

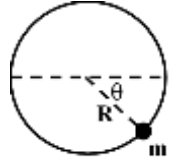
BIRLA INSTITUTE OF TECHNOLOGY AND SCIENCE PILANI, RAJASTHAN

Comprehensive Examination (Closed-Book): 2017-18, 2nd Semester

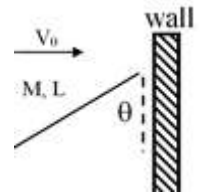
Mechanics Oscillations and Waves (MEOW): PHY F111, 5th May 2018, Duration: 3 hrs., Full Marks: 105

Instruction(s): All questions are compulsory. Answer all parts of a question together.
Write your final answer of each sub-part inside a box.

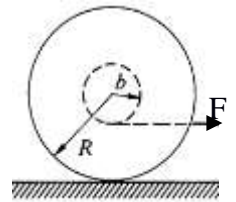
Q.1(A) A bead of mass 'm' is at rest at the top of a fixed frictionless hoop of radius R that lies in a vertical plane. The bead is given infinitesimal push so that it slides down and around the hoop in clockwise direction. Find all points on the hoop (angular position as θ w.r.t. horizontal) where the bead's acceleration is horizontal. (10)



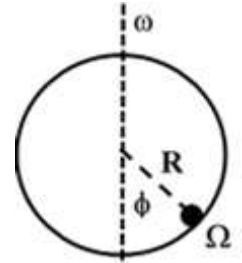
Q.1(B) On a frictionless floor (x-y plane) a uniform stick of mass M and length L moves with a linear speed V_0 (without rotation) in the direction perpendicular to a wall. It makes an angle θ with the wall as shown. If the stick collide elastically with a rigidly fixed wall, what should $\cos\theta$ be so that the speed of the center of mass of the stick immediately after the collision is zero? (10)



Q.2(A) A Yo-Yo of mass M has an axle of radius b and a spool of radius R. Its moment of inertia can be taken to be $MR^2/2$. The Yo-Yo is placed upright on a table and the string is pulled with a horizontal force F as shown. The coefficient of friction between the Yo-Yo and the table is μ . What is the maximum value of F in terms of given parameters for which the Yo-Yo will roll without slipping? (7)



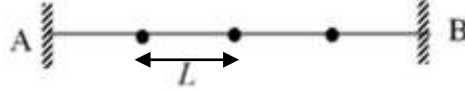
Q.2(B) Consider a thin hoop of radius R, spinning with constant angular speed ω about its diameter. A small bug of mass m walks with constant angular speed Ω on its inside surface as shown in figure. Let \vec{F} be the total force the hoop applies on the bug when it is making an angle ϕ with the vertical. Taking the plane of hoop to coincide with the plane of paper at this instant as shown in figure, answer the following questions. (Neglect gravity).



- (a) Find (\vec{F}_r) and (\vec{F}_θ) (with directions) in terms of m, R, ϕ , ω and Ω . [Hint. Resolve \vec{F} into components towards the axis of rotation (\vec{F}_r) and perpendicular to the plane of hoop (\vec{F}_θ)] write down the differential equation of motion for bug using polar coordinates. (6)
- (b) Using the result of (a) find the torque applied by the hoop on the bug about the axis of rotation. (2)
- (c) Find the $d\vec{L}/dt$ about the axis of rotation. (2)
- (d) Now analyze the problem from the rotating frame of hoop and answer the following questions.
- (i) What is the net acceleration of the bug in the rotating frame of hoop? (4)
- (ii) What are the magnitudes and describe the directions of fictitious forces. (4)

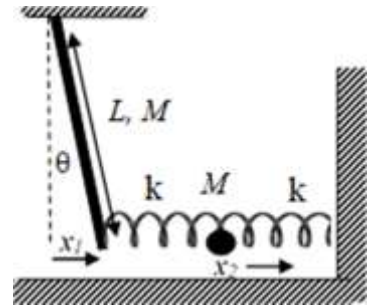
Q.3(A) The phase velocity of waves is given by $V_{ph} = \sqrt{(a/k) + (k/b)}$, where a and b are constants and k is the wave number. **(a)** Calculate the group velocity V_g of the waves. **(b)** Find the value of k at which phase and group velocities are equal. **(3+2)**

Q.3(B) An elastic string of negligible mass, stretched so as to have a tension T , is attached to fixed points A and B, a distance $4L$ apart, and carries three equally spaced particles of mass M , as shown in figure below.

