

STRESS & STRAIN

- # Strain is "measurable" ~ easily.

• Stress can be something "relative".

* STRAIN IS MORE FUNDAMENTAL

→ Stress can be calculated by strain

→ Material response is the relation between "stress & strain"
↳ The output depends on it

⇒ HOOKE'S LAW

HOOKES LAW

Stress & Strain ————— "perceived quantity"
————— measurable

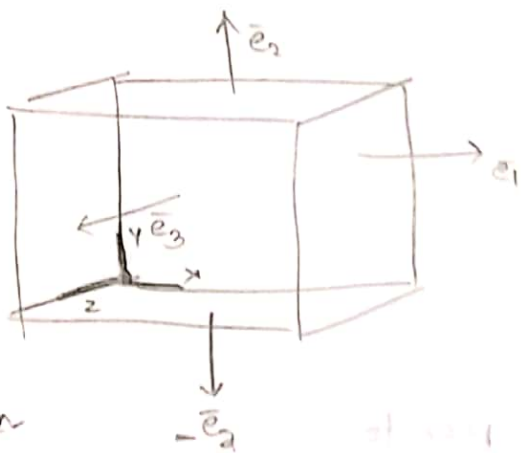
$$V \propto e \quad (\text{linear proportion})$$

$$|\sigma = Fe|$$

Young's Modulus Elasticity

~~Q~~ BOTH STRESS & STRAIN ARE "TENSORS" (SECOND ORDER)

"There's nothing called stress at point" on a infinitesimal cube but not point




$$\begin{bmatrix} \overline{0_{xx}} & \overline{0_{xy}} & \overline{0_{xz}} \\ \overline{0_{yx}} & \overline{0_{yy}} & \overline{0_{yz}} \\ \overline{0_{zx}} & \overline{0_{zy}} & \overline{0_{zz}} \end{bmatrix}$$

All the 9 components determine stress

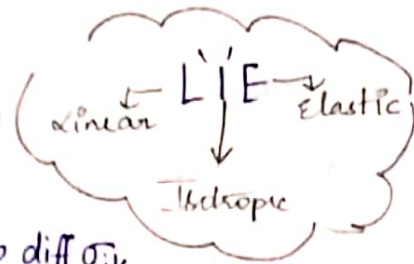
My Grain

$$\begin{bmatrix} e_{xx} & e_{xy} & e_{xz} \\ e_{yx} & e_{yy} & e_{yz} \\ e_{zx} & e_{zy} & e_{zz} \end{bmatrix}$$


 direction
 plane orientation

Generalised Hooke's Law :-

$$\sigma_{ij} \propto \epsilon_{kl}$$



All the strains are dep contribute to diff σ_{ik}

P₂

$$\sigma_{ij} = C_{ijkl} \epsilon_{kl}$$

9 components 9 components

STIFFNESS TENSOR (4th order Tensor)

$1/C \rightarrow$ Compliance Tensor

$$\epsilon_{ij} = S \sigma_{kl}$$

81 unique attributes required to differentiate materials.

with Isotropic — 2 are only req. — [Young's Modulus
Any other like Bulk's Modulus
Poisson's Ratio
Plasticity
Viscoplasticity]

Along way down the properties are same

All the others are dependent on them

Major Symmetry

$$\begin{bmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{yx} & \sigma_{yy} & \sigma_{yz} \\ \sigma_{zx} & \sigma_{zy} & \sigma_{zz} \end{bmatrix}$$

Not ferro magnetic material then top & bottom are symmetric

$$\sigma_{ij} = \sigma_{ji}$$

\Rightarrow dilatational - distortive
Size Shape



* finite deformation (large)
infinite " (small) \Rightarrow Assumption for "linearity"

