

5. 静电场与静磁场

$$\text{由 } \nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

$$\nabla \times \vec{E} = 0$$

可以定义电势:

$$\vec{E} = -\nabla \phi$$

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$$\therefore \nabla^2 \phi = -\frac{\rho}{\epsilon_0} \quad (\text{泊松方程})$$

Poisson.

$$\text{当 } \rho = 0,$$

$$\nabla^2 \phi = 0, \quad (\text{Laplace 方程})$$

在线性介质中:

$$\nabla \cdot \vec{D} = \rho_f$$

$$\therefore \nabla \cdot \epsilon \vec{E} = \rho_f$$

$$\therefore \nabla^2 \phi = \frac{\rho_f}{\epsilon}$$

$$\text{边界条件: } \vec{E}_{\text{up}}^{\perp} = \vec{E}_{\text{down}}^{\perp}$$

$$\begin{cases} D_{\text{up}}^{\perp} - D_{\text{down}}^{\perp} = \sigma_f \Rightarrow -\epsilon_{\text{up}} \frac{\partial \phi_{\text{up}}}{\partial n} + \epsilon_{\text{down}} \frac{\partial \phi_{\text{down}}}{\partial n} = \sigma_f \\ \phi = \phi_1 = \phi_2 \end{cases}$$

$$\text{对于导体, 等势体有: } \phi_{\text{down}} = \text{const}, \quad D_{\text{up}}^{\perp} = -\epsilon_{\text{up}} \frac{\partial \phi_{\text{up}}}{\partial n} = \sigma_f$$

$$\text{总能量: } W = \frac{1}{2} \int_V \vec{E} \cdot \vec{D} dV$$

$$= \frac{1}{2} \int_V \rho \phi dV$$

5.2 唯一性定理.

V 内电势要确定, 即给出 $\nabla^2 \varphi = \frac{\rho}{\epsilon_0}$ 的 φ 有唯解 (差常数).

除了给出 ρ 外, 还给给出 $\varphi|_{\partial V}$ 或 $\frac{\partial \varphi}{\partial n}|_S$.

若区域中包含导体, 若在上述条件上加上给定每个导体总电荷, 电势唯一确定.

在球坐标中, Laplace 方程解为:

$$\varphi(R, \theta, \varphi) = \sum (A_{nm} R^n + \frac{b_{nm}}{R^{n+1}}) P_n^m(\cos \theta) \cdot \cos m\varphi + \sum (C_{nm} R^n + \frac{b_{nm}}{R^{n+1}}) P_n^m(\cos \theta) \sin m\varphi.$$

轴对称时, $\varphi(R, \theta) = \sum (A_n R^n + \frac{b_n}{R^{n+1}}) P_n(\cos \theta)$

$$\begin{cases} P_1 = x \\ P_2(x) = (3x^2 - 1) \cdot \frac{1}{2} \end{cases}$$

5.3 电多极矩:

$$\varphi = \frac{1}{4\pi\epsilon_0} \int \rho(\vec{r}') \frac{1}{|\vec{r} - \vec{r}'|} dV.$$

$\frac{1}{|\vec{r} - \vec{r}'|}$ 在 $\vec{r}' = 0$ 附近展开.

$$\frac{1}{|\vec{r} - \vec{r}'|} = \frac{1}{r} + \frac{\vec{r}' \cdot \nabla}{r^2} + \frac{1}{2!} \frac{\partial^2}{\partial r^2} (\vec{r}' \cdot \vec{r}) + \dots$$

$$\varphi^{(0)} = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}')}{r} dV.$$

$$\varphi^{(1)} = \frac{1}{4\pi\epsilon_0} \int \rho(\vec{r}') \frac{\vec{r}' \cdot \nabla}{r^2} dV = \frac{\nabla \cdot \vec{r}}{4\pi\epsilon_0} \int \rho(\vec{r}') \frac{(\vec{r}') \cdot \nabla}{r^2} dV$$