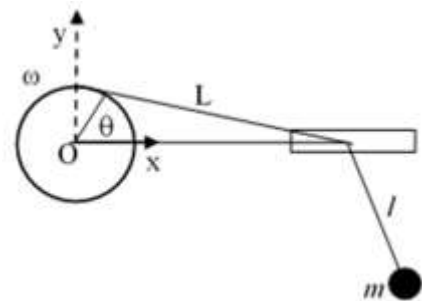


(a) Suppose that the particles have small transverse displacements y_1, y_2 and y_3 respectively, at some instant. Derive the differential equations of motion for each particle. **(b)** The appearance of the normal modes can be found by the sine curves that pass through A and B. Find the relative values and sign of the amplitudes A_1, A_2 and A_3 in each of the possible modes of the system. **(c)** Now take $T = 10 \text{ N}$, $M = 4 \text{ grams}$ and $4L = 1 \text{ m}$. Find the angular frequencies of all the possible modes by putting $y_1 = A_1 \sin \omega t$, $y_2 = A_2 \sin \omega t$ and $y_3 = A_3 \sin \omega t$ in the equations obtained in part **(a)** and using the ratios $A_1 : A_2 : A_3$ obtained in part **(b)**. **(3+3+4)**

Q.3(C) W.r.t. the figure, in a coupled oscillator system, a physical pendulum of length L and mass M pivoted at the upper end, and its lower end is coupled to a mass-spring system. A particle of mass M is at the junction of two identical springs of spring constant k . The springs are massless. Using the approximation of small oscillations. **(a)** Write down the total energy of the system at some instant when angular displacement of physical pendulum is θ , linear displacement of lowest end of the pendulum is x_1 and linear displacement of the particle is x_2 . **(b)** Derive the differential equations of motion for x_1 and x_2 . **(c)** If $g/L = 2k/M$, determine the angular frequencies of the normal modes of the system in terms of g and L . Neglect damping in this problem. **(4+5+6)**.



Q.4(A) The point of suspension of a simple pendulum of mass m length l is driven (strictly horizontally) by a motor which consists of a disk of radius R rotating with an angular speed ω as shown. The point of suspension of the pendulum is connected to the motor-wheel by a rod of negligible mass of length L ($L \gg R$) and also it can move inside a fixed frictionless groove. In free vibration (when the motor is off) the amplitude of the pendulum drops by a factor of 'e' in 50 swings. Assume the whole system is in x - y plane. **(a)** Write down the equation of motion of the pendulum for forced oscillation choosing the origin of your axis at the center of the disc. **(b)** Find out $A(\omega)$ and $\delta(\omega)$ from **(a)**. **(c)** Calculate the maximum mean power absorbed by the forced oscillator for the given data: $m=0.1 \text{ kg}$, $l=0.5 \text{ m}$, $R=0.01 \text{ m}$ (Take, $g=10 \text{ m/s}^2$). **(8+8+4)**



Q.4(B). In the forced oscillation, the power input to maintain forced vibration can be calculated by recognizing that this power is the mean rate of doing work against the resistive force $-b\dot{u}$. **(a)** Find out the instantaneous rate

of doing work against the resistive force. **(b)** Find out the mean rate of doing work in the forced oscillation assuming $x = A \cos(\omega t - \delta)$. **(4+6)**

End

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MODEL SOLUTIONS

Q.1(A). Let at an angle θ below w.r.t horizontal diameter the acceleration is purely horizontal.

While moving along the circle, the bead has two acceleration always, tangential acceleration (a_θ) and radial acceleration (a_r). At a position θ below horizontal:

From energy conservation: $\frac{1}{2}mv^2 = mg(R + R \sin \theta) \Rightarrow \boxed{v = \sqrt{2gR(1 + \sin \theta)}} \text{-----(3)}$

\therefore Radial acceleration: $a_r = \frac{v^2}{R} = 2g(1 + \sin \theta) \text{-----(1)}$

and tangential acceleration: $a_\theta = g \cos \theta \text{-----(1)}$

Condition for horizontal acceleration: Vertical components of a_r and a_θ must be equal and opposite.

i.e. $a_r \sin \theta = a_\theta \cos \theta \Rightarrow 2g(1 + \sin \theta) \sin \theta = g \cos \theta \cos \theta \Rightarrow 3 \sin^2 \theta + 2 \sin \theta - 1 = 0$

