

# 1 The theta term and Axion Electrodynamics

In relativistic notation, the Maxwell action for electromagnetism takes a wonderfully compact form:

$$S_{Max} = \int d^4x -\frac{1}{4}F^{\mu\nu}F_{\mu\nu} = \int d^4x \frac{1}{2}(\mathbf{E}^2 - \mathbf{B}^2) \quad (1)$$

Here:

$$\left\{ \begin{array}{l} F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu \\ = \begin{pmatrix} 0 & -E_x & -E_y & -E_z \\ E_x & 0 & -B_z & B_y \\ E_y & B_z & 0 & -B_x \\ E_z & -B_y & B_x & 0 \end{pmatrix} \\ E_i = F_{0i} \\ F_{ij} = -\epsilon_{ijk}B_k \end{array} \right. \quad (2)$$

One reason that the Maxwell action is so simple is that there is very little else we can write down that is both gauge invariant and Lorentz invariant. There is, however one term that we can add to the Maxwell action is also both gauge invariant and Lorentz invariant, which we call it theta term:

$$S_\theta = \frac{\theta e^2}{4\pi^2} \int d^4x \frac{1}{4} \star F^{\mu\nu} F_{\mu\nu} = -\frac{\theta e^2}{4\pi^2} \int d^4x \mathbf{E} \cdot \mathbf{B} \quad (3)$$

which  $\star F^{\mu\nu}$  is the dual tensor:

$$\begin{aligned} \star F^{\mu\nu} &= \frac{1}{2} \epsilon^{\mu\nu\rho\sigma} F_{\rho\sigma} \\ &= \begin{pmatrix} 0 & -B_x & -B_y & -B_z \\ B_x & 0 & E_z & -E_y \\ B_y & -E_z & 0 & E_x \\ B_z & E_y & -E_x & 0 \end{pmatrix} \end{aligned} \quad (4)$$

and  $\theta$  is a constant parameter.

However, the theta term is simple to check that it can be written as a total derivative:

$$\begin{aligned} \int d^4x \frac{1}{4} \star F^{\mu\nu} F_{\mu\nu} &= \frac{1}{2} \int d^4x \cdot \frac{1}{4} \epsilon^{\mu\nu\rho\sigma} F_{\rho\sigma} F_{\mu\nu} \\ &= \frac{1}{8} \int d^4x \cdot \epsilon^{\mu\nu\rho\sigma} (\partial_\rho A_\sigma - \partial_\sigma A_\rho) (\partial_\mu A_\nu - \partial_\nu A_\mu) \\ &= \frac{1}{2} \int d^4x \cdot \epsilon^{\mu\nu\rho\sigma} (\partial_\mu A_\nu \partial_\rho A_\sigma) \end{aligned} \quad (5)$$

Now let's proof a useful equation:

$$\epsilon^{\mu\nu\rho\sigma} A_\nu \partial_\mu \partial_\rho A_\sigma = 0 \quad (6)$$

Because:

$$\begin{aligned}
 \epsilon^{\mu\nu\rho\sigma} A_\nu \partial_\mu \partial_\rho A_\sigma &= \epsilon^{\rho\nu\mu\sigma} A_\nu \partial_\rho \partial_\mu A_\sigma \\
 &= \epsilon^{\rho\nu\mu\sigma} A_\nu \partial_\mu \partial_\rho A_\sigma \\
 &= -\epsilon^{\mu\nu\rho\sigma} A_\nu \partial_\mu \partial_\rho A_\sigma
 \end{aligned} \tag{7}$$

So we can get:

$$2\epsilon^{\mu\nu\rho\sigma} A_\nu \partial_\mu \partial_\rho A_\sigma = 0 \Rightarrow \epsilon^{\mu\nu\rho\sigma} A_\nu \partial_\mu \partial_\rho A_\sigma = 0 \tag{8}$$

So the quadratic part of theta term become:

$$\begin{aligned}
 \int d^4x \frac{1}{4} \star F^{\mu\nu} F_{\mu\nu} &= \frac{1}{2} \int d^4x \epsilon^{\mu\nu\rho\sigma} (\partial_\mu A_\nu \partial_\rho A_\sigma + A_\nu \partial_\mu \partial_\rho A_\sigma) \\
 &= \frac{1}{2} \int d^4x \epsilon^{\mu\nu\rho\sigma} \partial_\mu (A_\nu \partial_\rho A_\sigma)
 \end{aligned} \tag{9}$$

The total derivative form of theta term is:

$$S_\theta = \frac{\theta e^2}{8\pi^2} \int d^4x \partial_\mu (\epsilon^{\mu\nu\rho\sigma} A_\nu \partial_\rho A_\sigma) \tag{10}$$

We say that the theta term is topological. It depends only on boundary information. However, when deriving using the principle of least action to derive the field's equation of motion, the values of field are fixed on the infinite boundary. Therefore, the upshot is that the theta term does not change the equations of motion and, it would seem, can have little effect on the physics.

