# Derivation of Newton Euler Algorithm

#### Advanced Robotic Systems – MANU2453

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**RMIT Classification: Trusted** 

#### Content

Derivation of the Newton-Euler Formulation



## Recap: Velocities

- Last week, we learnt about the notation of velocities:
  - Derivative of a position vector, with respect to reference frame {B}:

$${}^{B}V_{Q} = \frac{d}{dt} ({}^{B}Q)$$

 Linear velocity (differentiation done in {B}), but expressed in frame {A}:

$${}^{A}({}^{B}V_{Q}) = {}^{A}_{B} R \cdot {}^{B} V_{Q}$$

 Linear velocity of the origin of a frame {C}, relative to an understood universal frame {U}.

$$u_C = V_{CORG}$$

$$\upsilon_C = {}^{U} V_{CORG}$$
  ${}^{A} \upsilon_C = {}^{A} ({}^{U} V_{CORG})$ 

Rotation of frame {B} relative to {A}:

$$^{A}\Omega_{B}$$

 Angular velocity of frame {B} relative to {A}, but expressed in {C}:

$$C(A\Omega_B)$$

 Angular velocity of a frame {C}, relative to an understood universal frame {U}.

$$\omega_{c} = \Omega_{c}$$

$$\omega_C = {}^U \Omega_C$$
  ${}^A \omega_C = {}^A ({}^U \Omega_C)$ 



## **Acceleration of Rigid Body**

Now we can write the definition of acceleration:

$${}^{B}\dot{V}_{Q} = \frac{d}{dt} ({}^{B}V_{Q})$$

$${}^{A}\dot{\Omega}_{B} = \frac{d}{dt} \Big( {}^{A}\Omega_{B} \Big)$$

- That means, the accelerations are the derivatives of the velocities.
- Similar to the velocities, if the reference frame of differentiation is understood as the universal reference frame {U}, then we write:  $\dot{v}_C = \dot{V}_{CORG}$   $\dot{\phi}_C = \dot{V}_{CORG}$ 
  - e.g.  ${}^{A}\dot{\mathcal{U}}_{C}={}^{A}\left({}^{U}\dot{V}_{CORG}\right)$ 
    - Acceleration of origin of frame {C},
    - where differentiation is done in frame {U}.
    - The whole  ${}^{U}\dot{V}_{CORG}$  is then expressed in {A}.



#### **Linear Acceleration**

The velocity of a point can be obtained by differentiating its position

$$^{A}Q = ^{A}P_{BORG} + _{B}^{A}R \cdot ^{B}Q$$

which gives:

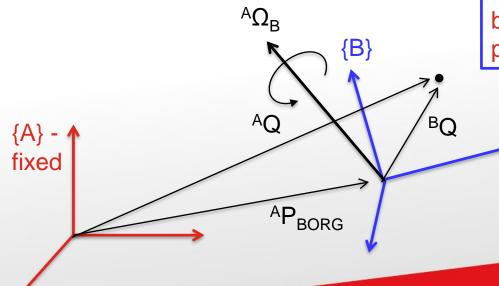
$${}^{A}V_{Q} = {}^{A}V_{BORG} + {}^{A}_{B}R \cdot {}^{B}V_{Q} + {}^{A}_{B}\dot{R} \cdot {}^{B}Q$$

• It can be shown that the last term is:  ${}^{A}\Omega_{R} \times_{R}^{A} R \cdot_{R}^{B} Q$ 

$${}^{A}\Omega_{B} \times_{B}^{A} R \cdot {}^{B} Q$$

Note: the "dot" is not a dot product. Just standard matrix-vector multiplication

Note: Multiplication before cross product





#### **Linear Acceleration**

- We now differentiate  ${}^{A}V_{Q} = {}^{A}V_{BORG} + {}^{A}_{B}R \cdot {}^{B}V_{Q} + {}^{A}\Omega_{B} \times {}^{A}_{B}R \cdot {}^{B}Q$  once more.
- This leads to (using chain rule for red and green underlined terms):

$${}^{A}\dot{V_{Q}} = {}^{A}\dot{V_{BORG}} + {}^{A}_{B}\dot{R} \cdot {}^{B}V_{Q} + {}^{A}_{B}R \cdot {}^{B}\dot{V_{Q}} + {}^{A}_{B}\dot{\Omega}_{B} \times {}^{A}_{B}R \cdot {}^{B}Q + {}^{A}\Omega_{B} \times {}^{A}_{B}\dot{R} \cdot {}^{B}Q + {}^{A}\Omega_{B}\dot{R} \cdot {}^{A}\Omega_{B}\dot{R} \cdot {}^{B}Q + {}^{A}\Omega_{B}\dot{R} \cdot {}^{A}\Omega$$

