

CHEM F111: General Chemistry Semester II: AY 2017-18

Lecture-09, 31-01-2018

Summary: Lecture - 08



Hydrogen atom wavefunction:

$$\Psi_{n,l,m_l}(r,\theta,\varphi) = N R_{n,l}(r) \Theta_{l,m_l}(\theta) \Phi_{m_l}(\varphi)$$

Orbital	n	l	$R_{n,l}$	l	m_l	Y_{l,m_l}
1s	1	0	$2\left(\frac{Z}{z}\right)^{3/2}e^{-\rho/2}$	0	0	$\left(\frac{1}{4\pi}\right)^{1/2}$
2s	2	0	$\frac{1}{2(2)^{1/2}} \left(\frac{Z}{a_0}\right)^{3/2} \left(2 - \frac{1}{2}\rho\right) e^{-\rho/2}$	4 1	0	$\left(\frac{3}{4\pi}\right)^{1/2}\cos\theta$
2p	2	1	$ \frac{1}{2(2)^{1/2}} \left(\frac{Z}{a_0}\right)^{3/2} \left(2 - \frac{1}{2}\rho\right) e^{-\rho/2} \\ \frac{1}{4(6)^{1/2}} \left(\frac{Z}{a_0}\right)^{3/2} \rho e^{-\rho/4} $		±1	$\mp \left(\frac{3}{8\pi}\right)^{1/2} \sin\theta e^{\pm i\phi}$

$$\rho = 2Zr/a_0$$

Summary: Lecture - 08



H-atom/Hydrogen like atom

- Consider an electron in an energy state represented by n and orbital angular momentum I.
- An atom is a centrally symmetric system.
- Three spatial coordinates are expressed in: r, θ , Φ .
- The θ and Φ angular dependence is represented by $Y_l^{m_l}$
- The spatial part of the wave function is represented by

$$\Psi_{n,l,m_l}(r,\theta,\varphi) = N R_{n,l}(r) Y_l^{m_l}(\theta,\varphi)$$

 $\Psi_{n,l,m_l}(r,\theta,\varphi)$ simultaneous eigen functions of \widehat{H} , \widehat{L}^2 , and \widehat{L}_z

$$\widehat{L^2} = \widehat{L_x^2} + \widehat{L_y^2} + \widehat{L_z^2}$$

Summary: Lecture - 08



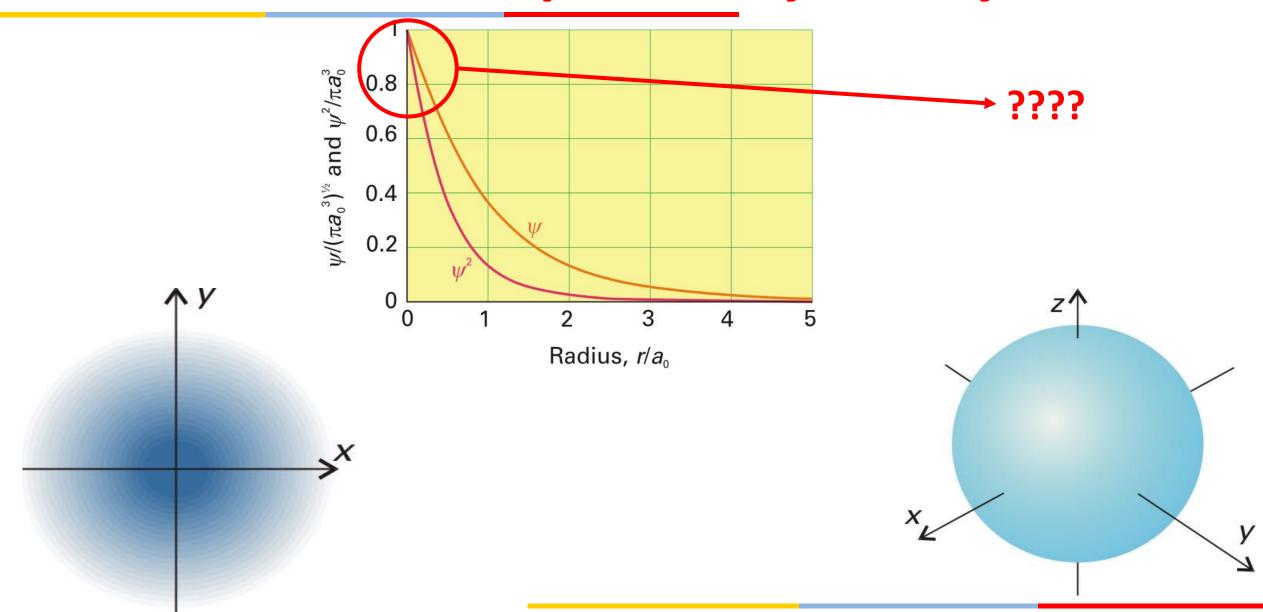
Work out: Determine the (i) energy, (ii) angular momentum, and (iii) Z-component of angular momentum for an in H-atom having the following wave function $\{a_0 \text{ is } \}$

$$\Psi_{n,l,m}(r,\theta,\varphi) = \frac{1}{\sqrt{32\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \left(\frac{Zr}{a_0}\right) e^{-Zr/2a_0} \cos\theta$$

Work out: Calculate the probability at the nucleus for an electron with n=1, l=0, and $m_l=0$.

