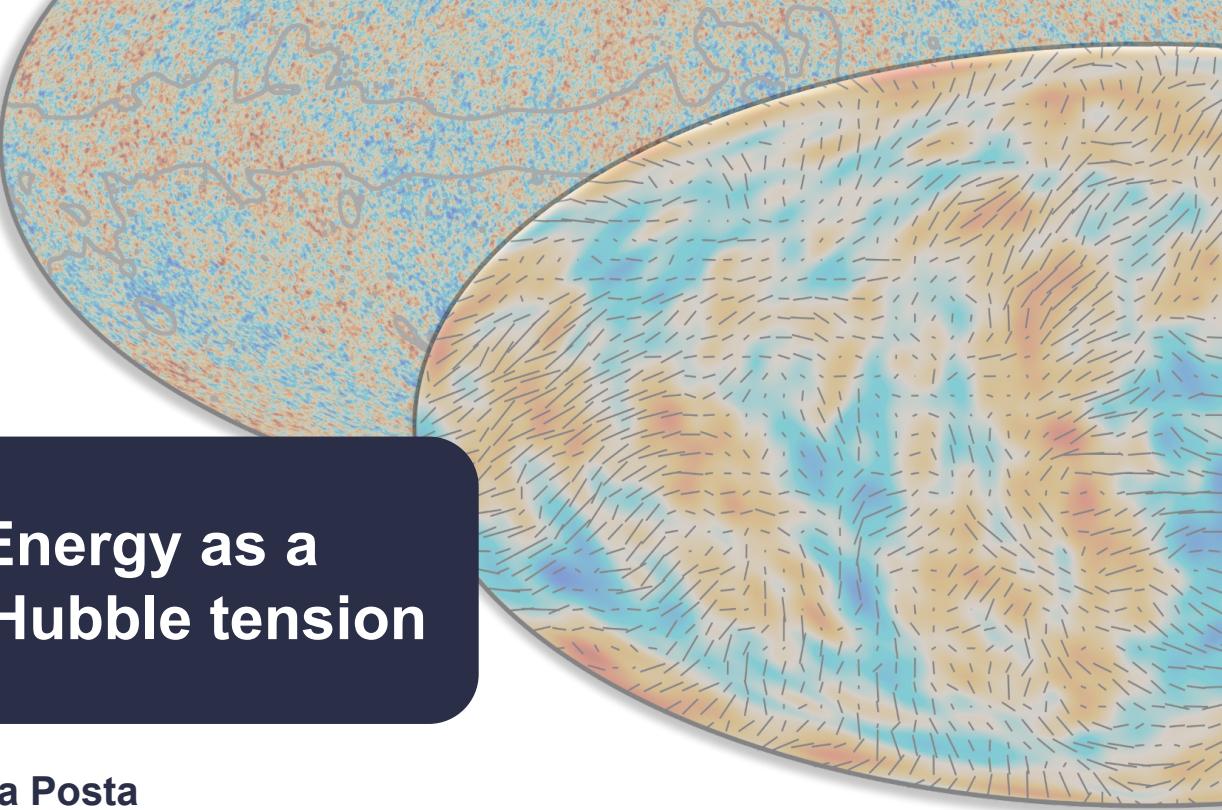




# Early Dark Energy as a solution to the Hubble tension

Adrien La Posta  
IJC lab



[link](#)

CosPT meeting - 06/12/22

laposta@ijclab.in2p3.fr

## The standard model of cosmology – $\Lambda$ CDM model

1

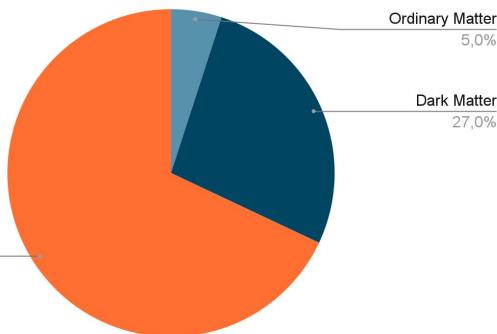
**FLRW metric**

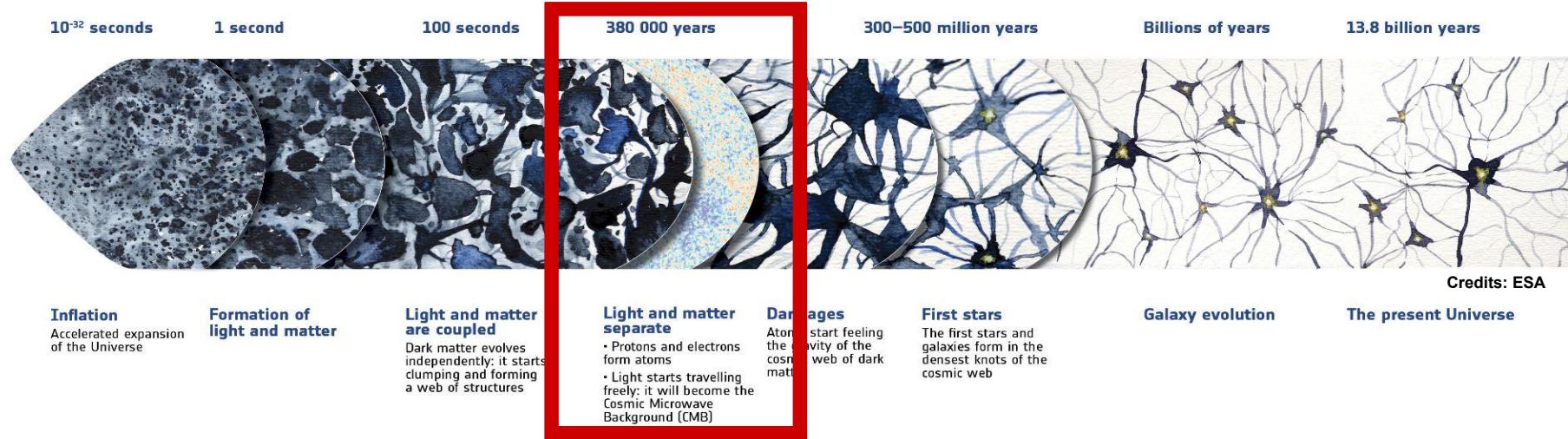
$$ds^2 = -c^2 dt^2 + a^2(t) \left( \frac{dr^2}{1 - Kr^2} + r^2 d\Omega^2 \right)$$

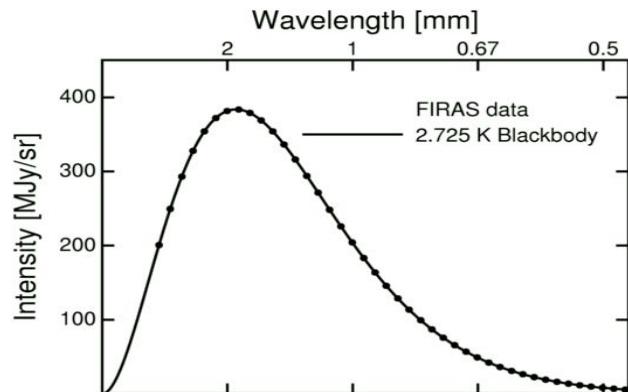
**Friedmann equation**

$$H^2(z) = \left( \frac{\dot{a}}{a} \right)^2 = \frac{8\pi G}{3} [(\rho_b^0 + \rho_c^0)(1+z)^3 + \rho_r^0(1+z)^4 + \rho_\Lambda]$$

baryon      CDM      radiation      dark energy



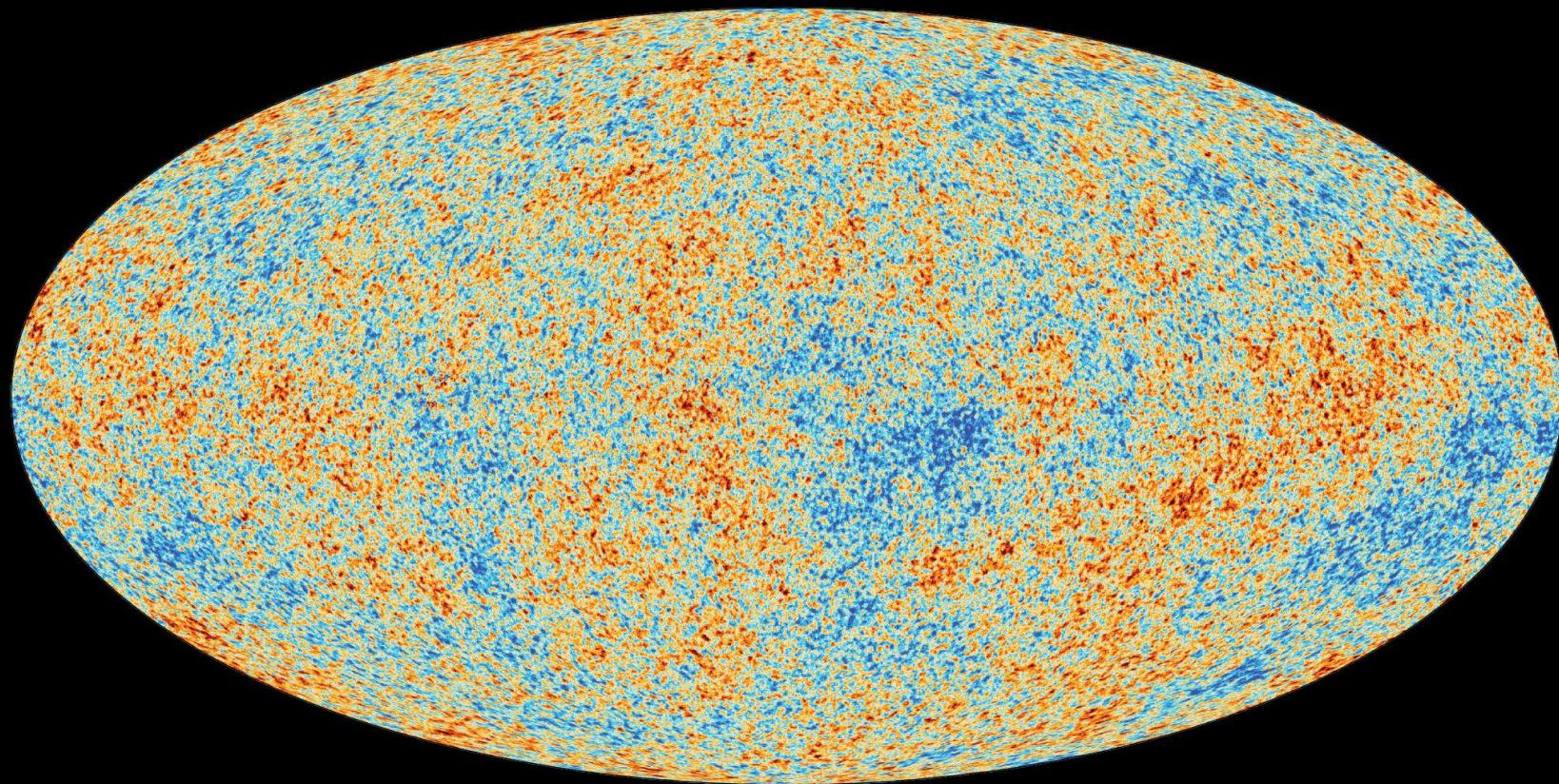




Nearly isotropic blackbody spectrum at T = 2.725 K

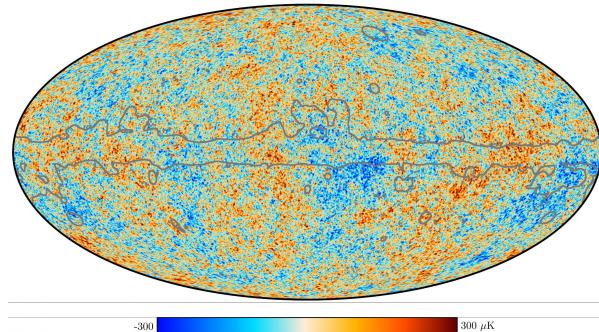
$$\frac{\delta T}{T} \sim 10^{-5}$$

# CMB temperature as measured by the Planck satellite

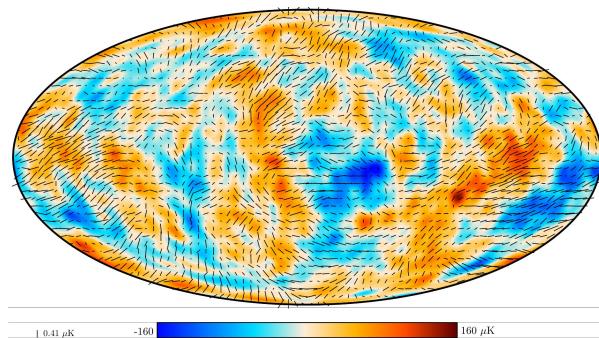


# Measuring $H_0$ from the CMB

Temperature



Polarization E-modes



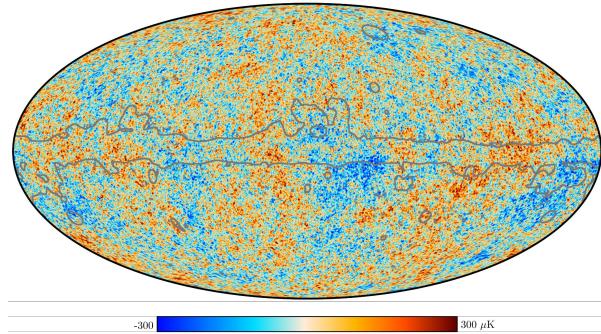
Spherical harmonics

$$\delta T(\hat{n}) = \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} a_{\ell m}^T Y_{\ell}^m(\theta, \phi)$$

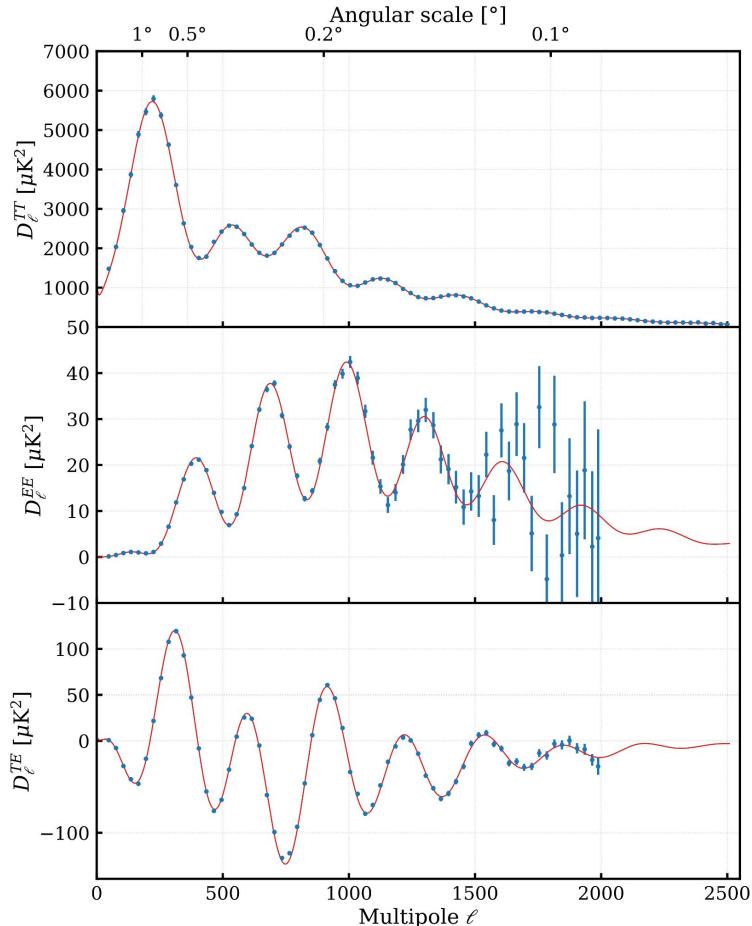
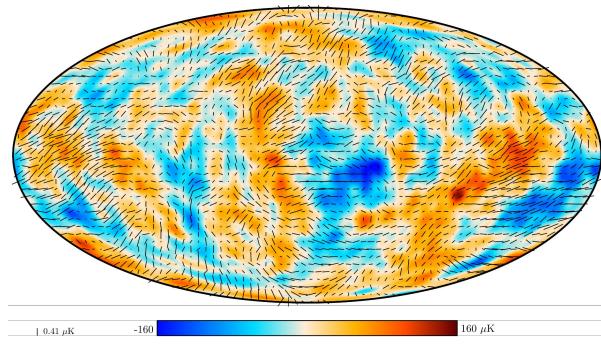
$$\langle a_{\ell m}^T a_{\ell' m'}^{T*} \rangle = \delta_{\ell \ell'} \delta_{mm'} C_{\ell}^{TT}$$

