

Exponential Integrators

for ODEs

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In our discussion of IMEX
(implicit-explicit) schemes we
considered ODEs of the form
(these often arise when discretizing
PDEs with the "method of lines"):

$$u'(t) = A u(t) + B(u(t))$$

stiff linear
with A constant
(e.g. diffusion)

non-stiff
but non-linear
(e.g. advection)

①

If $B \equiv 0$ then $u' = Au \Rightarrow$

$$u_A(t) = e^{At} u(0)$$

Requires action of matrix exponential

For small normal A , use
eigenvalue or SVD decomposition:

$$A = U \Sigma V^* \Rightarrow$$

$$\exp(A) = U \exp(\Sigma) V^*$$

scalar exponential of singular values
Trivial for diagonal A !

(2)

If A is large but sparse, can use Krylov / Lanczos / Arnoldi to compute $\exp(At)u$.

An example

Pseudo spectral discretization in space for KDV (homeworks 2 & 4)

$$\frac{d\hat{\psi}}{dt} = \underbrace{ik^3 \hat{\psi}}_{A \hat{\psi}} - \underbrace{3ik \mathcal{F}((\mathcal{F}^{-1}\hat{\psi})^2)}_{B(\hat{\psi})}$$

$$A = \text{Diag}(ik^3) \quad (\text{easy!})$$

③

For more general stiff ODEs,
Rosenbrock methods simply
linearize the ODE around the
current solution. Take the
general autonomous ODE:

$$u' = f(u(t))$$

$$u' = \left(\frac{\partial f}{\partial u} \right)^n (u - u^n) + f^n \} \text{linear part}$$

$$+ f(u) - \left[\left(\frac{\partial f}{\partial u} \right)^n (u - u^n) + f^n \right]$$

(add and subtract)

④

$$\Rightarrow u' = A^n u + B^n(u)$$

$$A^n = \left. \frac{\partial f}{\partial u} \right|_{u=u^n}$$

Jacobian

$$B^n(u) = f(u) - \left(\frac{\partial f}{\partial u} \right)^n u$$

This is an exact splitting into stiff and (hopefully!) non-stiff terms. Note that

$$\frac{\partial B}{\partial u}(u^n) = \left. \frac{\partial f}{\partial u} \right|_{u^n} - A^n = 0$$

(5)

which means $B(n)$ is almost constant during the time step!

$$u' = Au + B(n) = f(n)$$

Variation of constants / Duhamel's methods + Picard iteration give:

$$u^{n+1} = \left(e^{A \Delta t} \right)^n u + \int_{t^n}^{t^{n+1}} e^{A(t-\bar{z})} B(u(\bar{z})) d\bar{z}$$

Exact!

Approximate integral using a quadrature rule as for RK / multistep (6)

If $A \equiv 0$, $\exp(A \Delta t) = \text{Identity}$,
 we want to get **explicit RK**.

If $B \equiv 0$, we want to get
exact solution $\exp(A \Delta t) \cdot u(t)$.

For example, for RK1 (forward
 Euler), approximate/extrapolate:
 $B^n(u(\tau)) \simeq B^n(u^n)$

and use

$$\int_{t_n}^{t_{n+1}} \exp(A^n(t_{n+1} - \tau)) d\tau = A_n^{-1} (e^{A_n \Delta t} - I) \quad (7)$$

to get

$$u^{n+1} = e^{A_n \Delta t} + A_n^{-1} (e^{A_n \Delta t} - I) B_n(u^n)$$

Exponential Euler scheme:

$$u^{n+1} = u^n + (A^n)^{-1} (e^{A^n \Delta t} - I) \underbrace{f(u^n)}_{\text{r.h.s.}}$$

If $\|A_n\|_2 \Delta t \ll 1$, then

$$u^{n+1} \approx u^n + \Delta t f(u^n)$$

which is the usual forward Euler, with bad stability properties.

(8)

Local truncation error (LTE):

$$\tau^n = \left(\frac{u(t_{n+1}) - u(t_n)}{\Delta t} \right) - \frac{1}{\Delta t} (A^n)^{-1} (e^{A_n \Delta t} - I) u'(t_n)$$

from ODE

$$= \frac{\Delta t}{2} \left(\frac{\partial}{\partial u} (u(t_n)) - A_n \right) u'(t_n) + O(\Delta t^2)$$

\Rightarrow { second-order accurate if $A_n = \left(\frac{\partial}{\partial u} \right)^n$
 (Rosenbroch)
 first-order in general
 (but much better than
 standard Euler for stiff ODEs)

this method will be accurate
if $B(n)$ varies slowly over the
time step.

