

M392C NOTES: INDEX THEORY

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These notes were taken in UT Austin's M392C (Index theory) class in Spring 2018, taught by Dan Freed. I live-TeXed them using vim, so there may be typos; please send questions, comments, complaints, and corrections to a.debray@math.utexas.edu. Any mistakes in the notes are my own. Thanks to Rok Gregoric for fixing a few errors, and to Riccardo Pedrotti for providing the notes for §§3 and 20.

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Lecture 1.

Overview, History, and some Linear Algebra: 1/17/18

"This formula should look fake if you haven't seen it before."

We'll start with an overview and some history of index theory. The overview will use a little bit of complex geometry, but if you don't know it that's okay; the rest of the class will not depend on it.

One of the earliest manifestations of index theory was in the theory of algebraic curves. Let M be a compact smooth connected complex curve, i.e. a Riemann surface, and let D be a divisor on M , a finite formal sum of points of M with integer coefficients. For example, if $p_1, p_2, p_3 \in M$, one divisor is $4p_1 - 2p_2 + 7p_3$.

Definition 1.1. Let f be a meromorphic function on M ; then, its divisor $\text{div}(f)$ is the zeros of f minus the poles of f , where both are counted with multiplicity. For $f = 0$, we let $\text{div}(0) = 0$.

For example, if $M = \mathbb{CP}^1 = \mathbb{C} \cup \{\infty\}$, then a meromorphic function on M is a rational function. If we took $f(z) = (z-1)^2/(z+2)$, then $\text{div}(f) = 2 \cdot 1 - 1 \cdot (-2) - 1 \cdot \infty$: f has a double zero at 1 and a single pole at -2 , and at ∞ there is a simple pole.¹

A divisor has a *degree* which is the sum of its terms.

Theorem 1.2. *The degree of the divisor of a meromorphic function is zero.*

This is a consequence of the Cauchy integral formula.

A divisor specifies the zeros and poles of a meromorphic function, and it's a classical problem to, given a degree-zero divisor D on a Riemann surface, construct a function whose divisor is D . More generally, let $\mathcal{L}(D)$ denote the set of meromorphic f such that $\text{div}(f) + D \geq 0$.² $\mathcal{L}(D)$ is a vector space, and if $\deg(D) < 0$, $\mathcal{L}(D) = 0$; we also have $\mathcal{L}(0) = \mathbb{C}$, given by constant functions.

Another classical question is to compute $\dim \mathcal{L}(D)$. Riemann provided an estimate:

$$(1.3) \quad \dim \mathcal{L}(D) \geq 1 - g + \deg(D),$$

where g is the *genus* of M , defined to be

$$(1.4) \quad g := \frac{1}{2} \text{rank} H_1(X).$$

The next natural question is to identify the discrepancy, and Riemann's student Roch found the answer.

Theorem 1.5 (Riemann-Roch). *here is a canonical divisor K_M such that*

$$(1.6) \quad \dim \mathcal{L}(D) - \dim \mathcal{L}(K_M - D) = 1 - g + \deg D.$$

We won't say much about K_M , though $\deg(K_M) = 2g - 2$.

Corollary 1.7. *The genus is an integer.*

A more modern interpretation of this story is that D determines a holomorphic line bundle $L \rightarrow M$, and $\mathcal{L}(D)$ is the vector space of holomorphic sections of L , i.e. $\mathcal{L}(D) \cong H^0(M; L)$. If s is any smooth section of L , s is holomorphic iff $\bar{\partial}s = 0$. That is, in local coordinates $z = x + iy$, and

$$(1.8) \quad \bar{\partial} = \frac{\partial}{\partial \bar{z}} = \frac{1}{2} \left(\frac{\partial}{\partial x} - i \frac{\partial}{\partial y} \right).$$

Thus, $\bar{\partial}s = 0$ is a first-order differential equation, and computing $\dim \mathcal{L}(D)$ is asking for the dimension of the space of solutions to the equation. Thus one way you might prove Theorem 1.5 is to analyze the differential operator $\bar{\partial}$, which is a linear operator

$$\bar{\partial} : \Omega^{0,0}(M; L) \longrightarrow \Omega^{0,1}(M; L).$$

Then, $\mathcal{L}(D) = \ker(\bar{\partial})$ and $\mathcal{L}(K_M - D) \cong \text{coker}(\bar{\partial})$.

Definition 1.9. The *index* of $\bar{\partial}$ is $\text{ind}(\bar{\partial}) := \dim \ker(\bar{\partial}) - \dim \text{coker}(\bar{\partial})$.

Broadly speaking, this course will be about indices of this sort, and their applications: for example, the Riemann-Roch theorem from this perspective is about computing the index of $\bar{\partial}$.

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For a simpler case, let V and W be finite-dimensional vector space and $T : V \rightarrow W$ be a linear map. Then, $\ker(T) \subset V$ and $\text{coker } T := W/T(V)$. Computing the index is a fundamental theorem in linear algebra.

Theorem 1.10.

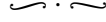
$$\text{ind}(T) := \dim(\ker T) - \dim(\text{coker } T) = \dim V - \dim W.$$

In particular, it's independent of T ! One way you might prove this is to observe that it's true when $T = 0$ and then try to prove that it's locally constant.

In this class, we're interested in operators between infinite-dimensional vector spaces, such as $\Omega^{p,q}(M; L)$, whose kernels and cokernels are finite-dimensional (such that the definition of an index makes sense). There will be no nice formula like Theorem 1.10, but some aspects stay the same: though the dimension of the kernel or cokernel may jump along a continuous path, their difference is constant.

¹To see this, use the change-of-variables $z = 1/w$ and evaluate f at $w = 0$.

²This is missing a zero element, so one needs to adjoin 0 for everything to work.



Another classical subject that relates to index theory is that of the Euler number of a compact smooth n -manifold M . Betti defined *Betti numbers* b_0, \dots, b_n associated to M , and Noether realized they can be identified with ranks of abelian groups (or dimensions of certain real vector spaces).³

Definition 1.11. The *Euler characteristic* of M is

$$\chi(M) := \sum_{i=0}^n (-1)^i b_i.$$

The Betti numbers are defined via simplices, and how M is built out of cells. Since M is a smooth manifold, one might want to compute them in another way, using the smooth structure of the manifold. To do this, one introduces the *de Rham complex*

$$(1.12) \quad 0 \longrightarrow \Omega^0(M) \xrightarrow{d} \Omega^1(M) \xrightarrow{d} \Omega^2(M) \longrightarrow \cdots \longrightarrow \Omega^n(M) \longrightarrow 0,$$

with linear maps d such that $d^2 = 0$. Unlike in the previous example, this is built out of real functions and real differential forms.

Definition 1.13. The *de Rham cohomology* of M is the sequence of real vector spaces

$$H_{\text{dR}}^i(M) := \frac{\ker(d: \Omega^i(M) \rightarrow \Omega^{i+1}(M))}{\text{Im}(d: \Omega^{i-1}(M) \rightarrow \Omega^i(M))}.$$

Theorem 1.14 (de Rham). *There is an isomorphism $H_{\text{dR}}^i(M) \cong H^i(M; \mathbb{R})$, and therefore $\dim H_{\text{dR}}^i(M) = b_i$.*

From this perspective, the Euler characteristic looks more like an index, where we stack together the pieces of the de Rham complex:

$$(1.15) \quad \bigoplus_{i \text{ even}} \Omega^i(M) \longrightarrow \bigoplus_{i \text{ odd}} \Omega^i(M).$$

However, the index of this is *not* the Euler characteristic! The issue is that the de Rham cohomology groups are a subquotient, not just a subspace or just a quotient. To compute the Euler characteristic as an index, we'll need some way of turning them into pure subspaces or quotients. One way to do this is to use an inner product and take orthogonal complements.

Let M be a Riemannian manifold. Then, there is a Laplace operator $\Delta: \Omega^i(M) \rightarrow \Omega^i(M)$, which is a linear second-order elliptic differential operator.

Remark 1.16. There are three basic kinds of differential operators studied in a typical differential equations course: elliptic, parabolic, and hyperbolic. The Laplacian is the basic example of an elliptic operator; the heat operator is the basic example of a parabolic operator; and the Schrödinger operator is the basic example of a hyperbolic operator. We will focus on elliptic operators in this course, but both the heat equation and the Schrödinger equation will appear. ◀

Example 1.17. Let \mathbb{E}^n denote n -dimensional Euclidean space with coordinates x^1, \dots, x^n . Then, the Laplacian on \mathbb{E}^n is

$$\Delta = \left(\frac{\partial}{\partial x^1} \right)^2 + \cdots + \left(\frac{\partial}{\partial x^n} \right)^2.$$

◀

For more general Riemannian manifolds, the definition of the Laplacian is more complicated, but not much more so.

Definition 1.18. If M is a Riemannian manifold, there is an L^2 inner product on $\Omega^i(M)$ defined by

$$\langle \alpha, \beta \rangle_{L^2} := \int_M \langle \alpha(M), \beta(M) \rangle \, \text{dvol}_m.$$

Using these inner products, we can let $d^*: \Omega^{i+1}(M) \rightarrow \Omega^i(M)$ be the formal adjoint to d .

³These days, this would be called *categorification*: it can often be useful to identify a number as the dimension of some vector space attached to your object.

Fact. d^* exists and is a first-order differential operator. ◀

Definition 1.19. The Laplace operator on M is $\Delta := dd^* + d^*d$.

A form in the kernel of Δ is called *harmonic*, and the space of harmonic forms is denoted $\mathcal{H}^i(M) \subset \Delta^i(M)$.

Theorem 1.20 (Hodge theorem). *The natural map $\mathcal{H}^i(M) \rightarrow H_{\text{dR}}^i(M)$ is an isomorphism. In particular, $\dim \mathcal{H}^i(M) = b_i$.*

This is how index theory enters the picture: if we can access the space of harmonic forms as kernels and cokernels of operators, we could compute the Euler characteristic as an index. And indeed, we can fix (1.15) as follows:

$$(1.21) \quad \bigoplus_{i \text{ even}} \Omega^i(M) \xrightarrow{d+d^*} \bigoplus_{i \text{ odd}} \Omega^i(M).$$

The index of this operator is the Euler characteristic.

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A third example of index theory is the higher-dimensional Riemann-Roch theorem. Let M be a compact complex manifold; then, the $\bar{\partial}$ operator defines a *Dolbeault complex* analogous to the de Rham complex. If M is 2-(complex-)dimensional, the Euler characteristic satisfies a formula

$$(1.22) \quad \chi(M) = \frac{1}{12} (c_1^2(M) + c_2(M)) [M].$$

Here c_1 and c_2 are examples of *characteristic classes*, which we'll start on in the next few lectures. In particular, the right-hand side is an integer. In higher dimensions, there are similar expressions with larger denominators and more characteristic classes.

These were studied by Todd and his student Egger, by Weyl, and others. But the general forms remained conjectures until 1954, when Hirzebruch proved these generalizations of the Riemann-Roch theorem, and an additional, similar result called the signature theorem. He wove together two very new pieces of mathematics: the cobordism theory of René Thom (published only earlier that year!) and the theory of sheaves.

Hirzebruch and others in this field introduced a rational combination of different characteristic numbers, called *Pontrjagin numbers*, called the \hat{A} -genus (said "A-hat genus"). This is defined on closed oriented manifolds, and on a spin manifold is an integer.

That $\hat{A}(M)$ is an integer is a suggestion that it's a dimension of something, and when Singer went to visit Oxford in 1963, Atiyah asked him what object has the \hat{A} -genus as its dimension, and this is the problem that they solved: they constructed a differential operator called the Dirac operator on a spin manifold, and showed that its index is the \hat{A} -genus.

The Dirac operator

$$D := \gamma^\mu \frac{\partial}{\partial x^\mu}$$

for some γ^μ (this notation means the index μ is implicitly summed over) is a first-order linear differential operator. We'd like this to be a square root of the Laplacian operator.

Exercise 1.23. Show that $D^2 = \Delta$ iff

$$(1.24) \quad \gamma^\mu \gamma^\nu + \gamma^\nu \gamma^\mu = -2\delta^{\mu\nu}.$$

Here, $\delta^{\mu\nu}$ means 1 if $\mu = \nu$ and 0 otherwise.

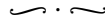
So the operator has to satisfy n^2 equations. If you try to solve this for functions on \mathbb{E}^n , you can show that no such γ^μ exist, but one could instead ask for vector-valued functions which satisfy (1.24), and indeed we will spend some time studying the abstract theory of matrices which satisfy this condition, rephrased as the algebraic theory of Clifford modules. In particular, we will be able to show that a spin structure is precisely what one needs to be able to construct the Dirac operator on a Riemannian manifold.

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Before Atiyah and Singer told this story, Grothendieck took the Hirzebruch-Riemann-Roch theorem and generalized it still further, and Atiyah and Hirzebruch saw how to translate his ideas from algebraic geometry to topology, and replace sheaves with vector bundles. They then defined K -theory and rapidly developed it from 1958 to 1962. When Atiyah asked Singer his question, it was in this context.

At the same time, parallel work was undertaken in the Soviet Union under Gelfand and his students. He observed that the index sometimes can be computed topologically, and asked whether this is true in general, and Atiyah-Singer's answer also incorporates this question.

Subsequently, in the 1970s, Gilkey, Patodi, and others were able to provide more rigid, simpler proofs with analytic methods, and in the 1980s Getzler made another important simplifying step to what's now called the heat equation proof of the index theorem, which we'll follow.

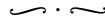


We'll use John Roe's book in this course. It's analytic in flavor, but also treats many other nice results, and if we go quickly enough, we'll get to see some of them, including Witten's physical treatment of Morse theory, the Lefschetz theorem, the Hodge theorem, and more.

In this class, the students will give lectures, two each week, and we hope to go through two chapters a week. You don't have to use all three hours!

On the course website (<https://www.ma.utexas.edu/users/dafr/M392C/>), there will be some useful information, including some old course notes, some historical background, and more to come. These will be there so that you do not forget the beauty of the material amongst all the details in the lectures.

Not everybody may know all of the prerequisites for this course, since it draws in lots of different parts of mathematics. One can ask the professor for references or talk to other students in the course.



The second half of the first day is on the first chapter of the book, reviewing some of the basics of Riemannian geometry.

Let's first start with some linear algebra and differential forms. Let V be an n -dimensional real vector space. Eventually, V will be a tangent space at a point to a manifold, and if the manifold has a Riemannian metric, V picks up an inner product.

Associated to V are several canonical vector spaces built from it: its wedge powers $\Lambda^2 V, \dots, \Lambda^n V$, and $\Lambda^0 V$, which is canonically \mathbb{R} . The top exterior power is also called the *determinant line*, $\text{Det } V := \Lambda^n V$. Dually, there are the exterior powers of the dual space V^* : $\mathbb{R}, V^*, \Lambda^2 V^*, \dots, \text{Det } V^*$.

An inner product on V canonically induces inner products on all of these exterior powers. One way to see this is to let e_1, \dots, e_n be an orthonormal basis of V ; then, there is a dual basis e^1, \dots, e^n of V^* , defined by the relation

$$(1.25) \quad e^\mu(e_\nu) = \delta^\mu_\nu,$$

i.e. 1 if $\mu = \nu$ and 0 otherwise.

We specify the inner product on V^* by declaring this dual basis orthonormal, which suffices, though you have to check that if you change the orthonormal basis of V you started with, you'll end up with the same inner product nonetheless.

We also obtain bases for the exterior powers of V and V^* : for $\Lambda^q V$, the basis is

$$(1.26) \quad \{e_{i_1} \wedge \dots \wedge e_{i_q} : 1 \leq i_1 < \dots < i_q \leq n\},$$

and for $\Lambda^q V^*$, it's

$$(1.27) \quad \{e^{i_1} \wedge \dots \wedge e^{i_q} : 1 \leq i_1 < \dots < i_q \leq n\}.$$

Again we define the inner products on $\Lambda^q V$ and $\Lambda^q V^*$ by asking for these bases to be orthonormal, and again the inner product in question does not depend on the specific choice of orthonormal basis of V .

Definition 1.28. An *orientation* of V is an orientation of its determinant line. That is, $\text{Det } V \setminus 0$ has two components, and an orientation is a choice of one of them.

Given n vectors $e_1, \dots, e_n \in V$, we can wedge them together to an $e_1 \wedge \dots \wedge e_n \in \text{Det } V$; $\{e_1, \dots, e_n\}$ is a basis iff $e_1 \wedge \dots \wedge e_n \neq 0$. Thus a basis singles out one of the two rays in $\text{Det } V \setminus 0$, hence defines an orientation. Since $(\text{Det } V)^* = \text{Det}(V^*)$ canonically, then this also defines an orientation on $(\text{Det } V)^*$: the duality pairing implies there's a single $\theta \in \text{Det } V^*$ which sends $e_1 \wedge \dots \wedge e_n \mapsto 1$; we call it the *volume form* and denote it vol .

On an oriented Riemannian n -manifold, this is a differential n -form, hence can integrate it to determine the volume of the manifold. If it's not oriented, there are two at each point, which may twist globally into something called a density. Nonetheless, this can be integrated, and the volume of, e.g. \mathbb{RP}^2 still makes sense.

The pairing $\Lambda^q V^* \otimes \Lambda^{n-q} V^* \rightarrow \text{Det } V^*$ defined by

$$(1.29) \quad \alpha, \beta \mapsto \alpha \wedge \beta$$

is nondegenerate. An orientation of V defines a trivialization of $\text{Det } V^*$ (where $\text{vol} = 1$), so this pairing is \mathbb{R} -valued. Therefore we obtain an isomorphism $\Lambda^q V^* \cong \Lambda^{n-q}(V)$, though it depends on the inner product and the orientation.

Example 1.30. In three dimensions, we use this frequently, to shift from the perspective of vector fields and scalars and div , ∇ , and curl to differential forms. \blacktriangleleft

There's also an isomorphism $\star: \Lambda^q V^* \rightarrow \Lambda^{n-q} V^*$ which only uses the inner product; this is called the *Hodge star*. Putting everything together, the Hodge star is defined uniquely by the stipulation that

$$(1.31) \quad \alpha_1 \wedge \star \alpha_2 = \langle \alpha_1, \alpha_2 \rangle \text{vol}$$

for any $\alpha_1, \alpha_2 \in \Lambda^q V^*$.

Exercise 1.32. For example, check that $\star(e_{i_1} \wedge \cdots \wedge e_{i_q})$ is the wedge of all of the e_j not in (i_1, \dots, i_q) , possibly multiplied by -1 .

Exercise 1.33. Show that $\star^2 = (-1)^{q(n-q)}$.

Here, “inner product” means nondegenerate inner product; much of this story still goes through for a Lorentz-signature metric, but not all of it.

Exercise 1.34. Show that on a closed, oriented Riemannian manifold M , $d^* = \pm \star d \star$, and determine the sign (which depends on n and q).

You can type-check that the right-hand side is a first-order differential operator which lowers the degree by 1. Solving the exercise boils down to checking that

$$\int_M \langle d\alpha, \beta \rangle \text{vol} = \pm \int_M \langle \alpha, d\star\beta \rangle \text{vol}.$$

You'll end up using Stokes' theorem.

$\smile \cdot \smile$

Now let's think about parallelism. Let \mathbb{A}^n be n -dimensional affine space (no distinguished origin), where we learn calculus. This has parallel transport: if $\xi \in \mathbb{R}^n$ is a tangent vector at some point, we can translate it everywhere to a vector field. This allows us to define differentiation: if $f: U \rightarrow \mathbb{R}$, where $U \subset \mathbb{A}^n$ is open, then we define the derivative of f at p in the direction of ξ to be

$$(1.35) \quad \xi_p f := \lim_{h \rightarrow 0} \frac{f(p + h\xi) - f(p)}{h}.$$

This uses parallelism in the expression $p + h\xi$.

More generally, if M is a smooth manifold, we don't always have a canonical parallel transport between tangent spaces for different points of the manifold, so we can't compare tangent vectors in different places and differentiate.

For example, if $\gamma: [a, b] \rightarrow M$ is a curve, its tangent vectors at two different points can't be compared (without extra structure), so there's no way to make the subtraction in (1.35). We'll introduce the structure that allows us to do this.

Definition 1.36. Let $V \rightarrow M$ be a vector bundle and $C^\infty(M; V)$ denote its space of smooth sections, which is a real vector space. A *covariant derivative* is a bilinear operator

$$\nabla: C^\infty(M; TM) \times C^\infty(M; V) \longrightarrow C^\infty(M; V),$$

denoted

$$X, s \mapsto \nabla_X s,$$

such that

- (1) $\nabla_{fX} s = f \nabla_X s$, and
- (2) $\nabla_X (fs) = (X \cdot f)s + f \nabla_X s$,

where $(X \cdot f)$ is the usual directional derivative associated to a vector field.

For $V = TM$, we have the usual Lie bracket

$$[-, -]: C^\infty(M; TM) \times C^\infty(M; TM) \longrightarrow C^\infty(M; TM)$$

sending $X, Y \mapsto [X, Y]$; if $f, g: M \rightarrow \mathbb{R}$ are functions, then

$$[fX, gY] = f g [X, Y] + f (X \cdot g) Y - g (Y \cdot f) X.$$

This operator is the commutator of an infinitesimal flow of X and an infinitesimal flow of Y .

Definition 1.37. Let ∇ be a covariant derivative for the tangent bundle. Its *torsion* is

$$\tau(X, Y) := \nabla_X Y - \nabla_Y X - [X, Y].$$

Exercise 1.38. Show that $\tau(fX, gY) = f g \tau(X, Y)$ and $\tau(X, Y) = -\tau(Y, X)$.

Let's write this out in local coordinates. There are two things we could mean – coordinates on M or on V . Since V is a vector bundle, we can use for its coordinates the coordinates of M and a (local) basis of sections s_1, \dots, s_r . (Global nonvanishing sections might not exist at all, e.g. $TS^2 \rightarrow S^2$). In this case, you can differentiate s_j , obtaining some linear combination of the sections depending on x in a neighborhood U :

$$\nabla_X s_j = \Gamma_j^i(x) s_i.$$

This is just parameterized linear algebra. These Γ_j^i are 1-forms on U . We can also obtain coordinates for these 1-forms: if we let

$$\nabla_{\partial/\partial x^\mu} s_j = \Gamma_{j\mu}^i s_i,$$

then $\Gamma_j^i = \Gamma_{j\mu}^i dx^\mu$.

If $V \rightarrow M$ has an inner product (metric), a positive definite pairing $C^\infty(M; V) \times C^\infty(M; V) \rightarrow C^\infty(M)$ sending $s_1, s_2 \mapsto \langle s_1, s_2 \rangle$, we can ask how a covariant derivative interacts with it.

Definition 1.39. A covariant derivative is *compatible with the metric* if for all $X \in C^\infty(M; TM)$ and $s_1, s_2 \in C^\infty(M; V)$,

$$X \cdot \langle s_1, s_2 \rangle = \langle \nabla_X s_1, s_2 \rangle + \langle s_1, \nabla_X s_2 \rangle.$$

Definition 1.40. A section $s \in C^\infty(M; V)$ is *parallel* if $\nabla_X s = 0$ for all X .

Parallel sections exist in \mathbb{A}^n but not in general; the obstruction is called the curvature.

Definition 1.41. The *curvature* of a covariant derivative ∇ is

$$K(X, Y) := [\nabla_X, \nabla_Y] - \nabla_{[X, Y]},$$

i.e.

$$K(X, Y)(s) := \nabla_X \nabla_Y s - \nabla_Y \nabla_X s - \nabla_{[X, Y]} s.$$

If M is Riemannian, there's a beautiful theorem about how all of these structures interact.

Theorem 1.42 (Levi-Civita). *Let M be a Riemannian manifold. Then, there is a unique connection on $TM \rightarrow M$ which is torsion-free and compatible with the metric.*

Exercise 1.43. Prove this theorem. The way you do this is to compute $\langle \nabla_X Y, Z \rangle$, because if you know this for all Z , you know $\nabla_X Y$. Using the torsion-free and metric compatibility conditions, you can expand it out, and after some number of steps, you'll get the answer.

This local but non-global parallelism is an important property of Riemannian manifolds.

Next we will write a local formula for this connection. Suppose we have local coordinates x^1, \dots, x^n on an open set $U \subset M$; then, we obtain the symbols $\Gamma_{jk}^i: U \rightarrow \mathbb{R}$. If we define the inner product and the Lie bracket, we can write down formulas for them. Namely, if we let

$$g_{ij} := \left\langle \frac{\partial}{\partial x^i}, \frac{\partial}{\partial x^j} \right\rangle,$$

and since

$$\left[\frac{\partial}{\partial x^i}, \frac{\partial}{\partial x^i} \right] = 0,$$

then we can determine equations that the Γ_{jk}^i must satisfy. These can be encoded in the Riemann curvature tensor $R(X, Y)Z$, and in coordinates, on lets

$$R_{jkl}^i \frac{\partial}{\partial x^i} = R\left(\frac{\partial}{\partial x^k}, \frac{\partial}{\partial x^\ell}\right) \frac{\partial}{\partial x^j}.$$

This tensor has a bunch of important symmetries. The curvature is a 2-form on the manifold, but valued in $\text{End}(TM)$: X and Y are the two directions you're testing, and are the 2 components of the 2-form.

The symmetry $R_{jkl}^i = -R_{jlk}^i$ means that $R(X, Y)-$ is a skew-symmetric endomorphism of TM .

You can also lower an index by defining

$$(1.44) \quad R_{ijkl} = \langle R(\partial_k, \partial_\ell) \partial_j, \partial_i \rangle,$$

and skew-symmetry means

$$R_{ijkl} = -R_{jikl}.$$

These are the two "easier" symmetries, in that they don't use much specifically about R . A more interesting one is

$$R_{ijkl} + R_{iklj} + R_{iljk} = 0,$$

and the fourth identity, which follows from the other three, is

$$R_{ijkl} = R_{klij}.$$

Exercise 1.45. Compute the dimension of the vector space of tensors which satisfy these identities, as a subspace of $(V^*)^{\otimes 4}$.

Lecture 2.

Principal G -bundles: 1/24/18

The first part of today's talk was given by George Torres, corresponding to the first part of Chapter 2 of Roe's book.

Definition 2.1. A *Lie group* G is a group that is also a smooth manifold, and such that multiplication $m: G \times G \rightarrow G$ and inversion $i: G \rightarrow G$ are smooth maps.

Associated to any Lie group G is its *Lie algebra* $\mathfrak{g} := T_e G$. There is a *Lie bracket* operation

$$[-, -]: \mathfrak{g} \times \mathfrak{g} \longrightarrow \mathfrak{g}$$

defined like last time (use multiplication on G to extend tangent vectors to G -invariant vector fields, then take their commutators).

Definition 2.2. Let G be a Lie group and M be a smooth manifold. A *principal G -bundle* is a map of smooth manifolds $\pi: E \rightarrow M$ together with a smooth right action of G on E whose orbits are fibers of π , such that G acts freely and transitively on each fiber.

This implies that for each $m \in M$, $\pi^{-1}(M) \cong G$ noncanonically, and $E/G \cong M$.

Connections on principal bundles are analogous to those on vector bundles; the goal is to define a horizontal subspace of the bundle, and use that and its G -translates to define parallel transport. Then, one must show that connections always exist, but this turns out to be true. This definition of connections is sometimes called an *Ehresmann connection*.

In order to define connections, we'll need a few preliminary definitions.

Definition 2.3. Let G be a Lie group. The *adjoint representation* of G is the map $\text{Ad}: G \rightarrow \text{End}(\mathfrak{g})$ which sends a $g \in G$ to $d|_e \psi_g$, where $\psi_g: G \rightarrow G$ is conjugation by G and $d|_e$ is differentiating at the identity.

Definition 2.4. Let $\pi: E \rightarrow M$ be a principal G -bundle and $\rho: G \rightarrow \text{Aut}(F)$ be a (real, finite-dimensional) representation of G . The *associated vector bundle* to E and F is the $E \times_G F := E \times F / \sim$, where $(e \cdot g, f) \sim (e, g \cdot f)$ for all $e \in E$, $f \in F$, and $g \in G$.

Exercise 2.5. Show that $E \times_G F \rightarrow M$ is indeed a vector bundle.

Definition 2.6. Let $\pi: E \rightarrow M$ be a principal G -bundle. A *vertical vector field* is a vector field v on E such that for all $p \in M$, $v_p = d_e A_p(u)$ for some $u \in \mathfrak{g}$, where $A_p: G \rightarrow E$ is an identification of the fiber E_p with G .⁴

⁴TODO: I might have gotten this wrong.

An equivalent, more intuitive, definition is that a vertical vector field is contained within the *vertical subbundle* of E , i.e. the kernel of $d\pi$. A third equivalent definition is that V is vertical if for all $p \in M$, there is some $u \in \mathfrak{g}$ such that

$$(2.7) \quad v_p = \left. \frac{d}{dt}(p \exp(tu)) \right|_{t=0}.$$

Let $R_g : E \rightarrow E$ denote the right action of G , sending $e \mapsto e \cdot g$. We can pushforward by this map: let

$$(Rg)_* v := \left. \frac{d}{dt}(p \exp(tu)g) \right|_{t=0}.$$

In particular, the pushforward of a vertical vector field is still vertical.

Horizontal differential forms are dual to vertical vector fields.

Definition 2.8. Let $\alpha \in \Omega^p(E)$. Then α is *horizontal* if for all vertical vector fields X_1, \dots, X_p , $\alpha(X_1, \dots, X_p) = 0$. If in addition $(Rg)^* \omega = \omega$, we say α is *invariant* under the G -action.

Example 2.9. For any $\beta \in \Omega^p(M)$, $\pi^* \beta$ is an invariant horizontal form:

$$(R_g)^* \pi^* \beta = (\pi \circ R_g)^* \beta = \pi^* \beta. \quad \blacktriangleleft$$

More generally, we can consider G -equivariant forms.

Definition 2.10. Let $\rho : G \rightarrow \text{Aut}(F)$ be a representation and $f : E \rightarrow F$ be a smooth map. Then, f is ρ -equivariant if for all $e \in E$ and $g \in G$,

$$f(eg) = \rho(g^{-1})f(e).$$

Invariance is the same thing as equivariance for the trivial representation.

Lemma 2.11 (Correspondence lemma). *Let $f : E \rightarrow F$ be as in the previous definition. There is a bijective correspondence between ρ -equivariant maps $f : E \rightarrow F$ and sections of $E \times_G F$.*

Proof. Let f be a ρ -equivariant map; then, we define a section s_f to send $x \in M$ to $(x, f(x)) \in E \times F$. To check that this is indeed a section, we need it to commute with the G -action on $E \times F$, and this follows because

$$\begin{aligned} s_f(x) \cdot g &= (x, f(x)) \cdot g = (xg, \rho(g^{-1})f(x)) \\ &= (xg, f(xg)) \\ &= s_f(x \cdot g). \end{aligned}$$

Conversely, let $s : M \rightarrow E \times_G F$ be a section, and consider the diagram

$$(2.12) \quad \begin{array}{ccccc} E & \xrightarrow{\sigma} & E \times F & \xrightarrow{\pi_1} & E \\ \downarrow \pi & & \downarrow p & & \downarrow \pi \\ M & \xrightarrow{s} & E \times_G F & \longrightarrow & M, \end{array}$$

where p is the quotient map, π_1 is projection onto the first factor, and σ is defined such that the composition of the maps across the top is the identity. The key observations are

- (1) the right-hand square is a pullback square, and
- (2) $f = \pi_2 \sigma$ is ρ -equivariant.

The first property is true because for any $g \in G$,

$$\begin{aligned} \pi_1(\sigma(xg)) &= xg \\ &= \pi_1(\sigma(x)) \cdot g = \pi_1(\sigma(x \cdot g)), \end{aligned}$$

and along the other corner,

$$\begin{aligned} p(\sigma(xg)) &= s\pi(xg) = s\pi(x) \\ &= p\sigma(x) \\ &= p\sigma(x \cdot g), \end{aligned}$$

so $\sigma(xg) = \sigma(x) \cdot g$. **TODO:** I missed the last part, that f is ρ -equivariant (and why these two properties suffice). \square

Exercise 2.13. Finish the proof by checking that these assignments are mutual inverses.