$\sin \theta = \frac{-2 \pm \sqrt{16}}{6} = \frac{-2 \pm 4}{6} = \frac{1}{3} \text{ or } -1$

$\therefore \sin \theta = \frac{1}{3}$ is only the physically acceptable solution. i.e. $\theta = 19.47^\circ$ or 167.53° -----(5)

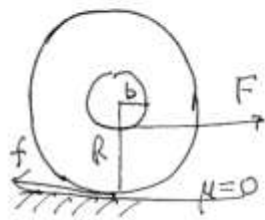
Q.1(B). Method 1: W.r.t. the collision point, $L_i = L_f \Rightarrow MV_0 \frac{L}{2} \cos \theta = I\omega = \frac{ML^2}{12} \omega \Rightarrow \boxed{\omega = \frac{6V_0 \cos \theta}{L}} \text{-----(5)}$

From energy conservation: $\frac{1}{2}MV_0^2 = \frac{1}{2}I\omega^2 = \frac{1}{2} \cdot \frac{ML^2}{12} \cdot \left(\frac{6V_0 \cos \theta}{L} \right)^2 \Rightarrow \boxed{\cos \theta = \frac{1}{\sqrt{3}}} \text{-----(5)}$

Method 2: $\frac{1}{2}MV_0^2 = \frac{1}{2}I\omega^2 \Rightarrow MV_0^2 = \frac{ML^2}{12} \omega^2 \Rightarrow \boxed{\omega = \frac{2\sqrt{3}V_0}{L}} \text{-----(3)}$

Now, $L_i = L_f \Rightarrow MV_0 \frac{L}{2} \cos \theta = I\omega = \frac{ML^2}{12} \cdot \frac{2\sqrt{3}V_0}{L} \Rightarrow \boxed{\cos \theta = \frac{1}{\sqrt{3}}} \text{-----(8)}$

Q-2 A



Translational
equation of motion

$$F - f = ma \quad (1)$$

It is clear that
Yo-Yo will translate
to the right. For
it to roll without

slipping, the torque which makes it
roll to the right ($+R$) will determine
the sign of angular acceleration α .
Thus

$$+R - bF = \frac{mR^2}{2} \alpha \quad (2) \quad (2)$$

Rolling without slipping $\rightarrow a = R\alpha \quad (3)$

Solving (1), (2), (3), we get,

$$F_{\max} = \frac{3\mu mgR}{2b + R} \quad (2) \quad \text{(For } F_{\max}, \text{ we need } f_{\max} = \mu mg).$$

$$c) \vec{L} = I\omega = mR^2 \sin^2 \phi \omega$$

$$\boxed{\frac{dL}{dt} = 2mR^2 \omega \sin \phi \cos \phi \Omega} \quad [2]$$

This matches torque in part b).

d) Equation of motion in rotating frame

i) of hoop is:

$$\vec{F}' = \vec{F} - m\vec{\omega} \times (\vec{\omega} \times \vec{r}) - \cancel{2m\vec{\omega} \times \vec{v}_{rot}} - 2m\vec{\omega} \times \vec{v}_{rot} = m\vec{a}'$$

Here,

\vec{F}' = Forces on bug in rotating frame

\vec{F} = Vector sum of real forces on bug

\vec{a}' = acceleration of bug in rot. frame

The rest are fictitious forces.

$$\vec{F} = \vec{F}_r + \vec{F}_\theta$$

$$|\vec{F}_r| = mR \sin \phi \omega^2 \text{ axially inwards}$$

$$|F_\theta| = 2mR \cos \phi \Omega \omega \text{ pointing into the page.}$$

$$v = R \sin \phi \Omega$$

$$\rightarrow |2m\vec{\omega} \times \vec{v}| = |2m\omega v \sin(90^\circ - \phi)|$$

$$= 2m\omega R \sin \phi \cos \phi$$

$$\vec{F}_{cor} \rightarrow -2m\vec{\omega} \times \vec{v} = 2m\omega R \sin \phi \cos \phi \text{ pointing out of page} \quad [2]$$

$$\vec{\omega} \times (\vec{\omega} \times \vec{r}) = m\omega^2 R \sin \phi \text{ pointing axially outwards.} \quad [2]$$

$$\text{Thus } \vec{F}' = 0 = m\vec{a}'$$

$$\Rightarrow \boxed{\vec{a}' = 0} \quad [4]$$

Here we have neglected $-R\Omega^2 \hat{r}$.
= acceleration of bug.

