

①

$$E_n^{(0)} = \frac{\pi^2 \hbar^2}{2m a^2}, \quad \psi_n^{(0)} = \sqrt{\frac{2}{a}} \sin \frac{n\pi x}{a}$$

$E_n^{(1)}$ (First Order Correction)

$$= \langle \psi_n^{(0)} | H' | \psi_n^{(0)} \rangle$$

$$\text{where } H' = V_0 \cos\left(\frac{\pi x}{a}\right)$$

$$= \int_0^{a/2} \frac{2}{a} \sin^2\left(\frac{\pi x}{a}\right) V_0 \cos\left(\frac{\pi x}{a}\right) dx$$

$$= \frac{2V_0}{a} \int_0^{a/2} \sin^2\left(\frac{\pi x}{a}\right) \cos\left(\frac{\pi x}{a}\right) dx$$

$$= \frac{2V_0}{a} \left[\sin^3\left(\frac{\pi x}{a}\right) \right]_0^{a/2} \Rightarrow \frac{2V_0}{a} \Delta E = 2V_0$$

②

③ Perturbation theory is used to calculate small disturbance in Hamiltonian of the system.

$$H = H_0 + H' \quad \text{where} \quad \begin{array}{l} H \rightarrow \text{perturbed Hamiltonian} \\ H_0 \rightarrow \text{unperturbed} \end{array}$$

~~say take~~ * For non-degenerate case,

$$\text{say take} \quad u = a + \epsilon u^2$$

if ϵ is small,

$$u = u_0 + \epsilon u_1 + \epsilon^2 u_2 + \dots$$

Substitute series in eq(1).

$$u = a + \epsilon a^2 + 2\epsilon^2 a^3 + \dots$$

exact solution

$$u = \frac{1}{2\epsilon} \left[1 - \left(1 - 4\epsilon \right)^{1/2} \right]$$

$$= a + \epsilon a^2 + 2\epsilon^2 a^3 + \dots$$

$$\boxed{H = H_0 + g H'}$$

$$\psi_n = \psi_n^{(0)} + g \psi_n^{(1)} + g^2 \psi_n^{(2)} + \dots$$

$$W_n = W_n^{(0)} + g W_n^{(1)} + g^2 W_n^{(2)} + \dots$$

$$\begin{aligned} & (H_0 + g H') (\psi_n^{(0)} + g \psi_n^{(1)} + g^2 \psi_n^{(2)} + \dots) \\ &= \left(W_n^{(0)} + g W_n^{(1)} + g^2 W_n^{(2)} + \dots \right) \left(\psi_n^{(0)} + g \psi_n^{(1)} + g^2 \psi_n^{(2)} + \dots \right) \end{aligned}$$

$$\hookrightarrow H_0 \psi_n^{(0)} = W_n^{(0)} \psi_n^{(0)}$$

$$H_0 \psi_n^{(1)} + H' \psi_n^{(0)} = W_n^{(0)} \psi_n^{(1)} + \psi_n^{(0)} W_n^{(1)}$$

$$H_0 \psi_n^{(2)} + H' \psi_n^{(1)} = W_n^{(0)} \psi_n^{(2)} + W_n^{(1)} \psi_n^{(1)} + W_n^{(2)} \psi_n^{(0)}$$

First Order

$$\psi_n^{(1)} = \sum_m a_m^{(1)} u_m$$

$$\sum_m a_m^{(1)} E_m u_m + H' u_n = E_n \sum_m a_m^{(1)} u_m + W_n^{(1)} u_n$$

$$a_k^{(1)} (E_n - E_k) + W_n^{(1)} \delta_{kn} = H'_{kn}$$

$k \neq n$

$$W_n^{(1)} = H'_{nn} = \int u_k^* H' u_n d\tau = \langle n | H' | n \rangle$$

$$a_k^{(1)} = \left(\frac{H'_{kn}}{E_n - E_k} \right), \quad k \neq n$$

Second order

$$\psi_n^{(2)} = \sum_m a_m^{(2)} \psi_m$$

$$\sum_m a_m^{(2)} E_m \delta_{km} + \sum_m a_m^{(1)} H'_{km} = E_n \sum_m a_m^{(2)} \delta_{km} + W_n^{(2)} \delta_n + H'_{nn} \sum_m a_m^{(1)} \delta_{km}$$

KOH

$$W_n^{(2)} = \sum_m a_m^{(1)} - a_n^{(1)} H'_{nn}$$

$$= \sum_{m \neq n} a_m^{(1)} H'_{nm} = \sum_{m \neq n} \frac{|H'_{mn}|^2}{E_n - E_m}$$

