

Introduction to GR
Assignment - 3
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Problems from D'Inverno

7.1 Let ϕ be a scalar density of weight +1.

Then $\phi' = J\phi$

Using eq (7.2) in the book,

$$\nabla_c \phi = \partial_c \phi - \Gamma_{ac}^a \phi$$

7.2 Given,

$$\nabla_a [\sqrt{-g} T^a] = \partial_a [\sqrt{-g} T^a]$$

$$\Rightarrow \nabla_a [\sqrt{-g}] T^a + \sqrt{-g} \nabla_a T^a = \partial_a (\sqrt{-g}) T^a + \sqrt{-g} \partial_a T^a$$

$$\Rightarrow \nabla_a [\sqrt{-g}] T^a + \cancel{\sqrt{-g} \partial_a T^a} + \sqrt{-g} \Gamma_{ac}^a T^c = \partial_a (\sqrt{-g}) T^a + \cancel{\sqrt{-g} \partial_a T^a}$$

$$\Rightarrow \nabla_a [\sqrt{-g}] T^a + \sqrt{-g} \Gamma_{ac}^a T^c = \partial_a [\sqrt{-g}] T^a$$

$$\Rightarrow \{ \nabla_a [\sqrt{-g}] + \sqrt{-g} \Gamma_{ca}^c \} T^a = \partial_a [\sqrt{-g}] T^a$$

$$\Rightarrow \nabla_a [\sqrt{-g}] = \partial_a [\sqrt{-g}] - \sqrt{-g} \Gamma_{ca}^c$$

which is indeed consistent with the answer of the previous problem.

7.6

We start off by considering the Euler-Lagrange equations written in a slightly modified way (in terms of \mathcal{L}^2), [eq. (7.40)],

$$\frac{d}{du} \left[\frac{\partial \mathcal{L}^2}{\partial \dot{x}^a} \right] - \frac{\partial \mathcal{L}^2}{\partial x^a} = \frac{2 \partial \mathcal{L}}{\partial \dot{x}^a} \frac{d\mathcal{L}}{du}$$

Here, $\mathcal{L}^2 = g_{ab} \dot{x}^a \dot{x}^b = 2K$

Since we are looking for spacelike geodesics, we can always choose an affine parametrization, and $\frac{d\mathcal{L}}{du} = 0$.

Substituting this fact and $2K = \mathcal{L}^2$,

$$2 \frac{d}{du} \left[\frac{\partial K}{\partial \dot{x}^a} \right] - 2 \frac{\partial K}{\partial x^a} = 0$$

$$\Rightarrow \frac{d}{du} \left[\frac{\partial K}{\partial \dot{x}^a} \right] - \frac{\partial K}{\partial x^a} = 0$$

We can always choose s to be our affine parameter in this case,

$$ds^2 = 2K = g_{ab} dx^a dx^b > 0$$

$$\therefore 1 = g_{ab} \dot{x}^a \dot{x}^b$$

7.7

$$2K = g_{ab} \dot{x}^a \dot{x}^b$$

$$= e^\nu \dot{t}^2 - e^\lambda \dot{r}^2 - r^2 \dot{\theta}^2 - r^2 \sin^2 \theta \dot{\phi}^2$$

$$\frac{\partial K}{\partial \dot{t}} = e^\nu \dot{t} \quad ; \quad \frac{\partial K}{\partial \dot{r}} = -e^\lambda \dot{r}$$

$$\frac{\partial K}{\partial \dot{\theta}} = -r^2 \dot{\theta} \quad ; \quad \frac{\partial K}{\partial \dot{\phi}} = -r^2 \sin^2 \theta \dot{\phi}$$

$$\frac{\partial K}{\partial t} = \frac{e^\nu \dot{t}^2 \bar{\nu}}{2} - \frac{e^\lambda \bar{\lambda} \dot{r}^2}{2}$$

$$\frac{\partial K}{\partial r} = \frac{e^\nu \nu' \dot{t}^2}{2} - \frac{e^\lambda \lambda' \dot{r}^2}{2} - r [\dot{\theta}^2 + \sin^2 \theta \dot{\phi}^2]$$

$$\frac{\partial K}{\partial \theta} = -r^2 \sin \theta \cos \theta \dot{\phi}^2 \quad ; \quad \frac{\partial K}{\partial \phi} = 0$$

