Ch. 8: Non Inertial Systems and Fictitious Forces

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Principle of Equivalence

There is no way to distinguish locally between a uniform gravitational acceleration \vec{g} and an acceleration of the coordinate system $\vec{A} = -\vec{g}$. Locally here means a sufficiently small region like that of an elevator) where you can assume \vec{g} to be practically constant.

For instance if your elevator is in a state of free fall under earth's gravity and you drop an apple, then the apple will float in front of $you(m_i = m_g)$

$$ec{F'} = \underbrace{-m_g ec{g}}_{real\ force} + \underbrace{m_i ec{g}}_{Fictitious\ force} = 0$$



Principle of Equivalence

That means inside of a sufficiently giant elevator since \vec{g} would vary from point to point, I can see departures from this equivalence as cancellation won't be exact.



Earth as a giant elevator

Earth is in a state of free-fall towards Sun.



$$\vec{G}_0 = G M_s rac{\hat{n}}{r_s^2}$$

If $\vec{G}(\vec{r})$ is the gravitational field of Sun at some location \vec{r} on earth. Then,

$$\vec{F} = m \, \vec{G}(\vec{r})$$

The apparent force to an earthbound observer is

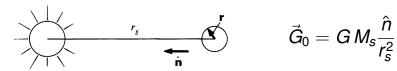
$$\vec{F}' = F - m\vec{A} = m[\vec{G}(\vec{r}) - \vec{G}_0] = m\vec{G}'(\vec{r})$$





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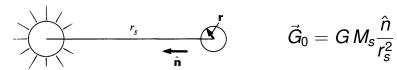
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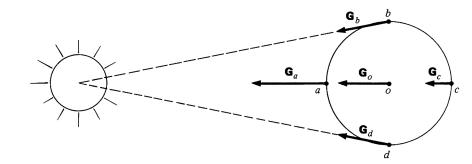
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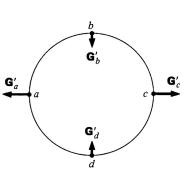


Sun's gravitational field $\vec{G}(\vec{r})$ at different points on Earth





Apparent field $\vec{G}'(\vec{r})$ at different points on earth's surface



$$G_a = rac{G \, M_s}{(r_s - R_e)^2}$$

Calculation of G'_a and G'_c :

$$G'_{a} = G_{a} - G_{0}$$

$$= \frac{G M_{s}}{(r_{s} - R_{e})^{2}} - \frac{G M_{s}}{r_{s}^{2}}$$

$$= \frac{G M_{s}}{r_{s}^{2}} \left[\frac{1}{[1 - (R_{e}/r_{s})^{2}]} - 1 \right]$$





Apparent field $G'(\vec{r})$ at different points on earth's surface

Since
$$\frac{R_e}{r_s} = \frac{6.4 \times 10^3 \text{ km}}{1.5 \times 10^8 \text{ km}} = 4.3 \times 10^{-3} << 1$$
, we have

$$G'_{a} = G_{0} \left[\left(1 - \frac{R_{e}}{r_{s}} \right)^{-2} - 1 \right]$$

$$= G_{0} \left[1 + 2 \frac{R_{e}}{r_{s}} + \dots - 1 \right]$$

$$= 2 G_{0} \frac{R_{e}}{r_{s}}$$

where we have neglected terms of order $(R_e/r_s)^2$ and higher.

Apparent field $\vec{G}'(\vec{r})$ at different points on earth's surface

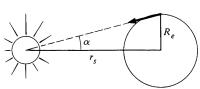
Similarly for point c, the distance from the Sun changes to $r_s + R_e$. Hence,

$$G_c'=-2\,G_0rac{R_e}{r_s}$$



Apparent field $G'(\vec{r})$ at different points on earth's surface

Calculation of G'_h and G'_d :



Points b and d are approximately same distance from the Sun as the center of the earth. However, $\alpha \approx R_e/r_s = 4.3 \times 10^{-5} << 1.$

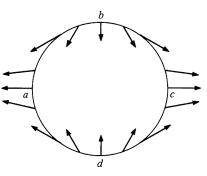


$$egin{array}{ll} G_b' &pprox & G_0 \, lpha \ &pprox & G_0 rac{R_e}{r_s} \end{array}$$

 $G_b' \approx G_0 \alpha$ $\approx G_0 \frac{R_e}{r_s}$ By Symmetry G_d' is equal but opposite to G_b' .