With these definitions in mind, we can define connections.

Definition 2.14. Let $\pi: E \rightarrow M$ be a principal G -bundle. A *connection* on E is a subbundle $H \subset TE$ (H for “horizontal”) such that

- (splitting) there is another subbundle $V \subset TE$ such that for all $u \in E$, $T_u E = V_u \oplus H_u$, and
- (G -invariance) $d|_e R_g(H_u) = H_{u \cdot g}$.

Using the splitting lemma for vector bundles, the first condition is equivalent to the existence of a split short exact sequence

$$0 \longrightarrow V \longrightarrow TE \longrightarrow \pi^* TM \longrightarrow 0.$$

The splitting is determined by a section $\pi^* TM \rightarrow TE$ (which defines what we call “horizontal”) or by a section $TE \rightarrow V$. This leads to an equivalent definition of a connection on a principal bundle, which is also useful: a connection on E is a \mathfrak{g} -valued 1-form ω on E , called the *connection one-form*, such that

- (G -invariance) for any $\xi \in T_u E$, $\omega(\xi \cdot g) = \text{Ad}(g^{-1})\omega(\xi)$, and
- (splitting) for any $v \in \mathfrak{g}$, $\omega(X_v) = v$.

The splitting lemma guarantees you can always split short exact sequences of vector bundles. But to show that connections exist, we need to address G -invariance, which is not as immediate.

Lemma 2.15. If G is a Lie group and $\pi: E \rightarrow M$ is a principal G -bundle, then there is a connection on E .

Proof. Let $\underline{\mathfrak{g}} := E \times \mathfrak{g}$ and $V := E \times_G \mathfrak{g}$, where G acts on \mathfrak{g} by the adjoint action. Then consider the diagram

$$(2.16) \quad \begin{array}{ccccccc} 0 & \longrightarrow & \underline{\mathfrak{g}} & \longrightarrow & TE & \longrightarrow & \pi^* TM \longrightarrow 0 \\ & & \downarrow & & \downarrow & & \downarrow \\ 0 & \longrightarrow & V & \longrightarrow & Q & \longrightarrow & TM \longrightarrow 0. \end{array}$$

Proving that this commutes takes a while, so we won’t delve into the details; one reference is Atiyah, “Complex analytic connections on fibre bundles.”

The point of introducing (2.16) is that if we can lift a splitting from the bottom row to the top row, it will be a G -invariant splitting, hence a connection. So choose a splitting $\sigma: TM \rightarrow Q$, which splits the bottom row of (2.16). Then we have a diagram

$$(2.17) \quad \begin{array}{ccc} \pi^* TM & \longrightarrow & TM \\ \downarrow & & \downarrow \sigma \\ TE & \longrightarrow & Q \\ \downarrow & & \downarrow \\ E & \xrightarrow{\pi} & M. \end{array}$$

The bottom rectangle is a pullback, and the total rectangle ($\pi^* TM$, TM , E , and M) is a pullback. Therefore by the universal property of pullbacks, the top rectangle also is a pullback, and this implies that σ lifts across it to something G -invariant. \square

There is also a more geometric proof.

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Next we’ll talk about exterior derivatives.

Definition 2.18. Let $\pi: E \rightarrow M$ be a principal G -bundle with connection H and let $p_\omega: \Omega^*(E) \rightarrow \Omega^*(H)$ denote projection onto the horizontal subspace. The *exterior covariant derivative* is the composition

$$\Omega^p(E) \xrightarrow{d} \Omega^{p+1}(E) \xrightarrow{p_\omega} \Omega^{p+1}(H).$$

Proposition 2.19. Let $\rho: G \rightarrow \mathrm{GL}(F)$ be a representation, $\alpha \in \Omega_E^p(F)$ be an F -valued, ρ -equivariant horizontal p -form on E , and ω denote the (\mathfrak{g} -valued) connection one-form on E . Then

$$p_\omega(d\alpha) = d\alpha + \rho_* \omega \wedge \alpha.$$

Definition 2.20. Let ω be a differential 1-form. Then, the *curvature* Ω of ω is the 2-form

$$\Omega(X_1, X_2) = d\omega(X_1, X_2) + [\Omega(X_1), \Omega(X_2)].$$

Exercise 2.21. With notation as above,

$$(p_\omega d)^2 \alpha = \rho_* \Omega \wedge \alpha.$$

The point is that acting on α is the same as wedging with $\rho_* \Omega$, and this tells you something about what Ω is doing.

Remark 2.22. In some of these formulas, it's important to be careful about what the wedge products are doing. For example, we once or twice saw $\omega \wedge \alpha$, where $\omega \in \Omega_E^1(\mathfrak{g})$ and $\alpha \in \Omega_E^1(F)$. If \mathfrak{g} is a matrix algebra, we can organize the components of ω into a matrix of differential forms, and F is a representation of the matrix algebra \mathfrak{g} , so α is a vector. In this case, the wedge product is a combination of matrix multiplication and the wedge product:

$$(2.23) \quad \begin{pmatrix} \omega_1^1 & \omega_2^1 \\ \omega_1^2 & \omega_2^2 \end{pmatrix} \wedge \begin{pmatrix} \alpha^1 \\ \alpha^2 \end{pmatrix} = \begin{pmatrix} \omega_1^1 \wedge \alpha^1 + \omega_2^1 \wedge \alpha^2 \\ \omega_1^2 \wedge \alpha^1 + \omega_2^2 \wedge \alpha^2 \end{pmatrix}.$$

The same is true for computing $\omega \wedge \omega$; in particular, this is not automatically zero.

Therefore one sometimes sees the formula for curvature written

$$(2.24) \quad \Omega := d\omega + [\omega \wedge \omega].$$

What does this mean? We have $\omega \wedge \omega \in \Omega_E^2(\mathfrak{g} \otimes \mathfrak{g})$, and the Lie bracket $[-, -]: \mathfrak{g} \otimes \mathfrak{g} \rightarrow \mathfrak{g}$. We implement this on differential forms by

$$(\omega^i \wedge \omega^j) e_i \otimes e_j \mapsto (\omega^i \wedge \omega^j) [e_i, e_j].$$

In particular,

$$\begin{aligned} [\omega \wedge \omega](X_1, X_2) &= [\omega(X_1), \omega(X_2)] - [\omega(X_2), \omega(X_1)] \\ &= 2[\omega(X_1), \omega(X_2)], \end{aligned}$$

which is why (2.24) has an extra 1/2 in it compared to the first definition. \triangleleft

Remark 2.25. One can even define differential forms valued in vector bundles: $\Omega_M^*(E \times_G F)$ consists of sections of the exterior powers of $E \times_G F$. Alternatively, you can think of these as valued in $(E \times_G F)_p$ at a point p ; the vector spaces changes as p moves, but that's okay. The quotient map $\pi: E \rightarrow G$ defines a pullback $\Omega_M^*(E \times_G F) \rightarrow \Omega_E^*(F)$. This provides yet another interpretation of the definition of a connection.

- Invariance is that $\alpha \in \mathrm{Im}(\pi^*)$ iff $R_g^* \alpha = \rho(g)^{-1} \alpha$ for all $g \in G$.
- Splitting comes from the fact that $\iota_\zeta \alpha = 0$ when ζ is vertical. (This denotes *contraction*: $\iota_\zeta \alpha(X) := \alpha(\zeta, X)$).

It's a good exercise to check, to get practice manipulating these vector- or bundle-valued forms. But principal bundles make some of these computations easier, by turning some bundle-valued forms into constant vector space-valued forms. \triangleleft

Lecture 3.

Characteristic classes: 1/24/18

The second part of today's lecture was given by Riccardo Pedrotti, on characteristic classes from a geometric perspective.

The theory of characteristic classes comes from the simple question: how can we tell two vector bundles apart? For instance, how do we know that the tangent bundle to the 2-sphere is non-trivial? Characteristic classes gives a systematic approach.

Definition 3.1. A *characteristic class* c is a natural transformation which to each vector bundle V over a manifold M associates an element $c(V)$ of the cohomology group $H^*(M)$, with property that if $V_1 \simeq V_2$ then $c(V_1) = c(V_2)$.

The idea of Chern-Weil theory is the following: suppose that our bundle V is equipped with a connection. In some sense, the curvature of this connection measures the local deviation of V from flatness. Now if V is flat, and the base manifold M is simply connected, then V is trivial. This suggests that there may be a link between curvature and characteristic classes, which measure the global deviation of V from triviality. Such a link is provided by the theory of invariant polynomials.

By polynomial function we mean the following:

Definition 3.2. Let $\mathfrak{gl}_m(\mathbb{C})$ denote the Lie algebra of $m \times m$ matrices over \mathbb{C} . A *homogeneous polynomial function* P on $\mathfrak{gl}_m(\mathbb{C})$ is a function such that there exists a $\tilde{P} \in \text{Sym}^k(\mathbb{C}^m)^*$ such that $P(A) = \tilde{P}(A, A, \dots, A)$. A *polynomial function* is a sum of homogeneous ones.

Definition 3.3. An *invariant polynomial* on $\mathfrak{gl}_m(\mathbb{C})$ is a polynomial function $P: \mathfrak{gl}_m(\mathbb{C}) \rightarrow \mathbb{C}$ such that for all $X, Y \in \mathfrak{gl}_m(\mathbb{C})$, $P(XY) = P(YX)$. An *invariant formal power series* is a formal power series over $\mathfrak{gl}_m(\mathbb{C})$ each of whose homogeneous components is an invariant polynomial.

For example, the determinant and the trace are invariant polynomials.

Lemma 3.4. The ring of invariant polynomials on $\mathfrak{gl}_m(\mathbb{C})$ is a polynomial ring generated by the polynomials

$$c_k(X) = (-2\pi i)^{-k} \text{tr}(\Lambda^k X),$$

where $\Lambda^k X$ denotes the transformation induced by X on $\Lambda^k \mathbb{C}^m$.

Proof. Let P be any invariant polynomial. Restricting P to diagonal matrices, we see that P must be a polynomial function of the diagonal entries. Since these diagonal entries can be interchanged by conjugation, P must in fact be a *symmetric* polynomial function. Now since P is invariant under conjugation, it must be a symmetric polynomial function of the eigenvalues for all matrices with distinct eigenvalues, since by elementary linear algebra such matrices are conjugate to diagonal matrices. The set of such matrices is dense in $\mathfrak{gl}_m(\mathbb{C})$, so a continuity argument shows that P is just a symmetric polynomial function in the eigenvalues. Now it is easy to see that $\text{tr}(\Lambda^k X)$ is the k^{th} elementary symmetric function in the eigenvalues of X . The main theorem on symmetric polynomials states that the ring of symmetric polynomials is itself a polynomial ring generated by the elementary symmetric functions, and this now completes the proof. \square

Example 3.5. To make the idea of the proof more concrete, let $m = 4$. A 4×4 diagonal matrix is of the form

$$X = \begin{pmatrix} a & 0 & 0 & 0 \\ 0 & b & 0 & 0 \\ 0 & 0 & c & 0 \\ 0 & 0 & 0 & d \end{pmatrix}$$

for some $a, b, c, d \in \mathbb{R}$. Then, $\Lambda^2 X$ acts on $e_1 \wedge e_2$ by $Xe_1 \wedge Xe_2 = abe_1 \wedge e_2$. Therefore $\Lambda^2 X$ is a 6×6 diagonal matrix with diagonal terms ab, ac, ad, bc, bd , and cd , and therefore its trace is

$$\text{tr}(\Lambda^2 X) = ab + ac + ad + bc + bd + cd.$$

This is a quadratic polynomial, and is symmetric; it's an example of an *elementary symmetric polynomial*. There's a theorem that the ring of all symmetric polynomials are generated by these elementary symmetric polynomials. \blacktriangleleft

Now let V be a complex vector bundle over M with connection ∇ and curvature $K \in \Omega_M^2(\text{End}(V))$. Choosing a local framing for V , we may locally identify K with a matrix of ordinary 2-forms. Hence, if P is an invariant polynomial, we can apply P to this matrix to get an even-dimensional differential form $P(K)$. *A priori*, this depends on the choice of local framing, but since P is invariant, $P(K)$ doesn't depend on the choice, and is therefore globally defined.

In terms of the principal $\text{GL}_m(\mathbb{C})$ -bundle E associated to V , this construction may be phrased as follows. Let Ω be the curvature form of the induced connection on E ; Ω is a horizontal, equivariant 2-form on E with values in $\mathfrak{gl}_m(\mathbb{C})$, so $P(\Omega)$ is a horizontal invariant form on E . Such a form is the lift to E of a form on M , and this form is $P(K)$.

Since 2-forms are nilpotent elements in the exterior algebra $\Omega_M^*(\mathfrak{gl}_m(\mathbb{C}))$, all formal power series with 2-form-valued variables in fact converge. Thus, this construction makes good sense if P is merely an invariant formal power series.

Proposition 3.6. *For any invariant polynomial (or formal power series) P , the differential form $P(K)$ is closed, and its de Rham cohomology class is independent of the choice of connection ∇ on V .*

Proof. For the purposes of this proof call an invariant formal power series P as *respectable* if the conclusion of the proposition holds for P . Clearly the sum and product of respectable formal power series are respectable. Thus, it is enough to prove that the generators defined in Lemma 3.4 are respectable. Equivalently, since

$$\det(1 + qK) = \sum q^k \operatorname{tr}(\Lambda^k K),$$

it is enough to prove that $\det(1 + qK)$, considered as a formal power series depending on the parameter q , is respectable.

If P is a respectable formal power series with constant term a , and g is a function holomorphic in a neighborhood of a , then $g \circ P$ is also a respectable formal power series. Hence, $\det(1 + qK)$ is respectable if and only if $\log \det(1 + qK)$ is respectable. We will now prove directly that $\log \det(1 + qK)$ is respectable.

For this purpose we will work in the associated principal $\mathrm{GL}_n(\mathbb{C})$ -bundle E of frames for V , with matrix-valued connection 1-form ω and corresponding curvature 2-form Ω . Recall the formula

$$\Omega = d\omega + \omega \wedge \omega$$

where the product in the ring of matrix-valued forms is obtained by tensoring exterior product and matrix multiplication as in (2.23).

Now suppose that ω depends on a parameter t ; then Ω also depends on t , and if we use a dot to denote differentiation with respect to t , then

$$\dot{\Omega} = d\dot{\omega} + \omega \wedge \dot{\omega} + \dot{\omega} \wedge \omega.$$

Consider

$$(3.7) \quad \frac{d}{dt} \log \det(1 + q\Omega) = q \operatorname{tr}(\dot{\Omega}(1 + q\Omega)^{-1})$$

$$(3.8) \quad = \sum_{\ell=0}^{\infty} (-1)^{\ell} q^{\ell+1} \operatorname{tr}(\Omega^{\ell}(d\dot{\omega} + \omega \wedge \dot{\omega} + \dot{\omega} \wedge \omega))$$

where (3.7) is justified by the formula

$$\frac{d}{dt} \det A(t) = (\det A(t)) \cdot \operatorname{tr}(\dot{A}(t)A(t)^{-1}),$$

and (3.8) is justified by the power series expansion

$$\frac{1}{1+z} = \sum_{i=0}^{\infty} (-1)^i z^i.$$

We also need the *second Bianchi identity*

$$d\Omega = \Omega \wedge \omega - \omega \wedge \Omega,$$

which can be proven directly from the definition of the exterior derivative. Using this, plus the fact that trace is symmetric, we have that

$$\begin{aligned} \operatorname{tr}(\Omega^{\ell}(\omega \wedge \dot{\omega} + \dot{\omega} \wedge \omega)) &= \operatorname{tr}(\Omega^{\ell} \wedge \omega \wedge \dot{\omega} - \omega \wedge \Omega^{\ell} \wedge \dot{\omega}) \\ &= \operatorname{tr}((d\Omega^{\ell}) \wedge \dot{\omega}). \end{aligned}$$

Therefore

$$\begin{aligned} \operatorname{tr}(\Omega^{\ell} \wedge (d\dot{\omega} + \omega \wedge \dot{\omega} + \dot{\omega} \wedge \omega)) &= \operatorname{tr}((d\Omega^{\ell}) \wedge \dot{\omega} + \Omega^{\ell} \wedge d\dot{\omega}) \\ &= d \operatorname{tr}(\Omega^{\ell} \wedge \dot{\omega}), \end{aligned}$$

so (3.7) simplifies to

$$(3.9) \quad \frac{d}{dt} \log \det(1 + q\Omega) = d \sum_{\ell=0}^{\infty} (-1)^{\ell} q^{\ell+1} \operatorname{tr}(\Omega^{\ell} \wedge \dot{\omega}),$$

and in particular is an exact form on E .

In fact, it is the exterior derivative of a horizontal and invariant form on E : $\dot{\omega}$ is horizontal and G -equivariant since it is a 1-form on M (it follows from Lemma 2.11 that the space of connections is an affine space modeled on

the vector space of V -valued forms on M), Ω is horizontal and G -equivariant as well. Hence $\text{tr}(\Omega^\ell \wedge \dot{\omega})$ is invariant since the trace is an invariant polynomial, and is horizontal since $\Omega^\ell \wedge \dot{\omega}$ is.

Therefore, the projection to the base manifold

$$\frac{d}{dt} \log \det(1 + qK)$$

is also exact. Now the result follows; for since any connection can be deformed locally to flatness (i.e. $K = 0$), we see that $\log \det(1 + qK)$ is locally exact, hence closed, and since any two connections can be connected by a smooth path, the cohomology class of $\log \det(1 + qK)$ is independent of the choice of connection, since their difference is an exact form. \square

It follows from the proposition that any invariant formal power series P defines a characteristic class for complex vector bundles, by the recipe “pick any connection and apply P to the curvature.”

Definition 3.10. The k^{th} Chern class is the characteristic class corresponding to the generators c_k defined in Lemma 3.4.

Remark 3.11. We immediately see from the definition of Chern classes that if a complex vector bundle has rank m , then $c_k = 0$ for $k > m$: $\Lambda^k K$ is the linear transformation induced by K on $\Lambda^k \mathbb{C}^m$, and for $k > m$, the latter is trivial. Naturality comes from the fact that if on a local patch U_i , E has the local connection form ω_i , then on $f^{-1}(U_i)$, the curvature is $f^* \Omega_i$. \triangleleft

Lemma 3.12. Let V be a real vector bundle and $V_{\mathbb{C}}$ denote its complexification. Then, $c_{2k+1}(V_{\mathbb{C}}) = 0$.

Proof. We can give V a metric and compatible connection. The curvature of such a connection is skew (i.e. $\mathfrak{o}(m)$ -valued), so

$$\text{tr}(\Lambda^k F) = (-1)^k \text{tr}(\Lambda^k F).$$

To see this, recall that the coefficients of the characteristic equation for F are exactly $\text{tr}(\Lambda^k F)$ up to a sign.⁵ If λ is an eigenvalue of a skew-symmetric matrix, then $-\lambda$ is too, and on \mathbb{C} this means that the characteristic polynomial is up to a constant the product of polynomials $(z^2 - \lambda^2)$, so there are no coefficients of odd index, hence proving that for k odd, $\Lambda^k F$ is traceless. \square

Genera. Holomorphic functions can be used to build important combinations of characteristic classes. Let $f(z)$ be any function holomorphic near $z = 0$. We can use f to construct an invariant formal power series Π_f by defining

$$\Pi_f(X) := \det\left(f\left(-\frac{1}{2\pi i}X\right)\right).$$

Again, to make sense of this, we need to sidestep convergence issues! But since we'll just be applying this to differential forms, which are nilpotent, this is okay.

The associated characteristic class is called the *Chern f -genus*. It has a few nice properties.

Lemma 3.13. If $L \rightarrow M$ is a complex line bundle, $\Pi_f(L) = f(c_1(L))$.

Proof. This comes from the fact that in this case the curvature is a $\mathfrak{gl}_1(\mathbb{C})$ -valued 2-form, so

$$\begin{aligned} \Pi_f(L) &= \Pi_f(K_L) = \det\left(f\left(-\frac{1}{2\pi i}K_L\right)\right) \\ &= f\left(-\frac{1}{2\pi i}K_L\right) \\ &= f\left(\text{tr}\left(-\frac{1}{2\pi i}K_L\right)\right) = f(c_1(L)). \end{aligned} \quad \square$$

Lemma 3.14. For any complex vector bundles V_1 and V_2 , $\Pi_f(V_1 \oplus V_2) = \Pi_f(V_1)\Pi_f(V_2)$.

Proof sketch. Compute using a direct sum connection, which gives rise to a curvature matrix which is a block matrix. \square

⁵There's a sign convention here; this is true using our definition $\det(1 - qK)$. An alternative choice is to use $\det(q - K)$, in which case one must swap the indices to preserve evenness.

Now observe a very useful property: if the eigenvalues of the matrix $(-1/2\pi i)X$ are $\{x_j\}$, then

$$(3.15) \quad \Pi_f(X) = \prod f(x_j)$$

is a symmetric formal power series in the x_j , and can therefore be expressed in terms of the elementary symmetric functions of the x_j . But these elementary symmetric functions are just the Chern classes. Thus in the literature the genus $\Pi_f(V)$ is often written just as in (3.15), where x_1, \dots, x_m are formal variables subject to the relations

$$\begin{aligned} x_1 + x_2 + \dots + x_m &= c_1, \\ x_1 x_2 + \dots + x_{m-1} x_m &= c_2, \end{aligned}$$

and so on.

Example 3.16. The genus associated to $f(z) = 1 + z$ is the *total Chern class*

$$c(V) := 1 + c_1(V) + c_2(V) + c_3(V) + \dots.$$

To see this, consider the power expansion of the determinant:

$$\det\left(1 - \frac{1}{2\pi}X\right) = \sum_k \left(-\frac{1}{2\pi}\right)^k \text{tr}(\Lambda^k X) = \sum_k c_k(X).$$

From this we immediately get that $c(V_1 \oplus V_2) = c(V_1)c(V_2)$. ◀

Definition 3.17. Let $V \rightarrow M$ be a real vector bundle and g be a holomorphic function near 0, with $g(0) = 1$. Let $V_{\mathbb{C}}$ be the complexification of V . Denote by f be the branch of $z \mapsto (g(z^2))^{1/2}$ which has $f(0) = 1$; we call the genus associated to f the *Pontrjagin g -genus* of V .

Since f is an even function of z , the associated genus involves only the even Chern classes.

Lemma 3.18. Let g be as above. Then for a real vector bundle V , the Pontrjagin g -genus is equal to

$$\prod_j g(y_j)$$

for some formal variables y_j .

Definition 3.19. Let V be a real vector bundle. Its k^{th} Pontrjagin class $p_k(V)$ is the k^{th} elementary symmetric function in the formal variables y_j .

Proof of Lemma 3.18. Regard this as an identity between invariant polynomials over $\mathfrak{o}(n)$. Any matrix in $\mathfrak{o}(n)$ is similar to one in block diagonal form, where the blocks are 2×2 and are of the form

$$X = \begin{pmatrix} 0 & \lambda \\ -\lambda & 0 \end{pmatrix}$$

with eigenvalues $\pm i\lambda$. Since both sides of the desired identity are multiplicative for direct sums, it is enough to prove it for this block X , whose first two elementary symmetric functions are

$$\begin{aligned} c_1(X) &= \left(-\frac{1}{2\pi i}\right) \text{tr}(X), \\ c_2(X) &= \left(-\frac{1}{2\pi i}\right)^2 \text{tr}(\Lambda^2 X). \end{aligned}$$

Since X is skew, then its trace vanishes, so $c_1(X) = 0$. By looking at the characteristic polynomial of X we see that $\text{tr}(\Lambda^2 X) = \lambda^2$, giving

$$c_2(X) = -\frac{\lambda^2}{4\pi^2}.$$

Thus

$$y = p_1(X) = \frac{\lambda^2}{4\pi^2}.$$

On the other hand, X is similar over \mathbb{C} to

$$\begin{pmatrix} -i\lambda & 0 \\ 0 & i\lambda \end{pmatrix},$$

so

$$\Pi_f(X) = f\left(-\frac{\lambda}{2\pi}\right) f\left(\frac{\lambda}{2\pi}\right) = g\left(\frac{\lambda^2}{4\pi^2}\right) = g(y)$$

as required. \square

Two important examples are the \hat{A} -genus, which is the Pontrjagin genus associated to the holomorphic function

$$z \mapsto \frac{\sqrt{z}/2}{\sinh(\sqrt{z}/2)}$$

and the Hirzebruch L -genus, which is the Pontrjagin genus associated with the holomorphic function

$$z \mapsto \frac{\sqrt{z}}{\tanh(\sqrt{z})}.$$

Lecture 4.

Clifford algebras, Clifford bundles, and Dirac operators: 1/31/18

Ricky spoke today about Clifford algebras and Clifford bundles.

Let k be a field with characteristic not equal to 2. If V is a vector space over k , its *tensor algebra* is

$$(4.1) \quad T(V) := \bigoplus_{k \geq 0} V^{\otimes k},$$

where $V^{\otimes 0} := k$.

Definition 4.2. Let V be a k -vector space with a quadratic form $Q: V \times V \rightarrow k$. Let $I_Q \subset T(V)$ denote the two-sided ideal generated by elements of the form $v \otimes v + Q(v)$ for $v \in V$. Then, the quotient algebra

$$\mathcal{Cl}(V, Q) := T(V)/I_Q$$

is called the *Clifford algebra* of V and Q .

Example 4.3. The zero function is a quadratic form, so $\mathcal{Cl}(V, 0)$ is $T(V)/(v \otimes v = 0)$, which is just the exterior algebra $\Lambda(V)$ of V . \blacktriangleleft

There is a natural map $i: V \rightarrow \mathcal{Cl}(V, Q)$ which is the composition

$$V = V^{\otimes 1} \hookrightarrow T(V) \xrightarrow{\pi_Q} \mathcal{Cl}(V, Q).$$

Lemma 4.4. $i: V \rightarrow \mathcal{Cl}(V, Q)$ is injective.

This is not too hard to check.

Moreover, 1 and V generate $\mathcal{Cl}(V, Q)$, subject to the relations $v^2 = q(v, v)$. To get a smaller set of generators, we can choose a basis of V . From now on, we assume Q is positive definite and choose an orthonormal basis e_1, \dots, e_n of V . In this case, $\mathcal{Cl}(V, Q)$ is generated by $1, e_1, \dots, e_n$ subject to the relations

$$(4.5) \quad \begin{aligned} e_i^2 &= 1 \\ e_j \cdot e_j &= -e_j \cdot e_i, \end{aligned}$$

because $e_i^2 = -Q(e_i) = -1$. This implies the following fact.

Proposition 4.6. The set

$$\{e_{i_1} \cdots e_{i_k} \mid 1 \leq i_1 < \cdots < i_k \leq n, 1 \leq k \leq n\}$$

is a basis for $\mathcal{Cl}(V, Q)$ as a vector space, and hence

$$\dim \mathcal{Cl}(V, Q) = 2^n = \sum_{k=0}^n \binom{n}{k}.$$

Example 4.7. Let $\langle \cdot, \cdot \rangle$ denote the standard inner product on \mathbb{R}^n ; then, $\mathcal{Cl}_n := \mathcal{Cl}(\mathbb{R}^n, \langle \cdot, \cdot \rangle)$. In low dimensions these are familiar.

- $\mathcal{Cl}_1 = \langle 1 \rangle \oplus \langle e \rangle$ with $e^2 = 1$, hence $\mathcal{Cl}_1 \cong \mathbb{C}$.
- $\mathcal{Cl}_2 = \langle 1 \rangle \oplus \langle e_1 \rangle \oplus \langle e_2 \rangle \oplus \langle e_1 \cdot e_2 \rangle$ with $e_1^2 = e_2^2 = (e_1 e_2)^2 = -1$, so as \mathbb{R} -algebras, $\mathcal{Cl}_2 \cong \mathbb{H}$, the quaternions.

There is a sense in which real Clifford algebras are 8-fold periodict, which is an instance of *Bott periodicity*. We won't delve into this, but see Atiyah-Bott-Shapiro, "Clifford modules," for more information. ◀

Clifford algebras satisfy a universal property.

Proposition 4.8. *Let A be a k -algebra and $\varphi : V \rightarrow A$ be a map of vector spaces such that $\varphi(v)^2 = -Q(v) \cdot 1_A$. Then, there is a unique algebra map $\widehat{\varphi} : \mathcal{Cl}(V, Q) \rightarrow A$ making the following diagram commute:*

$$\begin{array}{ccc} V & \xrightarrow{\varphi} & A \\ & \searrow i & \uparrow \widehat{\varphi} \\ & & \mathcal{Cl}(V, Q) \end{array}$$

Proof. By the universal property of $T(V)$, there's a unique map $\psi : T(V) \rightarrow A$ sending

$$(4.9) \quad v_1 \otimes \cdots \otimes v_k \longmapsto \varphi(v_1)\varphi(v_2)\cdots\varphi(v_k).$$

The claim follows because $\psi(I_Q) = 0$, hence factors through the quotient, which is $\mathcal{Cl}(V, Q)$. ◻

Let \mathbf{QVect}_k denote the category of *quadratic spaces* over k , i.e. vector spaces together with quadratic forms; the morphisms $\varphi : (V_1, Q_1) \rightarrow (V_2, Q_2)$ are data of a linear map $\varphi : V_1 \rightarrow V_2$ such that for all $v, w \in V_1$,

$$(4.10) \quad Q_1(v, w) = Q_2(\varphi(v), \varphi(w)).$$

Using Proposition 4.8, one can show that $\mathcal{Cl} : \mathbf{QVect}_k \rightarrow \mathbf{Alg}_k$ is a functor.

Another use of the universal property is to define a representation

$$\rho_{\mathcal{Cl}} : \mathcal{O}_n \rightarrow \text{Aut}(\mathcal{Cl}_n);$$

a map $\mathbb{R}^n \rightarrow \mathbb{R}^n$ respecting the inner product defines a map $\mathcal{Cl}_n \rightarrow \mathcal{Cl}_n$, and the space of these maps is \mathcal{O}_n .

Definition 4.11. Let (V, Q) be a quadratic space. A *Clifford module* over (V, Q) is a k -vector space S together with a k -linear map $\varphi : \mathcal{Cl}(V, Q) \rightarrow \text{End}_k(S)$.

So it's just a module over the algebra $\mathcal{Cl}(V, Q)$.

Remark 4.12. When S is a complex vector space and V is a real vector space, then we will instead ask for the action map to be an \mathbb{R} -algebra homomorphism $\mathcal{Cl}(V, Q) \rightarrow \text{End}_{\mathbb{C}}(S)$. This is equivalent to having a module over $\mathcal{Cl}(V, Q) \otimes_{\mathbb{R}} \mathbb{C}$. ◀

Remark 4.13. By the universal property, it suffices to specify a map $\varphi : V \rightarrow \text{End}_k(S)$ with $\varphi(v)^2 = -Q(v) \cdot \text{id}$. ◀

Example 4.14.

- (1) Let's consider \mathcal{Cl}_1 as a module over itself. This is the same data as a map $\varphi : \mathcal{Cl}_1 \rightarrow \text{End}_{\mathbb{C}}(\mathbb{C}) \cong \mathbb{C}$; one choice is $1 \mapsto \text{id}$ and $e \mapsto i$.
- (2) We can also make \mathbb{R}^4 into a \mathcal{Cl}_2 -module by having it act on itself by left multiplication. For example, e_1 acts by the matrix

$$\begin{pmatrix} 0 & -1 & & \\ 1 & 0 & & \\ & & 0 & -1 \\ & & 1 & 0 \end{pmatrix}.$$

Now, we apply this to geometry. Let (M, g) be a Riemannian manifold, so for each $p \in M$, $\mathcal{Cl}(T_p M, g_p) \cong \mathcal{Cl}_n$.