We have used the notation $\bar{a} = \frac{\partial a}{\partial t}$ and $a' = \frac{\partial a}{\partial r}$

For t : $e^\nu [\bar{\nu} \dot{t} + \nu' \dot{r}] \dot{t} + e^\nu \ddot{t} = \frac{e^\nu \dot{t}^2 \bar{\nu}}{2} - \frac{e^\lambda \bar{\lambda} \dot{r}^2}{2}$

$$\therefore \ddot{t} = -\frac{\bar{\nu}}{2} \dot{t}^2 - \nu' \dot{r} \dot{t} - \frac{e^\lambda \bar{\lambda} \dot{r}^2 e^{-\nu}}{2}$$

For r : $-e^\lambda [\ddot{r} + \bar{\lambda} \dot{t} \dot{r} + \lambda' \dot{r}^2] = \frac{e^\nu \nu' \dot{t}^2}{2} - \frac{e^\lambda \lambda' \dot{r}^2}{2} - r [\dot{\theta}^2 + \sin^2 \theta \dot{\phi}^2]$

$$\therefore \ddot{r} = -\frac{e^{\nu-\lambda} \nu' \dot{t}^2}{2} - \frac{\lambda' \dot{r}^2}{2} - \bar{\lambda} \dot{t} \dot{r} + r e^{-\lambda} [\dot{\theta}^2 + \sin^2 \theta \dot{\phi}^2]$$

For θ : $-\kappa^2 \ddot{\theta} - 2\kappa \dot{\kappa} \dot{\theta} = -\kappa^2 \sin\theta \cos\theta \dot{\phi}^2$

$$\Rightarrow \ddot{\theta} = -\frac{2\dot{\kappa}}{\kappa} \dot{\theta} + \sin\theta \cos\theta \dot{\phi}^2$$

For ϕ : $-\left[2\kappa \dot{\kappa} \sin\theta \dot{\phi} + 2\kappa^2 \cos\theta \dot{\phi} \dot{\theta} + \kappa^2 \sin\theta \ddot{\phi}\right] = 0$

$$\Rightarrow \ddot{\phi} = -\frac{2\dot{\kappa}}{\kappa} \dot{\phi} - 2\cot\theta \dot{\phi} \dot{\theta}$$

Now we read off the Γ 's,

$$\Gamma_{\kappa\phi}^{\phi} = \Gamma_{\phi\kappa}^{\phi} = +\frac{1}{\kappa} ; \quad \Gamma_{\theta\phi}^{\phi} = \Gamma_{\phi\theta}^{\phi} = +\cot\theta$$

$$\Gamma_{\kappa\theta}^{\theta} = \Gamma_{\theta\kappa}^{\theta} = +\frac{1}{\kappa} ; \quad \Gamma_{\phi\phi}^{\theta} = -\sin\theta \cos\theta$$

$$\Gamma_{tt}^{\kappa} = \frac{+e^{\nu-\lambda} \nu'}{2} ; \quad \Gamma_{\kappa\kappa}^{\kappa} = +\frac{\lambda'}{2} ; \quad \Gamma_{\kappa t}^{\kappa} = \Gamma_{t\kappa}^{\kappa} = +\frac{\bar{\lambda}}{2}$$

$$\Gamma_{\theta\theta}^{\kappa} = -\kappa e^{-\lambda} ; \quad \Gamma_{\phi\phi}^{\kappa} = -\kappa e^{-\lambda} \sin^2\theta$$

$$\Gamma_{tt}^t = +\frac{\bar{\nu}}{2} ; \quad \Gamma_{\kappa t}^t = \Gamma_{t\kappa}^t = +\frac{\nu'}{2}$$

$$\Gamma_{\kappa\kappa}^t = +\frac{e^{\lambda} \bar{\lambda}}{2} e^{-\nu}$$

All the Γ 's not mentioned above are zero.

