

Superconductivity : It is the ability of certain materials to conduct electric current with practically zero resistance.

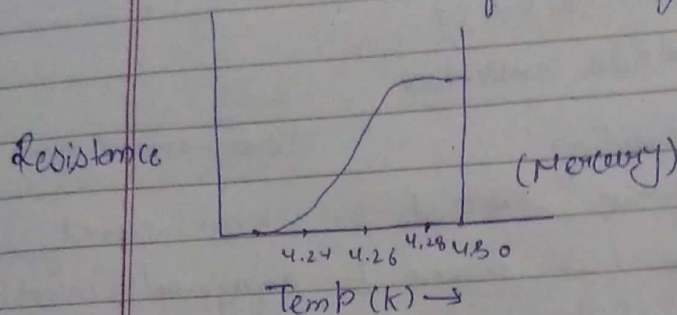
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CMP (Friday class)

Superconductivity

Properties: (1) zero resistance
(2) Perfect diamagnetic (have low (H₀) resistance)



$T_c \rightarrow$ Transition Temp

$T_c = 4.2\text{K}$

$\text{La-Ba-Cu-O} \rightarrow T_c = 34\text{K}$

* Meissner Effect : The expulsion of magnetic flux from a superconducting material when it is cooled below the transition temperature when it is in a magnetic field is called Meissner effect.

→ When we decrease the temp. below T_c then, magnetic flux away from the sample.

→ When we increase the temp. above the T_c then, magnetic flux is penetrated.

$B \rightarrow$ magnetic field,
 $H \rightarrow$ strength of M.F.
 $M \rightarrow$ Magnetization or intensity of magnetic field.

$$B = \mu_0 (H + M)$$

$$H = -M$$

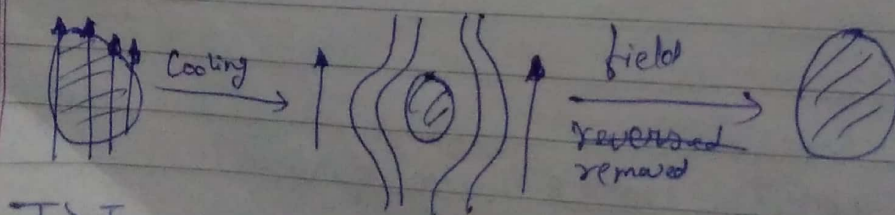
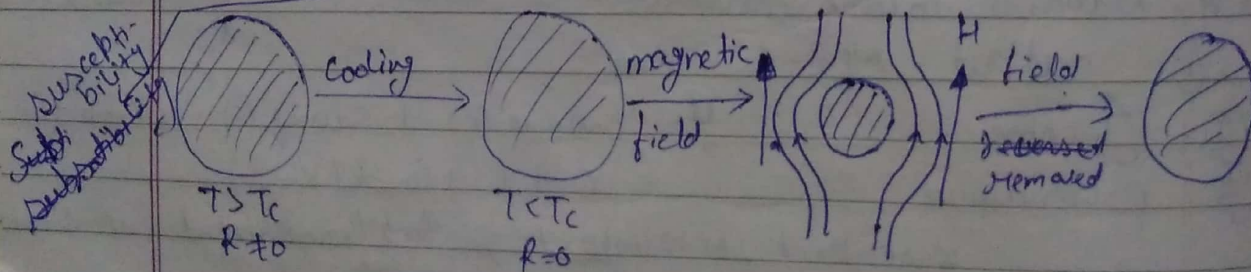
$$\chi = \frac{M}{H} = -1$$

A.C.T. Meissner effect,

($\therefore B = 0$)

so, $\mu_0 \neq 0$

$\therefore H + M = 0$



! Magnetic Behaviour of Superconductor.

$\rho = 0$ resistivity \rightarrow The value of magnetic field at which it will destroy superconductivity even temp is below T_c .