1s wave function and probability density





1s – radial distribution function



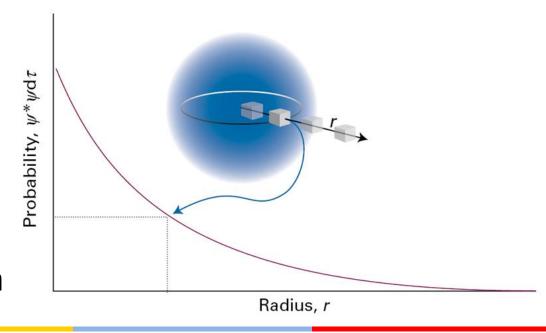
$$|\psi|^2 \Rightarrow |\Psi_{n,l,m_l}(r,\theta,\varphi)|^2 \Rightarrow$$
 Probability of finding an electron in any region

- $|\psi|^2 \Rightarrow$ Dimensions: 1/volume
- $|\psi|^2$ multiplied by infinitesimal volume element \Rightarrow dimensionless, probability

$$\Psi_{1,0,0}(r,\theta,\varphi) = \frac{1}{\sqrt{\pi}} \left(\frac{Z}{a_0}\right)^{3/2} e^{-Zr/a_0}$$

$$|\Psi_{1,0,0}(r,\theta,\varphi)|^2 \alpha e^{-2Zr/a_0}$$

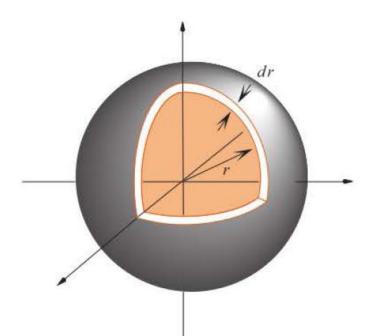
- Imagine an electron probe with infinitesimal small volume $d\tau$ (dx dy dz)
- Gives the greatest reading at the nucleus.
- Reading decreases along any radius vector.
- Reading is constant if the probe is moved on a circle of constant radius.



1s – radial distribution function



How do we measure the probability of finding the electron anywhere between the two walls of a spherical shell of thickness dr at radius r?



The sensitive volume of the probe: $4\pi r^2 dr$ Probability:

$$|\psi|^2 \times 4\pi r^2 dr = P(r) dr$$

$$P(r) = 4\pi r^2 |\psi|^2$$

P(r) is called the radial distribution function – when P(r) is multiplied by the dr, it gives the probability of finding the electron anywhere between the two walls of a spherical shell of thickness dr at radius r

1s – radial distribution function

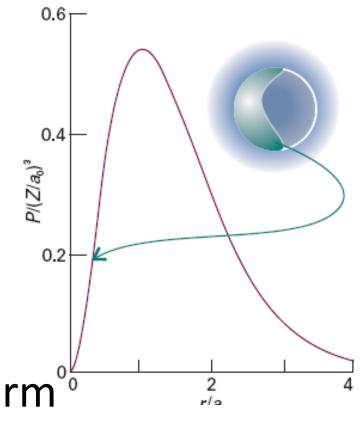


For, H-atom 1s orbital:

$$P(r) = \frac{4Z^3}{a_0^3} r^2 e^{-2Zr/a_0}$$

Interpretation:

- 1. $r^2 = 0$ at the nucleus, P(0) = 0
- 2. As $r \to \infty$, P(r) $\to 0$ due to exponential term $^{\circ}$
- 3. P(r) passes through a maxima.



All orbitals are not spherically symmetric!!