7.10

Any Killing vector must satisfy

$$\nabla_a X_b + \nabla_b X_a = 0$$

We are working in flat cartesian 3-space,
Hence all the Γ 's are zero identically.

$$\Rightarrow \nabla_b X_a = \partial_b X_a$$

$$\Rightarrow \partial_a X_b + \partial_b X_a = 0$$

$$\therefore \partial_c \partial_a X_b + \partial_c \partial_b X_a = 0 \rightarrow \textcircled{1}$$

Permuting indices, we get,

$$\rightarrow \partial_b \partial_c X_a + \partial_b \partial_a X_c = 0 \rightarrow \textcircled{2}$$

$$\rightarrow \partial_a \partial_b X_c + \partial_a \partial_c X_b = 0 \rightarrow \textcircled{3}$$

$\textcircled{1} + \textcircled{2} - \textcircled{3}$ gives,

$$\partial_b \partial_c X_a = 0$$

$$\Rightarrow X_a = \omega_{ab} x^b + t_a$$

$$\therefore X^a = \omega^a_b x^b + t^a$$

where ω^a_b and t^a are (tensor) constants of integration.

Since $\partial_b X_a = -\partial_a X_b$, $\omega_{ba} = -\omega_{ab}$

Hence ω^a_b is antisymmetric and has only 3 independent components. t^a also has three independent components, giving 6 constants of integration in total.

Moreover, we can see that these transformations can be split into 3 translations (along x, y, z), rotations around three axes.

$$\begin{aligned} \therefore X^a = & \lambda_1 (\partial_x) + \lambda_2 (\partial_y) + \lambda_3 (\partial_z) \\ & + \lambda_4 [y\partial_x - x\partial_y] + \lambda_5 [x\partial_z - z\partial_x] \\ & + \lambda_6 [z\partial_y - y\partial_z] \end{aligned}$$

7.11 If X^a and Y^a are Killing vectors,

$$\nabla_a X_b + \nabla_b X_a = 0 \rightarrow \textcircled{1}$$

$$\nabla_a Y_b + \nabla_b Y_a = 0 \rightarrow \textcircled{2}$$

$\lambda \times \textcircled{1} + \mu \times \textcircled{2}$ gives,

$$\nabla_a [\lambda X_b + \mu Y_b] + \nabla_b [\lambda X_a + \mu Y_a] = 0$$

which shows that $[\lambda X_b + \mu Y_b]$ is a Killing vector.

7.14 $(\nabla_b \nabla_a - \nabla_a \nabla_b) X_c = R_{cdba} X^d \rightarrow \textcircled{1}$

$$\Rightarrow g^{bc} \nabla_b \nabla_a X_c - g^{bc} \nabla_a \nabla_b X_c = R_{cdca} X^d$$

If X_c is Killing vector, $\nabla_b X_c = -\nabla_c X_b$
ie. $\nabla_c X_b$ is antisymmetric

Then $\underbrace{g^{bc}}_{\text{symmetric}} \nabla_a \underbrace{[\nabla_b X_c]}_{\text{antisymmetric}} = 0$

$$\therefore g^{bc} \nabla_b \nabla_a X_c = R_{ad} X^d \rightarrow \textcircled{2}$$

Consider $\textcircled{1}$ again

$$(\nabla_b \nabla_a - \nabla_a \nabla_b) X_c = R_{cdba} X^d$$

$$\begin{aligned} \therefore \nabla_b \nabla_a X_c &= \nabla_a \nabla_b X_c + R_{cdba} X^d \\ &= \nabla_a \nabla_b X_c - (R_{cbad} + R_{cadb}) X^d \\ &= -R_{cadb} X^d + (\nabla_a \nabla_b X_c - R_{cbad} X^d) \end{aligned}$$

