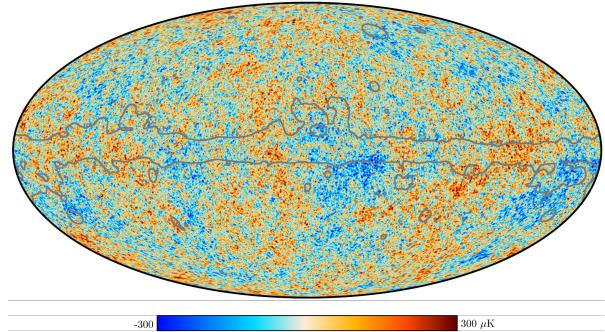


The Hubble tension : a CMB perspective

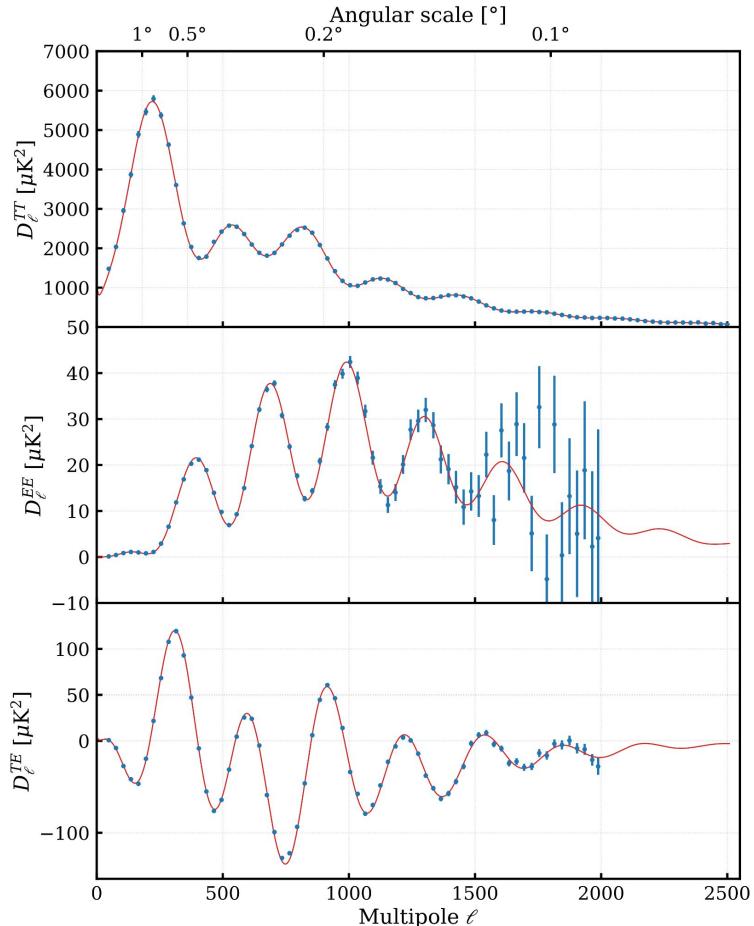
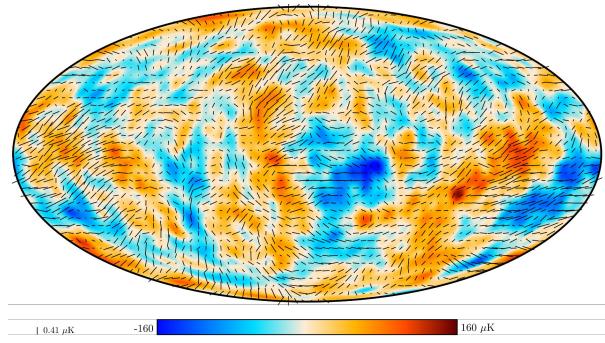
Adrien La Posta
IJClab