Definition 4.15. The *Clifford tangent bundle* is $\mathcal{Cl}(TM) := \mathcal{P}_{\mathcal{O}}(M) \times_{\mathcal{O}_n} \mathcal{Cl}_n$, where $\mathcal{P}_{\mathcal{O}}(M)$ is the principal \mathcal{O}_n -bundle of orthonormal frames on M .

More generally, if $S \rightarrow M$ is any complex vector bundle, we can equip S with a Clifford action $c : \mathcal{Cl}(TM) \rightarrow \text{End}_{\mathbb{C}}(S)$ in a similar way.

Definition 4.16. Let $S \rightarrow M$ be a complex vector bundle with a Hermitian metric $\langle \cdot, \cdot \rangle$ and a connection $\nabla : C^\infty(TM) \otimes C^\infty(S) \rightarrow C^\infty(S)$. This data $(S, \langle \cdot, \cdot \rangle, \nabla)$ defines a *Clifford bundle* if

- for all $X \in C^\infty(TM)$ of unit norm and $s_1, s_2 \in C^\infty(S)$, $\langle X \cdot s_1, X \cdot s_2 \rangle = \langle s_1, s_2 \rangle$ iff $\langle X \cdot s_1, s_2 \rangle + \langle s_1, X \cdot s_2 \rangle = 0$ and

- for all $X, Y \in C^\infty(TM)$ and $s \in C^\infty(S)$, $\nabla_X^L(Y \cdot s) = (\nabla_X Y) \cdot s + Y \cdot \nabla_X s$, where ∇^L denotes the Levi-Civita connection.⁶

Now we need to take a brief detour into something called synchronous frames.

Definition 4.17. Let (M, g) be an n -dimensional Riemannian manifold and $\varphi: U \rightarrow \mathbb{R}^n$ be a chart for M containing some $y \in U$. Let e_1, \dots, e_n be an orthonormal basis of $T_y M$.

TODO: I couldn't figure out what happened here. Sorry. I'll have to fix this later.

Remark 4.18. The exponential map gives a canonical choice for a local neighborhood on a Riemannian manifold. \blacktriangleleft

Dirac operators.

Definition 4.19. Let $S \rightarrow M$ be a Clifford bundle. The Dirac operator $D: C^\infty(S) \rightarrow C^\infty(S)$ is the composition

$$C^\infty(S) \xrightarrow{\nabla} C^\infty(T^*M \otimes S) \xrightarrow{g} C^\infty(TM \otimes S) \xrightarrow{\text{Cl}} C^\infty(S).$$

In a neighborhood of a point $x \in M$, choose a local orthonormal frame e_1, \dots, e_n . Let $e^i := g(e_i, -)$ be the dual frame. Then, the Dirac operator in coordinates looks like

$$(4.20) \quad s \mapsto \nabla_{(\cdot)} s = \sum e^i \otimes \nabla_{e_i} s \mapsto \sum e_i \otimes \nabla_{e_i} s \mapsto \sum_i e_i \cdot \nabla_{e_i} s.$$

Example 4.21. Let \mathbb{R}^n denote Euclidean space, i.e. \mathbb{R}^n with the usual flat metric. If V is a complex vector space, it canonically defines a complex vector bundle $\underline{V} \rightarrow \mathbb{R}^n$ by translation. Let e_1, \dots, e_n be the standard orthonormal frame on $T\mathbb{R}^n$.

Let $\gamma := c(e_i)$, where $c: \text{Cl}(TM) \rightarrow \text{End}_{\mathbb{C}}(\underline{V})$ denotes the Clifford bundle action. Then, the Dirac operator $D: C^\infty(\underline{V}) \rightarrow C^\infty(\underline{V})$ is

$$(4.22) \quad D = \sum_i \gamma_i \cdot \partial_i,$$

where ∂_i is the usual partial derivative operator. The γ_i satisfy the *anticommutation relations*

$$(4.23) \quad \{\gamma_i, \gamma_j\} := \gamma_i \cdot \gamma_j + \gamma_j \cdot \gamma_i = -2\delta_{ij}.$$

Specifically, if $V = \mathbb{C} = \text{Cl}_1$, then $e_1 \mapsto i$, so $\gamma = i$. Therefore $D = i \frac{\partial}{\partial x}$. \blacktriangleleft

The Dirac operator is self-adjoint. Well, it's *formally self-adjoint*: $C^\infty(S)$ is not a Hilbert space, so we can't talk about self-adjointness strictly speaking. One way to abrogate this problem is to take some kind of L^2 completion, but then it's an unbounded operator, so things are still a little complicated. Anyways, we'll talk about this in a bit.

Definition 4.24. Let M be a closed Riemannian manifold and V be its volume form. Then, there is an inner product on $C^\infty(S)$ defined by

$$\langle s_1, s_2 \rangle := \int_M \langle s_1(x), s_2(x) \rangle dV.$$

Theorem 4.25. The Dirac operator on a closed manifold is formally self-adjoint.

Proof. That is, we want to prove that

$$(4.26) \quad \langle Ds_1, s_2 \rangle = \langle s_1, Ds_2 \rangle.$$

We will compute this locally in a synchronous frame e_1, \dots, e_n for a chart in X . Then

$$\begin{aligned} (Ds_1, s_2) - (s_1, Ds_2) &= \sum_i ((e_i \nabla_i s_1, s_2) - (s_1, e_i \nabla_i s_2)) \\ &= \sum_i \left(\nabla_i (e_i \cdot s_1) - \underbrace{(\nabla_i e_i)}_{=0} \cdot s_1, s_2 \right) - (s_1, e_i \nabla_i s_2) \\ &= \sum_i \nabla_i (e_i \cdot s_1, s_2) \\ &= \sum_i \partial_{e_i} (e_i \cdot s_1, s_2) \\ &= d^* \omega, \end{aligned}$$

⁶**TODO:** I am not completely sure I wrote this down correctly.

where

$$\omega_x := - \sum (e_i \cdot s_1, s_2) e^i.$$

Hence the difference is $\langle 1, d^* \omega \rangle = \langle d1, \omega \rangle = 0$. \square

There's a local-vs.-neighborhood argument to make here, but this is the idea.

Lecture 5.

The Weitzenbock formula: 1/31/18

The next talk was by Ivan, on more Clifford bundles and Dirac operators.

Definition 5.1. Let $S \rightarrow M$ be a Clifford bundle on a Riemannian manifold (M, g) , $A \in \Omega_M^2(\text{End } S)$, and $\{e_i\}$ be a local synchronous orthonormal frame for M . The *Clifford contraction* of A is $\mathbb{A} \in \Omega_M^0(\text{End } S)$ defined by the local formula

$$\mathbb{A} \cdot s := \sum_{i < j} c(e_i) c(e_j) A(e_i, e_j) \cdot s.$$

One should check this is independent of the choice of frame, but that is true.

Let \mathbb{K} denote the Clifford contraction of the curvature for ∇ on S .

Theorem 5.2 (Weitzenbock formula). *Let D denote the Dirac operator of $S \rightarrow M$. Then*

$$D^2 = \nabla^* \nabla + \mathbb{K}.$$

$\nabla^* \nabla$ is called the *covariant Laplacian*, and D^2 the *Dirac Laplacian*.

Proof. Let $p \in M$. Then

$$\begin{aligned} D^2 s|_p &= \sum_{i,j} e_i \cdot \nabla_{e_i} (e_j \cdot \nabla_{e_j} s) \Big|_p \\ &= \sum_{i,j} e_i \cdot e_j \cdot \nabla_{e_i} \nabla_{e_j} s|_p. \end{aligned}$$

Splitting this into the cases $i = j$ and $i \neq j$, we get

$$\begin{aligned} &= - \sum_i \nabla_{e_i}^2 s|_p + \sum_{i < j} e_i e_j (\nabla_{e_i} \nabla_{e_j} - \nabla_{e_j} \nabla_{e_i})|_p \\ &= \nabla^* \nabla s|_p + \mathbb{K} s|_p. \end{aligned}$$

This uses the fact that we're on a synchronous frame, so $\nabla_{[e_i, e_j]} = 0$, and therefore the curvature simplifies to $\nabla_{e_i} \nabla_{e_j} - \nabla_{e_j} \nabla_{e_i}$ as we used above. \square

This formula will be crucial for us, allowing us to supplant some of the general theory of elliptic operators in the proof of the index theorem.

Corollary 5.3 (Bochner theorem). *Let $S \rightarrow M$ and \mathbb{K} be as above. If $\langle \mathbb{K} s, s \rangle > 0$ at some point then there are no nontrivial solutions to $D^2 s = 0$.*

Positivity makes sense because \mathbb{K} is a Hermitian operator on a bundle which is fiberwise Hermitian.

Proof. Suppose that $D^2 s = 0$ and $s \neq 0$. Then, $\nabla^* \nabla s + \mathbb{K} s = 0$, so

$$(5.4) \quad 0 = \underbrace{\|\nabla s\|^2}_{\geq 0} + \underbrace{\int_M \langle \mathbb{K} s, s \rangle dV}_{> 0},$$

which is a contradiction. \square

Theorem 5.5. *Let $S \rightarrow M$ be as above, and K denote the curvature of ∇ on S . Then, $K = R^s + F^s$, where*

$$R^s(X, Y) := \frac{1}{4} \sum_{i,j} c(e_i) c(e_j) \langle R(X, Y) e_i, e_j \rangle$$

$$[F^s(X, Y), c(Z)] = 0.$$

R^s is usually called the *Riemann endomorphism*, and only depends on the Riemannian metric of the base manifold. F^s is called the *twisting curvature*.

We'll prove Theorem 5.5 in a series of lemmas. First, we find an obstruction for k being a Clifford module endomorphism.

Lemma 5.6.

$$[K(X, Y), c(Z)] = c(R(X, Y)Z).$$

Proof. Let $\{e_i\}$ denote a synchronous frame at a $p \in M$. Then it suffices to prove the lemma for $X = e_i$, $Y = e_j$, and $Z = e_k$. Since we're in a synchronous frame, the $\nabla_{[e_i, e_j]}$ component of the curvature vanishes, so

$$\begin{aligned} K(e_i, e_j)e_k \cdot s|_p &= \left(\nabla_{e_i} \nabla_{e_j} - \nabla_{e_j} \nabla_{e_i} \right) (e_k \cdot s) \Big|_p \\ &= \left(R(e_i, e_j)e_k \right) \cdot s + e_k \cdot K(e_i, e_j)s \Big|_p. \end{aligned}$$

The result follows because⁷

$$\nabla_{e_i}(e_k)\nabla_{e_j}s|_p = 0. \quad \square$$

Lemma 5.7.

$$[R^s(X, Y), c(Z)] = c(R(X, Y)Z).$$

Proof. Again let $\{e_i\}$ be an orthonormal frame, $X = e_i$, $Y = e_j$, and $Z = e_k$. Then

$$R^s(e_i, e_j)e_k \cdot s = \frac{1}{4} \sum_{\ell, m} c(e_\ell e_m e_k) \underbrace{\langle R(e_i, e_j)e_\ell, e_m \rangle}_{R_{m\ell ij}} s,$$

and similarly

$$c(e_k)R^s(e_i, e_j)s = \frac{1}{4} \sum_{\ell, m} c(e_k e_\ell e_m) R_{m\ell ij} s.$$

Hence when we put these together, we get

$$(5.8) \quad [R^s(e_i, e_j), c(e_k)] = \frac{1}{4} \sum_{\ell, m} c([e_\ell e_m, e_k]) R_{m\ell ij} s.$$

If $\ell = m$ and ℓ, m , and k are distinct, then $[e_\ell e_m, e_k] = 0$, so we only care about the cases $\ell = k \neq m$ and $k = m \neq \ell$.

Both the commutator and $R_{m\ell ij}$ are antisymmetric under the exchange of m and ℓ , so (5.8) reduces to

$$\begin{aligned} [R^s(e_i, e_j), c(e_k)] &= \frac{1}{2} \sum_{\ell} c([e_\ell e_k, e_k]) R_{k\ell ij} s \\ &= \sum_{\ell} c(e_\ell) R_{\ell kij} s \\ &= c \left(\sum_{\ell} R_{\ell kij} s \right). \end{aligned}$$

Since we're working in an orthonormal frame, the metric looks like the identity matrix in coordinates, so

$$\begin{aligned} &= c \left(\sum_{\ell} R_{kij}^\ell s \right) \\ &= c(R(e_i, e_j)e_k)s. \end{aligned} \quad \square$$

These two lemmas suffice to prove Theorem 5.5.

Remark 5.9. Before we go on, let's review Ricci and scalar curvature, which we'll need. Let (M, g) be a Riemannian manifold. Its *Ricci curvature* is the map $\text{Ric}: TX \times TX \rightarrow \mathbb{R}$ defined by

$$\text{Ric}(X, Y) := \text{tr}(Z \mapsto R(Z, X)Y).$$

Why this trace? You could try others, but they all vanish or give you the Ricci curvature up to a sign!

⁷TODO: maybe I missed something.

Raising an index, define $\mathcal{R}ic: TX \rightarrow TX$ by

$$\text{Ric}(X, Y) = g(X, \mathcal{R}ic(Y)).$$

Then, the scalar curvature of (M, g) is $\kappa := \text{tr}(\mathcal{R}ic)$. In an orthonormal frame, it has the formula

$$\kappa = \sum_j \text{Ric}_{jj}. \quad \blacktriangleleft$$

Theorem 5.10 (Improvement on Theorem 5.2). *With notation as in Theorem 5.2, let \mathbb{F}^s denote the Clifford contraction of F^s and κ denote the scalar curvature of (M, g) . Then*

$$D^2 = \nabla^* \nabla + \mathbb{F}^s + \frac{\kappa}{4} \mathbf{1}_{\text{End } S}.$$

Proof. By Theorem 5.5, $\mathbb{K} = \mathbf{R}^s + \mathbb{F}^s$, where \mathbf{R}^s is the Clifford-contracted Riemann endomorphism. So all we have to show is that $\mathbf{R}^s = (\kappa/4) \mathbf{1}_{\text{End } S}$. Again we compute in an orthonormal basis:

$$\begin{aligned} \mathbf{R}^s &= \sum_{i < j} c(e_i) c(e_j) R^s(e_i, e_j) s \\ &= \frac{1}{2} \sum_{i, j} c(e_i e_j) R^s(e_i, e_j) s \\ &= \frac{1}{8} \sum_{i, j, k, \ell} c(e_i e_j e_k e_\ell) \langle R(e_i, e_j) e_k, e_\ell \rangle s \\ &= \frac{1}{8} \sum_{i, j, k, \ell} c(e_i c_j e_k e_\ell) R_{\ell k i j} s. \end{aligned}$$

If you decompose this into parts where various subsets of $\{i, j, k, \ell\}$ are equal to each other, the Bianchi identities allow you to simplify this sum:

$$\begin{aligned} &= \frac{1}{4} \sum_{\ell, i, j} c(e_i c_j e_i e_\ell) R_{\ell i i j} s \\ &= \frac{1}{4} \sum_{\ell, j} c(e_j e_\ell) \left(\underbrace{\sum_i R_{\ell i j}^i}_{\text{Ric}_{\ell j}} \right) \cdot s. \end{aligned}$$

If $\ell \neq j$, the Ricci tensor piece is antisymmetric, so does not contribute to the sum. Hence we only get the case where $k = \ell$:

$$= \frac{1}{4} \sum_j \text{Ric}_{jj} \mathbf{1}_{\text{End } S},$$

and this is indeed the scalar curvature. \square

Now we'll give an example of a Clifford bundle on a non-flat space.

Example 5.11. Let (M, g) be a closed Riemannian manifold and $S := \Lambda^*(T^*M) \otimes \mathbb{C}$. The Riemannian metric on M induces a Riemannian metric on $\Lambda^* T^* M$, hence a Hermitian metric on its complexification; similarly, the Levi-Civita connection induces a connection on $\Lambda^* T^* M$ and therefore also on its complexification. Since the Levi-Civita connection is compatible with the metric on M , our induced connection is compatible with the Hermitian metric on S .

We define the Clifford action $c: TM \rightarrow \text{End } S$ to satisfy

$$(5.12) \quad c(e)^2 = -g(e, e) \mathbf{1}_{\text{End } S},$$

which characterizes it uniquely. Namely, if ω is a k -form and $e \in \Gamma_M(TM)$,

$$c(e) \cdot \omega = \tilde{e} \wedge \omega - e \lrcorner \omega,$$

where the first term is a $(k+1)$ -form and the second is a $(k-1)$ -form. Then

$$\begin{aligned} c(e)^2 \cdot \omega &= \underbrace{\tilde{e} \wedge \tilde{e} \wedge \omega}_{=0} - \tilde{e} \wedge (e \lrcorner \omega) - e \lrcorner (\tilde{e} \wedge \omega) + \underbrace{e \lrcorner (e \lrcorner \omega)}_{=0} \\ &= \underbrace{-e \wedge (e \lrcorner \omega)}_{=0} - \underbrace{(e \lrcorner \tilde{e}) \wedge \omega}_{g(e,e)} + \underbrace{\tilde{e} \wedge (e \lrcorner \omega)}_{=0} \\ &= -g(e, e)\omega. \end{aligned}$$

There are more things to check, including

$$(5.13) \quad g(e \cdot \omega_1, \omega_2) + g(\omega_1, e \cdot \omega_2) = 0,$$

which is left as an exercise, and the fact that

$$(5.14) \quad \nabla_X(e \cdot \omega) = (\nabla_X e) \cdot \omega + e \cdot (\nabla_X \omega).$$

One relatively quick way to prove it is to establish that

$$(5.15) \quad e \lrcorner \omega = (-1)^2 \star(\tilde{e} \wedge \star \omega).$$

This implies

$$\begin{aligned} \nabla_X(e \lrcorner \omega) &= (-1)^2 \star \nabla_X(\tilde{e} \wedge \star \omega) \\ &= (-1)^2 \star (\nabla_X(\tilde{e}) \wedge \star \omega + \tilde{e} \wedge \star \nabla_X \omega) \\ &= (\nabla_X e) \lrcorner \omega + e \lrcorner (\nabla_X \omega). \end{aligned}$$

Hence

$$\begin{aligned} \nabla_X(\tilde{e} \wedge \omega - e \lrcorner \omega) &= (\nabla_X \tilde{e}) \wedge \omega + \tilde{e} \wedge \nabla_X \omega - \nabla_X(e \lrcorner \omega) \\ &= \nabla_X \tilde{e} \wedge \omega + e \wedge \nabla_X \omega - (\nabla_X e) \lrcorner \omega - e \lrcorner \nabla_X \omega \\ &= (\nabla_X e) \cdot \omega + e \cdot (\nabla_X \omega), \end{aligned}$$

proving (5.14). Neither side of (5.14) depends on the metric, and in fact there should be a proof that doesn't use it either.

Now we compute the Dirac operator. Let $\{e_i\}$ be a synchronous frame at p . Then

$$\begin{aligned} D\omega|_p &= \sum_i e_i \cdot \nabla_{e_i} \omega \\ &= \sum_i e^i \wedge \nabla_{e_i} \omega \Big|_p - e_i \lrcorner \nabla_{e_i} \omega \Big|_p \\ &= d\omega|_p - \sum_i (-1)^2 \star(e^i \wedge \nabla_{e_i} \star \omega)|_p \\ &= d\omega|_p + d^* \omega|_p. \end{aligned}$$

This implies $D = d + d^*$, which you can again show in a more abstract way. The Laplacian is

$$\Delta := D^2 = dd^* + d^*d,$$

called the *Hodge Laplacian*. This in particular exists on any Riemannian manifold, without any reference to Clifford bundles. \blacktriangleleft

There are a few more theorems proven in similar ways to the other ones above.

Theorem 5.16. *Restricted to $\Lambda^1 T^*M$, $\Delta = \nabla^* \nabla + \text{Ric}$.*

Theorem 5.17 (Brchner vanishing theorem). *If (M, g) is a Riemannian manifold, $\text{Ric} \geq 0$, and $\text{Ric} > 0$ at some point, then $H_{\text{dR}}^1(M) = 0$.*

This one uses some Hodge theory (e.g. identifying harmonic representatives of de Rham cohomology classes).

Lecture 6.

Spin groups: 2/7/18*“The moment you might begin to think it’s wrong is when you see ‘the proof is obvious.’”*

Today, Sebastian spoke about superalgebra and spin groups.

Superalgebra. Superalgebra, or more generally supermathematics, is the process of adding the prefix “super” to things to denote their $\mathbb{Z}/2$ -graded counterparts with the Koszul sign rule.

Definition 6.1. An algebra A (over \mathbb{R} or \mathbb{C}) is called a *superalgebra* if there is a decomposition as vector spaces $A = A_0 \oplus A_1$ such that $A_i \cdot A_j \subset A_{i+j \bmod 2}$. We call A_0 the *even elements*, A_1 the *odd elements*, and $A_0 \cup A_1$ the *homogeneous elements*.

The degree of a homogeneous element $x \in A_0 \cup A_1$ is denoted $|x| \in \mathbb{Z}/2$.

Remark 6.2. This is equivalent data to an involutive automorphism $\varepsilon: A \rightarrow A$ such that $\varepsilon(a_0 + a_1) = a_0 - a_1$; A_0 and A_1 are the ± 1 -eigenspaces of ε . ◀

Definition 6.3. Let A be a superalgebra.

- The *supercommutator* of $x, y \in A$ is

$$[x, y]_s := xy - (-1)^{|x||y|}yx.$$

- The *super center* of A is

$$Z_s(A) := \{x \in A \mid [x, y]_s = 0 \text{ for all } y \in A\}.$$

Exercise 6.4. Let (V, Q) be a quadratic space. Check that $\mathcal{Cl}(V, Q)$ is a superalgebra with

$$\mathcal{Cl}(V, Q)_0 := \text{span}\{e_{i_1} \cdots e_{i_k} \mid k \text{ even}\}$$

$$\mathcal{Cl}(V, Q)_1 := \text{span}\{e_{i_1} \cdots e_{i_k} \mid k \text{ odd}\}$$

Proposition 6.5. If V is a real inner product space, $Z_s(\mathcal{Cl}(V)) \cong \mathbb{R}$ and $Z_s(\mathcal{Cl}(V) \otimes \mathbb{C}) \cong \mathbb{C}$.

Proof. Suppose $x = a + e_1 b$, where $a = a_0 + \cdots + a_{k-1}$ and $b = b_0 + \cdots + b_{k-1}$. Suppose that $x \in Z_s(\mathcal{Cl}(V))$. Then

$$\begin{aligned} 0 &= [x, e_1]_s \\ &= x e_1 - \sum_{i=1}^{k-1} (-1)^i e_i (a_i - e_1 b_i) \\ &= x e_1 - \sum_{i=0}^{k-1} (-1)^i e_1 a_i - \sum_{i=0}^{k-1} (-1)^i b_i \\ &= \underbrace{\sum_{i=0}^{k-1} (-1)^i e_1 a_i - \sum_{i=0}^{k-1} (-1)^i e_1 a_i}_{=0} + \sum_{i=0}^{k-1} (-1)^{i+1} b_i - \sum_{i=0}^{k-1} (-1)^{i+1} b_i \\ &= -2 \sum_{i=0}^{k-1} (-1)^i b_i. \end{aligned}$$

This forces $b = 0$, so x cannot have any e_1 -component. But e_1 was arbitrary. ◻

The proof boils down to checking minus signs, but is not hard *per se*.

Exercise 6.6. Let V be an oriented inner product space with a positively oriented orthonormal basis $\{e_1, \dots, e_k\}$. Let

$$\omega := e_1 \cdots e_k.$$

- (1) Show that ω does not depend on the choice of basis.
- (2) Show that $\omega^2 = (-1)^{k(k+1)/2}$.
- (3) Show that for $v \in V$, $\omega v = (-1)^{k-1} v \omega$.

Hence if k is odd, $\omega \in Z(\mathcal{Cl}(V))$ (the ordinary center, not the super center). If $k = 2m$, then $\omega^{-1} = (-1)^m \omega$. Hence

$$(6.7) \quad \varepsilon(x) = \omega x \omega^{-1} = -x \omega \omega^{-1}.$$

Hence ε is an inner automorphism.

Pin and spin groups. From now on, fix \mathcal{Cl}_k to be the Clifford algebra of \mathbb{R}^k with its usual inner product, and let $\{e_1, \dots, e_k\}$ be the standard basis for \mathbb{R}^k . Any $v \in \mathbb{R}^k$ satisfies $v \cdot v = -\|v\|^2$, so $V \setminus 0 \subset \mathcal{Cl}_k^\times$.

Definition 6.8.

- The *pin⁻ group* $\text{Pin}_k^- \subset \mathcal{Cl}_k^\times$ is the Lie subgroup generated by the norm-1 elements of \mathbb{R}^k .
- The *spin group* is $\text{Spin}_k := \text{Pin}_k^- \cap (\mathcal{Cl}_k)_0$.

Example 6.9. Let $k = 1$.

- $\mathcal{Cl}(1) = \mathbb{R} \oplus \mathbb{R} \cdot e_1$ where $e_1^2 = -1$, so $\mathcal{Cl}_1 \cong \mathbb{C}$.
- $\text{Pin}_1^- = \{\pm 1, \pm i\} \subset \mathbb{C}^\times$.
- $\text{Spin}_1 = \{\pm 1\} \subset \mathbb{C}^\times$.

◀

We want to study the representation theory of these groups. One way to produce representations would be to find actions of \mathcal{Cl}_k on itself that preserve \mathbb{R}^k .

Recall that $vw + wv = -2\langle v, w \rangle$, and therefore if $\|v\| = 1$, $v^{-1} = -v$. Hence

$$\begin{aligned} -vxv^{-1} &= vxv \\ &= (-xv - 2\langle x, v \rangle)v \\ &= x - 2\langle x, v \rangle v. \end{aligned}$$

Geometrically, this is the reflection of x through the hyperplane $\langle v \rangle^\perp$, as you may remember from linear algebra. This extends to a group representation

$$(6.10) \quad \begin{aligned} \rho : \text{Pin}_k^- &\longrightarrow \text{GL}_k(\mathbb{R}) \\ \rho(y)(x) &:= yx\varepsilon(y^{-1}). \end{aligned}$$

This is called the *twisted adjoint representation*. We will also let ρ denote the restriction of this representation to Spin_k . If $y = u_1 \cdots u_\ell$ for $u_1, \dots, u_\ell \in \mathbb{R}^k$, then

$$yx\varepsilon(y^{-1}) = \pm(u_1 \cdots u_\ell)x(u_\ell^{-1} \cdots u_1^{-1})$$

is also a composition of hyperplane reflections in a similar way as above. This means in particular that if $y \in \text{Spin}_k$, this is a composition of an even number of hyperplane reflections, so $\rho : \text{Spin}_k \rightarrow \text{SO}_k$, not just O_k .

Proposition 6.11. *There is a short exact sequence*

$$(6.12) \quad 1 \longrightarrow \{\pm 1\} \longrightarrow \text{Spin}_k \xrightarrow{\rho} \text{SO}_k \longrightarrow 1.$$

Proof. First, ρ is surjective, because every element of SO_k can be expressed as an even number of hyperplane reflections.

Next, let $y \in \ker(\rho)$, so that $yx\varepsilon(y^{-1}) = x$ for all x . Hence $y \in Z_s(\mathcal{Cl}(k))$, so by Proposition 6.5, $y \in \text{Spin}_k \cap \mathbb{R}$, which is $\{\pm 1\}$. ◻

This implies that Spin_k is a compact simple Lie group, and a double cover of SO_k . We want to know which one it is — in particular, is it a trivial double cover? For $k > 1$, no.

Proposition 6.13. *For $k \geq 2$, Spin_k is connected. For $k \geq 3$, Spin_k is simply connected.*

Proof. The short exact sequence (6.12) is in particular a fiber sequence, hence induces a long exact sequence on homotopy groups.

$$\underbrace{\pi_1(\mathbb{Z}/2)}_1 \longrightarrow \pi_1(\text{Spin}_k) \longrightarrow \pi_1(\text{SO}_k) \longrightarrow \underbrace{\pi_0(\mathbb{Z}/2)}_{\mathbb{Z}/2} \longrightarrow \pi_0(\text{Spin}_k) \longrightarrow \underbrace{\pi_0(\text{SO}_k)}_1$$

The first assertion follows because ± 1 are connected in Spin_k , as $t \mapsto \cos t + e_1 e_2 \sin t$ connects them. Then, the second assertion follows because when $k \geq 3$, $\pi_1(\text{SO}_k) \cong \mathbb{Z}/2$, which can be checked with another long exact sequence of a fibration. \square

Since $\text{Spin}_k \twoheadrightarrow \text{SO}_k$ is a finite cover, it induces an isomorphism of the Lie algebras $\mathfrak{spin}_k \cong \mathfrak{so}_k = \{A \in M_k(\mathbb{R}) \mid A = -A^T\}$.

Lemma 6.14. \mathfrak{spin}_k can be identified with the subspace $\text{span}\{e_i e_j \mid i \neq j\} \subset \mathcal{Cl}_k$ via the map

$$A = (a_{ij}) \mapsto \frac{1}{4} \sum_{i,j} a_{ij} e_i e_j.$$

Representation theory. Let $E_k < \text{Pin}_k^-$ denote the finite subgroup $\{\pm e_1^{i_1} \cdots e_k^{i_k} \mid i_j \in \{0, 1\}\}$. Then $|E_k| = 2^{k+1}$. Let $\nu := -1 \in E_k$.

Proposition 6.15. There's a one-to-one correspondence between the irreducible representations of \mathcal{Cl}_k and the irreducible representations of E_k in which ν acts by -1 .

Remark 6.16. Warning: Roe's book has a stronger version of this theorem, which is **not** true! \triangleleft

Since ν is involutive and central, it acts by ± 1 . If it acts by ν on some representation V , then V is a representation for $E_k/\langle \nu \rangle$. This corresponds to ignoring signs, and in particular is an abelian group of order 2^k . Therefore it has 2^k conjugacy classes, hence 2^k irreducible representations.

Therefore the conjugacy classes in E_k must be either of the form $\{g\}$, where g is central, or $\{g, g\nu\}$. Since there are 2^k conjugacy classes, you can figure out how many of each there have to be.

Lemma 6.17.

- If k is even, $Z(E_k) = \{1, \nu\}$.
- If k is odd, $Z(E_k) = \{1, \nu, \omega, \nu\omega\}$.

Recall that $\omega := e_1 \cdots e_k$. This in particular implies

- If k is even, there are $2^k + 1$ irreducible representations of E_k , hence 1 of \mathcal{Cl}_k .
- If k is odd, there are $2^k + 2$ irreducible representations of E_k , hence 2 of \mathcal{Cl}_k .

For now, assume $k = 2m$ is even, so there's a unique irreducible Δ . Since

$$(6.18) \quad |E_k| = 2^{k+1} = 2^k \cdot 1 + (\dim \Delta)^2,$$

then $\dim \Delta = 2^m$. Explicitly, it arises through a \mathcal{Cl}_k -action on an exterior algebra.

So any finite-dimensional \mathcal{Cl}_k -representation W is of the form $W \otimes_{\mathbb{C}} V$, where V is some “coefficient” vector space, or a bunch of copies of the trivial representation. It's possible to recover V from W : using Schur's lemma,⁸

$$(6.19) \quad V \cong \text{Hom}_{\mathcal{Cl}_k \otimes \mathbb{C}}(\Delta, \Delta) \cong (\Delta^* \otimes_{\mathcal{Cl}_k \otimes \mathbb{C}} \Delta) \otimes_{\mathbb{C}} V \cong \Delta^* \otimes_{\mathcal{Cl}_k \otimes \mathbb{C}} W \cong \text{Hom}_{\mathcal{Cl}_k \otimes \mathbb{C}}(\Delta, W).$$

A similar calculation constructs a natural isomorphism $\beta: \text{End}_{\mathcal{Cl}_k \otimes \mathbb{C}}(W) \xrightarrow{\cong} \text{End}_{\mathbb{C}}(V)$.