② also means $g^{bc} [\nabla_b \nabla_a X_c - R_{cdba} X^d] = 0$
which in turn implies $\nabla_b \nabla_a X_c = R_{cdba} X^d = R_{cbad} X^d$.

$$\therefore \nabla_b \nabla_a X_c = -R_{cadb} X^d$$

$$\therefore \nabla_c \nabla_b X_a = -R_{abdc} X^d$$

$$\nabla_c \nabla_b X_a = +R_{abcd} X^d$$

Problems not from d'Inverno

Problem 1 If λ is non-affine, and $t^\alpha = \frac{dx^\alpha}{d\lambda}$

$$\begin{aligned}\frac{d}{d\lambda} [t^\alpha \xi_\alpha] &= t^\beta \nabla_\beta [t^\alpha \xi_\alpha] \\ &= t^\beta \xi_\alpha \nabla_\beta t^\alpha + \underbrace{t^\beta t^\alpha \nabla_\beta}_{\text{symmetric}} \underbrace{\xi_\alpha}_{\text{antisymmetric}}\end{aligned}$$

$$\frac{d}{d\lambda} [t^\alpha \xi_\alpha] = t^\beta \xi_\alpha \nabla_\beta t^\alpha$$

$$\frac{d}{d\lambda} [t^\alpha \xi_\alpha] = \kappa t^\alpha \xi_\alpha \implies \frac{d}{d\lambda} [p] = \kappa p$$

$$\int \frac{dp}{p} = \int \kappa(\lambda) d\lambda \implies p \propto \exp \left[\int \kappa(\lambda) d\lambda \right]$$

Part (b)

where κ is $\left(\frac{d^2 z}{d\lambda^2} \right) / \left(\frac{dz}{d\lambda} \right)^2$.

Part (a) $\epsilon = -t^\alpha t_\alpha$

$$\begin{aligned}\frac{d\epsilon}{d\lambda} &= \frac{D\epsilon}{d\lambda} = \frac{-2Dt^\alpha}{d\lambda} t_\alpha \\ &= -2\kappa t^\alpha t_\alpha\end{aligned}$$

$$\frac{d\epsilon}{d\lambda} = 2\epsilon\kappa$$

$$\therefore \epsilon \propto \exp \left[\int 2\kappa(\lambda) d\lambda \right]$$

Problem continued on next page

Part (c) $\mathcal{L}_b g_{\alpha\beta} = 2c g_{\alpha\beta}$, $q = b_\alpha u^\alpha$

$$\frac{Dq}{d\lambda} = b_\alpha \frac{Du^\alpha}{d\lambda} + u^\alpha \frac{Db_\alpha}{d\lambda}$$

$$\frac{dq}{d\lambda} = b_\alpha (\kappa t^\alpha) + u^\alpha u^\beta \nabla_\beta b_\alpha$$

$$\text{Since, } \mathcal{L}_b g_{\alpha\beta} = \nabla_\alpha b_\beta + \nabla_\beta b_\alpha = 2c g_{\alpha\beta}$$

$$\begin{aligned} \text{Then, } u^\alpha u^\beta \nabla_\beta b_\alpha &= u^\alpha u^\beta \nabla_\alpha b_\beta \\ &= u^\alpha u^\beta (2c g_{\alpha\beta}) - u^\alpha u^\beta (\nabla_\beta b_\alpha) \end{aligned}$$

$$2u^\alpha u^\beta \nabla_\beta b_\alpha = -2c$$

$$\Rightarrow u^\alpha u^\beta \nabla_\beta b_\alpha = -c$$

$$\text{And so, } \frac{dq}{d\lambda} = \kappa b_\alpha t^\alpha - c = \kappa q - c$$

$$\text{If } \lambda = \tau, \kappa(\lambda) = 0, \text{ and hence}$$

$$\frac{dq}{d\lambda} = -c \Rightarrow q = -c\lambda + C_2$$

Problem 3

$$\nabla_\alpha \xi^\beta = \partial_\alpha \xi^\beta + \Gamma_{\alpha\delta}^\beta \xi^\delta$$