But under some situations that involve subtle interplay between quantum mechanics and topology, there are a number of interesting phenomena of physics which are led by theta term.

Also, we can look at the situations where  $\theta$  affects the dynamics classically. This occurs when  $\theta$  is not constant, but instead varies in space and time.

$$\theta = \theta(x, t) \tag{11}$$

So we can write down the axion electrodynamics action:

$$S = \int d^4x \left( -\frac{1}{4} F^{\mu\nu} F_{\mu\nu} + \frac{e^2}{16\pi^2} \theta(x, t) \star F^{\mu\nu} F_{\mu\nu} \right) \tag{12}$$

We use Euler-Lagrange equations:

$$\frac{\partial \mathcal{L}}{\partial A_\mu} - \partial_\nu \left( \frac{\partial \mathcal{L}}{\partial (\partial_\nu A_\mu)} \right) = 0 \tag{13}$$

to get the equations of axion electrodynamics

Because  $\mathcal{L}$  doesn't depend on  $A_\mu$ , so the equations of motion are:

$$\frac{\partial \mathcal{L}}{\partial A_\mu} = 0 \Rightarrow \partial_\nu \left[ \frac{\partial \mathcal{L}}{\partial (\partial_\nu A_\mu)} \right] = 0 \tag{14}$$

and the  $\mathcal{L}$  of axion electrodynamics is ( $\alpha = \frac{e^2}{4\pi}$ ):

$$\begin{aligned}\mathcal{L} &= \mathcal{L}_{Max} + \mathcal{L}_\theta \\ &= -\frac{1}{4}F^{\mu\nu}F_{\mu\nu} + \frac{\alpha}{4\pi}\theta(\mathbf{x}, t)\star F^{\mu\nu}F_{\mu\nu}\end{aligned}\quad (15)$$

First, let's see  $\mathcal{L}_\theta$ :

$$\begin{aligned}\frac{\partial \mathcal{L}}{\partial(\partial_\nu A_\mu)} &= \frac{\partial}{\partial(\partial_\nu A_\mu)} \cdot \frac{\alpha}{4\pi}\theta(\mathbf{x}, t) \cdot \frac{1}{2}\epsilon^{\alpha\beta\rho\sigma}F_{\alpha\beta}F_{\rho\sigma} \\ &= \frac{\alpha}{8\pi}\left\{\left[\frac{\partial}{\partial(\partial_\nu A_\mu)}\theta(\mathbf{x}, t)\right] \cdot \epsilon^{\alpha\beta\rho\sigma}F_{\alpha\beta}F_{\rho\sigma} + \theta(\mathbf{x}, t)\left[\frac{\partial\epsilon^{\alpha\beta\rho\sigma}F_{\alpha\beta}F_{\rho\sigma}}{\partial(\partial_\nu A_\mu)}\right]\right\}\end{aligned}\quad (16)$$

let's see the first term:

$$\epsilon^{\alpha\beta\rho\sigma}F_{\alpha\beta}F_{\rho\sigma}\frac{\partial}{\partial(\partial_\nu A_\mu)}\theta(\mathbf{x}, t) = 0 \quad (17)$$

then let's see the second term:

$$\begin{aligned}\theta(\mathbf{x}, t) \cdot \frac{\partial}{\partial(\partial_\nu A_\mu)} \cdot \epsilon^{\alpha\beta\rho\sigma}F_{\alpha\beta}F_{\rho\sigma} &= 4\theta(\mathbf{x}, t)\left(\frac{\partial}{\partial(\partial_\nu A_\mu)}\partial_\alpha A_\beta\partial_\rho A_\sigma\right)\epsilon^{\alpha\beta\rho\sigma} \\ &= 4\theta(\mathbf{x}, t)(\delta_\alpha^\nu\delta_\beta^\mu\partial_\rho A_\sigma + \delta_\rho^\nu\delta_\sigma^\mu\partial_\alpha A_\beta) \cdot \epsilon^{\alpha\beta\rho\sigma} \\ &= 4\theta(\mathbf{x}, t) \cdot \epsilon^{\nu\mu\rho\sigma}(\partial_\rho A_\sigma - \partial_\sigma A_\rho) \\ &= 8\theta(\mathbf{x}, t)\star F^{\nu\mu}\end{aligned}\quad (18)$$

next, we can get:

$$\partial_\nu[\theta(\mathbf{x}, t)\star F^{\nu\mu}] = \star F^{\nu\mu}\partial_\nu\theta(\mathbf{x}, t) + \theta(\mathbf{x}, t) \cdot \partial_\nu\star F^{\nu\mu} \quad (19)$$

using the Bianchi identities:

$$\partial_\mu\star F^{\mu\nu} = 0 \quad (20)$$

so we can get:

$$\partial_\nu\theta(\mathbf{x}, t)\star F^{\nu\mu} = \star F^{\nu\mu}\partial_\nu\theta(\mathbf{x}, t) \quad (21)$$

