

If the functional integration over p is Gaussian, we can exactly integrate out P . For example, $H = \frac{P^2}{2m} + V$

$$\text{so } \mathcal{L} = \dot{p}\dot{x} - H = (\dot{p}\dot{x} - \frac{P^2}{2m} - V(x))$$

$$iG = \int Dp Dx \exp \left[\frac{i}{\hbar} \int_{t_0}^t \mathcal{L} dt \right]$$

$$= \int Dp Dx \exp \left[\frac{i}{\hbar} \int_{t_0}^t \right]$$

$$iG = \int Dp Dx \exp \left[\frac{i}{\hbar} \sum_{t=1}^N (P_t \cdot \dot{x}_t - \frac{P_t^2}{2m} - V(x)) \Delta t \right]$$

$$\text{let } \vec{P} = \begin{bmatrix} P_1 \\ \vdots \\ P_N \end{bmatrix}, \quad \vec{x} = \begin{bmatrix} \dot{x}_1 \\ \vdots \\ \dot{x}_N \end{bmatrix}$$

$$\therefore iG = \int Dx \cdot e^{\frac{i}{\hbar} \sum (-V(x)) \Delta t} \cdot Dp \cdot \exp \left[\frac{i}{2m\hbar} (-\vec{P}^T \vec{P} + 2m \vec{P}^T \cdot \vec{x}) \right]$$

we have an useful formula: Gaussian integral:

$$\int_{n=1}^N dx_n e^{-\frac{1}{2} \vec{x}^T A \vec{x} - \vec{x}^T \vec{y}} = (2\pi)^{\frac{N}{2}} (\det A)^{-\frac{1}{2}} e^{\frac{1}{2} \vec{y}^T A^{-1} \vec{y}}$$

where, x, y are real vectors, A is real symmetric matrix

let.

$$I = \int Dp \exp \left[\frac{i}{2m\hbar} (-\vec{P}^T \vec{P} + 2m \vec{P}^T \cdot \vec{x}) \right]$$

$$= \left(\frac{1}{2\pi\hbar} \right)^N \cdot \int_{t=1}^N dp \cdot \exp \left[-\frac{1}{2} \vec{P}^T A \vec{P} - \vec{P}^T \vec{x}' \right]$$

where $A = \frac{i\omega}{m\hbar} \hat{I}_{N \times N}$, $\dot{x}' = -\frac{i\omega}{\hbar} \dot{x}$, \hat{I} is a $N \times N$ identity matrix.

∴ So we can get,

$$\left\{ \begin{array}{l} (\text{Det } A)^{\frac{1}{2}} = \left(\frac{i\omega}{m\hbar}\right)^{-\frac{N}{2}}, \\ A^{-1} = \frac{m\hbar}{i\omega} \hat{I}_{N \times N} \end{array} \right.$$

So,

$$\begin{aligned} \frac{1}{2} \dot{x}'^T A^{-1} \dot{x}' &= \frac{1}{2} \cdot \frac{m\hbar}{i\omega} \cdot \left(-\frac{i\omega}{\hbar}\right)^2 \sum_{\ell=1}^N \dot{x}_\ell^2 \\ &= \frac{i}{\hbar} \sum_{\ell=1}^N \frac{m}{2\pi\hbar} \left(\frac{x_\ell - x_{\ell-1}}{\Delta t} \right)^2 \Delta t \end{aligned}$$

$$\begin{aligned} \text{So: } I &= \left(\frac{1}{2\pi\hbar}\right)^N \cdot (2\pi)^{\frac{N}{2}} \cdot \left(\frac{i\omega}{m\hbar}\right)^{-\frac{N}{2}} \cdot e^{\frac{i}{\hbar} \sum_{\ell=1}^N \frac{m}{2} \left(\frac{x_\ell - x_{\ell-1}}{\Delta t} \right)^2 \Delta t} \\ &= \left(\frac{m}{i2\pi\hbar\Delta t}\right)^{\frac{N}{2}} e^{\frac{i}{\hbar} \sum_{\ell=1}^N \frac{m}{2} \left(\frac{x_\ell - x_{\ell-1}}{\Delta t} \right)^2 \Delta t} \end{aligned}$$

So we can get the integration without P .

$$\begin{aligned} IG &= \left(\frac{m}{i2\pi\hbar\Delta t}\right)^{\frac{N}{2}} \int Dx \exp \left[\frac{i}{\hbar} \sum_{\ell=1}^N \left(\frac{m}{2} \left(\frac{x_\ell - x_{\ell-1}}{\Delta t} \right)^2 - V(x) \right) \Delta t \right] \\ &= \left(\frac{m}{i2\pi\hbar\Delta t}\right)^{\frac{N}{2}} \int Dx \exp \left[\frac{i}{\hbar} \int_{t_0}^t d\tau \cdot L(x, \dot{x}) \right]. \end{aligned}$$

$$S \propto \int Dx \cdot e^{\frac{i}{\hbar} S[x(t)]}$$

where, $S[x(t)] = \int_{t_0}^t dt \cdot L(x, \dot{x})$ is action,

and $L(x, \dot{x}) = \frac{m\dot{x}^2}{2} - V(x)$ is Lagrangean

From the Path-integral in real space-time, we can get some information about Physics Picture:

(1) Each path is weighted with a. (11) phase factor $e^{\frac{i}{\hbar} S}$
The quantum interference between different paths.

(2). Since $\hbar \sim 10^{-34} \text{ J}\cdot\text{s}$, any "small change" in S (we change S to $S + \delta S$), will drastically lead to quantum destructive interference. So the only the paths that satisfy $\delta S = 0$ make dominant contributions to the path-integral.