Measuring H_0 from the CMB

Temperature

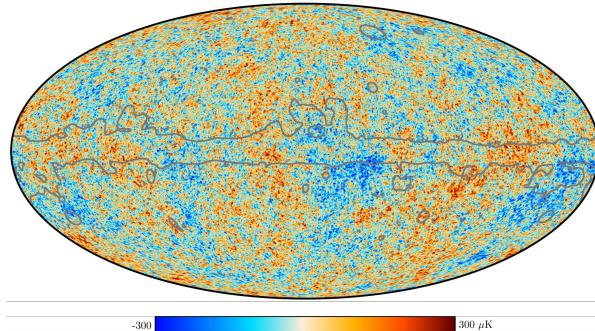


Polarization E-modes

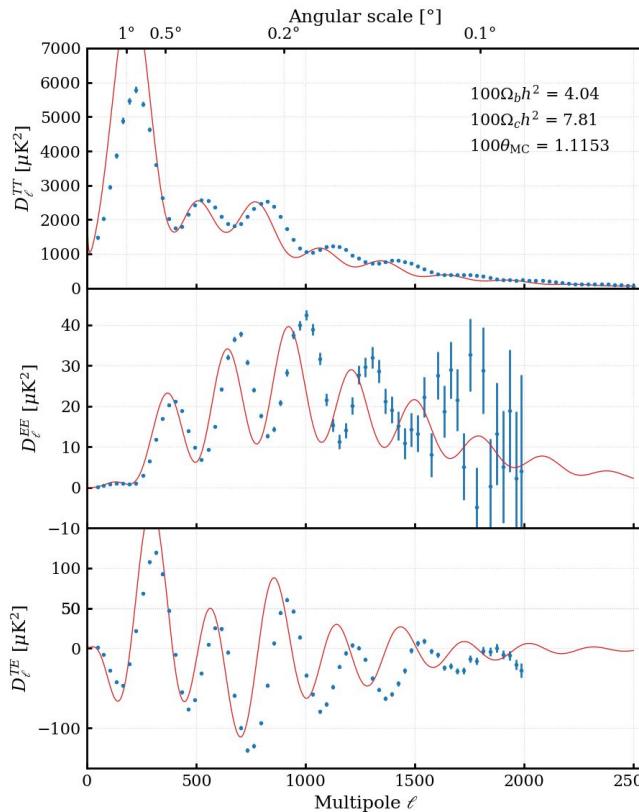
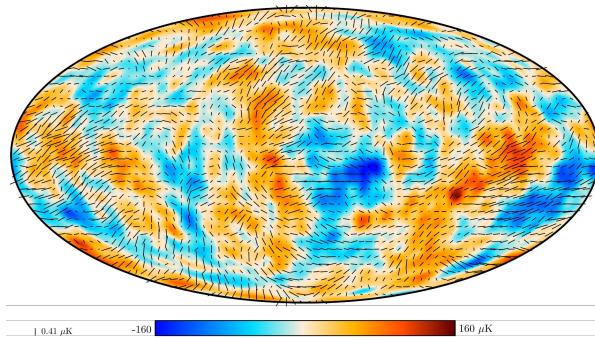