ξ 's have nonzero components only in θ, ϕ dirn.

$$\begin{aligned} \therefore \nabla_\alpha \xi^\beta &= \partial_\alpha [\xi^\theta \delta_\theta^\beta + \xi^\phi \delta_\phi^\beta] + \\ &\quad + \left(\Gamma_{\alpha\delta}^\theta \delta_\theta^\beta + \Gamma_{\alpha\delta}^\phi \delta_\phi^\beta \right) (\xi^\theta \delta_\theta^\alpha + \xi^\phi \delta_\phi^\alpha) \\ &= \left(\delta_\alpha^\theta \partial_\theta + \delta_\alpha^\phi \partial_\phi \right) \times [\xi^\theta \delta_\theta^\beta + \xi^\phi \delta_\phi^\beta] \\ &\quad + \left(\Gamma_{\alpha\delta}^\theta \delta_\theta^\beta + \Gamma_{\alpha\delta}^\phi \delta_\phi^\beta \right) (\xi^\theta \delta_\theta^\alpha + \xi^\phi \delta_\phi^\alpha) \end{aligned}$$

Now the task is to find non-trivial $\Gamma_{\alpha\beta}^\theta, \Gamma_{\alpha\beta}^\phi$

$$\text{For } \theta : \frac{2d}{dz} (r^2 \dot{\theta}) = 2r^2 \sin\theta \cos\theta \dot{\phi}^2$$

$$2r^2 \ddot{\theta} + 4r \dot{r} \dot{\theta} = 2r^2 \sin\theta \cos\theta \dot{\phi}^2$$

$$\Rightarrow \Gamma_{r\theta}^\theta = \frac{1}{r}, \quad \Gamma_{\phi\phi}^\theta = \sin\theta \cos\theta$$

$$\text{For } \phi : 2 \frac{d}{dz} (r^2 \sin^2\theta \dot{\phi}) = 0$$

$$\Rightarrow r^2 \sin^2\theta \ddot{\phi} + 2r \sin^2\theta \dot{r} \dot{\phi} + 2r^2 \sin\theta \cos\theta \dot{\theta} \dot{\phi} = 0$$

$$\Rightarrow \Gamma_{r\phi}^\phi = \frac{1}{r}, \quad \Gamma_{\theta\phi}^\phi = \cot\theta$$

$$\text{Now we note } \xi_{(1)}^\alpha = (0, 0, \sin\phi, \cot\theta \cos\phi)$$

$$\xi_{(2)}^\alpha = (0, 0, -\cos\phi, \cot\theta \sin\phi)$$

$$\begin{aligned}\nabla_\phi \xi_{(1)}^\theta &= \partial_\phi \xi_{(1)}^\theta + \Gamma_{\phi\phi}^\theta \xi_{(1)}^\phi \\ &= \sin\phi + \cancel{\sin\theta} \cos\theta \cot\theta \sin\phi \\ &= \sin\phi [1 + \sin\theta \cos\theta]\end{aligned}$$

$$\begin{aligned}\nabla_\phi \xi_{(1)}^\phi &= -\cot\theta \sin\phi + \Gamma_{\phi\phi}^\phi \xi_{(1)}^\eta + \Gamma_{\theta\phi}^\phi \xi_{(1)}^\theta \\ &= -\cot\theta \sin\phi + \cot\theta \sin\phi = 0\end{aligned}$$

$$\nabla_\theta \xi_{(1)}^\theta = 0 + \Gamma_{\theta\eta}^\theta \xi_{(1)}^\eta = 0$$

$$\begin{aligned}\nabla_\theta \xi_{(1)}^\phi &= -\operatorname{cosec}^2\theta \cos\phi + \Gamma_{\theta\phi}^\phi (\cot\theta \cos\phi) \\ &= -\operatorname{cosec}^2\theta \cos\phi + \cot\theta \cos\phi\end{aligned}$$

