## Solenoid electron lenses

Fundamentals and design

22.06.2020

## Structure

1 Motivation

2 Magnetic lens overview

3 Summary

### Motivation

- Magnetic lenses important components of particle accelerators, microscopes
- Electromagnet solenoids physical foundation of more advanced magnetic lenses
- Design constraints:
  - Power consumption
  - Physical size, material usage
  - Characteristic parameters for interaction with other machine components



Figure: Schematic<sup>1</sup> of AREAL, an electron bunch-research oriented linac<sup>2</sup>

<sup>&</sup>lt;sup>1</sup>Grigoryan et al., "Status of AREAL RF Photogun Test Facility".

<sup>&</sup>lt;sup>2</sup>Grigoryan et al., "Advanced Research Electron Accelerator Laboratory Based on Photocathode RE Gun"

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# Electron optics

Definition

### Classic optics

- Light ray photons
- •Can pass through optically transparent solids
- •Bends due to refractive index difference between media

### Electron optics

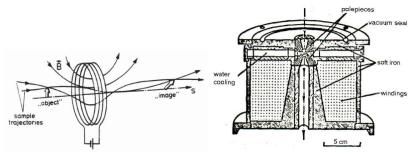
- •Electron beam electrons
- Gets absorbed/loses energy due to interactions with atoms in media
- •Bends due to Coulomb and Lorentz forces in the presense of external EM-fields

Requirements

#### A lens must have the following properties:

- deflection increases with increasing deviation of the beam from the optic axis;
- electron energy should not change, or change negligebly;
- symmetry of deflection on all sides of the optical axis;
- methods and laws of classical optics (such as thin lens formula and approximation, matrix formalism) are assumed to be applicable.

Solenoids



(a) A solenoid with S as optical axis<sup>1</sup>

**(b)** Cross-section of a magnetic lens<sup>2</sup>

<sup>&</sup>lt;sup>1</sup>Rossbach and Schmüser, Basic course on accelerator optics.

<sup>&</sup>lt;sup>2</sup>Egerton et al., *Physical principles of electron microscopy*.

Solenoids

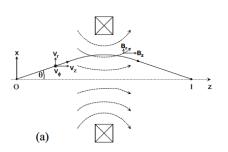


Figure: A solenoid cross-section along the optic axis<sup>1</sup>

$$F_{\varphi} = e \left( B_z v_r - v_z B_r \right) \tag{1}$$

$$\gamma m \left( \ddot{r} - r \dot{\varphi}^2 \right) = -e r \dot{\varphi} B_z \tag{4}$$

$$F_r = -e(v_z B_z)$$

$$F_z = e(v_\varphi B_r)$$
(2) 
$$\gamma m \frac{d}{dt} (r^2 \dot{\varphi}) = -er(\dot{r} B_z + B_r \dot{z})$$
(5)

$$\gamma m\ddot{z} = er\dot{\varphi}B_r \tag{6}$$

<sup>&</sup>lt;sup>1</sup>Egerton et al., Physical principles of electron microscopy.

#### Electron path equations

•From (5) follows:

$$\dot{\varphi} = \frac{e}{2\gamma m} B_z \tag{9}$$

$$B_z(z,r) = \sum_n \frac{(-1)^n}{n! n!} \left(\frac{r}{2}\right)^{2n} \frac{\partial^{2n} B_{z, axis}}{\partial z^{2n}} \tag{7}$$
•From (4) and (6) follows:

$$B_{r}(z,r) = \sum_{n} \frac{(-1)^{n}}{n! (n-1)!} \left(\frac{r}{2}\right)^{2n-1} \frac{\partial^{2n-1} B_{z, axis}}{\partial z^{2n-1}}$$
(8) 
$$\ddot{r} = -\left(\frac{e}{2\gamma m}\right)^{2} r B_{z}^{2}$$
(10)

$$\ddot{r} = -\left(\frac{e}{2\gamma m}\right)^2 r B_z^2 \tag{10}$$

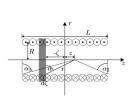
$$\ddot{z} = -\left(\frac{e}{2\gamma m}\right)^2 r^2 B_z B_z^2 \tag{11}$$

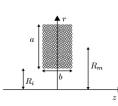
$$\ddot{r} = r''\dot{z}^2 \approx r'' (\beta c)^2 \Rightarrow r'' = \left(\frac{e}{2p_z}\right)^2 rB_z^2 \tag{12}$$

