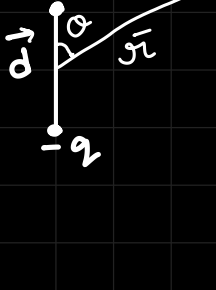


* LECTURE 9



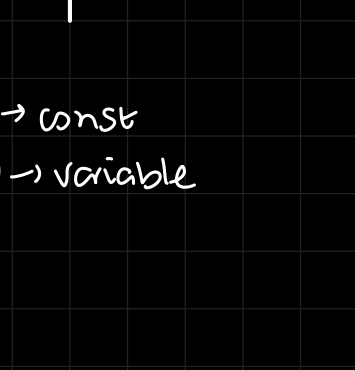
$$\vec{p} = q\vec{d}$$

$$V(r, \theta) = \frac{p \cos \theta}{4\pi\epsilon_0 r^2} = \left(\frac{p}{4\pi\epsilon_0} \cdot \frac{\hat{r}}{r^2} \right)$$

net dipole moment due to induced and inherent dipole moment in the direction of electric field.

Instead of dipole moment for individual particles, we express dipole moment as polarization dipole moment per volume for a surface.

\vec{P} = polarization dipole moment per unit volume



$$\frac{\partial f}{\partial z} = \frac{\partial f}{\partial (x-x')} \cdot \frac{\partial (x-x')}{\partial x}$$

$x \rightarrow \text{const}$
 $x' \rightarrow \text{variable}$

$$= \frac{\partial f}{\partial (x-x')}$$

$$\frac{\partial f}{\partial x'} = \frac{\partial f}{\partial (x-x')} \cdot \frac{\partial (x-x')}{\partial x'}$$

$$= (-1) \frac{\partial f}{\partial (x-x')}$$

$$\text{So, } \vec{\nabla}' f = -\vec{\nabla} f$$

$$\vec{\nabla} \left(\frac{1}{r} \right) = -\vec{\nabla}' \left(\frac{1}{r} \right) = \frac{\hat{r}}{r^2}$$

* remember to use the vector notation (\rightarrow) for quiz

$$\vec{p} = \vec{P} dv'$$

$$dv = \frac{1}{4\pi\epsilon_0} \left(\frac{\vec{P} \cdot \hat{r}}{r^2} \right) dv'$$

$$V = \frac{1}{4\pi\epsilon_0} \int_V \frac{\vec{P} \cdot \hat{r}}{r^2} dv'$$

$$V = \frac{1}{4\pi\epsilon_0} \int_V \vec{P} \cdot \vec{\nabla} \left(\frac{1}{r} \right) dv'$$

remember last lecture \rightarrow

$$\int_V \vec{A} (\vec{\nabla} f) dv = \underbrace{\oint_S \vec{A} d\vec{s}}_{\text{surface integral}} - \int_V f (\vec{\nabla} \cdot \vec{A}) dv$$

here $\vec{P} = \vec{A}$ and $f = \frac{1}{r}$ } magnitude

$$V = \frac{1}{4\pi\epsilon_0} \left[\oint_S \frac{\vec{P} \cdot d\vec{s}}{r} - \int_V \frac{1}{r} (\vec{\nabla} \cdot \vec{P}(r')) dv' \right]$$

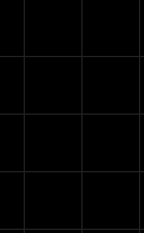
$$V = \frac{1}{4\pi\epsilon_0} \left[\oint_S \frac{\vec{P} \cdot d\vec{s}}{r} - \int_V \frac{1}{r} (\vec{\nabla} \cdot \vec{P}(r')) dv' \right]$$

Potential due to a surface charge density:



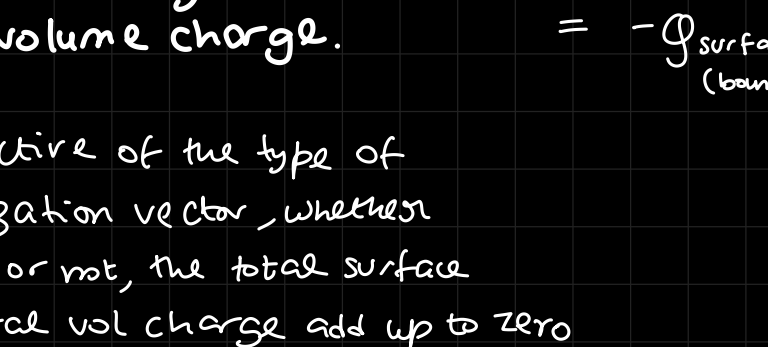
$$= \frac{1}{4\pi\epsilon_0} \int \frac{\sigma_s ds'}{r}$$

