

Numerics of incompressible flows 1 (SS 25)
Homework 2

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Homework 2

Exercise 1

$$v(x, y, z) = v_r(r) \cdot e_r(r, \varphi, z)$$

$$\rho: \Omega \rightarrow \mathbb{R} \quad p(x, y, z) = p(r)$$

a)

i) Prove $\operatorname{div} v = 0$

Proof With the Note from Ex 1 we have

$$\operatorname{div} v = \frac{1}{r} \partial_r (r v_r) + \frac{1}{r} \partial_\varphi v_\varphi + \partial_z v_z$$

Since $v = v_r(r) \cdot e_r$ we can say ~~$v_\varphi = v_z = 0$~~ $v_r = v_z = 0$.

$$\Rightarrow \operatorname{div} v = \frac{1}{r} \partial_\varphi v_\varphi$$

Since v_φ ~~does~~^{is} not dependent to φ , the derivative in respect to φ is 0 ($\partial_\varphi v_\varphi = 0$)

$$\Rightarrow \operatorname{div} v = 0$$

ii) according to i) $V_z = 0$

$$\text{rot } v = \begin{pmatrix} \partial_y V_z - \partial_z V_y \\ \partial_z V_x - \partial_x V_z \\ \partial_x V_y - \partial_y V_x \end{pmatrix} = \begin{pmatrix} -\partial_z V_y \\ \partial_z V_x \\ \partial_x V_y - \partial_y V_x \end{pmatrix}$$

we further know

$$V = v_p(r) \begin{pmatrix} -\sin \varphi \\ \cos \varphi \\ 0 \end{pmatrix} = -\sin \varphi \cdot v_p(r) e_x + \cos \varphi \cdot v_p(r) e_y + 0 e_z$$

$$= -\frac{y}{r} v_p(r) e_x + \frac{x}{r} v_p(r) e_y = v_x e_x + v_y e_y + 0 e_z$$

$$\Rightarrow v_x = -\frac{y}{r} v_p(r) \quad v_y = \frac{x}{r} v_p(r)$$

→ these components are independent of z , thus:

$$\text{rot}(v) = \begin{pmatrix} 0 \\ 0 \\ \partial_x v_y - \partial_y v_x \end{pmatrix} = (\partial_x v_y - \partial_y v_x) e_z$$

$$\rightarrow \partial_x v_y = \left(\frac{d}{dx} \frac{x}{r(x,y)} \right) v_p(r) + \frac{x}{r} \partial_r v_p(r) \frac{d}{dx} r(x,y)$$

$$= \frac{1 \cdot r - \frac{x}{r^2} \frac{d}{dx} r(x,y)}{r^2} v_p(r) + \frac{x}{r} v_p(r) \frac{x}{r}$$

$$= \frac{r - \frac{x^2}{r}}{r^2} v_p(r) + \frac{x^2}{r^2} \partial_r v_p(r)$$

$$\rightarrow \partial_y v_x = -\frac{r - \frac{y^2}{r}}{r^2} v_p(r) + \frac{y^2}{r^2} \partial_r v_p(r)$$

$$\Rightarrow \partial_x v_y - \partial_y v_x = \left(\frac{r - \frac{x^2}{r}}{r^2} + \frac{r - \frac{y^2}{r}}{r^2} \right) v_p(r) + \frac{x^2 + y^2}{r^2} \partial_r v_p(r)$$

$$= \frac{2r - \frac{x^2 + y^2}{r}}{r^2} v_p(r) + \partial_r v_p(r) = \frac{1}{r} v_p(r) + \partial_r v_p(r)$$

→ with $\frac{1}{r} \partial_r (r v_p(r)) e_z = \left(\frac{1}{r} v_p(r) + \partial_r v_p(r) \right) e_z$ we conclude

$$\text{rot}(v) = (\partial_x v_y - \partial_y v_x) e_z = \left(\frac{1}{r} v_p(r) + \partial_r v_p(r) \right) e_z \quad \square$$

iii)

