

Chapter 5 (A.P.French)

Following is the list of problems of chapter 5 of A P French

Lecture : 9, 10, 12, 16; Tutorial : 7, 11, 13; Suggested : 2, 4, 6, 8, 14, 15, 17

5.7. Eq. of motion of A : $m\ddot{x}_A = -k_0x_A - k_c x_A + k_c x_B = -k_0x_A - k_c(x_A - x_B)$

$$\Rightarrow \ddot{x}_A + \left(\frac{k_0}{m} + \frac{k_c}{m}\right)x_A - \frac{k_c}{m}x_B = 0; \text{ Let : } \omega_0 = \sqrt{\frac{k_0}{m}} \text{ and } \omega_c = \sqrt{\frac{k_c}{m}}$$

$$\text{So, } \boxed{\ddot{x}_A + (\omega_0^2 + \omega_c^2)x_A - \omega_c^2 x_B = 0} \text{-----(1)}$$

$$\text{Similarly : } \boxed{\ddot{x}_B + (\omega_0^2 + \omega_c^2)x_B - \omega_c^2 x_A = 0} \text{-----(2)}$$

$$(1) + (2) \Rightarrow (\ddot{x}_A + \ddot{x}_B) + (\omega_0^2 + \omega_c^2)(x_A + x_B) - \omega_c^2(x_A + x_B) \Rightarrow (\ddot{x}_A + \ddot{x}_B) + \omega_0^2(x_A + x_B) = 0$$

$$\Rightarrow \ddot{q}_1 + \omega_0^2 q_1 = 0; \text{ where, } q_1 = x_A + x_B \text{-----(3)}$$

$$(1) - (2) \Rightarrow \boxed{(\ddot{x}_A - \ddot{x}_B) + (\omega_0^2 + 2\omega_c^2)(x_A - x_B) = 0 \Rightarrow \ddot{q}_2 + (\omega_0^2 + 2\omega_c^2)q_2 = 0} : \text{ where, } q_2 = x_A - x_B \text{-----(4)}$$

$$\therefore \text{Two normal frequencies are : } \boxed{\omega_1 = \omega_0 \text{ and } \omega_2 = \sqrt{\omega_0^2 + 2\omega_c^2}}$$

If B is clamped : $x_B = 0$; From eq.(1) : Angular frequency of A = $\omega_A = \sqrt{\omega_0^2 + \omega_c^2}$

$$\therefore \nu_A = \frac{\omega_A}{2\pi} = 1.181 \text{ sec}^{-1} \text{ and } \nu_1 = \frac{\omega_1}{2\pi} = 1.14 \text{ sec}^{-1} = \nu_0$$

$$\therefore \omega_0 = 2\pi\nu_0; \text{ Now : } \omega_A = \sqrt{\omega_0^2 + \omega_c^2} \Rightarrow \omega_c = \sqrt{\omega_A^2 - \omega_0^2}$$

$$\Rightarrow \omega_c^2 = 4\pi^2(\nu_A^2 - \nu_0^2); \text{ Similarly : } \omega_2 = \sqrt{\omega_0^2 + 2\omega_c^2} \Rightarrow 2\pi\nu_2 = \sqrt{\omega_0^2 + 4\pi^2(\nu_A^2 - \nu_0^2)} \Rightarrow 2.29 \text{ sec}^{-1}$$

$$(d) \omega_A = \sqrt{\omega_0^2 + \omega_c^2} \Rightarrow \omega_c^2 = \omega_A^2 - \omega_0^2 = \frac{k_c}{m} \text{ and } \omega_0 = \frac{k_0}{m}$$

$$\therefore \frac{\omega_c^2}{\omega_0^2} = \frac{k_c}{k_0} = \frac{\omega_A^2 - \omega_0^2}{\omega_0^2} = \frac{\nu_A^2 - \nu_0^2}{\nu_0^2} = 1.52$$

5.8.(a) When the force F is applied at a distance ' a ' to hold the pendulum at an angle θ , the total

torque acting on the pendulum must be zero. So, $Fa \cos \theta = mgL \sin \theta \Rightarrow \tan \theta \approx \theta = \frac{Fa}{mgL}$.

When, the force F' applied at the position of the bob to hold the same angle, then $a = L$

$$\therefore \theta = \frac{F'a}{mgL} = \frac{F'L}{mgL} = \frac{F'}{mg} \therefore \frac{F'}{mg} = \frac{Fa}{mgL} \Rightarrow \boxed{F' = \frac{Fa}{L}}$$

(b) If the angular displacements of the two pendulums are θ_1 and θ_2 , the equation of motion of the 1st pendulum is :

$I\ddot{\theta}_1 = -mgL \sin \theta_1 + (ka \sin \theta_2 - ka \sin \theta_1)a \cos(\theta_1 - \theta_2) \approx -mgL\theta_1 + ka(\theta_2 - \theta_1)a \cos(\theta_1 - \theta_2)$ assuming $\sin \theta_1 \approx \theta_1$ and $\sin \theta_2 \approx \theta_2$ for small values of the angles.