Definition 6.20. The *relative trace* $\text{tr}^{W/\Delta} F$ of an $F \in \text{End}_{\mathcal{Cl}(k) \otimes \mathbb{C}}(W)$ is the trace of $\beta(F) \in \text{End}_{\mathbb{C}}(V)$.

Since the elements of Pin_k^- generate \mathcal{Cl}_k , then Δ is also an irreducible representation of Pin_k^- . One can use this to show that Spin_k is an index-2 subgroup of Pin_k^- , hence is automatically normal. Moreover, $\text{Pin}_k^- \twoheadrightarrow \text{O}_k$ is a double cover map.

Since $\text{Spin}_k \triangleleft \text{Pin}_k^-$, then Δ is also a spin representation. Since Spin_k is index 2, then either Δ is irreducible or splits as $\Delta = \Delta^+ \oplus \Delta^-$.

Proposition 6.21. As a Spin_k -representation Δ splits as $\Delta = \Delta^+ \oplus \Delta^-$.

Proof. Recall that $\omega = e_1 \cdots e_{2m}$, so $\omega^2 = (-1)^m$ and $\omega x = \varepsilon(x)\omega$. Clearly $\omega \in \text{Pin}_k^-$, and since k is even, it's also in Spin_k . Suppose $\omega v = \lambda v$. Then

$$(6.22) \quad \omega^2 v = (-1)^m v = \lambda^2 v.$$

Consider $i^m \omega \in \mathfrak{spin}_k \otimes \mathbb{C}$. Then $(i^m \omega)^2 v = v = \lambda^2 v$, so $\lambda \in \{\pm 1\}$. Let Δ_{\pm} denote the ± 1 -eigenspace of $i^m \omega$, so $\Delta = \Delta^+ \oplus \Delta^-$ as vector spaces.

⁸This requires semisimplicity of \mathcal{Cl}_n , which is true but we haven't proven yet.

Let's consider how $\mathcal{Cl}_k \otimes \mathbb{C}$ acts on this splitting. Since

$$(i^m \omega) x v_{\pm} = \varepsilon(x) i^m \omega v_{\pm} = \pm \varepsilon(x) v_{\pm},$$

then if x is even, $x v_{\pm} \in \Delta_{\pm}$, and if x is odd, $x v_{\pm} \in \Delta_{\mp}$. Therefore Spin_k preserves Δ_+ and Δ_- , so they're Spin_k -representations of dimension 2^{m-1} . \square

Δ^+ (resp. Δ^-) is called the *positive* (resp. *negative*) *half-spin representation* of Spin_n . In particular, $\Delta = \Delta^+ \oplus \Delta_-$ is a $\mathbb{Z}/2$ -graded representation of \mathcal{Cl}_{2m} .

Lecture 7.

Spin geometry: 2/7/18

These are Arun's notes for his lecture on spin structures on manifolds. Some errors have been corrected.

Tangential structures. Let $\rho: H \rightarrow G$ be a homomorphism of Lie groups and $\pi: P \rightarrow M$ be a principal G -bundle. Recall that a *reduction of the structure group* of P to H is data $(\pi': Q \rightarrow M, \theta)$ such that

- $\pi': Q \rightarrow M$ is a principal H -bundle, and
- $\theta: Q \times_H G \rightarrow P$ is an isomorphism of principal G -bundles, where H acts on G through ρ .

An equivalence of reductions $(Q_1, \theta_1) \rightarrow (Q_2, \theta_2)$ is a map $\psi: Q_1 \rightarrow Q_2$ intertwining θ_1 and θ_2 .

Definition 7.1. Let M be a smooth n -manifold and $\rho: H \rightarrow \text{GL}_n(\mathbb{R})$ be a homomorphism of Lie groups. If $\mathcal{B}(M) \rightarrow M$ denotes the principal $\text{GL}_n(\mathbb{R})$ -bundle of frames on M , an H -structure on M is an equivalence class of reductions of the structure group of $\mathcal{B}(M)$ to H .

Example 7.2. Let $\rho: \text{O}_n \hookrightarrow \text{GL}_n(\mathbb{R})$ be inclusion. A reduction of the structure group of $\mathcal{B}(M)$ to O_n is a smoothly varying choice of which bases of $T_x M$ are orthonormal, i.e. a smoothly varying inner product on $T_x M$. Hence it's equivalent data to a Riemannian metric. The space of Riemannian metrics on M is connected, which implies that all reductions are equivalent; a manifold has a single O_n -structure. \triangleleft

Example 7.3. Let $\rho: \text{SO}_n \hookrightarrow \text{GL}_n(\mathbb{R})$ be inclusion. In this case, a reduction of the structure group of $\mathcal{B}(M)$ to SO_n specifies which bases of $T_x M$ are oriented at every point, and therefore defines an orientation on M . Two reductions are equivalent iff they define the same orientation. Therefore an SO_n -structure on M is equivalent data to an orientation. \triangleleft

In particular: an H -structure is data, and it need not always exist.

Definition 7.4. A *spin structure* on a manifold M is an H -structure for $H = \text{Spin}_n$ along the map $\rho: \text{Spin}_n \twoheadrightarrow \text{SO}_n \hookrightarrow \text{GL}_n(\mathbb{R})$. A *spin manifold* is a manifold with a specified spin structure.

Example 7.3 immediately implies that a spin structure determines an orientation.

Example 7.5 (Spin structures on the circle). The map $\text{Spin}_1 \rightarrow \text{GL}_1(\mathbb{R})$ factors through $\text{SO}_1 = \{e\}$, so a trivialization of $\mathcal{B}(S^1)$ defines an orientation of S^1 , and a spin structure is a lift of the trivial $\{e\}$ -bundle to a principal $\mathbb{Z}/2$ -bundle. This is the same as a double cover, and there are two isomorphism classes of these.

The connected double cover defines a spin structure which extends over the disc, and is hence called the *bounding spin structure* S_b^1 . In physics, it's also called the *Ramond circle*.

The disconnected double cover defines a spin structure which does not extend over the disc, and is hence called the *nonbounding spin structure* S_{nb}^1 . In physics, this is sometimes called the *Neveu-Schwarz circle*. \triangleleft

The spinor bundle. Recall from Sebastian's talk the spin representation $\rho: \text{Spin}_n \rightarrow \text{GL}(\Delta)$. Throughout this section, (M, g) is a Riemannian manifold with a spin structure implemented by a principal Spin_n -bundle of frames $\mathcal{B}_{\text{Spin}}(M) \rightarrow M$.

Definition 7.6. The *spin bundle* or *spinor bundle* of M is $\mathcal{B}_{\text{Spin}}(M) \times_{\text{Spin}_n} \Delta \rightarrow M$, which is a complex vector bundle of rank n .

The spinor bundle has a lot of structure.

- The spin representation is $\mathbb{Z}/2$ -graded, and therefore the spinor bundle is as well: $S = S^0 \oplus S^1$.
- The spinor bundle has a canonical connection D on it: the Levi-Civita connection lifts from $\mathcal{B}(M)$ to $\mathcal{B}_{\text{Spin}}(M)$, and therefore passes to the associated bundle.

- The spin representation factors through the unitary group $U(\Delta) \hookrightarrow GL(\Delta)$, and therefore there is an induced Hermitian metric on S .

Proposition 7.7. *The spin connection D is compatible with h , and therefore the spinor bundle is a Clifford bundle.*

Proof. If $(E, \nabla, g) \rightarrow M$ is a real vector bundle with connection and metric, compatibility with the metric means that the connection one-forms are valued in $\mathfrak{o}(E, g)$; if E is complex with a Hermitian metric h , compatibility with the metric means the connection one-forms are valued in $\mathfrak{u}(E, g)$. So this is a fun exercise in what “induced connection” actually means.

Since the Levi-Civita connection is metric-compatible, its connection one-form lives in $\Omega^1_{\mathcal{B}(M)}(\mathfrak{o}_n)$. The pullback connection on $\mathcal{B}_{\text{Spin}}(M)$ has connection one-form in $\Omega^1_{\mathcal{B}_{\text{Spin}}(M)}(\mathfrak{spin}_n)$. Since the spinor representation factors through $U(\Delta) \hookrightarrow GL(\Delta)$, the connection on the associated bundle factors through $\mathfrak{u}(\Delta) \hookrightarrow \mathfrak{gl}(\Delta)$; in particular, its connection one-form lives in $\Omega_S(\mathfrak{u}(\Delta))$, so D is compatible with the metric. \square

Proposition 7.8. *The twisting curvature of the spinor bundle is zero.*

Proof. Recall from Sebastian’s talk that there’s a natural identification of \mathfrak{spin}_n with the vector subspace of $\mathcal{C}\ell_n$ spanned by products $e_i e_j$ for $i \neq j$ by the assignment

$$(7.9) \quad a_{ij} \mapsto \frac{1}{4} \sum a_{ij} e_i e_j.$$

Let $\{e_k\}$ be a local orthonormal frame for TM . Therefore the curvature of the Levi-Civita connection is an \mathfrak{so}_n -valued 2-form whose (k, ℓ) entry is $\langle Re_k, e_\ell \rangle$, where R is the Riemann curvature tensor. By (7.9), the curvature 2-form of the spin connection is

$$(7.10) \quad K_{\text{Spin}} = \frac{1}{4} \sum_{k, \ell} \langle Re_k, e_\ell \rangle e_k e_\ell.$$

This acts on the spinor representation through $e_i \mapsto c(e_i)$, i.e. through

$$(7.11) \quad \frac{1}{4} \sum_{k, \ell} \langle Re_k, e_\ell \rangle c(e_k) c(e_\ell).$$

With notation as last time, this is exactly $R^s(e_j, e_i)$, so since the twisting curvature F^s satisfies $F^s = K_{\text{Spin}} - R^s$, $F^s = 0$. \square

Spin^c structures.

Definition 7.12. The Lie group Spin_n^c is defined to be $\text{Spin}_n \times_{\mathbb{Z}/2} U_1$, where $\mathbb{Z}/2 \hookrightarrow \text{Spin}_n$ as $\{\pm 1\} \subset \mathcal{C}\ell_n$ and $\mathbb{Z}/2 \hookrightarrow U_1$ as $\{\pm 1\} \subset \mathbb{C}$.

This means in particular there is a short exact sequence

$$(7.13) \quad 1 \longrightarrow \mathbb{Z}/2 \longrightarrow \text{Spin}_n^c \longrightarrow \text{SO}_n \times U_1 \longrightarrow 1.$$

A *spin^c structure* on a manifold M is an H -structure for $H = \text{Spin}_n^c$ along the map to SO_n . This comes with some extra structure.

- The map $\text{Spin}_n^c \rightarrow \text{SO}_n \times U_1$ means that a Spin_n^c -structure on M determines an orientation and a principal U_1 -bundle $P \rightarrow M$. The associated complex line bundle $L := P \times_{U_1} \mathbb{C} \rightarrow M$ is called the *fundamental line bundle* associated to the spin^c structure; its Chern class is called the *Chern class* of the spin^c structure.
- Since the spin representation is a representation of the complexified Clifford algebra, it also induces a representation of Spin_n^c . In the same manner as above, a spin^c structure on M defines an associated real vector bundle $E^c \rightarrow M$, again called the *spinor bundle* S^c . Since the spin representation is unitary, the spinor bundle acquires a Hermitian metric.

However, we don’t get a connection for free: instead we have to choose a connection ∇_L on L ; together with the Levi-Civita connection on $\mathcal{B}_{\text{SO}}(M)$, this induces a connection ∇^c on $\mathcal{B}_{\text{Spin}_n^c}(M) \rightarrow M$.

Proposition 7.14. *In this situation, the twisting curvature of ∇^c is $(1/2)K_L$, where K_L is the curvature of ∇_L .*

Characteristic classes. We can use characteristic classes to determine whether a manifold has a spin or spin^c structure.

First we mention a general form of the *splitting principle*, a way of making computations about characteristic classes for principal G -bundles.

Let G be a compact Lie group, T be a maximal torus for G , and $P \rightarrow X$ be a principal G -bundle over a paracompact base X . The quotient map induces a map $f : P/T \rightarrow X$, and we can pull P back along f :

$$(7.15) \quad \begin{array}{ccc} f^*P & \longrightarrow & P \\ \downarrow G & & \downarrow G \\ P/T & \xrightarrow{f} & X. \end{array}$$

Then, there is a canonical reduction of the structure group of f^*P along $T \hookrightarrow G$. In general, T splits as a product of copies of S^1 , so f^*P similarly splits. The key fact is that the map $f^* : H^*(X) \rightarrow H^*(P/T)$ is injective. Therefore any question about characteristic classes of P may be solved for f^*P , which splits, with no information lost.

Definition 7.16. If E is a real vector bundle, $\mathfrak{S}(E)$ will denote the Pontrjagin genus of E associated to

$$g(z) = \cosh\left(\frac{1}{2}\sqrt{z}\right).$$

Proposition 7.17. The Chern character of the spinor bundle on an even-dimensional spin manifold M is

$$\text{ch}(S) = 2^m \mathfrak{S}(TM).$$

Proof. We will prove a stronger result. A *spin structure* on a real vector bundle $E \rightarrow M$ is an (equivalence class of) reduction of the structure group of the principal $\text{GL}_n(\mathbb{R})$ -bundle of local frames of E across $\text{Spin}_n \rightarrow \text{GL}_n(\mathbb{R})$. Given a spin structure on E , one can define its spinor bundle S_E in the same manner as before. We will show that if E has even rank, then

$$(7.18) \quad \text{ch}(S_E) = 2^{\dim E/2} \mathfrak{S}(E).$$

If $E_1, E_2 \rightarrow X$ are even-dimensional spin vector bundles, $S_{E_1 \oplus E_2} \cong S_{E_1} \otimes S_{E_2}$, both sides of (7.18) are multiplicative under direct sum. Therefore we may apply the splitting principle for principal Spin_n -bundles to E . One choice of maximal torus in Spin_n is

- $\text{Spin}_2 \times \cdots \times \text{Spin}_2 \hookrightarrow \text{Spin}_n$ if n is even, or
- $\text{Spin}_2 \times \cdots \times \text{Spin}_2 \times \{1\} \hookrightarrow \text{Spin}_n$ if n is odd.

Thus over E/T , $\mathcal{B}_{\text{Spin}}(E)$ splits as a direct sum of principal Spin_2 -bundles, so E splits as a sum of rank-2 vector bundles with spin structure. Since (7.18) is additive, it suffices to assume $\text{rank } E = 2$.

Recall that a *classifying space* BG for a Lie group G is a space with a principal G -bundle $EG \rightarrow BG$ whose total space is contractible. This determines BG up to homotopy, and also implies that every principal G -bundle $P \rightarrow X$ is the pullback of EG along a map $X \rightarrow BG$. Since characteristic classes are natural under pullback, verifying (7.18) over $B\text{Spin}_2$ will prove it everywhere.

Given a representation V of G , we get a vector bundle $EG \times_G V \rightarrow BG$. The universal choice of $E \rightarrow \text{Spin}_2$ is the one induced from the representation $\text{Spin}_2 \rightarrow \text{SO}_2 \hookrightarrow \text{GL}_2(\mathbb{R})$, and its spinor bundle is the complex vector bundle $\Delta \rightarrow B\text{Spin}_2$ associated to the spinor representation.

Since the map $\text{Spin}_2 \rightarrow \text{SO}_2$ is isomorphic to the map $\mathbb{T} \rightarrow \mathbb{T}$ which is multiplication by 2, the map $H^*(BSO_2; \mathbb{Q}) \rightarrow H^*(B\text{Spin}_2; \mathbb{Q})$ is also multiplication by 2. Therefore if $x \in H^1(BSO_2; \mathbb{Q})$ denotes the generator, which is the Euler class of the defining representation, $H^*(BSO_2; \mathbb{Q}) \cong \mathbb{Q}[x]$ and $H^*(B\text{Spin}_2; \mathbb{Q}) \cong \mathbb{Q}[x/2]$. Since $E \rightarrow \text{Spin}_2$ comes from the defining representation of SO_2 , $c_1(E_{\mathbb{C}}) = x$.

Next we compute $\text{ch}(\Delta)$. This spinor representation is $\text{Spin}_2 \hookrightarrow \mathcal{C}\ell_2 \otimes \mathbb{C} \cong M_2(\mathbb{C}) \cong \text{End}(\mathbb{C}^{1|1})$, which acts on $\mathbb{C}^{1|1}$ and therefore induces a representation of the spin group. Since $\text{Spin}_2 \cong \text{U}_1$ abstractly, and irreducible representations of U_1 are $z \mapsto z^n$ for an $n \in \mathbb{Z}$, we will decompose Δ into characters and compare them to E .

The identification $\mathcal{C}\ell_2 \otimes \mathbb{C} \cong M_2(\mathbb{C})$ sends

$$(7.19a) \quad e_1 \mapsto \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

$$(7.19b) \quad e_2 \mapsto \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix},$$

and therefore

$$(7.19c) \quad e_1 e_2 \mapsto \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

This is the direct sum of Δ^+ (first entry) and Δ^- (second entry). In particular, the Δ^+ is the character $z \mapsto z$, and Δ^- is the character $z \mapsto z^{-1}$. Via the identification $H^*(B\text{Spin}_2; \mathbb{Q}) \cong \mathbb{Q}[x/2]$, $c_1(\Delta^\pm) = \pm x/2$, so $\text{ch}(\Delta^\pm) = e^{\pm x/2}$. Therefore

$$(7.20) \quad \text{ch}(\Delta) = \text{ch}(\Delta^+) - \text{ch}(\Delta^-) = e^{x/2} - e^{-x/2} = 2 \cosh\left(\frac{x}{2}\right).$$

We also have $p_1(E) = e(E)^2 = x^2$, so

$$(7.21) \quad \mathfrak{S}(E) = \cosh\left(\frac{1}{2}\sqrt{p_1(E)}\right) = \cosh\left(\frac{x}{2}\right). \quad \boxtimes$$

Definition 7.22. Let BSO_n denote the classifying space of SO_n ; there's a natural bijection between isomorphism classes of principal SO_n -bundles $P \rightarrow X$ and homotopy classes of maps $X \rightarrow BSO_n$. One can show that when $n \geq 2$, $H^2(BSO_n; \mathbb{Z}/2) \cong \mathbb{Z}/2$, with generator denoted w_2 . Given a principal SO_n -bundle $P \rightarrow X$, the pullback of w_2 by the classifying map $X \rightarrow BSO_n$ is called the *second Stiefel-Whitney class* of P , denoted $w_2(P) \in H^2(X; \mathbb{Z}/2)$.

Proposition 7.23. Let M be a closed, oriented manifold. Then, M has a spin structure iff $w_2(M) = 0$. In this case, the spin structures of M extending its given orientation are a torsor over $H^1(M; \mathbb{Z}/2)$.

Proof. We follow the proof in Roe's exercises. We'll assume $n \geq 2$; otherwise small modifications must be made, but the result is still true. Let M be a closed, oriented, manifold and $\pi: \mathcal{B}_{SO}(M) \rightarrow M$ be its principal SO_n -bundle of frames. This is a fiber bundle with fiber SO_n , hence has an associated spectral sequence

$$(7.24) \quad E_2^{p,q} = H^p(M; H^q(SO_n; \mathbb{Z}/2)) \implies H^{p+q}(\mathcal{B}_{SO}(M); \mathbb{Z}/2).$$

Associated a first-quadrant spectral sequence converging to $H^n(A)$ is its *five-term exact sequence*

$$(7.25) \quad 0 \longrightarrow E_2^{1,0} \longrightarrow H^1(A) \longrightarrow E_2^{0,1} \longrightarrow E_2^{2,0} \longrightarrow H^2(A),$$

which in this case specializes to

$$(7.26) \quad 0 \longrightarrow H^1(M; \mathbb{Z}/2) \xrightarrow{\pi^*} H^1(\mathcal{B}_{SO}(M); \mathbb{Z}/2) \xrightarrow{i^*} H^1(SO_n; \mathbb{Z}/2) \xrightarrow{\delta} H^2(M; \mathbb{Z}/2) \longrightarrow H^2(\mathcal{B}_{SO}(M); \mathbb{Z}/2).$$

We claim, and will prove in a bit, that δ sends the generator of $H^1(SO_n; \mathbb{Z}/2)$ to $w_2(M)$.

A spin structure lifting the specified orientation of M is the data of an isomorphism class of double covers $\mathcal{B}_{\text{Spin}}(M) \rightarrow \mathcal{B}_{SO}(M)$ which induce the connected double cover $\text{Spin}_n \rightarrow SO_n$ when restricted to each fiber. Isomorphism classes of covers of $\mathcal{B}_{SO}(M)$ are naturally identified with $H^1(\mathcal{B}_{SO}(M); \mathbb{Z}/2)$, and their restriction to a fiber is the map i^* in (7.26), so the set of spin structures on M is in bijection with the subset of $H^1(\mathcal{B}_{SO}(M); \mathbb{Z}/2)$ which is *not* in $\ker(i^*)$.

Suppose that $w_2(M)$ vanishes, meaning $\delta = 0$, so i^* is a surjective map onto $\mathbb{Z}/2$. Therefore $|(i^*)^{-1}(0)| = |(i^*)^{-1}(1)|$ and π^* is an injection whose image is $(i^*)^{-1}(0)$. Thus we have a noncanonical bijection $H^1(M; \mathbb{Z}/2) \rightarrow (i^*)^{-1}(1)$, the set of spin structures on M . The torsor structure arises because $(i^*)^{-1}(0) + (i^*)^{-1}(1) \subset (i^*)^{-1}(1)$.

If instead $w_2(M) \neq 0$, then δ is injective. Then $i^* = 0$, so there can be no spin structures on M .

To finish this, we had better check that δ sends the generator of $H^1(SO_n; \mathbb{Z}/2)$ to w_2 . It suffices to check this for the universal bundle $ESO_n \rightarrow BSO_n$, as all other principal SO_n -bundles are pullbacks of this one, but since ESO_n is contractible, (7.26) simplifies to

$$(7.27) \quad 0 \longrightarrow H^1(BSO_n; \mathbb{Z}/2) \longrightarrow 0 \longrightarrow H^1(SO_n; \mathbb{Z}/2) \xrightarrow{\delta} H^2(BSO_n; \mathbb{Z}/2) \longrightarrow 0.$$

so δ has to send the nonzero element to the generator of $H^2(BSO_n; \mathbb{Z}/2)$, which is w_2 , as desired. \boxtimes

Proposition 7.28. Let M be a closed manifold. Then, M has a spin^c structure iff $w_2(M)$ is the reduction of an integral class. In this case, the spin^c structures of M are a torsor over $H^2(M; \mathbb{Z})$.

Corollary 7.29. A spin structure determines a spin^c structure; the converse is not true: \mathbb{CP}^2 is a spin^c manifold which is not spin.

Partial proof. We'll prove the last part, by showing $w_2(\mathbb{CP}^2) \neq 0$. It's equal to the mod 2 reduction of $c_1(\mathbb{CP}^2)$, and $c_1(\mathbb{CP}^2) = (1+x)^3$, where x is the generator dual to a \mathbb{CP}^1 , so $c_1(\mathbb{CP}^2)$ is odd and therefore $w_2(\mathbb{CP}^2) \neq 0$, so \mathbb{CP}^2 is not spin, but it is spin^c . \square

Lecture 8.

Analytic properties of the Dirac operator: 2/14/18

Today, Gill and Kenny spoke about analytic questions related to the Dirac operator $\not{D}: C^\infty(S) \rightarrow C^\infty(S)$, where $S \rightarrow M$ is the spinor bundle on a compact spin Riemannian manifold M .

Since $C^\infty(S)$ is in general an infinite-dimensional vector space, we would like to place a good topology on it. The default is the Fréchet topology, but this is kind of unpleasant to work with: it's not a Hilbert space, for example. Maybe we can relax the regularity of things we feed to \not{D} , but since we're interested in the Dirac Laplacian \not{D}^2 , which is a second-order differential operator, we want everything to be at least twice differentiable.

Relatedly, one asks how smooth solutions to $\not{D}x = y$ can be. For example, in \mathbb{R}^2 , harmonic functions, those functions $u: \mathbb{R}^2 \rightarrow \mathbb{C}$ satisfying

$$(8.1) \quad \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0,$$

are smooth and in fact analytic, because $f := u + i\bar{u}$ is holomorphic. This is an instance of *elliptic regularity*, a very general phenomenon.

Thus the general place to start would be to use the induced Hermitian metric h on S and take the space $L^2(S)$, the Hilbert space of L^2 sections of $S \rightarrow M$ with respect to h . Since smooth sections are dense in $L^2(S)$, we know how to make sense of $\not{D}: L^2(S) \rightarrow L^2(S)$.

Exercise 8.2. Show that $\not{D}: L^2(S) \rightarrow L^2(S)$ is unbounded, i.e. there is no $C \in \mathbb{R}$ such that for all $X \in L^2(S)$, $\|\not{D}X\|_{L^2} \leq C\|X\|_{L^2}$.

This is a little unfortunate; we would like something whose functional analysis is more tractable. For this reason, we'll change the Hilbert space we're working with slightly.

Recall that if $f, g \in C^\infty(S)$, we defined

$$(8.3) \quad \langle f, g \rangle_{L^2} := \int_M h(f(x), g(x)) dV,$$

where dV is the volume form on M . Then $L^2(S)$ is defined to be the completion of $\{f \in C^\infty(S) \mid \langle f, f \rangle < \infty\}$ in the norm induced by this inner product. Therefore by changing the inner product we can change what Hilbert space we're using.

Definition 8.4. The k^{th} Sobolev inner product of smooth sections f and g is

$$\langle f, g \rangle_{H^k} := \int_M \sum_{i=0}^k h(\nabla^i f, \nabla^i g) dV,$$

where $\nabla^0 := \text{id}$, and h denotes both the Hermitian metric on S and the induced metrics on $(T^*M)^{\otimes i} \otimes S$ using h and the Riemannian metric on T^*M defined by taking the product of all the metrics.

The completion of $C^\infty(S)$ under $\langle -, - \rangle_{H^k}$ is denoted $H^k(S)$ and called a *Sobolev space*.

Proposition 8.5. Any choices of connection and metric on M induce equivalent Sobolev k -norms.

That is, the underlying topological vector spaces of these Hilbert spaces are isomorphic.

Let T^n be the n -torus. It suffices to understand the proposition on all charts U_α of M for an atlas which trivializes S , and we can map these charts onto T^n , thus reducing the question to T^n .

On T^n we use Fourier theory: for any $L^2 f: T^n \rightarrow \mathbb{C}$, its *Fourier series* is

$$(8.6) \quad f(x) = \sum_{\nu \in \mathbb{Z}^n} \hat{f}(\nu) e^{i\nu x},$$

where

$$(8.7) \quad \hat{f}(\nu) := \frac{1}{(2\pi)^n} \int_{T^n} f(x) e^{-i\nu x} dx.$$

It's a standard result that sending $f \mapsto \widehat{f}$ converts differentiation into multiplication: if f is C^1 , then

$$(8.8) \quad \left(\frac{\partial}{\partial x_j} \right)^\wedge (\nu) = \frac{1}{(2\pi)^n} \int_{T^n} \left(\frac{\partial}{\partial x_j} f \right)(x) e^{-i\nu x} dx.$$

Hence for an arbitrary L^2 f , we can define $\partial_j f$ through its Fourier series:

$$(8.9) \quad \frac{\partial f}{\partial x_j} := \sum_{\nu \in \mathbb{Z}^n} i \nu_j \widehat{f}(\nu) e^{i\nu x}.$$

Definition 8.10. Let $H^k(T^n)$, $k \in \mathbb{R}$. The k^{th} Sobolev space on T^n , denoted $H^k(T^n)$, is the completion of the subspace of $L^2(T^n)$ for which the norm

$$\|f\|_k^2 := \sum_{\nu \in \mathbb{Z}^n} |\widehat{f}(\nu)|^2 (1 + |\nu|^2)^k$$

is finite, in this norm. Here $|\nu| = \sum \nu_i^2$.

For example,

$$(8.11) \quad \|f\|_1^2 = \|f\|_{L^2}^2 + \sum_j \left\| \frac{\partial f}{\partial x_j} \right\|_{L^2}^2,$$

which uses the fact that the Fourier series functional is an isometry.

Proposition 8.12. *There exist constants $c_1, c_2 \in \mathbb{R}$ such that*

$$c_1 \sum_{|\alpha| \leq k} \left\| \frac{\partial f}{\partial x_\alpha} \right\|_{L^2} \leq \|f\|_k \leq c_2 \sum_{|\alpha| \leq k} \left\| \frac{\partial f}{\partial x_\alpha} \right\|_{L^2}.$$

This implies in particular that if D is a linear ℓ^{th} -order differential operator, then it's a linear map $D: H^k \rightarrow H^{k-\ell}$.

Proposition 8.13. *Let $k_1 \leq k_2$. There are continuous embeddings*

- $H^{k_1}(S) \hookrightarrow H^{k_2}(S)$ and
- $C^k(T^n) \hookrightarrow H^k(T^n)$.

Most of these propositions are true in general Sobolev spaces, since they're all proven locally.

Theorem 8.14 (Sobolev embedding theorem). *If $p > n/2$, $H^{k+p} \hookrightarrow C^k$ continuously.*

Recall that a map of topological vector spaces is *compact* if it carries closures of bounded spaces to compact spaces. (By linearity, it suffices to consider the closed unit ball.)

Theorem 8.15 (Rellich). *If $k_1 \leq k_2$, the embedding $H^{k_1} \hookrightarrow H^{k_2}$ is compact.*

~ ~ ~

With these preliminaries out of the way, we proceed to the actual analysis of the Dirac operator. The estimates we discuss apply more generally, to any first-order operator D on $C^\infty(S)$ such that

$$(8.16) \quad D^2 = \nabla^* \nabla + B$$

for some first-order operator B on S . Since M is compact, $C^\infty(S) \hookrightarrow L^2(S)$, so D is a bounded map $H^1(S) \rightarrow L^2(S)$, i.e. there's a $C \in \mathbb{R}$ such that for all $s \in C^\infty(S)$, $\|Ds\|_0 \leq C\|s\|_1$. We will call such a D a *generalized Dirac operator*.

Proposition 8.17 (Gårding inequality). *Let D be a generalized Dirac operator. Then there is a C such that for any $s \in C^\infty(S)$,*

$$\|s\|_1 \leq C(\|s\|_0 + \|Ds\|_0).$$

This is an analogue of the smoothness of harmonic functions we discussed above, buying us unexpected extra regularity.

Proof. The proof follows fairly standard methods in PDE. It suffices to assume that S is the trivial bundle, because we can work locally: let \mathfrak{U} be an open cover of M which trivializes S and $\{\psi_U : U \in \mathfrak{U}\}$ be a partition of unity subordinate to it. Then

$$(8.18) \quad \|s\|_0^2 = \sum_{U \in \mathfrak{U}} \int_U \langle \sqrt{\psi_U} s, \sqrt{\psi_U} s \rangle dx.$$

Since M is compact we can take \mathfrak{U} to be finite, and therefore take the maximum of the constants we got on each $U \in \mathfrak{U}$.