# Measuring $H_0$ from the CMB

Temperature

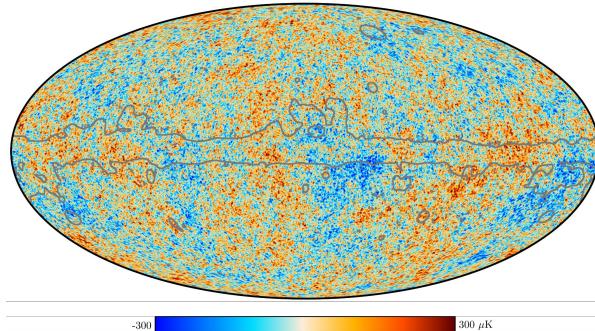


Polarization E-modes

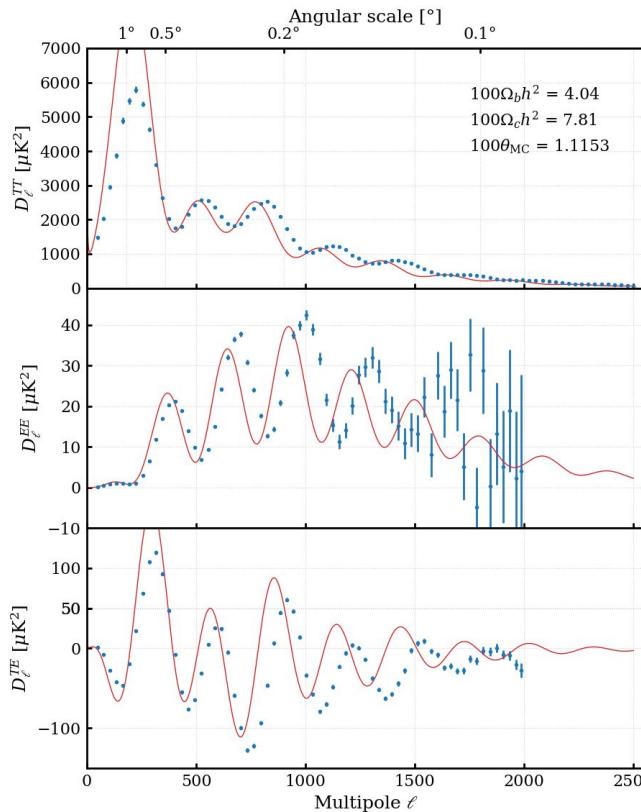
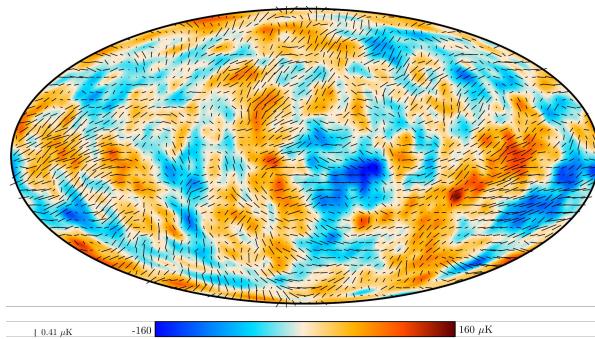


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Polarization E-modes



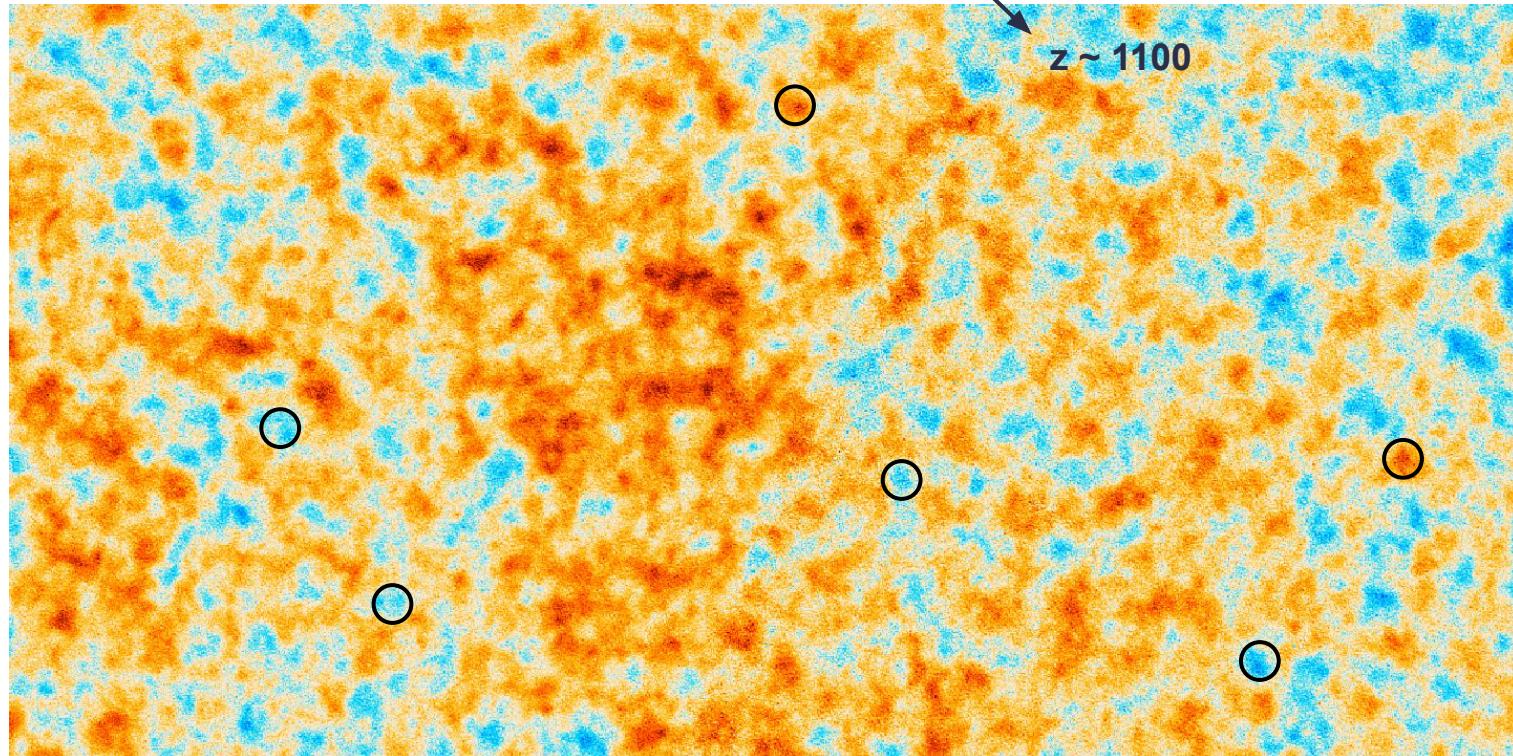
→  $\theta_*$   $\rho_b^0$   $\rho_c^0$

# Measuring $H_0$ from the CMB

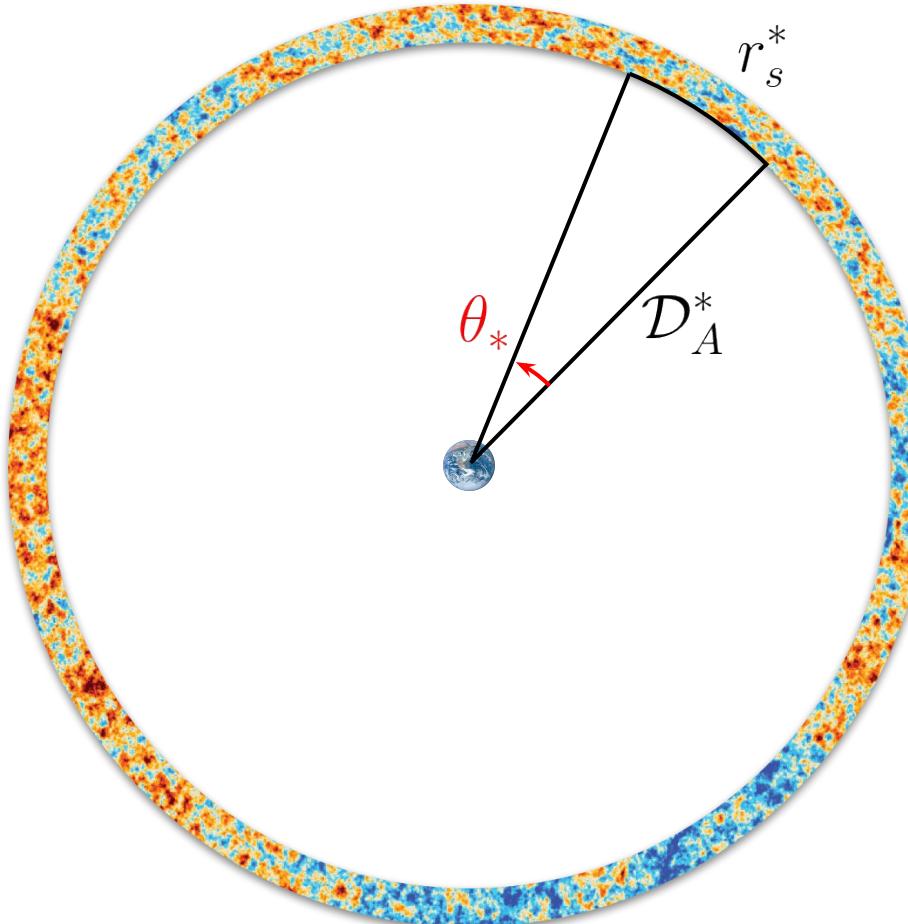
CMB standard ruler : **size of the sound horizon at decoupling** imprinted in the CMB radiation

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CMB standard ruler : size of the sound horizon at decoupling imprinted in the CMB radiation

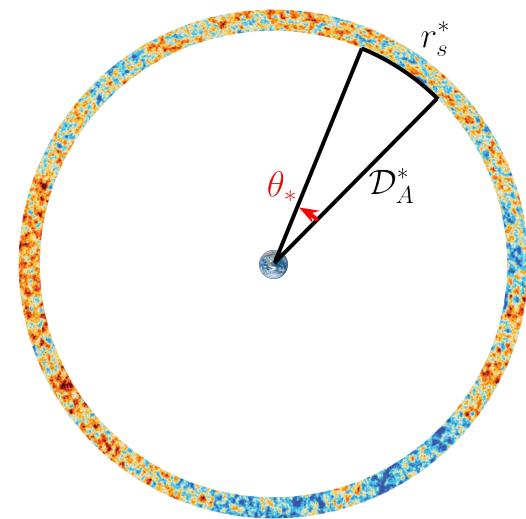


# How to measure $H_0$ from the CMB ?



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$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*} \longrightarrow r_s^* = \int_{z^*}^{\infty} \frac{dz}{H(z)} c_s(z)$$



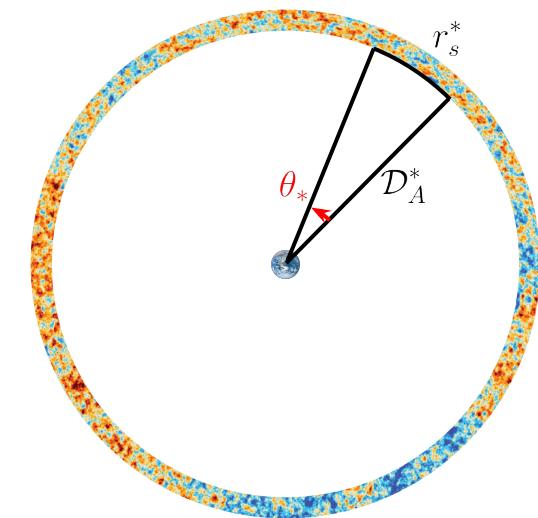
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↓

$$c_s(z) = c \sqrt{\frac{1}{3 [1 + 3\rho_b^0/4\rho_\gamma^0(1+z)^{-1}]}}$$

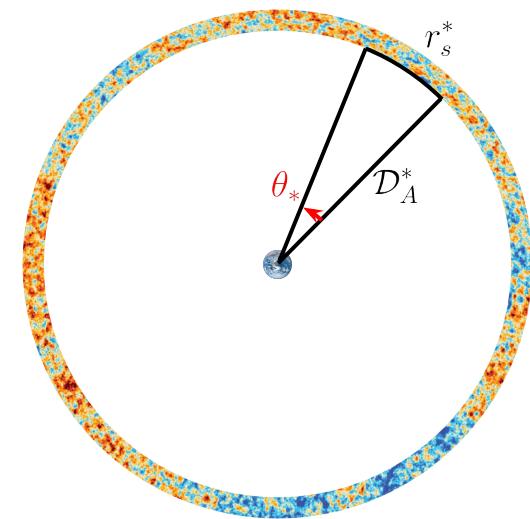
$$H_{\text{early}}^2(z) = \frac{8\pi G}{3} [\rho_r^0(1+z)^4 + (\rho_b^0 + \rho_c^0)(1+z)^3]$$



# How to measure $H_0$ from the CMB ?

Now  $\mathcal{D}_A^*$  is known

$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*}$$



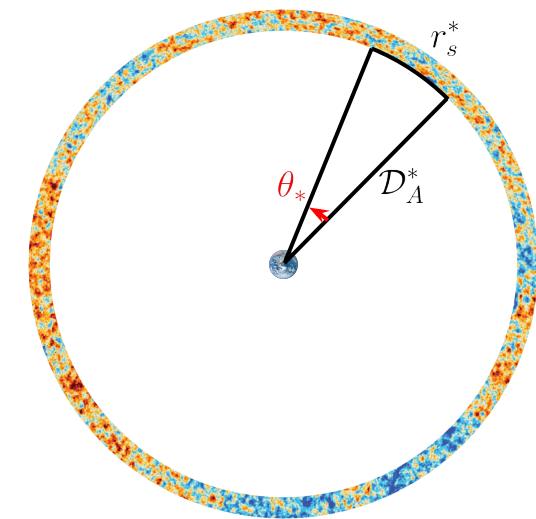
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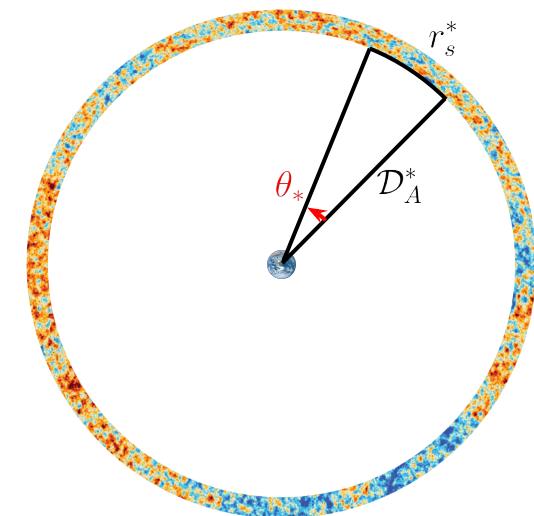
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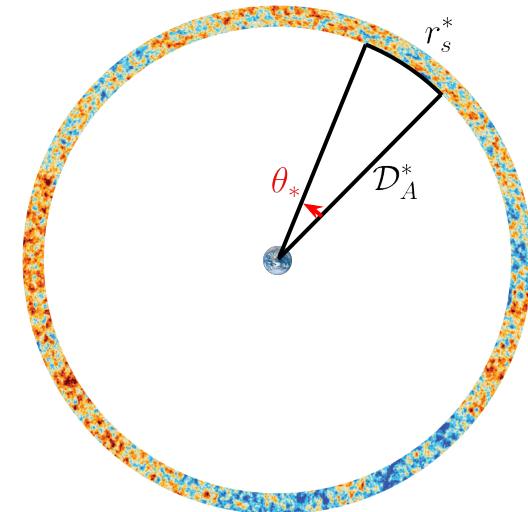
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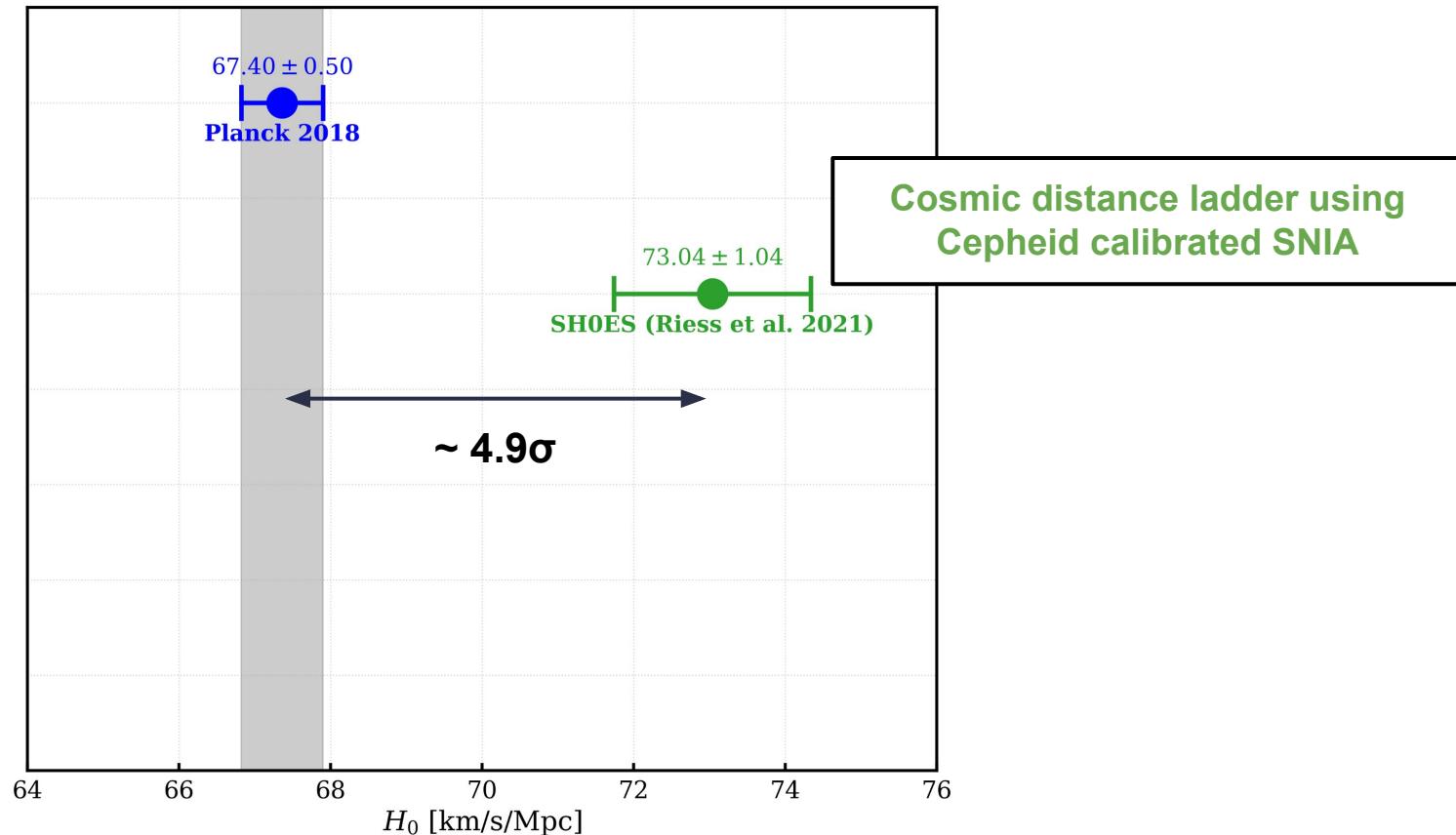
$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*} \quad \rightarrow \quad \mathcal{D}_A^* = c \int_0^{z^*} \frac{dz}{H(z)}$$

$$H_{\text{late}}^2(z) = \frac{8\pi G}{3} [(\rho_b^0 + \rho_c^0)(1+z)^3 + \rho_\Lambda]$$