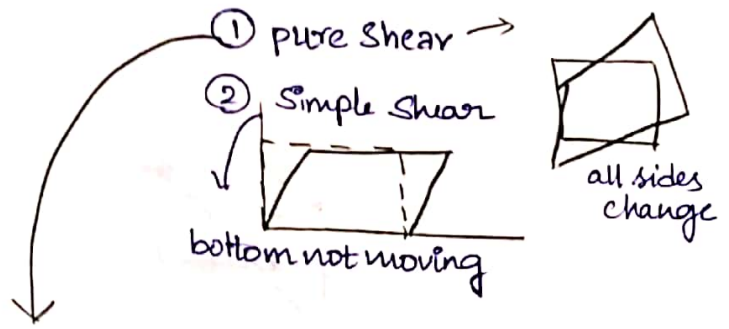
$$\epsilon_{ij} = \frac{1}{2} \left[\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right] = \epsilon_{ji}$$

$$\epsilon_{11} = \frac{\partial u_1}{\partial x_1}$$

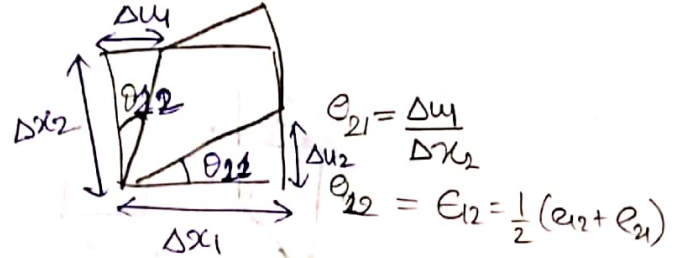
Stretch

$$\epsilon_{12} = \frac{1}{2} \left[\frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right]$$

distortion.



$$\epsilon_{ij} = \frac{1}{2} \left[\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right]$$



⇒ due to "linearity" — Superposition allowed (Retailer)
Non linear — Superposition not allowed (Whole sale)

⇒ Stress is symmetric unless no body movements

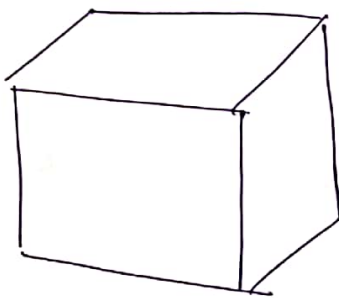
Strain is "SYMMETRIC"

$$\begin{array}{|c|c|c|} \hline 2 & 2 & 2 \\ \hline 1 & \underline{a} \otimes \underline{b} = \underline{c} & \hline \hline \end{array}$$

3⁰ scalar
3¹ vector
3² 2nd order Tensor
3⁴ 4th order Tensor

$\underline{a} \cdot \underline{b} = 0$

dot product due by 1

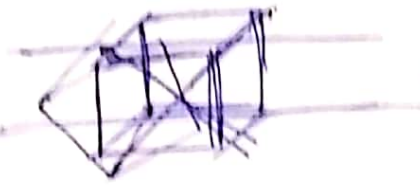


$e_i e_j$ → plane whose normal is in 'i' direction
It is in 'j' direction in the plane

$$\begin{aligned} \sigma &= \sigma_{ij} e_i \otimes e_j \quad \text{tensor product} \\ &= \sigma_{11} e_1 \otimes e_1 + \sigma_{12} e_1 \otimes e_2 + \sigma_{21} e_2 \otimes e_1 + \dots \end{aligned}$$

stress

$$\underline{C} = C_{ijkl} e_i \otimes e_j \otimes e_k \otimes e_l$$



P_m



$$\begin{aligned} \gamma &= \gamma_1 e_1 + \gamma_2 e_2 + \gamma_3 e_3 \\ &= \gamma_1^1 e_1^1 + \gamma_2^1 e_2^1 + \gamma_3^1 e_3^1 \end{aligned}$$