Since

$$(8.19) \quad \langle D^2 s, s \rangle = \langle \nabla^* \nabla s, s \rangle + \langle B s, s \rangle,$$

then

$$(8.20) \quad \|D s\|_0^2 = \|\nabla s\|_0^2 + \langle B s, s \rangle,$$

so by the Cauchy-Schwarz inequality, and the fact that B is bounded in H^1 ,

$$(8.21) \quad \|\nabla s\|_0^2 \leq c_1 (\|s\|_0 \|s\|_1 + \|D s\|_0^2).$$

We will bound the left side of (8.21) below. Since we're working locally, $\nabla_i s = \partial_i s + T_i s$, where $\nabla s := dx^i \otimes \nabla_i s$, and therefore

$$\begin{aligned} \langle \nabla s, \nabla s \rangle &= \left\langle \sum_i dx^i \otimes \nabla_i s, \sum_j dx^j \otimes \nabla_j s \right\rangle \\ &= \sum_{i,j} \underbrace{\langle dx^i, dx^j \rangle}_{g^{ij}} \langle \nabla_i s, \nabla_j s \rangle \\ &= \int \sum_{i,j} g^{ij} (\partial_i s, \partial_j s) + 2 \operatorname{Re}(\partial_i s, T_j s) + (T_i s, T_j s) dx. \end{aligned}$$

Since

$$(8.22) \quad \|D s\|_0^2 := \int \langle \nabla s, \nabla s \rangle dx \geq C_2 \|s\|_1^2 - C_3 \|s\|_0 \|s\|_1$$

(the reasoning for which is a little unclear), then

$$(8.23) \quad C_2 \|s\|_1^2 - C_3 \|s\|_0 \|s\|_1 \leq C_1 (\|s\|_0 \|s\|_1 + \|D s\|_0^2),$$

so there are constants C_4, C_5 such that

$$(8.24) \quad \|D s\|_0^2 \geq C_4 \|s\|_1^2 - C_5 \|s\|_0 \|s\|_1.$$

It's a general fact (proven maybe by a convexity argument?) that for any $\varepsilon > 0$, there's a $K > 0$ such that for all $a, b > 0$,

$$(8.25) \quad ab \leq \varepsilon a^2 + K b^2.$$

Therefore there's a C_6 such that

$$(8.26) \quad C_5 \|s\|_0 \|s\|_1 \leq \frac{1}{2} C_4 \|s\|_1^2 + C_6 \|s\|_0^2,$$

so

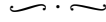
$$(8.27) \quad \|D s\|_0^2 \geq \frac{1}{2} C_4 \|s\|_1^2 - C_6 \|s\|_0^2,$$

which can be rearranged into what we wanted to show. \square

The Gårding inequality says that we can control all of the first derivatives of s if we can control $D s$ (we need the extra term in case D has kernel – if the first term weren't there, this would've been a proof that D were injective). Something elliptic is going on here, and we can bootstrap it into stronger estimates.

Theorem 8.28 (Elliptic estimate). *For any $k > 0$, there's a constant C_k such that for all $s \in C^\infty(S)$,*

$$\|s\|_{k+1} \leq C_k (\|s\|_k + \|D s\|_k).$$



We now turn to spectral theory. This is particularly nice for Hilbert spaces, so when one wants to study the spectral theory of an operator, part of the story is to find a Hilbert space on which its spectrum is nice.

Definition 8.29. By an *unbounded operator* on a Hilbert space H we mean a linear operator $f : U \rightarrow H$, where U is a dense subset of H .

We make no assumptions on continuity; hence unbounded operators are a (in general strict) superset of bounded operators.

Example 8.30. We defined $L^2(S)$ to be the completion of $C^\infty(S)$ under the L^2 norm, and hence $D : C^\infty(S) \rightarrow L^2(S)$ is an unbounded operator. \triangleleft

Proposition 8.31. Let D be a generalized Dirac operator. Then the closure of the graph of D is the also graph of an unbounded operator \overline{D} .

Proof. Recall that D has a formal adjoint, i.e. an operator D^\dagger such that $\langle Dx, y \rangle = \langle x, D^\dagger y \rangle$. Let Γ denote the graph of D and suppose that $\overline{\Gamma}$ isn't the graph of a function; then, $\overline{\Gamma}$ fails the “vertical line test,” so there are $x, y_1, y_2 \in H$ such that $(x, y_1), (x, y_2) \in \overline{\Gamma}$. Since D is linear, then $(0, y_1 - y_2) \in \overline{\Gamma}$, and $y_1 - y_2 \neq 0$.

Hence there are $x_i \in C^\infty(S)$ such that $x_i \rightarrow 0$ and $Dx_i \rightarrow z \neq 0$ in L^2 . Hence for all $s \in C^\infty(S)$, $\langle Dx_i, s \rangle \rightarrow \langle z, s \rangle$ and $\langle x_i, D^\dagger s \rangle \rightarrow 0$, but since $\langle Dx_i, s \rangle = \langle x_i, D^\dagger s \rangle$, $\langle z, s \rangle = 0$ for all s . Hence $z = 0$.

From this and the fact that $\overline{\Gamma}$ is a linear subspace of $H \oplus H$, one can infer that it's the graph of a *linear operator*. \square

What's the domain of \overline{D} ? It's all $x \in L^2(S)$ for which there exists a sequence $x_i \rightarrow x$ converging in L^2 such that $Dx_i \rightarrow \overline{D}x$ in L^2 . Hence by Proposition 8.17, the domain contains $H^1(S)$, and one can show that the domain is $H^1(S)$.

Definition 8.32. Let $x, y \in L^2(S)$. We say $Dx = y$ *weakly* if for all $s \in C^\infty(S)$, $\langle x, D^\dagger s \rangle = \langle y, s \rangle$ (equivalently, $\langle Dx, s \rangle = \langle y, s \rangle$).

Denote by $S \boxtimes S^* \rightarrow M \times M$ the bundle $\pi_1^* S \otimes \pi_2^* S^*$, where $\pi_i : M \times M \rightarrow M$ is projection onto the i^{th} factor.

Definition 8.33. Let $A : H \rightarrow H$ be a bounded operator. Then, A is a *smoothing operator* if there is a smooth section $k(p, q)$ of $S \boxtimes S^*$, called the *kernel* of A , such that

$$(8.34) \quad A(s)(p) = \int_M k(p, q)s(q) dq.$$

The assertion that the integral in (8.34) converges depends both on the compactness of M and a nontrivial element.

Definition 8.35. A *Friedrichs' mollifier* for a section s is a family F_ε of self-adjoint smoothing operators on $L^2(S)$, indexed by $\varepsilon \in (0, 1)$, such that

- (1) $\{F_\varepsilon\}$ is bounded in $B(L^2(S))$,
- (2) for any first-order differential operator B , $\{[B, F_\varepsilon]\}$ extends to a bounded family on $B(L^2(S))$, and
- (3) $F_\varepsilon \rightarrow 1$ as $\varepsilon \rightarrow 0$ (i.e. in the weak topology on $B(L^2(S))$).

Exercise 8.36. Friedrichs' mollifiers exist.

Theorem 8.37. Suppose that $x, y \in L^2(X)$ and $Dx = y$ weakly. Then $x \in H^1$ and $\overline{D}x = y$.

Proof. Let $\{F_\varepsilon\}$ be a Friedrichs' mollifier and $x_\varepsilon := F_\varepsilon x$. Then for any $s \in C^\infty(S)$,

$$\begin{aligned} \langle Dx_\varepsilon, s \rangle &= \langle x_\varepsilon, D^\dagger s \rangle \\ &= \langle x, F_\varepsilon D^\dagger s \rangle \\ &= \langle x, D^\dagger F_\varepsilon s \rangle + \underbrace{\langle x, [F_\varepsilon, D^\dagger] s \rangle}_{\leq \|x\|_0 \| [F_\varepsilon, D^\dagger] \|_{B(L^2(S))} \|s\|_0} \\ &= \langle y, F_\varepsilon s \rangle \leq C' \|y\| \|F_\varepsilon\|_B \|s\|_0 \\ &\leq K \|s\|_0 \end{aligned}$$

for some K . Using the first two properties of Friedrichs' mollifiers, $|\langle Dx_\varepsilon, s \rangle| \leq C' \|s\|_0$; since $\|Dx_\varepsilon\|_0 \leq C$, then by Proposition 8.17, $\|x_\varepsilon\|_{H^1} \leq C'$.

This implies there exists a subsequence $\varepsilon_j \rightarrow 0$ such that $x_{\varepsilon_j} \rightarrow z$ in H^1 .

Exercise 8.38. Show that compact operators send weakly convergent sequences to strongly convergent sequences.

Thus we can conclude that $x = z$: $\langle x_{\varepsilon_j}, s \rangle \rightarrow \langle x, s \rangle = \langle z, s \rangle$, so $\langle x, D^\dagger s \rangle = \langle y, D^\dagger s \rangle$, and after an argument that I missed, this implies that $\bar{D}x = y$. \square

The next theorem tells us that eigenvalues of a (generalized) Dirac operator have better regularity than one might expect.

Theorem 8.39. Let D be a generalized Dirac operator. Then, the kernel of \bar{D} consists of smooth sections.

Proof. Let $s \in \ker(\bar{D})$; we'll show that $s \in H^k$ for all k , which implies by Theorem 8.14 that it's smooth. We'll induct by assuming $s \in H^{k-1}$. Let $\{F_\varepsilon\}$ be a Friedrichs' mollifier, so that $\{F_\varepsilon\}$ and $\{[D, F_\varepsilon]\}$ are bounded families of operators on H^{k-1} .

Using the elliptic estimate,

$$(8.40) \quad \|F_\varepsilon s\|_k \leq C_k(\|F_\varepsilon s\|_{k-1} + \|DF_\varepsilon s\|_{k-1}) = C_k(\|F_\varepsilon s\|_{k-1} + \|[D, F_\varepsilon]s\|_{k-1}).$$

Since $Ds = 0$, this is bounded above. \square

Lecture 9.

Spectral theory: 2/21/18

"There's a reason to erase the boards at the beginning of a lecture rather than the end: you have an incentive to do it well."

This part of today's lecture was given by Dan.

Let V be a finite-dimensional complex inner product space and $T : V \rightarrow V$ be a self-adjoint operator. Then we can diagonalize T : there is an orthogonal direct sum

$$(9.1) \quad V = \bigoplus_{\lambda \in \mathbb{R}} V_\lambda$$

where $T|_{V_\lambda} = \lambda$. Obviously only finitely many V_λ are nonzero.

Geometrically, you could think of this as a family of vector spaces over a base \mathbb{R} , and in fact it's a sheaf (a finite sum of skyscraper sheaves).

Remark 9.2. If you don't have an inner product and T isn't self-adjoint, you can still get something useful, the Jordan decomposition. In this case the sheaf has nilpotent information if there are nontrivial Jordan blocks, but we're not going to encounter this in this course. \triangleleft

Given a function $f : \mathbb{R} \rightarrow \mathbb{R}$, one obtains a new operator $f(T)$ defined by $f(T)|_{V_\lambda} := f(\lambda)$. This is an excellent way to make new operators out of old. For example, the exponential of a matrix is defined through its diagonalization:

$$(9.3) \quad \exp \begin{pmatrix} 4 & 2 \\ 1 & 5 \end{pmatrix} = \begin{pmatrix} -2 & 1 \\ 1 & 1 \end{pmatrix} \exp \begin{pmatrix} 3 & 0 \\ 0 & 6 \end{pmatrix} \begin{pmatrix} -1/3 & 1/3 \\ 1/3 & 2/3 \end{pmatrix} = \begin{pmatrix} -2 & 1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} e^3 & 0 \\ 0 & e^6 \end{pmatrix} \begin{pmatrix} -1/3 & 1/3 \\ 1/3 & 2/3 \end{pmatrix}.$$

We are interested in generalizing this story to infinite dimensions. Let \mathcal{H} be a separable⁹ complex Hilbert space and $T : \mathcal{H} \rightarrow \mathcal{H}$ be a self-adjoint, positive, compact operator.¹⁰

Theorem 9.4 (Spectral theorem). There is an orthogonal basis $\{e_n\}_{n=1}^\infty$ of \mathcal{H} and $\mu_n \in \mathbb{R}_{>0}$ such that

- $\mu_1 \geq \mu_2 \geq \mu_3 \geq \dots$,
- $Te_n = \mu_n e_n$, and
- $\lim_{n \rightarrow \infty} \mu_n = 0$.

⁹If you don't know what this is, don't worry about it; it's a point-set topological axiom. All Hilbert spaces we will encounter in this class are separable.

¹⁰By writing that the domain of T is \mathcal{H} , and not a dense subset of it, we imply that T is already a bounded operator.

Suppose $s \in \mathcal{H}$ is such that $\langle s, e_n \rangle = 0$ for all n . Then $s = 0$. Hence \mathcal{H} decomposes as

$$(9.5) \quad \mathcal{H} = \overline{\bigoplus_{\lambda \in \mathbb{R}} \mathcal{H}_\lambda},$$

where $\dim \mathcal{H}_\lambda$ is finite. This is again an orthogonal direct sum, and again \mathcal{H}_λ defines a sheaf of Hilbert spaces over \mathbb{R} . Because $\{\mu_i\}$ has an accumulation point at 0, this is not a finite sum of skyscraper sheaves.

Another way to think about that it defines a measure of sorts on \mathbb{R} . For $\lambda \in \mathbb{R}$, let $\pi_\lambda: \mathcal{H} \rightarrow \mathcal{H}_\lambda \hookrightarrow \mathcal{H}$ be orthogonal projection onto \mathcal{H}_λ followed by inclusion; hence $\pi_\lambda^2 = \pi_\lambda$. Then

$$(9.6) \quad T = \sum_{\lambda \in \mathbb{R}} \lambda \pi_\lambda,$$

though one must make sense of this uncountable sum! In particular, spectral calculus is the same: given $f: \mathbb{R} \rightarrow \mathbb{R}$, one defines

$$(9.7) \quad f(T) := \sum f(\lambda) \pi_\lambda,$$

where again we have to make sense of this sum, but this is possible.

We won't get into the guts of the proof of Theorem 9.4, but the idea is to study the quadratic form $Q: S(\mathcal{H}) \rightarrow \mathbb{R}$ sending $\xi \mapsto \langle \xi, T\xi \rangle$, where $S(\mathcal{H})$ is the unit sphere inside \mathcal{H} .

Remark 9.8. Theorem 9.4 is the only spectral theorem we'll need in this class, but is not the only one you'll need in your life. ◀

We will apply this theory to the Dirac operator $D: C_M^\infty(S) \rightarrow C_M^\infty(S)$ associated to a Clifford bundle $S \rightarrow M$ over a closed Riemannian manifold M . (More generally, we could take an operator with $D^2 = \nabla^* \nabla + K$, as discussed last time.)

Example 9.9. The first example is $M = S^1$ with length 2π , $S = \mathbb{C}$, and $D = -i \frac{d}{dx}: C^\infty(S^1) \rightarrow C^\infty(S^1)$. This operator is diagonalized by Fourier series: for any $n \in \mathbb{Z}$,

$$(9.10) \quad -i \frac{d}{dx} e^{inx} = n e^{inx}.$$

Unfortunately, this is not a positive operator. $D^2 = -\frac{d^2}{dx^2}$ isn't either, because it has 0 for an eigenvalue (its eigenvalues are n^2 for $n \geq 0$). So maybe we should replace it with $D^2 + 1$, which is positive, but where is $(D^2 + 1)^{-1}$ defined? This, and our desire to live in the world of Hilbert spaces, are why we have to think about Sobolev spaces. ◀

Recall that we defined a Sobolev space $W^k(S)$ for $k = 0, 1, 2, \dots$, denoting the space of L^2 functions which have k derivatives that are also in L^2 . (Well, this is not the precise definition, but that's the point.) Then $W^0(S) = L^2(S)$.

Remark 9.11. One can more generally use L^p norms for any $1 \leq p \leq \infty$, and define $L_k^p(\mathbb{R}^n)$ to be the Sobolev space of L^p functions on \mathbb{R}^k which have k derivatives in L^p . ◀

The Sobolev embedding theorem (Theorem 8.14) shows that $L_k^p(\mathbb{R}^n) \hookrightarrow C^\ell$ iff $k - n/p > \ell$. One can remember this inequality by taking a Sobolev function and scaling it by $x \mapsto \lambda x$. This is the inequality which makes sense as $\lambda \rightarrow 0$ or $\lambda \rightarrow \infty$.

It's also true that $\bigcap_k W^k = C^\infty$. This produces another interpretation of Sobolev spaces: we want to study smooth functions, but their topology is not the best (it's Fréchet, not Hilbert), so we weaken our regularity a little bit to get a better function space.

If $D = \nabla^* \nabla + K$ is a generalized Dirac operator, then

$$(9.12) \quad \int |Ds|^2 = \int |\nabla s|^2 + \int \langle Ks, s \rangle,$$

so already this looks like a Sobolev norm, which is another hint that Dirac operators naturally live on Sobolev spaces.

We then discussed the Rellich theorem (Theorem 8.15), that if $k' < k$, $W_k \hookrightarrow W_{k'}$ is compact. So W_k has a finer topology, but compactness is a really convenient thing to have. This implies that $D: W^1(S) \rightarrow W^0(S)$ is bounded, and that $1 + D^2: W^2(S) \rightarrow W^1(S)$ is bounded. If we could in addition find an inverse R to $1 + D^2$, then $T: W^0 \rightarrow W^2 \hookrightarrow W^0$ using R , then including, T would be a compact operator, and as the inverse of a positive,

self-adjoint operator, T is also self-adjoint. Therefore Theorem 9.4 applies, and by diagonalizing T we're almost all of the way to diagonalizing D .

Roughly, to invert $1 + D^2$, we need to show that

$$(9.13) \quad \|s\|_{W^2} \leq \|(1 + D^2)s\|_{W^0},$$

which shows $1 + D^2$ is injective and has closed image. Then, it only remains to prove it's surjective.

To prove (9.13), we used Proposition 8.17 a few times. The way Roe does surjectivity is to look at the graph of D as an operator. This allows surjectivity to relate to a direct construction of R . Since W^1 is a dense subspace of L^∞ , the graph Γ of $D: W^1 \rightarrow L^2$ is a subspace of $L^2 \times L^2$, and we may take its closure. Using self-adjointness of D , we proved Proposition 8.31, that $\bar{\Gamma}$ is also a graph of some operator \bar{D} . We can use this to define a map $Q: L^2 \rightarrow L^2$ to its orthogonal projection onto $\bar{\Gamma}$ (all inside $L^2 \times L^2$), then projecting that back onto the domain L^2 . Then for all Y ,

$$\begin{aligned} 0 &= \langle Qx - x, y \rangle + \langle \bar{D}Qx, \bar{D}y \rangle \\ &= \langle Qx - x + \bar{D}^2 Qx, y \rangle, \end{aligned}$$

and therefore $(D^2 + 1)^{-1}Qx = x$, so we have not just surjectivity, but an inverse to D .

We want the eigenvectors of the Dirac operator to be smooth, but *a priori* they're just W^2 . This is where the argument involving elliptic estimates and elliptic regularity came in. The upshot is the following spectral theorem for Dirac operators.

Theorem 9.14. *There exists an orthonormal basis $\{s_n\}_{n=1}^\infty$ for $L^2(S)$ and $\lambda_n \in \mathbb{R}_{\geq 0}$ such that*

- $0 \leq \lambda_1 \leq \lambda_2 \leq \dots$,
- $D^2 s_n = \lambda_n s_n$,
- $\lambda_n \rightarrow \infty$, and
- $s_n \in C^\infty(S)$.

Later we will prove a theorem about how fast λ_n grows; for the Dirac operator on the trivial bundle over S^1 , $\lambda_n \approx n^2$.

For any section s , we have the projection formula

$$(9.15) \quad s = \sum_{n=1}^\infty \langle e_n, s \rangle e_n = \sum_{\lambda} s_{\lambda},$$

where $s_{\lambda} := \pi_{\lambda} s$. On the circle, this is just Fourier series: large eigenvalues tell you about large-scale behavior, and small eigenvalues tell you about small-scale behavior. In order for this to converge, larger-scale behavior has to be tamed, and in particular, as you might remember from Fourier theory, better local behavior or regularity is converted to faster decrease of $\|s_{\lambda}\|$ as $\lambda \rightarrow \infty$.

One way to think of this is that $D^2 = -\frac{d}{dx^2}$ has units of $1/L^2$, where L is length. Therefore the units of the eigenvalues are also $1/L^2$, hence small at large scales.

Theorem 9.16. *An L^2 section s is smooth iff $\|s_{\lambda}\| = O(\lambda^{\ell})$ for all $\ell \in \mathbb{Z}_{>0}$.*

Proof. Since

$$(9.17) \quad (D^2)^{\ell} s = \sum_{\lambda} \lambda^{\ell} s_{\lambda},$$

then

$$(9.18) \quad \|(D^2)^{\ell} s\|_0 \leq \sum_{\lambda} \lambda^{\ell} \|s_{\lambda}\| \leq \sum_{\lambda} \frac{C}{\lambda^2} < \infty.$$

TODO: I missed something. \(\square\)

In a similar vein, there's a result for spectral calculus.

Theorem 9.19. *Let $f: \mathbb{R} \rightarrow \mathbb{R}$ be such that $f(\lambda) = O(\lambda^{-\ell})$ for all $\ell \in \mathbb{Z}_{>0}$. Then $\text{Im}(f(D))$ consists of smooth sections.*

That is, to get smooth sections, we use functions which vanish more quickly than $1/p$ where p is any polynomial. The proof idea is to show that $f(D)$ is an integral kernel, i.e.

$$(9.20) \quad f(D)s(x) = \int_M K(x, y)s(y) dy$$

for all $x \in M$ and $s \in L^2(S)$. This is smooth iff $K \in C_{M \times M}^\infty(S \boxtimes S^*)$, and we can use elliptic estimates to prove this. We need an explicit formula for $K(x, y)$ in order to do this, and the answer turns out to be

$$(9.21) \quad K(x, y) = \sum_\lambda f(\lambda) \langle s_\lambda(y), - \rangle s_\lambda(x) : S_y \rightarrow S_x.$$

This is explicit enough that you can do analysis with it, showing that it's bounded in every Sobolev norm, and hence is smooth.

Lecture 10.

Dirac complexes and the Hodge theorem: 2/21/18

In this part of the lecture, Adrian spoke about Dirac complexes and the Hodge theorem. Everything will be worked out in finite dimensions first, where things are really easy, and then cover the infinite-dimensional case, whose proof looks similar at large scales, but requires checking a few more things.

A Dirac complex can be viewed as a refinement of a Clifford bundle.

Definition 10.1. Let M be a compact, oriented Riemannian manifold. A *Dirac complex* on M consists of

- a sequence S^0, \dots, S^n of Hermitian vector bundles over M ,
- differential operators $d^i : C^\infty(S^i) \rightarrow C^\infty(S^{i+1})$ for $i = 0, \dots, n-1$, and
- a Clifford bundle structure on

$$S := \bigoplus_{i=0}^n S^i,$$

such that $D = d + d^*$, where

- D is the Dirac operator on S and
- d^* is the adjoint of $d : S \rightarrow S$ with respect to the inner product

$$(10.2) \quad \langle s, s' \rangle := \int_M \langle s, s' \rangle d\text{vol}.$$

Now let's do some finite-dimensional Hodge theory. Let $d : V^0 \rightarrow V^1$ be a map between finite-dimensional inner product spaces (real or complex) and $d^* : V^1 \rightarrow V^0$ denote its adjoint.

Lemma 10.3. $\ker(d) = \ker(d^*d)$.

Proof. That $\ker(d) \subseteq \ker(d^*d)$ is obvious, so suppose $v \in \ker(d^*d)$, meaning

$$(10.4) \quad 0 = \langle v, d^*dv \rangle = \langle dv, dv \rangle,$$

and therefore $dv = 0$. □

The operators $d^*d : V^0 \rightarrow V^0$ and $dd^* : V^1 \rightarrow V^1$ are self-adjoint, so they have well-behaved spectral theory. Let $\lambda_0, \dots, \lambda_n$ be the eigenvalues of d^*d and v_0, \dots, v_n be eigenvectors for them. Then, $\lambda_0, \dots, \lambda_n$ are the eigenvalues of dd^* , and the eigenvectors are dv_0, \dots, dv_n .

More generally, consider a finite sequence of finite-dimensional inner product spaces together with adjoint operators

$$(10.5) \quad 0 \longrightarrow V^0 \xrightleftharpoons[d^*]{d} V^1 \xrightleftharpoons[d^*]{d} V^2 \xrightleftharpoons[d^*]{d} \dots \xrightleftharpoons[d^*]{d} V^{n-1} \xrightleftharpoons[d^*]{d} V^n \longrightarrow 0.$$

Lemma 10.3 generalizes to another, almost as easy, lemma.

Lemma 10.6. In the setting of (10.5), $\ker(d^*) \cap \ker(d) = \ker(dd^* + d^*d)$.

Proof. Again the forward inclusion is obvious, so suppose $v \in \ker(dd^* + d^*d)$, so $\langle d^*dv + dd^*v, v \rangle = 0$.

TODO: something I didn't understand happened here, using eigenvectors. □

In particular, this leads to a canonical identification of $\ker(dd^* + d^*d)$ with the homology of the complex (10.5).

General Hodge theory. In this case we have a complex of vector bundles S^i , and let $V^i := C^\infty(S^i)$, so we have an induced complex of topological vector spaces.¹¹

... something happened here regarding how the finite-dimensional statement changes because certain projections exist...

Theorem 10.7. *We have a projection $C^\infty(S^i) \rightarrow \mathcal{H}^i$ such that $H^i(C^\infty(S^\bullet)) \cong \mathcal{H}^i$.*

In fact, you can promote this into an isomorphism of the cohomology of $C^\infty(S^\bullet)$ with \mathcal{H}^\bullet , where the latter's differentials are all the zero map.

Proof. There exists an orthogonal projection $P: L^2(S^i) \rightarrow \mathcal{H}^i$ because \mathcal{H}^i is finite-dimensional. Let $i: \mathcal{H}^i \hookrightarrow L^2(S^i)$ denote inclusion. Let δ_{ij} denote the Kronecker delta and $f(\lambda) = 1 - \delta_{\lambda 0}$. Then $iP = 1 \cdot f(0)$.

If

$$(10.8) \quad g(\lambda) := \begin{cases} \lambda^{-2}, & \lambda \neq 0 \\ 0, & \lambda = 0, \end{cases}$$

then g is bounded on D because of (TODO) something I didn't understand. If $G := g(D)$ and $H := d^*G$, then

$$(10.9) \quad (d^*d + dd^*)G = f(D) = Hd + dH = 1 - (1 - f(D)) = 1 - iP$$

which means it's a chain homotopy between $(C^\infty(S^\bullet), d)$ and $(\mathcal{H}^\bullet, 0)$, so their cohomology groups agree. \square

The case we're interested in is the Dirac complex of differential forms: $S^i := \Lambda^i(T^*M) \otimes \mathbb{C}$ and $d^i := d$, with the Clifford action defined as follows: if $v \in TM$, $v \cdot \varphi := v^\vee \wedge \varphi - v \lrcorner \varphi$ and extending to the rest of the Clifford algebra.

The inner product is the usual inner product on differential forms induced from the Riemannian metric, which we've seen a few times. The whole story follows from a few formal facts.

- The Hodge star is characterized by $\varphi \wedge \star \psi = \langle \varphi, \psi \rangle \text{vol}$ for any k -forms φ, ψ .
- $\star \star \varphi = (-1)^{k(n-k)} \varphi$.
- Contraction is a derivation:

$$(10.10) \quad \lrcorner (\varphi \wedge \psi) = \lrcorner \varphi \wedge \psi + (-1)^{|\varphi|} \varphi \wedge \lrcorner \psi.$$

- The relation

$$(10.11) \quad v \lrcorner \varphi = (-1)^{k(n-k)} \star (\star \varphi \wedge v^\vee).$$

- ∇ commutes with the Hodge star.
- The relation

$$(10.12) \quad \nabla_X (v \lrcorner \varphi) = (\nabla_X v) \lrcorner \varphi + X \lrcorner \nabla_v \varphi.$$

You can check all the properties for a Dirac bundle using these, and using asynchronous frames, you can also show that the Dirac operator is the Hodge operator.

Lecture 11.

Applications of Hodge theory: 2/21/18

Rok spoke in this part of the talk. Throughout M is a closed, oriented Riemannian manifold of dimension n .

Corollary 11.1. *If S is a Dirac complex, the cohomology groups $H^i(S)$ are finite.*

This follows because we constructed an isomorphism above. Compactness of M is crucial here, entering in the form of ellipticity.

Theorem 11.2 (Poincaré duality). *The integration pairing*

$$(11.3) \quad H^k(M; \mathbb{C}) \otimes_{\mathbb{C}} H^{n-k}(M; \mathbb{C}) \xrightarrow{\sim} H^n(M; \mathbb{C}) \xrightarrow[\cong]{\int_M} \mathbb{C}$$

is a perfect pairing, i.e. it defines an isomorphism $H^{n-k}(M; \mathbb{C}) \cong H^k(M; \mathbb{C})^$.*

¹¹TODO: is this actually a Dirac complex?

Proof. Stokes' theorem implies that, since M is closed, the pairing is well-defined. Fix an $[\alpha] \in H^k(M; \mathbb{C})$ and suppose $\int_M \alpha \wedge \beta = 0$ for all β . By the Hodge theorem, there is a unique harmonic representative for $[\alpha]$, which we'll also call α , and $d\alpha = 0$ and $d^*\alpha = 0$. Hence $(-1)^2 \star d\star \alpha = 0$, so $d(\star \alpha) = 0$. Taking $\beta = \star \alpha$, we get

$$(11.4) \quad 0 = \int_M \alpha \wedge \star \alpha = \|\alpha\|_{L^2}^2,$$

i.e. $\alpha = 0$. □

You can use this to set up intersection theory, in the form of finding Poincaré duals to closed, oriented submanifolds.

Let $C \subset M$ be a closed, oriented submanifold of dimension k . Then integrating k -forms restricted to C defines a linear map $\varphi: H^k(M; \mathbb{C}) \rightarrow \mathbb{C}$, well-defined by Stokes' theorem. By the Hodge theorem, $H^k(M; \mathbb{C}) \cong \mathcal{H}^k(M)$, so we may complete and extend $\tilde{\varphi} := \varphi \circ P$, where $P: \Omega_{L^2}^k(M) \rightarrow \mathcal{H}^k(M)$ is projection onto the closed subspace of harmonic forms. In particular, $\tilde{\varphi}$ is a continuous linear functional, so by the Riesz representation theorem, $\tilde{\varphi} = \langle -, \beta \rangle$ for some $\beta \in \Omega_{L^2}^k(M)$.

Since P is a projector, it's self-adjoint, and

$$\begin{aligned} \langle \alpha, \beta \rangle &= \int_C P\alpha = \int_C P^2\alpha = \langle P\alpha, \beta \rangle \\ &= \langle \alpha, P^*\beta \rangle = \langle \alpha, P\beta \rangle. \end{aligned}$$

Therefore in particular β is harmonic!

Definition 11.5. The Poincaré dual to C is $P_C := [\star \beta] \in H^{n-k}(M; \mathbb{C})$.

There are other ways to define this, but the cleanliness of the Hodge star is nice.

We can therefore do intersection theory: if C and C' are closed, oriented, transverse submanifolds of M such that $\dim C + \dim C' = n$, then their intersection number is

$$(11.6) \quad I(C, C') := \int_C P_{C'} = (-1)^{(\dim C)(\dim C')} \int_M P_{C'} \wedge P_C = (-1)^{(\dim C)(\dim C')} I(C', C).$$

So this is commutative up to a sign, just as in cohomology.

Example 11.7. For any point $p \in M$, $P_p \in \Omega^n(M)$ must be λ times the volume form for some $\lambda \in \mathbb{C}$. Then

$$(11.8) \quad f(p) = \int_M f \cdot P - p = \lambda \int_M f \text{ vol}.$$

This seems quite strange until you realize that it means $\lambda = 1/\text{vol}(M)$, after which (11.8) is a tautology. ◀

We'll use this to provide another proof of Corollary 5.3, that a compact manifold with nonnegative Ricci curvature, and positive Ricci curvature somewhere, has vanishing first cohomology.

