

Lecture

$$\frac{E}{H} = \sqrt{\frac{\mu}{\epsilon}} = \eta : \Omega$$

$$W_e = \frac{1}{2} \epsilon |E|^2$$

$$W_m = \frac{1}{2} \mu |H|^2$$

$$W_e + W_m = \frac{1}{2} (\epsilon |E|^2 + \mu |H|^2)$$

$$W_e = \frac{1}{2} \epsilon \int |E|^2 dz$$

$$W_m = \frac{1}{2} \mu \int |H|^2 dz$$

$$\bar{\nabla} \cdot (\bar{E} \times \bar{H}) = \bar{H} \cdot (\bar{\nabla} \times \bar{E}) - \bar{E} \cdot (\bar{\nabla} \times \bar{H})$$

$$= \bar{H} \cdot \left(-\mu \frac{\partial \bar{H}}{\partial t} \right) - \bar{E} \cdot (\dots)$$

$$= -\mu \bar{H} \cdot \frac{\partial \bar{H}}{\partial t} - \bar{E} \cdot \left(\bar{J} + \epsilon \frac{\partial \bar{E}}{\partial t} \right)$$

$$= -\frac{\mu}{2} \frac{\partial |H|^2}{\partial t} - \bar{E} \cdot \bar{J} - \frac{\epsilon}{2} \frac{\partial |E|^2}{\partial t}$$

$$= -\frac{\partial}{\partial t} \left\{ \frac{1}{2} (\mu |H|^2 + \epsilon |E|^2) - \bar{E} \cdot \bar{J} \right\}$$

$$-\frac{\partial}{\partial t} (w_e + w_m) = \overline{\mathbf{E}} \cdot \overline{\mathbf{J}} + \overline{\nabla} \cdot (\overline{\mathbf{E}} \times \overline{\mathbf{H}})$$

(heat) (outgoing)

↪ decay rate of stored energy
in an electromagnetic field

we said in the last class

$$\overline{\mathbf{E}} \equiv \text{Voltage}$$

$$\overline{\mathbf{B}} \equiv \text{current}$$

so, $\overline{\mathbf{E}} \cdot \overline{\mathbf{J}} \approx \text{Power}$

$$dw \overline{\mathbf{F}} \cdot d\overline{\mathbf{e}} = q(\overline{\mathbf{E}} + \underbrace{(\overline{\mathbf{v}} \times \overline{\mathbf{B}})}_0) \cdot \overline{\mathbf{v}} dt$$

$$dw \overline{\mathbf{F}} \cdot d\overline{\mathbf{e}} = q \overline{\mathbf{E}} \cdot \overline{\mathbf{v}} dt$$

$$\frac{dw}{dt} = \oint d\mathbf{z} \overline{\mathbf{E}} \cdot \overline{\mathbf{v}}$$

$$\frac{dw}{dt} = \overline{\mathbf{E}} \cdot \overline{\mathbf{J}} d\mathbf{z}$$

The energy is either used by doing
some work ($\overline{\mathbf{E}} \cdot \overline{\mathbf{J}}$) or comes out
of the surface ($\overline{\nabla} \cdot (\overline{\mathbf{E}} \times \overline{\mathbf{H}})$)

$$\begin{aligned} -\frac{\partial}{\partial t} \int (w_e + w_m) d\tau &= \int \bar{\mathbf{E}} \cdot \bar{\mathbf{J}} d\tau + \int \bar{\nabla} \cdot (\bar{\mathbf{E}} \times \bar{\mathbf{H}}) d\tau \\ &= \int \bar{\mathbf{E}} \cdot \bar{\mathbf{J}} d\tau + \oint (\bar{\mathbf{E}} \times \bar{\mathbf{H}}) d\bar{\mathbf{S}} \end{aligned}$$

$$\text{So, } P_{\text{outgoing}} = \oint (\bar{\mathbf{E}} \times \bar{\mathbf{H}}) d\bar{\mathbf{S}}$$

$\underbrace{\qquad\qquad\qquad}_{\vec{S}}$
 (Poynting vector)

$$\bar{\mathbf{S}} = \bar{\mathbf{E}} \times \bar{\mathbf{H}}$$

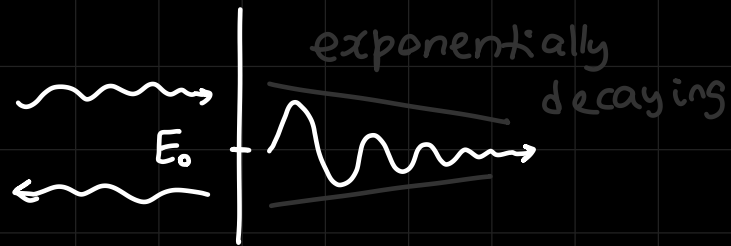
$$|\bar{\mathbf{S}}| = |\mathbf{E}| |\mathbf{H}| = |\mathbf{E}| \cdot \frac{|\mathbf{E}|}{\mu} = \frac{|\mathbf{E}|^2}{\mu} : \approx \text{power per unit area}$$