Note that if $u^{n+1} = u^n$ (fixed point)

then $f^n = 0$. This means

that this method can be used
to quickly find stable fixed
points of ODEs / PDEs.

This is not true for another class

of "Integrating Factor" Exponential
Integrators which are used in
some fields (10)

We can improve the order of accuracy with the same ideas we used for RK methods:

$$u^{n+1} = \left(e^{A \Delta t} \right)^n u + \int_{t^n}^{t^{n+1}} e^{A(t - \bar{z})} B(u(\bar{z})) d\bar{z}$$

$$B^n(u(\bar{z})) \approx B^n + \frac{\bar{z} - t^n}{\Delta t} (B^n - B^{n-1})$$

Now do the integral (assume A is normal & diagonalize it)

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Exponential Time Differencing **ETD2**:

$$u^{n+1} = e^{A_n \Delta t} u^n + A_n (e^{A_n \Delta t} - I) B^n + \underbrace{\frac{A_n^{-2}}{\Delta t} (e^{A_n \Delta t} - I - A_n \Delta t)}_{O(\Delta t^2) \text{ 2nd order correction}} (B^n - B^{n-1})$$

This is an **exponential multistep scheme** (two step like AB2)

If we want an RK2 scheme instead, say explicit midpoint:

① Midpoint ETD RK1 (exponential Euler) predictor step:

$$u^{n+1/2,*} = \left(e^{A_n \Delta t / 2} \right) u^n + A_n^{-1} \left(e^{A_n \Delta t / 2} - I \right) B^n$$

corrector:

② ETD RK2

$$u^{n+1} = \left(e^{A_n \Delta t} \right) u_n + A_n^{-1} \left(e^{A_n \Delta t} - I \right) B^n$$

$$+ 2 A_n^{-2} \left(\underbrace{e^{A_n \Delta t} - I - A_n \Delta t}_{\text{ETD2}} \right) \begin{pmatrix} u^{n+1/2,*} \\ B^n - B^n \end{pmatrix}$$

②

Cox and Matthews also derive a set of ETD methods based on Runge-Kutta time-stepping, which they call ETDRK schemes. In this report we consider only the fourth-order scheme of this type, known as ETDRK4. According to Cox and Matthews, the derivation of this scheme is not at all obvious and requires a symbolic manipulation system. The Cox and Matthews ETDRK4 formulae are:

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$$a_n = e^{\mathbf{L}h/2}u_n + \mathbf{L}^{-1}(e^{\mathbf{L}h/2} - \mathbf{I})\mathbf{N}(u_n, t_n),$$

$$b_n = e^{\mathbf{L}h/2}u_n + \mathbf{L}^{-1}(e^{\mathbf{L}h/2} - \mathbf{I})\mathbf{N}(a_n, t_n + h/2),$$

$$c_n = e^{\mathbf{L}h/2}a_n + \mathbf{L}^{-1}(e^{\mathbf{L}h/2} - \mathbf{I})(2\mathbf{N}(b_n, t_n + h/2) - \mathbf{N}(u_n, t_n)),$$

$$\begin{aligned} u_{n+1} = & e^{\mathbf{L}h}u_n + h^{-2}\mathbf{L}^{-3}\{[-4 - \mathbf{L}h + e^{\mathbf{L}h}(4 - 3\mathbf{L}h + (\mathbf{L}h)^2)]\mathbf{N}(u_n, t_n) \\ & + 2[2 + \mathbf{L}h + e^{\mathbf{L}h}(-2 + \mathbf{L}h)](\mathbf{N}(a_n, t_n + h/2) + \mathbf{N}(b_n, t_n + h/2)) \\ & + [-4 - 3\mathbf{L}h - (\mathbf{L}h)^2 + e^{\mathbf{L}h}(4 - \mathbf{L}h)]\mathbf{N}(c_n, t_n + h)\}. \end{aligned}$$

FROM PAPER BY KASSAM & TREFFETHEN
(linked on course webpage)

PROBLEM: Roundoff in
 $|z| < 1$

$$\frac{e^z - 1}{z}$$

(use complex contour integration)