$$\varphi^{(2)} = \frac{1}{4\pi\epsilon_0} \int \rho(\vec{r}') \frac{1}{2!} (\vec{r}') \cdot \nabla^2 \frac{1}{r} dV =$$



$$\therefore Q = \int \rho(\vec{r}) dV$$

$$\vec{P} = \int_V \rho(\vec{r}) \cdot \vec{r} dV$$

$$\vec{D} = \int_V 3 \vec{r} \vec{r} \rho(\vec{r}) dV$$

$$\therefore \varphi^{(2)} = \frac{1}{4\pi\epsilon_0} \cdot \frac{1}{6} \cdot \vec{D} : \nabla \nabla \frac{1}{R}$$

$$\varphi^{(1)} = \frac{1}{4\pi\epsilon_0} \vec{P} \cdot \nabla \frac{1}{R} = \frac{\vec{P} \cdot \vec{R}}{4\pi\epsilon_0 R^3}$$

$$\varphi^{(2)} = \frac{1}{4\pi\epsilon_0} \cdot \frac{1}{6} \sum_{ij} D_{ij} \frac{\partial^2}{\partial x_i \partial x_j} \frac{1}{R}, \quad D_{ij} = \int 3x_i x_j \rho(\vec{r}) dV$$

能量

$$W = \int \rho \varphi_e dV$$

$\varphi_e(\vec{r})$ 在 0 处展开:

$$\varphi_e(\vec{r}) = \varphi_e(0) + \vec{r} \cdot \nabla \varphi_e|_{\vec{r}=0} + \frac{1}{2!} \sum_{ij} 3x_i x_j \frac{\partial^2}{\partial x_i \partial x_j} \varphi_e|_{\vec{r}=0}$$

$$\therefore W^{(0)} = \int \rho(x) \varphi_e(0) dV = Q \varphi_e(0)$$

$$\begin{aligned} W^{(1)} &= \int \rho(x) \cdot \vec{r} \cdot \nabla \varphi_e|_{\vec{r}=0} dV \\ &= \vec{P} \cdot \nabla \varphi_e|_{\vec{r}=0} = -\vec{P} \cdot \vec{E}(0) \end{aligned}$$

$$\begin{aligned} W^{(2)} &= \frac{1}{6} \int \sum 3x_i x_j \frac{\partial^2}{\partial x_i \partial x_j} \varphi_e|_{\vec{r}=0} \rho(\vec{r}) dV \\ &= \frac{1}{6} \vec{D} : \nabla \nabla \varphi_e|_{\vec{r}=0} \end{aligned}$$

5.4. 磁矢势

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{B} = \mu_0 \vec{J}$$

$$\nabla \times \vec{B} = \nabla \times (\nabla \times \vec{A})$$

$$= \nabla (\nabla \cdot \vec{A}) - \nabla^2 \vec{A} = \mu_0 \vec{J}$$

$$\therefore \nabla \times \nabla \varphi = 0$$

$$\vec{A}' = \vec{A} + \nabla \varphi$$

$$\vec{B}' = \vec{B}$$

$$\nabla \cdot \vec{A}' + \nabla^2 \varphi = 0$$

找到 $\nabla \cdot \vec{A}' = -\nabla^2 \varphi$ 时,

$$\nabla \cdot (\nabla \cdot \vec{A}') = 0$$

$$\therefore \nabla^2 \vec{A}' = \mu_0 \vec{J}$$

分量满足泊松方程.

$$\therefore \vec{A} = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}')}{|\vec{r} - \vec{r}'|} dV'$$

对线电流

$$\vec{A} = \frac{\mu_0}{4\pi} \int \frac{I \cdot d\vec{r}'}{|\vec{r} - \vec{r}'|}$$

面电流

$$\vec{A} = \frac{\mu_0}{4\pi} \int \frac{\vec{K} \cdot d\vec{s}}{|\vec{r} - \vec{r}'|}$$