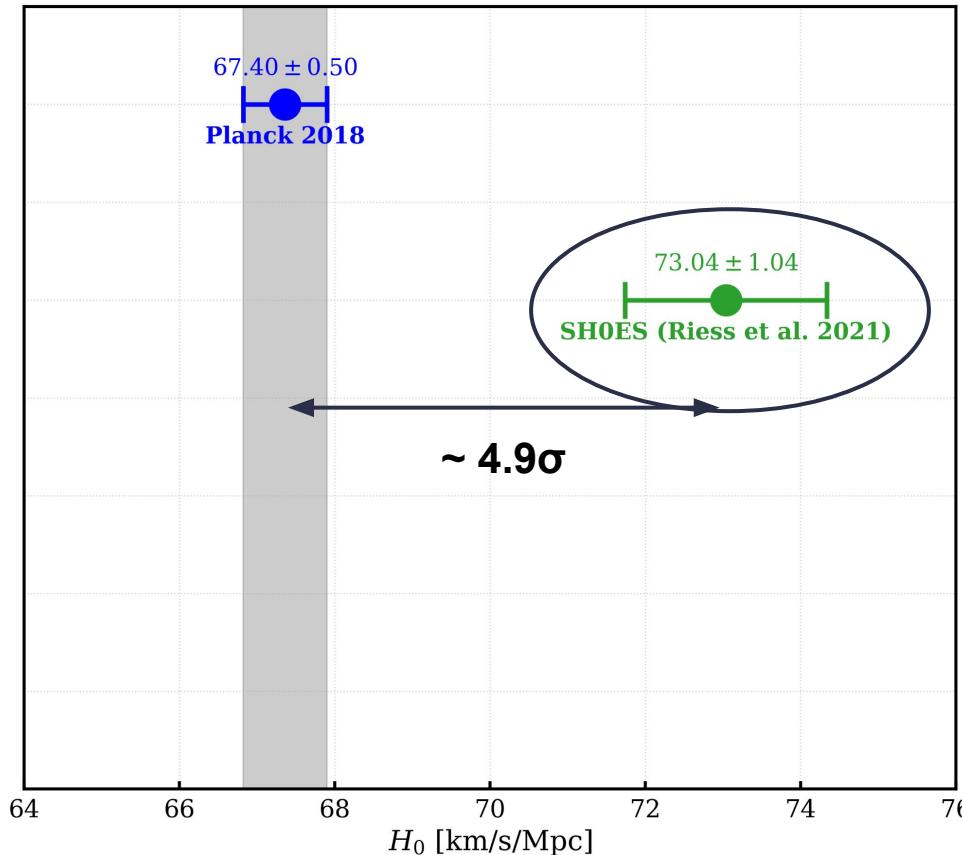
$$H_0^2 = \frac{8\pi G}{3} [\rho_b^0 + \rho_c^0 + \rho_\Lambda]$$



# The Hubble tension as of today



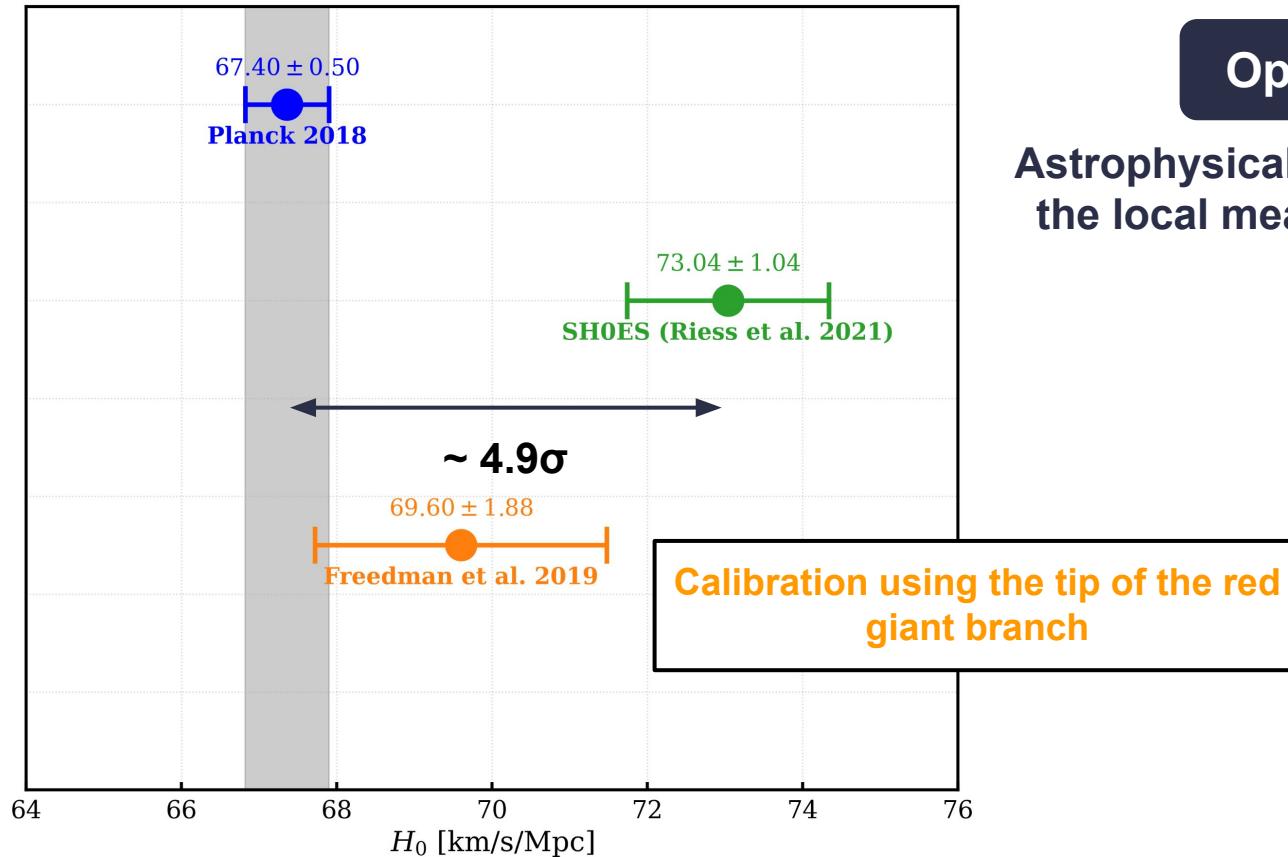
# The Hubble tension as of today



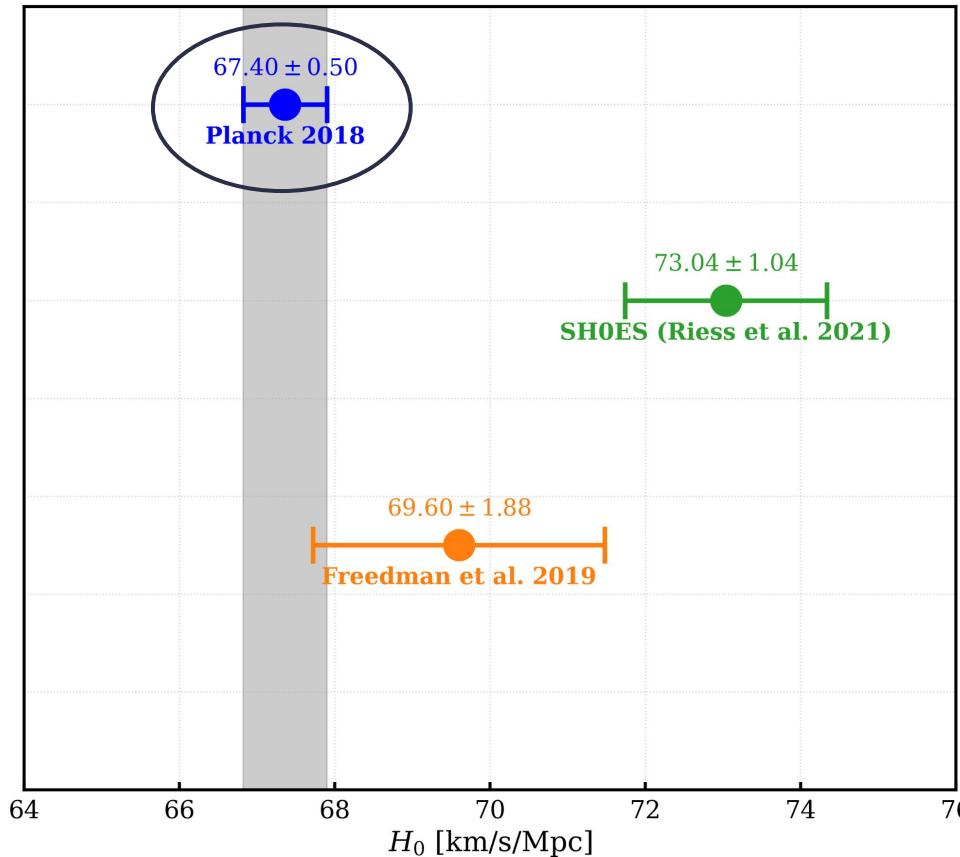
## Option 1

Astrophysical biases affecting the local measurement of  $H_0$

# The Hubble tension as of today



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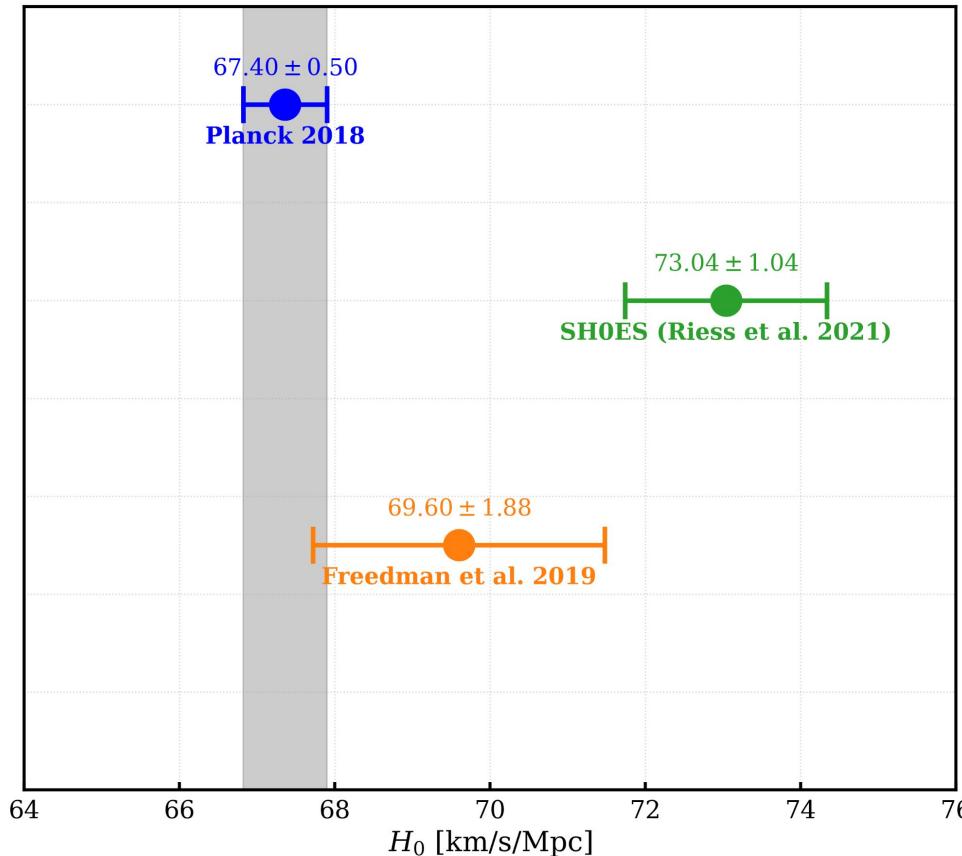
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Astrophysical biases affecting  
the local measurement of  $H_0$

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Instrumental systematic effect  
biasing the value of  $H_0$   
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# The Hubble tension as of today



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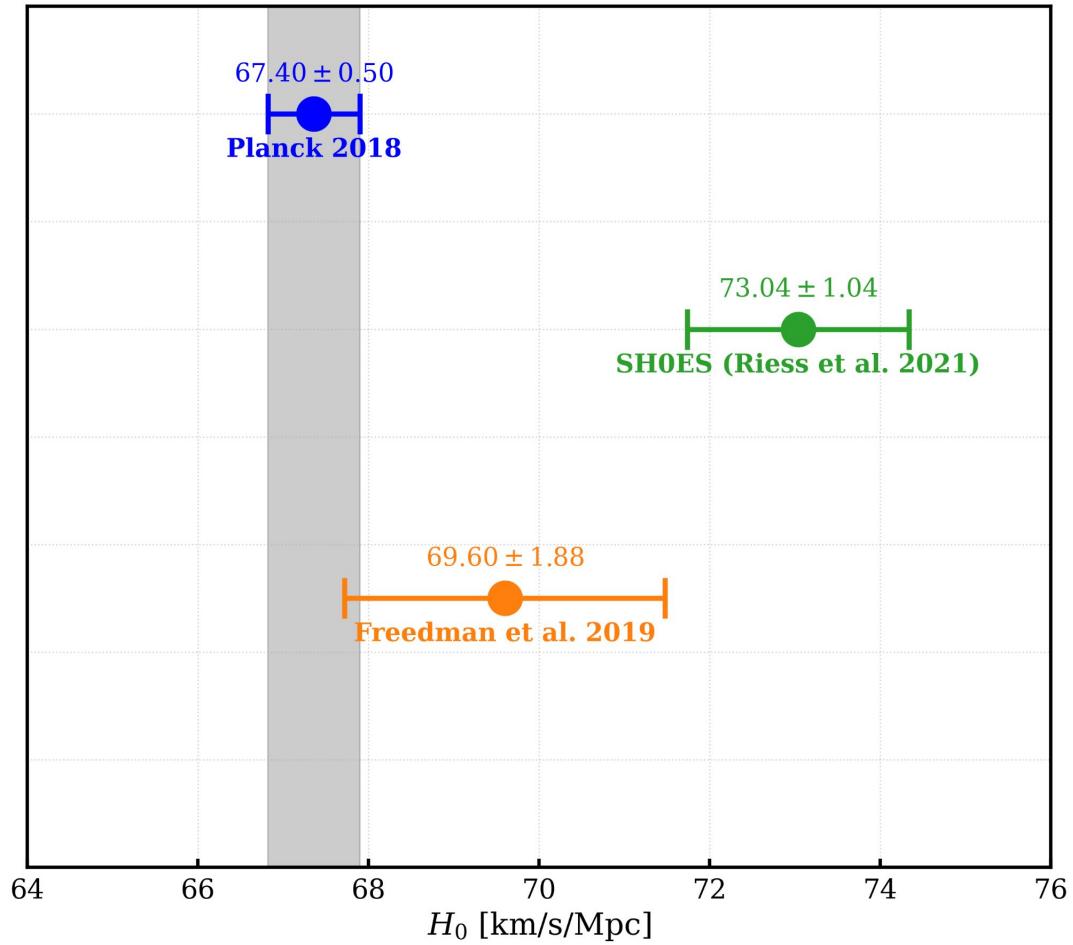
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Physics beyond  $\Lambda$ CDM