Prove $\exists g = g(r) : v \times \text{rot}(v) = \nabla g$

Proof

$$\begin{aligned}
 v \times \text{rot} v &= \begin{pmatrix} v_\varphi \frac{1}{r} \partial_r (r v_\varphi(r)) - 0 \\ 0 \\ -v_\varphi \frac{1}{r} \partial_r (r v_\varphi(r)) \end{pmatrix} \\
 &= \left(v_\varphi(r) \frac{x}{r} \cdot \frac{1}{r} \partial_r (r v_\varphi(r)), v_\varphi(r) \frac{y}{r} \cdot \frac{1}{r} \partial_r (r v_\varphi(r)), 0 \right)^T \\
 &= \left(\frac{x}{r^2} v_\varphi(r) \partial_r (r v_\varphi(r)), \frac{y}{r^2} v_\varphi(r) \partial_r (r v_\varphi(r)), 0 \right)^T \\
 &= \frac{1}{r^2} v_\varphi(r) \partial_r (r v_\varphi(r)) \begin{pmatrix} x \\ y \\ 0 \end{pmatrix}
 \end{aligned}$$

$$v = v_\varphi(r) e_\varphi(r, \varphi, z), \quad \text{rot}(v) = \frac{1}{r} \partial_r (r v_\varphi(r)) e_z$$

For simplicity we define $\lambda(r) := \frac{1}{r} \partial_r (r v_\varphi(r))$

Therefore

$$\begin{aligned}
 v \times \text{rot} v &= v_\varphi(r) e_\varphi(r, \varphi, z) \times \lambda(r) e_z(r, \varphi, z) \\
 &= v_\varphi(r) \lambda(r) \cdot (e_\varphi(r, \varphi, z) \times e_z(r, \varphi, z)) \\
 &= v_\varphi(r) \lambda(r) \cdot (-e_r(r, \varphi, z))
 \end{aligned}$$

Note:
 $e_\varphi = e_z \times e_r$
 $\Rightarrow e_r = e_\varphi \times e_z$

We get for $g = g(r) : \nabla g(r) = \frac{dg}{dr} e_r$

So we get

$$v \times \text{rot} v = \nabla g(r)$$

$$\text{for } \frac{dg}{dr} = -v_\varphi(r) \lambda(r)$$

b.) Prove: If solution has assumed symmetry, then:

$$-\frac{1}{Re} \Delta v + \nabla \Pi = 0 \quad \text{with } \Pi = p + \frac{1}{2} \|v\|^2 - g$$

Proof:

We use Homework 1b:

$$(v \cdot \nabla) v = \frac{1}{2} \nabla \|v\|^2 - v \times \text{rot } v$$

and put in the results from Ex 1ciii:

$$(v \cdot \nabla) v = \frac{1}{2} \nabla \|v\|^2 - \nabla g \quad \text{with } \frac{dg}{dr} = -v_\varphi(r) \lambda(r)$$

The Navier Stokes equation with 0 on the right side, can be written as

$$-\frac{1}{Re} \Delta v + (v \cdot \nabla) v + \nabla p = 0 \quad \text{where } p \text{ is pressure.}$$

Therefore we get

$$-\frac{1}{Re} \Delta v + \left(\frac{1}{2} \nabla \|v\|^2 - \nabla g \right) + \nabla p = 0$$

This can be rewritten

as:

$$-\frac{1}{Re} \Delta v + \nabla \left(\frac{1}{2} \|v\|^2 - g + p \right) = 0$$

and which is

$$-\frac{1}{Re} \Delta v + \nabla \Pi = 0 \quad \text{where } \Pi = p + \frac{1}{2} \|v\|^2 - g$$

$$\text{and } \frac{dg}{dr} = -v_\varphi(r) \lambda(r)$$

Ex 2

a.)

$$V'_\varphi = v_\varphi$$

Prove \exists scalar-valued $\psi(r)$: $\nabla \Delta V_\varphi = \psi e_r$

The Laplacian of a vector field in cylindrical coordinates

$$\begin{aligned} \Delta v &= (\nabla \cdot \nabla) v = \nabla(\nabla \cdot v) - \nabla \times (\nabla \times v) \\ &= \nabla \operatorname{div} v - \nabla \times \operatorname{rot} v \end{aligned}$$