- The orange term is:  ${}^{A}_{B}\dot{R}\cdot{}^{B}Q={}^{A}\Omega_{B}\times_{B}^{A}R\cdot{}^{B}Q$  (from previous page)
- Analogously, the blue term would be:

$${}_{B}^{A}\dot{R}\cdot {}^{B}V_{Q}={}^{A}\Omega_{B}\times_{B}^{A}R\cdot {}^{B}V_{Q}$$

Therefore:

$${}^{A}\dot{V_{Q}} = {}^{A}\dot{V_{BORG}} + {}^{A}\Omega_{B} \times_{B}^{A} R \cdot {}^{B}V_{Q} + {}^{A}_{B}R \cdot {}^{B}\dot{V_{Q}} + {}^{A}\dot{\Omega}_{B} \times_{B}^{A}R \cdot {}^{B}Q + {}^{A}\Omega_{B} \times_{B}^{A}\Omega_{B} \times_{B}^{A}R \cdot {}^{B}Q + {}^{A}\Omega_{B} \times_{B}^{A}Q + {}^$$

Combining similar terms, we get:

Note: 2<sup>nd</sup> cross product before the 1<sup>st</sup>



$${}^{A}\dot{V_{Q}} = {}^{A}\dot{V_{BORG}} + 2{}^{A}\Omega_{B} \times_{B}^{A} R \cdot {}^{B}V_{Q} + {}^{A}_{B}R \cdot {}^{B}\dot{V_{Q}} + {}^{A}\dot{\Omega}_{B} \times_{B}^{A}R \cdot {}^{B}Q + {}^{A}\Omega_{B} \times {}^{A}\Omega_{B} \times_{B}^{A}R \cdot {}^{B}Q$$

• Note: For revolute joint:

$$^{B}V_{Q} = ^{B}\dot{V}_{Q} = 0$$



## **Angular Acceleration**

If {C} rotates relative to {B} with <sup>B</sup>Ω<sub>C</sub>, and {B} rotates relative to {A} with <sup>A</sup>Ω<sub>B</sub>, then {C} rotates relative to {A} with:

$${}^{A}\Omega_{C} = {}^{A}\Omega_{B} + {}^{A}_{B}R \cdot {}^{B}\Omega_{C}$$

Differentiating this gives: (using chain rule)

$${}^{A}\dot{\Omega}_{C} = {}^{A}\dot{\Omega}_{B} + {}^{A}_{B}\dot{R} \cdot {}^{B}\Omega_{C} + {}^{A}_{B}R \cdot {}^{B}\dot{\Omega}_{C}$$

 Again, the second term on the RHS can be written using cross product formulation.

$${}_{B}^{A}\dot{R}\cdot{}^{B}\Omega_{c}={}^{A}\Omega_{B}\times_{B}^{A}R\cdot{}^{B}\Omega_{C}$$

• Therefore:



$${}^{A}\dot{\Omega}_{C} = {}^{A}\dot{\Omega}_{B} + {}^{A}\Omega_{B} \times_{B}^{A} R \cdot_{B}^{B} \Omega_{C} +_{B}^{A} R \cdot_{B}^{B} \dot{\Omega}_{C}$$



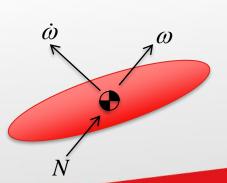
## **Newton's and Euler's Equations**

- Consider each link as a rigid body.
- Its mass distribution is characterized by the location of center of mass, and the inertia tensor of the link, both assumed known by now.
- Newtons' law and Euler's equation describe how the links move (accelerate) when given external forces.
  - Newtons' Law: F acting on the center of mass, causes a linear acceleration v-dot.

$$F = m\dot{v}_C$$

• Euler's equation: Moment N causes an angular acceleration of the rigid body:

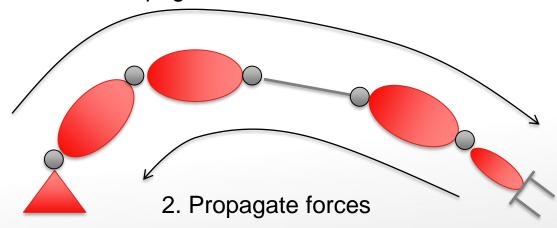
$$N = {^C} I\dot{\omega} + \omega \times {^C} I\omega$$





 Now, to calculate the joint torques (for all joints) which are required to move the manipulator along a given trajectory, we can use the iterative Newton-Euler Formulation:







- Outward iteration (1): Rotational velocities of links:
  - Start with  ${}^0\omega_0 = 0$  and  ${}^0\dot{\omega}_0 = 0$ .
  - Calculate velocity iteratively:

$$\stackrel{i+1}{\longrightarrow} \omega_{i+1} = \begin{pmatrix} i+1 & R \cdot i & \omega_i \end{pmatrix} + \begin{pmatrix} \dot{\theta}_{i+1} \cdot i+1 & \hat{Z}_{i+1} \end{pmatrix}$$