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# Independent measurements of $H_0$ from the ground

8



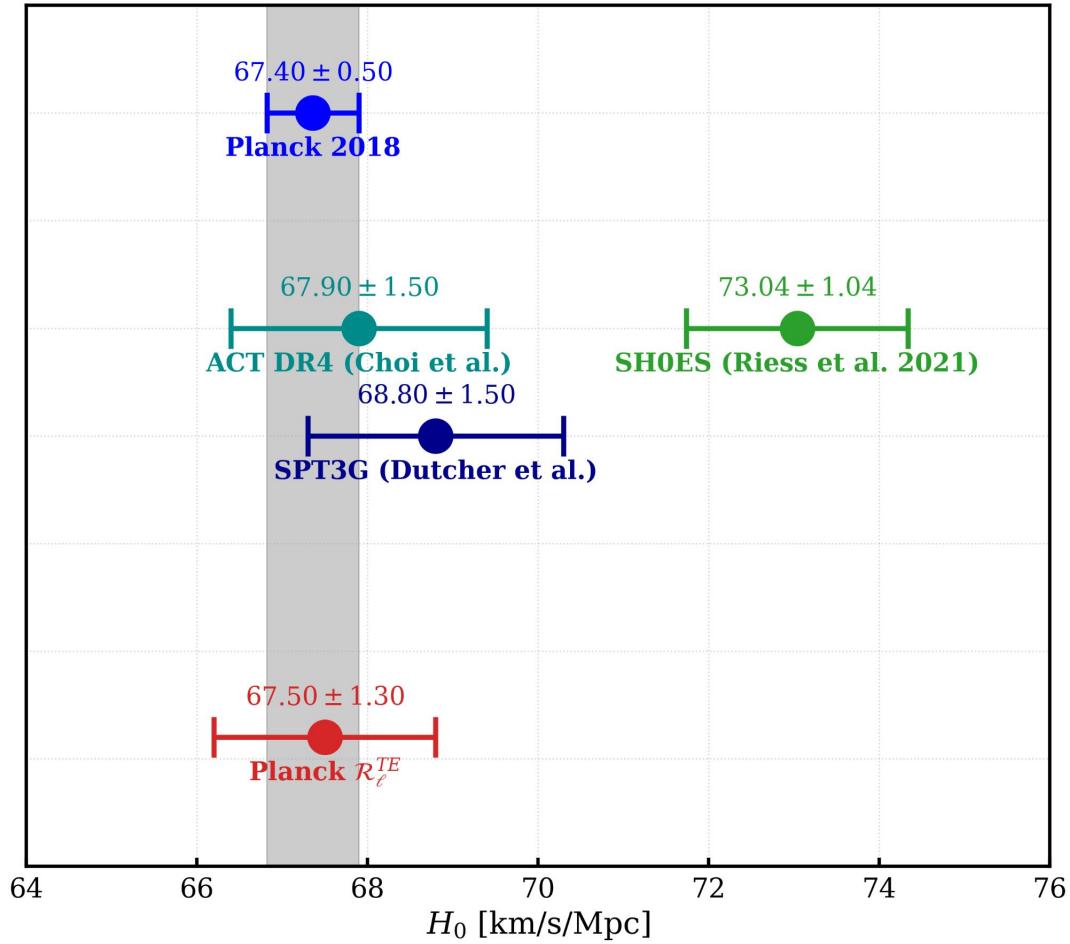
## Option 2

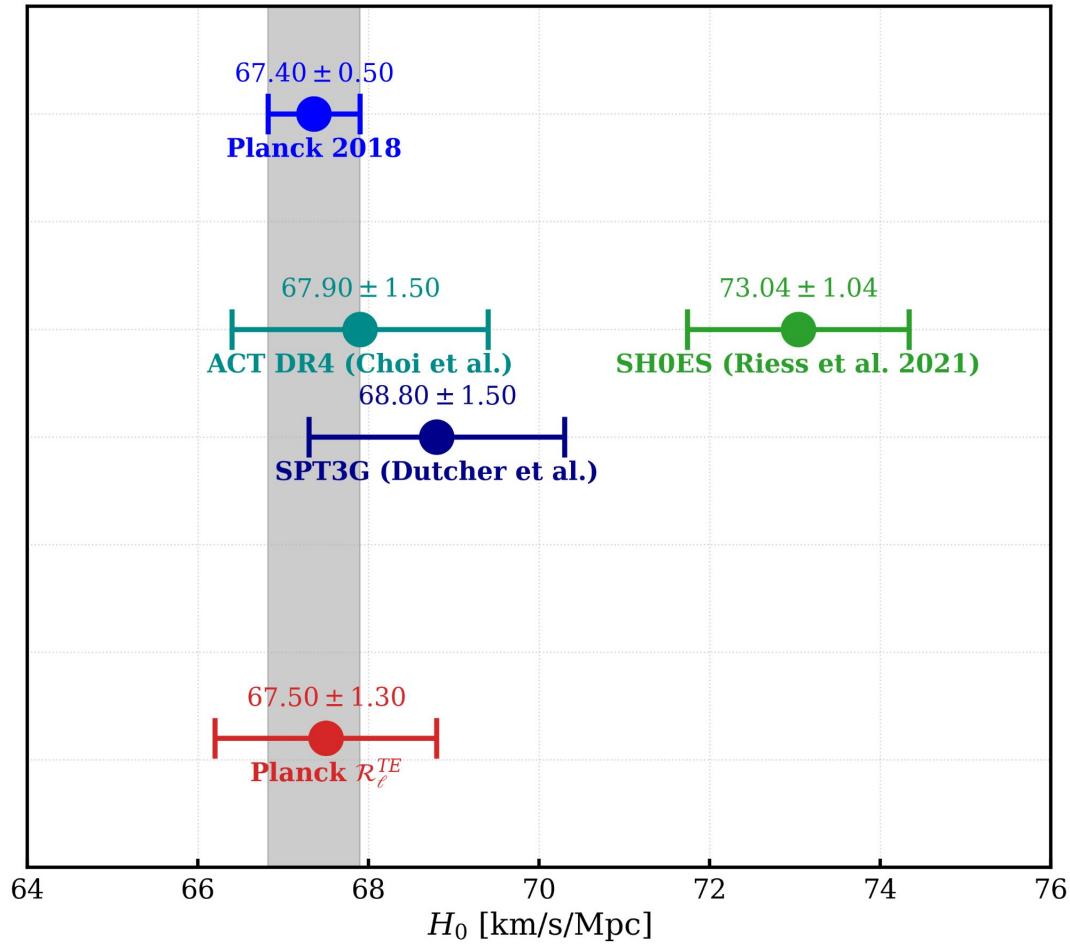
Instrumental systematic effect  
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inferred from the CMB



Hard to shift the CMB inferred  $H_0$   
with a systematic effect :

- Independent measurements from Planck, ACT and SPT
- Constraint from the correlation coefficient, robust against multiplicative systematics





### Option 1

Astrophysical biases affecting the local measurement of  $H_0$

### Option 2

Instrumental systematic effect biasing the value of  $H_0$  inferred from the CMB

### Option 3

Physics beyond  $\Lambda$ CDM

# Early-time modification to $\Lambda$ CDM

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$\downarrow$

$$\frac{3H_{\text{early}}^2(z)}{8\pi G} = \rho_r(z) + \rho_m(z)$$

**Fixed by observations**

# One proposed solution : Early Dark Energy

**Motivation** : obtain a higher value of  $H_0$  from the CMB  $\longrightarrow$  lower  $\mathcal{D}_A^*$

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Fixed by observations

$$\frac{3H_{\text{early}}^2(z)}{8\pi G} = \rho_r(z) + \rho_m(z) + \rho_{\text{EDE}}(z)$$

# EDE phenomenology

The EDE component is described as a scalar field  $\phi$  (Poulin+ 2019, Smith+ 2019)

**Background evolution :**  $\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0$

$$V(\phi) = m^2 f^2 \left[ 1 - \cos \left( \frac{\phi}{f} \right) \right]^3$$

axion-like potential

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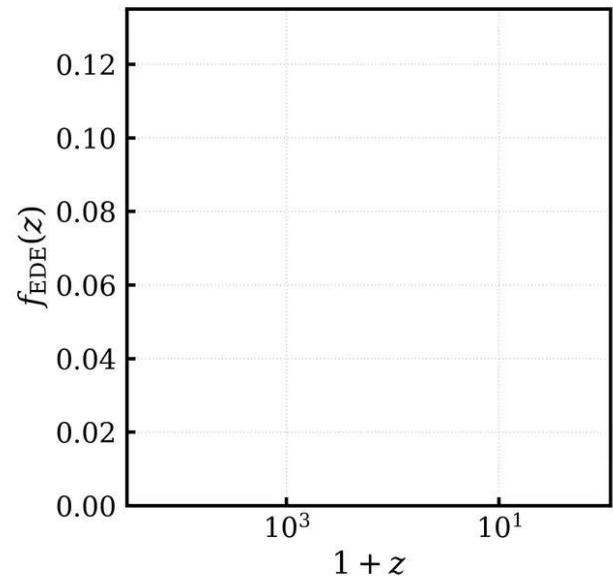
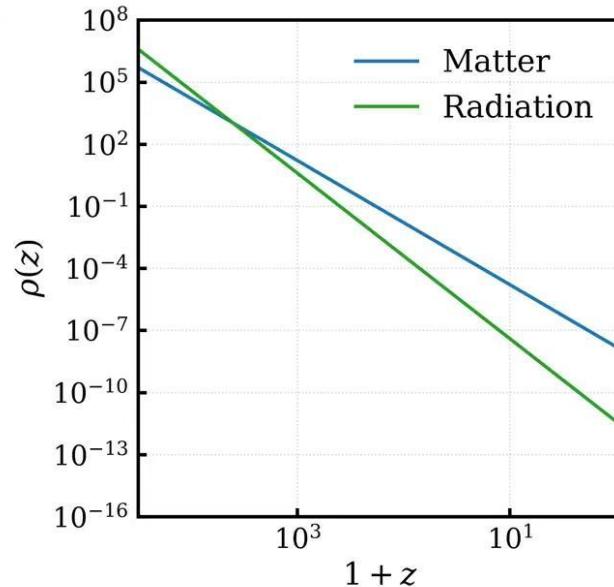
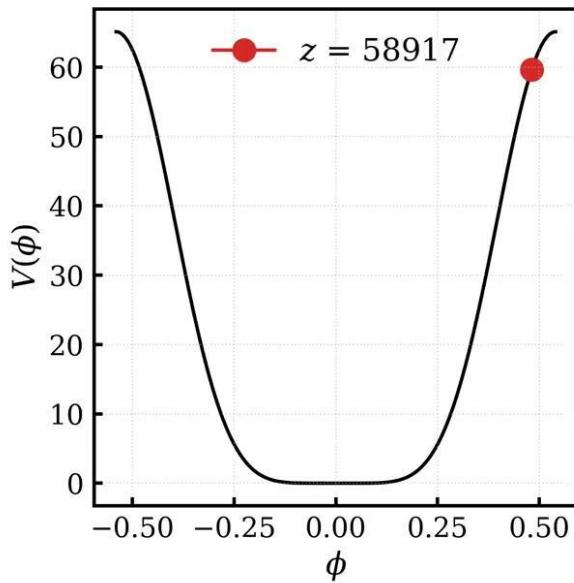
**$\phi_i$**  : initial field value

# Early Dark Energy : frozen at early times

$$\ddot{\phi} + \boxed{3H\dot{\phi}} + V'(\phi) = 0$$

The field is initially frozen due to  
Hubble friction ( $H \gg m$ )

acts as dark energy ( $w = -1$ )

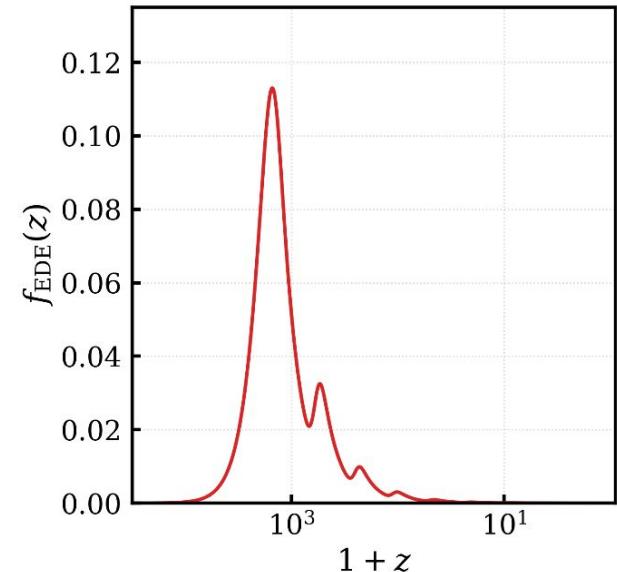
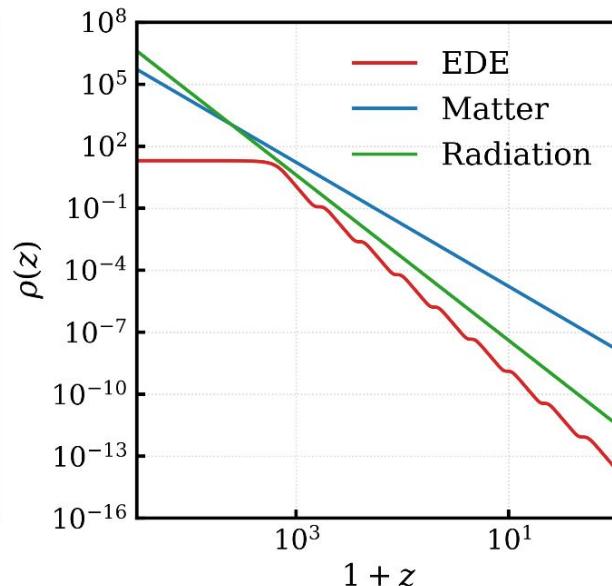
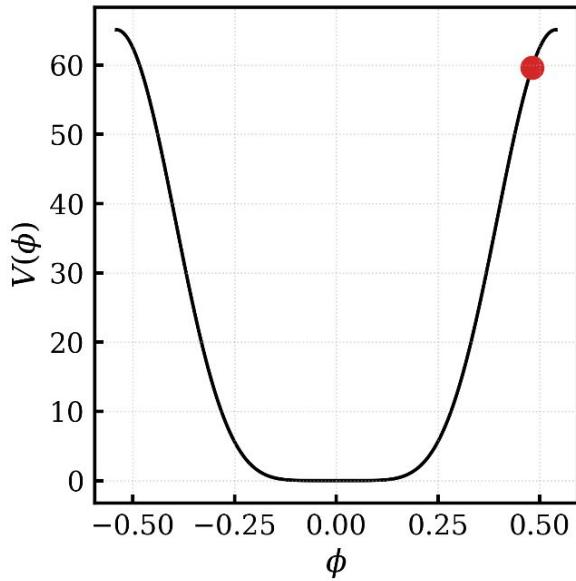


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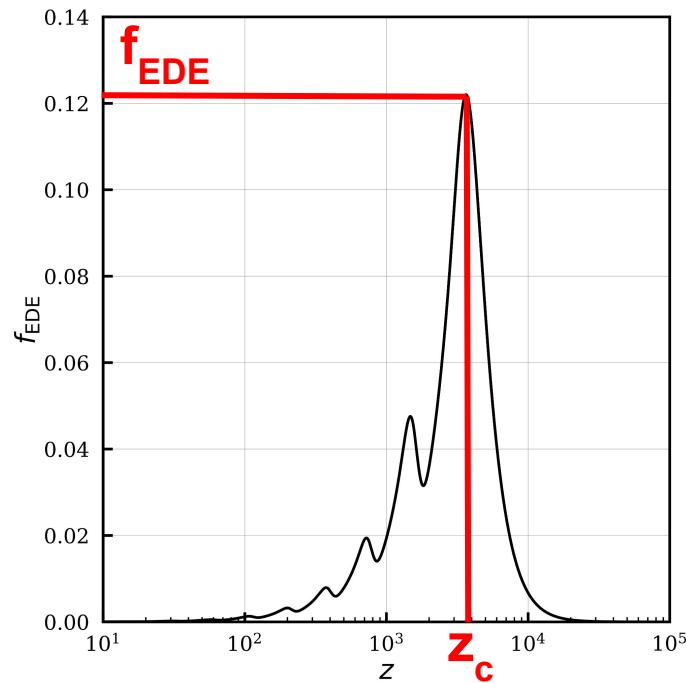


