

AFTER CHIRAL ANOMALIES WE AT LAST TRY TO GRAND UNIFIED THEORIES OR GUT.

THE FIRST EXAMPLE IS BASED ON SO(5).

WHAT IS A GUT THEORY? A CHOICE OF GAUGE FIELDS + MATRICES REPRESENTATIONS

STANDARD MODEL : $SU(3) \otimes SU(2) \otimes U(1)$

TAKE QUARK DOUBLET $Q_L = \begin{pmatrix} u_L \\ d_L \end{pmatrix} \sim (3, 2, \frac{1}{6})_L$

TAKE ANTIQUARK $Q_R \sim (3, 1, 2/3)_R$

SO ONE FAMILY IS 15 FIELDS (IN COMPLEXES)

$(3, 2, \frac{1}{6})_L, (3, 1, 2/3)_R, (3, 1, -1/3)_R, (1, 2, -1/2)_L, (1, 1, -1)_R$

THESE ARE THE BUILDING BLOCKS OF THE SM (3 TIMES AS THERE ARE 3 FAMILIES)

NOW $G = SU(3) \otimes SU(2) \otimes U(1)$ WITH (HOPEFULLY) FEWER BLOCKS?

FIRST MAKE THEM ALL L-HADED (WE GET NO L SUBSCRIPT):

$(3, 2, \frac{1}{6}), (\bar{3}, 1, -\frac{2}{3}), (3, 1, \frac{1}{3}), (1, 2, -\frac{1}{2}), (1, 1, 1)$

$\bar{3}$ = COMPLEX CONJUGATE OF 3

SMALLEST SIMPLE GROUP THAT CONTAINS $SU(3) \otimes SU(2) \otimes U(1)$ IS SO(5).

THE HAVE THE SAME RANK = NUMBER OF DIAGONAL GENERATORS.

2 FOR $SU(2)$ + 1 FOR $SU(3)$ + 1 FOR $U(1)$ (TRIVIALLY)

AND 4 FOR SO(5)

CLEARLY, THE RANK OF $SO(N)$ GROUPS IS $N-1$

Q3: RANK? GENERATORS ARE HERMITIAN \Rightarrow CAN BE DIAGONALIZED WITH REAL EIGENVALUES.

NOT POSSIBLE FOR ALL OF THEM (OR THE GROUP GENERATORS WOULD COMMUTE). THE LARGEST SET HAS DIMENSION

c RAYIL

THE ARE ALSO TRACELESS $\Rightarrow \sum_i \lambda_i^{(A)} = 0$ FOR A GIVEN (A) DIAGONALIZED GENERATORS.

SU(N): T^A ARE $N \times N \Rightarrow N$ EIGENVALUES - 1

THUS THE SPACE OF COMMUTING GENERATORS HAS DIM $N-1$

NR: RANK OF $SO(2N)$ GROUPS IS N . SHOW THAT FOR $SO(8)$

$SO(8)$ HAS $4^2 = 16$ GENERATORS. IN THE FUNDAMENTAL REP,
5x5 MATRICES; 8 ARE OF THE FORM

$$\begin{pmatrix} 3 \times 3 & \begin{matrix} 0 & 0 \\ 0 & 0 \\ 0 & 0 \end{matrix} \\ \begin{matrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{matrix} & \vdots \\ 0 & 0 & 0 \end{pmatrix}$$

AND 8 OF THE FORM

$$\begin{pmatrix} \begin{matrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{matrix} & \begin{matrix} 0 & 0 \\ 0 & 0 \\ 0 & 0 \end{matrix} \\ \begin{matrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{matrix} & \boxed{2 \times 2} \end{pmatrix}$$

THEY COMMUTE WITH EACH OTHERS \Leftrightarrow SUBGROUPS $\sim SO(3) \otimes SO(4)$
FURTHERMORE, THEY COMMUTE WITH

$$\frac{y}{z} = \begin{pmatrix} -\frac{1}{3} & -\frac{1}{3} & 0 \\ -\frac{1}{3} & -\frac{1}{3} & 0 \\ 0 & 0 & \frac{1}{2} \end{pmatrix}$$

WE CALL THE HYPERCUBE.

THAT MAKES 4 DIAGONAL GENERATORS USING

$$X^3 \sim \frac{1}{i} \begin{pmatrix} 1 & 0 \\ 0 & -1 \\ 0 & 0 \end{pmatrix} \quad \text{AND} \quad X^8 \sim \frac{1}{2\sqrt{2}} \begin{pmatrix} 1 & 0 \\ 0 & 1 \\ 0 & -1 \end{pmatrix} \quad \text{FOR } SU(3)$$

$$\text{AND} \quad S^3 \sim \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad \text{FOR } SO(2)$$

Rmk: $T_n \left[(\gamma_i)^2 \right] = \frac{1}{e}$ ie. we use the normalization explained in the previous lecture.