Note:  $\dot{\theta}_{i+1}$  is a scalar. It is multiplied with  $\dot{Z}_{i+1}$  to become vector

- Outward iteration (2): Rotational acceleration of links:
  - Just now, we derived the following equation for angular acceleration:

$${}^{A}\dot{\Omega}_{C} = {}^{A}\dot{\Omega}_{B} + {}^{A}\Omega_{B} \times_{B}^{A} R \cdot_{B}^{B} \Omega_{C} + {}^{A}_{B} R \cdot_{B}^{B} \dot{\Omega}_{C}$$

• Set A = 0, B = i and C = i+1, we have:

$$\begin{split} {}^{0}\dot{\Omega}_{i+1} &= \left({}^{0}\dot{\Omega}_{i}\right) + \left({}^{0}\Omega_{i} \times_{i}^{0} R \cdot_{i}^{i} \Omega_{i+1}\right) + \left({}^{0}_{i} R \cdot_{i}^{i} \dot{\Omega}_{i+1}\right) \\ &= \left({}^{0}\dot{\Omega}_{i}\right) + \left({}^{0}\Omega_{i} \times_{i}^{0} R \cdot_{i+1}^{i} R \cdot_{i+1}^{i+1} \Omega_{i+1}\right) + \left({}^{0}_{i} R \cdot_{i+1}^{i} R \cdot_{i+1}^{i+1} \dot{\Omega}_{i+1}\right) \end{split}$$



• Last week, we had the following (see section Velocity Propagation):

$$^{i+1}\Omega_{i+1} = \dot{\theta}_{i+1} \cdot ^{i+1} \hat{Z}_{i+1}$$

Putting this into the equation from last page gives:

$$\begin{split} {}^{0}\dot{\Omega}_{i+1} &= \left({}^{0}\dot{\Omega}_{i}\right) + \left({}^{0}\Omega_{i} \times_{i}^{0} R \cdot_{i+1}^{i} R \cdot \dot{\theta}_{i+1} \cdot_{i+1}^{i+1} \hat{Z}_{i+1}\right) + \left({}^{0}_{i} R \cdot_{i+1}^{i} R \cdot \ddot{\theta}_{i+1} \cdot_{i+1}^{i+1} \hat{Z}_{i+1}\right) \\ &= \left({}^{0}\dot{\Omega}_{i}\right) + \left({}^{0}\Omega_{i} \times_{i+1}^{0} R \cdot \dot{\theta}_{i+1} \cdot_{i+1}^{i+1} \hat{Z}_{i+1}\right) + \left({}^{0}_{i+1} R \cdot \ddot{\theta}_{i+1} \cdot_{i+1}^{i+1} \hat{Z}_{i+1}\right) \end{split}$$

• Multiply both sides with  ${0 \atop 0}^{i+1}R$ , and using the identity  $R(a \times b) = Ra \times Rb$  where R is a rotation matrix gives:

$$\begin{split} &\stackrel{i+1}{\circ} R \cdot {}^{0} \ \dot{\Omega}_{i+1} = \begin{pmatrix} i+1 \\ 0 \end{pmatrix} R \cdot {}^{0} \ \dot{\Omega}_{i} \end{pmatrix} + \begin{pmatrix} i+1 \\ 0 \end{pmatrix} R \cdot {}^{0} \ \Omega_{i} \times_{0}^{i+1} R \cdot {}^{0} R \cdot \dot{\theta}_{i+1} \cdot {}^{i+1} \ \dot{Z}_{i+1} \end{pmatrix} + \begin{pmatrix} i+1 \\ 0 \end{pmatrix} R \cdot \ddot{\theta}_{i+1} \cdot {}^{i+1} \ \dot{Z}_{i+1} \end{pmatrix} \\ &= \begin{pmatrix} i+1 \\ i \end{pmatrix} R \cdot {}^{i} R \cdot {}^{0} \ \dot{\Omega}_{i} \end{pmatrix} + \begin{pmatrix} i+1 \\ i \end{pmatrix} R \cdot {}^{i} R \cdot {}^{0} R \cdot {}^{0} \Omega_{i} \times_{\underbrace{i+1}}^{i+1} R \cdot \dot{\theta}_{i+1} \cdot {}^{i+1} \ \dot{Z}_{i+1} \end{pmatrix} + \begin{pmatrix} i+1 \\ \underbrace{i+1} R \cdot \ddot{\theta}_{i+1} \cdot {}^{i+1} \ \dot{Z}_{i+1} \end{pmatrix} \\ &= \begin{pmatrix} i+1 \\ i \end{pmatrix} R \cdot {}^{i} R \cdot {}^{0} \Omega_{i} \end{pmatrix} + \begin{pmatrix} i+1 \\ i \end{pmatrix} R \cdot {}^{i} R \cdot {}^{0} \Omega_{i} \times \dot{\theta}_{i+1} \cdot {}^{i+1} \ \dot{Z}_{i+1} \end{pmatrix} + \begin{pmatrix} \ddot{\theta}_{i+1} \cdot {}^{i+1} \ \dot{Z}_{i+1} \end{pmatrix} \end{split}$$