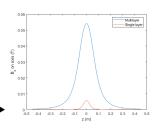
$$-\frac{r'}{r} = \frac{1}{f} \underset{integrate (12)}{\Rightarrow} \frac{1}{f} = \left(\frac{e}{2p_z}\right)^2 \int_{-\infty}^{\infty} B_z^2 dz := \left(\frac{e}{2p_z}\right)^2 F_2 \tag{13}$$

(11)

#### Magnetic field of solenoid







- (a) Single-wind solenoid<sup>1</sup>
- **(b)** Solenoid with multilayred windings<sup>1</sup>
- (c) Fields produced by two coils with same N and L=b

Approximate field of the solenoid (b):

Field of the solenoid (a):

$$B_{Z}(z) = \frac{\mu_{0}nI}{2} \left( \frac{\triangle z}{\sqrt{R^{2} + \triangle z^{2}}} - \frac{\triangle z^{*}}{\sqrt{R^{2} + \triangle z^{*2}}} \right)$$
(14)

$$\triangle z = z - L/2 \tag{15}$$

$$B_z(z) \approx \frac{\mu_0 N I}{4} \left( \frac{Rc^2}{\left(z^2 + Rc^2\right)^{3/2}} + \frac{Rc^{*2}}{\left(z^2 + Rc^{*2}\right)^{3/2}} \right)$$
 (16)

$$Rc = R_{sq} + c$$
,  $R_{sq} = R_m \left( 1 + \frac{a^2}{24R_m^2} \right)$ ,  $c^2 = \frac{b^2 - a^2}{12}$  (17)

<sup>&</sup>lt;sup>1</sup>Gehrke, "Design of Permanent Magnetic Solenoids for REGAE".

## Lens impefections

There are 3 main limitations to consider when designing a solenoid lens:

Sc	111	CE
	Jui	CC

•Chromatic aberration Spread of electron energies

•RMS Emittance growth

Spread of eletron coordinates in position-and-momentum phase space

•Spherical aberration Real lens  $\rightarrow$  different refraction based on distance from axis

Emittance:

$$\epsilon_{n,rms} = \frac{1}{mc} \sqrt{\left\langle x^2 \right\rangle \left\langle \bar{p}_X^2 \right\rangle - \left\langle x \bar{p}_X \right\rangle^2} = \frac{1}{mc} \left( \frac{e^2 \sigma^4}{3\sqrt{2} p_{z,0}} F_3 + \frac{e^4 \sigma^4}{24\sqrt{2} p_{z,0}^3} F_4 \right) \tag{18}$$

Chromatic aberration:

$$r_c \approx \alpha f \left(\triangle E_0 / E_0\right) \approx \alpha C_c \left(\triangle E_0 / E_0\right)$$
 (19)

# Lens impefections

#### Spherical aberration<sup>1</sup>

$$\triangle f = c \cdot x^2 \tag{20}$$

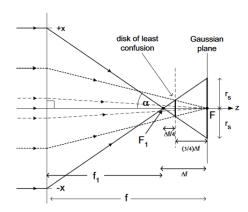
$$x \approx f \cdot tan(\alpha)$$
 (21)

$$r_s = \triangle f \cdot tan(\alpha) \approx c \cdot f^2 tan(\alpha)^3$$
 (22)

$$= C_{s} \cdot \left(\frac{r_{in}}{f - \frac{C_{s}r_{in}^{2}}{f^{2}}}\right)^{3} \tag{23}$$

$$F_3 = -\int \frac{B_Z B_Z^{\prime\prime}}{2} dz \tag{24}$$

$$F_4 = \int B_z^4 dz \tag{25}$$



$$C_{s} = \frac{e}{96m\tilde{U}} \int \left(\frac{2e}{m\tilde{U}}B_{z}^{4} + 5\left(B_{z}^{\prime}\right)^{2} - B_{z}B_{z}^{\prime\prime}\right)R^{4}dz \tag{26}$$

$$=\frac{e^2R^4}{4\rho_{z,0}^2}F_3+\frac{e^4R^4}{12\rho_{z,0}^4}F_4\tag{27}$$

Figure: Illustration of beam radius change due to spherical aberration

<sup>&</sup>lt;sup>1</sup>Egerton et al., Physical principles of electron microscopy.

# Summary

Subtitle

### References I



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