$$Q.3(A)(a) : v_{ph} = v/k \Rightarrow v = kv_{ph};$$

$$v_g = \frac{dv}{dk} = v_{ph} + k \frac{dv_{ph}}{dk} = \sqrt{\frac{a}{k} + \frac{k}{b}} + k \frac{d}{dk} \left(\sqrt{\frac{a}{k} + \frac{k}{b}} \right)$$

$$\Rightarrow v_g = \frac{1}{v_{ph}} \left(\frac{a}{2k} + \frac{3k}{2b} \right) \text{-----}(3)$$

$$(b) v_{ph} = v_g \Rightarrow v_{ph} = \frac{1}{v_{ph}} \left(\frac{a}{2k} + \frac{3k}{2b} \right) \Rightarrow v_{ph}^2 = \frac{a}{k} + \frac{k}{b} = \frac{a}{2k} + \frac{3k}{2b} \Rightarrow k = \pm \sqrt{ab} \text{-----}(2)$$

Q.3(B)(a) From Fig.:

$$\begin{aligned} m\ddot{y}_1 &= -T \sin \theta_1 + T \sin \theta_2 = -T \frac{y_1}{l} + T \frac{y_2 - y_1}{l} = -\frac{2T}{l} y_1 + \frac{T}{l} y_2 \\ m\ddot{y}_2 &= -T \sin \theta_2 - T \sin \theta_3 = -T \frac{y_2 - y_1}{l} - T \frac{y_2 - y_3}{l} = -\frac{2T}{l} y_2 + \frac{T}{l} y_1 + \frac{T}{l} y_3 \\ m\ddot{y}_3 &= T \sin \theta_3 - T \sin \theta_4 = T \frac{y_2 - y_3}{l} - \frac{T}{l} y_3 = -\frac{2T}{l} y_3 + \frac{T}{l} y_2 \end{aligned} \text{-----}(3)$$

$$(b) A_{pn} = C_n \sin \left(\frac{pn\pi}{N+1} \right); p = \text{principle number}, n = \text{normal mode number}; \text{Here : } N = 3$$

$$\therefore A_{pn} = C_n \sin \left(\frac{pn\pi}{4} \right); \text{For, } n = 1 : A_{p1} = C_1 \sin \left(\frac{p\pi}{N+1} \right)$$

$$\begin{aligned} \text{For, } n = 1, A_{11} : A_{21} : A_{31} &= 1 : \sqrt{2} : 1 \\ \text{For } n = 2, A_{12} : A_{22} : A_{32} &= 1 : 0 : -1 \\ \text{For, } n = 3, A_{13} : A_{23} : A_{33} &= 1 : -\sqrt{2} : 1 \end{aligned} \text{-----}(3)$$

(c) Using the results of (a) and substituting $y_1 = A_1 \cos \omega t$, $y_2 = A_2 \cos \omega t$ and $y_3 = A_3 \cos \omega t$ we have :

$$\begin{aligned} \text{In 1st mode, } \frac{A_1}{A_2} &= \frac{T/ml}{\frac{2T}{ml} - \omega^2} \Rightarrow \frac{1}{\sqrt{2}} = \frac{T/ml}{\frac{2T}{ml} - \omega_1^2} \Rightarrow \omega_1 = 76.54 \text{ rad/sec} \\ \text{In 2nd mode, } \frac{A_2}{A_1} &= 0 = \frac{\frac{2T}{ml} - \omega_2^2}{T/ml} \Rightarrow \omega_2 = 141.4 \text{ rad/sec} \\ \text{In 3rd mode, } \frac{A_1}{A_2} &= -\frac{1}{\sqrt{2}} = \frac{T/ml}{\frac{2T}{ml} - \omega_3^2} \Rightarrow \omega_3 = 184.78 \text{ rad/sec} \end{aligned} \text{-----}(4)$$

Q.3(C)(a) Total energy E is constant.

$$E = \frac{1}{2} I \dot{\theta}^2 + Mg \frac{L}{2} (1 - \cos \theta) + \frac{1}{2} k (x_1 - x_2)^2 + \frac{1}{2} M \dot{x}_2^2 + \frac{1}{2} k x_2^2 \quad \text{----- (4)}$$