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\rightarrow At zero resistance \rightarrow when mag. flux varies \rightarrow then we calculated by Maxwell equation

$$\nabla \times E = -\frac{\partial B}{\partial t} \quad (\text{Maxwell equation})$$

$$\rho = 0$$

$$E = \oint \mathbf{J}$$

$\mathbf{J} \rightarrow$ current density

$$\therefore E = 0$$

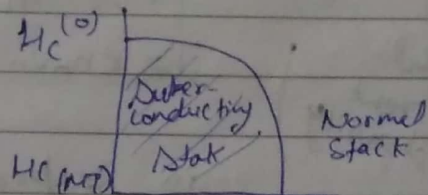
$$\therefore \frac{\partial B}{\partial t} = 0 \rightarrow B \rightarrow \text{constant}$$

* Critical Magnetic field (H_c): The tem. at which the magnetic field/resistance caused to zero or superconducting vanishes below the transition temp.

$H_c \rightarrow$ Max.

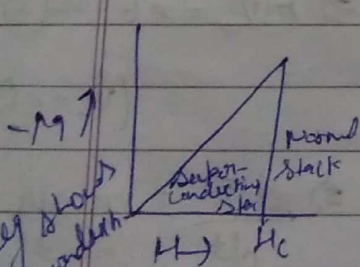
Critical Strength

$$H_c(T) = H_c(0) \left[1 - \left(\frac{T}{T_c} \right)^2 \right]$$

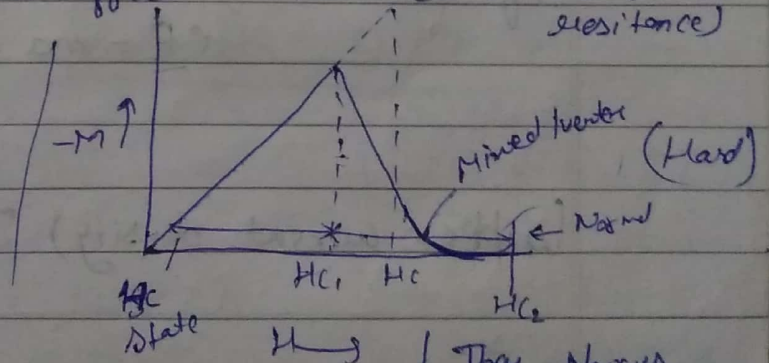


* Superconductor classified in 2 ways:

- ① Hard \rightarrow Completely Meissner effect \rightarrow soft
- ② Soft \rightarrow Partially Meissner effect \rightarrow Hard (have electrical resistance)



They show superconductivity upto $H_c(T_c)$ and after that they will show normal



- upto $H_{c1} \rightarrow$ completely Meissner effect show
- " H_{c1} to $H_{c2} \rightarrow$ Penetrating Meissner effect
- after $H_{c2} \rightarrow$ Normal State (Conductor)

They show superconductivity even after $H_c(T_c)$ at

Silsbee effect \rightarrow if H_c is less than 0.1 T, high magnetic field cannot be produced

- Types of Superconductor
- Maxwell's eq.
- Dynamics of Superconductor

At any tem. sum of superconductor and normal conductor is Total conductor

London Equations. (Statement → Book)

$n_n \rightarrow$ normal current density no. of normal e^- 's

$n_s \rightarrow$ Super " " no. of superconductor e^- 's

(Total no. of e^-) $n = n_n + n_s$ — (1) $v_n =$ velocity of normal e^-
 $J = J_n + J_s$ $v_s =$ velocity of Super electrons

(Total current density) $J = en_n v_n + en_s v_s$ — (2)

in Superconductor, there is no resistance,

if voltage is apply
 eq. of Motion for Superconductivity. e^- 's

$$eE = ma = m \left(\frac{dv_s}{dt} \right) \Rightarrow m \dot{v}_s = eE \quad (3)$$

$$(\dot{v}_s = \frac{eE}{m})$$

(Super current density) $J_s = en_s v_s \Rightarrow \frac{dJ_s}{dt} = n_s \cdot \frac{eE}{dt}$

$$\Rightarrow \frac{dJ_s}{dt} = \frac{n_s e^2 E}{m}$$

1st London eq.

$$\frac{dJ_s}{dt} = \frac{n_s e^2 E}{m}$$

— (A)

it express zero resistance property

Critical current (I_c) \rightarrow Minimum current that passes in a sample without destroying its superconductivity.