边界条件:

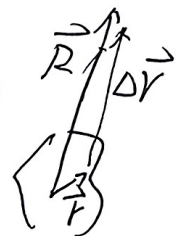
$$\begin{cases} B_{up} - B_{down} = 0 \\ \vec{B}_{up}'' - \vec{B}_{down}'' = \mu_0 \vec{K} \times \vec{n} \\ (\vec{H}_{up}'' - \vec{H}_{down}'') = \vec{K}_f \times \vec{n} \end{cases}$$

$$\Rightarrow \begin{cases} \vec{A}_{up} = \vec{A}_{down} \\ \epsilon_1 \times (\frac{1}{\mu_2} \nabla \times \vec{A}_2 - \frac{1}{\mu_1} \nabla \times \vec{A}_1) = \vec{\alpha} \end{cases}$$

5.5 磁矩矩.

$$\vec{A} = \frac{\mu_0}{4\pi} \int \vec{J}(\vec{r}') \cdot d\vec{V} \cdot \frac{1}{r}$$

$$= \frac{\mu_0}{4\pi} \int \vec{J}(\vec{r}') dV \left[\frac{1}{r} + (-\vec{r}') \cdot \nabla \frac{1}{r} + \frac{1}{2!} \sum_{i,j} x_i' x_j' \frac{\partial^2}{\partial x_i \partial x_j} \frac{1}{r} + \dots \right]$$



$$\vec{A}^{(0)} = \frac{\mu_0}{4\pi R} \int \vec{J}(\vec{r}') dV$$

$$= \frac{\mu_0}{4\pi R} \oint I d\vec{\ell}$$

$$= 0$$



$$\vec{A}^{(1)} = \frac{\mu_0}{4\pi} \int \vec{J}(\vec{r}') \cdot d\vec{V} \cdot (-\vec{r}') \cdot \nabla \frac{1}{r}$$

$$\nabla \frac{1}{r} = \frac{\vec{r}}{r^2} \cdot \frac{1}{r^3}$$

$$\vec{A}^{(1)} = \frac{\mu_0}{4\pi} \cdot \frac{\vec{R}}{R^3} \cdot \int \vec{J}(\vec{r}') dV \cdot (-\vec{r}')$$

$$0 = \oint d\vec{l} (\vec{r} \cdot \vec{R}) \vec{r}' = \oint_L (\vec{r}' \cdot \vec{R}) d\vec{l} + \oint_L (d\vec{l} \cdot \vec{R}) \cdot \vec{r}'$$

$$\int \vec{r} \cdot \vec{R} \cdot d\vec{l} = \frac{1}{2} \int [(\vec{r} \cdot \vec{R}) d\vec{l} - (d\vec{l} \cdot \vec{R}) \cdot \vec{r}]$$

$$= \frac{1}{2} \oint (\vec{r} \times d\vec{l}) \times \vec{R}$$

$$\therefore \vec{A}^{(1)} = \frac{\mu_0}{4\pi R^3} (-\vec{r}) \cdot \vec{R} \vec{J}(\vec{r}) dV = \frac{\mu_0}{4\pi} \cdot \vec{R} \times \frac{1}{2} \int \vec{r} \times \vec{J}(\vec{r}) dV$$

磁矩偶极矩:

$$= \frac{\mu_0}{4\pi} \cdot \frac{\vec{R} \times \vec{m}}{R^3}$$

$$\vec{m} = \frac{1}{2} \int_V \vec{r} \times \vec{J}(\vec{r}) \cdot dV$$

$$= \frac{1}{2} \int \vec{r} \times \vec{J} d\vec{r}$$

$$= \frac{I}{2} \int \vec{r} \times d\vec{r}$$

$$\Delta \vec{S} = \frac{1}{2} \int \vec{r} \times d\vec{r}$$

$$\therefore \vec{m} = I \Delta \vec{S}$$

$$\therefore \vec{B}^{(1)} = \nabla \times \vec{A}^{(1)}$$

$$= \frac{\mu_0}{4\pi} \nabla \times \left(\vec{m} \times \frac{\vec{R}}{R^3} \right)$$

