Lecture after midsem J(r') - A/m = dI $\overline{B}(\gamma) = \underbrace{\text{U.}}_{\text{4t}} \int \overline{J}(\gamma') \times \hat{x}$ $\nabla f = -\nabla f$ note: assuming steady current $\overline{\nabla} \cdot \overline{B} = \underline{u_0} \int \nabla (\overline{J} \times \hat{J}_1) = \underline{u_0} \int \underline{J}_1 (\overline{\nabla} \times \overline{J}_{(1)} - \overline{J}_1 (\overline{\nabla} \times \hat{J}_1))$ Zero because div $\nabla - (\bar{P} \times \bar{g}) = \overline{g} (\bar{\nabla} \times \bar{P}) - \bar{P} (\bar{\nabla} \times \bar{g})$ wrt r but J(r1) and since $\hat{x} = \nabla(\frac{1}{y})$, curl = 0

So, [V·B = 0]

Curl of magnetic field

$$\overline{B} = \overline{\nabla} \times \overline{A}$$

5) magnetic vector potential

(need not be unique)

Anew = A + PO 2 gouse treedom

DXAren = DXA+O = B always

if B gives, find the magnetic vector potential or given, J, find the magnetic field

Since $B \propto 1$ -> $A \propto 1$ because $B = \nabla \times A$

TXA = Uo S (J(r) x si) dV'

We know $\frac{\hat{x}}{x^2} = -P(\frac{1}{x})$

 $\nabla \times \overline{A} = u_0 \int \left(\overline{T}(r') \times \overline{P}(\underline{L}) \right) dv'$

(PX Qf)

$$\nabla x(fP) = f(\nabla xP) - Px(\nabla f)$$

$$\neg P_{X}(\overline{\nabla}f) = f(\overline{\nabla}x\overline{P}) - \overline{\nabla}x(f\overline{P})$$

her
$$\nabla wrt & but T(r')$$

So, $f(\nabla x P) = 0$

$$\nabla \times A = M_0 \int \nabla \times (\overline{\mathcal{I}}(r)) dv'$$

$$\overline{A} = U_0 \int \overline{J(Y')} dV'$$

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Current carrying flamentary conductor

$$f = I S(x) S(y)$$
 $z = 0$ $J = dI$ da_1

$$\overline{A} = \underbrace{J_0}_{4\pi} \iiint \overline{J(r')} dx' dy' dz'$$

 $J(r') \times dz'dy' = I$ for $x, y \rightarrow infinitesimally small$

$$\overline{A} = \underbrace{M_0}_{\text{UIT}} \underbrace{\int_0^\infty S(x)S(y)dxdy}_{\text{L}} \cdot \underline{I}\hat{z} \underbrace{\int_0^1 dz}_{(\alpha-\pi^i)^2 + (y-y^i)^2 + (z-z^i)^2}$$

= function of x,y,z

curl (E) = 0 but div (E) => Gauss's Law

div(B) = 0 but cure of B => Ampere's Law

 $\overline{B} = \frac{u_0}{4\pi} \int_{0}^{\infty} \overline{J(r')} \times \frac{\hat{y_1}}{\hat{y_1}^2} dV'$

V×B = 2

笥

problem - connot apply chain rule to Px(Pxg)

$$\nabla \times (P \times \bar{g}) = (\bar{g} \cdot \nabla) \bar{P} - (\bar{P} \cdot \nabla) g + \bar{P} (\bar{v} \cdot \bar{g}) - \bar{g} (\bar{v} \cdot \bar{P})$$

$$\left(g_{\lambda} \frac{\partial P}{\partial y} + g_{y} \frac{\partial P}{\partial y} + g_{z} \frac{\partial P}{\partial z}\right)$$

non zero term:

$$\nabla \times (P \times \bar{g}) = P(\bar{v} \cdot \bar{p}) - (\bar{v} \cdot \bar{v}) \bar{g}$$

$$\overline{\nabla} \times \overline{B} = \underbrace{u_0 \left(\overline{J} (\overline{\nabla} \cdot (\overline{S})) dV' + \int (\overline{J} \cdot \overline{\nabla}') \widehat{S} dV' \right)}_{\text{MT}} dV'$$

$$\underbrace{becouse}_{\text{becouse}} \overline{\nabla} f = -\overline{\nabla} f$$

assuming that 2nd term is zero

$$\int (\nabla \times \vec{B}) \cdot ds = u_0 \int \vec{J} \cdot d\vec{s}$$

Proving the previous assumption

Steady current
$$\rightarrow \nabla \overline{f}(r') = 0$$

So, $\nabla ' \overline{f}(r') = 0$ also

So,
$$\nabla' f(r') = 0$$
 also

We ned to prove =
$$\int (\overline{J}(ri) \cdot \overline{D}) \frac{\hat{J}}{\hat{J}} dv$$

$$\hat{x} = \hat{y} = \hat{x}(x-x')+\hat{y}(y-y')+\hat{z}(z-z')$$

We need contribution of entire system but we will be checking for individual components $\int_{V} \left(\overline{J(Y')} \cdot \overline{\nabla} \right) \left(2c - 2c' \right) dv'$ $\overline{A} \cdot (\nabla f) = \nabla' (f \overline{A}) - \overline{f} (\overline{\nabla}' \overline{A})$ 0 peronso D. J=0 $= \int \nabla' \cdot (\Im(r)(x-x)) dv'$ $= \int_{S'} \overline{J}(v')(x-x') ds$ Since 5(1) =0 outsi de v' this I holds for v' x > as well \$\frac{1}{5(\(\cappa('))\) \surface \(\cappa(')\) \ds\\
\$5\ta D - Assumption correct 7250