# Early Dark Energy : phenomenological parametrization <sup>12</sup>

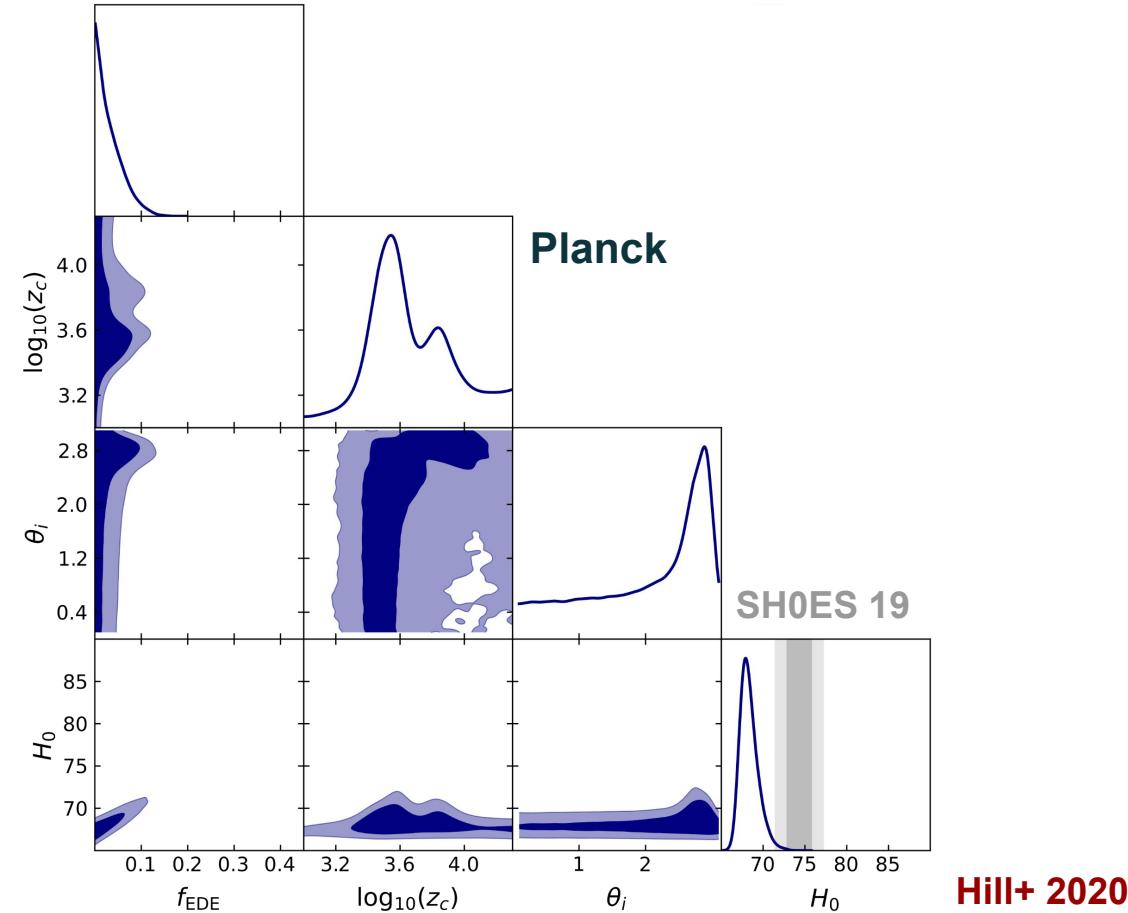
$$(m, f, \phi_i) \longrightarrow (f_{\text{EDE}}, z_c, \phi_i)$$

Field parameters

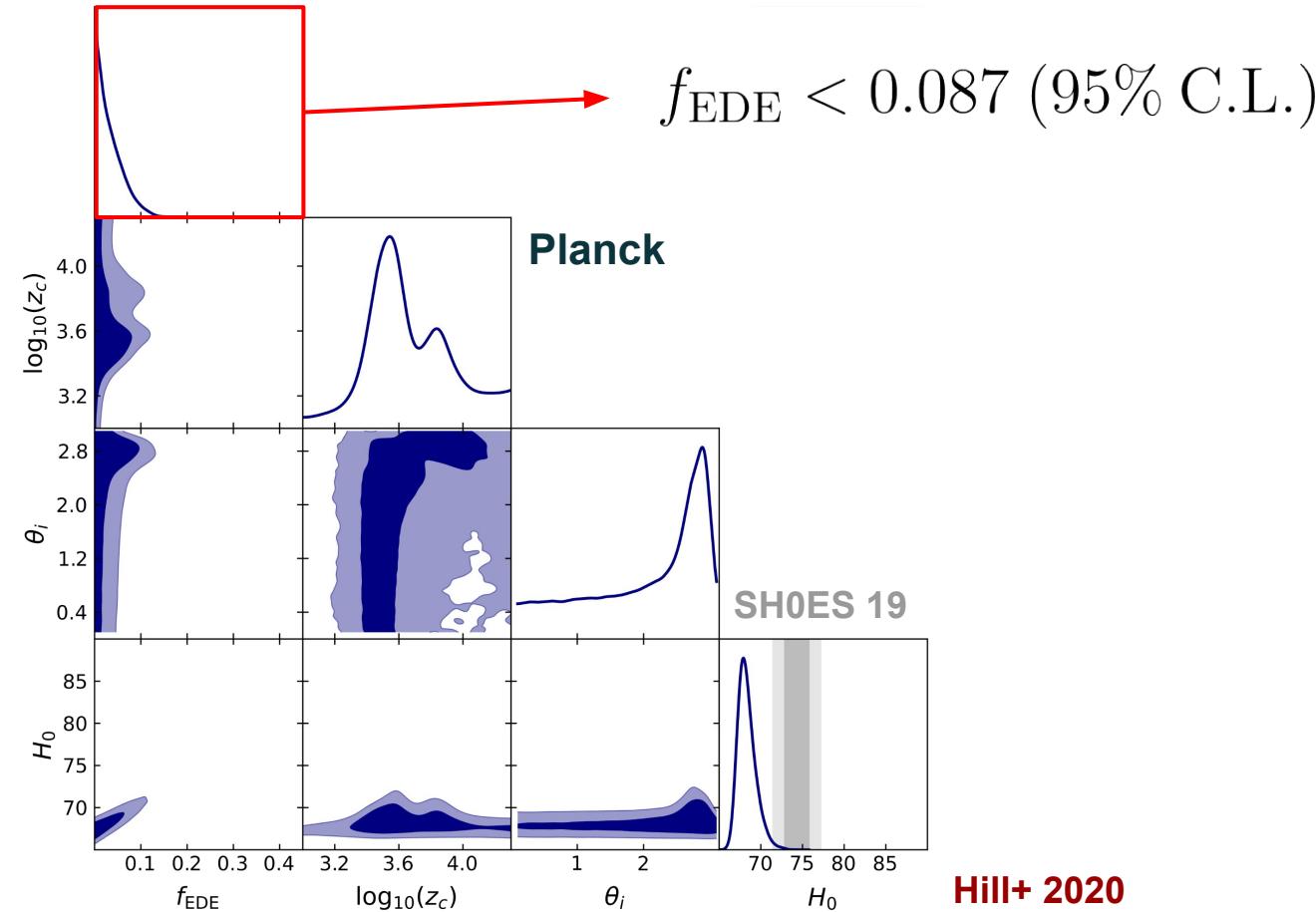
Phenomenological  
parametrization



# Constraints on EDE from Planck data

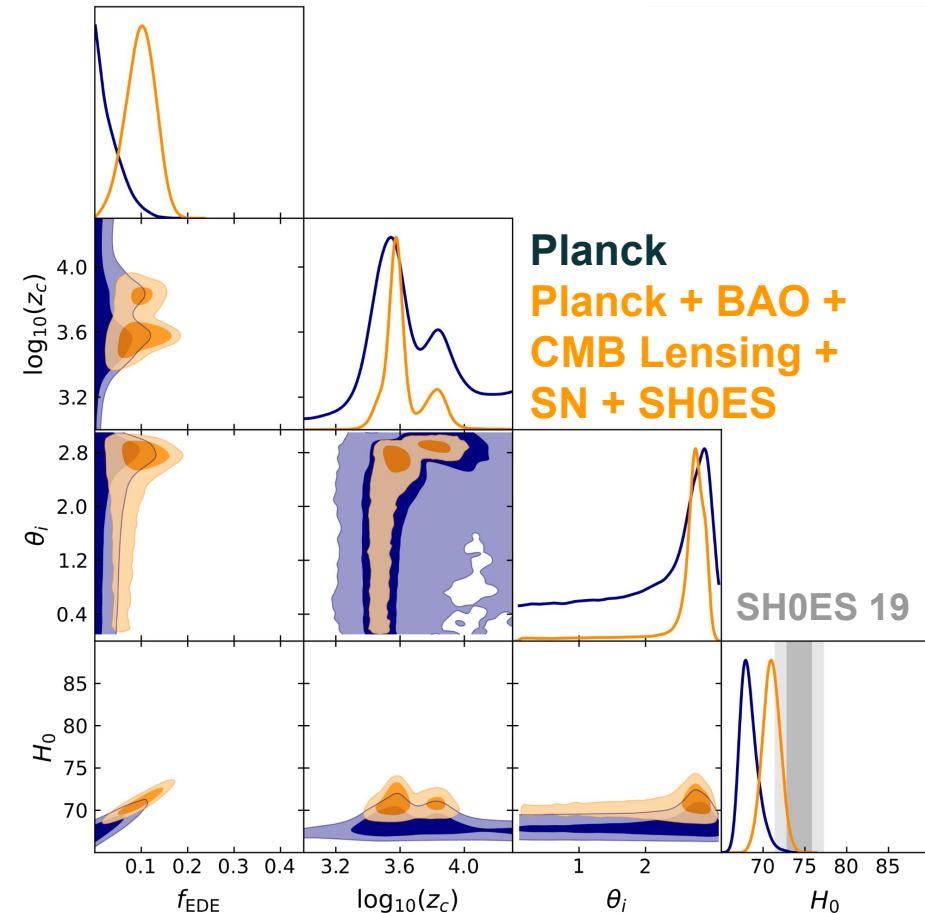


# Constraints on EDE from Planck data



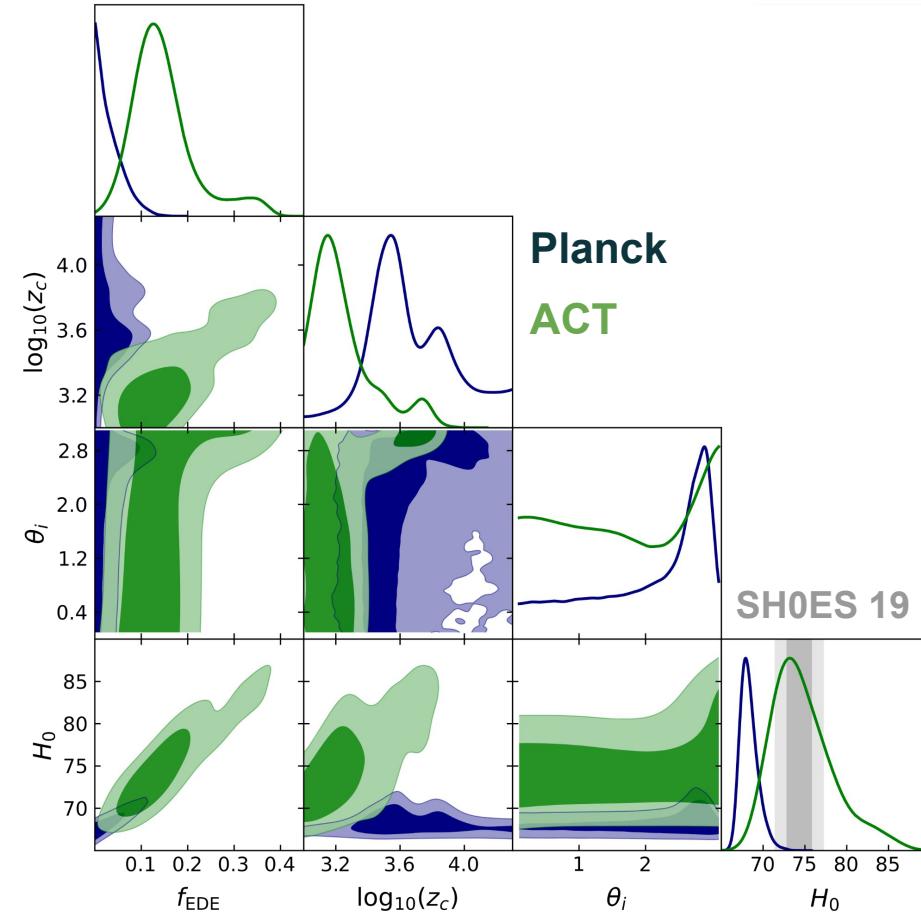
# Results for a combination of Planck and SH0ES data