Apparent field $\vec{G}'(\vec{r})$ at different points on earth's surface



Tidal Forces

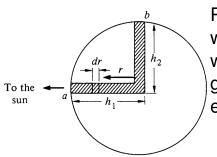
Forces at *a* and *c* tend to lift the ocean

Forces at *b* and *d* tend to depress them.

Can you see why there are two tides in a day?







Two orthogonal wells

Pressure due to short column of water of height dr is $\rho g(r)dr$ where g(r) is the effective gravitational field at r. For equilibrium:

$$\int_0^{h_1} \rho \, g_1(r) dr = \int_0^{h_2} \rho \, g_2(r) dr$$

The idea is to calculate $\Delta h = h_1 - h_2$, the height of tide due to Sun.





Effective field toward the center of earth in column 1:

$$g_1(r) = \underbrace{g(r)}_{\textit{Earth's gravitational field}} - \underbrace{G_1'(r)}_{\textit{Apparent field of Sun}}$$

Borrowing from G'_a (substitute r for R_e):

$$G'_1(r) = \frac{2 G M_s r}{r_s^3}$$
$$= 2 C r$$

where
$$C = G M_s / r_s^3$$
.



Thus,

$$g_1(r) = g(r) - 2 C r$$

 $g_2(r) = g(r) + G'_2(r)$
 $= g(r) + C r$

Condition for equilibrium is:

$$\int_{0}^{h_{1}} [g(r) - 2cr] dr = \int_{0}^{h_{2}} [g(r) + Cr] dr$$

$$\int_{0}^{h_{1}} g(r) dr - \int_{0}^{h_{2}} g(r) dr = \int_{0}^{h_{1}} 2Cr dr + \int_{0}^{h_{2}} Cr dr$$

$$\int_{h_{2}}^{h_{1}} g(r) dr = \int_{0}^{h_{1}} 2Cr dr + \int_{0}^{h_{2}} Cr dr$$

Since $h_1 \approx h_2 \approx R_e$, $g(r) \approx g(R_e) = g$ above equation reduces to

$$g\Delta h_s = \frac{3}{2}CR_e^2$$

$$\Delta h_s = rac{3}{2}rac{M_s}{M_e}\left(rac{R_e}{r_s}
ight)^3 R_e \quad \left[g = rac{GM_s}{R_e^2} \quad C = rac{GM_s}{r_s^3}
ight]$$

Using the data

$$M_s = 1.99 imes 10^{33} \, {
m g} \qquad r_s = 1.49 imes 10^{13} \, {
m cm}$$

$$M_e = 5.98 \times 10^{27} \, \mathrm{g}$$
 $R_e = 6.37 \times 10^8 \, \mathrm{cm}$,

we obtain

$$\Delta h_s = 24.0$$
 cm





An identical calculation for moon yields:

$$\Delta h_m = \frac{3}{2} \frac{M_m}{M_e} \left(\frac{R_e}{r_m}\right)^3 R_e$$
$$= 53.5 cm$$

Since $\Delta h \to 1/r^2$, the distance factor more than kills whatever advantage Sun has due to its mass. Strongest tides (spring tides) occur when moon and Sun act along the same line. Weak (neap tides) occur midway between, at the quarters of the moon.

$$rac{\Delta h_{spring}}{\Delta h_{neap}} = rac{\Delta h_m + \Delta h_s}{\Delta h_m - \Delta h_s} pprox 3$$