Potential due to a volume charge density



$$= \frac{1}{4\pi\epsilon_0} \int \frac{\rho dv'}{r}$$

for a material kept in an electric field \rightarrow



We are assuming a uniform polarization vector

So, potential of the surface charge is $\sigma_B = \vec{P} \cdot \hat{n} = P \cos \theta$

$$V = \frac{1}{4\pi\epsilon_0} \left[\oint_S \frac{\vec{P} \cdot \hat{n}}{r} ds' + \int_V \frac{1}{r} (-\vec{\nabla} \cdot \vec{P}(r')) dv' \right]$$

$$V(r) = \frac{1}{4\pi\epsilon_0} \left[\oint_S \frac{\sigma_B ds'}{r} + \int_V \frac{\rho_B dv'}{r} \right]$$

$$\left[\text{where } \underbrace{\sigma_B = \vec{P} \cdot \hat{n}}_{\text{total surface charge?}} \text{ and } \underbrace{\rho_B = -\vec{\nabla} \cdot \vec{P}}_{\text{total volume charge?}} \right]$$

B: bound in σ_B, ρ_B

$$\begin{aligned} \rho_{\text{volume (bound)}} &= \int \rho_B(r) dv \\ &= \int (-\vec{\nabla} \cdot \vec{P}) dv \\ &= -\oint \vec{P} \cdot d\vec{s} \\ &= -\oint (\vec{P} \cdot \hat{n}) ds \end{aligned}$$

So, total surface charge is equal to negative of total volume charge. $= -\oint \sigma_B ds = -\rho_{\text{surface (bounded)}}$

irrespective of the type of polarization vector, whether uniform or not, the total surface and total vol charge add up to zero

Note: We are considering bounded charge

\downarrow
none of them are mobile
they are tied to the molecules
and molecules are tied to their pos.
they are bound to the material.

$$\text{Remember } \rightarrow \vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho_{\text{free}} + \rho_{\text{bound}}}{\epsilon_0}$$

$$\vec{\nabla} (\epsilon_0 E) = \rho_{\text{free}} - \vec{\nabla} \cdot \vec{P}$$

$$\rho_{\text{free}} = \vec{\nabla} \cdot (\vec{P} + \epsilon_0 E)$$

$\chi_e \Rightarrow$ Electric susceptibility

easier to polarize if χ_e / ϵ_0 is higher

$$\vec{P} = \underbrace{\alpha}_{\text{atomic polarizability}} \vec{E} = \epsilon_0 \chi_e \vec{E}$$

$$\text{So, } \rho_{\text{free}} = \vec{\nabla} \cdot (\epsilon_0 \chi_e \vec{E} + \epsilon_0 \vec{E}) = \epsilon_0 (1 + \chi_e) \vec{\nabla} \cdot \vec{E}$$

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho_{\text{free}}}{\epsilon_0 (1 + \chi_e)} = \frac{\rho_{\text{free}}}{\epsilon_0 \epsilon_r} = \frac{\rho_{\text{free}}}{\epsilon}$$

$$\epsilon_r = 1 + \chi_e$$

for free space $\rightarrow \chi_e = 0$

so $\epsilon_r = 1$

$$\boxed{\vec{\nabla} \cdot \vec{E} = \frac{\rho_{\text{free}}}{\epsilon}}$$

\vec{D} : displacement density

\vec{I}_D : displacement current

$$\vec{I}_D = \epsilon_0 \frac{\partial \vec{E}}{\partial t} = \frac{\partial \vec{D}}{\partial t}$$

$$\vec{D} = \epsilon_0 \vec{E}$$

displacement current is the reason current flows between the plates of the capacitor.

$$\rightarrow i_c = 0 \text{ because } Z_c = \frac{1}{j\omega C} \therefore \omega \rightarrow \infty \text{ for DC}$$

there is no movement of charge
dipoles re-orienting/rotating themselves due to \vec{E} and there is a change of charge density and so there is an illusion of charge movement and that's how current moves