14



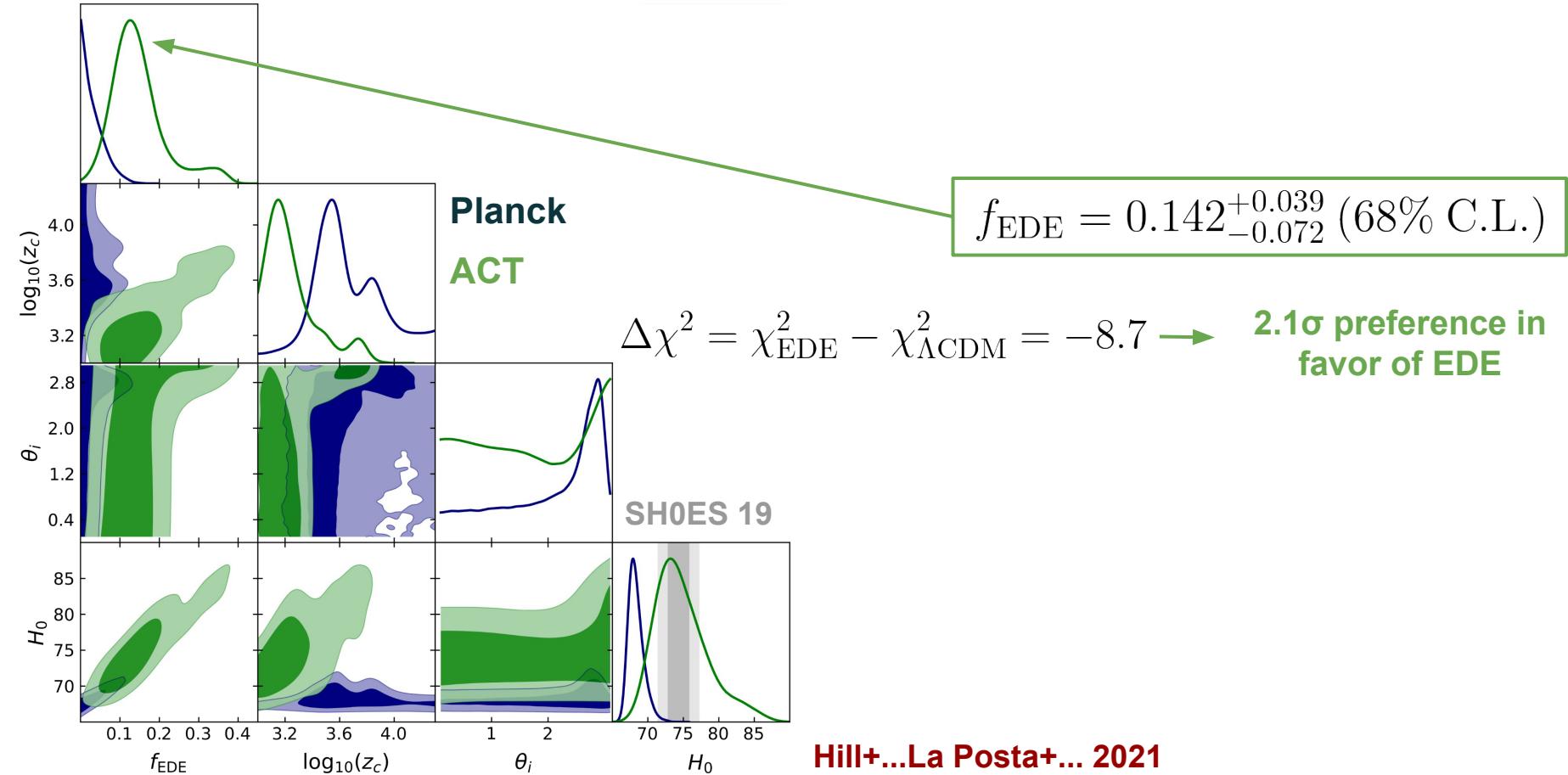
Poulin+ 2019, Smith+ 2019, Hill+ 2020

# Additional constraints from ACTPol



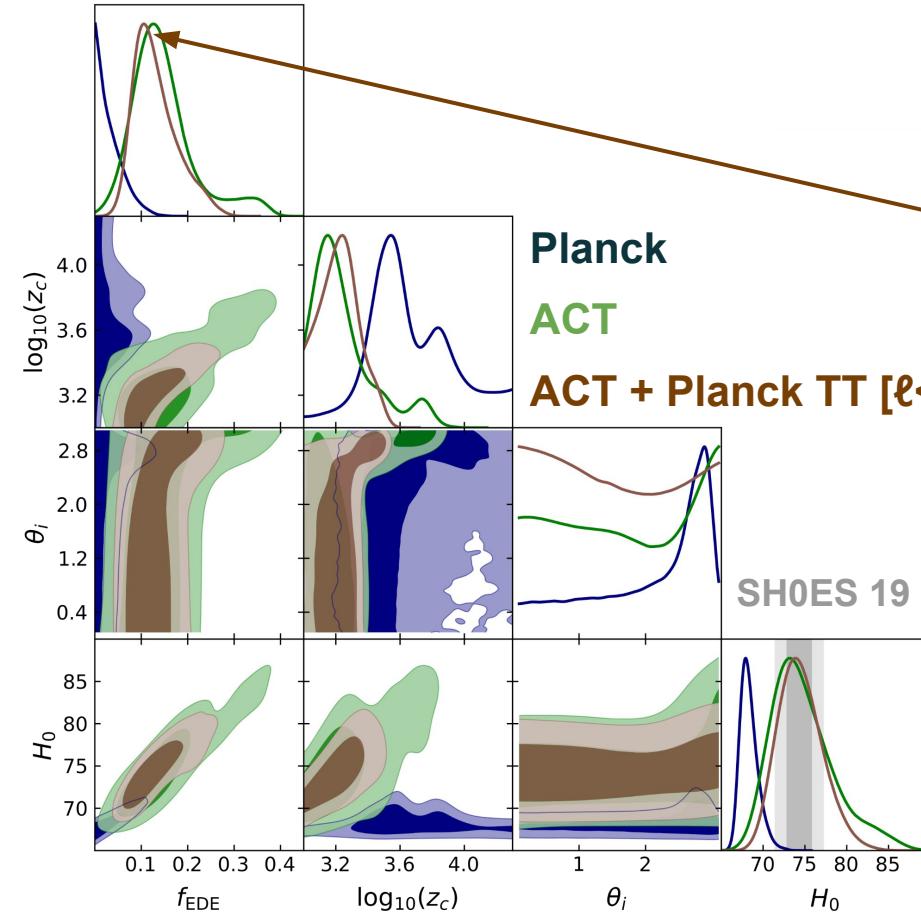
# Additional constraints from ACTPol

15



# Additional constraints from ACTPol

15



$$f_{\text{EDE}} = 0.129^{+0.028}_{-0.055} \text{ (68\% C.L.)}$$

$\Delta\chi^2 = -15.4 \rightarrow$  **3.2 $\sigma$  preference in favor of EDE**

Hill+...La Posta+... 2021

# Summary of EDE results

- Planck data alone don't favor high  $f_{\text{EDE}}$  values (Hill+ 2020)
- Planck data in combination with SH0ES show a preference for non-zero  $f_{\text{EDE}}$  (Poulin+ 2019, Smith+ 2019)
- ACT data alone favors EDE over  $\Lambda$ CDM (Hill+...La Posta+... 2021)

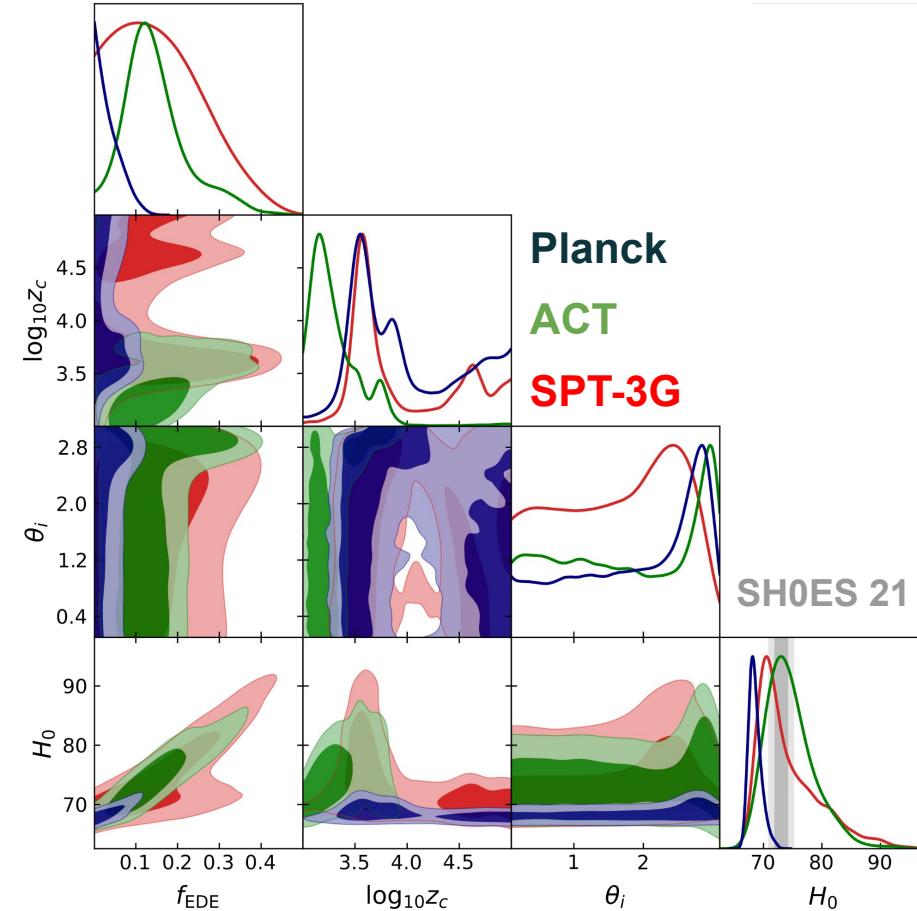
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→ Motivates an analysis of EDE with public SPT-3G data

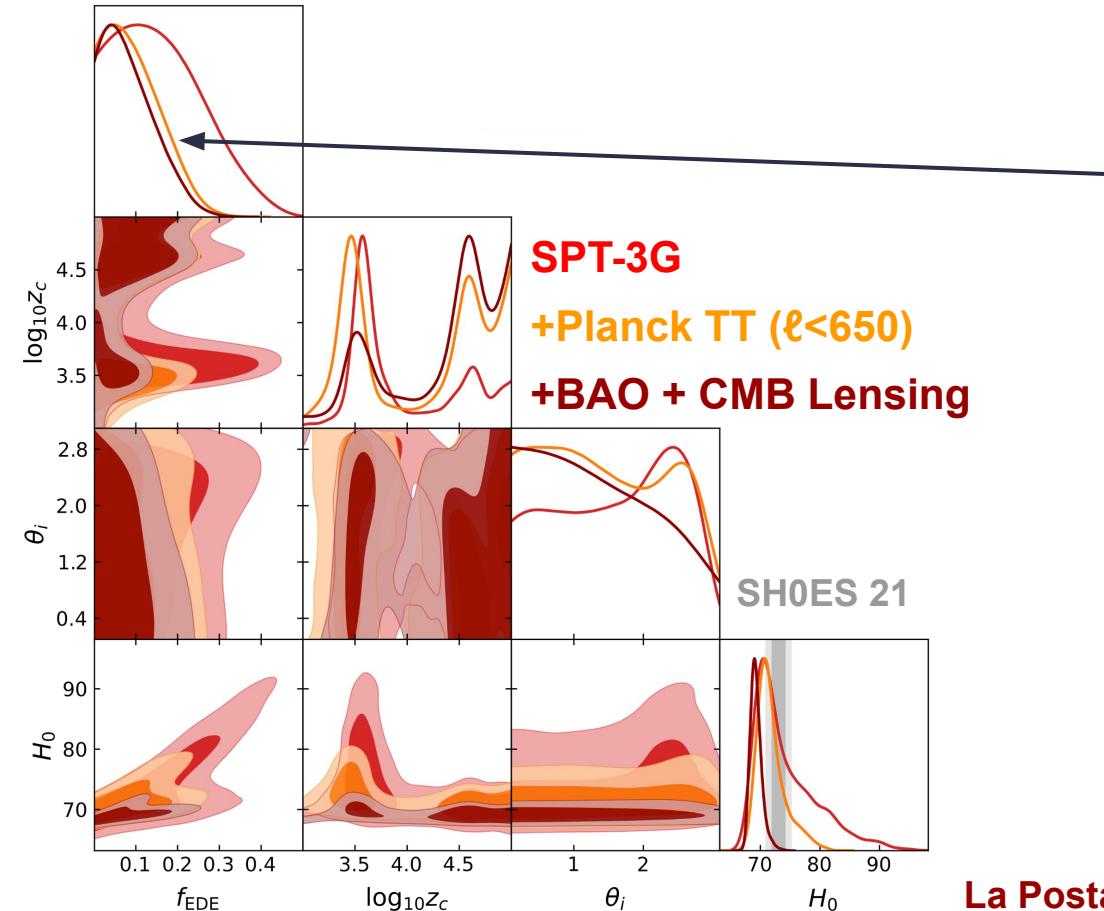
# Constraints from SPT-3G public data

17



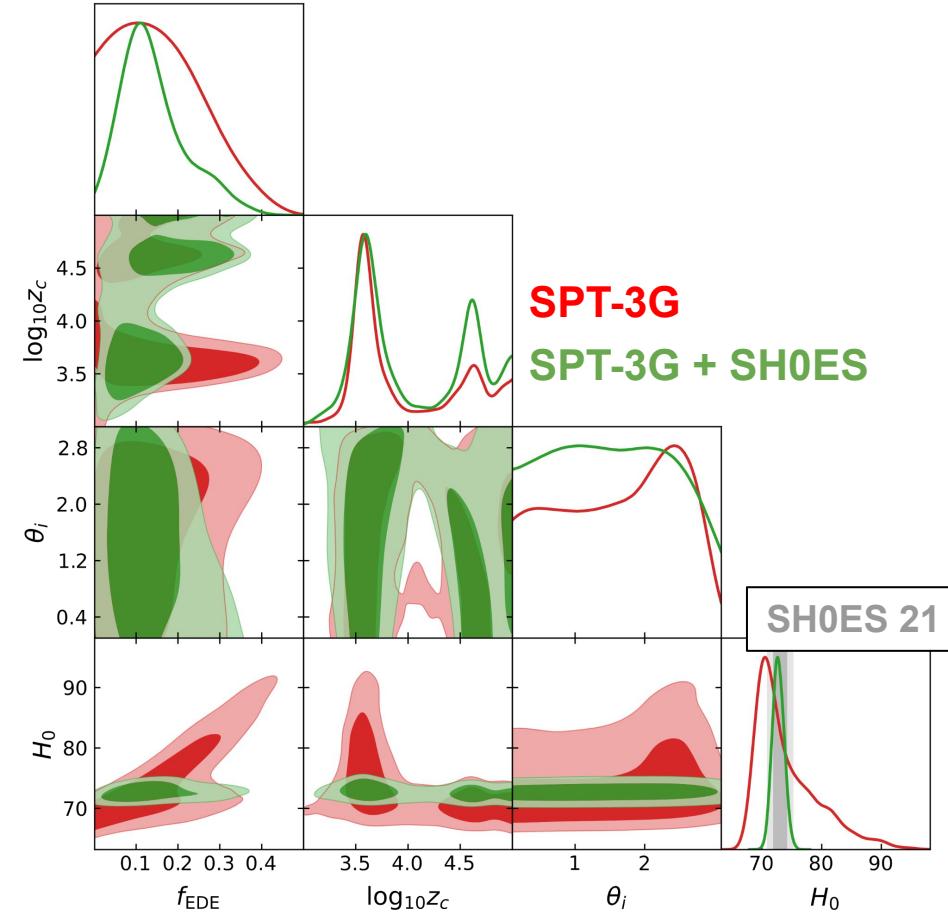
# Constraints from SPT-3G public data

17



We tighten the constraint on  $f_{\text{EDE}}$  when we combine SPT3G and Planck TT ( $\ell < 650$ ) or when we add LSS probes

# Combining with SH0ES constraint



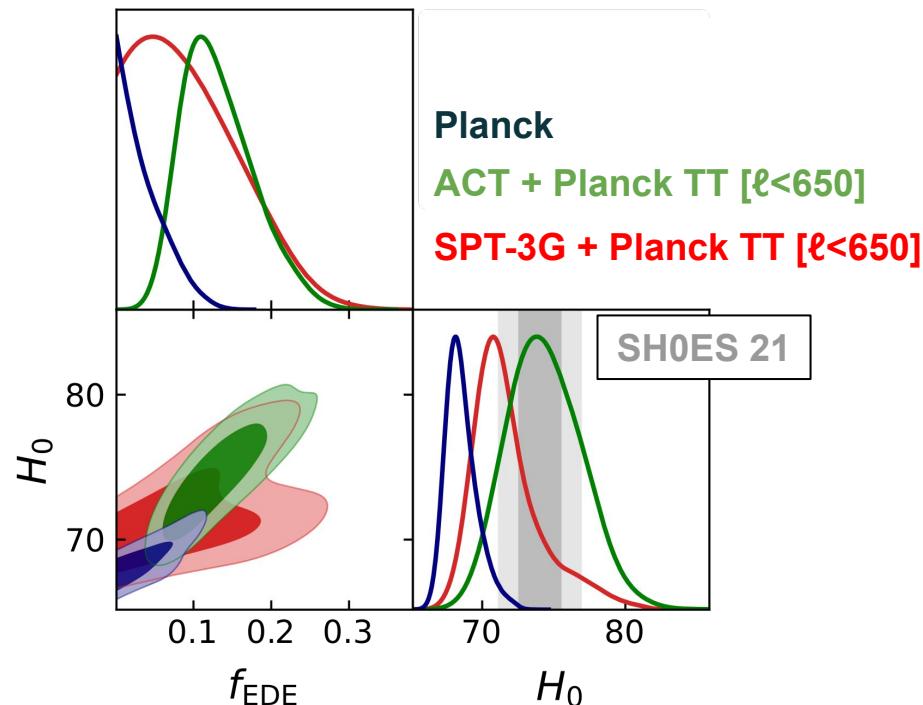
$$\Delta\chi^2_{\text{SPT-3G}} = -6.3$$

improvement of the fit to  
SPT-3G data (with  
respect to  $\Lambda$ CDM)

# Conclusions

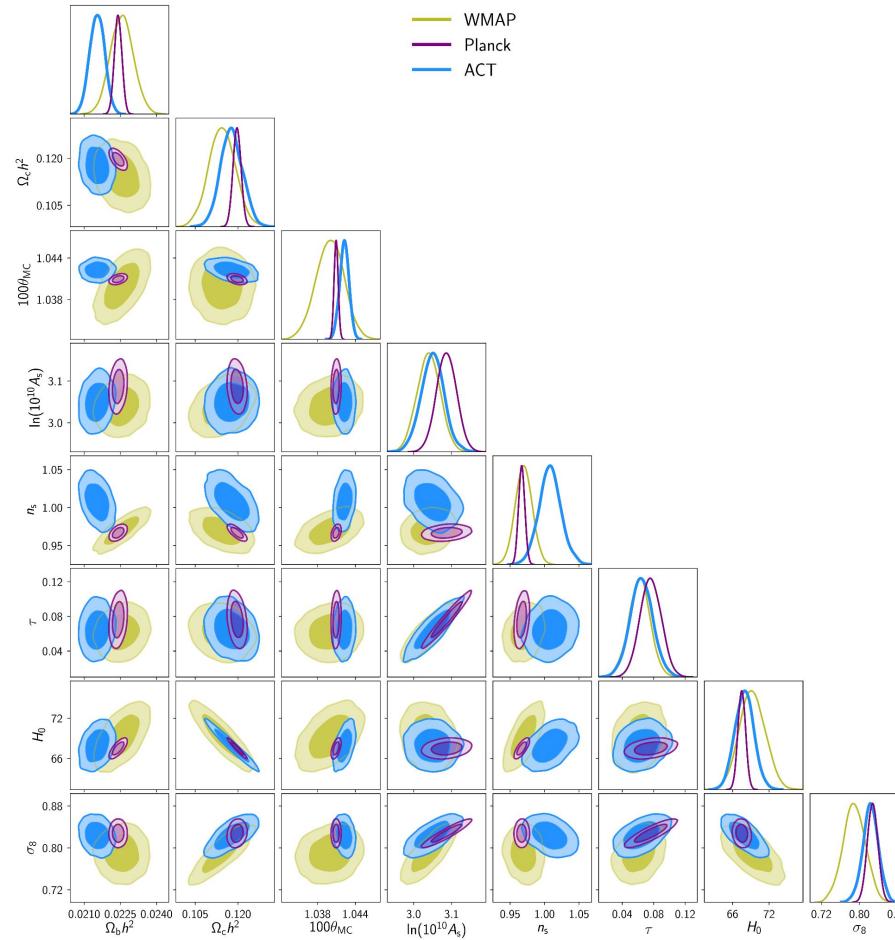
19

- Planck data alone do not favor high  $f_{\text{EDE}}$  values
- Planck + SH0ES show a preference for  $f_{\text{EDE}} \sim 10\%$
- ACT DR4 data favors EDE over  $\Lambda\text{CDM}$  (with  $f_{\text{EDE}} \sim 10\%$ )
- SPT-3G is not as constraining as ACT and Planck : but sees some degree of EDE when combined with SH0ES

