

If the functional integration over p is Gaussian, we can exactly integrate out P . For example, $H = \frac{P^2}{2m} + V$

$$\text{so } \mathcal{L} = \dot{p}\dot{x} - H = \dot{p}\dot{x} - \frac{p^2}{2m} - V(x)$$

$$iG = \int \mathcal{D}p \mathcal{D}x \exp \left[\frac{i}{\hbar} \int_{t_0}^t \mathcal{L} dt \right]$$

$$= \int \mathcal{D}p \mathcal{D}x \exp \left[\frac{i}{\hbar} \int_{t_0}^t \mathcal{L} dt \right]$$

$$iG = \int \mathcal{D}p \mathcal{D}x \exp \left[\frac{i}{\hbar} \sum_{t=1}^N (P_t \cdot \dot{x}_t - \frac{P_t^2}{2m} - V(x)) \Delta t \right]$$

$$\text{let } \begin{matrix} \vec{P}_t \\ \vec{P}_t \\ \vdots \\ \vec{P}_1 \end{matrix} = \begin{bmatrix} P_t \\ \vdots \\ P_1 \end{bmatrix}, \quad \begin{matrix} \dot{x}_t \\ \vdots \\ \dot{x}_1 \end{matrix} = \begin{bmatrix} \dot{x}_t \\ \vdots \\ \dot{x}_1 \end{bmatrix}$$

$$\therefore iG = \int \mathcal{D}x \cdot e^{\frac{i}{\hbar} \sum (-V(x)) \Delta t} \cdot \mathcal{D}p \cdot \exp \left[\frac{i}{2m\hbar} (-P^T P + 2m P^T \cdot \dot{x}) \right]$$

we have an useful formula: Gaussian integral:

$$\int \prod_{n=1}^N dx_n e^{-\frac{1}{2} x^T A x - x^T y} = (2\pi)^{\frac{N}{2}} (\text{Det} A)^{-\frac{1}{2}} e^{\frac{1}{2} y^T A^{-1} y}$$

where, x, y are real vectors, A is real symmetric matrix

let.

$$\begin{aligned} \mathcal{I} &= \int \mathcal{D}p \exp \left[\frac{i}{2m\hbar} (-P^T P + 2m P^T \cdot \dot{x}) \right] \\ &= \left(\frac{1}{2\pi\hbar} \right)^N \cdot \int \prod_{t=1}^N dP_t \cdot \exp \left[-\frac{1}{2} P^T A P - P^T \cdot \dot{x}' \right] \end{aligned}$$

where $A = \frac{i\Delta t}{m\hbar} \hat{I}_{N \times N}$ $\dot{x}' = -\frac{i\Delta t}{\hbar} \dot{x}$, \hat{I} is a $N \times N$ identity matrix.

So we can get,

$$\begin{cases} (\text{Det } A)^{-\frac{1}{2}} = \left(\frac{i\Delta t}{m\hbar}\right)^{-\frac{N}{2}}, \\ A^{-1} = \frac{m\hbar}{i\Delta t} \hat{I}_{N \times N} \end{cases}$$

So,

$$\begin{aligned} \frac{1}{2} \dot{x}'^T A^{-1} \dot{x}' &= \frac{1}{2} \cdot \frac{m\hbar}{i\Delta t} \cdot \left(-\frac{i\Delta t}{\hbar}\right)^2 \sum_{l=1}^N \dot{x}_l^2 \\ &= \frac{i}{\hbar} \sum_{l=1}^N \frac{m}{2\Delta t} (\frac{x_l - x_{l-1}}{\Delta t})^2 \Delta t \end{aligned}$$

$$\begin{aligned} \text{So: } I &= \left(\frac{1}{2\pi\hbar}\right)^N \cdot (2\pi)^{\frac{N}{2}} \cdot \left(\frac{i\Delta t}{m\hbar}\right)^{-\frac{N}{2}} \cdot e^{\frac{i}{\hbar} \sum_{l=1}^N \frac{m}{2} (\frac{x_l - x_{l-1}}{\Delta t})^2 \Delta t} \\ &= \left(\frac{m}{i2\pi\hbar\Delta t}\right)^{\frac{N}{2}} e^{\frac{i}{\hbar} \sum_{l=1}^N \frac{m}{2} (\frac{x_l - x_{l-1}}{\Delta t})^2 \Delta t} \end{aligned}$$