Proof of Corollary 5.3. Recall that $D^2\alpha = 0$ has no nonzero solutions, and by the Hodge theorem, $H^1(M; \mathbb{C}) \cong \mathcal{H}^1(M) = 0$. □

Lecture 12.

The heat and wave equations: 2/28/18

In this part of the lecture, Ravi Mohan spoke about the heat and wave equations on a compact Riemannian manifold (M, g) .

Let $S \rightarrow M$ be a Clifford bundle on M and s be a smooth section of S varying smoothly in time t . The *heat equation* is the second-order differential equation

$$(12.1) \quad \frac{\partial s}{\partial t} + D^2s = 0.$$

In physics, this is known as the *diffusion equation*: it takes a point distribution and spreads its mass out over M , modeling heat flow.

The wave equation is the first-order equation

$$(12.2) \quad \frac{\partial s}{\partial t} - iDs = 0.$$

Given a wave as initial data, this propagates the wave.

Proposition 12.3. *Let $s_0 \in \Gamma(M, S)$ be a smooth section.*

- (1) *The heat equation (12.1) has a unique solution s_t corresponding to the initial data s_0 , which exists for all $t \geq 0$. Moreover, $\|s_t\|^2 \leq \|s_0\|^2$.*
- (2) *The wave equation (12.2) has a unique solution s_t corresponding to the initial data s_0 , which exists for all $t \in \mathbb{R}$. Moreover, $\|s_t\|^2 = \|s_0\|^2$.*

Proof. First we prove (1). For existence, consider $e^{-tD^2} : C^\infty(S) \rightarrow C^\infty(S)$, and let

$$(12.4) \quad s_t := e^{-tD^2} s_0.$$

Then

$$(12.5) \quad \frac{\partial}{\partial t} s_t = -D^2 s_t$$

$$(12.6) \quad \implies \left(\frac{\partial}{\partial t} + D^2 \right) s_t = 0.$$

For the norm, let $\{e_i\}$ be an orthonormal basis of eigenfunctions for D^2 , so that $a_0 = \sum_i a_i e_i$ and $D^2 e_i = n_i e_i$. Then

$$\begin{aligned} s_t &= e^{-tD^2} s_0 \\ &= e^{-tD^2} \sum_i a_i e_i \\ &= \sum_i a_i e^{-tn_i} e_i. \end{aligned}$$

Therefore

$$(12.7) \quad \|s_t\|^2 = \sum_i \|a_i\|^2 e^{-2tn_i},$$

which decays with t .

For the wave equation, the proof is very similar, except we use e^{iDt} in place of e^{-tD^2} . For (12.7), $|e^{itn_i}| = 1$, allowing us to get $=$ instead of \leq . \square

Remark 12.8. Looking at the heat and wave equations in the eigenbasis allows us to understand what they're doing. For example, the heat equation dampens Fourier modes, but dampens the larger Fourier modes faster. In fact, even if your initial condition s_0 isn't smooth, s_ε is smooth for any $\varepsilon > 0$, because the large Fourier modes which cause s_0 to not be smooth are dampened to something acceptable. \blacktriangleleft

There exists a smooth family of smooth sections k_t , $t > 0$, of $S \boxtimes S^* \rightarrow M \times M$ such that for all smooth sections s ,

$$(12.9) \quad e^{-tD^2} s(p) = \int_M k_t(p, q) s(q) dq.$$

This k_t is called the *heat kernel* for S .

Proposition 12.10. *We have that*

$$(12.11) \quad \left(\frac{\partial}{\partial t} - D_p^2 \right) k_t(p, q) = 0.$$

$$(12.12) \quad \lim_{t \rightarrow 0} k_t = \delta(p - q).$$

Proof of part (12.11). It suffices to compute: since

$$(12.13) \quad \left(\frac{\partial}{\partial t} + D_p^2 \right) e^{-tD^2} s(p) = \left(\frac{\partial}{\partial t} + D_p^2 \right) \int_M k_t(p, q) s(q) dq,$$

then

$$(12.14) \quad 0 = \int \left(\frac{\partial}{\partial t} + D_p^2 \right) k_t(p, q) s(q) dq.$$

Since this is true for all s , (12.11) follows. \square

The heat kernel is for smooth sections. If we have weaker regularity, we need something slightly different.

Definition 12.15. An approximate heat kernel of order m is a time-dependent section $k'_t(p, q)$ of $S \boxtimes S^* \rightarrow M \times M$ which approximately satisfies the heat equation, in that there is a C^m section r_t^m of $S \boxtimes S^* \rightarrow M \times M$ such that

$$\left(\frac{\partial}{\partial t} + D_p^2 \right) k'_t(p, q) = t^m r_t^m(p, q).$$

Proposition 12.16. Let s_t be a continuous family of C^2 sections of S . Then, there is a unique smooth family of smooth sections \tilde{s}_t of S with $\tilde{s}_0 = 0$ and satisfying the heat equation

$$\left(\frac{\partial}{\partial t} + D^2 \right) \tilde{s}_t = s_t.$$

Lemma 12.17. The asymptotic behavior of the heat kernel in Euclidean space \mathbb{E}^n is

$$(12.18) \quad k_t(x, y) \sim \frac{1}{(4\pi t)^{n/2}} \exp\left(-\frac{|x-y|^2}{4t}\right).$$

On a Riemannian n -manifold (M, g) , if p and q are given in geodesic coordinates on M , the asymptotic behavior of the heat kernel is

$$(12.19) \quad k_t(p, q) \sim h := \frac{1}{(4\pi t)^{n/2}} \exp\left(-\frac{d(p, q)^2}{4t}\right).$$

Here $d(p, q)$ is the geodesic distance, i.e. the infimum of the lengths of geodesics between p and q .

Lemma 12.20. Let h be as in (12.19). Then

$$(12.21) \quad \nabla h = -\frac{h}{2t} r \frac{\partial}{\partial r}$$

$$(12.22) \quad \frac{\partial h}{\partial t} + \Delta h = \frac{rh}{4gt} \frac{\partial g}{\partial r}.$$

TODO: proof was given, and I missed it.

Definition 12.23. Let $f: \mathbb{R} \rightarrow \mathbb{C}$ be a function and $S := \sum c_n t^n$ be a formal power series. Then S is an asymptotic series for f if for all N ,

$$\lim_{t \rightarrow 0} t^{-N} \left| f(t) - \sum_{n=0}^N c_n t^n \right| = 0.$$

Remark 12.24. We know that as $t \rightarrow \infty$, the heat operator converges very strongly to projection onto the kernel. So when we use it to understand the index, we'll find a formula for the index in terms of the heat kernel at all $t \geq 0$. At infinity, it simplifies to where we'll recognize it as an index; then at 0, we'll be able to identify it with something else.

If $x \neq y$, (12.18) has exponential decay as $t \rightarrow 0$. On the diagonal, it's just the function $1/(4\pi t)^{n/2}$, which doesn't converge, but we'll be able to use an asymptotic series to compare what happens on a curved manifold with what happens on \mathbb{E}^n . Even though these series don't converge, they're still useful for computing something! The way this is accomplished is by taking polynomials instead of formal power series. \leftarrow

Lecture 13.

Estimates with the heat kernel: 2/28/18

In this part of the lecture, Cameron Darwin spoke.
First, we'll generalize Definition 12.23.

Definition 13.1. Let E be a Banach space and $f : \mathbb{R}_+ \rightarrow E$ be a function. Let $S = \sum_{k=0}^{\infty} a_k(t)$ be a power series valued in E (i.e. each a_k is a function $\mathbb{R}_+ \rightarrow E$). We say that S is an *asymptotic expansion* for f if for all n , there exists an ℓ_n such that if $\ell \geq \ell_n$, there's a constant $C_{\ell,n}$ such that for t sufficiently small,

$$\left\| f(t) - \sum_{k=0}^{\ell} a_k(t) \right\| \leq C_{\ell,n} |t|^n.$$

Proposition 13.2. Let k_t be the heat kernel and h_t be as in (12.18).

(1) There is an asymptotic expansion of for k_t of the form

$$(13.3) \quad k_t(p, q) \sim h_t(p, q) (\Theta_0 + t\Theta_1 + t^2\Theta_2 + \cdots),$$

where the Θ_n are smooth sections of $S \boxtimes S^* \rightarrow M \times M$.

(2) (13.3) is valid in the Banach space of C^r sections of $S \boxtimes S^*$ for all $r \geq 0$, and may be formally differentiated in p, q , and t to obtain asymptotic expansions for derivatives of k_t .

(3) The values of $\Theta_j(p, p)$ can be computed from algebraic expressions in the metric and connection coefficients of S , and the first term is $\Theta_0(p, p) = 1$.

If k'_t is an approximate heat kernel of order m' , then there's a C such that for small t ,

$$(13.4) \quad \|k_t - k'_t\|_{H^{m'}} \leq C t^{m'+1},$$

so by the Sobolev embedding theorem, if $m' > m + n/2$, we get a bound for the C^m norm:

$$(13.5) \quad \|k_t - k'_t\|_{C^m} \leq \|k_t - k'_t\|_{H^{m'}} \leq C t^{m'+1}.$$

A similar analysis can be performed with

$$(13.6) \quad \|\nabla_{x_1} \cdots \nabla_{x_n} k_t - \nabla_{y_1} \cdots \nabla_{y_n} k_t^j\|$$

to bound it above.

Let $\Theta \in \Gamma(M \times M, S \boxtimes S^*)$ and $s \in \Gamma(M, S)$. Then we will use the notation

$$(13.7) \quad (\theta \vdash s)(p) := \int_M \theta(p, q) s(q) dq.$$

Lemma 13.8. For all sections $\Theta_0, \dots, \Theta_j$ of $S \boxtimes S^*$,

$$(13.9) \quad (h_t(\Theta_0 + t\Theta_1 + \cdots + t^j\Theta_j) \vdash s)(p) \xrightarrow{t \rightarrow 0} \Theta_0(p, p) s(p)$$

and the convergence is uniform in p .

Lemma 13.10. For all $\varepsilon, \delta > 0$, there's a t_0 such that for $0 < t < t_0$ and all $p \in \mathbb{R}^N$,

$$(13.11) \quad 1 - \varepsilon < \int_{B_\delta(p)} h_t(p, q) dq < 1.$$

This says that h_t looks like a delta function near 0.

Corollary 13.12. For all $\alpha, \beta, \delta > 0$, there's a t_0 such that for $0 < t < t_0$,

$$\left| \int_{B_\delta(p)} h_t(p, q) dq - 1 \right| < \alpha$$

$$\left| \int_{B_\delta(p)^c} h_t(p, q) dq \right| < \beta.$$

Given a continuous section $s \in \Gamma(M, S)$, we'll define a section \tilde{s} of $S \times M \rightarrow M \times M$ by

$$(13.13) \quad |\tilde{s}(p, q)| = |\Theta(p, q)s(q) - \Theta(p, p)s(p)|.$$

TODO: what happened here? Why are we defining \tilde{s} via its norm?

Anyways, by Lemma 13.8, \tilde{s} is uniformly continuous in p and q . In particular, for all $\omega > 0$, there's a $\delta > 0$ such that if $d(p, q) < \delta$, then $|\tilde{s}(p, q)| < \omega$. This implies that for all $\varepsilon > 0$, there's a t_0 independent of p such that for $t \in (0, t_0)$,

$$(13.14) \quad \left| \int_{B_\delta(p)^c} h_t(p, q) \Theta(p, q) s(q) dq \right| < \frac{\varepsilon}{2}.$$

Finally, for all $\varepsilon > 0$, and independently of p and R , if

$$(13.15a) \quad \left| \int_{B_R(p)} h_t(p, q) dq - 1 \right| < \alpha$$

and

$$(13.15b) \quad |\Theta(p, q) - \Theta(p, p)s(p)| < \varepsilon,$$

then

$$(13.16) \quad \left| \int_{B_R(p)} h_t(p, q) \Theta(p, q) s(q) dq - \Theta(p, p)s(p) \right| < \frac{\varepsilon}{2}.$$

The upshot of all of this is that we know what conditions to put on Θ_0 such that $h_t(\Theta_0 + t\Theta_1 + \dots + t^j\Theta_j)$ to converge to a delta function. We'll next discuss what else we need to ask of the Θ_i in order to ensure we get an approximate heat kernel.

Lemma 13.17. Suppose we have two finite sequences $\Theta_0, \Theta_1, \dots$, and $\Theta'_0, \Theta'_1, \dots$ and that there's an open neighborhood U of the diagonal in $M \times M$ such that for all i , $\Theta_i|_U = \Theta'_i|_U$. Then $k_t \sim h_t(\Theta_0 + t\Theta_1 + \dots)$ iff $k_t \sim h_t(\Theta'_0 + t\Theta'_1 + \dots)$.

Proof sketch. Since M is compact, we can assume U is the set of points $(p, q) \in M \times M$ such that $d(p, q) < \varepsilon$. Therefore

$$(13.18) \quad h_t(p, q) = \frac{1}{(4\pi t)^{n/2}} \exp\left(-\frac{d(p, q)^2}{4t}\right)$$

decays faster than any polynomial on U^c , so only the stuff inside U matters. \square

Therefore it suffices to construct these Θ_i in a neighborhood of the diagonal.

Definition 13.19. Let (M, g) be a compact Riemannian manifold.

- At a $p \in M$, the *injectivity radius* is the supremum of the set of radii r such that the exponential map $\exp: T_p M \rightarrow M$ is injective when restricted to $B_r(0) \subset T_p M$.
- The *injectivity radius of M* is the infimum of the injectivity radii at all $p \in M$.

Let R be less than the injectivity radius of M ; then the neighborhood $N_R(\Delta)$ of $\Delta \subset M \times M$ admits a cover by charts which are in geodesic coordinates. Using a system of cutoff functions for these charts, we can reduce the problem to working in a single chart in geodesic coordinates.

TODO: something happened here that I missed and didn't understand.

Using these estimates, we get

$$\begin{aligned} \frac{1}{h} \left(\frac{\partial}{\partial t} + D^2 \right) (hs) &= \frac{1}{h} \left(h \frac{\partial s}{\partial t} + h D^2 s + \left(\frac{\partial h}{\partial t} + \Delta h \right) s - 2 \nabla_{\nabla h} s \right) \\ &= \frac{1}{h} \left(h \frac{\partial s}{\partial t} + h D^2 s + \frac{r h}{4g t} \frac{\partial g}{\partial r} s + \frac{h}{t} \nabla_{r \partial / \partial r} s \right) \\ &= \frac{\partial s}{\partial t} + D^2 s + \frac{r}{4g t} \frac{\partial g}{\partial r} s + \frac{1}{r} \nabla_{r \partial / \partial r} s. \end{aligned}$$

Then **TODO:** ???, and we're trying to determine u , which is **TODO**. If

$$(13.20) \quad u \sim u_0 + t u_1 + t^2 u_2 + \dots,$$

then **TODO** the estimate was erased before I could follow it or write it down. I'll pick up taking notes where I can next follow what's going on...

Let $\varphi := g^{1/4} D^2 u_j$ and

$$(13.21) \quad \psi(r, s) := \int_0^r \varphi(\rho, s) d\rho.$$

Then, the function [TODO: erased before I could write it down] is smooth, which is somewhat mysterious: Roe asserts it, but it's not immediately clear, and there doesn't seem to be a way around it.

Remark 13.22. So now we have sections Θ_j in a neighborhood of the diagonal, constructed from an inductive sequence of ODEs. The coefficients of the ODEs are smooth everywhere, including on the diagonal, so the theorem we're relying on asserts that we can solve a family of ODEs with parameters and obtain smooth solutions. In fact, since the family is linear, we can write down the solutions explicitly with integration.

There's a geometric picture of this, which relates asymptotic expansions to genuine Taylor expansions, but this is harder. \blacktriangleleft

The last thing we'll do is show that these asymptotic sums are actually approximate heat kernels. We can compute

$$(13.23) \quad \left(\frac{\partial}{\partial t} + D^2 \right) (h_t(p, q) (\Theta_0(p, q) + \cdots + t^j \Theta_j(p, q))) = h \left(\sum_{i=0}^j i t^{i-1} u_i + t^i D^2 u_i + t^{i-1} \frac{r}{4g} \frac{\partial g}{\partial r} u_i + t^{i-1} \nabla_{r\partial/\partial r} u_i \right)$$

which is equal to TODO and therefore we can conclude that it's an approximate heat kernel.

The last thing we asserted is that on Δ , $\Theta_i(p, p)$ can be computed using algebraic expressions in the metric, connection coefficients, and finitely many derivatives of these things. This is obvious for $\Theta_0 = \mathbf{1}$, and only slightly less obvious for Θ_1 : in synchronous coordinates, it's also possible to see it explicitly. For general Θ_j , the argument is inductive.

Suppose that x^p and x^q are in geodesic coordinates; then, the Taylor expansion for the metric is

$$(13.24) \quad g \sim 1 + \frac{1}{3} \sum_{i,p,q} R_{ipqi} x^p x^q + O(|x|^3).$$

Since $K = \sum_{i,p} R_{ippi}$, we have

$$(13.25) \quad g^{1/4} \sim 1 - \frac{1}{12} \sum_{i,p,q} R_{ipqi} x^p x^q - \frac{1}{6} \sum_{i,p} K + \cdots$$

and therefore can give a few examples of the Θ_j terms.

$$(13.26a) \quad \Theta_0 = \mathbf{1}$$

$$(13.26b) \quad \Theta_1(p, p) = \frac{1}{6} \kappa(p) - \mathbb{K}(p).$$

Here \mathbb{K} is the Clifford-contracted curvature.

Lecture 14.

Asymptotic expansion, I: 3/20/18

Dan spoke about the asymptotic expansion, in order to provide an additional perspective.

Let M be a closed Riemannian manifold. Then the heat equation for $t > 0$ and $x, y \in M$ is

$$(14.1) \quad \left(\frac{d}{dt} + \Delta_x \right) p_t(x, y) = 0.$$

The kernel $p_t(x, y)$ is called the *heat kernel*. As $t \rightarrow 0$, $p_t(-, y) \rightarrow \delta_y$, so it forms an approximation to the identity:

$$(14.2) \quad \lim_{t \rightarrow 0} \int_M f(x) p_t(x, y) |dx| = f(y).$$

There's something strange about (14.1): we're mixing units of time and of length squared. In physics, this is solved by inserting a constant, the *specific heat*, which has units that make everything type-check. For us, this says that, approximately, $T \sim L^2$ in scaling. A similar phenomenon occurs in special relativity, where the constant is the speed of light.

Example 14.3. Let \mathbb{E}^n denote n -dimensional Euclidean space. Then the heat kernel is

$$(14.4) \quad p_t(x, y) = (4\pi t)^{-n/2} e^{-d_E(x, y)^2/4t}.$$

Any exponentiated quantity should be dimensionless, so length squared must be comparable to time. \leftarrow

For a Dirac operator $\Delta = D^2$, where D acts on sections of a spinor bundle $S \rightarrow M$, the heat kernel is section-valued, a one-parameter family of sections of $\text{Hom}(S, S)$.

If M is compact, we're going to scale it (which is why we care about which units are comparable). In a neighborhood of a point y that we're interested in, we'll consider an $\varepsilon > 0$ and scale distances near y by $1/\varepsilon$, and correspondingly time by $1/\sqrt{\varepsilon}$. This has the effect of pushing everything else in M far away as $\varepsilon \rightarrow 0$, so the geometry only sees things near y .

To make this precise, we can choose a neighborhood of y on which the exponential map is an isometry, then replace its image with a much larger ball around $0 \in T_y M$ and glue. This defines a family of manifolds over $(0, \infty)$: topologically, the fiber is just M , but its Riemannian metric has changed.

Example 14.5. For a simple example, suppose $M = S^1$; then, as $\varepsilon \rightarrow 0$ we get longer and longer circles. \leftarrow

When $\varepsilon \rightarrow 0$, the limit is precisely $T_y M$, with Euclidean structure induced from the Riemannian metric on M . The bundle has a section given by y , and this means we obtain an origin on $T_y M$, so it's an inner product space, not just a Euclidean space.

We will then consider heat flow of a specific $x \in M$ on each fiber of this bundle. This generates a path that blows up as $\varepsilon \rightarrow 0$ (since we don't see any $x \neq y$ at 0), so we have a flow $p_\varepsilon^t(x, y)$. If we look at the diagonal and let $t = 1$, this turns out to be a smooth function, hence has a Taylor expansion. In fact, this function is even, so the Taylor expansion is only in terms of even powers of ε .

Then, after scaling back, we'll recover the asymptotic expansion of the heat kernel.

Then, we can make the problem slightly fancier: we don't just want y to be fixed. Instead, we consider the family $\mathcal{Y} \rightarrow [0, \infty) \times M$, where over $(-, y)$ we blow the metric up near y . There's a canonical section of this bundle sending $(\varepsilon, y) \mapsto y$ in the blowup by ε around y .

The map $\mathcal{Y} \rightarrow [0, \infty) \times M$ is not a fiber bundle, because the fiber over 0 isn't diffeomorphic (or even homeomorphic) to the fiber at ε for any $\varepsilon > 0$. However, it is a surjective submersion, which is still nice. Nonetheless, this class of families of manifolds contains a little weirdness: there's a surjective submersion $X \rightarrow \mathbb{R}$ whose fiber is S^1 over a negative number and $S^1 \amalg S^1$ over a positive number. Over 0, it's two lines.

Thus surjective submersions allow you to connect manifolds via paths. This perspective is very important in algebraic geometry; there's a conjecture that Calabi-Yau 3-folds are connected under a similar kind of paths. There's a process called *deformation to the normal cone* which is analogous to studying how a small circle around y blows up as $\varepsilon \rightarrow 0$; this is a useful technique to prove analogues of index theorems in algebraic geometry.

Remark 14.6. Here we sketch a construction of the bundle $\mathcal{X} \rightarrow \mathbb{R}$ for a fixed y . Let Σ_y denote the one-point compactification of $T_y M$; then, there's a map $F: M \rightarrow \Sigma_y$ which is the inverse to the exponential map inside a small neighborhood of y and sends things far from y to the added point ∞ . Then, we define $\mathcal{X} \subset \mathbb{R} \times \Sigma_y \times M$ with coordinates (ε, a, x) to be cut out by the equation $F(x) = \varepsilon a$ away from $\{0\} \times \{\infty\} \times M$.

One must show this is a regular value, but it is, so this is a smooth manifold. We have metrics on M and \mathbb{R} , and get one on Σ_y induced from the inner product on $T_y M$, so \mathcal{X} inherits a Riemannian metric — and in particular, so do its fibers. \leftarrow

To show that the function is smooth and even, one has to bring in analysis, making estimates such as

$$(14.7) \quad \left| k_t^1(a) - |\varepsilon|^{-n} k_{t/\varepsilon^2}^\varepsilon(a/\varepsilon) \right| \leq C |\varepsilon|^{-n} e^{-c/4t}.$$

This also gives you some concrete information about the asymptotic expansion, which is the case $k_t^1(0)$. Specifically,

$$(14.8) \quad k_t^1(0) = |\varepsilon|^{-n} k_{t/\varepsilon^2}^\varepsilon(0) + O(|\varepsilon|^{-n} e^{-c/t}).$$

Now set $\varepsilon = t^{1/2}$; in particular,

$$(14.9) \quad k_t^1(0) = t^{-n/2} k_{t^{1/2}}^{t^{1/2}}(0) + O(t^{-n/2} e^{-c/t}).$$

This is the thing whose Taylor expansion we want.

TODO: after this I ceased to follow.

Lecture 15.

Asymptotic expansion, II: 3/20/18

This part of the lecture was by Cameron. I'm really out of it right now, so these notes might not be so great. Sorry about that.

Ultimately, to understand index theory on noncompact manifolds, we'll need to study the heat equation as it relates to the wave equation. Therefore we're interested in the same kind of analysis we've been doing, but for the wave equation, and will prove a similar localization result.

Proposition 15.1. *Let $\mathcal{R}(\mathbb{R})$ denote the space of functions vanishing at infinity and D be a Dirac operator. Then the function from $\mathcal{R}(\mathbb{R})$ to the space of C^∞ integral kernels sending $f \mapsto f(D)$ is continuous.*

The topology on the codomain is a Fréchet topology induced by all of the C^k norms.

Proof. Let λ be an eigenvalue for D and P_λ be projection onto its eigenspace. If K_λ denotes the smoothings kernel for P_λ , then the kernel for $f(D)$ is given by

$$(15.2) \quad \sum_{\lambda} f(\lambda) K_{\lambda}.$$

We next claim that for all k , there's a $c_k > 0$ and an $\ell(k) \in \mathbb{Z}^{>0}$ such that for all λ ,

$$(15.3) \quad \|K_{\lambda}\|_k \leq c_k (1 + |\lambda|)^{\ell(k)},$$

and $\ell(k) \geq \dim M/2 + k$. The proof of this was not particularly enlightening and wasn't presented.

Something more enlightening to prove: if $\{\psi_n\}$ is an orthonormal basis for $L^2(S)$ such that $D^2\psi_n = \lambda_n\psi_n$, then if $\ell > \dim M/2$, then

$$(15.4) \quad \sum_n \frac{1}{(1 + \lambda_n)^\ell} < \infty.$$

To prove this claim, we fix an $N > 0$. Then

$$\begin{aligned} \sum_{n \leq N} \frac{1}{(1 + \lambda_n)^\ell} &= \sum_{n \leq N} \frac{\|\psi_n\|^2}{(1 + \lambda_n)^\ell} \\ &= \sum_{n \leq N} \frac{1}{(1 + \lambda_n)^\ell} \int_M \|\psi_n(x)\|^2 dx \\ &= \int_M \sum_i \sum_{n \leq N} \frac{|\langle \psi_n(x), \sigma_i(x) \rangle|^2}{(1 + \lambda_n)^2} dx, \end{aligned}$$

where at each x , $\{\sigma_i(x)\}$ is an orthonormal basis for S_x .

Now suppose $\sum a_n \psi_n \in W^\ell$. By the Sobolev embedding theorem, $W^\ell \hookrightarrow C^0$ with the operator norm. An elliptic estimate tells us $\|\psi_n\|_{W^\ell} \leq (1 + \lambda_n)^\ell$, and so

$$(15.5) \quad \left| \sum_n a_n \psi_n(x) \right|_x \leq \left\| \sum_n a_n \psi_n \right\|_{C^0} \leq C \left(\sum_n |a_n|^2 (1 + \lambda_n)^\ell \right)^{1/2}.$$

Let $\sigma_1, \dots, \sigma_m$ be an orthonormal basis for S_x , and fix an $N > 0$. Then let

$$(15.6) \quad a_n := \begin{cases} \frac{\langle \psi_n(x), \sigma_i(x) \rangle_x}{(1 + \lambda_n)^\ell}, & n \leq N \\ 0, & n > N. \end{cases}$$

Then

$$\begin{aligned} (15.7) \quad \left\langle \sigma_i, \sum_{n \leq N} a_n \psi_n(x) \right\rangle_x &= \sum_{n \leq N} a_n \langle \sigma_i, \psi_n(x) \rangle \\ &= \sum_{n \leq N} \frac{|\langle \sigma_i, \psi_n(x) \rangle_n|^2}{(1 + \lambda_n)^\ell}. \end{aligned}$$

Because $\|\sigma_i\| = 1$, the Cauchy-Schwarz inequality buys us

$$(15.8) \quad \left| \left\langle \sigma_i, \sum_{n \leq N} a_n \psi_n(x) \right\rangle_x \right| \leq \left| \sum_{n \leq N} a_n \psi_n(x) \right|$$

and because $|a_n|^2 = |\langle \sigma_i, \psi_n(x) \rangle|^2 / (1 + \lambda n)^{2\ell}$, then one of our previous estimates (TODO: which? probably (15.5)) implies that

$$(15.9) \quad \left| \sum_{n \leq N} a_n \psi_n(x) \right| \leq C \left(\sum_{n \leq N} \frac{|\langle \sigma_i, \psi_n(x) \rangle_x|^2}{(1 + \lambda n)^\ell} \right).$$

Therefore

$$(15.10) \quad \left| \left\langle \sigma_i, \sum_{n \leq N} a_n \psi_n(x) \right\rangle_x \right| = \sum_{n \leq N} \frac{|\langle \sigma_i, \psi_n(x) \rangle_x|^2}{(1 + \lambda n)^\ell} \leq c \sqrt{\sum_{n \leq N} \frac{|\langle \sigma_i, \psi_n(x) \rangle_x|^2}{(1 + \lambda n)^\ell}}.$$

Therefore

$$(15.11) \quad \frac{1}{C} \sqrt{\sum_{n \leq N} \frac{|\langle \sigma_i, \psi_n(x) \rangle_x|^2}{(1 + \lambda n)^2}} \leq 1,$$

so summing over i ,

$$(15.12) \quad \sum_i \sum_{n \leq N} A_{i,n}(x) \leq (\text{rank } S) C^2,$$

and therefore the infinity sum is bounded by $(\text{vol } M)(\text{rank } S) C^2$. \square

Next we'll discuss finite propagation speed for the wave equation. We need to choose a sign convention in the wave equation, and we choose (TODO: erased before I could write it down).

Remark 15.13. For the wave equation, we need to equate units of length and time. You can either judiciously choose units, or end up with a constant with units length over time. This is the speed of the wave, and the proposition below says that the waves travel at that finite speed.

The heat equation, by contrast, flows compactly supported functions into noncompactly supported functions after any $\varepsilon > 0$. The upshot is that, according to the heat equation, heat flow travels infinitely fast (though it does decay exponentially away from the heat source at small times). \blacktriangleleft

Proposition 15.14. *For any $s \in C_c^\infty(S)$, the support of $e^{itD}s$ lies within distance $|t|$ from the support of s .*

So we've normalized constants such that the speed of the wave is 1.

To prove Proposition 15.14, we'll need an energy estimate. Recall that if ω is the 1-form defined by $\omega(X) = (c(X)s_t, s_t)$, where $c(\cdot)$ is the Clifford action, then

$$(15.15) \quad (iDs_t, s_t) + (s_t, iDs_t) = -id^*\omega.$$

Lemma 15.16 (Energy estimate). *Let $m \in M$, R be less than the injectivity radius of M , and s_t be a solution of the wave equation. Then for $t \geq 0$,*

$$\frac{d}{dt} \int_{B_{R-t}(m)} |s_t|^2 \leq 0.$$

Proof. First observe that at $x \in B_R(m)$,

$$\begin{aligned} \frac{\partial}{\partial t} (|s_t|^2) &= \left(\frac{\partial}{\partial t} s_t, s_t \right) + \left(s_t, \frac{\partial}{\partial t} s_t \right) \\ &= (iDs_t, s_t) + (s_t, iDs_t) = -id^*\omega. \end{aligned}$$

Therefore

$$(15.17) \quad \frac{d}{dt} \int_{B_{R-t}(m)} |s_t|^2 = - \int_{B_{R-t}(m)} id^*\omega - \int_{\partial B_{R-t}(m)} (s_t, s_t) i_{n^*} \text{vol},$$

where i_{n_*} denotes contracting with the normal vector field on $\partial B_{R-t}(m)$. This is an instance of Stokes' theorem. Applying the Divergence theorem, this is equal to

$$(15.18) \quad = \int_{\partial B_{R-t}(m)} |-i(c(N)s_t, s_t)| i_{n_*} \text{vol} - \int_{\partial B_{R-t}(m)} (s_t, s_t).$$

Cauchy-Schwarz shows the first term is bounded above by $|s_t|^2$, so the entire expression is nonpositive, as desired. \square

Proof of Proposition 15.14. Since $\exp(i(t_1 + t_2)D) = \exp(it_1D)\exp(it_2D)$, it suffices to prove this for small $t_0 > 0$; specifically, we'll choose t_0 less than the injectivity radius of M , and suppose that $d(m, \text{supp}(s)) > t_0$. Then $\int_{B_{t_0}(m)} |s_t|^2 \text{dvol} = 0$, so by Lemma 15.16, $\int_{B_{t_0-t}(m)} s_t = 0$ for all $t \in (0, t_0)$. \square

Definition 15.19. A C^∞ function $f : \mathbb{R} \rightarrow \mathbb{R}$ is called *Schwarz* if f and all of its derivatives are rapidly decreasing. The Fréchet space of Schwarz functions is denoted $\mathcal{S}(\mathbb{R})$.

