Classical Mechanics: Assignment #2

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Problem 1

Solution

The Lagrangian for the given system can be written as,

$$L = T + V = \frac{1}{2}mx^2\omega^2 + \frac{1}{2}m\left(\dot{x}^2 + \dot{y}^2\right) - mgy$$

From the problem, we know that $y = k\left(\frac{x}{l}\right)^\alpha$, which means that $\dot{y} = k\alpha\frac{x^{\alpha-1}}{l^\alpha}\dot{x}$. Substituting these into the form of the Lagrangian and simplifying, we get,

$$L = \frac{1}{2}m\left(-2gk\left(\frac{x}{l}\right)^\alpha + \dot{x}^2\left(\frac{\alpha^2k^2x^{2\alpha-2}}{l^\alpha} + 1\right) + x^2\omega^2\right)$$

The equation of motion can be written as,

$$\boxed{\alpha gkx^2\left(\frac{x}{l}\right)^\alpha + (\alpha-1)\alpha^2k^2\dot{x}^2\left(\frac{x}{l}\right)^{2\alpha} + \alpha^2k^2x\ddot{x}\left(\frac{x}{l}\right)^{2\alpha} - x^4\omega^2 + x^3\ddot{x} = 0}$$

The equlibrium points will satisfy $\dot{x} = \ddot{x} = 0$. This means that the equilibrium point will be,

$$x_0 = \left(\frac{\omega^2l^\alpha}{gk\alpha}\right)^{\frac{1}{\alpha-2}}$$

We substitute $x = x_0 + \epsilon$

$$\ddot{y} + y\frac{(\alpha-2)\omega^2}{\alpha^2k^2\left(\frac{x_0}{l}\right)^{2\alpha} + 1} = 0$$

For small oscillations, the coefficient of y in the above equation should be positive. Hence,

$$\frac{(\alpha-2)\omega^2}{\alpha^2k^2\left(\frac{x_0}{l}\right)^{2\alpha} + 1} > 0$$

$$\therefore \alpha - 2 > 0 \implies \boxed{\alpha > 2}$$

The frequency of oscillations ω_0 is simply the square root of the coefficient of y,

$$\boxed{\omega_0 = \sqrt{\frac{\alpha-2}{\alpha^2k^2\left(\frac{x_0}{l}\right)^{2\alpha} + 1}}\omega} \quad \text{where} \quad x_0 = \left(\frac{\omega^2l^\alpha}{gk\alpha}\right)^{\frac{1}{\alpha-2}}$$

Problem 2

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Problem 3

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Problem 4

Part (a)
The Schrodinger equation is given by

$$-\frac{\hbar^2}{2m}\frac{\partial^2\psi}{\partial x^2} + V\psi = i\hbar\frac{\partial\psi}{\partial t} \quad \text{and} \quad -\frac{\hbar^2}{2m}\frac{\partial^2\psi^*}{\partial x^2} + V\psi^* = -i\hbar\frac{\partial\psi^*}{\partial t}$$

We choose ψ and ψ^* as our generalized coordinates, and (t, x) as the dependent coordinates. One should be able to write the equations of motion in a compact form as follows,

$$\partial_\mu\left(\frac{\partial L}{\partial(\partial_\mu\psi)}\right) = \frac{\partial L}{\partial\psi} \quad \text{and} \quad \partial_\mu\left(\frac{\partial L}{\partial(\partial_\mu\psi^*)}\right) = \frac{\partial L}{\partial\psi^*}$$

where the index μ goes over (t, x). Let's analyze every term in the Schrodinger Equation and figure out what corresponding term in the Lagrangian will give rise to that term,

- The first term on the LHS is a double x derivative and will come from some single derivative term of the form $L_1 = \psi'\psi'^*$
- The second term on the LHS has no derivatives and will come from a term of the form $L_2 = \psi\psi^*$
- The RHS is a single t derivative and will come from some term of the form $\dot{\psi}\psi^*$. To make it symmetric, let's consider $L_3 = -\dot{\psi}\psi^* + \dot{\psi}^*\psi$

So our final Lagrangian will be of the form $L = a_1L_1 + a_2L_2 + a_3L_3$. Substituting our ansatz, we find our constants, and then the final Lagrangian can be written as,

$$\boxed{L = -\frac{\hbar^2}{2m}\psi'\psi'^* + V\dot{\psi}\psi^* + i\hbar(-\dot{\psi}\psi^* + \dot{\psi}^*\psi)}$$

Part (b)
Kinetic energy of the wire is zero. The Lagrangian can be written as,

$$L = -\int ds\ \rho g y = -\int \sqrt{dx^2 + dy^2}\ \rho g y = -\int dx y\sqrt{1 + y'^2}\ \rho g$$

Writing down the equation of motion for the Lagrangian density instead of the Lagrangian, one gets,

$$\frac{\mathrm{d}}{\mathrm{d}x}\left(\frac{yy'}{\sqrt{1+y'^2}}\right) - \sqrt{1+y'^2} = 0$$

$$\frac{yy'' + y'^2}{\sqrt{1+y'^2}} - \frac{yy'^2y''}{1+y'^2} - \sqrt{1+y'^2} = 0$$

Expanding this out and simplifying a bit, one gets,

$$\frac{yy''}{(1+y'^2)^{\frac{3}{2}}} - \frac{1}{\sqrt{1+y'^2}} = 0 \implies \frac{\mathrm{d}}{\mathrm{d}x}\left(\frac{y}{\sqrt{1+y'^2}}\right) = 0$$

$$\therefore \frac{y}{\sqrt{1+y'^2}} = \alpha \implies \boxed{y = \alpha\cosh\left(\frac{x}{\alpha} + \beta\right)}$$

Problem 4 continued on next page...

One can get the constants α and β by imposing the end point conditions for the curve.

Part (c)

The distance metric on a sphere spherical polar coordinates is given by,

$$ds^2 = r^2(d\theta^2 + \sin^2\theta d\phi^2) = d\theta^2[r^2(1 + \sin^2\theta\phi'^2)]$$

$$\therefore ds = d\theta\sqrt{r^2(1 + \sin^2\theta\phi'^2)}$$

From the ansatz $S = Ld\tau$, we can identify that the Lagrangian $L = \sqrt{r^2(1 + \sin^2\theta\phi'^2)}$. For finding the *equations of motion*, it is fine and also easier to work with L^2 rather than L in this problem. Writing down the equations of motion for $\phi(\theta)$,

$$\frac{\mathrm{d}\sin^2\theta\phi'}{\mathrm{d}\theta} = 0 \implies \phi' = \alpha\csc^2(\theta) \implies \boxed{\phi(\theta) = a\cot\theta + b}$$

where α, a, b are constants. If the distance is to be found out between two points (ϕ_1, θ_1) and (ϕ_2, θ_2), then,

$$\phi_1 = a\cot\theta_1 + b \quad \text{and} \quad \phi_2 = a\cot\theta_2 + b$$

which gives,

$$a = \frac{\phi_1 - \phi_2}{\cot\theta_1 - \cot\theta_2} \quad \text{and} \quad b = \frac{\phi_1\tan\theta_1 - \phi_2\tan\theta_2}{\tan\theta_1 - \tan\theta_2}$$