Here, $-mgL \sin \theta_1$ is the torque due to gravity. $(ka \sin \theta_2 - ka \sin \theta_1)a \cos \theta$ is the torque due to the resultant compression $(a(\theta_2 - \theta_1))$ of the spring.

$\Rightarrow I\ddot{\theta}_1 + mgL\theta_1 - ka(\theta_2 - \theta_1)a \Rightarrow mL^2\ddot{\theta}_1 + mgl\theta_1 + ka^2(\theta_1 - \theta_2) = 0$ for small displacements.

$$\Rightarrow \ddot{\theta}_1 + \frac{g}{L}\theta_1 + \frac{ka^2}{mL^2}(\theta_1 - \theta_2) = 0 \Rightarrow \boxed{\ddot{\theta}_1 + \left(\frac{g}{L} + \frac{ka^2}{mL^2}\right)\theta_1 - \frac{ka^2}{mL^2}\theta_2 = 0} \text{-----(1)}$$

$$\text{similarly : } \boxed{\ddot{\theta}_2 + \left(\frac{g}{L} + \frac{ka^2}{mL^2}\right)\theta_2 - \frac{ka^2}{mL^2}\theta_1 = 0} \text{-----(2)}$$

In terms of x : $x_1 = a \sin \theta_1 \approx a\theta_1$ and $x_2 \approx a\theta_2$.

$$\therefore \text{Eq. becomes : } \boxed{\ddot{x}_1 + \left(\frac{g}{L} + \frac{ka^2}{mL^2}\right)x_1 - \frac{ka^2}{mL^2}x_2 = 0} \text{ and } \boxed{\ddot{x}_2 + \left(\frac{g}{L} + \frac{ka^2}{mL^2}\right)x_2 - \frac{ka^2}{mL^2}x_1 = 0}$$

After : (1) + (2) and (1) - (2) we have :

$$\boxed{\frac{d^2}{dt^2}(\theta_1 + \theta_2) + \left(\frac{g}{L} + \frac{ka^2}{mL^2}\right)(\theta_1 + \theta_2) - \frac{ka^2}{mL^2}(\theta_1 + \theta_2) = 0 \Rightarrow \ddot{q}_1 + \frac{g}{L}q_1 = 0} \text{-----(3); where } q_1 = \theta_1 + \theta_2 \text{ or } x_1 + x_2$$

$$\text{and } \boxed{\ddot{q}_2 + \left(\frac{g}{L} + \frac{2ka^2}{mL^2}\right)q_2 = 0} \text{-----(4); where, } q_2 = \theta_1 - \theta_2 \text{ or } x_1 - x_2$$

$$\therefore \boxed{\omega_1 = \sqrt{\frac{g}{L}} \text{ and } \omega_2 = \sqrt{\frac{g}{L} + \frac{2ka^2}{mL^2}}}$$

5.11.

$$(a) m_1 \ddot{x}_1 = -kx_1 + T \sin \theta \Rightarrow -kx_1 + m_2 g \cos \theta \sin \theta \text{ where : } T = m_2 g \cos \theta$$

$$\text{For small angles : } \cos \theta = 1; \text{ Here, } \sin \theta = \frac{x_2 - x_1}{l} \approx \theta \Rightarrow m_1 \ddot{x}_1 = -kx_1 + m_2 g \frac{x_2 - x_1}{l}$$

$$\text{Similarly, for } m_2 : m_2 \ddot{x}_2 = -T \sin \theta = -m_2 g \frac{x_2 - x_1}{l}$$

$$(b) \text{ If } m_1 = m_2 = m,$$

$$m_1 \ddot{x}_1 = -kx_1 + m_2 g \frac{x_2 - x_1}{l} \Rightarrow \ddot{x}_1 + \left(\frac{k}{m} + \frac{g}{l} \right) x_1 - \frac{g}{l} x_2 = 0 \text{ --- (1)}$$

$$\text{Similarly } m_2 \ddot{x}_2 = -m_2 g \frac{x_2 - x_1}{l} \Rightarrow \ddot{x}_2 + \frac{g}{l} x_2 - \frac{g}{l} x_1 = 0 \text{ --- (2)}$$