• Using the notation  $\omega_C = \Omega_C$  and  $\omega_C = \Omega_C$  and  $\omega_C = \Omega_C$ , which in this case are:

$$\omega_i = {}^0 \Omega_i$$
 $\omega_i = {}^k \Omega_i = {}^k R \cdot {}^0 \Omega_i$ 

• gives:

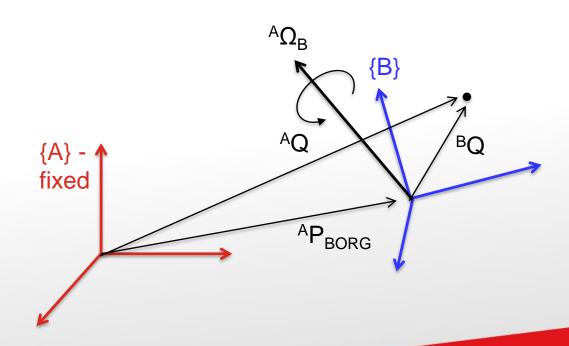
Note: This is allowed because of the property  $R(a \times b) = Ra \times Rb$ 

$$\stackrel{i+1}{\dot{\omega}_{i+1}} = \stackrel{(i+1}{i}R \cdot \stackrel{i}{\dot{\omega}_{i}}) + \stackrel{(i+1}{i}R \cdot \stackrel{i}{\omega}_{i} \times \dot{\theta}_{i+1} \cdot \stackrel{i+1}{\dot{\omega}_{i+1}} \hat{Z}_{i+1} + \stackrel{(i+1)}{\dot{\omega}_{i+1}} \hat{Z}_{i+1}$$



- Outward iteration (3): Linear acceleration of link-frame origin
  - Just now, we derived the following equation for angular acceleration:

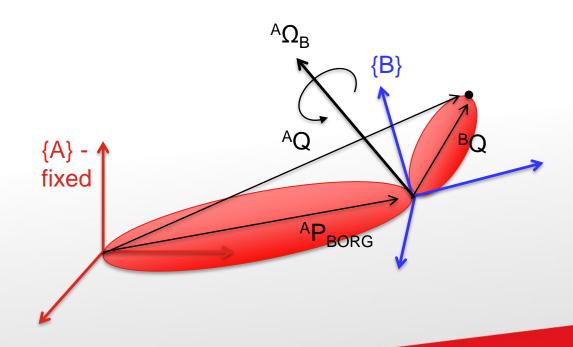
$${}^{A}\dot{V_{Q}} = \left({}^{A}\dot{V_{BORG}}\right) + \left(2{}^{A}\Omega_{B} \times_{B}^{A} R \cdot {}^{B}V_{Q}\right) + \left({}^{A}R \cdot {}^{B}\dot{V_{Q}}\right) + \left({}^{A}\dot{\Omega}_{B} \times_{B}^{A} R \cdot {}^{B}Q\right) + \left({}^{A}\Omega_{B} \times {}^{A}\Omega_{B} \times_{B}^{A} R \cdot {}^{B}Q\right)$$





• Set A = 0, B = i and Q = i+1, we have:

$${}^{0}\dot{V}_{i+1} = \left({}^{0}\dot{V}\right)_{i} + \left(2{}^{0}\Omega_{i} \times_{i}^{0} R \cdot {}^{i} V_{i+1}\right) + \left({}^{0}R \cdot {}^{i} \dot{V}_{i+1}\right) + \left({}^{0}\dot{\Omega}_{i} \times_{i}^{0} R \cdot {}^{i} Q\right) + \left({}^{0}\Omega_{i} \times {}^{0} \Omega_{i} \times_{i}^{0} R \cdot {}^{i} Q\right)$$





• Multiply both sides with  ${0 \atop 0}^{i+1}R$ , and using the identity  $R(a \times b) = Ra \times Rb$  where R is a rotation matrix gives:

$$\begin{split} & \stackrel{i+1}{_{0}}R \cdot {}^{0} \dot{V}_{i+1} = \begin{pmatrix} i^{+1}R \cdot {}^{0} \dot{V}_{i} \end{pmatrix} + \begin{pmatrix} 2^{i+1}R \cdot {}^{0} \Omega_{i} \times {}^{i+1}R \cdot {}^{0}_{i} R \cdot {}^{i} V_{i+1} \end{pmatrix} + \begin{pmatrix} i^{+1}R \cdot {}^{0}_{i} R \cdot {}^{i} \dot{V}_{i+1} \end{pmatrix} \\ & \quad + \begin{pmatrix} i^{+1}R \cdot {}^{0} \dot{\Omega}_{i} \times {}^{i+1}R \cdot {}^{0}_{i} R \cdot {}^{i} Q \end{pmatrix} + \begin{pmatrix} i^{+1}R \cdot {}^{0} \Omega_{i} \times {}^{i+1}R \cdot {}^{0} \Omega_{i} \times {}^{i+1}R \cdot {}^{0}_{i} R \cdot {}^{i} Q \end{pmatrix} \\ & \quad = \begin{pmatrix} i^{+1}R \cdot {}^{0} \dot{V}_{i} \end{pmatrix} + \begin{pmatrix} 2^{i+1}R \cdot {}^{i}_{0} R \cdot {}^{0} \Omega_{i} \times {}^{i+1}R \cdot {}^{0}_{0} R \cdot {}^{i} V_{i+1} \end{pmatrix} + \begin{pmatrix} i^{+1}R \cdot {}^{0}_{i} R \cdot {}^{i} \dot{V}_{i+1} \end{pmatrix} \\ & \quad + \begin{pmatrix} i^{+1}R \cdot {}^{i}_{0} R \cdot {}^{0} \dot{\Omega}_{i} \times {}^{i+1}R \cdot {}^{0}_{i} R \cdot {}^{i} Q \end{pmatrix} + \begin{pmatrix} i^{+1}R \cdot {}^{0} \Omega_{i} \times {}^{i+1}R \cdot {}^{0} \Omega_{i} \times {}^{i+1}R \cdot {}^{0}_{i} R \cdot {}^{i} Q \end{pmatrix} \\ & \quad = \begin{pmatrix} i^{+1}R \cdot {}^{i}_{0} R \cdot {}^{0} \dot{V} \end{pmatrix}_{i} + \begin{pmatrix} 2^{i+1}R \cdot {}^{i}_{0} R \cdot {}^{0} \Omega_{i} \times {}^{i+1} V_{i+1} \end{pmatrix} + \begin{pmatrix} i^{+1}\dot{V}_{i+1} \end{pmatrix} \\ & \quad + \begin{pmatrix} i^{+1}R \cdot {}^{i}_{0} R \cdot {}^{0} \dot{\Omega}_{i} \times {}^{i+1}R \cdot {}^{0}_{i} R \cdot {}^{i} Q \end{pmatrix} + \begin{pmatrix} i^{+1}R \cdot {}^{i}_{0} R \cdot {}^{0} \Omega_{i} \times {}^{i+1}R \cdot {}^{0}_{i} R \cdot {}^{0} R \cdot {}^{0} \Omega_{i} \times {}^{i+1}R \cdot {}^{0}_{i} R \cdot {}^{0} \Omega_{i} \times {}^{0} \Omega_{i} \times {}^{i+1}R \cdot {}^{0}_{i} R \cdot {}^{0} \Omega_{i} \times {}^{0} \Omega$$



• Using the notation  $\omega_C = {}^U \Omega_C$ ,  ${}^A \omega_C = {}^A ({}^U \Omega_C)$ ,  ${}^D \omega_C = {}^U V_{CORG}$  and  ${}^A \nu_C = {}^A ({}^U V_{CORG})$  which in this case are:

$$\omega_j = {}^{0} \Omega_j \qquad {}^{k} \omega_j = {}^{k} \Omega_j = {}^{k} R \cdot {}^{0} \Omega_j$$

$$\upsilon_{j} = {}^{0} V_{j} \qquad {}^{k} \upsilon_{j} = {}^{k} \left( {}^{0} V_{j} \right) = {}^{k}_{0} R \cdot {}^{0} V_{j}$$

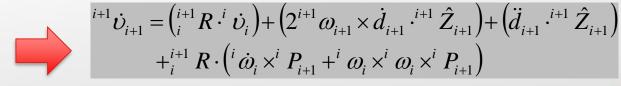
• gives:

$$\underbrace{\begin{bmatrix} i+1 \\ 0 \end{bmatrix} R \cdot \underbrace{0}_{i+1} \dot{V}_{i+1}}_{i} = \begin{bmatrix} i+1 \\ i \end{bmatrix} R \cdot \underbrace{0}_{i} R \cdot \underbrace{0}_{i} \dot{V}_{i} \\ \vdots \dot{v}_{i} \end{bmatrix}}_{i \dot{v}_{i}} + \underbrace{\begin{bmatrix} 2i+1 \\ 2i \end{bmatrix} R \cdot \underbrace{0}_{i} R \cdot \underbrace{0}_{i} \dot{Q}_{i}}_{i \dot{w}_{i}} \times \underbrace{\begin{bmatrix} i+1 \\ i+1 \end{bmatrix} \dot{V}_{i+1}}_{i \dot{w}_{i+1}} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix} \dot{V}_{i+1} \dot{Z}_{i+1}}_{i \dot{w}_{i}} + \underbrace{\begin{bmatrix} i+1 \\ i \end{bmatrix}$$