$$\nabla_\theta \xi_{(1)}^\phi = -\cos\phi$$

Since $g_{\alpha\beta}$ is diagonal, and none of the terms of type $\delta_\beta^\alpha \nabla_\beta \xi_{(1)}^\alpha$ are nonzero, we conclude that

$$g_{\theta\theta} \nabla_\phi \xi_{(1)}^\theta + g_{\phi\phi} \nabla_\theta \xi_{(1)}^\phi = 0$$

for $\xi_{(1)}$ to be Killing vector

$$\sin\phi [1 + \sin\theta \cos\theta] + \sin^2\theta [-\cos\phi]$$

Similarly,

$$\nabla_\theta \xi_{(2)}^\theta = 0 + \Gamma_{\theta\eta}^\theta \xi_{(2)}^\eta = 0$$

$$\begin{aligned}\nabla_\phi \xi_{(2)}^\phi &= \cot\theta \cos\phi + \cot\theta \cos\phi (-1) \\ &= 0\end{aligned}$$

$$\begin{aligned}\nabla_\theta \xi_{(2)}^\phi &= -\operatorname{cosec}^2\theta \sin\phi + \Gamma_{\theta\phi}^\phi \xi_{(2)}^\phi \\ &= -\operatorname{cosec}^2\theta \sin\phi + \cot^2\theta \sin\phi \\ &= -\sin\phi\end{aligned}$$

$$\begin{aligned}
 \nabla_\phi \xi_{(2)}^\theta &= \sin\phi + \Gamma_{\phi\phi}^\theta (\cot\theta \sin\phi) \\
 &= \sin\phi - \sin\theta \cos\theta \cot\theta \sin\phi \\
 &= \sin\phi [1 - \cos^2\theta]
 \end{aligned}$$

$$\begin{aligned}
 &g_{\phi\phi}(-\sin\phi) + g_{\theta\theta}(1 - \cos^2\theta) \\
 &= -\sin^2\theta \sin\phi + \sin\phi [1 - \cos^2\theta] \\
 &= \sin\phi [-\sin^2\theta - \cos^2\theta + 1] = 0
 \end{aligned}$$

Hence $\xi_{(2)}^\theta$ is a Killing vector.

Problem 4

Given metric $g_{\alpha\beta}$, vector potential A_α .

Equations of motion $\rightarrow u^\beta \nabla_\beta u_\alpha = e F_{\alpha\beta} u^\beta$

and $F_{\alpha\beta} = \nabla_\alpha A_\beta - \nabla_\beta A_\alpha$, $d_\xi g_{\alpha\beta} = d_\xi A_\alpha = 0$

To Prove: $(u_\alpha + e A_\alpha) \xi^\alpha$ is constant along the worldline. i.e. $u^\beta \nabla_\beta [\xi^\alpha (u_\alpha + e A_\alpha)] = 0$

The Lie Derivative conditions give,

$$\xi^\nu \nabla_\nu g_{\alpha\beta} = 0 \quad \text{and} \quad \xi^\nu \nabla_\nu A_\alpha + A_\alpha \nabla_\nu \xi^\nu = 0$$

Consider contracting EOMs with ξ^α

$$\begin{aligned} \xi^\alpha u^\beta \nabla_\beta u_\alpha &= e \xi^\alpha F_{\alpha\beta} u^\beta \\ &= e [\xi^\alpha \nabla_\alpha A_\beta - \xi^\alpha \nabla_\beta A_\alpha] u^\beta \end{aligned}$$

$$\xi^\alpha u^\beta [\nabla_\beta (u_\alpha + e A_\alpha)] = e \xi^\alpha (\nabla_\alpha A_\beta) u^\beta$$

$$\begin{aligned} \therefore u^\beta \nabla_\beta [\xi^\alpha (u_\alpha + e A_\alpha)] - u^\beta (u_\alpha + e A_\alpha) \nabla_\beta \xi^\alpha \\ = e \xi^\alpha (\nabla_\alpha A_\beta) u^\beta \end{aligned}$$

$$\therefore u^\beta \nabla_\beta [\xi^\alpha (u_\alpha + e A_\alpha)] = \cancel{u^\beta u_\alpha \nabla_\beta \xi^\alpha} + e u^\beta [A_\alpha \nabla_\beta \xi^\alpha + \xi^\alpha \nabla_\alpha A_\beta]$$