(Watt/m²)

last topic for EM waves →
 next: transmission lines

SOMETHING4IDK123

$$\epsilon_c = \epsilon \left(1 - \frac{i\sigma}{\omega\epsilon}\right)$$



$$k = \omega \sqrt{\mu\epsilon_c}$$

$$k = k' - ik''$$

$$e^{-ikz} = e^{-i(k' - ik'')z} = e^{-ik'z} e^{-k''z}$$

Skin depth of a medium:

$$\delta = \frac{1}{k''} \quad \rightarrow \quad 0 \text{ for a PEC } (\infty \times)$$

distance for the signal to decay

for $z = \delta$,

$$\bar{E} = E_0 e^{-ik'z} e^{-k''/k''} = \frac{E_0}{e} e^{-ik'z}$$

$$k^2 = \omega^2 \mu \epsilon \left(1 - \frac{i\sigma}{\omega\epsilon}\right)$$

$$(k' - ik'')^2 = k'^2 - k''^2 - 2ik'k'' = \omega^2 \mu \epsilon - i\omega\sigma\mu$$

$$2k'k'' = \omega\sigma\mu$$

$$k'^2 - k''^2 = \omega^2 \mu \epsilon \quad \text{--- ①}$$

for simplicity, we find

$$\begin{aligned}(k'^2 + k''^2)^2 &= (k'^2 - k''^2)^2 + 4(k'k'')^2 \\ &= \omega^4 \mu^2 \epsilon^2 + \omega^2 \mu^2 \sigma^2\end{aligned}$$

$$(k'^2 + k''^2) = \omega^2 \mu \epsilon \left[1 + \frac{\sigma^2}{\omega^2 \epsilon^2} \right]^{1/2} \quad \text{--- (2)}$$

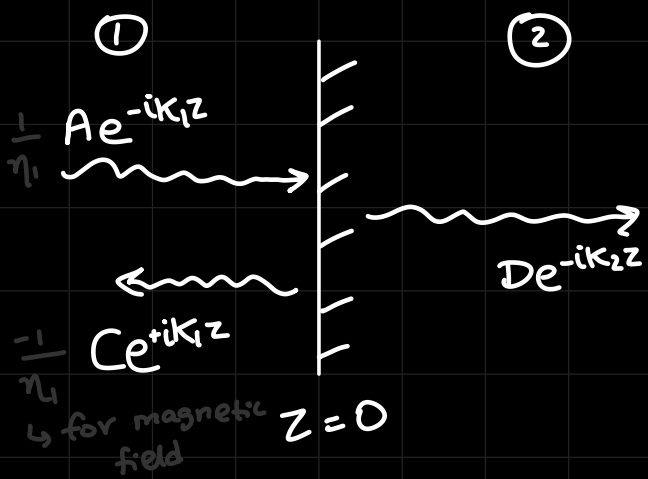
using ① and ② \rightarrow

$$k' = \omega \sqrt{\frac{\mu \epsilon}{2}} \left[1 + \left(1 + \frac{\sigma^2}{\omega^2 \epsilon^2} \right)^{1/2} \right]^{1/2}$$

$$k'' = \omega \sqrt{\frac{\mu \epsilon}{2}} \left[\left(1 + \frac{\sigma^2}{\omega^2 \epsilon^2} \right)^{1/2} - 1 \right]^{1/2}$$

Note: Lossless medium \rightarrow no decay

We figure out reflection and transmission of the EM waves



assuming SRC \equiv medium ①
transmitted \equiv medium ②

$$\Gamma(z) = \frac{Ce^{-ik_1z}}{Ae^{-ik_1z}}$$

Reflection coefficient:

(gamma)

$$\Gamma(z) = \frac{C}{A} e^{i2k_1z}$$

at the interface $\rightarrow z=0$

$$\Gamma_0 = \frac{C}{A}$$

$$\Gamma(z) = \Gamma_0 e^{i2k_1z}$$

lossless medium \equiv no
conduction
current
(pure dielectric)

for a lossless medium $\rightarrow \Gamma_0$ decreases
as we move away from the interface
and $|\Gamma(z)| = |\Gamma_0|$

periodicity:

$$e^{i2\pi} = 1 \rightarrow 2k_1 z = 2\pi$$

$$2 \cdot \frac{2\pi}{\lambda} \cdot z = 2\pi \rightarrow$$

$$z = \frac{\lambda}{2}$$

Transmission coefficient : $T = \frac{D}{A}$

defined at the interface only,

not all points because there exists only the transmitted wave in medium 2.

$$\text{at } z=0: \quad \frac{E_{t1}}{H_{t1}} = \frac{E_{t2}}{H_{t2}}$$

$$\frac{A+C}{A/\eta_1 + C/(-\eta_1)} = \eta_2$$

$$\frac{1 + \Gamma_0}{\frac{1}{\eta_1}(1 - \Gamma_0)} = \eta_2$$

$$\frac{\eta_2}{\eta_1} = \frac{1 + \Gamma_0}{1 - \Gamma_0}$$

$$A + C = D$$

$$1 + C/A = D/A$$

$$\boxed{1 + \Gamma_0 = T}$$

$$T = 1 + \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1}$$

$$\rightarrow \boxed{T = \frac{2\eta_2}{\eta_1 + \eta_2}}$$