\Rightarrow For the vector \Rightarrow invariant: the magnitude

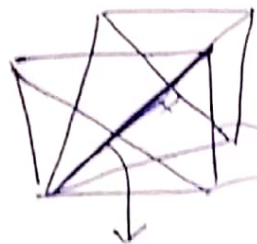
2nd For tensors (2nd order) \Rightarrow invariants $(\lambda_1, \lambda_2, \lambda_3)$ eigenvalues, order

$$\text{trace } I_1 = \lambda_1 + \lambda_2 + \lambda_3$$

$$\text{trace of Adj matrix } I_2 = \lambda_1 \lambda_2 + \lambda_2 \lambda_3 + \lambda_3 \lambda_1$$

$$\text{determinant } I_3 = \lambda_1 \lambda_2 \lambda_3$$

* 6 independent stress components



Rhombohedron

here both are \parallel

both diagonal coincide



$$Ax = b$$

In this transformation few transformations follow this relation: "few directions retain their direction"

$$|Ax = \lambda x|$$

$$\Rightarrow (A - \lambda I)x = 0 \Rightarrow |A - \lambda I| = 0$$

↳ characteristic equations

$$\lambda^3 - I_1 \lambda^2 + I_2 \lambda - I_3 = 0$$

characteristic polynomial

Caley Hamilton Theorem

$$A^3 - I_1 A^2 + I_2 A - I_3 I = 0$$

Finding inverse:

$$A^2 - I_1 A + I_2 I = I_3 A^{-1}$$

physical meaning of eigen values: there are certain orientation
i.e. in this case 3, where
it "stretches" but not distorted

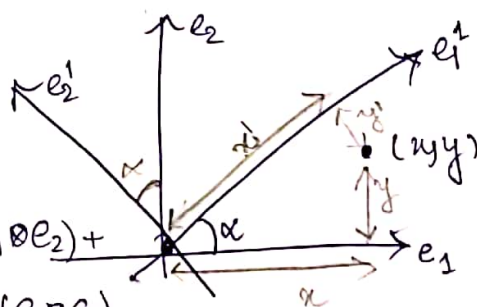
⇒ 6 independent stress components: $\lambda_1, \lambda_2, \lambda_3$ & x_1, x_2, x_3

↓
6 are ~~needed~~ present but only 3 are needed.

↳ only for the orientations which "stretch" [Eigenvalues]

$$\underline{\sigma} = \begin{bmatrix} \sigma_{11} & \sigma_{12} \\ \sigma_{21} & \sigma_{22} \end{bmatrix}$$

$$\begin{aligned} \underline{\sigma} &= \sigma_{11}(e_1 \otimes e_1) + \sigma_{12}(e_1 \otimes e_2) + \\ &\quad \sigma_{21}(e_2 \otimes e_1) + \sigma_{22}(e_2 \otimes e_2) \\ &= \sigma_{11}'(e_1' \otimes e_1') + \sigma_{12}'(e_1' \otimes e_2') + \\ &\quad \sigma_{21}'(e_2' \otimes e_1') + \sigma_{22}'(e_2' \otimes e_2') \end{aligned}$$



$$\begin{aligned} * e_1' &= e_1 \cos \alpha + e_2 \sin \alpha \\ e_2' &= e_2 \cos \alpha - e_1 \sin \alpha \end{aligned}$$

$$\begin{aligned} \sigma_2 = & \sigma_{11}^{-1} [(e_1 C_\alpha + e_2 S_\alpha) \otimes (e_1 C_\alpha + e_2 S_\alpha)] + \\ & \sigma_{12}^{-1} [(e_1 C_\alpha + e_2 S_\alpha) \otimes (e_2 C_\alpha - e_1 S_\alpha)] + \\ & \sigma_{21}^{-1} [(e_2 C_\alpha - e_1 S_\alpha) \otimes (e_1 C_\alpha + e_2 S_\alpha)] + \\ & \sigma_{22}^{-1} [(e_2 C_\alpha - e_1 S_\alpha) \otimes (e_2 C_\alpha - e_1 S_\alpha)] \end{aligned}$$

$$\Rightarrow e_1 \otimes e_1 [C_\alpha^2 \sigma_{11}^{-1} - \sigma_{12}^{-1} C_\alpha S_\alpha - \sigma_{21}^{-1} C_\alpha S_\alpha + \sigma_{22}^{-1} S_\alpha^2]$$

$$\star \sigma_{12}^{-1} = \sigma_{21}^{-1} \text{ [Symmetric Matrix]}; e_1 \otimes e_2 = e_2 \otimes e_1$$

$$\Rightarrow \sigma_{11} = [C_\alpha^2 \sigma_{11}^{-1} + \sigma_{22}^{-1} S_\alpha^2 - 2\sigma_{12}^{-1} S_\alpha C_\alpha]$$