(b) For small angle : $\theta = \frac{x_1}{L}$ and $1 - \cos \theta = \frac{x_1^2}{2L^2}$

$$\therefore E = \frac{1}{6} M \dot{x}_1^2 + \frac{Mg x_1^2}{4L} + \frac{1}{2} k (x_1 - x_2)^2 + \frac{1}{2} M \dot{x}_2^2 + \frac{1}{2} k x_2^2$$

$$\text{Now, } \frac{\delta E}{\delta t} = \frac{\delta E}{\delta x_1} \cdot \frac{\delta x_1}{\delta t} = 0 \Rightarrow \left[\frac{1}{3} \ddot{x}_1 + \frac{g}{2L} x_1 + \frac{k}{M} (x_1 - x_2) \right] = 0 \quad \text{----- (3)}$$

$$\text{Similarly, } \left[\ddot{x}_2 + \frac{k}{M} x_2 - \frac{k}{M} (x_1 - x_2) \right] = 0 \quad \text{----- (2)}$$

(c) Given : $\frac{2k}{M} = \frac{g}{L}$; So, above two equations become :

$$\left[\ddot{x}_1 + \frac{3g}{2L} x_1 + \frac{3g}{2L} (x_1 - x_2) \right] = 0 \quad \text{and} \quad \left[\ddot{x}_2 - \frac{g}{2L} x_1 + \frac{g}{L} x_2 \right] = 0$$

In normal modes: $x_1 = A \cos \omega t$, $x_2 = B \cos \omega t$

$$\therefore \ddot{x}_1 + \frac{3g}{2L} x_1 + \frac{3g}{2L} (x_1 - x_2) = 0 \Rightarrow \left(\frac{3g}{L} - \omega^2 \right) A - \frac{3g}{2L} B = 0$$

$$\ddot{x}_2 - \frac{g}{2L} x_1 + \frac{g}{L} x_2 = 0 \Rightarrow -\frac{g}{2L} A + \left(\frac{g}{L} - \omega^2 \right) B = 0$$

$$\text{Solving : } \omega = \sqrt{\frac{g}{2L} (4 \pm \sqrt{7})} \quad \text{----- (6)}$$

Solution:

Q4 (A) (a) The point of suspension is driven by the projection of rotating motion of the contact point of the disc. Let it be $x_p = A \cos \omega t$ from the equilibrium position (from the position in groove when the contact point is at $\theta = 90^\circ$).

The equation of motion is;

$$F = -mg \frac{(x-x_p)}{l} - b \dot{x} = -\frac{mg x}{l} + \frac{mg R \cos \omega t}{l} - b \dot{x}$$

Finally we get equation for simple pendulum;

$$\ddot{x} + \gamma \dot{x} + \omega_0^2 x = R \omega_0^2 \cos \omega t, \text{ here } \omega_0^2 = g/l \text{ and } \gamma = \frac{b}{m}. \quad (8)$$

$$(b) E = E_0 e^{-\gamma t} = E_0 e^{-1}, \text{ which gives } \gamma = \frac{1}{t} = \frac{\omega_0}{n2\pi} = \frac{\sqrt{\frac{g}{l}}}{100\pi}. \quad (2)$$

$$A(\omega) = \frac{F_0/m}{\sqrt{(\omega_0^2 - \omega^2)^2 + (\gamma\omega)^2}} = \frac{gR/l}{\sqrt{\left(\frac{g}{l} - \omega^2\right)^2 + \left(\frac{\sqrt{g}}{100\pi\sqrt{l}}\omega\right)^2}}$$

$$\delta(\omega) = \tan^{-1} \left(\frac{\gamma\omega}{\omega_0^2 - \omega^2} \right) = \tan^{-1} \left(\frac{\frac{\sqrt{g}}{l} \frac{\omega}{100\pi}}{\frac{g}{l} - \omega^2} \right) \quad (6)$$

$$(c) P_{max} = \frac{Q F_0^2}{2 m \omega_0} = \frac{2 \pi n \left(\frac{mg l}{R}\right)^2}{2 \times 0.1 \times \sqrt{\frac{g}{l}}} = 0.14 \text{ Watts} \quad (4)$$

Q4 (B) (a) Instantaneous Work done:

$$dW = F dx = F v dt = b v v dt = b v^2 dt. \text{ Therefore, } \frac{dW}{dt} = b v^2. \quad (4)$$

(b) Mean rate of doing work

$$W = \frac{1}{T} \int dW = \frac{1}{T} \int_0^T b v^2 dt = \frac{1}{T} \int_0^T b A^2 \omega^2 \sin^2 (\omega t - \delta) dt = \frac{b \omega^2 A^2}{2} \quad (6)$$