④

$$\langle E \rangle = \frac{\hbar^2}{2m} \langle \psi | \frac{\partial^2}{\partial x^2} | \psi \rangle + g \langle \psi | |x| | \psi \rangle$$

$$\text{as } \psi(x) = \begin{cases} \sqrt{\frac{c}{a^5}} (a^2 - x^2), & x < |a| \\ 0, & x > |a| \end{cases}$$

$$\langle \psi | \frac{\partial^2}{\partial x^2} | \psi \rangle = \int_{-a}^a \sqrt{\frac{c}{a^5}} (a^2 - x^2) \frac{\partial^2}{\partial x^2} \left(\sqrt{\frac{c}{a^5}} (a^2 - x^2) \right) dx$$

$$= \frac{c}{a^5} \int_{-a}^a (a^2 - x^2) (-2) dx$$

$$= \frac{-4c}{a^5} \left[a^2 x - \frac{x^3}{3} \right]_0^a$$

$$= \frac{-4c}{a^5} \left[\frac{2a^3}{3} \right] = \frac{-8c}{3a^2}$$

$$\Rightarrow \frac{-\hbar^2}{2m} \langle \psi | \frac{\partial^2}{\partial x^2} | \psi \rangle = \frac{4c \hbar^2}{3ma^2}$$

$$g \langle \psi | |x| | \psi \rangle = \frac{c}{a^5} \int_{-a}^a (a^2 - x^2)^2 |x| dx.$$

$$= \frac{c}{a^5} \left[\int_{-a}^0 (-x) (a^2 - x^2)^2 dx + \int_0^a x (a^2 - x^2)^2 dx \right]$$

$$= \frac{2c}{a^5} \int_0^a x (a^2 - x^2)^2 dx$$

$$= \frac{2c}{a^5} \int_0^a (a^4 x + x^5 - 2a^2 x^3) dx$$

$$= \frac{2c}{a^5} \left[\frac{6a^6 - 2a^6 - 6a^6}{12} \right] = \frac{ac}{3} g.$$

~~EE~~ = Normalization $\int_{-a}^a \psi^* \psi = 1$

$$\Rightarrow \frac{c}{a^5} \int_{-a}^a (a^2 - x^2)^2 dx = 1$$

$$\frac{2c}{a^5} \left[\frac{8a^5}{15} \right] = 1$$

$$\boxed{c = \frac{15}{16}}$$

$$\langle E \rangle = \frac{4\hbar^2}{3ma^2} \left(\frac{15}{16} \right) + \frac{ga}{3} \left(\frac{15}{16} \right)$$

$$\frac{d\langle E \rangle}{da} = 0; \quad -\frac{5}{4} \frac{\hbar^2}{m} \cdot 2a^{-3} + \frac{g}{16} = 0$$

$$a^3 = \frac{8\hbar^2}{mg} = a = \left(\frac{8\hbar^2}{mg} \right)^{1/3}$$

$$\begin{aligned}\langle E \rangle &= \frac{5\hbar^2}{4m} \left(\frac{mg}{8\hbar^2} \right) + \frac{5g^2}{16} = \frac{5g^2}{16} \left(\frac{3}{2} \right) \left(\frac{8\hbar}{mg} \right)^{1/3} \\ &= \frac{15}{16} \left(\frac{\hbar^2 g^2}{m} \right)^{1/3}.\end{aligned}$$

$$(2) \quad E_n^{(1)} = \langle n | H' | n \rangle = 0$$

$$\underline{E_n^{(2)}} = \sum_{m \neq n} \frac{|\langle m | H' | n \rangle|^2}{E_n - E_m} = \frac{\hbar^2 \lambda^2}{2m\omega} \sum_{m \neq n} \frac{|\langle m | \hat{a}^2 + \hat{a}^{\dagger 2} | n \rangle|^2}{E_n - E_m}$$

$$= \frac{\hbar^2 \lambda^2}{2m\omega} \left| \frac{\sqrt{n+1} \langle m | n+1 \rangle + \sqrt{n} \langle m | n-1 \rangle}{E_n - E_m} \right|^2$$

$$= (n) \left| \frac{\sqrt{n+1} \delta_{m, n+1} + \sqrt{n} \delta_{m, n-1}}{E_n - E_m} \right|^2$$

$$= \frac{\hbar^2 \lambda^2}{2m\omega} \frac{1}{\hbar\omega} [(-n+1) + n] = \frac{-\lambda^2}{2m\omega^2}.$$

$$\underline{|\psi_n^{(1)}\rangle} = \sum_{m \neq n} \frac{\langle m | H' | n \rangle}{E_n - E_m} |m\rangle$$

$$= \sqrt{\frac{\hbar}{2m\omega}} \lambda \left(\sum \frac{\sqrt{n+1} \delta_{m, n+1}}{E_n - E_m} |m\rangle + \sum \frac{\sqrt{n} \delta_{m, n-1}}{E_n - E_m} |m\rangle \right)$$

$$= \sqrt{\frac{\hbar}{2m\omega}} \lambda \left(\frac{\sqrt{n+1}}{-\hbar\omega} |n+1\rangle + \frac{\sqrt{n}}{\hbar\omega} |n-1\rangle \right)$$