So we can get the integration without p .

$$\begin{aligned} iG &= \left(\frac{m}{i2\pi\hbar\Delta t}\right)^{\frac{N}{2}} \int D_x \exp \left[\frac{i}{\hbar} \sum_{l=1}^N \left(\frac{m}{2} (\frac{x_l - x_{l-1}}{\Delta t})^2 - V(x) \right) \Delta t \right] \\ &= \left(\frac{m}{i2\pi\hbar\Delta t}\right)^{\frac{N}{2}} \int D_x \exp \left[\frac{i}{\hbar} \int_{t_0}^t dt \cdot L(x, \dot{x}) \right] \end{aligned}$$

$$\text{So } i\hbar \propto \int \mathcal{D}x \cdot e^{\frac{i}{\hbar} S[x(t)]},$$

where, $S[x(t)] = \int_{t_0}^t dt \cdot \mathcal{L}(x, \dot{x})$ is action,

and $\mathcal{L}(x, \dot{x}) = \frac{m\dot{x}^2}{2} - V(x)$ is Lagrangean

From the Path-integral in real space-time, we can get some information about Physics Picture:

(1) Each path is weighted with a unit phase factor $e^{\frac{i}{\hbar} S}$
The Quantum interference between different paths

(2). Since $\hbar \sim 10^{-34} \text{ J}\cdot\text{s}$, any "small change" in S (we change S to $S + \delta S$), ~~it~~ will drastically lead to quantum destructive interference. So ~~the~~ only the paths that ~~se~~ satisfy $\delta S = 0$ make dominant contributions to the path-integral.

(3). Remarkably, $\delta S = 0$ is exactly Hamilton's Principle in classical mechanics. So the classical paths ($\delta S = 0$) dominate the path integral in the limit $\hbar \rightarrow 0$. In other words, in classical mechanics, ~~it~~ as $\hbar \rightarrow 0$, it neglects the contribution of the integral over all other paths near the path ~~with~~ ~~those~~ with $\delta S = 0$. So we can get the conclusion:

Quantum system $\xrightarrow{\hbar \rightarrow 0}$ classical system.

Example: Computing the path-integral of free particles:

Free particles' Hamiltonian, $H = \frac{p^2}{2m}$.

$$\text{So } iG = \left(\frac{m}{i2\pi\hbar\Delta t}\right)^{\frac{N}{2}} \int \mathcal{D}x \exp \left[\frac{i}{\hbar} \sum_{t=1}^N \frac{m}{2} \frac{(x_t - x_{t-1})^2}{\Delta t} \right] \Delta t.$$

we let:

$$\mathcal{I} = \int \mathcal{D}x \exp \left[\frac{im}{2\hbar\Delta t} \sum_{t=1}^N (x_t^2 - 2x_t x_{t-1} + x_{t-1}^2) \right]$$

In order to use Gaussian integral:

$$\int \prod_{n=1}^N dx_n e^{-\frac{1}{2} x^T A x - x^T y} = (2\pi)^{\frac{N}{2}} (\text{Det } A)^{-\frac{1}{2}} e^{\frac{1}{2} y^T A^{-1} y}.$$

we should rewrite the form of $\exp \left[\frac{im}{2\hbar\Delta t} \sum (x_t^2 - 2x_t x_{t-1} + x_{t-1}^2) \right]$