Finally, we obtain:

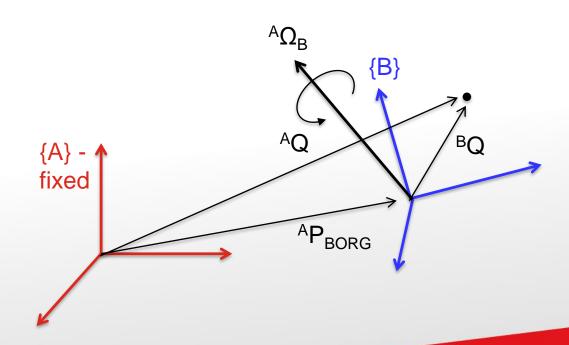
- The second bracket and the third brackets exist only if joint i+1 is prismatic.
- Since  $i^{i+1}\omega_{i+1} = \binom{i+1}{i}R \cdot i^i \omega_i + \underbrace{(\dot{\theta}_{i+1} \cdot i^{i+1} \hat{Z}_{i+1})}_{0 \text{ if prismatic}} = \binom{i+1}{i}R \cdot i^i \omega_i$
- the linear acceleration equation can be written as:





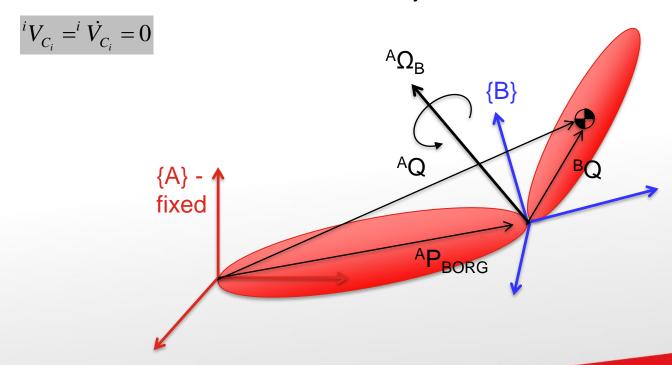
- Outward iteration (4): Linear acceleration of link's center of mass
  - Just now, we derived the following equation for angular acceleration:

$${}^{A}\dot{V_{Q}} = \left({}^{A}\dot{V_{BORG}}\right) + \left(2{}^{A}\Omega_{B} \times_{B}^{A} R \cdot {}^{B}V_{Q}\right) + \left({}^{A}_{B}R \cdot {}^{B}\dot{V_{Q}}\right) + \left({}^{A}\dot{\Omega}_{B} \times_{B}^{A} R \cdot {}^{B}Q\right) + \left({}^{A}\Omega_{B} \times {}^{A}\Omega_{B} \times_{B}^{A} R \cdot {}^{B}Q\right)$$





- Set A = 0, B = i and Q =  $C_i$ , we have:  ${}^0\dot{V}_{C_i} = \left({}^0\dot{V}_i\right) + \left(2{}^0\Omega_i \times_i^0 R \cdot {}^i V_{C_i}\right) + \left({}^0R \cdot {}^i \dot{V}_{C_i}\right) + \left({}^0\dot{\Omega}_i \times_i^0 R \cdot {}^i Q\right) + \left({}^0\Omega_i \times^0 \Omega_i \times_i^0 R \cdot {}^i Q\right)$
- Now, imagine the frame of C<sub>i</sub> has the same orientation as {B}. Also, the center of mass is constant from the ith joint. Therefore:





The equation thus simplifies to:

$${}^{0}\dot{V}_{C_{i}} = \left({}^{0}\dot{V}_{i}\right) + \left({}^{0}\dot{\Omega}_{i} \times_{i}^{0} R \cdot_{i}^{i} Q\right) + \left({}^{0}\Omega_{i} \times_{i}^{0} \Omega_{i} \times_{i}^{0} R \cdot_{i}^{i} Q\right)$$

• Multiply both sides with  ${}^{i}_{0}R$ , and using the identity  $R(a \times b) = Ra \times Rb$  where R is a rotation matrix gives:

$$\begin{split} & \stackrel{i}{_{0}}R \cdot {}^{0} \dot{V}_{C_{i}} = \left( \stackrel{i}{_{0}}R \cdot {}^{0} \dot{V}_{i} \right) + \left( \stackrel{i}{_{0}}R \cdot {}^{0} \dot{\Omega}_{i} \times \underbrace{\stackrel{i}{_{0}}R \cdot {}^{0}R}_{I} \cdot {}^{0}R \cdot {}^{0}Q \right) + \left( \stackrel{i}{_{0}}R \cdot {}^{0}\Omega_{i} \times \stackrel{i}{_{0}}R \cdot {}^{0}\Omega_{i} \times \underbrace{\stackrel{i}{_{0}}R \cdot {}^{0}R}_{I} \cdot {}^{0}Q \right) \\ & = \left( \stackrel{i}{_{0}}R \cdot {}^{0} \dot{V}_{i} \right) + \left( \stackrel{i}{_{0}}R \cdot {}^{0} \dot{\Omega}_{i} \times {}^{i}Q \right) + \left( \stackrel{i}{_{0}}R \cdot {}^{0}\Omega_{i} \times {}^{i}R \cdot {}^{0}\Omega_{i} \times {}^{i}Q \right) \end{split}$$