Finally, we can get:

$$\partial_\nu\left[\frac{\partial \mathcal{L}_\theta}{\partial(\partial_\nu A_\mu)}\right] = \frac{\alpha}{\pi}\star F^{\nu\mu} \cdot \partial_\nu\theta(\mathbf{x}, t) \quad (22)$$

And the Lagrangian of Maxwell part:

$$\partial_\nu\left[\frac{\partial \mathcal{L}_{Max}}{\partial(\partial_\nu A_\mu)}\right] = -\partial_\nu F^{\nu\mu} \quad (23)$$

So the equations of axion electrodynamics is:

$$\begin{cases} \partial_\nu F^{\nu\mu} = \frac{\alpha}{\pi}\star F^{\nu\mu} \cdot \partial_\nu\theta(\mathbf{x}, t) \\ \partial_\mu\star F^{\mu\nu} = 0 \end{cases} \quad (24)$$

In the end, the deformed Maxwell equations which are the equations of axion electrodynamics are:

$$\begin{cases} \nabla \cdot \mathbf{E} = \frac{\alpha}{\pi} (\nabla \theta) \cdot \mathbf{B} \\ \nabla \times \mathbf{B} = \frac{\partial \mathbf{E}}{\partial t} - \frac{\alpha}{\pi} (\dot{\theta} \mathbf{B} + \nabla \theta \times \mathbf{E}) \\ \nabla \cdot \mathbf{B} = 0 \\ \nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \end{cases} \quad (25)$$

The first equation tells us that in regions of space where  $\theta$  varies, a magnetic field  $\mathbf{B}$  acts like an electric charge density. The second equation tells us that the combination  $(\dot{\theta} \mathbf{B} + \nabla \theta \times \mathbf{E})$  acts like a current density

## 1.1 The periodicity of theta and its values under Parity or Time - Reversal Symmetry

In quantum theory,  $\theta$  is a periodic variable: it lies in the range:

$$\theta \in [0, 2\pi) \quad (26)$$

After imposing appropriate boundary conditions,  $S_\theta$  can only take values of the form:

$$S_\theta = \theta N \quad \text{with } N \in \mathbb{Z} \quad (27)$$

This means that the theta angle contributes to the partition function as

$$e^{iS_\theta} = e^{iN\theta} \quad (28)$$

To show that  $S_\theta$  must take the form, we consider a compact Euclidean spacetime which we take to  $T^4$  and we take each of the circles in the torus to have radii  $R$ .

We consider easy case:

$$\begin{cases} \mathbf{E} = (0, 0, E_z), \quad E_z = F_{03} = \partial_0 A_3 - \partial_3 A_0 \\ \mathbf{B} = (0, 0, B_z), \quad B_z = F_{21} = \partial_2 A_1 - \partial_1 A_2 \end{cases} \quad (29)$$

The integral of  $S_\theta$  is:

$$I = \int_{T^4} d^4x E_z \cdot B_z = \int_{T^4} dx^0 dx^3 \cdot E_z \cdot \int_{T^4} dx^1 dx^2 \cdot B_z \quad (30)$$

The gauge field  $A_\mu$  must be well defined on the underlying torus, which will put restrictions on the allowed values of  $E_z$  and  $B_z$ . So the integral can't take any value.

Now let consider the restrictions of  $A_\mu$  if  $A_\mu$  be defined on torus. When a direction of space, say  $x^1$ , is periodic with radius  $R$ , the physics state in original point ( $x^1 = 0$ ) must be identical to the final physics state:

$$\text{state}[A_1(0)] = \text{state}[A_1(2\pi R)] \quad (31)$$

Because  $B_z = \partial_2 A_1$ ,  $(A_2 = 0) = \text{constant}$ , so  $A_1 = B_z x^2 + C$ . The simplest case is  $A_1 = 0$  but this case is trivial. So we should find a gauge transformations to connect  $A(0)$  and  $A(2\pi R)$ , which will make two different values of  $A_1$  correspond to the same physics state.

$$A_1(2\pi R) = A_1(0) + G \quad (32)$$

let's find the gauge transformation that we are interested:

$$A_1 \rightarrow A_1 + \partial_1 \chi \quad (33)$$

We do not insist that  $\chi(x)$  is single valued. Instead, we require only that  $e^{ie\chi}$  is single-valued, since this is what acts on the wavefunction. So  $\chi$  can have the simple form followed:

$$\chi = \frac{nx^1}{eR} \quad \text{with } n \in \mathbb{Z} \quad (34)$$

Because in the same point  $(2\pi NR + X)$ ,  $e^{ie\chi}$  is the same state:

There are sometimes called large gauge transformations, a name which reflects the fact that they cannot be continuously deformed to the identity. Under such a gauge transformation, we can see that:

$$A_1 \rightarrow A_1 + \partial_1 \chi = A_1 + \frac{n}{eR} \quad (35)$$

the two different values of  $A_1$  correspond the same physics state. From this transformation, we can determine the expression of  $A_1$ ,

$$\begin{aligned} A_1(0) &= A_1(2\pi R) + \frac{n}{eR} \\ \Rightarrow A_1 &= \frac{nx^2}{2\pi eR^2} \end{aligned} \quad (36)$$

And we can write down the expression of  $B_z$ :

$$B_z = \partial_2 A_1 = \frac{n}{2\pi eR^2} \quad (37)$$

So the integral of magnetic field is:

$$\int_{T^2} dx^1 dx^2 B_z = \frac{2\pi n}{e} \quad \text{with } n \in \mathbb{Z} \quad (38)$$

We can also apply exactly the same argument to the electric field

$$E_z = \partial_0 A_3 \quad (A_0 = 0) \quad (39)$$

So the integral of electric field is:

$$\int_{T^2} dx^0 dx^3 E_z = \frac{2\pi n'}{e} \quad \text{with } n' \in \mathbb{Z} \quad (40)$$

Armed with two integrals above, the  $S_\theta$  becomes:

$$\begin{aligned}
 S_\theta &= -\frac{\theta e^2}{4\pi^2} \int_{T^4} d^4x \mathbf{E} \cdot \mathbf{B} \\
 &= -\frac{\theta e^2}{4\pi^2} \frac{4\pi^2}{e^2} \cdot nn' \quad (\text{let } N = -nn') \\
 &= N\theta \quad \text{with } N \in \mathbb{Z}
 \end{aligned} \tag{41}$$

It's pretty straightforward to generalise it to non-constant  $\mathbf{E}$  and  $\mathbf{B}$  fields:

$$\begin{aligned}
 \int_{T^2} dx^1 dx^2 B_z &= \int_{T^2} dx^1 dx^2 \cdot \partial_2 A_1 \\
 &= 2\pi R [A_1(2\pi R) - A_1(0)] \\
 &= \frac{2\pi n}{e}
 \end{aligned} \tag{42}$$

So we can see that  $\theta$  is a periodic variable in quantum theory.

Although the theta term is gauge invariant and Lorentz invariant, it is not invariant under certain discrete symmetries, and does not preserve the same symmetries as the Maxwell term. These discrete symmetries are parity  $\mathcal{P}$  and time reversal  $\mathcal{T}$ .

$$\begin{cases} \mathcal{P} : x \rightarrow -x \\ \mathcal{T} : t \rightarrow -t \end{cases} \tag{43}$$

If we act on both  $\mathbf{E}$  and  $\mathbf{B}$

$$\begin{cases} \mathcal{P} : \mathbf{E}(x, t) \rightarrow -\mathbf{E}(-x, t) \text{ and } \mathcal{P} : \mathbf{B}(x, t) \rightarrow \mathbf{B}(-x, t) \\ \mathcal{T} : \mathbf{E}(x, t) \rightarrow \mathbf{E}(x, -t) \text{ and } \mathcal{T} : \mathbf{B}(x, t) \rightarrow -\mathbf{B}(x, -t) \end{cases} \tag{44}$$

So  $\mathbf{E}$  is odd under parity and even under time reversal.  $\mathbf{B}$  is even under parity and odd under time reversal.

This means that, the theta term  $S_\theta \propto \int d^4x \mathbf{E} \cdot \mathbf{B}$  breaks both parity and time reversal invariance, which means  $S_\theta \rightarrow -S_\theta$  under  $\mathcal{P}$  and  $\mathcal{T}$ . But because of the periodicity of  $\theta$ , there are two exceptions:

$$S_\theta = -S_\theta + 2\pi N \tag{45}$$

In order to make the equation hold for any integer values of  $N$ ,  $\theta$  must equal to 0 or  $\pi$  ( $\theta \in [0, 2\pi)$ ).

So we can know that when  $\theta = 0$  or  $\theta = \pi$ , the theory is invariant under  $\mathcal{P}$  and  $\mathcal{T}$  symmetries.

In subsequent sections, we can see that the topological insulator said to have a  $\mathbb{Z}_2$  classification has  $\theta = \pi$  because these materials are defined to be time-reversal invariant. Most materials have  $\theta = 0$ . But some materials have a band structure which is twisted in a particular way, which results in  $\theta = \pi$ .

## **1.2 The Witten effect**