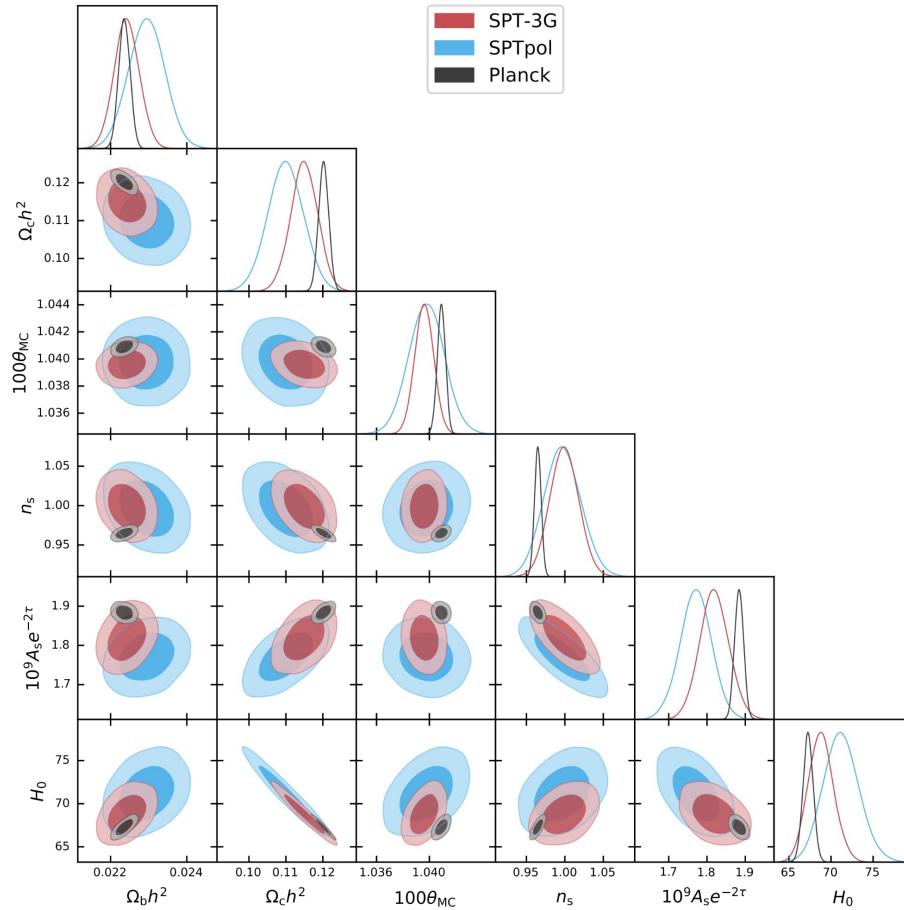


# Extra-slides

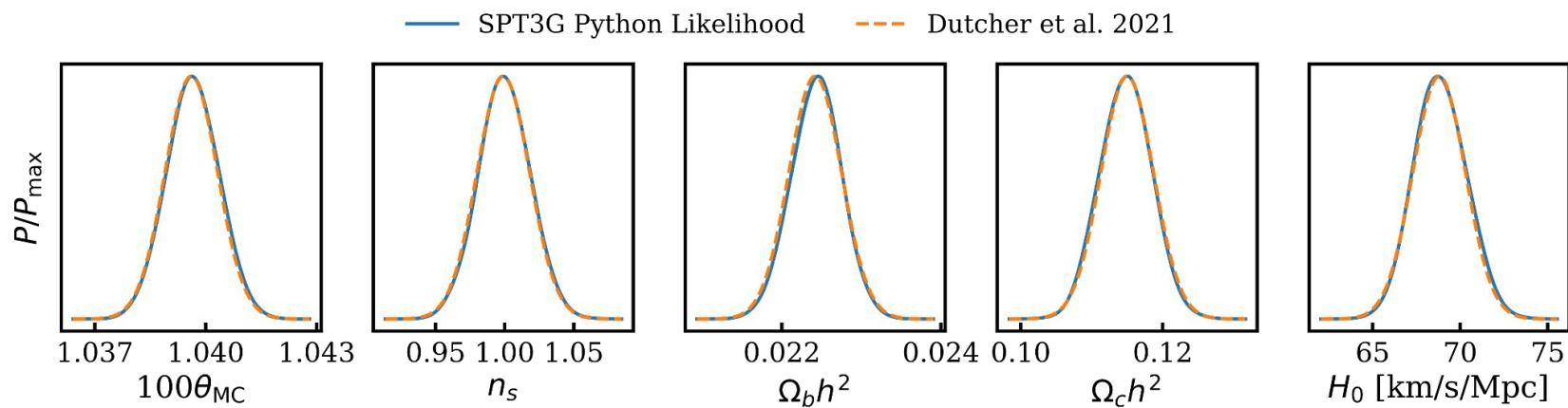
# ACTPol cosmology



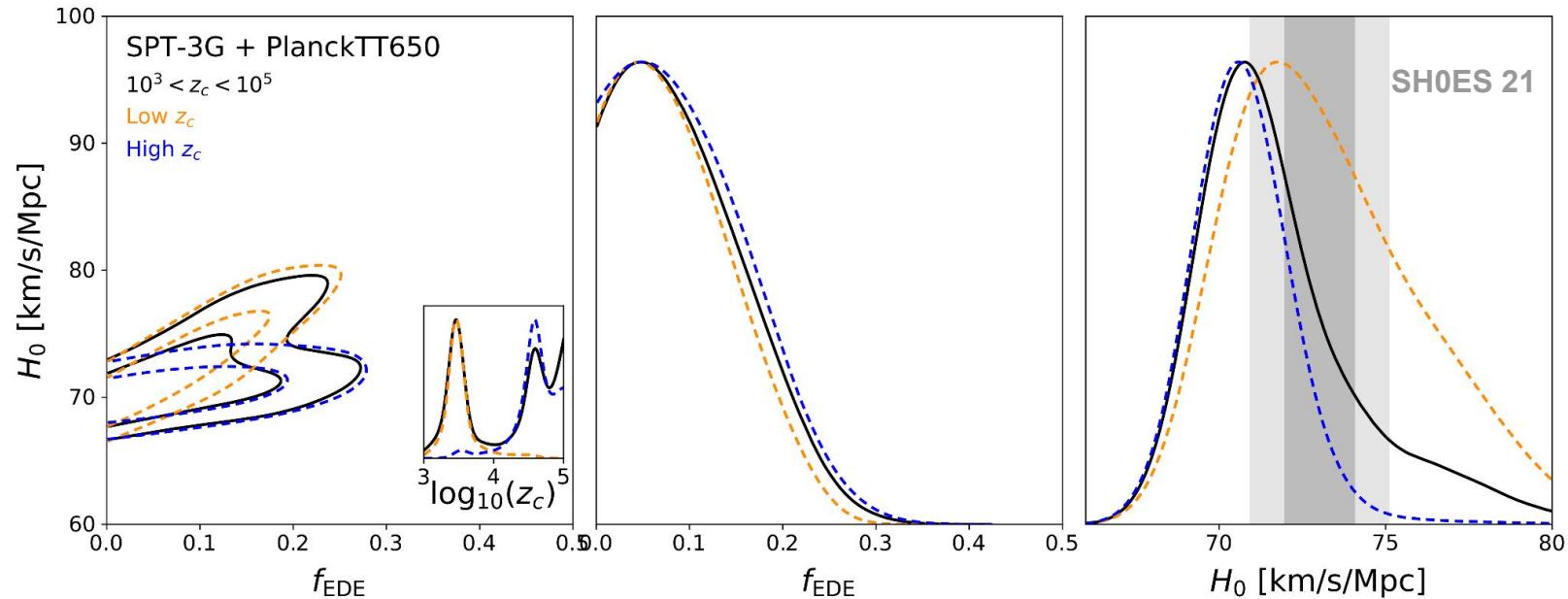
# SPT-3G cosmology



# SPT-3G python implementation

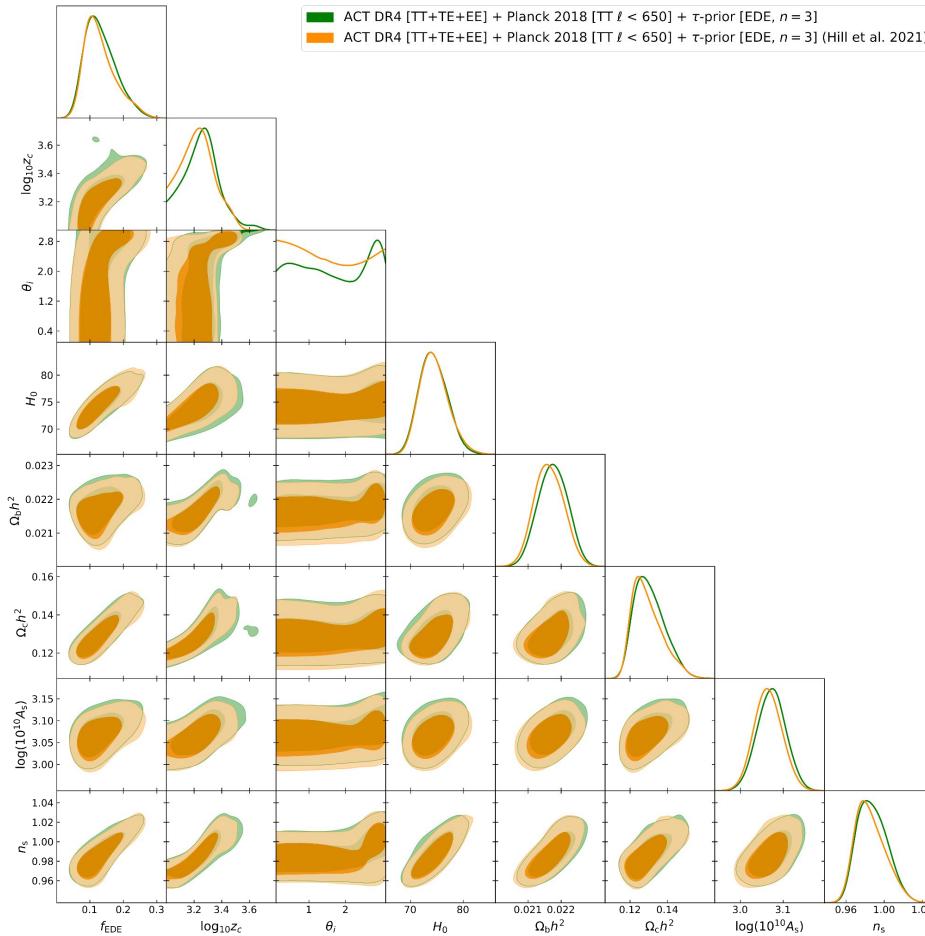


# Impact of the $z_c$ prior

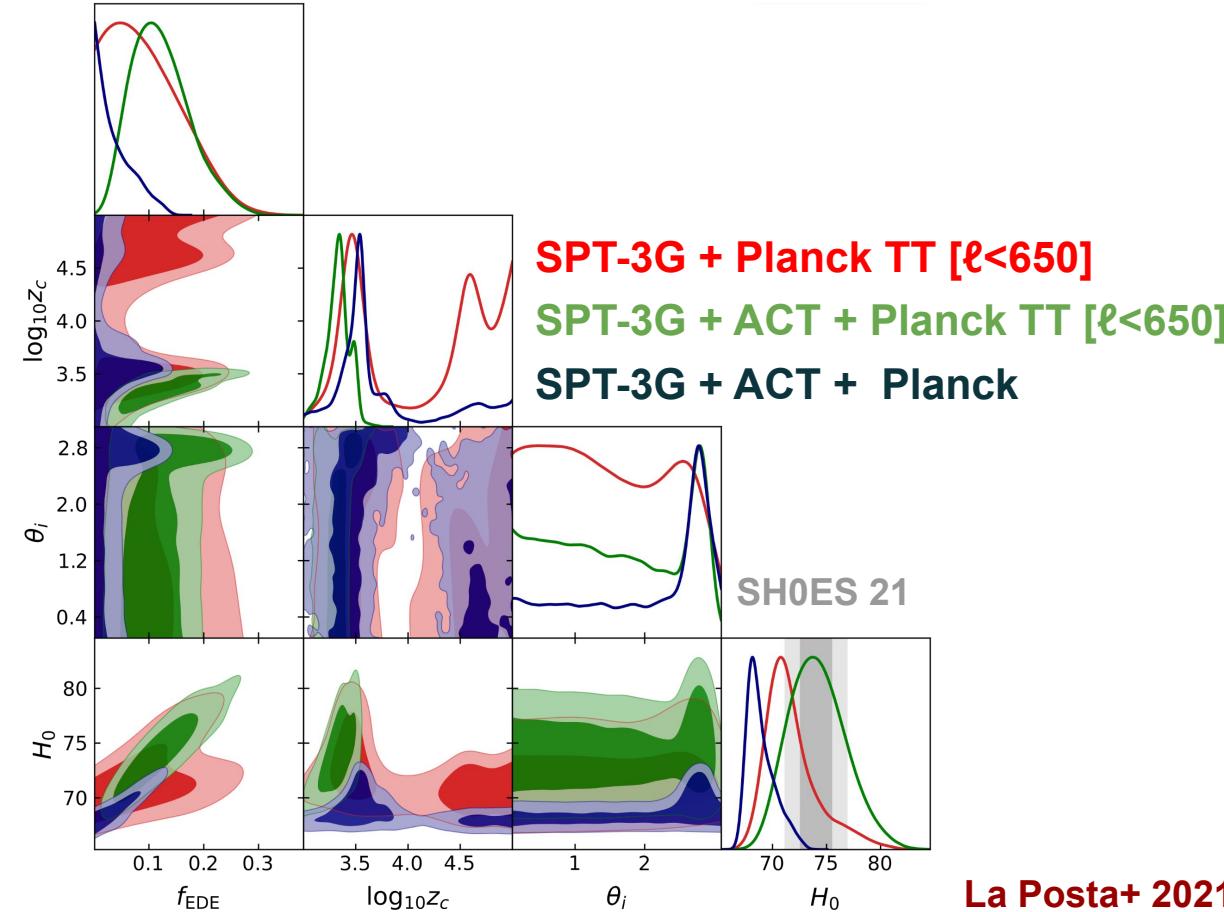


La Posta+ 2021 [arXiv:2112.10754]

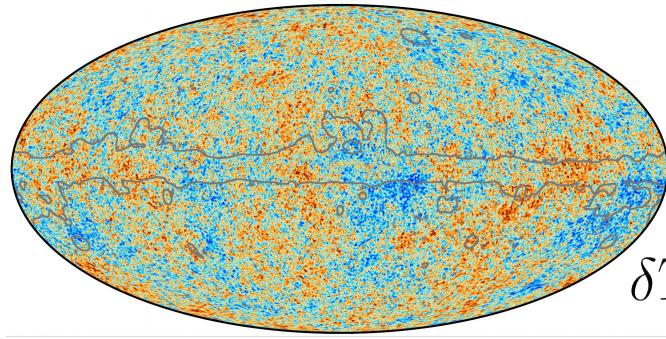
# CAMB/CLASS EDE models



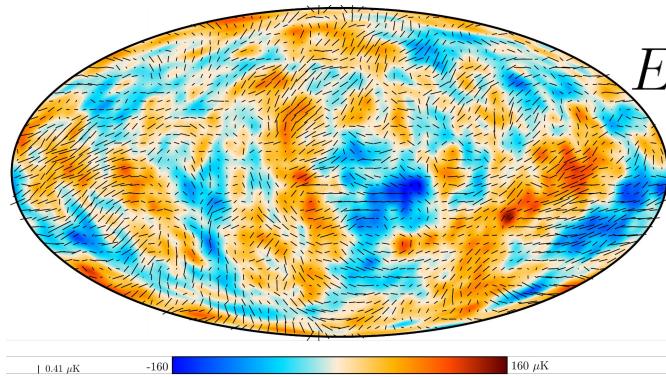
# Combining with other CMB datasets



# Systematics in the CMB



$$\delta T^{\text{sky}}(\hat{n}) \xrightarrow{\mathcal{I}_{T/E}} E^{\text{sky}}(\hat{n})$$



$$\begin{aligned} \delta T^{\text{obs}}(\hat{n}) &= \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n}) \\ E^{\text{obs}}(\hat{n}) &= \mathcal{I}_E * E^{\text{sky}}(\hat{n}) \end{aligned}$$

# Systematics in the CMB

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

$$\mathcal{I}_T = \mathcal{F}_T * c * B_T$$

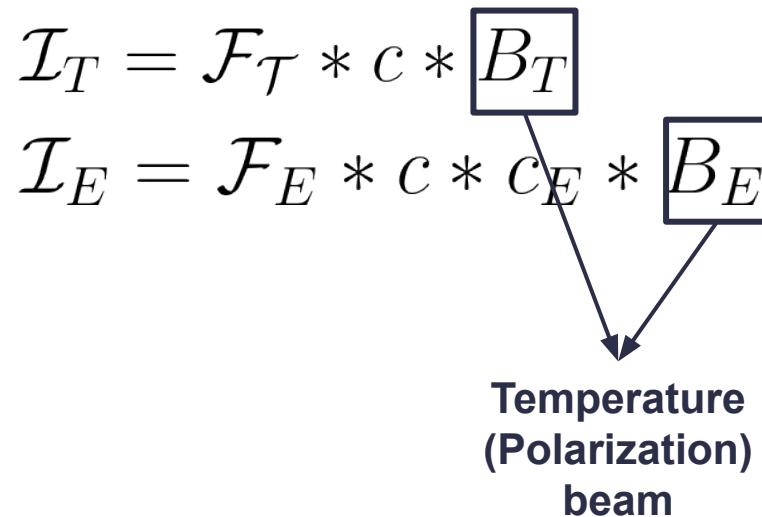
$$\mathcal{I}_E = \mathcal{F}_E * c * c_E * B_E$$