Radial distribution function – general form



$$\Psi_{2,1,0}(r,\theta,\varphi) = \frac{1}{\sqrt{32\pi}} \left(\frac{z}{a_0}\right)^{3/2} \left(\frac{zr}{a_0}\right) e^{-zr/2a_0} \cos\theta$$

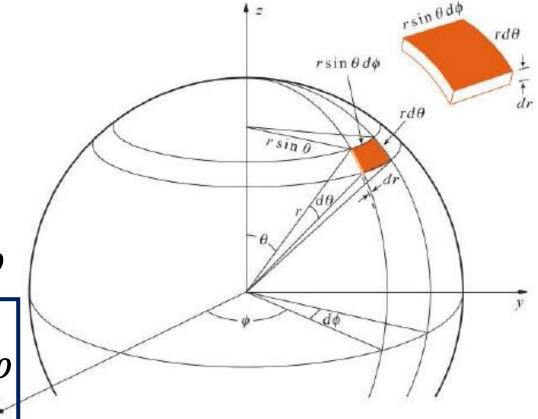
Wave function: $R(r) Y_l^{m_l}$ $d\tau = r^2 dr \sin\theta d\theta d\phi$ P(r) dr

$$= \int_0^\pi \int_0^{2\pi} R(r)^2 |Y_l^{m_l}|^2 r^2 dr \sin\theta d\theta d\phi$$

$$= \int_{0}^{\pi} \int_{0}^{2\pi} R(r)^{2} |Y_{l}^{m_{l}}|^{2} r^{2} dr \sin \theta d\theta d\varphi$$

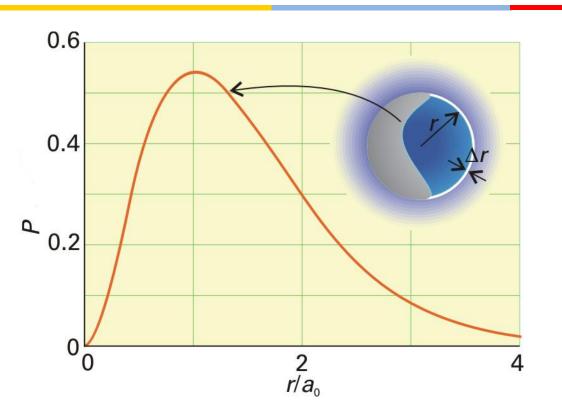
$$= r^{2} |R(r)|^{2} dr \int_{0}^{\pi} \int_{0}^{2\pi} |Y_{l}^{m_{l}}|^{2} \sin \theta d\theta d\varphi$$

$$= r^2 |R(r)|^2 dr$$

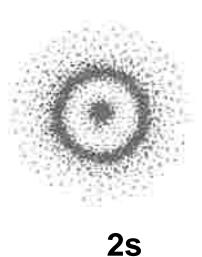


Most probable radius









$$\frac{dP(r)}{dr} = 0$$

The maximum of P(r) marks the most probable radius at the electron will be found!! For a 1s orbital in hydrogen, P(r) is maximum at $r = a_0$ (Bohr radius)

Most probable radius

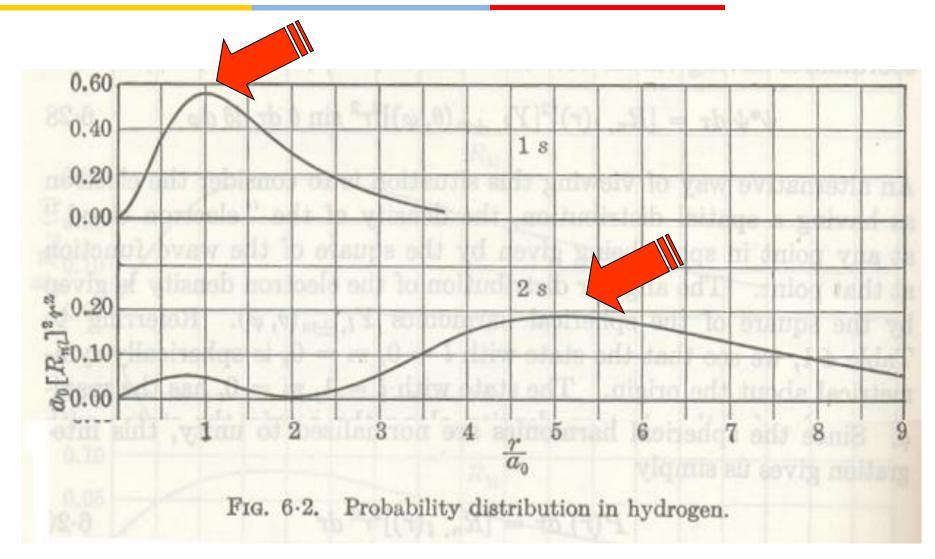


Calculate the most probable radius, r^* , at which an electron will be found when it occupies a 1s orbital of a hydrogenic atom of atomic number Z (H to Ne⁹⁺)

