complementary avenue for discovery at current and future colliders
- A 100 TeV collider will open up many new physics opportunities, including finding a dark matter candidate particle.
- Machine learning will likely play an increasing role in particle physics in the future.

Thank you!
(Questions?)

The Hierarchy Problem

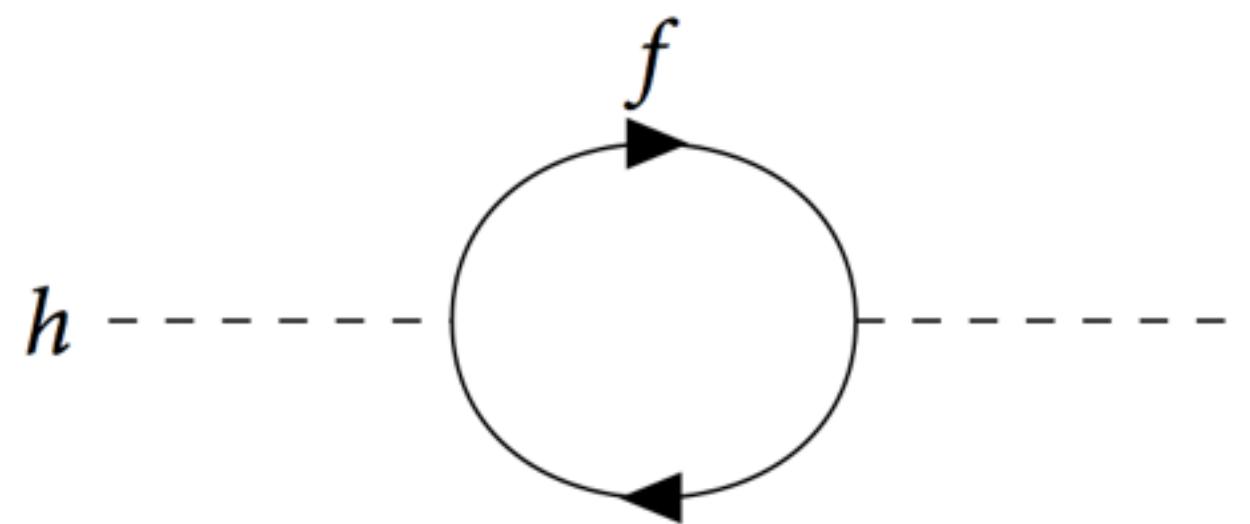


Figure 2.2: Feynman diagram for the one-loop fermionic correction to the SM Higgs mass