• Using the notation  $\omega_C = {}^U \Omega_C$ ,  ${}^A \omega_C = {}^A ({}^U \Omega_C)$ ,  ${}^D \omega_C = {}^U V_{CORG}$  and  $^{A}v_{C} = ^{A}(^{U}V_{CORG})$  which in this case are:

$$\omega_j = {}^0 \Omega_j \quad {}^k \omega_j = {}^k \Omega_j = {}^k R \cdot {}^0 \Omega_j \quad \upsilon_j = {}^0 V_j \quad {}^k \upsilon_j = {}^k ({}^0 V_j) = {}^k R \cdot {}^0 V_j$$

$${}^{i}\dot{\mathcal{U}}_{C_{i}} = ({}^{i}\dot{\mathcal{U}}_{i}) + ({}^{i}\dot{\omega}_{i} \times {}^{i}P_{C_{i}}) + ({}^{i}\omega_{i} \times {}^{i}\omega_{i} \times {}^{i}P_{C_{i}})$$

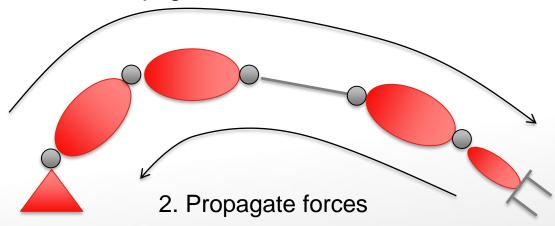
- Outward iteration (5): Force acting on a link
  - Using Newton's law, we have the force acting at the center of mass of each link:  $F_i = m\dot{\upsilon}_{C_i}$
- Outward iteration (6):Torque acting on a link

$$N_i = {^{C_i}} I\dot{\omega}_i + \omega_i \times {^{C_i}} I\omega_i$$



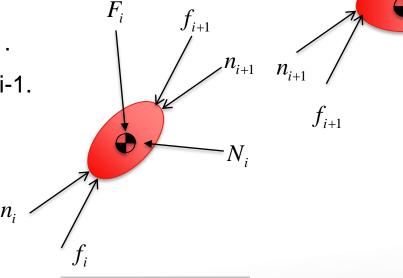
- So far, we have obtained the equations for outward iterations.
- Now, it's time to do the inward iterations.







- Define:
  - fi = force exerted on link i by link i-1.
  - ni = torque exerted on link i by link i-1.



- Summing the forces acting on link i gives:  ${}^{i}F_{i} = {}^{i}f_{i} {}^{i}_{i+1}R^{i+1}f_{i+1}$
- Summing the torques about center of mass gives:

$${}^{i}N_{i} = {}^{i}n_{i} - {}^{i}n_{i+1} + (-{}^{i}P_{C_{i}}) \times {}^{i}f_{i} - ({}^{i}P_{i+1} - {}^{i}P_{C_{i}}) \times {}^{i}f_{i+1}$$



• Substituting  ${}^{i}F_{i} = {}^{i}f_{i} - {}^{i}_{i+1}R^{i+1}f_{i+1}$  into the torque equation gives:

$$\begin{split} {}^{i}N_{i} &= {}^{i}n_{i} - {}^{i}n_{i+1} + \left( -{}^{i}P_{C_{i}} \right) \times {}^{i}f_{i} - \left( {}^{i}P_{i+1} - {}^{i}P_{C_{i}} \right) \times {}^{i}f_{i+1} \\ &= {}^{i}n_{i} - {}^{i}n_{i+1} + \left( -{}^{i}P_{C_{i}} \right) \times \left( {}^{i}F_{i} + {}^{i}_{i+1}R \cdot {}^{i+1}f_{i+1} \right) - \left( {}^{i}P_{i+1} - {}^{i}P_{C_{i}} \right) \times {}^{i}f_{i+1} \\ &= {}^{i}n_{i} - {}^{i}_{i+1}R \cdot {}^{i+1}n_{i+1} + \left( -{}^{i}P_{C_{i}} \right) \times \left( {}^{i}F_{i} + {}^{i}_{i+1}R \cdot {}^{i+1}f_{i+1} \right) - \left( {}^{i}P_{i+1} - {}^{i}P_{C_{i}} \right) \times {}^{i}_{i+1}R \cdot {}^{i+1}f_{i+1} \end{split}$$