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When E is 0 $\therefore J_s$ is constant
i.e., after removal of electric field, there is a current

So it enforces the zero resistance property

Take both side curl in eq. (A)

$$\nabla \times \frac{\partial J_s}{\partial t} = -\frac{nse^2}{m} \nabla \times E$$

$$\nabla \times \frac{\partial J_s}{\partial t} = -\frac{nse^2}{m} \frac{\partial B}{\partial t}$$

using Maxwell eq

$$\nabla \times E = -\frac{\partial B}{\partial t}$$

integrating both sides

$$\nabla \times J_s = -\frac{nse^2}{m} B$$

$$B = \frac{-m}{nse^2} \text{curl } J_s$$

2nd eq.
London eq.

it explains

when $B=0 \therefore J_s$ is conservative

* Derivation for the Penetration Depth

$$\nabla \times B = \mu_0 J_s$$

Take curl both sides

$$\nabla \times \nabla \times B$$

$$= (\nabla \cdot B) \nabla - (\nabla \cdot \nabla) B$$

$$= 0 - \nabla^2 B$$

$$\nabla \times \nabla \times B = \mu_0 \nabla \times J_s$$

$$-\nabla^2 B = \mu_0 \nabla \times J_s$$

$$-\nabla^2 B = -\mu_0 \frac{nse^2}{m} B$$

$$\left(\begin{aligned} \nabla(\nabla \cdot B) - \nabla^2 B \\ = 0 - \nabla^2 B \end{aligned} \right) = -\nabla^2 B$$

in case of 1-D,

$$\nabla^2 B = \mu_0 \frac{nse^2}{m} B = \frac{1}{\lambda^2} B$$

formula for Penetration depth (λ)

$$\text{if in 1D } \frac{\partial^2 B}{\partial x^2} = \frac{1}{\lambda^2} B$$

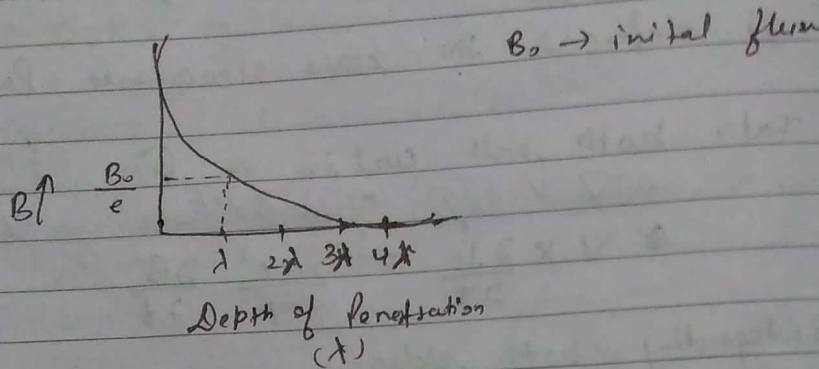
$$\therefore \lambda = \left(\frac{m}{\mu_0 nse^2} \right)^{\frac{1}{2}}$$

$\lambda \rightarrow$ Penetration Depth

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(Soln for 1-0)

$$B_z(x) = B_z(0) \exp\left(-\frac{x}{\lambda}\right)$$

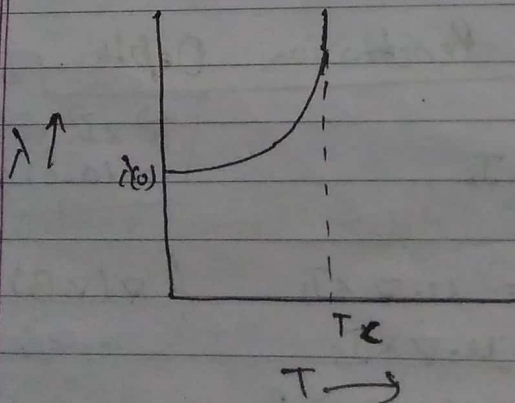


$$\text{curl } \vec{J}_s = -\frac{\vec{B}}{\mu_0 \lambda^2}$$

$$\vec{J}_s = \frac{\vec{E}}{\mu_0 \lambda^2}$$

How Penetration Depth varies with temperature.