# Systematics in the CMB

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)



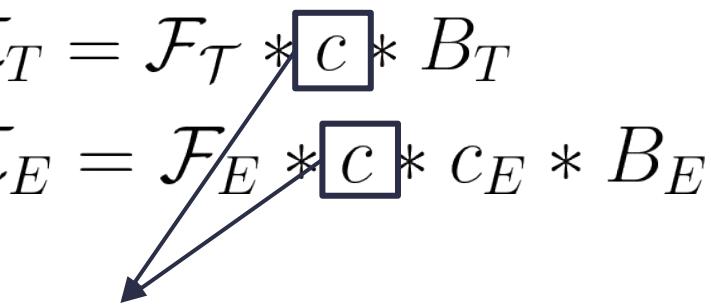
# Systematics in the CMB

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)
- Calibration

$$\mathcal{I}_T = \mathcal{F}_T * \boxed{c} * B_T$$
$$\mathcal{I}_E = \mathcal{F}_E * \boxed{c} * c_E * B_E$$



Global  
calibration

# Systematics in the CMB

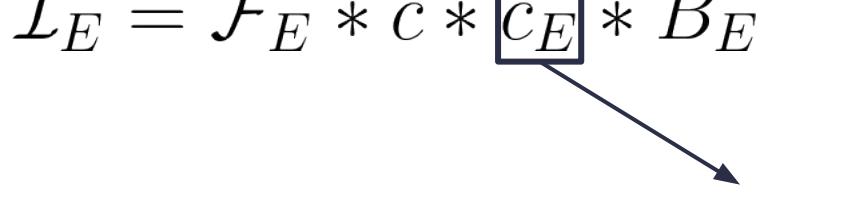
$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)
- Calibration
- Polarization efficiency

$$\mathcal{I}_T = \mathcal{F}_T * c * B_T$$

$$\mathcal{I}_E = \mathcal{F}_E * c * \boxed{c_E} * B_E$$



Polarization  
efficiency

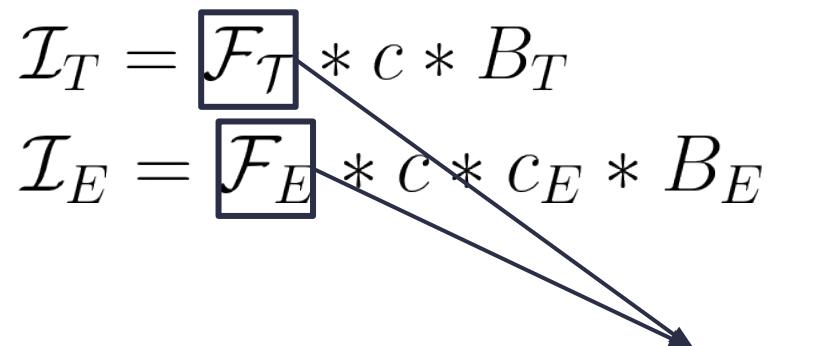
# Systematics in the CMB

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)
- Calibration
- Polarization efficiency
- Transfer functions (map-making)

$$\mathcal{I}_T = \boxed{\mathcal{F}_T} * c * B_T$$
$$\mathcal{I}_E = \boxed{\mathcal{F}_E} * c * c_E * B_E$$



Temperature  
(polarization)  
transfer function

# Systematics in the CMB

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)
- Calibration
- Polarization efficiency
- Transfer functions (map-making)

These instrumental effects are multiplicative in harmonic space

$$C_{\ell}^{TT,\text{obs}} = (\mathcal{F}_{\ell}^T)^2 c^2 (B_{\ell}^T)^2 C_{\ell}^{TT}$$

$$C_{\ell}^{EE,\text{obs}} = (\mathcal{F}_{\ell}^E)^2 c^2 c_E^2 (B_{\ell}^E)^2 C_{\ell}^{EE}$$

$$C_{\ell}^{TE,\text{obs}} = \mathcal{F}_{\ell}^T \mathcal{F}_{\ell}^E c^2 c_E B_{\ell}^T B_{\ell}^E C_{\ell}^{EE}$$

# Correlation coefficient of T and E modes

$$\mathcal{R}_\ell^{TE} = \frac{\langle a_{\ell m}^T a_{\ell m}^{E*} \rangle}{\sqrt{\langle a_{\ell m}^T a_{\ell m}^{T*} \rangle \langle a_{\ell m}^E a_{\ell m}^{E*} \rangle}} = \frac{C_\ell^{TE}}{\sqrt{C_\ell^{TT} C_\ell^{EE}}}$$

# Correlation coefficient of T and E modes

$$\mathcal{R}_\ell^{TE} = \frac{\langle a_{\ell m}^T a_{\ell m}^{E*} \rangle}{\sqrt{\langle a_{\ell m}^T a_{\ell m}^{T*} \rangle \langle a_{\ell m}^E a_{\ell m}^{E*} \rangle}} = \frac{C_\ell^{TE}}{\sqrt{C_\ell^{TT} C_\ell^{EE}}}$$

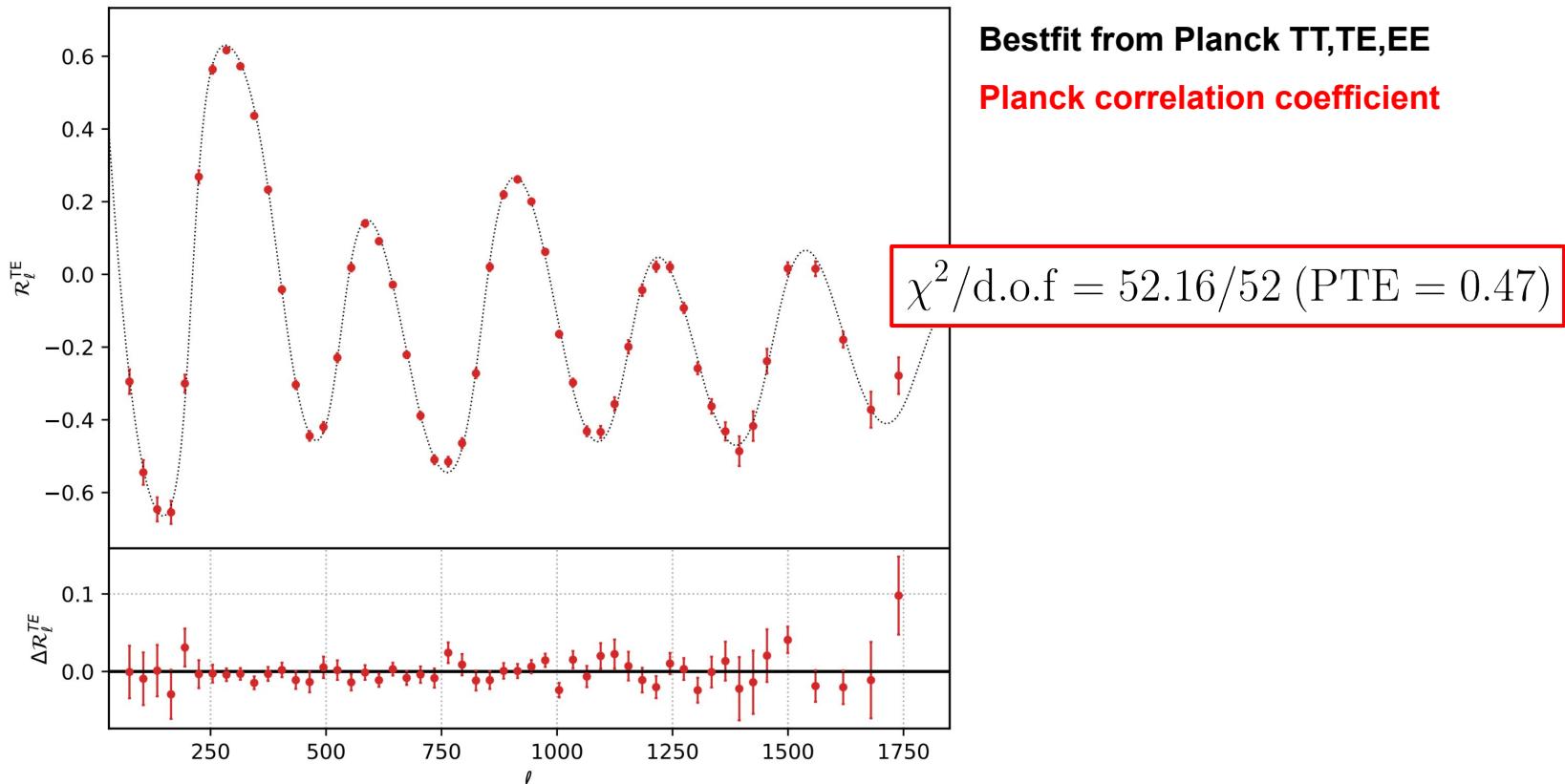
$$\mathcal{R}_\ell^{TE,\text{obs}} = \frac{\mathcal{F}_\ell^T \mathcal{F}_\ell^E c^2 c_E B_\ell^T B_\ell^E C_\ell^{TE}}{\sqrt{(\mathcal{F}_\ell^T)^2 c^2 (B_\ell^T)^2 C_\ell^{TT} \times (\mathcal{F}_\ell^E)^2 c^2 c_E^2 (B_\ell^E)^2 C_\ell^{EE}}}$$

# Correlation coefficient of T and E modes

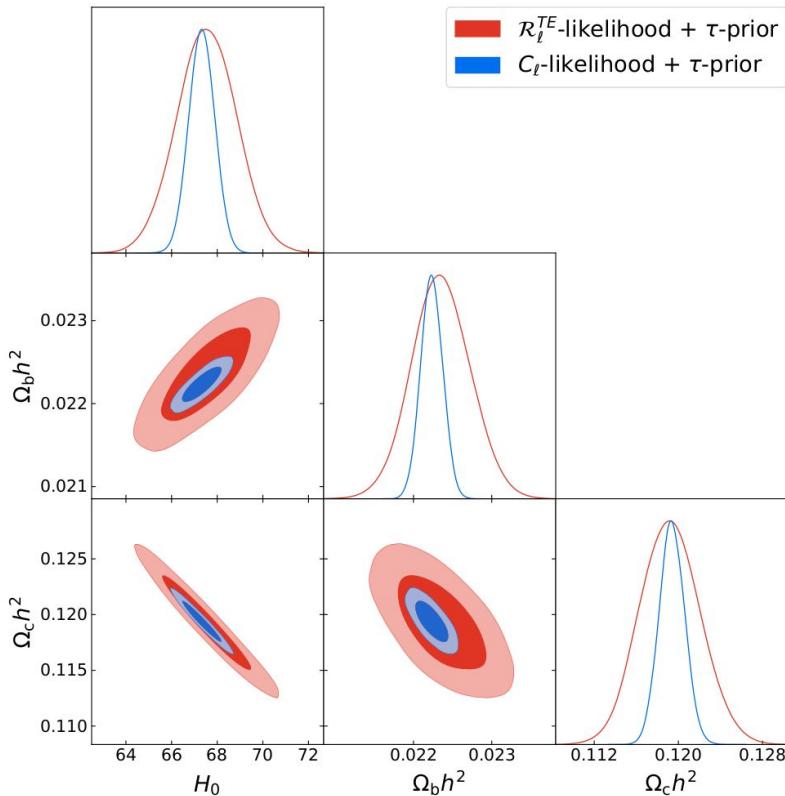
$$\mathcal{R}_\ell^{TE} = \frac{\langle a_{\ell m}^T a_{\ell m}^{E*} \rangle}{\sqrt{\langle a_{\ell m}^T a_{\ell m}^{T*} \rangle \langle a_{\ell m}^E a_{\ell m}^{E*} \rangle}} = \frac{C_\ell^{TE}}{\sqrt{C_\ell^{TT} C_\ell^{EE}}}$$

$$\mathcal{R}_\ell^{TE,\text{obs}} = \frac{\cancel{\mathcal{F}_\ell^T \mathcal{F}_\ell^E c^2 c_E B_\ell^T B_\ell^E} C_\ell^{TE}}{\sqrt{\cancel{(\mathcal{F}_\ell^T)^2 c^2 (B_\ell^T)^2} C_\ell^{TT} \times \cancel{(\mathcal{F}_\ell^E)^2 c^2 c_E^2 (B_\ell^E)^2} C_\ell^{EE}}} = \mathcal{R}_\ell^{TE}$$

# Planck correlation coefficient



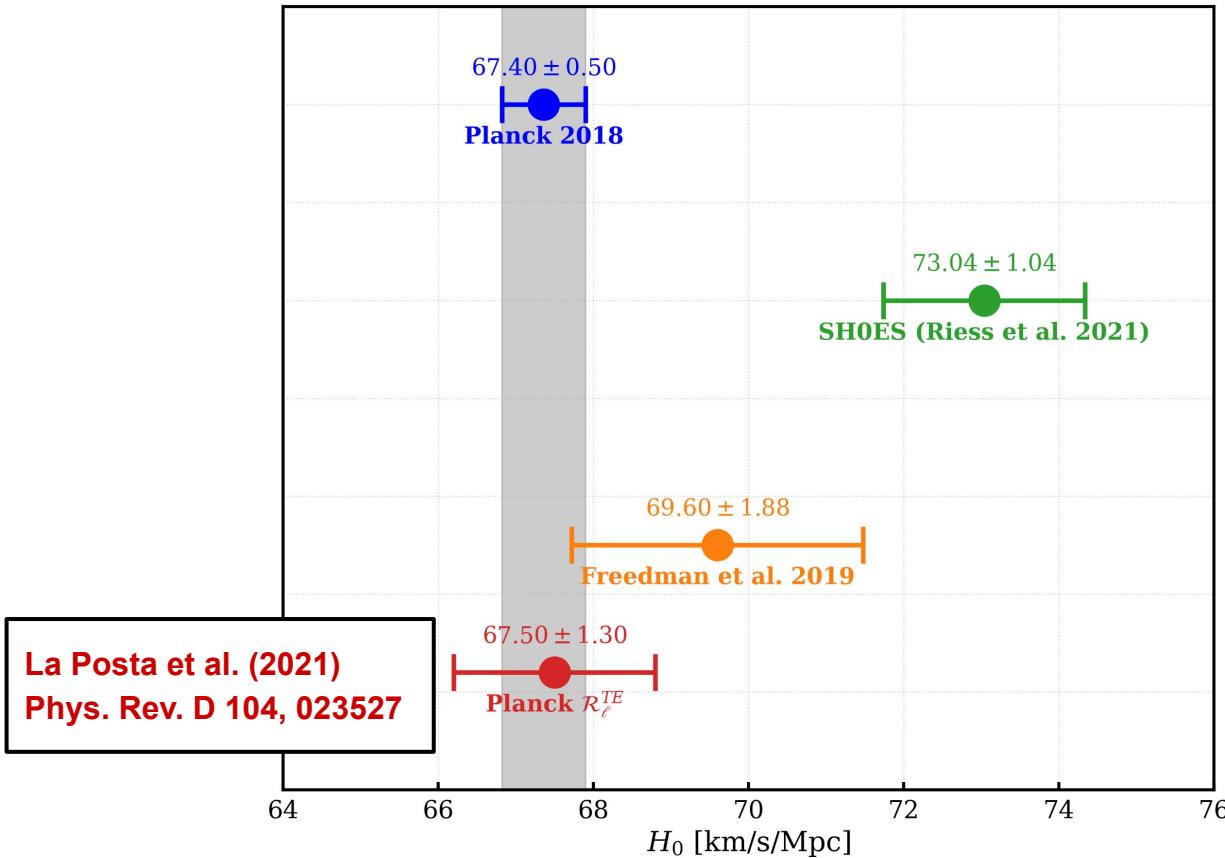
# Cosmological results from R<sup>TE</sup>



3.3 $\sigma$  away from the latest  
SH0ES measurement

$$H_0 = 67.5 \pm 1.3 \text{ [km/s/Mpc]}$$

# Hubble tension



# Independent measurements of $H_0$ from the ground

## Atacama Cosmology Telescope

6m telescope in the Atacama desert  
(Chile ~5000m high)

ACT DR4 ([Choi+ 2020](#), [Aiola+ 2020](#))

data collected from 2013 to 2016

Cosmological analysis on ~5400 deg<sup>2</sup>

observed at 98 and 150 GHz

## South Pole Telescope

10m primary mirror  
(South Pole ~2800m high)

SPT-3G results ([Dutcher+ 2021](#))

4 month period in 2018

Cosmological analysis on ~1500 deg<sup>2</sup>

observed at 95, 150 and 220 GHz