LET US FIRST ACT ON A 5-DIMENSIONAL OBJECT ψ^a WITH $a=1\dots 5$. WE SPLIT IT INTO $a = \sum_i a_i$ WITH $i=1, 2, 3$ AND $a_i \sim \text{SO}(2)$

i.e.

$$\psi_a \left(\psi^a \right) = \begin{pmatrix} \psi^1 \\ \psi^2 \\ \psi^3 \\ \psi^4 \\ \psi^5 \end{pmatrix} = \begin{pmatrix} \psi^I \\ \dots \\ \psi^i \end{pmatrix} \quad \begin{cases} \text{3 numbers} \\ \text{(well, L-spinors)} \end{cases} \quad \begin{cases} \sim 2 \end{cases}$$

IT IS NATURAL TO TAKE $\psi^I \sim 3$, A TRIPLET OF $\text{SO}(3)$, AND $\psi^i \sim 2$, A DOUBLET OF $\text{SO}(2)$.
OR $\psi^I \sim 1$, ORDER $\text{SO}(2)$ TRANSFORMATIONS AND $\psi^i \sim 1$, ORDER $\text{SU}(3)$! HENCE

$$5 \rightarrow (3, 1, -\frac{1}{3}) \oplus (1, 2, \frac{1}{2})$$

$$\text{OR } \bar{5} \rightarrow (\bar{2}, 1, \frac{1}{3}) \oplus (1, 2, -\frac{1}{2}) \sim \text{OLC} \quad L$$

WE COULD FIT 2 OUT OF 5 REPRESENTATIONS OF THE SM!
ALSO, OUT OF 24 GENERATORS, 8+3+1 DON'T MIX
THE OL AND L COMPONENTS, BUT 12 DO!
THEY TRANSFORM LEPTONS INTO (ANTI)QUARKS!
(THEY ARE SOMETIMES CALLED LEPTO-QUARKS)

WE FITTED 5 FIELDS; WE ARE LEFT WITH 10.

$$\underbrace{(3, 2, \frac{1}{6})}_{6} \oplus \underbrace{(\bar{3}, 1, -\frac{2}{3})}_{3} \oplus \underbrace{(1, 1, 1)}_{1} = 10$$

The ψ^a are $\text{dim}=5$. The most one if dimension is a 10.
indeed, all antisymmetric matrix $\gamma^{ab} = -\gamma^{ba}$ has

$$\frac{N(N-1)}{2} = 10 \text{ components!}$$

THAT SOUNDS GOOD BUT THE QUANTUM NUMBER MUST COME
RIGHT.

WE KNOW NOW THAT $\psi^a \sim 5 \rightarrow (3, 1, -\frac{1}{3}) \oplus (1, 2, \frac{1}{2})$

so $\gamma^{ab} \sim \psi^a \otimes_A \psi^b$ (TRANSFORMS AS)
 \hookrightarrow antisymmetric tensor product.

$$\rightarrow 5 \otimes_A 5 \rightarrow ((3, 1, -\frac{1}{3}) \oplus (1, 2, \frac{1}{2})) \otimes_A ((3, 1, -\frac{1}{3}) \oplus (1, 2, \frac{1}{2}))$$

here $(3, 1, -\frac{1}{3}) \otimes (3, 1, -\frac{1}{3}) = (\bar{3}, 1, -\frac{2}{3}) \sim u^c$

(REMEMBER $u^i \bar{u}^j - \bar{u}^i u^j \sim \epsilon^{ijk} N_R \sim \bar{s}$)

This state is the ANTI-QUARK OP (singlet of $SU(2)$)!
LEFT-HANDED,
CHARGE $-2/3$)

ALSO $(1, 2, \frac{1}{2}) \otimes (1, 2, \frac{1}{2}) = (1, 1, 1) \sim e^c$

HOW ABOUT

$$\underbrace{(\bar{2}, 1, -\frac{1}{3})}_{\psi^I} \otimes_A (\bar{1}, 2, \frac{1}{2}) ?$$

\otimes_A $\psi^i \sim 6$ components

$$\sim (\bar{3}, 2, \frac{1}{6}) \sim Q \quad (\text{TRIPLET OF SU(3)} \\ \text{DOUBLET OF SU(2)})$$

HYPERCARGE = $\frac{1}{6}$

$$so \quad 10 \rightarrow (\bar{3}, 2, \frac{1}{6}) \oplus (\bar{3}, 1, -\frac{2}{3}) \oplus (1, 1, 1) \\ 6 \quad + \quad 4 \quad + \quad 1 \quad = 10$$