Let ω is the normal mode frequency and

$$x_1(t) = A \cos(\omega t) = A e^{i\omega t} \text{ and } x_2 = B e^{i\omega t}; \therefore \ddot{x}_1 = -\omega^2 A e^{i\omega t} \text{ and } \ddot{x}_2 = -\omega^2 B e^{i\omega t}$$

Putting these values in eq.(1) and (2) we have :

$$-\omega^2 A e^{i\omega t} + \left(\frac{k}{m} + \frac{g}{l} \right) A e^{i\omega t} - \frac{g}{l} B e^{i\omega t} = 0 \Rightarrow \left(\frac{k}{m} + \frac{g}{l} - \omega^2 \right) A - \frac{g}{l} B = 0 \text{ ---- (3)}$$

$$\text{And : } -\omega^2 B e^{i\omega t} + \frac{g}{l} B e^{i\omega t} - \frac{g}{l} A e^{i\omega t} = 0 \Rightarrow -\frac{g}{l} A + \left(\frac{g}{l} - \omega^2 \right) B = 0 \text{ ---- (4)}$$

These equations only will have a non-trivial solution if :

$$\begin{vmatrix} \frac{k}{m} + \frac{g}{l} - \omega^2 & -\frac{g}{l} \\ -\frac{g}{l} & \frac{g}{l} - \omega^2 \end{vmatrix} = 0 \Rightarrow \omega^4 - \left(\frac{k}{m} + \frac{2g}{l} \right) \omega^2 + \frac{gk}{ml} = 0$$

$$\therefore \omega^2 = \frac{\left(\frac{2g}{l} + \frac{k}{m} \right) \pm \sqrt{\left(\frac{2g}{l} + \frac{k}{m} \right)^2 - \frac{4gk}{ml}}}{2} = \frac{\left(\frac{2g}{l} + \frac{k}{m} \right) \pm \sqrt{\frac{4g^2}{l^2} + \frac{k^2}{m^2}}}{2}$$

$$(c) \frac{g}{l} \gg \frac{k}{m} \Rightarrow \omega^2 = \frac{\left(\frac{2g}{l} + \frac{k}{m} \right) \pm \sqrt{\frac{4g^2}{l^2} + \frac{k^2}{m^2}}}{2} \approx \frac{\left(\frac{2g}{l} + \frac{k}{m} \right) \pm \frac{2g}{l}}{2} = \left(\frac{2g}{l} + \frac{k}{2m} \right) \text{ and } \frac{k}{2m}$$

$$\omega = \pm \sqrt{\left(\frac{2g}{l} + \frac{k}{2m} \right)} \text{ and } \pm \sqrt{\frac{k}{2m}}; \text{ '-' frequency has no physical meaning.}$$

$$\text{Therefore acceptable frequencies are : } \sqrt{\left(\frac{2g}{l} + \frac{k}{2m} \right)} \text{ and } \sqrt{\frac{k}{2m}}$$

5.13. This type of problems are known as beaded string. For one and two beaded string the solutions are as follows:

(a) Single beaded string: Let the horizontal displacement is x . So, eq. of motion of the mass :

$$m\ddot{x} = -T(\sin\alpha_1 + \sin\alpha_2) \approx -T\left(\frac{x}{l} + \frac{x}{2l}\right) = -\frac{3T}{2l}x$$

$$\Rightarrow \ddot{x} + \frac{3T}{2ml}x = 0 \Rightarrow \omega = \sqrt{\frac{3T}{2ml}}$$

(b) Two beaded string : Eq. of motion of the 1st mass with displacement x_1 : $m\ddot{x}_1 = -T(\sin\alpha_1 + \sin\alpha_2)$;

$$\sin\alpha_1 = \frac{x_1}{l} \text{ and } \sin\alpha_2 = \frac{x_1 - x_2}{l}$$

$$\therefore m\ddot{x}_1 = -T\left(\frac{2x_1 + x_2}{l}\right) \Rightarrow \ddot{x}_1 + \frac{T}{ml}(2x_1 + x_2) = 0 \text{ --- (1)}$$

Similarly, for 2nd mass the if the displacement is x_2 , the the equation of motion is:

$$\ddot{x}_2 + \frac{T}{ml}(2x_2 + x_1) = 0 \text{ --- (2) [Change the suffix only]}$$

$$(1) + (2) \Rightarrow (\ddot{x}_1 + \ddot{x}_2) + \frac{T}{ml}(x_1 + x_2) = 0 \Rightarrow \omega_1 = \sqrt{\frac{T}{ml}}$$

and

$$(1) - (2) \Rightarrow (\ddot{x}_1 - \ddot{x}_2) + \frac{3T}{ml}(x_1 - x_2) = 0 \Rightarrow \omega_2 = \sqrt{\frac{3T}{ml}} \text{ Higher mode}$$