* Effect of temp. on Penetration.



$T_c \rightarrow$ critical Temperature
 \rightarrow Penetration is infinite at T_c

The mo
to whic
of Pair
to Pseud
is call

$$\lambda(T) = \lambda_0 \left(\frac{1 - T^4}{T_c^4} \right)^{-\frac{1}{2}}$$

at $T = T_c$, ^{Normal} conductor is convert into superconductor.

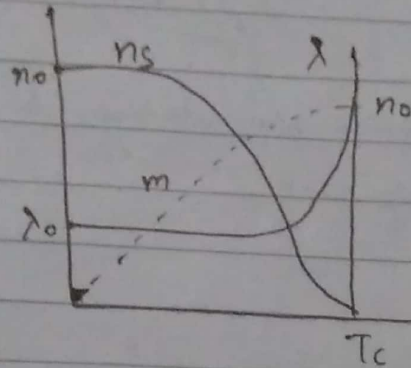
at low temp, it is independent of temp.

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When $\lambda \ll \lambda_0$

$$W = \frac{n_s}{n_0} = \left(1 - \frac{T^4}{T_c^4}\right)$$



$\therefore n_0$ is \propto Penetration depth increases rapidly and approaches infinity as the temp. approaches the T_c

When λ is below T_c
Normal λ_0

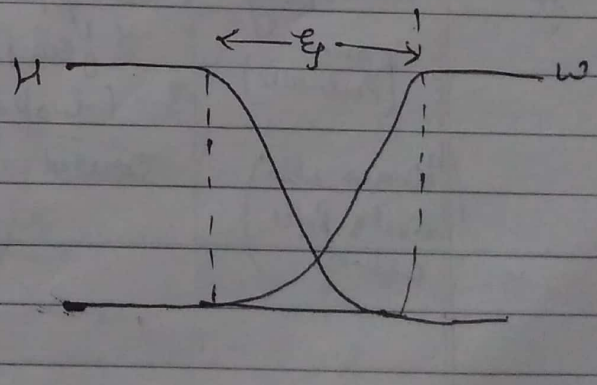
* Characteristic of Superconductor.
Related to energy gap &

VF & energy gap

Coherence length:

$$\xi_0 = \frac{2\hbar V_F}{\pi \Delta}$$

The max. distance up to which the states of pair e are correlated to produced superconductivity is called coherence length



It is a distance within which the order parameter is changes from its max. value in the bulk superconductor region to zero in the normal region. The distance is known as coherence length.

CMP (Friday class)

* Application of Superconductor

$T_c \rightarrow$ Transition temp.
High $T_c \rightarrow$ high Transition

(1) Stewart SQUIDS

(2) High T_c Superconductor

(3) Magnetic levitation

(4) Cryotron Application

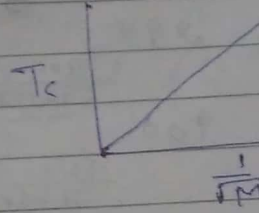
(1) SQUIDS \rightarrow Superconductor Quantum

Isotope effects:

$$T_c \propto M^{-1/2}$$

$$T_c \propto \frac{1}{\sqrt{M}}$$

$T_c M^{1/2} = \text{Constant}$, This is isotope effect

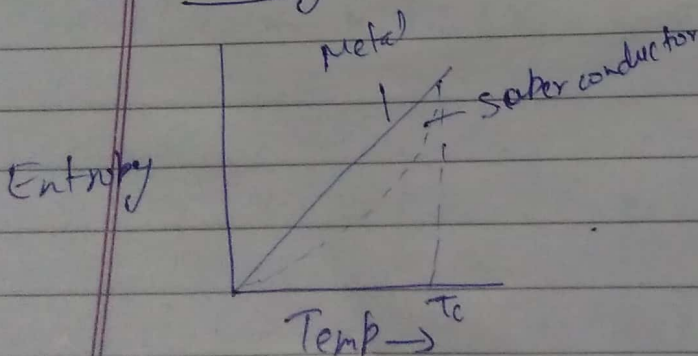


if there is a change in 199.5 to 203.4 mass then there is change in Transition temp. at 4.185K to 4.146.

