

Lecture 3

* Wave Function

$\Psi(x)$: varies with time
(Born interpretation) or $\Psi(r, t)$

for 1 particle system: $\Psi = \Psi(x, y, z, t)$

for 2 particle system: $\Psi = \Psi(x_1, y_1, z_1, x_2, y_2, z_2, t)$

wave function contains all info about the system

wave function \longleftrightarrow classical trajectory
(quantum mechanics) (newtonian mechanics)

Probability density

$$\hookrightarrow P(r) = |\Psi|^2 = \int \Psi^* \Psi d\tau \quad \text{--- (1)}$$

(1) COPENHAGEN/BORN'S INTERPRETATION

The probability that a particle can be found

at a point x at time t is given by (1)

Ψ can be represented as a linear combination of other Ψ s

(2) To every physical property, ^{observable in classical mechanics} there corresponds a linear, Hermitian operator in quantum mechanics

	SYMBOL	OPERATOR	
Position	x	\hat{X}	$x \cdot x$
	y	\hat{Y}	$x \cdot y$
Momentum	p_x	\hat{P}_x	$-i\hbar \frac{\partial}{\partial x}$
	p_y	\hat{P}_y	$-i\hbar \left(i \frac{\partial}{\partial x} + j \frac{\partial}{\partial y} + k \frac{\partial}{\partial z} \right)$
	p_z	\hat{P}_z	$-i\hbar \left(i \frac{\partial}{\partial x} + j \frac{\partial}{\partial y} + k \frac{\partial}{\partial z} \right)$
KE	T_x	\hat{T}_x	$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2}$
	T	\hat{T}	$-\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right)$
			$= -\frac{\hbar^2}{2m} \nabla^2$

PE	$V(x)$	$\hat{V}(x)$	$x \cdot V(x)$
	$V(x, y, z)$	$\hat{V}(x, y, z)$	$x \cdot V(x, y, z)$

Total Energy $= KE + PE$	E	\hat{H}	$-\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) + V(x, y, z)$
			$= -\frac{\hbar^2}{2m} \nabla^2 + V(x, y, z)$

NOTE: Hermitian matrix:

if $A = A^T$

(3) $\hat{A}\Psi = Q\Psi$ --- (2)

In any measurement of the observable associated

with operator \hat{A} , the only values that will be

observed are the eigenvalues 'a' which satisfy: (2)

$$\hat{A}f(x) = kf(x)$$

$f(x)$: eigenfunction of \hat{A}

with eigenvalue k

the eigenvalue

gives the real

observable

eg: e^{ikx} is an eigenfunction

of the operator $\hat{P}_x = -i\hbar \frac{\partial}{\partial x}$

if I operate \hat{P}_x on Ψ , i.e. $\left(-i\hbar \frac{\partial}{\partial x}\right)\Psi$

and we get: $\# \Psi$

\hookrightarrow some number

then $\#$ is eigenvalue

here $\Psi = e^{ikx}$

$$-i\hbar \frac{\partial}{\partial x} e^{ikx} = -i\hbar k^2 e^{ikx} = \hbar k^2 e^{ikx}$$

$k \equiv$ wave number

$$k = \frac{2\pi}{\lambda}$$

$\lambda \equiv$ wave length

$$p = \frac{h}{\lambda} = \frac{h}{2\pi/k} = \hbar \cdot k$$

\hookrightarrow reduced plank constant

\hookrightarrow wave number

$$\hbar = \frac{h}{2\pi}$$

$$p = \hbar \cdot k$$

hence, Momentum space is also called k -space

and we can swop between the

two using a real number \hbar

$\Psi = e^{ikx}$ ^{momentum space} \rightarrow real space

time independent wave func

$\Psi = e^{i(kx - \omega t)}$

time dependent wave func

(4) SCHRODINGER EQUATION

$$i\hbar \frac{\partial \Psi}{\partial t} = \hat{H}\Psi \quad \text{time dependent}$$

$$E = \hat{H}\Psi \quad \text{time independent}$$

$\hat{H} \equiv$ Hamiltonian

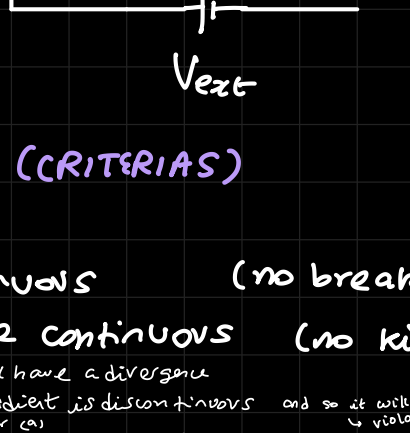
$$\hat{H} = \hat{T} + \hat{V} + V_{ext}$$

$TE = KE + PE$

total energy

Asymmetric

Symmetric



* BORN'S Interpretation (CRITERIAS)

- (a) Ψ must be continuous (no breaks)
- (b) $\nabla \Psi = \frac{\partial \Psi}{\partial x}$ must be continuous (no kinks)
gradient of Ψ
- (c) Ψ must have a single value at any pt. in space
we will have a divergence if gradient is discontinuous and so it will violate (b)
- (d) Ψ must be finite everywhere
- (e) Ψ cannot be zero everywhere

Quantization of the wave function

These restrictions on Ψ ensures that

only certain wavefunctions and only

certain energies of system are allowed

\therefore Quantization of $\Psi \Rightarrow$ Quantization of E

Examples

(1) Particle moving in 1 dimension

Cases:

- $PE = 0 \quad \forall x$
- $PE = 0$ for certain x
- $PE = \infty \quad \forall x$

Case 1: $PE \equiv 0$

$$\hat{H}\Psi = E\Psi; \hat{H} = \hat{T} + \hat{V} \xrightarrow{\text{zero}} \hat{H} = \hat{T}$$

$$\hat{H}\Psi = E\Psi \Rightarrow \hat{T}\Psi = E\Psi \Rightarrow -\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} \right) \Psi = E\Psi$$

$$\hat{T} = \frac{\hat{p}^2}{2m}; \hat{p} = -i\hbar \frac{\partial}{\partial x}; \hat{p}^2 = -\hbar^2 \frac{\partial^2}{\partial x^2}$$

$$\hat{T} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2}$$

$$\Psi = A \sin(kx) + B \cos(kx) = A e^{ikx}$$

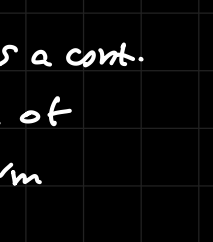
$$\frac{\partial^2 \Psi}{\partial x^2} = -k^2 (A \sin(kx) + B \cos(kx)) = -k^2 \Psi$$

Energy Dispersion Relation OR E-K Relation

$$E = \frac{k^2 \hbar^2}{2m}$$

(E depends on k)

$$E \propto k^2$$



We know $p = \hbar k$

Energy is a cont.

function of

momentum

$$\text{So, } E = \frac{p^2}{2m}$$

\therefore we can find out E from k -space

wave num \rightarrow momentum \rightarrow position

we can represent E in terms of x, y, z

if we have boundary condition like

$\lim_{x \rightarrow 5}$, we are also restricting k