$$\Delta m_H^2 = -\frac{|\lambda_f|^2}{8\pi^2} \Lambda_{\text{UV}}^2 + \dots$$

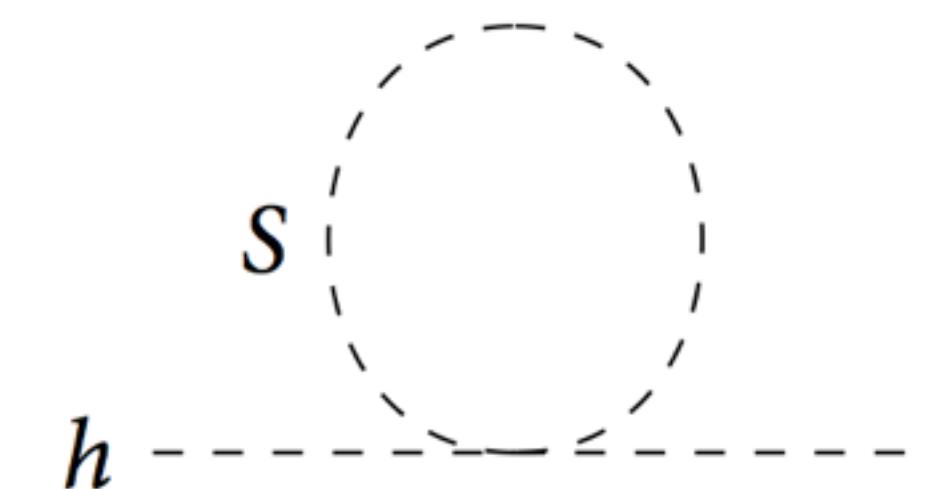
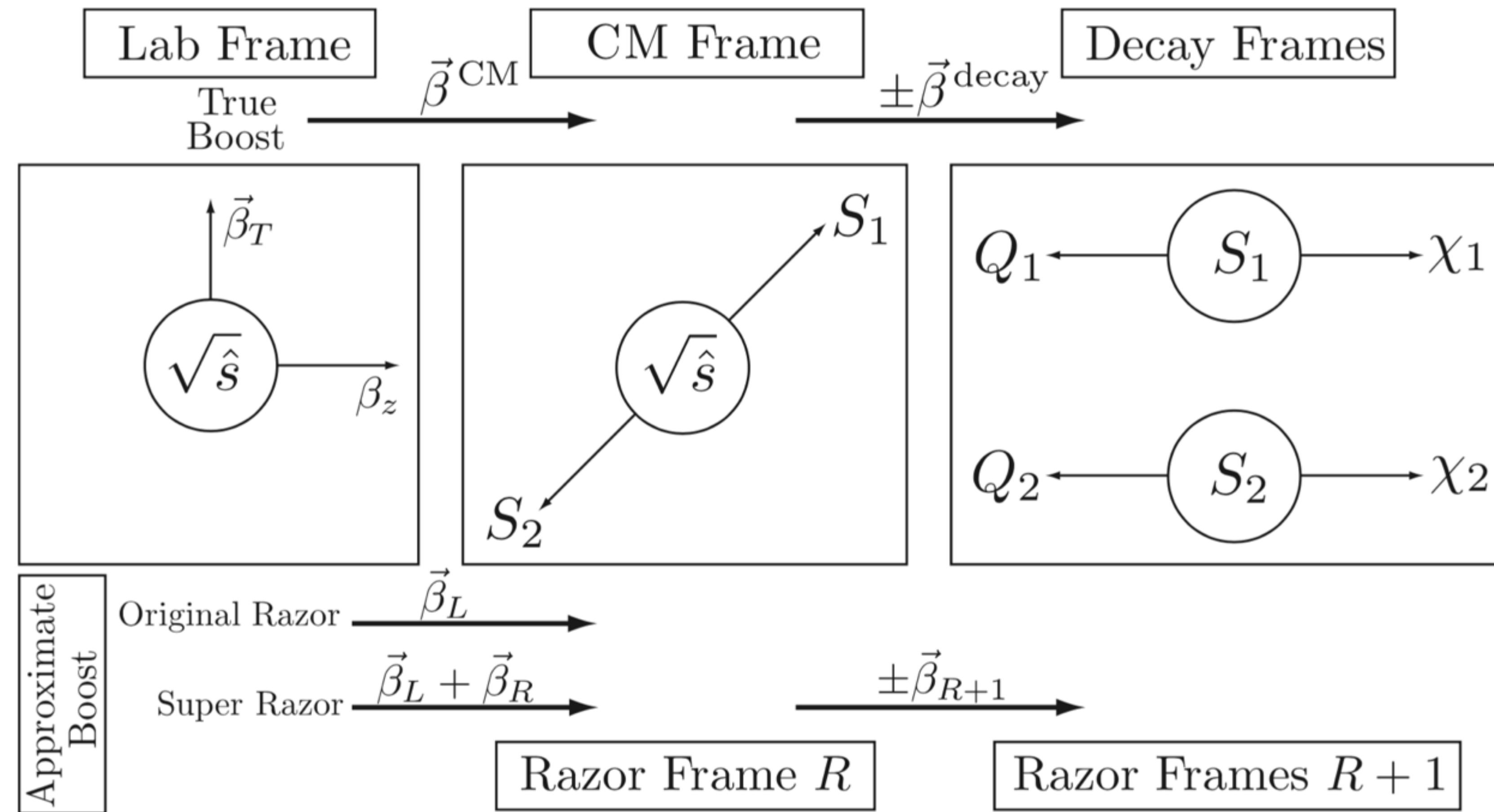


Figure 2.3: Feynman diagram for the one-loop scalar correction to the SM Higgs mass

$$\Delta m_H^2 = \frac{\lambda_s}{16\pi^2} [\Lambda_{\text{UV}}^2 + \dots]$$

Razor variables

arXiv:1006.2727



Razor variables

arXiv:1006.2727

$$M_R = \sqrt{(E_1 + E_2)^2 - (q_1^z + q_2^z)^2}$$

$$M_T^R = \sqrt{\frac{1}{2} \left[\cancel{E}_T (|\vec{q}_{1T}| + |\vec{q}_{2T}|)^2 - \cancel{E}_T \cdot (\vec{q}_{1T} + \vec{q}_{2T}) \right]}$$