* Thermal Property.

(1) Entropy (Randomness)

As temp. rise, e^- starts vibration at random motion. That's randomness is known as Entropy.



Entropy is decrease below Transition temperature
S.C has more order than Metal

Superconductivity affects \bar{e} 's mainly, so the lattice vibration part remains unaffected. That means the lattice part is same in both normal & superconducting states.

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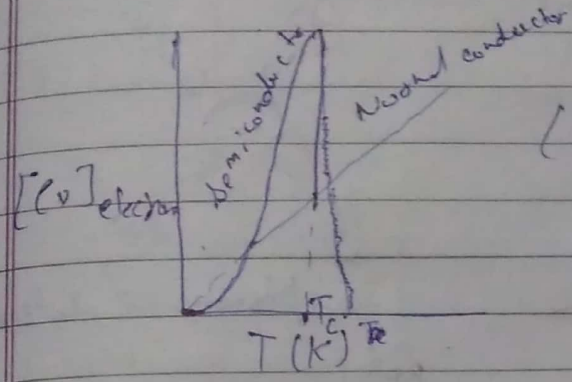
② Specific Heat : release from 2 source, 1st is lattice & 2nd is electronic

Shows at T_c Jump

$$C_N(T) = \alpha T + (BT^3) - \text{lattice contribution}$$

electronic contribution.

$\therefore BT^3$ is same in both normal & superconducting states.



$$(C_{el})_s = A \exp\left(-\frac{\Delta}{k_B T}\right)$$

$$\Delta = b k_B T_c, \quad b \rightarrow \text{constant}$$

$$C = C_{\text{lattice}} + C_{\text{electron}}$$

(Normal electron) $C_N = (C_{\text{lattice}})_N + (C_{el})_N$

el \rightarrow electron

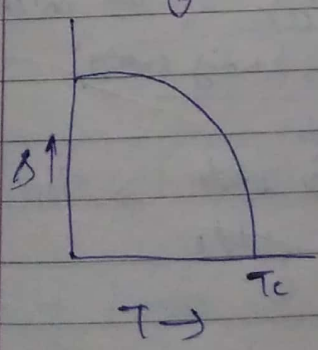
~~S-S~~ $C_S = (C_{\text{lattice}})_S + (C_{el})_S$

S-S specific heat

$$(C_{\text{lattice}})_S = (C_{\text{lattice}})_N$$

$$C_N - C_S = (C_{el})_N - (C_{el})_S$$

③ Energy Gap



In superconductor, Energy gap vary with temp.

When $E-G=0$, supercond. convert into normal conductor.

E.G. exists at fermi gas

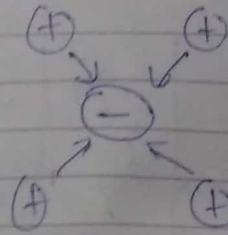
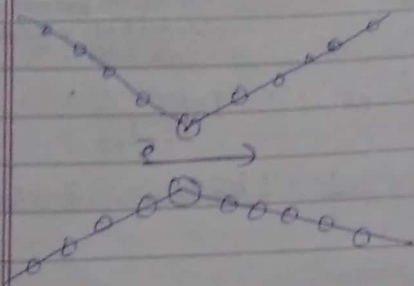
E.G. has no affect on special \bar{e} , that carry current in a superconductor. superconductor contains normal \bar{e} 's as well and these \bar{e} 's are affected by E.G.

✓ Mermin based on isotop eq.
Penetration from magnetic field
London eq.

* BCS Theory: (Bardeen, Cooper, Schrieffer) - It divides into 3 parts

① Electron-phonon and electron interaction.