WE HAVE FITTED ALL THE SM REPS (1 FAMILY) INTO

$$\boxed{\bar{5} \oplus 10}$$

!!!

in components

$$\psi_a \equiv \bar{5} = \begin{pmatrix} \psi_I \\ \psi_i \end{pmatrix} = \begin{pmatrix} \bar{d}_L \\ \bar{d}_S \\ \bar{d}_B \\ \bar{e} \\ e \end{pmatrix}$$

$$\psi^{ab} \equiv 5 = \{ \psi^{IJ}, \psi^{Ij}, \psi^{ij} \}$$

WITH

$$\gamma^{ab} = \begin{pmatrix} 0 & u^c & -u^c & u & d \\ -u^c & 0 & u^c & u & d \\ u^c & -u^c & 0 & u & d \\ -u & -u & -u & 0 & e^c \\ -d & -d & -d & -e^c & 0 \end{pmatrix}$$

FURTHER INSIGHTS OFFERED BY GRAND UNIFICATION

* CHARGE QUANTIZATION

- in QED EM FIELD
- $\mathcal{L}_2 : i\bar{\psi}_1(\phi - iA - m_1)\psi_1 + i\bar{\psi}_2(\phi - i\epsilon A - m_2)\psi_2$
is no problem with ϵ ARBITRARY NUMBER.
- in SO(5), A_μ COUPLES TO A GENERATOR!
 \iff COUPLINGS ARE QUANTIZED!
i.e. eigenvalues of $\frac{\epsilon}{2} \stackrel{3}{=} \pm \frac{1}{2}$
IDM WITH λ^3 AND λ^8

- * WHY IS THE ELECTRIC CHARGE OF THE PROTON OPPOSITE TO THAT OF THE ELECTRON?
OPERATORS ARE TRACELESS!

$$\Rightarrow T_\mu [Q] = 3Q\bar{a} + Q_e = 0$$

(or \bar{b})

$$Q\bar{a} = -\frac{1}{3}Q_e$$

if $Q_e = -1$

$Q\bar{a} = +\frac{1}{3}$

The up quarks are in the 10 of $SU(3)$; their charge follow from group theory, so $Q_u = -2 Q_d$, etc.

$$\text{Then } Q_p = 2Q_u + Q_d = -3Q_d = -Q_e \quad \text{CQFD}$$

* Weinberg Angle

$$\text{in } SU(3) \times SU(2) \times U(1) \quad \tan \theta_W = g'/g$$

$$\text{From } g W_\mu^A T^A + g' \frac{Y}{2} B_\mu$$

NORMALIZATION OF $SU(3)$ GENERATORS TO

$$T_3 [T^a T^b] = \frac{1}{2} \delta^{ab}$$

$\Rightarrow g$ AND g' MUST BE RELATED

$$T_3 [T_3^2] = \left(\frac{1}{2}\right)^2 + \left(\frac{1}{2}\right)^2 = \frac{1}{2}$$

$$T_3 \left[\left(\frac{Y}{2}\right)^2\right] = 3 \times \left(-\frac{1}{3}\right)^2 + 2 \left(\frac{1}{3}\right)^2 = \frac{1}{3} + \frac{1}{2} = \frac{5}{6}$$

\Rightarrow THUS T_3 AND $\sqrt{\frac{3}{5}} \frac{Y}{2}$ ARE CORRECTLY NORMALIZED.

$$\Rightarrow g_s A_\mu^3 T^3 + g_s \sqrt{\frac{3}{5}} \frac{Y}{2} B_\mu$$

$$\Rightarrow g A_\mu^3 T^3 + g' \frac{Y}{2} B_\mu \Rightarrow$$

$$\boxed{\frac{g'}{g} = \sqrt{\frac{3}{5}}}$$

$$\left| \sin^2 \theta_W \right| = \frac{\frac{g'^2}{g^2 + g'^2}}{SU(3)} = \frac{\frac{3/5}{1+3/5}}{2} = \frac{3}{8} = 0.375 \quad \text{TO BE COMPARED}$$

To 0.323 OR
SO AT LOW ENERGIES

OF COURSE $\sin^2 \theta_W = 3/8$ IS THE SUPPOSED VALUE AT SOME HIGH SCALE. HOW HIGH SHOULD IT BE?
TO BE DISCUSSED SOON.

- FURTHER INSIGHTS FROM ANOMALY CANCELLATION?
TO BE A GOOD THEORY, WE NEED TO CHECK THE ABSENCE OF ANOMALIES!
- WHAT IS $A^{abc} = \text{Tr}[\{T^a\} T^b, T^c]$ FOR THE S AND LO?
(ALL L-HANDED
SO IT IS A CHIRAL THEORY)

\Rightarrow CHECK P&S?