• Because cross product is distributive over addition, i.e.  $a \times (b+c) = a \times b + a \times c$  we have:

$$\begin{array}{l}
^{i}N_{i} = ^{i}n_{i} - ^{i}_{i+1}R \cdot ^{i+1}n_{i+1} + \left( - ^{i}P_{C_{i}}\right) \times \left( ^{i}F_{i} + ^{i}_{i+1}R \cdot ^{i+1}f_{i+1}\right) - \left( ^{i}P_{i+1} - ^{i}P_{C_{i}}\right) \times ^{i}_{i+1}R \cdot ^{i+1}f_{i+1} \\
= ^{i}n_{i} - ^{i}_{i+1}R \cdot ^{i+1}n_{i+1} - ^{i}P_{C_{i}} \times ^{i}F_{i} - ^{i}P_{C_{i}} \times ^{i}_{i+1}R \cdot ^{i+1}f_{i+1} - ^{i}P_{i+1} \times ^{i}_{i+1}R \cdot ^{i+1}f_{i+1} + ^{i}P_{C_{i}} \times ^{i}_{i+1}R \cdot ^{i+1}f_{i+1} \\
= ^{i}n_{i} - ^{i}_{i+1}R \cdot ^{i+1}n_{i+1} - ^{i}P_{C_{i}} \times ^{i}F_{i} - ^{i}P_{i+1} \times ^{i}_{i+1}R \cdot ^{i+1}f_{i+1} \\
= ^{i}n_{i} - ^{i}_{i+1}R \cdot ^{i+1}n_{i+1} - ^{i}P_{C_{i}} \times ^{i}F_{i} - ^{i}P_{i+1} \times ^{i}_{i+1}R \cdot ^{i+1}f_{i+1} \\
= ^{i}n_{i} - ^{i}_{i+1}R \cdot ^{i+1}n_{i+1} - ^{i}P_{C_{i}} \times ^{i}F_{i} - ^{i}P_{i+1} \times ^{i}_{i+1}R \cdot ^{i+1}f_{i+1} \\
= ^{i}n_{i} - ^{i}_{i+1}R \cdot ^{i+1}n_{i+1} - ^{i}P_{C_{i}} \times ^{i}F_{i} - ^{i}P_{i+1} \times ^{i}_{i+1}R \cdot ^{i+1}f_{i+1} \\
= ^{i}n_{i} - ^{i}_{i+1}R \cdot ^{i+1}n_{i+1} - ^{i}P_{C_{i}} \times ^{i}F_{i} - ^{i}P_{i+1} \times ^{i}_{i+1}R \cdot ^{i+1}f_{i+1} \\
= ^{i}n_{i} - ^{i}_{i+1}R \cdot ^{i+1}n_{i+1} - ^{i}P_{C_{i}} \times ^{i}F_{i} - ^{i}P_{i+1} \times ^{i}F_{i+1}R \cdot ^{i+1}f_{i+1} \\
= ^{i}n_{i} - ^{i}_{i+1}R \cdot ^{i+1}n_{i+1} - ^{i}P_{C_{i}} \times ^{i}F_{i} - ^{i}P_{i+1} \times ^{i}F_{i+1}R \cdot ^{i+1}f_{i+1} \\
= ^{i}n_{i} - ^{i}_{i+1}R \cdot ^{i+1}n_{i+1} - ^{i}P_{C_{i}} \times ^{i}F_{i} - ^{i}P_{i+1} \times ^{i}F_{i+1}R \cdot ^{i+1}f_{i+1} \\
= ^{i}n_{i} - ^{i}_{i+1}R \cdot ^{i+1}n_{i+1} - ^{i}P_{C_{i}} \times ^{i}F_{i} - ^{i}P_{i+1} \times ^{i}F_{i+1}R \cdot ^{i+1}f_{i+1} \\
= ^{i}n_{i} - ^{i}_{i+1}R \cdot ^{i}F_{i} - ^{i}P_{i} - ^$$



Finally, we put those equations into forms suitable for iterations:

$${}^{i} f_{i} = {}^{i}_{i+1} R^{i+1} f_{i+1} + {}^{i} F_{i}$$

$${}^{i} n_{i} = {}^{i}_{i+1} R \cdot {}^{i+1} n_{i+1} + {}^{i} P_{C_{i}} \times {}^{i} F_{i} + {}^{i} P_{i+1} \times {}^{i}_{i+1} R \cdot {}^{i+1} f_{i+1} + {}^{i} N_{i}$$

- Finally, the torques at each joint are calculated as:



$$\tau_i = {}^i n_i^T \cdot {}^i \hat{Z}$$

• Prismatic joint:



Note that these inward iterations started with  $\int_{N+1}^{N+1} f_{N+1} = 0$  &  $\int_{N+1}^{N+1} f_{N+1} = 0$  if the robot is moving in free space, or the actual contact force / torque if the robot is in contact with some objects.



## Thank you!

Have a good evening.

