

Advanced Quantum Mechanics: Assignment #5

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Problem 1

Problem 2

We don't need to apply any perturbation theory in this problem, and it can be solved exactly. The Hamiltonian is $H = \lambda S_1 \cdot S_2 = \lambda(S^2 - S_1^2 - S_2^2)$. We consider the action of the Hamiltonian on the singlet state $|00\rangle = \frac{|+-\rangle - |-+\rangle}{\sqrt{2}}$ and $|10\rangle = \frac{|+-\rangle + |-+\rangle}{\sqrt{2}}$. We know $S^2|00\rangle = 0$ and $S^2|10\rangle = |10\rangle$. Initially the system is in $|+-\rangle = \frac{|00\rangle + |10\rangle}{\sqrt{2}}$. Then we know, by the usual rules of time-evolution,

$$\begin{aligned} |\psi_f(t)\rangle &= e^{iHt} |+-\rangle = \frac{e^{i\lambda t/4}}{\sqrt{2}} |10\rangle + \frac{e^{-i3\lambda t/4}}{\sqrt{2}} |00\rangle \\ &= \left(\frac{e^{i\lambda t/4} + e^{-i3\lambda t/4}}{2} \right) |+-\rangle + \left(\frac{e^{i\lambda t/4} - e^{-i3\lambda t/4}}{2} \right) |-+\rangle \\ \Rightarrow |\langle+-|\psi_f(t)\rangle|^2 &= \left| \left(\frac{e^{i\lambda t/4} + e^{-i3\lambda t/4}}{2} \right) \right|^2 = \frac{1 + \cos \lambda t}{2} = P(|+-\rangle) \\ \Rightarrow |\langle-+|\psi_f(t)\rangle|^2 &= \left| \left(\frac{e^{i\lambda t/4} - e^{-i3\lambda t/4}}{2} \right) \right|^2 = \frac{1 - \cos \lambda t}{2} = P(|-+\rangle) \\ \Rightarrow |\langle++|\psi_f(t)\rangle|^2 &= 0 = P(|++\rangle) \\ \Rightarrow |\langle--|\psi_f(t)\rangle|^2 &= 0 = P(|--\rangle) \end{aligned}$$

where $P(|\rangle)$ denotes probability of initial state to be in state $|\rangle$.

Problem 3

From (5.7.17) of Sakurai, we have the relations, (with $|i, t_0; t\rangle = \sum c_n(t) |n\rangle$)

$$c_n^0(t) = \delta_{ni} \quad , \quad c_n^1(t) = -i \int_{t_0}^t e^{i\omega_{ni}t'} V_{ni}(t') dt' \quad , \quad c_n^2(t) = - \sum_m \int_{t_0}^t dt' \int_{t_0}^{t'} e^{i\omega_{nm}t'} V_{nm}(t') e^{i\omega_{mi}t''} V_{mi}(t'') dt''$$

For our problem, we have $V = \lambda\delta(x - vt)$. We insert $1 = \int dx |x\rangle \langle x|$ such that $V_{ni}(t) = \int V(t) u_i^*(x) u_n(x) dx$. We have initial state $u_i(x)$ and final state $u_f(x)$. Hence, we can write the above coefficients as,

$$\begin{aligned} c_f^1(t) &= -i \int_{-\infty}^{\infty} dx \int_0^t dt' e^{i(E_i - E_f)t'} \delta(x - vt') u_i^*(x) u_n(x) \\ &= -i \int_{-\infty}^{\infty} dx e^{i(E_i - E_f)x/v} u_i^*(x) u_n(x) \end{aligned}$$

Hence the probability is just $|c_f^1|^2$

Problem 4

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Problem 5

We first note for $|\psi_I(t)\rangle = \sum_n c_n(t) |\alpha_n\rangle$

$$\begin{aligned} i \frac{\partial |\psi_I\rangle}{\partial t} &= i \frac{\partial (e^{iH_0 t} |\psi_S\rangle)}{\partial t} \\ &= i \left[e^{iH_0 t} \frac{\partial |\psi_S\rangle}{\partial t} + iH_0 e^{iH_0 t} |\psi_S\rangle \right] \\ &= -e^{iH_0 t} (H_0 + V) |\psi_S\rangle - H_0 e^{iH_0 t} |\psi_S\rangle \\ &= e^{iH_0 t} V |\psi_S\rangle \\ i \frac{\partial |\psi_I\rangle}{\partial t} &= V_I |\psi_I\rangle \\ i \frac{\partial \langle \alpha_n | \psi_I \rangle}{\partial t} &= \langle \alpha_n | V_I | \psi_I \rangle \\ \dot{c}_n &= -i \langle \alpha_n | V_I | \psi_I \rangle \\ \dot{c}_n &= -i \langle \alpha_n | V | \alpha_m \rangle e^{i(E_n - E_m)t} c_m \end{aligned}$$

So for the given problem, we have

$$\begin{aligned} |\psi_I(t)\rangle &= c_1(t) |1\rangle + c_2(t) e^{iEt} |2\rangle \\ \dot{c}_1 &= -iV_{11}c_1 - iV_{12}e^{-iEt}c_2 = -i\gamma e^{i(\omega - E)t}c_2 \quad \text{and} \quad \dot{c}_2 = -iV_{21}e^{iEt}c_1 - iV_{22}c_2 = -i\gamma e^{i(E - \omega)t}c_1 \end{aligned}$$

To solve the above equations, we make the substitution $c_1 = b_1 e^{i\Delta t}$ and $c_2 = b_2 e^{-i\Delta t}$, where $2\Delta = \omega - E$. We then have the equations in terms of b 's,

$$i\dot{b}_1 = \Delta b_1 + \gamma b_2 \quad \text{and} \quad i\dot{b}_2 = \gamma b_1 - \Delta b_2$$

These are coupled equations, and we can solve these by making the substitution $b_1 = A e^{i\Omega t}$ and $b_2 = B e^{i\Omega t}$. We then have,

$$\begin{aligned} -A\Omega &= \Delta A + \gamma B \quad \text{and} \quad -B\Omega = \gamma A - \Delta B \\ \text{For non-trivial solutions,} \quad -\frac{\gamma}{\Delta + \Omega} &= \frac{\Delta - \Omega}{\gamma} \implies \Omega = \pm \sqrt{\gamma^2 + \Delta^2} = \pm \Omega_0 \\ \implies c_1 &= A_1 e^{i(\Delta + \Omega_0)t} + A_2 e^{i(\Delta - \Omega_0)t} \quad \text{and} \quad c_2 = B_1 e^{i(-\Delta + \Omega_0)t} + B_2 e^{i(-\Delta - \Omega_0)t} \end{aligned}$$

We are told that at $t = 0$, the system is in state $|1\rangle \implies c_1(0) = 0, c_2(0) = 1 \implies A_1 = -A_2, B_1 = 1 - B_2$. We also know for a fact that $|c_1(0)|^2 + |c_2(0)|^2 = 1$. Using all these facts, we can write

$$4A_1^2 \cos^2 \Omega t + 4(B_1^2 - B_1) \sin^2 \Omega t + 1 = 1$$