$$= \frac{\mu_0}{4\pi} \left[(\nabla \cdot \frac{\vec{R}}{R^3}) \cdot \vec{m} - (\vec{m} \cdot \nabla) \cdot \frac{\vec{R}}{R^3} \right]$$

$$\nabla \cdot \frac{\vec{R}}{R^3} = 0 \quad R \neq 0$$

$$\therefore \vec{B}^{(1)} = -\frac{\mu_0}{4\pi} (\vec{m} \cdot \nabla) \cdot \frac{\vec{R}}{R^3}$$

$$\nabla \cdot (\vec{m} \cdot \frac{\vec{R}}{R^3}) = \underbrace{(\nabla \cdot \vec{m}) \cdot \vec{R}}_{=0} + \underbrace{(\vec{m} \cdot \nabla) \frac{\vec{R}}{R^3}}_{=0} + \frac{\vec{R}}{R^3} \times (\nabla \times \vec{m}) + \underbrace{(\frac{\vec{R}}{R^3} \cdot \nabla) \cdot \vec{m}}_{=0}$$

$$\therefore \nabla \cdot (\vec{m} \cdot \frac{\vec{R}}{R^3}) = (\vec{m} \cdot \nabla) \cdot \frac{\vec{R}}{R^3}$$

$$\therefore B^{(1)} = -\mu_0 \nabla \psi_m^{(1)}, \quad \psi_m^{(1)} = \frac{\vec{m} \cdot \vec{R}}{4\pi R^3}$$

磁场能量

$$W = \frac{1}{2} \int \vec{B} \cdot \vec{H} dV$$

$$\begin{aligned} \vec{B} \cdot \vec{H} &= (\nabla \times \vec{A}) \cdot \vec{H} = \nabla \cdot (\vec{A} \times \vec{H}) + \vec{A} \cdot (\nabla \times \vec{H}) \\ &= \cancel{\nabla \cdot (\vec{A} \times \vec{H})} + \vec{A} \cdot \vec{J}_f \end{aligned}$$

$$\begin{aligned} W &= \frac{1}{2} \int_V (\vec{A} \times \vec{J}) dV + \frac{1}{2} \int_V \vec{A} \cdot \vec{J}_f dV \\ &= \frac{1}{2} \int_S (\vec{A} \times \vec{J}) \cdot d\vec{s} \end{aligned}$$

无穷远处, $\vec{A} \times \vec{J} = 0$

$$\therefore W = \frac{1}{2} \int \vec{A} \cdot \vec{J} dV$$

$$\vec{J} = \vec{J}_I + \vec{J}_0, \quad \vec{A} = \vec{A}_I + \vec{A}_0$$

产生

$$\therefore W = \underbrace{\frac{1}{2} \int (\vec{J}_0 \cdot \vec{A}_I + \vec{J}_I \cdot \vec{A}_0) dV}_{\text{相互作用}} + \underbrace{\frac{1}{2} \int \vec{J}_I \cdot \vec{A}_I dV + \frac{1}{2} \int \vec{J}_0 \cdot \vec{A}_0 dV}_{\text{自作用}}$$

$$\vec{A}_I = \frac{\mu}{4\pi} \int \frac{\vec{J}_I(\vec{r}) dV}{r}$$

$$\vec{A}_I \cdot \vec{J}_0 = \frac{\mu}{4\pi} \int \frac{\vec{J}_I(\vec{r}) \cdot \vec{J}_0(\vec{r}) dV}{r}$$

同理:

$$\vec{A}_I \cdot \vec{J}_0 = \vec{A}_0 \cdot \vec{J}_I$$

$$\therefore W_{int} = \int \vec{J}_I \cdot \vec{A}_0 dV.$$