Proposition 15.20. *The Fourier transform of a Schwarz function is Schwarz, and in fact the Fourier transform is an invertible continuous linear map $\mathcal{S}(\mathbb{R}) \rightarrow \mathcal{S}(\mathbb{R})$.*

There are conventions in the definition of the Fourier transform; if we define

$$(15.21a) \quad \widehat{f}(\xi) := \int_{\mathbb{R}} e^{-ix\xi} f(x) dx,$$

then the inverse to the Fourier transform is

$$(15.21b) \quad f(x) = \frac{1}{2\pi} \int_{\mathbb{R}} e^{ix\xi} \widehat{f}(\xi) d\xi.$$

Proposition 15.22. *Let $f \in \mathcal{S}(\mathbb{R})$. Then for all $x, y \in L^2(S)$, we have*

$$\langle f(D)x, y \rangle = \frac{1}{2\pi} \int_{\mathbb{R}} \widehat{f}(\xi) \langle e^{i\xi D} x, y \rangle d\xi$$

and

$$f(D) = \frac{1}{2\pi} \int_{\mathbb{R}} \widehat{f}(\xi) e^{i\xi D} d\xi,$$

where we interpret these in a distributional sense.

Proof. Let s be an eigenfunction of D with eigenvalue λ . Then

$$\begin{aligned} \langle f(D)x, y \rangle &= f(\lambda) \langle x, y \rangle \\ &= \langle x, y \rangle \frac{1}{2\pi} \int_{\mathbb{R}} \widehat{f}(\xi) e^{i\lambda\xi} d\xi \\ &= \frac{1}{2\pi} \int_{\mathbb{R}} \widehat{f}(\xi) \langle e^{i\xi D} x, y \rangle d\xi. \end{aligned}$$

More generally, we could have a section $s = \sum_{\lambda} s_{\lambda}$, and have to justify interchanging the sum and the integral. The trick to do this is to integrate over small regions of the manifold, making everything finite enough to show that the integral converges uniformly in λ , allowing us to switch the sum and the integral. \square

Here's the main result.

Theorem 15.23. *Suppose $f \in \mathcal{S}(\mathbb{R})$ and \widehat{f} is supported in $[-c, c]$. Then $\langle f(D)x, y \rangle = 0$ for sections x, y of S whenever $d(\text{supp}(x), \text{supp}(y)) > c$. Consequently, the smoothing kernel of $f(D)$ is contained in a c -neighborhood of the diagonal in $M \times M$. Consequently, the smoothing kernel of $f(D)$ is contained in a*

Lecture 16.

Asymptotics of the counting function: 3/21/18

Kenny spoke about asymptotics of the counting function for eigenvalues of the Dirac operator.

Let M be a compact, oriented n -manifold with a Riemannian metric. We know its Laplacian has a discrete spectrum and its eigenvalues tend to infinity.

Definition 16.1. The *counting function* $\mathfrak{N}: \mathbb{R} \rightarrow \mathbb{N}$ sends λ to the number of eigenvalues of Δ less than or equal to λ .

Our first crude estimate will come from the Sobolev embedding theorem. Let λ' be an eigenvalue with $\lambda' \leq \lambda$, s_1, \dots, s_m be orthonormal eigenfunctions for λ' , and $s := \sum a_i s_i$.

Let $k := \lfloor n/2 \rfloor$. Then

$$\begin{aligned}
 |s(x)| &\leq \|s\|_{C^0} \leq C \|s\|_{H^k} \\
 &\leq C (\|s\|_{L^2} + \|D^k s\|_{L^2}) \\
 &\leq C (1 + \lambda^{k/2}) \|s\|_{L^2} \\
 &\leq C (1 + \lambda^{k/2}) \left(\sum_i |a_i|^2 \right)^{1/2} \\
 (16.2) \quad &\leq C (1 + \lambda)^{k/2} \left(\sum_i |a_i|^2 \right)^{1/2}.
 \end{aligned}$$

Supposing we take $a_i := \bar{s}_i(x)$, we can rearrange (16.2) into

$$(16.3) \quad \sum_i |s_i(x)|^2 \leq C^2 (1 + \lambda)^k.$$

Integrating out x over M , we conclude

$$(16.4) \quad j \leq C^2 (1 + \lambda)^k \text{vol}(M).$$

We will be able to do better, though, using traces of operators.

Trace-class operators Let H and H' be separable, infinite-dimensional Hilbert spaces, and choose orthonormal bases $\{e_i\}$ and $\{e'_i\}$ for H and H' , respectively.

Let $A: H \rightarrow H'$ be a bounded operator. If H and H' were finite-dimensional, we would have enough data to represent A as a matrix; in this setting, we use matrix coefficients, which are more general.

Definition 16.5. Associated to $i, j \in \mathbb{N}$, we define the *matrix coefficient*

$$c_{ij}(A) := \langle Ae_i, e'_j \rangle.$$

We will use these to define nice classes of operators.

Definition 16.6. With A, H, H' as above, the *Hilbert-Schmidt norm* of A is

$$\|A\|_{\text{HS}}^2 := \sum_{i,j} |c_{ij}(A)|^2,$$

which may be infinite.

Proposition 16.7. $\|A\|_{\text{HS}}$ is independent of the choice of bases for H and H' .

Proof. We're going to use the following identity a lot today:

$$(16.8) \quad (x, y) = \sum_j (x, e_j)(y, e_j).$$

For example, it implies

$$(16.9) \quad \sum_{i,j} |c_{ij}(A)|^2 = \sum_{i,j} \langle Ae_i, e'_j \rangle \langle e'_j, Ae_i \rangle = \sum_i \|Ae_i\|^2,$$

so it's independent of the basis of H' . Next we observe that, by definition, $c_{ij}(A) = \overline{c_{ji}(A^*)}$, so $\|A\|_{\text{HS}} = \|A^*\|_{\text{HS}}$. The same reasoning as above tells us this quantity also doesn't depend on the basis for H . \square

Definition 16.10. If $\|A\|_{\text{HS}}$ is finite, A is called a *Hilbert-Schmidt operator*.

The idea is that these are the operators which can be represented by matrices. These operators have a whole bunch of nice properties.

Proposition 16.11.

(1) The Hilbert-Schmidt norm is induced by the following inner product on $B(H, H')$:

$$(16.12) \quad \langle A, B \rangle_{\text{HS}} := \sum_{i,j} \overline{c_{ij}(A)} c_{ij}(B).$$

(2) The space of Hilbert-Schmidt operators is a Hilbert space with respect to this inner product.

(3) For any $A \in B(H, H')$, $\|A\|_{B(H, H')} \leq \|A\|_{\text{HS}}$.

(4) Hilbert-Schmidt operators are compact.

(5) If A is Hilbert-Schmidt and B is bounded, AB and BA are Hilbert-Schmidt.

You can follow your nose to prove some of these but compactness is tricky: the idea is that your infinite matrix is a limit of finite matrices, which are compact; then one translates this into rigorous language.

Definition 16.13. An operator $T \in B(H)$ is called *trace-class* if there are two Hilbert-Schmidt operators A and B whose product is T . In this case its *trace* is $\text{tr } T := \langle A^*, B \rangle_{\text{HS}}$.

Remark 16.14. Using (16.12), one can quickly compute that the trace is

$$(16.15) \quad \text{tr}(T) = \sum_j c_{jj}(T),$$

so it's the sum of diagonal elements, as expected. But the intrinsic perspective will also be helpful. ◀

One useful perspective is that trace-class operators are a subspace of Hilbert-Schmidt operators, which are a subspace of compact operators, which are a subspace of bounded operators. The matrix coefficients of a trace-class operator form an ℓ^1 sequence, of a Hilbert-Schmidt operator form an ℓ^2 sequence, of a compact operator form a c_0 sequence, and of a bounded operator form an ℓ^∞ sequence. Thus this follows from the inclusions

$$(16.16) \quad \ell^1 \subset \ell^2 \subset c_0 \subset \ell^\infty.$$

Proposition 16.17. Let T be a self-adjoint, trace-class operator with eigenvalues λ_i . Then $\text{tr}(T) = \sum \lambda_i$.

Again, this is a generalization of something we knew in finite dimension.

Proposition 16.18. Let $T, B: H \rightarrow H$ be bounded operators, and assume either

- T is trace-class, or
- both T and B are Hilbert-Schmidt.

In this case, $\text{tr}(TB) = \text{tr}(BT)$.

In other words, trace-class operators are an ideal inside the algebra of bounded operators under composition. Hilbert-Schmidt operators are also an ideal, as are compact operators.

Proof. We know TB and BT are trace-class from above. Choose an orthonormal basis $\{e_i\}$ for H . Then we can compute using (16.8):

$$\begin{aligned} \text{tr}(TB) &= \sum_i \langle TBe_i, e_i \rangle \\ &= \sum_i \langle Be_i, T^*e_i \rangle \\ &= \sum_{i,j} \overline{c_{ij}(B)} \overline{c_{ij}(T)} \\ &= \text{tr}(BT). \end{aligned}$$

◻

For us, the most important Hilbert-Schmidt and trace-class operators arise from integral operators on manifolds.

Proposition 16.19. *Let M be a compact manifold and ω be a smooth volume form on M . If $k: M \times M \rightarrow \mathbb{C}$ is continuous and $A: L^2(M) \rightarrow L^2(M)$ defined by*

$$(16.20) \quad Au(m) := \int_M k(m, m') u(m') \omega_{m'}$$

is bounded, then A is Hilbert-Schmidt and

$$(16.21) \quad \|A\|_{\text{HS}}^2 = \iint |k(m, m')|^2 \omega_m \omega_{m'}.$$

Proof. Let $\{e_i\}$ be an orthonormal basis for $L^2(M)$. Then

$$\begin{aligned} \|A\|_{\text{HS}}^2 &= \sum_j \|Ae_j\|_{L^2}^2 \\ &= \sum_j \int_M \left| \int_M k(m, m') e_j(m') \omega_{m'} \right|^2 \omega_m \\ &= \int_M \sum_j \left| \int_M k(m, m') e_j(m') \omega_{m'} \right|^2 \omega_m. \end{aligned}$$

Using (16.8),

$$(16.22) \quad \sum_j \left| \int_M k(m, m') e_j(m') \omega_{m'} \right|^2 = \int_M |k(m, m')|^2 \omega_{m'},$$

so

$$(16.23) \quad \|A\|_{\text{HS}}^2 = \iint |k(m, m')|^2 \omega_m \omega_{m'} < \infty. \quad \square$$

If we make k better, we make A better.

Proposition 16.24. *With M and A as in Proposition 16.19, if k is smooth, then A is trace-class, and*

$$\text{tr}(A) = \int_M k(m, m) \omega_m.$$

Proof. First suppose $A = BC$ for Hilbert-Schmidt operators B and C represented by continuous kernels k_B and k_C , respectively. Then

$$(16.25) \quad k(m, m'') = \int_M k_B(m, m') k_C(m', m'') \omega_{m'}.$$

(This is akin to the formula for matrix multiplication in coordinates.)

For an integral operator with continuous kernel, we've already found an explicit formula for the Hilbert-Schmidt norm, namely (16.23). Therefore, by a polarization identity, this expression determines the inner product on the space of integral operators with continuous kernel:

$$(16.26) \quad \text{tr}(A) = \langle B^*, C \rangle_{\text{HS}} = \iint k_B(m, m') k_C(m', m) \omega_m \omega_{m'} = \int_M k(m, m) \omega_m.$$

It therefore suffices to prove that every smoothing operator A factors as $A = BC$, where B and C are as above. We've already seen that $(1 + \Delta)^{-N}$ is an integral operator with continuous kernel for N sufficiently large, and that it's Hilbert-Schmidt. Therefore we can let $B := (1 + \Delta)^{-N}$ and $C := (1 + \Delta)^N A$. \square

Thus far we've only dealt with scalar functions, but for applications we'll need to consider operators on spaces of sections of the vector bundle $S \boxtimes S^* \rightarrow M$. Most of the results go through quite nicely.

Proposition 16.27. *Let A be a smoothing operator on $L^2(S)$ with kernel k , where S is a complex vector bundle with a Hermitian metric. Then A is trace-class and*

$$\text{tr}(A) = \int_M \text{tr} k(m, m) \omega_m.$$

Roe suggests proving this using a partition of unity.

Proposition 16.28. *Let $A: H \rightarrow H'$ be an operator and $\{e_i\}$ be a basis of H' . Then A is Hilbert-Schmidt iff $\{\|Ae_n\|_{H'}\} \in \ell^2(\mathbb{R})$.*

This is the idea of approximating by matrices we mentioned before. You can use this to prove that if T is a bounded operator, $\|TAe_n\|_{H'} \leq C\|Ae_n\|_{H'}$, so TA is also Hilbert-Schmidt, and this is a somewhat more flexible approach for sections of a bundle.

Lecture 17.

Geometry from the spectrum of the Laplacian: 3/21/18

In this part of the lecture, Ivan spoke. Throughout this section, M will be a closed, oriented n -manifold, Δ will denote the Laplacian on $L^2(M)$, and $\lambda_1, \lambda_2, \dots$ will be its eigenvalues, in nondecreasing order. The eigenvalue counting function for Δ will be denoted \mathfrak{N} . The theme of this lecture is the relationship between the spectrum of Δ and the geometry of M .

Recall that if $S \rightarrow M$ is a Clifford bundle with Dirac operator D and Laplacian $\Delta = D^2$, then the heat kernel $k_t \in C_M^\infty(S \boxtimes S^*)$ is the kernel for the operator $e^{-t\Delta}$. We've seen the asymptotic expansion of the heat kernel,

$$(17.1) \quad k_t(p, q) \sim_{t \rightarrow 0^+} h_t(p, q)(\Theta_0(p, q) + t\Theta_1(p, q) + \dots)$$

on $C_M^k(S \boxtimes S^*)$ for every k . In particular, on the diagonal,

$$(17.2) \quad k_t(p, p) \sim_{t \rightarrow 0^+} \frac{1}{(4\pi t)^{n/2}} \left(1 + \frac{t}{6}(\kappa(p)1 - \mathbb{K}(p)) + \dots \right).$$

Here κ is the scalar curvature and \mathbb{K} is the Clifford-contracted curvature.

Taking $S := \Lambda^*(T^*M \otimes \mathbb{C}) \rightarrow M$ with $D := d + d^*$, then restricting to $\Lambda^0(T^*M \otimes \mathbb{C}) \rightarrow M$, (17.2) simplifies to

$$(17.3) \quad k_t(p, p) \sim_{t \rightarrow 0^+} \frac{1}{(4\pi t)^{n/2}} \left(1 + \frac{t}{6}\kappa(p) + \dots \right).$$

This asymptotic expansion is uniform in p , so we can integrate over M , obtaining

$$(17.4) \quad \int_M k_t(p, p) dp \sim_{t \rightarrow 0^+} \frac{1}{(4\pi t)^{n/2}} \left(\text{vol}(M) + \frac{t}{6} \int_M \kappa(p) dp + \dots \right).$$

Proposition 17.5. *The spectrum of Δ determines $\text{vol}(M)$ and the total scalar curvature $\int_M \kappa(p) dp$. By the Gauss-Bonnet formula, if $\dim M = 2$, then it determines the topology.*

This follows from the previous calculation.

Remark 17.6. The question of whether the spectrum determines M up to isometry is the subject of a famous paper by Mark Kac called “Can you hear the shape of a drum?” The idea is that the frequencies produced by a drumhead arise as eigenvalues of its Laplacian. There is a rich history of research; the original question has a negative answer, but the counterexamples are beautiful too, relating to self-dual lattices. One can also get stronger statements by twisting the Laplacian. \blacktriangleleft

Theorem 17.7 (Weyl's asymptotic formula).

$$(17.8a) \quad \mathfrak{N}(\lambda) \sim_{\lambda \rightarrow \infty} \frac{\text{vol}(M)\lambda^{n/2}}{\Gamma(n/2 + 1)(4\pi)^{n/2}}.$$

Equivalently,

$$(17.8b) \quad \lim_{\lambda \rightarrow \infty} \lambda^{-n/2} \mathfrak{N}(\lambda) = \frac{\text{vol}(M)}{\Gamma(n/2 + 1)(4\pi)^{n/2}}.$$

The gamma function arises because one is looking at the volume of an n -dimensional unit ball.

Example 17.9. Let $M := S^1$ with length 2π . Then $\Delta = -\frac{\partial^2}{\partial \theta^2}$. The theory of Fourier series tells us that $\{e^{in\theta} \mid n \in \mathbb{Z}\}$ is an eigenbasis for $L^2(S^1)$, with eigenvalues $0, 1, 1, 4, 4, 9, 9, \dots$. Hence the counting function is

$$(17.10) \quad \mathfrak{N}(\lambda) = 1 + 2\lfloor \lambda^{1/2} \rfloor.$$

The right-hand side of (17.8a) is

$$(17.11) \quad \frac{\text{vol}(M)\lambda^{n/2}}{\Gamma(n/2+1)(4\pi)^{n/2}} = \frac{2\pi\lambda^{1/2}}{\Gamma(3/2)2\pi^{1/2}} = 2\sqrt{\lambda},$$

and clearly $1 + 2\lfloor\sqrt{\lambda}\rfloor \sim_{\lambda \rightarrow \infty} 2\sqrt{\lambda}$. ◀

Example 17.12. Let $M = \mathbb{R}^2/(2\pi\mathbb{Z})^2$, which is a flat torus; its Laplacian is

$$(17.13) \quad \Delta = -\frac{\partial^2}{\partial \theta_1^2} - \frac{\partial^2}{\partial \theta_2^2}.$$

The set $\{\exp(i(\theta_1 n + \theta_2 m)) \mid (n, m) \in \mathbb{Z}^2\}$ is an eigenbasis for $L^2(M)$, and its eigenvalues are $n^2 + m^2$ for $m, n \in \mathbb{Z}$. Therefore

$$(17.14) \quad \mathfrak{N}(\lambda) = \#\{m, n \in \mathbb{Z} \mid m^2 + n^2 \leq \lambda\}.$$

Finding a closed form for this function is a classical problem, called *Gauss' circle problem*. A closed form does exist, but it's messy. Nonetheless, we can see vaguely how it behaves.

The right-hand side of (17.8a) is

$$(17.15) \quad \frac{4\pi^2\lambda}{\Gamma(2)4\pi} = \pi\lambda,$$

which is in particular the area of a disk of radius $\sqrt{\lambda}$. One can show that $\mathfrak{N}(\lambda) = \pi\lambda + E(\lambda)$, for some $E(\lambda)$ which is $o(\lambda)$ (i.e. $\lim_{\lambda \rightarrow \infty} E(\lambda)/\lambda = 0$). ◀

Proof of Theorem 17.7. First, an outline of the proof.

(1) We'll study $\mathfrak{N}(\lambda)$ using $\sum_j e^{-t\lambda_j}$.

(2) If $f: [0, 1] \rightarrow \mathbb{R}$ is continuous, we can define $\varphi_f: \mathbb{R}^+ \rightarrow \mathbb{C}$ by

$$(17.16) \quad \varphi_f(t) := \sum_j f(e^{-t\lambda_j})e^{-t\lambda_j}.$$

(3) Then, we'll show that $\varphi_f(t)$ has a few nice properties: first, that for a suitable choice of f we can use it to bound \mathfrak{N} ; and second, that $\lim_{t \rightarrow 0^+} t^{n/2}\varphi_f(t)$ will exist and have a concrete expression.

(4) We will use this to show that $\lim_{\lambda \rightarrow \infty} \lambda^{-n/2}\mathfrak{N}(\lambda)$ exists and is given by (17.8a).

The fact that

$$(17.17a) \quad \sum_j e^{-t\lambda_j} \sim_{t \rightarrow 0^+} \frac{1}{(4\pi t)^{n/2}}(\text{vol}(M) + \dots)$$

implies

$$(17.17b) \quad \lim_{t \rightarrow 0^+} t^{n/2} \sum_j e^{-t\lambda_j} = \frac{\text{vol}(M)}{(4\pi)^{n/2}}.$$

Therefore if $A := \text{vol}(M)/(4\pi)^{n/2}$ and $\alpha := n/2$, we want to show that

$$(17.18) \quad \lim_{\lambda \rightarrow \infty} \lambda^{-\alpha} \mathfrak{N}(\lambda) = \frac{A}{\Gamma(\alpha+1)}.$$

Lemma 17.19.

$$\lim_{t \rightarrow 0^+} t^\alpha \varphi_f(t) = \frac{A}{\Gamma(\alpha)} \int_0^\infty f(e^{-s}) s^{\alpha-1} e^{-s} ds.$$

Proof. By the Stone-Weierstrass theorem, it suffices to prove the lemma when f is a polynomial. In fact, since the result is linear in f , we can just prove it for $f(x) = x^n$. In this case the left-hand side of Lemma 17.19 behaves as

$$(17.20) \quad t^\alpha \varphi_f(t) = t^\alpha \sum_j e^{-t(n+1)\lambda_j} \longrightarrow_{t \rightarrow 0^+} A(n+1)^{-\alpha},$$

and the right-hand side behaves as

$$(17.21) \quad \frac{A}{\Gamma(\alpha)} \int_0^\infty e^{-s(n+1)} s^{\alpha-1} ds = \frac{A}{\Gamma(\alpha)(n+1)^\alpha} \underbrace{\int_0^\infty e^{-\tilde{s}} \tilde{s}^{\alpha-1} d\tilde{s}}_{\Gamma(\alpha)} = A(n+1)^{-\alpha},$$

where $\tilde{s} := (n+1)s$. \(\square\)

For $0 < r < 1$, let $f_r : [0, 1] \rightarrow \mathbb{R}$ be defined to be 0 on $[0, r/e]$, $1/x$ on $[1/e, 1]$, linear on $[r/e, 1]$, and continuous. What can we say about $\varphi_{f_r}(t)$? First, the sum

$$(17.22) \quad \varphi_{f_r}(t) := \sum_j f_r(e^{-t\lambda_j}) e^{-t\lambda_j}$$

is finite!

If $e^{-t\lambda_j} \in [1/e, 1]$, then $f_r(e^{-t\lambda_j}) e^{-t\lambda_j} = 1$, and if $e^{-t\lambda_j} \in [r/e, 1/e]$, then $f_r(e^{-t\lambda_j}) e^{-t\lambda_j} \leq 1$. In particular,

$$(17.23) \quad \varphi_{f_r}\left(\frac{1}{\lambda}\right) = \sum_j f_r(e^{-\lambda_j/\lambda}) e^{-\lambda_j/\lambda},$$

which means the following are equivalent:

- $\lambda_j \leq \lambda$,
- $e^{-\lambda_j/\lambda} \geq 1/e$, and
- $f_r(e^{-\lambda_j/\lambda}) e^{-\lambda_j/\lambda} = 1$.

Therefore

$$(17.24) \quad \mathfrak{N}(\lambda) \leq \varphi_{f_r}\left(\frac{1}{\lambda}\right).$$

Now, for $r \in (0, 1)$, we consider $\varphi_{f_r}(1/r\lambda)$. A similar analysis shows the following are equivalent:

- $f_r(e^{-\lambda_j/r\lambda}) > 0$,
- $e^{-\lambda_j/r\lambda} > r/e = e^{-1+\ln r}$, and
- $\lambda > (1 - \ln r)r\lambda > \lambda_j$.

Since $(1 + \ln r)r < 1$, this combines with (17.24) to prove

$$(17.25) \quad \varphi_{f_r}\left(\frac{1}{r\lambda}\right) \leq \mathfrak{N}(\lambda) \leq \varphi_{f_r}\left(\frac{1}{\lambda}\right).$$

Therefore

$$\begin{aligned} \limsup_{\lambda \rightarrow \infty} \lambda^{-\alpha} \mathfrak{N}(\lambda) &\leq \lim_{\lambda \rightarrow \infty} \lambda^{-\alpha} \varphi_{f_r}\left(\frac{1}{\lambda}\right) \\ &= \frac{A}{\Gamma(\alpha)} \int_0^{1-\ln r} f_r(e^{-s}) s^{\alpha-1} e^{-s} ds \\ &\leq \frac{A}{\Gamma(\alpha)} \int_0^{1-\ln r} s^{\alpha-1} ds \\ &= \frac{A(1 - \ln r)^\alpha}{\Gamma(\alpha + 1)}. \end{aligned}$$

Therefore

$$(17.26) \quad \limsup_{\lambda \rightarrow \infty} \lambda^{-\alpha} \mathfrak{N}(\lambda) \leq \frac{A}{\Gamma(\alpha + 1)}.$$

Using the other inequality in (17.25),

$$\begin{aligned} \liminf_{\lambda \rightarrow \infty} \lambda^{-\alpha} \mathfrak{N}(\lambda) &\geq \lim_{\lambda \rightarrow \infty} \varphi_{f_r} \left(\frac{1}{r\lambda} \right) \\ &= \frac{r^\alpha A}{\Gamma(\alpha)} \int_0^{1-\ln r} f_r(e^{-s} s^{\alpha+1} e^{-s}) ds \\ &\geq \frac{r^\alpha A}{\Gamma(\alpha)} \int_0^1 s^{\alpha-1} ds \\ &= \frac{r^\alpha A}{\Gamma(\alpha+1)}. \end{aligned}$$

we conclude

$$(17.27) \quad \liminf_{\lambda \rightarrow \infty} \lambda^{-\alpha} \mathfrak{N}(\lambda) \geq \frac{A}{\Gamma(\alpha+1)},$$

so in fact

$$(17.28) \quad \lim_{\lambda \rightarrow \infty} \lambda^\alpha \mathfrak{N}(\lambda) = \frac{A}{\Gamma(\alpha+1)}$$

as desired. ⊠

If we're interested in sections of a Clifford bundle $S \rightarrow M$ rather than functions, the same techniques apply, and we obtain the asymptotic

$$(17.29) \quad \mathfrak{N}(\lambda) \sim_{\lambda \rightarrow \infty} \frac{\text{vol}(M) \text{rank}(S) \lambda^{n/2}}{(4\pi)^{n/2} \Gamma(n/2 + 1)}.$$

Lecture 18.

The harmonic oscillator: 3/28/18

Today, Sebastian spoke about the harmonic oscillator on \mathbb{R} .

Recall that the Dirac operator on $S^1 = \mathbb{R}/2\pi\mathbb{Z}$ is $D := i \frac{d}{dx}$. Fourier theory tells us this has a discrete spectrum and finite-dimensional eigenspaces. Everything is nice.

But if we try to set this up on a noncompact manifold, such as \mathbb{R} , the Fourier series is replaced with the Fourier transform. Therefore the Dirac operator has a continuous spectrum, and its “eigenfunctions” are harder to speak about (in particular, they're not L^2). But there is still a spectral decomposition.

From a physics perspective, the issue manifests as propagating waves that are noncompactly supported (e.g. infinitely many peaks), hence aren't L^2 . One way to fix this is to add a potential to the Dirac operator, which bounds the low-energy solutions. One can then hope to recover a discrete spectrum.

Definition 18.1. The *harmonic oscillator* is the unbounded operator $H: \mathcal{S}(\mathbb{R}) \rightarrow \mathcal{S}(\mathbb{R})$ sending

$$\psi \mapsto \left(-\frac{d^2}{dx^2} + a^2 x^2 \right) \psi,$$

for some $a > 0$.

One can describe this in terms of the *annihilation operator*

$$(18.2a) \quad A := ax + \frac{d}{dx}$$

and the *creation operator*

$$(18.2b) \quad A^* := ax - \frac{d}{dx}.$$

Using integration by parts, you can prove these are adjoints.

Namely,

$$\begin{aligned}
 AA^*\psi &= \left(ax + \frac{d}{dx}\right)(ax\psi - \psi') \\
 &= a^2x^2\psi + a\psi + ax\psi' - ax\psi' + \psi'' \\
 &= \left(a^2x^2 - \frac{d^2}{dx^2}\right)\psi + a\psi \\
 &= (H + a)\psi.
 \end{aligned}$$

This is one of several *canonical commutation relations* between these operators: others you can prove yourself are $A^*A = H - a$, $[A, A^*] = 2a$, $[H, A] = -2aA$, and $[H, A^*] = 2aA^*$. An algebra generated by operators satisfying these relations is called a *CCR algebra*.

Definition 18.3. The *ground state* of H is the $\psi_0 \in L^2(\mathbb{R})$ with $A\psi_0 = 0$ and $\|\psi_0\| = 1$.

This implies that $0 = A^*A\psi_0 = (H - a)\psi_0$, so

$$(18.4) \quad H\psi_0 = a\psi_0.$$

H is nice enough that we can explicitly solve this:

$$\begin{aligned}
 0 &= \left(ax + \frac{d}{dx}\right)\psi_0 \\
 0 &= ax\psi_0 + \psi_0' \\
 \psi_0 &= Ce^{-ax^2/2},
 \end{aligned}$$

and requiring $\|\psi_0\| = 1$, we get $C = a^{1/2}\pi^{1/4}$.

Definition 18.5. The *excited states* are

$$\psi_k = \frac{1}{(2ka)^{1/2}} A^k \psi_{k-1},$$

for $k \geq 1$.

Lemma 18.6. Each $\psi_k \in \mathcal{S}(\mathbb{R})$ and has norm 1. Moreover, $H\psi_k = (2k + 1)a\psi_k$.

Proof. You can just check inductively, e.g.

$$\begin{aligned}
 H\psi_k &= \frac{1}{(2ka)^{1/2}} HA^k \psi_{k-1} \\
 &= \frac{1}{(2ka)^{1/2}} (A^*H + 2aA^*) \psi_{k-1} \\
 &= \frac{1}{(2ka)^{1/2}} ((2k - 1)aA^* + 2aA^*) \psi_{k-1} \\
 &= (2k + 1)a \frac{1}{(2ka)^{1/2}} A^k \psi_{k-1} \\
 &= (2k + 1)a\psi_k.
 \end{aligned}$$

The base case is (18.4). ☒

Lemma 18.7. There is a degree- k polynomial $h_k(x)$ with positive leading coefficient such that

$$\psi_k(x) = h_k(x)e^{-ax^2/2}.$$

Again the proof is by induction. These h_k are (up to rescaling) the *Hermite polynomials*. This implies in particular that the ψ_k span the space \mathcal{P} of maps $x \mapsto p(x)e^{-ax^2/2}$, where $p \in \mathbb{R}[x]$.

Proposition 18.8. $\mathcal{P} \subset L^2(\mathbb{R})$ is dense.

Proof. Without loss of generality, we take $a = 1$, and let $f_j := x^j e^{-x^2/2}$. Then

$$\|f_j\|^2 = \int_{-\infty}^{\infty} x^{2j} e^{-x^2} dx.$$

Letting $y := x^2$,

$$\begin{aligned} &= \int_0^\infty y^j e^{-y} dy \\ &= \Gamma\left(j + \frac{1}{2}\right) \leq j!. \end{aligned}$$

Since

$$(18.9) \quad e^{i\lambda x - x^2/2} = \sum_{j=0}^{\infty} \frac{(i\lambda)^j}{j!} f_j(x),$$

and the j^{th} term has L^2 -norm bounded above by

$$(18.10) \quad |\lambda|^j \cdot (j!)^{-1} \cdot (j!)^{1/2} \leq |\lambda|^j \cdot (j!)^{-1/2}.$$

Since each $f_j \in \mathcal{D}$, $e^{i\lambda x - x^2/2} \in \overline{\mathcal{D}}$.