So we let:

$$A = \frac{-im}{\hbar\Delta t} \begin{bmatrix} 2 & -1 & & \\ -1 & 2 & -1 & \\ & -1 & 2 & -1 \\ & & & \ddots \end{bmatrix}_{(N-1) \times (N-1)}, \quad x = \begin{bmatrix} x_{N-1} \\ \vdots \\ x_1 \end{bmatrix}$$

$$y = \frac{im}{\hbar\Delta t} \begin{bmatrix} x_N \\ 0 \\ \vdots \\ \text{all zero} \\ 0 \\ x_0 \end{bmatrix}$$

$$\exp \left[\frac{im}{2\hbar\Delta t} \sum_{t=1}^N (x_t^2 - 2x_t x_{t-1} + x_{t-1}^2) \right] = \exp \left[-\frac{1}{2} x^T A x - x^T y \right] e^{\frac{im}{2\hbar\Delta t} (x_N^2 + x_0^2)}$$

because of $D_x = \prod_{t=1}^{N-1} dx_t$ without x_N and x_0 ,

$$\text{So } iG = \left(\frac{m}{i2\pi\hbar\Delta t}\right)^{\frac{N}{2}} \cdot e^{\frac{im}{2\hbar\Delta t}(x_N^2 + x_0^2)} \int D_x \cdot e^{-\frac{1}{2}x^T A x - x^T y}$$

so we can compute $|\text{Det } A| = N \cdot \left(\frac{-im}{\hbar\Delta t}\right)^N$

So we can get:

$$iG = \left(\frac{m}{i2\pi\hbar\Delta t}\right)^{\frac{N}{2}} \cdot e^{\frac{im}{2\hbar\Delta t}(x_N^2 + x_0^2)} \cdot (2\pi)^{\frac{N}{2}} \cdot N^{-\frac{1}{2}} \cdot \left(\frac{-im}{\hbar\Delta t}\right)^{-\frac{N}{2}} \cdot e^{\frac{1}{2}y^T A^{-1} y}$$

Although A^{-1} is difficult to compute, we notice that

$y = \begin{bmatrix} x_N \\ 0 \\ \vdots \\ 0 \\ x_0 \end{bmatrix}$ only have two non-zero elements,

which locate in the first row and the last row respectively

So we only need to calculate the first and last columns of matrix A^{-1} , denoted A_1^{-1} and A_{N-1}^{-1} , respectively

$$A_1^{-1} = \frac{i\hbar\Delta t}{m} \begin{bmatrix} \frac{N-1}{N} \\ \frac{N-2}{N} \\ \vdots \\ \frac{1}{N} \end{bmatrix}, \quad A_{N-1}^{-1} = \frac{i\hbar\Delta t}{mN} \begin{bmatrix} 1 \\ 2 \\ \vdots \\ N-1 \end{bmatrix}$$

$$\begin{aligned} \therefore \frac{1}{2}y^T A^{-1} y &= \frac{1}{2}y^T [x_N A_1^{-1} + x_0 A_{N-1}^{-1}] \\ &= \frac{im}{2\hbar\Delta t} \left[-(x_N^2 + x_0^2) + \frac{(x_N - x_0)^2}{N} \right] \end{aligned}$$

So we can get the complete integral \mathbb{I} :

$$\begin{aligned} \mathbb{I} &\equiv \int iG = \left(\frac{m}{i2\pi\hbar\Delta t}\right)^{\frac{N}{2}} (2\pi)^{\frac{N}{2}} N^{-\frac{1}{2}} \left(\frac{-im}{\hbar\Delta t}\right)^{\frac{N-1}{2}} e^{\frac{i}{\hbar} \cdot \frac{m(x_N-x_0)^2}{2N\Delta t}} \\ &= \left(\frac{m}{i2\pi\hbar\Delta tN}\right)^{\frac{1}{2}} e^{\frac{i}{\hbar} \cdot \frac{m(x_N-x_0)^2}{2N\Delta t}} \end{aligned}$$

where $t-t_0 = N\Delta t$, $x = x_N$

So:

$$iG = \left[\frac{m}{i2\pi\hbar(t-t_0)}\right]^{\frac{1}{2}} \cdot e^{\frac{i}{\hbar} \cdot \frac{m(x-x_0)^2}{2(t-t_0)}}$$