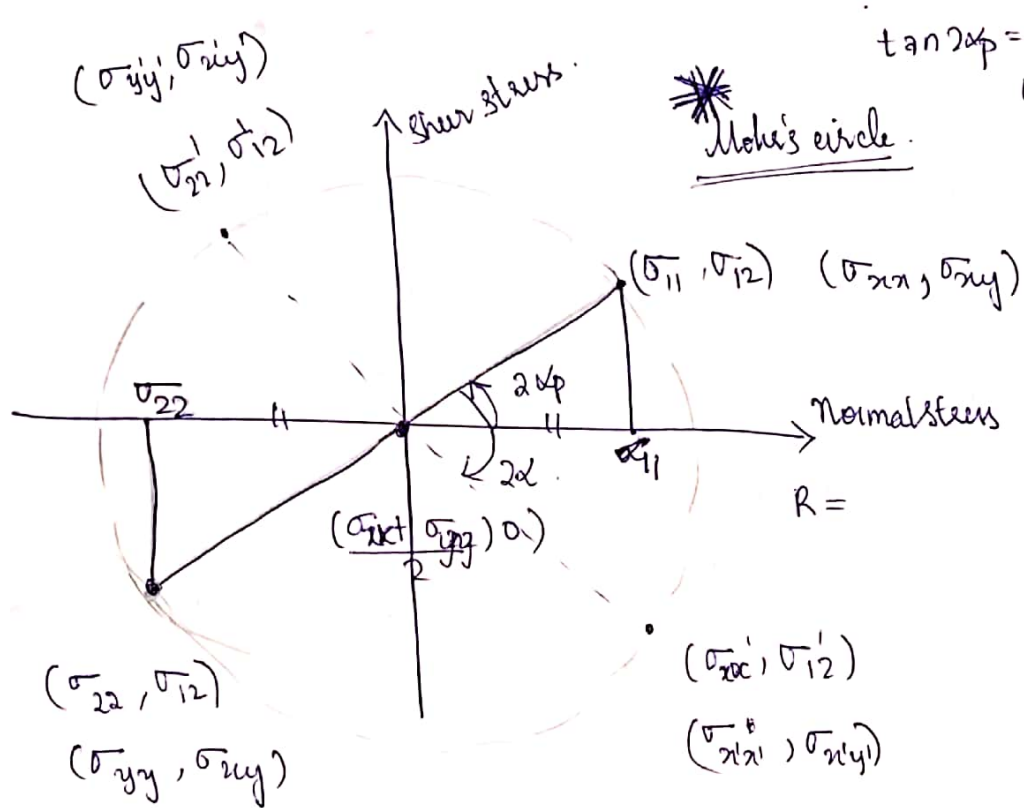
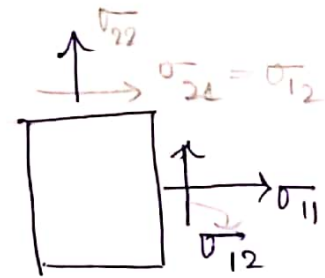
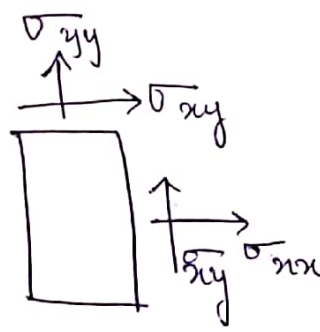
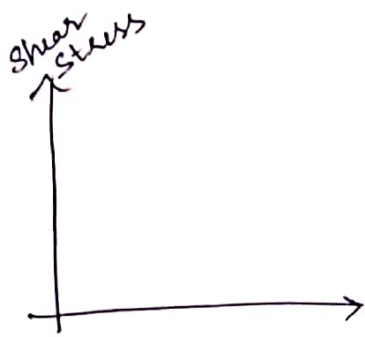
$$\sigma_{22} = [S_\alpha^2 \sigma_{11}^{-1} + 2\sigma_{12}^{-1} C_\alpha S_\alpha + C_\alpha^2 \sigma_{22}^{-1}]$$