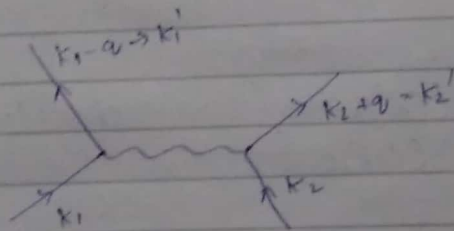
phonon: Quantize vibrated energy



phonon - e^- interaction

Book:

isotope effect: mass affects the



$$k_1 + k_2 = k_1' + k_2'$$

Phono energy dominates the e^- energy, e^- 's are interacting

② Cooper Pair \rightarrow (certain cooper pair & binding energy)
energy of e^- in pair state is less than
the energy of e^- in free state

Think \rightarrow If e^- doesn't make pair then there is no energy gap

Energy Binding
Bond energy

Energy required to separate to break the cooper pair

phonon pair
energy is
low

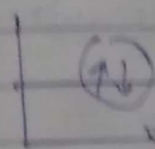
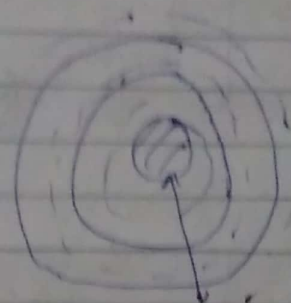
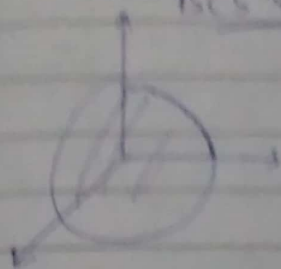
$$E_b = \frac{h v_f}{2}$$

$E_b \rightarrow$ Binding energy
 $v_f \rightarrow$ Fermi velocity

on any level thermal energy going to next state, there is no energy.

RCS Grand State

Superconductor

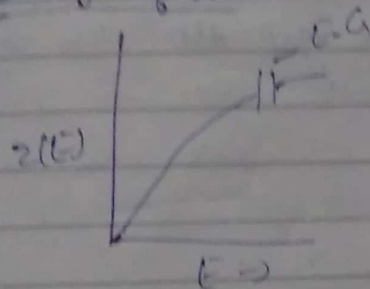


$k(\downarrow) \quad k(\uparrow)$

if any state is

Below the Fermi surface all the states are occupied by e^- 's

density of state



* To flux quantization

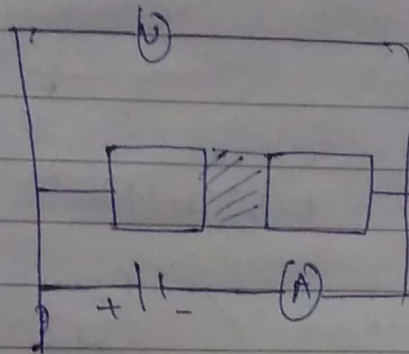
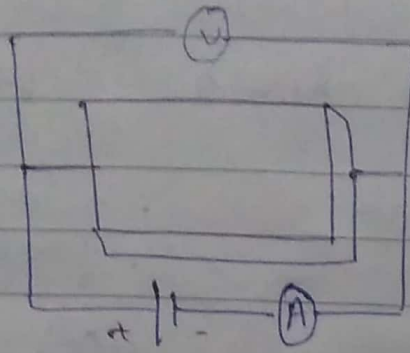
flux passing through $\Phi = n \Phi_0 = \frac{n h}{2e}$

flux is quantization

$$\Phi_0 = \frac{h}{2e} = 2.068 \times 10^{-15} \text{ Weber}$$

unit of flux is fluxoid.

* Josephson Tunneling.



Current is not flow thro the superconductor

if there is no voltage, then DC is flowing thro the superconductors

If current flows from the gap then this is AC

$$f = \frac{2eV}{h}$$

→ DC Josephson effect

A DC flows across the junction of 2 superconductors separated by thin insulator wire in the absence of any magnetic field called DC Josephson effect

→ A DC voltage applied across the junction of 2 superconductors separated by a thin insulator wire causes RF