We know from 1i and 1ii parts of it:

$$\Delta v = \nabla \operatorname{div} v - \nabla \times \operatorname{rot} v$$

$$\stackrel{1i}{\text{ii}} \rightarrow 0 - \nabla \times \left(\frac{1}{r} \partial_r (r v_\varphi(r)) \right) e_z$$

$$= - \nabla \times \lambda(r) \cdot e_z$$

Define $\lambda(r) = \dots$

$$= - \nabla \lambda(r) \times e_z$$

$$= - \frac{d\lambda(r)}{dr} e_r \times e_z$$

$$= - \frac{d\lambda(r)}{dr} (e_r \times e_z)$$

$$\stackrel{\text{Nok Exc. 1}}{\rightarrow} = - \frac{d\lambda(r)}{dr} e_\varphi$$

$$= - \frac{d}{dr} \left(\frac{1}{r} \partial_r (r v_\varphi(r)) \right) e_\varphi$$

$$\cancel{\frac{d}{dr} V_\varphi(r) = v_\varphi(r)} \quad v = \frac{d}{dr} V_\varphi(r) e_\varphi = V_\varphi(r) e_\varphi$$

So

$$\Delta v_\varphi = - \frac{d}{dr} \left(\frac{1}{r} \partial_r (r v_\varphi(r)) \right) e_\varphi$$

$$\Rightarrow \Delta V'_\varphi = - \frac{d}{dr} \left(\frac{1}{r} \partial_r (r V'_\varphi(r)) \right) e_\varphi$$

$$\Rightarrow \text{Define } \hat{\lambda} = - \frac{d}{dr} \left(\frac{1}{r} \partial_r (r V'_\varphi(r)) \right)$$

$$\text{Define } \hat{\lambda} = - \frac{d}{dr} \left(\frac{1}{r} \partial_r (r V'_\varphi(r)) \right)$$

$$\Rightarrow \nabla \Delta V'_\varphi = \nabla \hat{\lambda}$$

We see $\hat{\lambda} = \hat{\lambda}(r)$. Therefore the gradient $\nabla \hat{\lambda}(r)$ is only in e_r -direction

\Rightarrow There is a scalar valued funct. $\psi(r)$

Prove (*) equivalent to $-\frac{1}{\hbar c} e_2^* \nabla \Delta V_d(r) + \nabla \pi = 0$

Proof:

$$- \frac{1}{R_1} \Delta V + \Delta \pi = 0 \quad (*)$$

$$V = V_{cp} \cdot e_{cp}$$

$$\Delta V = \Delta V_{\varphi'}(r) e_{\varphi}$$

We have to show $\Delta V = \Delta V'_\varphi(r) e_\varphi \stackrel{!}{=} e_z \times \nabla \Delta V_\varphi(r)$

$$\Delta V = \Delta \frac{dV_{\psi}(r)}{dr} e_{\psi}$$

with $e_\varphi = e_z \times e_r$

$$\Delta V = \int \frac{dV_\varphi(r)}{dr} e_\varphi = \int \frac{dV_\varphi(r)}{dr} e_z \times e_r$$

$$= e_z \times \Delta \frac{dV_\phi(r)}{dr} e_r$$

$$= e_z \times \nabla \Delta V_y(r)$$

$$\Rightarrow \Delta V = e_z \times \nabla \Delta V_\phi(r)$$

$$\vec{r} = \frac{1}{\kappa_0} \vec{e}_z \times \nabla \Delta V_\varphi(r) + \sigma \vec{\pi} = 0$$

$$6.) -\frac{1}{\kappa c} e_z \times \nabla \Delta V_v(r) + \nabla \pi = 0$$

$$\stackrel{2a)}{\Rightarrow} -\frac{1}{\kappa c} e_z \times 4\kappa c e_r + \nabla \pi = 0$$

$$\stackrel{e_z = e_z \times e_r}{\Rightarrow} -\frac{1}{\kappa c} \psi(r) e_\varphi + \nabla \pi = 0$$

$$\Gamma \pi = p + \frac{1}{2} \|v\|^2 - g \quad \neg$$

$$\text{Since } p = p(r) \text{ and } g = g(r)$$

$$\Rightarrow \pi = \pi(r)$$

$$\neg \Rightarrow \nabla \pi = \frac{d\pi}{dr} e_r \quad \neg$$

$$\Rightarrow -\frac{1}{\kappa c} \psi(r) e_\varphi + \frac{d\pi}{dr} e_r = 0$$

$$\text{Since } e_\varphi \text{ and } e_r \text{ are linear independent}$$

$$\Rightarrow -\frac{1}{\kappa c} \psi(r) = 0 \quad \wedge \quad \frac{d\pi}{dr} = 0$$

$$\Rightarrow \psi = 0 \quad \wedge \quad \pi \equiv \text{const} //$$