1s orbital is drawn towards the nucleus as the nuclear charge increases

Most probable radius





For 2s orbital: Most probable radius is 5.2 $a_0 = 275 \text{ pm}$

Hydrogen atom wave function for n = 2



n = 2 state
$$\Rightarrow$$
 $l=0$, 1; $for l=1$, $m_l=0$ and 1, 0 , -1

$$\Psi_{2,0,0}(r,\theta,\varphi) = \frac{1}{\sqrt{32\pi}} \left(\frac{z}{a_0}\right)^{3/2} \left(2 - \frac{zr}{2a_0}\right) e^{-zr/2a_0}$$

$$\Psi_{2,1,0}(r,\theta,\varphi) = \frac{1}{\sqrt{32\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \left(\frac{Zr}{a_0}\right) e^{-Zr/2a_0} \cos\theta$$

$$\Psi_{2,1,\pm 1}(r,\theta,\varphi) = \frac{1}{\sqrt{64\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \left(\frac{Zr}{a_0}\right) e^{-Zr/2a_0} \sin\theta e^{\pm i\varphi}$$

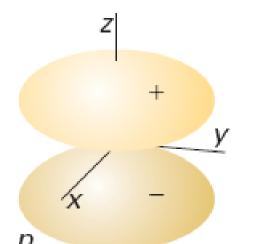
Hydrogen atom wave function for n = 2



2s orbital

$$\Psi_{2,0,0}(r,\theta,\varphi) = \frac{1}{\sqrt{32\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \left(2 - \frac{Zr}{2a_0}\right) e^{-Zr/2a_0} = f(r)$$

$$\Psi_{2,1,0}(r,\theta,\varphi) = \frac{1}{\sqrt{32\pi}} \left(\frac{z}{a_0}\right)^{3/2} \left(\frac{zr}{a_0}\right) e^{-zr/2a_0} \cos\theta$$



$$= r \cos \theta f(r) = z f(r)$$

$$\Psi_{p_z} = z f(r)$$

XY plane is nodal plane – angular node is one

Hydrogen atom wave function for n = 2



$$\Psi_{2,1,\pm 1}(r,\theta,\varphi) = \frac{1}{\sqrt{64\pi}} \left(\frac{Z}{a_0}\right)^{3/2} \left(\frac{Zr}{a_0}\right) e^{-Zr/2a_0} \sin\theta e^{\pm i\varphi}$$

What is your idea about p-orbitals?

$$\psi_{p_{x}} = -\frac{1}{2^{1/2}}(p_{+1} - p_{-1}) = r \sin \theta \cos \phi f(r) = xf(r)$$

$$\psi_{p_{y}} = \frac{i}{2^{1/2}}(p_{+1} + p_{-1}) = r \sin \theta \sin \phi f(r) = yf(r)$$

Also an energy eigen function with same energy and eigen function of the angular momentum operator

Orthonormal wave functions



$$\int \psi_m^* \, \psi_n \, d\tau = \delta_{mn} \qquad \begin{array}{l} \delta_{mn} = 1 \text{ if m = n} \\ = 0 \text{ if m} \neq n \end{array}$$

- **\$** Show that hydrogenlike atomic wave function $\Psi_{2 \ 1 \ 0}$ is normalized and it is orthogonal to $\Psi_{2 \ 0 \ 0}$.
- Show that the PIB wave functions satisfy the relation: $\int_0^a \psi_n^*(x) \, \psi_m(x) dx = 0 \ for \ m \neq n$

This concept would be used in bonding

What do we learn?



Do consider an experimental set-up in which you can populate $2p_x$ state of H-atom. Now, try to measure the following parameters:

- 1) Energy
- 2) Total angular momentum
- 3) Z-component of angular momentum

What would be the outcome?

$$\psi_{p_x} = -\frac{1}{2^{1/2}}(p_{+1} - p_{-1}) = r \sin \theta \cos \phi f(r) = xf(r)$$

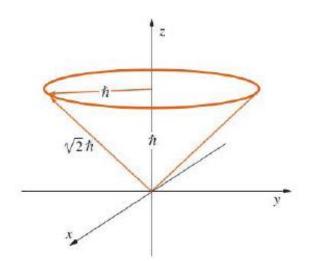
What do we learn?



- Atomic states
- Energy: 2s and 2p are having same energy
- Angular momentum: We can differentiate 2s and 2p
- There are three states in 2p??
- One of the component of angular momentum Z-component

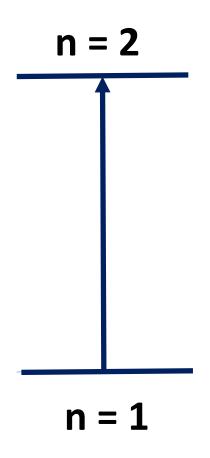
We can measure total angular momentum and any one of the

components: L_x , L_y , or L_z



A thought experiment





- We'll measure the z-component of angular momentum
- We'll apply a magnetic field
- We'll perturb the system.

We can only say about the state of the system when we are trying to measure