Suppose $f \in \overline{\mathcal{D}}^\perp$. Then

$$(18.11) \quad 0 = \langle f, e^{i\lambda x - x^2/2} \rangle = \int_{-\infty}^{\infty} \overline{f}(x) e^{i\lambda x - x^2/2} dx.$$

This is the Fourier transform of $\overline{f}(x)e^{-x^2/2}$, so by Plancherel's theorem, $\overline{f}(x)e^{-x^2/2}$ vanishes almost everywhere, and therefore so does $\overline{f}(x)$. \square

This all looked like a bunch of analysis, so what have we accomplished? Actually something cool: $L^2(\mathbb{R})$ admits a decomposition into orthogonal, one-dimensional subspaces for H with discrete spectrum tending to infinity. To the eyes of spectral theory it looks a lot like a Dirac operator on a compact manifold! So we can work through the same story about the heat kernel and all of that.

Lemma 18.12. *Let $u \in L^2(\mathbb{R})$. Then, $u \in \mathcal{S}(\mathbb{R})$ iff $a_k := \langle \psi_k, u \rangle$ is rapidly decreasing.*

Proposition 18.13. *Let f be a bounded function on $\sigma(H)$. Then $f(H)$ is defined and is a bounded operator on $L^2(\mathbb{R})$. The map $f \mapsto f(H)$ is a ring homomorphism from bounded functions on $\sigma(H)$ to $B(L^2(\mathbb{R}))$, and $f(H)$ maps $\mathcal{S}(\mathbb{R})$ to $\mathcal{S}(\mathbb{R})$.*

In this setting we have the *heat equation*

$$(18.14) \quad \frac{\partial u}{\partial t} + Hu = 0.$$

The above Hilbert space theory implies the existence of a solution operator e^{-tH} . Moreover, there is a *heat kernel* $k_t^H \in \mathcal{S}(\mathbb{R} \times \mathbb{R})$ such that

$$(18.15) \quad e^{-tH}u(x) = \int_{\mathbb{R}} k_t^H(x, y)u(y) dy,$$

as in the compact case, and again it's characterized by the following properties.

- (1) k_t^H is Schwarz.
- (2) k_t^H satisfies the heat equation.
- (3) k_t^H tends to a δ -function as $t \rightarrow 0$.

In our proof of the index theorem, we will need an expression for $u(x, t) := k_t^H(x, 0)$. We will guess that it's of the form $\alpha(t) \exp(-(1/2)\beta(t)x^2)$, which, through the heat equation, implies that

$$(18.16) \quad \begin{aligned} \dot{\beta} &= 2(\alpha^2 - \beta^2) \\ \dot{\alpha} &= -\beta\alpha, \end{aligned}$$

and $\beta(t) \sim 1/t$ as $t \rightarrow 0$. Therefore as $t \rightarrow 0$,

$$(18.17) \quad u(x, t) = \frac{1}{\sqrt{4\pi t}} e^{-x^2/4t},$$

and therefore

$$(18.18) \quad \frac{1}{\sqrt{4\pi t}} \int e^{-x^2/4t} s(x) dx \xrightarrow{t \rightarrow 0} s(0).$$

This, plus a lot more computation, yields the following result.

Proposition 18.19 (Mehler's formula).

$$u(x, t) = \sqrt{\frac{a}{2\pi \sinh(at)}} \exp\left(\frac{-ax^2 \coth(at)}{2}\right).$$

The heat kernel is the product of this and the term

$$(18.20) \quad \exp\left(\frac{-a(x^2 + y^2) \coth(2at) + 2 \operatorname{csch}(2at) \cdot xy}{a}\right).$$

It suffices to check that

$$(18.21) \quad e^{-ta} \psi_0(x) = \int k_t^H(x, y) \psi_0(y) dy,$$

which is involved but simple in concept. Then, one can address the ψ_k by the prefactor A^* , though there are details to be understood.

Remark 18.22. The expression in Proposition 18.19 makes sense for some complex a , except if a is a pole for $u(x, t)$. This analytic continuation is not immediate, but will be important in the proof of the index theorem. Concretely, there's an open neighborhood of \mathbb{R} in \mathbb{C} on which the function is holomorphic. ◀

Lecture 19.

Witten's deformation of the de Rham complex: 3/28/18

These are Arun's lecture notes for his lecture on the second part of Chapter 9 in Roe's book.

Throughout (M, g) is a complete Riemannian manifold and $h: M \rightarrow \mathbb{R}$ is a smooth function.

Remark 19.1. Recall that if M is compact, we can define a Clifford bundle and Dirac operator using Hodge theory: let $S := \Lambda^*(T^*M) \otimes \mathbb{C} \rightarrow M$, which is $\mathbb{Z}/2$ -graded by mod 2 degree. This has an action by $\mathcal{Cl}(TM)$: if $v \in \Gamma(TM)$ and $\varphi \in \Gamma(S)$, then

$$(19.2) \quad v \cdot \varphi := v^\vee \wedge \varphi - v \lrcorner \varphi.$$

The Dirac operator is $d + d^*$, and $D^2 = dd^* + d^*d$ is the Hodge Laplacian. ◀

Today, we're going to use h to deform this story.

Remark 19.3. Before we do so, we need to remark on a sign convention. Roe defines the interior product as

$$(19.4) \quad \iota_\xi \omega := (-1)^{n|\omega|+n+1} \star(\xi \wedge \star \omega),$$

while I'm used to the opposite sign in front. This has some unusual results; I tried to work out the differences and got confused. Using this definition, $[\iota_\xi, \varepsilon_\eta] = -\langle \xi, \eta \rangle$, which again is the opposite sign from what I'm used to. One consequence of this is that the adjoint of ε_ξ is $-\iota_\xi$.

TODO: double-check this if time. ◀

Definition 19.5. Let $s \in \mathbb{R}$ The perturbed exterior derivative $d_s: \Omega^*(M) \rightarrow \Omega^{*+1}(M)$ given by h is

$$(19.6) \quad d_s := d + \varepsilon_s dh.$$

It also has the formula $d_s \omega = e^{-sh} d(e^{sh} \omega)$. Its adjoint is

$$(19.7) \quad d_s^* := d^* - \iota_s dh,$$

which has the formula $d_s^* \omega = e^{sh} d^*(e^{-sh} \omega)$.

Hence the perturbed de Rham-Dirac operator is

$$(19.8) \quad D_s := d_s + d_s^* = D + sR,$$

where $R := \varepsilon_{dh} - \iota_{dh}$.

Recall that the *Hessian* of f , $H_f \in \text{Sym}^2(T^*M)$, is the form given in coordinates by the matrix of partial derivatives of h in a synchronous orthonormal frame.

Lemma 19.9. *For any $X, Y \in T_x M$,*

$$(19.10) \quad H_f(X, Y) = (\mathcal{L}_X \mathcal{L}_Y f)(x) - \mathcal{L}_{\nabla_X Y}(f).$$

Proof.

□

Definition 19.11. Define the operator $H_h \in \text{End}(\Lambda^*(T^*M))$ by

$$H_h := -\mathcal{L}_{dh} - (\mathcal{L}_{dh})^*.$$

If $e \in TM$, let $L_e, R_e \in \text{End}(\Lambda^*(T^*M))$ be defined by left and right Clifford multiplication by e , i.e. $e \cdot \omega$ and $(-1)^{|\omega|} \omega \cdot e$, so that $[L_e, R_{e'}] = 0$.

Lemma 19.12. *If $\{e_i\}$ is a local orthonormal frame, then*

$$H_h := \sum_{i,j} H_h(e_i, e_j) L_{e_i} R_{e_j}.$$

I wasn't able to show this — it should boil down to a computation.

Lemma 19.13.

$$(1) \quad R^2 = |dh|^2.$$

$$(2) \quad [R, D] = H_h.$$

Since R and D are odd, this is a supercommutator.

Proof. We make use of the fact that $[\iota_\xi, \varepsilon_\eta] = -\langle \xi, \eta \rangle$. Thus

$$\begin{aligned} R^2 &= (\varepsilon_{dh} - \iota_{dh})(\varepsilon_{dh} - \iota_{dh}) \\ &= \underbrace{\varepsilon_{dh}^2}_{=0} - \underbrace{(\iota_{dh} \varepsilon_{dh} + \varepsilon_{dh} \iota_{dh})}_{[\iota_{dh}, \varepsilon_{dh}]} + \underbrace{\iota_{dh}^2}_{=0} \\ &= \langle dh, dh \rangle = |dh|^2. \end{aligned}$$

For the second part,

$$\begin{aligned} [R, D] &= [\varepsilon_{dh} - \iota_{dh}, d + d^*] \\ &= [\varepsilon_{dh}, d] + [\varepsilon_{dh}, d^*] - [\iota_{dh}, d] - [\iota_{dh}, d^*]. \end{aligned}$$

Since dh is exact, $[\varepsilon_{dh}, d] = 0$, and therefore its adjoint $[\iota_{dh}, d^*] = 0$. Cartan's formula tells us that $[\iota_\eta, d] = \mathcal{L}_\eta$, and therefore $[\varepsilon_{dh}, d^*] = -(\mathcal{L}_\eta)^*$. □

Corollary 19.14.

$$D_s^2 = D^2 + s^2 |dh|^2 + sH_h.$$

Proof. Again, this is a quick computation:

$$(19.15) \quad D_s^2 = D^2 + \underbrace{sDR + sRD}_{s[R, D]} + s^2 R^2 = D^2 + sH_h + s^2 |dh|^2. \quad \square$$

Example 19.16. Consider Euclidean space \mathbb{E}^n (\mathbb{R}^n with the standard metric), and let

$$(19.17) \quad h := \frac{1}{2} \sum \lambda_j x_j^2.$$

In coordinates, $D^2 = -\sum \partial_j^2$, $|dh|^2 = \sum (\lambda_j x_j^j)^2$, and H_h is $\sum \lambda_j Z_j$, where $Z_j : \Lambda^*(T^*\mathbb{E}^n) \rightarrow \Lambda^*(T^*\mathbb{E}^n)$ is $-\iota_{dx^j} \varepsilon_{dx^j}$, the operator which on a basis element $dx_I := dx_{i_1} \wedge \cdots \wedge dx_{i_m}$ is equal to 1 if $j \in I$ and -1 otherwise. TODO: is that right?

Therefore by Corollary 19.14, the perturbed Laplacian is

$$(19.18) \quad D_s^2 = \sum_j \left(-\partial_j^2 + s^2(\lambda_j x^j)^2 + s\lambda_j Z_j \right).$$

The first two terms look just like the harmonic oscillator. ◀

Proposition 19.19. *With \mathbb{E}^n and h as above, then for $s > 0$, there's a basis of $L^2(\mathbb{R}^n)$ of smooth, rapidly decaying eigenfunctions for D_s^2 , and the eigenvalues are*

$$(19.20) \quad s \sum_j (|\lambda_j|(1 + 2p_j) + \lambda_j q_j),$$

where $p_j = 0, 1, 2, \dots$ and $q_j \in \{\pm 1\}$. Restricted to k -forms exactly k of the q_j act as $+1$.

Proof. Let

$$(19.21) \quad Y_j := -\partial_j^2 + s^2 \lambda_j^2 (x^j)^2,$$

which is a harmonic oscillator for x_j .

Lemma 19.22. $[Z_i, Y_j] = 0$, and in particular they can be simultaneously diagonalized.

As we saw in the first half of this chapter, $\sum_j Y_j$ is essentially self-adjoint with a discrete spectrum, and its eigenvalues are

$$(19.23) \quad s \sum_j |\lambda_j|(1 + 2p_j),$$

with multiplicity $\dim \Lambda^*(T^*\mathbb{E}^n) = 2^n$. Each Z_j acts on these eigenspaces by an involution, hence has ± 1 -eigenspaces. If Z_j acts by $q_j \in \{\pm 1\}$, we get a factor of $s\lambda_j q_j$, so the eigenvalue of something in this eigenspace for $\sum Y_j$ and the q_j -eigenspace for Z_j is the formula in (19.20).

TODO: what about if we act on k -forms? ◻

Functional calculus on noncompact manifolds Let D be a Dirac operator (the results of this section also apply to generalized Dirac operators) on M . We'll describe a way to build, from a suitably regular function f , an operator $f(D)$, such that the assignment $f \mapsto f(D)$ is a ring homomorphism with nice properties analogous to the compact case we've already discussed. One way to do this is to invoke the spectral theorem for self-adjoint unbounded operators, but following Roe, we'll use finite propagation of solutions to the wave equation

$$(19.24) \quad \begin{aligned} \frac{\partial s}{\partial t} &= iD_s \\ s(0) &= s_0. \end{aligned}$$

Proposition 19.25. *Let $s_0 \in C_0^\infty(M)$. Then (19.24) has a unique solution, and for all t , $s(t)$ is smooth and compactly supported.*

Proof. Existence: pick a $t_0 > 0$; we'll construct a solution for all t with $|t| < t_0$. Let $K := \text{supp}(s_0)$ and U be a t_0 -neighborhood of K . We can isometrically embed U into a closed manifold M' : let $X \subset M$ be a compact submanifold, possibly with boundary, containing U , and let M' be the double of X . This construction takes the Clifford bundle $S \rightarrow M$ and makes another Clifford bundle $S' \rightarrow M'$, such that the embedding $i: U \hookrightarrow M'$ induces an isomorphism of Clifford bundles $i^*S' \cong S|_U$. Thanks to hard work from previous sections, we know the wave equation has solutions $s(t)$ on M' with propagation speed 1, so for $|t| < t_0$, $\text{supp}(s(t)) \subset U$, and therefore $s(t)$ makes sense on M too.

Uniqueness follows in a similar way: if we have two solutions $s(t), s'(t)$, they must agree for $t < t_0$, because we proved uniqueness for the heat equation on closed manifolds. ◻

Therefore the solution operator $e^{itD}: C_c^\infty(S) \rightarrow L^2(S)$ exists and is unitary. By density, it extends to a unitary operator on all of $L^2(S)$.

Definition 19.26. For an $f \in \mathcal{S}(\mathbb{R})$, we define

$$f(D) := \frac{1}{2\pi} \int_{\mathbb{R}} \widehat{f}(t) e^{itD} dt.$$

Because f is Schwarz, this converges weakly, i.e.

$$(19.27) \quad \langle f(D)x, y \rangle = \frac{1}{2\pi} \int_{\mathbb{R}} \widehat{f}(t) \langle e^{itD}x, y \rangle dt :$$

it suffices to prove this on an eigenbasis for D , where it reduces to the formula for the inverse Fourier transform.

Proposition 19.28.

- (1) The assignment $f \mapsto f(D)$ is a ring homomorphism $\mathcal{S}(\mathbb{R}) \rightarrow B(L^2(S))$: the former as a ring under pointwise multiplication, and the latter under composition.
- (2) $\|f(D)\|_{B(L^2(S))} \leq \sup|f|$.
- (3) If $f(x) = xg(x)$, then $f(D) = Dg(D)$.

Proof. We reduce to closed manifolds, where this proposition is not new.

It suffices to assume \widehat{f} is compactly supported, because such f are dense in $\mathcal{S}(\mathbb{R})$, and let $\text{supp}(\widehat{f}) \subset [-t_0, t_0]$. Let $s \in C_0^\infty(S)$ and $K = \text{supp}(s)$. As in the previous proof, we can construct a closed manifold M' and an isometric embedding of a t_0 -neighborhood U of K into M' together with a Clifford bundle $S' \rightarrow M'$; let D' denote its Dirac operator.

We showed that the wave equation on M has propagation speed 1, and hence $f(D)s$ is supported in U and is equal to $f(D')s$. Since we know the proposition for $f(D')s$, since M' is closed, we also know it for $f(D)s$. Since such s are dense in $L^2(S)$, the result follows. \square

Remark 19.29. Part (2) implies that $f \mapsto f(D)$ extends uniquely to a map $C_0(\mathbb{R}) \rightarrow B(L^2(S))$; here $C_0(\mathbb{R})$ denotes the space of continuous functions vanishing at ∞ with the sup norm, which contains $\mathcal{S}(\mathbb{R})$ as a dense subspace. \blacktriangleleft

Lecture 20.

The Lefschetz fixed-point theorem: 4/4/18

These are Riccardo's notes on his talk on the Lefschetz formula.

In this chapter we will find our first example of a topological invariant defined by elliptic operators. The topological invariance will come from a pairwise cancellation of eigenspaces, which is sometimes called *supersymmetry*.

Prerequisites.

Definition 20.1. Let S be a bundle of Clifford modules over a Riemannian manifold M . S is a *Clifford bundle* if it is equipped with a Hermitian metric and compatible connection such that

- (1) the Clifford action of each vector $v \in T_m M$ on S_m is skew-adjoint, that is, $(v \cdot s_1, s_2) = -(s_1, v \cdot s_2)$, and
- (2) the connection on S is compatible with the Levi-Civita connection on M , in the sense that $\nabla_X(Ys) = (\nabla_X Y)s + Y \nabla_X s$ for all vector fields X, Y and sections $s \in C^\infty(S)$.

Definition 20.2. The *Dirac operator* D of a Clifford bundle S is the first-order differential operator on $C^\infty(S)$ defined by the following composition:

$$(20.3) \quad C^\infty(S) \rightarrow C^\infty(T^*M \otimes S) \rightarrow C^\infty(TM \otimes S) \rightarrow C^\infty(S),$$

where the first arrow is given by the connection, the second by the metric (identifying TM and T^*M), and the third by the Clifford action on each fiber.

Definition 20.4. Let M be an n -dimensional compact oriented Riemannian manifold, and let S_0, S_1, \dots, S_k be a sequence of vector bundles over M , equipped with Hermitian metrics and compatible connections. Suppose given differential operators $d_j: C^\infty(S_j) \rightarrow C^\infty(S_{j+1})$ in such a way that $d_{j+1}d_j = 0$, i.e. that

$$(20.5) \quad C^\infty(S_0) \xrightarrow{d} C^\infty(S_1) \xrightarrow{d} C^\infty(S_2) \rightarrow \dots \rightarrow C^\infty(S_k)$$

is a complex. It will be called a *Dirac complex* if $S = \bigoplus S_j$ is a Clifford bundle whose Dirac operator D equals $d + d^*$.

Proposition 20.6. Let T and B be bounded operators on a Hilbert space H , and suppose that either T is trace-class, or both T and B are Hilbert-Schmidt. Then TB and BT are trace-class, and $\text{tr}(TB) = \text{tr}(BT)$.

Theorem 20.7. Let A be a smoothing operator on $L^2(S)$, S vector bundle over M , compact smooth manifold, with kernel k . Then A is of trace-class, and

$$(20.8) \quad \text{tr}(A) = \int \tilde{\text{tr}}(k(m, m)) \text{vol}(m),$$

where $\tilde{\text{tr}}: S_m \otimes S_m^* \rightarrow \mathbb{C}$ denotes the canonical trace on the endomorphism of the finite-dimensional vector space S_m .

Lefschetz numbers. Let M be a manifold and $\phi: M \rightarrow M$ a map. Then ϕ induces an endomorphism ϕ^* of the (complex) cohomology of M , and the Lefschetz number $L(\phi)$ of ϕ is defined by

$$(20.9) \quad L(\phi) = \sum_q (-1)^q \text{tr}(\phi^* \text{ on } H^q(M))$$

The classical Lefschetz formula expresses $L(\phi)$ as a sum over fixed points of ϕ . In particular, if $L(\phi) \neq 0$, then ϕ has a fixed point!

We want to approach the Lefschetz formula analytically and we will do so in the more general context of Dirac complexes. Thus, let M be a compact oriented n -dimensional Riemannian manifold, and let (S, d) be a Dirac complex over M . Let ϕ be a smooth map from M to M . Then ϕ induces $\phi^*: C^\infty(S) \rightarrow C^\infty(\phi^*S)$ in the following way: given a section $s \in C^\infty(S)$, we can define the map $s \circ \phi: M \rightarrow S$. Observe that via the identification of $\phi^*S = \{(m, v) \in M \times S \mid \phi(m) = \pi(v)\}$, we can consider $s \circ \phi$ as a section of ϕ^*S . In fact given $m \in M$, $s(\phi(m)) \in S_{\phi(m)} = (\phi^*S)_m$.

In case S is the de Rham complex, i.e.

$$(20.10) \quad S = \bigoplus_k \Lambda^k T^*M,$$

there is a natural bundle map $\zeta = \Lambda^* T^* \phi$ (over ϕ) from ϕ^*S to S which sends sections of $\phi^* \bigoplus_k \Lambda^k T^*M$ to sections of $\bigoplus_k \Lambda^k T^*M$: in fact given a section $s: M \rightarrow \phi^* \bigoplus_k \Lambda^k T^*M$, i.e. a map $M \rightarrow \bigoplus_k \Lambda^k T^*M$ such that $\pi \circ s = \phi$, by composing with $\Lambda^* T^* \phi$ we have obtained a map over the identity.

As an example let

$$(20.11) \quad s(m) = \omega \in \left(\phi^* \bigoplus_k \Lambda^k T^*M \right)_m = \left(\bigoplus_k \Lambda^k T^*M \right)_{\phi(m)},$$

where $\zeta(s(m)) = \omega \circ T\phi$, which is a covector living in the fiber over m , i.e. $\zeta(s)$ is a section of $\bigoplus_k \Lambda^k T^*M$.

For a general Dirac complex there is no such map, and we must assume the existence of a bundle $\zeta: \phi^*S \rightarrow S$ as part of our data. Thus there is a composite map

$$(20.12) \quad F = \zeta \phi^*: C^\infty(S) \rightarrow C^\infty(S)$$

Definition 20.13. If F (as above) is a map of complexes (i.e. $Fd = dF$), one says that (ζ, ϕ) is a *geometric endomorphism* of the given Dirac complex. Its Lefschetz number $L(\zeta, \phi)$ is defined by

$$(20.14) \quad L(\zeta, \phi) = \sum_q (-1)^q \text{tr}(F^* \text{ on } H^q(S)).$$

This definition is arranged so that a smooth map ϕ induces a natural geometric endomorphism of the de Rham complex, and its Lefschetz number defined by (20.14) agrees with the classical definition (20.9).

To apply analysis to the calculation of the Lefschetz number, we use the Hodge theorem: $H^q(S)$ is represented by the space \mathcal{H}^q of harmonic sections of S_q . So if we define P_q to be the orthogonal projection $L^2(S_q) \rightarrow \mathcal{H}^q$, then

$$(20.15) \quad \text{tr}(F^* \text{ on } H^q(S)) = \text{tr}(FP_q),$$

where F is considered as a continuous linear operator from $C^\infty(S_q)$ to $L^2(S_q)$.

Linearity is clear since $F = \zeta \circ \phi^*$, where ζ is a bundle map, hence fiberwise linear, and ϕ^* is precomposition with ϕ . Continuity is done as follows: since we established linearity is enough to prove that for $s_n \rightarrow 0$ in the C^∞ sense (actually C^0 sense is enough), then $F(s_n) \rightarrow 0$ in the L^2 sense. In fact it's not hard to see that $\zeta \circ s_n \circ \phi \rightarrow 0$ if $s_n \rightarrow 0$, since both ζ and ϕ are continuous and s_n so tending to zero uniformly.

Remark 20.16. Beware that the operator F itself may not be continuous (i.e. bounded) on L^2 . In fact if $s_n \rightarrow 0$ in the L^2 sense, it might well be that our sections are nonzero on a set of measure zero and if ϕ is constantly taking a value in such subset on a non-empty open set, then the resulting integral (the norm in L^2) won't be zero. Clearly

F is continuous (hence bounded) from $C^0(S_q)$ to $L^2(S_q)$, and therefore the composition of F with any smoothing operator is bounded on L^2 . This suffices for the arguments that follows. \triangleleft

Lemma 20.17. *The operators P_q are smoothing operators. Moreover, if Δ_q denotes the restriction of D^2 to $C^\infty(S_q)$, then as $t \rightarrow \infty$ the smoothing kernel of $e^{-t\Delta_q}$ tends to the smoothing kernel of P_q in the C^∞ topology.*

Proof. P_q can be written as $f(\Delta_q)$, where $f(0) = 1$ and $f(\lambda) = 0$ for all other λ . In fact harmonic sections in $C^\infty(S_q)$ correspond to the trivial eigenvalue for (the restriction of) D^2 . Unfortunately, this function is not smooth. However, since Δ_q has discrete spectrum, there is a smooth function f of compact support equal to 1 at zero (and around it for sufficiently small values between 0 and the next eigenvalue) and equal to 0 at all other eigenvalues of Δ_q . Then $f \in \mathcal{S}(\mathbb{R})$ and $f(\Delta_q) = P_q$. If $g_t(x) = (1 - f(x))e^{-tx}\mathbf{1}_{\mathbb{R}_+}(x)$ then we have

$$(20.18) \quad g_t(x) = (1 - f(x))e^{-tx}\mathbf{1}_{\mathbb{R}_+}(x)$$

$$(20.19) \quad = e^{-tx} - f(x)e^{-tx}\mathbf{1}_{\mathbb{R}_+}(x)$$

$$(20.20) \quad g_t(\Delta_q) = e^{-t\Delta_q} - f(\Delta_q)e^{-t\Delta_q}$$

$$(20.21) \quad = e^{-t\Delta_q} - P_q,$$

where (20.21) is justified as follows: if $s \in C^\infty(S_q)$, we can write it over a basis given by the eigenfunctions of Δ_q , $s = \sum_i s_i$ (with respective eigenvalues denoted by λ_i). In particular $s = s_0 + \sum_{i \geq 1} s_i$. Therefore we have

$$(20.22) \quad g_t(\Delta_q(s)) = e^{-t\Delta_q}(s) - f(\Delta_q)e^{-t\Delta_q}(s)$$

$$(20.23) \quad = \sum_i e^{-t\lambda_i} s_i - \sum_i f(\lambda_i) e^{-t\lambda_i} s_i$$

$$(20.24) \quad = \sum_i e^{-t\lambda_i} s_i - f(0) e^0 s_0$$

$$(20.25) \quad = \sum_i e^{-t\lambda_i} s_i - s_0$$

$$(20.26) \quad = e^{-t\Delta_q}(s) - f(\Delta_q)(s),$$

where in (20.24) we used the fact that f is zero for nonzero eigenvalues. Therefore we showed that for our purposes $g_t(x) = e^{-tx} - f(x)$. Now notice that $g_t(x) \rightarrow 0$ in $\mathcal{S}(\mathbb{R})$ as $t \rightarrow \infty$. Hence, by functional calculus, we have $g_t(\Delta_q) \rightarrow 0$, and the result follows. \square

Therefore $\text{tr}(FP_q) = \lim_{t \rightarrow \infty} \text{tr}(Fe^{-t\Delta_q})$, so

$$(20.27) \quad L(\zeta, \phi) = \lim_{t \rightarrow \infty} \sum_q (-1)^q \text{tr}(Fe^{-t\Delta_q}).$$

Let us analyze this expression more closely.

Proposition 20.28. *For all values of $t > 0$,*

$$\sum_q (-1)^q \text{tr}(Fe^{-t\Delta_q}) = L(\zeta, \phi).$$

Proof. It is enough to prove that $\sum_q (-1)^q \text{tr}(Fe^{-t\Delta_q})$ is constant in t . We differentiate, getting

$$(20.29) \quad \sum_q (-1)^{q+1} \text{tr}(F(dd^* + d^*d)e^{-t\Delta_q}).$$

Now $dF = Fd$, so

$$(20.30) \quad \text{tr}(Fdd^*e^{-t\Delta_q}) = \text{tr}(dFd^*e^{-t\Delta_q}).$$

Assume for the moment that we have

$$(20.31) \quad \text{tr}(dFd^*e^{-t\Delta_q}) = \text{tr}(Fd^*e^{-t\Delta_q}d)$$

Then

$$(20.32) \quad \operatorname{tr}(dFd^*e^{-t\Delta_q}) = \operatorname{tr}(Fd^*e^{-t\Delta_q}d)$$

$$(20.33) \quad = \operatorname{tr}(Fd^*de^{-t\Delta_{q-1}}),$$

where (20.33) is justified by the fact that $\Delta_q d = d\Delta_{q-1}$. The terms in the sum for the derivative therefore cancel in pairs, giving 0. (This is sometimes called *supersymmetry*.)

It remains to prove the identity (20.31), which we do as follows:

$$(20.34) \quad \operatorname{tr}(dFd^*e^{-t\Delta_q}) = \operatorname{tr}(dFd^*e^{-t\Delta_q/2}e^{-t\Delta_q/2})$$

$$(20.35) \quad = \operatorname{tr}(e^{-t\Delta_q/2}dFd^*e^{-t\Delta_q/2})$$

$$(20.36) \quad = \operatorname{tr}(Fd^*e^{-t\Delta_q/2}e^{-t\Delta_q/2}d)$$

$$(20.37) \quad = \operatorname{tr}(Fd^*e^{-t\Delta_q}d),$$

where (20.35) is justified by applying Proposition 20.6 to $dFd^*e^{-t\Delta_q/2}$ and $e^{-t\Delta_q/2}$. Clearly the latter is a bounded Hilbert-Schmidt operator; the former is bounded since $dFd^*e^{-t\Delta_q/2} = Fdd^*e^{-t\Delta_q/2}$ and since dd^* a differential operator, $dd^*e^{-t\Delta_q/2}$ is a smoothing operator (move dd^* under the sign of integral, which can be done because M is smooth). As observed in Remark 20.16, F followed by a smoothing operator is bounded. Now by Theorem 20.7 we have that $e^{-t\Delta_q/2}$ is of trace-class, hence we can apply Proposition 20.6.

Similarly, (20.36) is justified by the fact that $e^{-t\Delta_q/2}d = de^{-t\Delta_{q-1}/2}$, which shows that the latter is a smoothing operator, hence bounded and trace-class. $Fd^*e^{-t\Delta_q/2}$ is bounded for the same reason as above. Therefore we can apply Proposition 20.6. \square

Proposition 20.38. *If ϕ has no fixed points, then the Lefschetz number $L(\zeta, \phi)$ is 0.*

Proof. By Proposition 20.28 we have

$$(20.39) \quad L(\zeta, \phi) = \sum_q (-1)^q \operatorname{tr}(Fe^{-t\Delta_q})$$

for any $t > 0$. Look at the behavior of this expression for small t . If $k_t^q(m_1, m_2)$ denotes the heat kernel corresponding to $e^{-t\Delta_q}$, then $Fe^{-t\Delta_q}$ is a smoothing operator with kernel

$$(20.40) \quad (m_1, m_2) \rightarrow {}_1\zeta \cdot k_t^q(\phi(m_1), m_2),$$

where ${}_1\zeta$ denotes ζ acting on the first variable in the tensor product bundle $S \boxtimes S^*$ that k is a section of. Therefore by Theorem 20.7,

$$(20.41) \quad \operatorname{tr}(Fe^{-t\Delta_q}) = \int_M \operatorname{tr} {}_1\zeta \cdot k_t^q(\phi(m), m) \operatorname{vol}(m).$$

Now the assumption that ϕ has no fixed points means that its graph $\{(\phi(m), m) \mid m \in M\}$ never meets the diagonal in $M \times M$. Therefore, by the asymptotic expansion of the Heat kernel, or by the localization result, $k_t^q(\phi(m), m)$ tends to 0, uniformly in m , as $t \rightarrow 0$. The result follows. \square

Example 20.42. Any holomorphic automorphism of \mathbb{CP}^n has a fixed point: since \mathbb{CP}^n is Kähler, its Dolbeault complex is a Dirac complex. It's known that the Dolbeault cohomology of \mathbb{CP}^n is

$$(20.43) \quad H_{\bar{\partial}}^q(M) = \begin{cases} \mathbb{C}, & q = 0 \\ 0, & q > 0. \end{cases}$$

A holomorphic automorphism of M induces a geometric endomorphism of the Dolbeault complex (as in the de Rham case), which must act as the identity on $H^0(\mathbb{CP}^n)$, hence has strictly positive Lefschetz number. \blacktriangleleft