$$\sigma_{21} = \sigma_{12} = [\sigma_{11}^{-1} C_\alpha S_\alpha + \sigma_{12}^{-1} (C_\alpha^2 - S_\alpha^2) - \sigma_{22}^{-1} S_\alpha C_\alpha]$$

$$\Rightarrow R = \begin{bmatrix} C_\alpha & S_\alpha \\ -S_\alpha & C_\alpha \end{bmatrix}$$

$$\boxed{\sigma = R^T \sigma R}$$

$$\star \alpha_p : \text{when } \sigma_{21} = \sigma_{12} = 0$$



$$\tan 2\phi = \frac{\sigma_{xy}}{\frac{\sigma_{xx} - \sigma_{yy}}{2}}$$

Mohr's circle.

$$I_{x'x'} = \left(\frac{I_{xx} + I_{yy}}{2} \right) + \left(\frac{I_{xx} - I_{yy}}{2} \right) C_{2\alpha} - I_{xy} S_{2\alpha}$$

$$I_{y'y'} = \left(\frac{I_{xx} + I_{yy}}{2} \right) + \left(\frac{I_{xx} - I_{yy}}{2} \right) C_{2\alpha} + I_{xy} S_{2\alpha}$$

$$I_{x'y'} = \left(\frac{I_{xx} - I_{yy}}{2} \right) S_{2\alpha} + I_{xy} C_{2\alpha}$$

⇒ ⇒

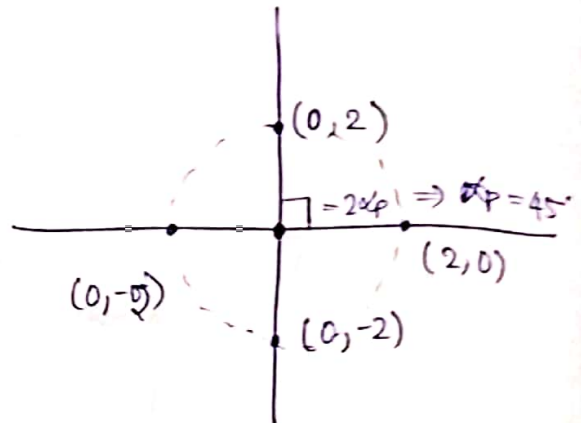
$$\begin{bmatrix} \sigma_{xx} & \sigma_{xy} \\ \sigma_{xy} & \sigma_{yy} \end{bmatrix}$$

Example

7

Pure stress

$$\underline{\sigma} = \begin{bmatrix} 0 & 2 \\ 2 & 0 \end{bmatrix}$$



$\sigma_1 > \sigma_2$

$$\det(A - \lambda I) = 0$$

$$\sigma_1 = 2$$

$$\sigma_2 = -2$$

$$\begin{bmatrix} 0 & 2 \\ 2 & 0 \end{bmatrix} - \lambda \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} -\lambda & 2 \\ 2 & -\lambda \end{bmatrix}$$

⇓

$\lambda = \pm 2 \Rightarrow$ this gives the principal values
 $(\sigma_1, \sigma_2) = (2, -2)$

Dynamics

⇒ CLOSED SYSTEM: Conservation of mass!

Conservation of linear momentum!

Conservation of Angular momentum!

Conservation of Energy!

} trivial

$$\text{Dissipation Energy} \geq 0$$

Law of motion: $F = \frac{d(mv)}{dt} = m \frac{dv}{dt} = ma$