$u^\beta u_\alpha$ is symmetric while $\nabla_\beta \xi^\alpha$ is antisymmetric. *Due to Killing eqn.*

Hence, $u^\beta \nabla_\beta [\xi^\alpha (u_\alpha + e A_\alpha)] = 0$

i.e. $\xi^\alpha (u_\alpha + e A_\alpha)$ is constant on the worldline.

$$\left[\frac{d[\xi^\alpha (u_\alpha + e A_\alpha)]}{d\lambda} = 0 \right]$$

Problem 5

$$\tilde{E} = -u_\alpha \tilde{Y}^\alpha \quad ; \quad \tilde{L} = u_\alpha \tilde{Y}^\alpha$$

$$\text{Also } u^\alpha = \gamma \left(\tilde{Y}^\alpha_{(t)} + \Omega \tilde{Y}^\alpha_{(\phi)} \right)$$

$$\tilde{Y}^\alpha_{(t)} u_\alpha = \gamma \left[g_{\alpha\beta} \tilde{Y}^\beta_{(t)} \tilde{Y}^\alpha_{(t)} + \Omega g_{\alpha\beta} \tilde{Y}^\alpha_{(t)} \tilde{Y}^\beta_{(\phi)} \right]$$

$$\tilde{E} = -\gamma \left[g_{tt} + \Omega g_{t\phi} \right]$$

$$\text{Also } \tilde{L} = \gamma \left[g_{t\phi} + \Omega g_{\phi\phi} \right]$$

$$\text{But } u^\alpha u_\alpha = -1 \Rightarrow \gamma^2 \left(\tilde{Y}^\alpha_{(t)} + \Omega \tilde{Y}^\alpha_{(\phi)} \right) \left(g_{\alpha\beta} \tilde{Y}^\beta_{(t)} + \Omega g_{\alpha\beta} \tilde{Y}^\beta_{(\phi)} \right) = -1$$

$$\Rightarrow \gamma^2 \left[g_{tt} + \Omega^2 g_{\phi\phi} + g_{t\phi} \Omega + g_{\phi t} \Omega \right] = -1$$

$$\text{Hence we get } -\tilde{E} + \Omega \tilde{L} = \gamma \left[g_{tt} + \Omega g_{t\phi} + \Omega^2 g_{\phi\phi} \right]$$

$$\therefore -\tilde{E} + \Omega \tilde{L} = -1/\gamma \Rightarrow -\frac{1}{\gamma^2} \delta\gamma = -\delta\tilde{E} + \Omega \delta\tilde{L}$$

$$\Rightarrow \frac{\delta\gamma}{\gamma} = -\gamma(\delta\tilde{E} + \Omega \delta\tilde{L})$$

$$\text{Now } \delta u^\alpha = \delta\gamma \frac{u^\alpha}{\gamma}$$

$$\delta u^\alpha = -\gamma(\delta\tilde{E} + \Omega \delta\tilde{L}) u^\alpha$$

$$\therefore u_\alpha \delta u^\alpha = \gamma(-\delta\tilde{E} + \Omega \delta\tilde{L})$$

$$\text{But, } \delta[u_\alpha u^\alpha] = 0 = 2u_\alpha \delta u^\alpha$$

$$\Rightarrow -\delta\tilde{E} + \Omega \delta\tilde{L} = 0$$

$$\Rightarrow \delta\tilde{E} = \Omega \delta\tilde{L}$$