(3). Remarkably, $\delta S = 0$ is exactly Hamilton's Principle in classical mechanics. So the classical paths ($\delta S = 0$) dominate the path integral in the limit $\hbar \rightarrow 0$. In other words, in classical mechanics, as $\hbar \rightarrow 0$, it neglects the contribution of the integral over all other paths near the path ~~with~~ with $\delta S = 0$. So we can get the conclusion:

Quantum system $\xrightarrow{\hbar \rightarrow 0}$ classical system.

Example: Computing the path-integral of free particles:

Free particles' Hamiltonian, $H = \frac{P^2}{2m}$.

$$\text{So } iG = \left(\frac{m}{2\pi\hbar t}\right)^{\frac{N}{2}} \int Dx \exp\left[\frac{i}{\hbar t} \sum_{l=1}^N \frac{m(x_l - x_{l+1})^2}{2} \right] dt.$$

we let:

$$I = \int Dx \exp\left[\frac{im}{2\hbar t} \left\{ \sum_{l=1}^N (x_l^2 - 2x_l x_{l+1} + x_{l+1}^2) \right\}\right]$$

In other to. use Gaussian integral:

$$\int_{n=1}^N dx_n e^{-\frac{1}{2} x^T A x - x^T y} = (2\pi)^{\frac{N}{2}} (\det A)^{-\frac{1}{2}} e^{\frac{1}{2} y^T A^{-1} y}.$$

We should rewrite the form of $\exp\left[\frac{im}{2\hbar t} \sum (x_l^2 - 2x_l x_{l+1} + x_{l+1}^2)\right]$

So we let:

$$A = \frac{-im}{\hbar t} \begin{bmatrix} 2 & -1 & & \\ -1 & 2 & -1 & \\ & -1 & 2 & -1 \\ & & \ddots & \ddots \end{bmatrix}_{(N-1) \times (N-1)}, \quad x = \begin{bmatrix} x_{N-1} \\ \vdots \\ x_1 \end{bmatrix}$$

$$y = \frac{im}{\hbar t} \begin{bmatrix} x_N \\ 0 \\ \vdots \\ 0 \\ x_0 \end{bmatrix} \text{ all zero}$$

$$- \exp\left[\frac{im}{2\hbar t} \sum_{l=1}^N (x_l^2 - 2x_l x_{l+1} + x_{l+1}^2)\right] = \exp\left(-\frac{1}{2} x^T A x - x^T y\right) e^{\frac{im}{2\hbar t} (x_N^2 + x_0^2)}$$

because of $Dx = \prod_{t=1}^N \lambda_t$ without x_N and x_0 ,

$$\text{So } |G| = \left(\frac{m}{i2\pi\hbar\omega}\right)^{\frac{N}{2}} \cdot e^{\frac{im}{2\hbar\omega}(x_N^2 + x_0^2)} \int Dx \cdot e^{-\frac{i}{\hbar\omega}x^T A x - x^T y}$$

$$\text{So we can compute } |\text{Det} A| = N \cdot \left(\frac{-im}{\hbar\omega}\right)^N$$

So we can get:

$$|G| = \left(\frac{m}{i2\pi\hbar\omega}\right)^{\frac{N}{2}} \cdot e^{\frac{im}{2\hbar\omega}(x_N^2 + x_0^2)} \cdot (2\pi)^{\frac{N}{2}} \cdot N^{-\frac{1}{2}} \cdot \left(\frac{-im}{\hbar\omega}\right)^{\frac{N}{2}} \cdot e^{\frac{1}{2}y^T A^T y}.$$

Although A^T is difficult to compute, we notice that

$$y = \begin{bmatrix} x_N \\ \vdots \\ 0 \\ x_0 \end{bmatrix} \text{ only have two non-zero elements,}$$

which locate in the first row and the last row respectively.

So we only need to calculate the first and last columns of matrix A^T , denoted A_1^T and A_{N-1}^T , respectively

$$A_1^T = \frac{i\hbar\omega}{m} \begin{bmatrix} \frac{N-1}{N} \\ \frac{N-2}{N} \\ \vdots \\ \frac{1}{N} \end{bmatrix}, \quad A_1^T = \frac{i\hbar\omega}{mN} \begin{bmatrix} N-1 \\ N-2 \\ \vdots \\ 1 \end{bmatrix}, \quad A_{N-1}^T = \frac{i\hbar\omega}{mN} \begin{bmatrix} 1 \\ 2 \\ \vdots \\ N-1 \end{bmatrix}$$

$$\therefore \frac{1}{2}y^T A^T y = \frac{1}{2}y^T [x_N A_1^T + x_0 A_{N-1}^T]$$

$$= \frac{im}{2\hbar\omega} \left[-(x_N^2 + x_0^2) + \frac{(x_N - x_0)^2}{N} \right]$$

So we can get the complete integral \mathcal{F} :

$$\begin{aligned} \mathcal{F} &= \int iG = \left(\frac{m}{i2\pi\hbar\Delta t} \right)^{\frac{1}{2}} (2\pi)^{\frac{N}{2}} N^{-\frac{1}{2}} \left(\frac{-im}{\hbar\Delta t} \right)^{\frac{N+1}{2}} e^{\frac{i}{\hbar} \frac{m(x_N-x_0)}{2N\Delta t}} \\ &= \left(\frac{m}{i2\pi\hbar\Delta t N} \right)^{\frac{1}{2}} e^{\frac{i}{\hbar} \frac{m(x_N-x_0)^2}{2N\Delta t}}. \end{aligned}$$

where $t-t_0=N\Delta t$, $x=x_N$

so:

$$iG = \left[\frac{m}{i2\pi\hbar(t-t_0)} \right]^{\frac{1}{2}} \cdot e^{\frac{i}{\hbar} \frac{m(x-x_0)^2}{2(t-t_0)}}.$$