AND ZEE'S TRICK?

$A^{abc}, A^{acb} = A^{cab}, A^{cba}$ etc... full symmetric.

$$\begin{aligned} \text{Tr}[\{T^a, \{T^b, T^c\}\}] &= \text{Tr}[T^a T^b T^c] + \text{Tr}[T^a T^c T^b] \\ &\quad - \text{Tr}[\{\{T^a, T^b\} T^c\}] \text{ etc...} \end{aligned}$$

A^{abc} IS OF IMPORTANCE OF $SU(3)$. COMES FROM

$$\{T^a, T^b\}_R = i \delta^{abc} \eta^c$$

THEY $\text{Tr}_{LR}[\{T^a_R, \{T^b_L, T^c_L\}\}] = A(r) \delta^{abc}$

JUST NEED THAT NUMBER!

GET IT FOR A SPECIFIC T.

A GOOD CHOICE IS

$$T_2 \frac{y}{z} = \begin{pmatrix} -\frac{1}{3} & -\frac{1}{3} & 0 \\ 0 & -\frac{1}{3} & -\frac{1}{3} \\ 0 & 0 & \frac{1}{2} \end{pmatrix}$$

$$\left| \begin{matrix} \left(\frac{y}{z} \right)^3 \end{matrix} \right|_{\frac{5}{3}} = 3 \times \left(+\frac{1}{3} \right)^3 + 2 \left(-\frac{1}{2} \right)^3 + \frac{1}{9} - \frac{1}{4}$$

$$\begin{aligned} \left| \begin{matrix} \left(\frac{y}{z} \right)^3 \end{matrix} \right|_{10} &= 3 \times 2 \times \left(\frac{1}{6} \right)^3 + 3 \times \left(-\frac{2}{3} \right)^3 + (+1)^3 \\ &= \frac{1}{36} - \frac{8}{9} + 1 = \frac{1+36-32}{36} \end{aligned}$$

$$= \frac{5}{36}$$

$$\boxed{\underbrace{\frac{5}{36} - \frac{5}{36}}_0 \approx 0}$$

very nice



Proton decay (in $SO(5)$)

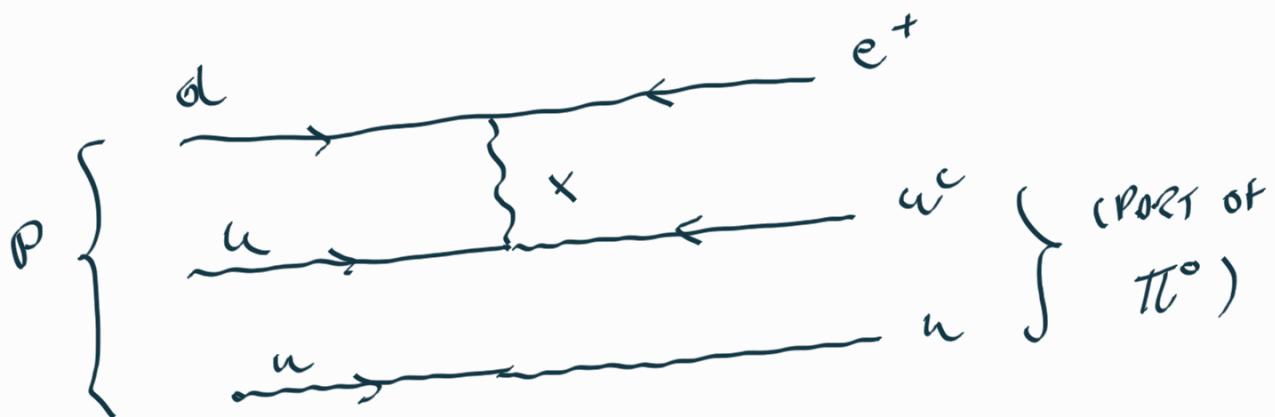
$$P \rightarrow \pi^0 e^+$$

if $\psi_a = \begin{pmatrix} d \\ d^c \\ d^c \\ u \\ e \end{pmatrix}$

THROUGH EXCHANGE
OF SV(S) GIVE BOSONS

in particular $sl \rightarrow e^+ + X (-4/3)$

A specific process is



This process is characterized by

$$\Delta B = -1 = B_f - B_i \quad \text{with } B_i = +1$$

(Proton is a baryon)

and $\Delta L = -1 = L_f - L_i$

with $L_f = -1$
(Postron is an anti-lepton)

so

$$\Delta(B+L) = -2$$

$$\Delta(B-L) = 0$$