Measuring H_0 from the CMB

Temperature



Polarization E-modes



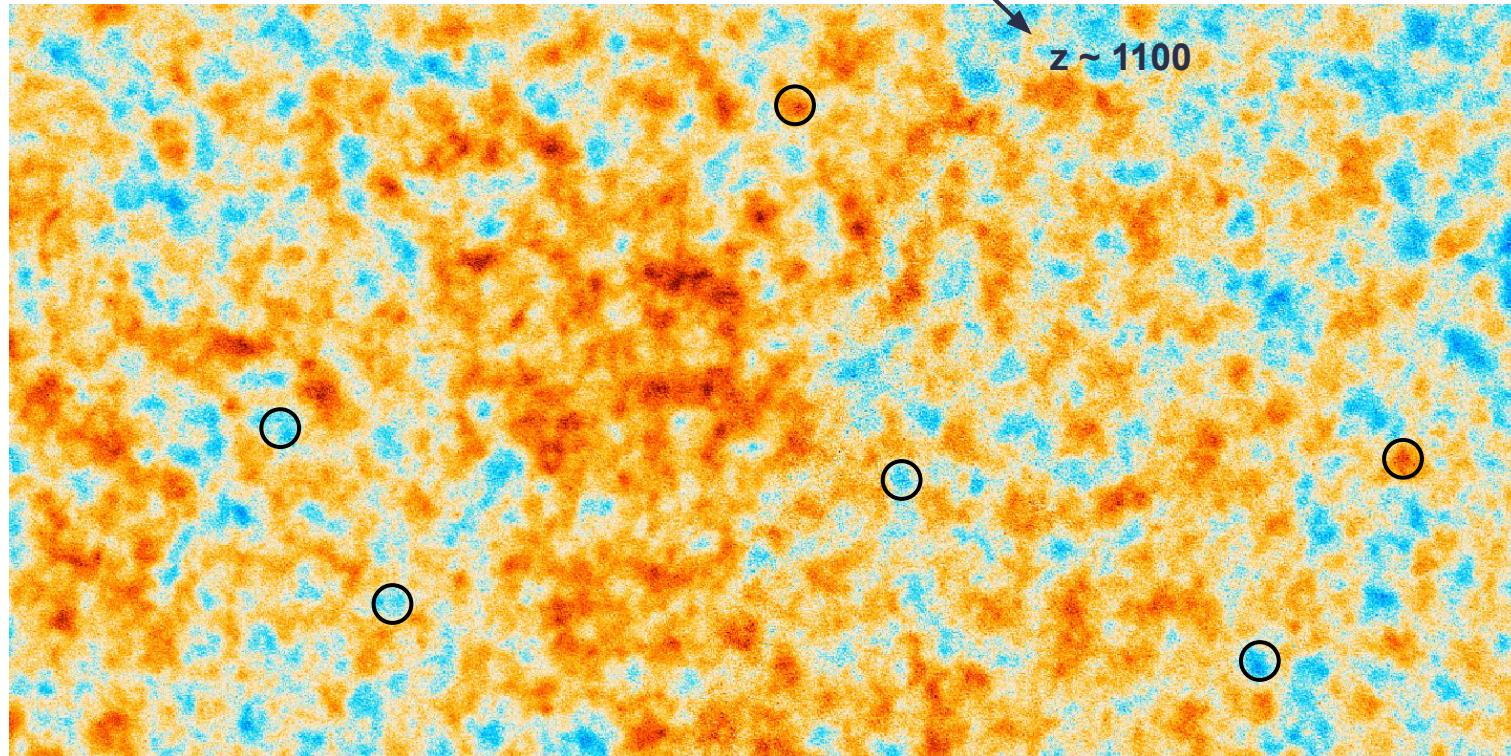
→ θ_* ρ_b^0 ρ_c^0

Measuring H_0 from the CMB

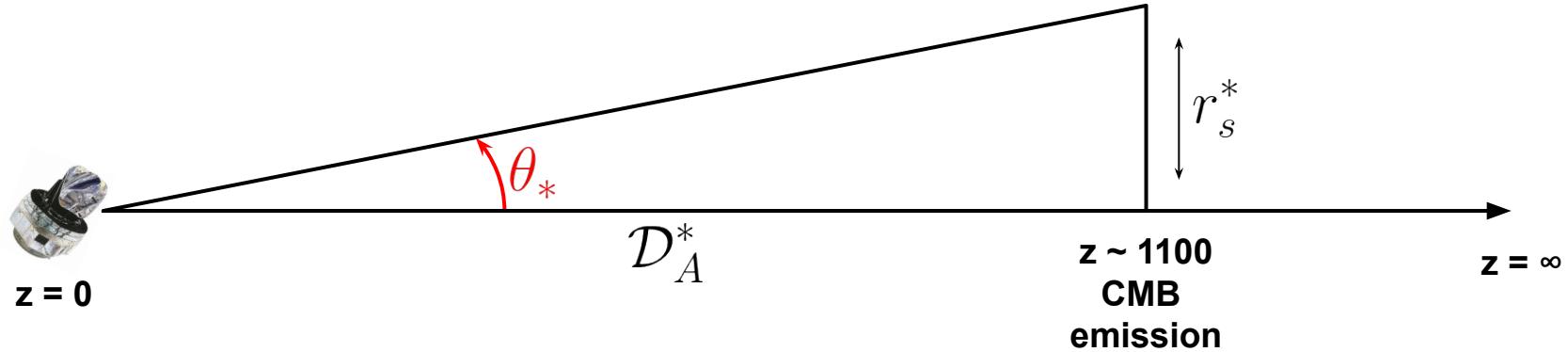
CMB standard ruler : **size of the sound horizon at decoupling** imprinted in the CMB radiation

Measuring H_0 from the CMB

CMB standard ruler : size of the sound horizon at decoupling imprinted in the CMB radiation

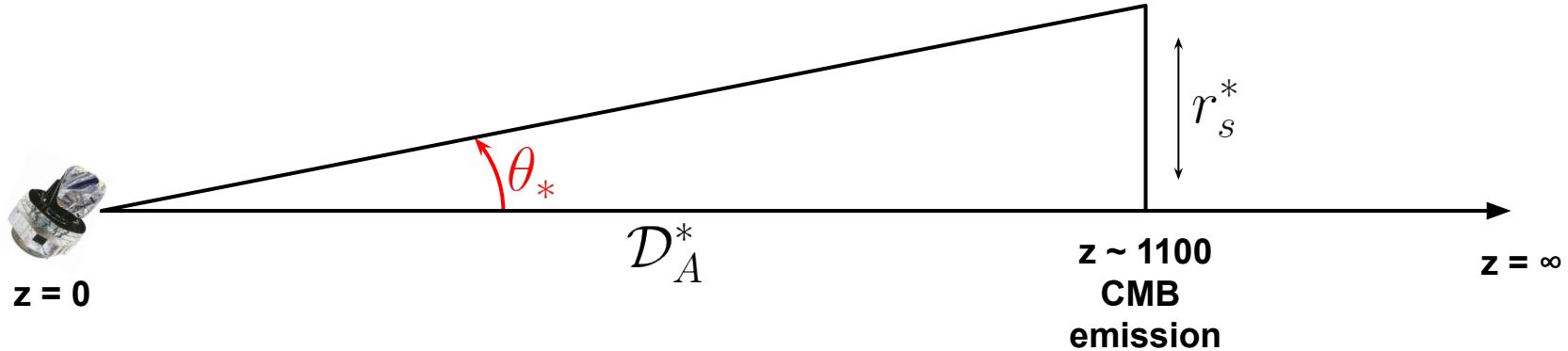


Measuring H_0 from the CMB



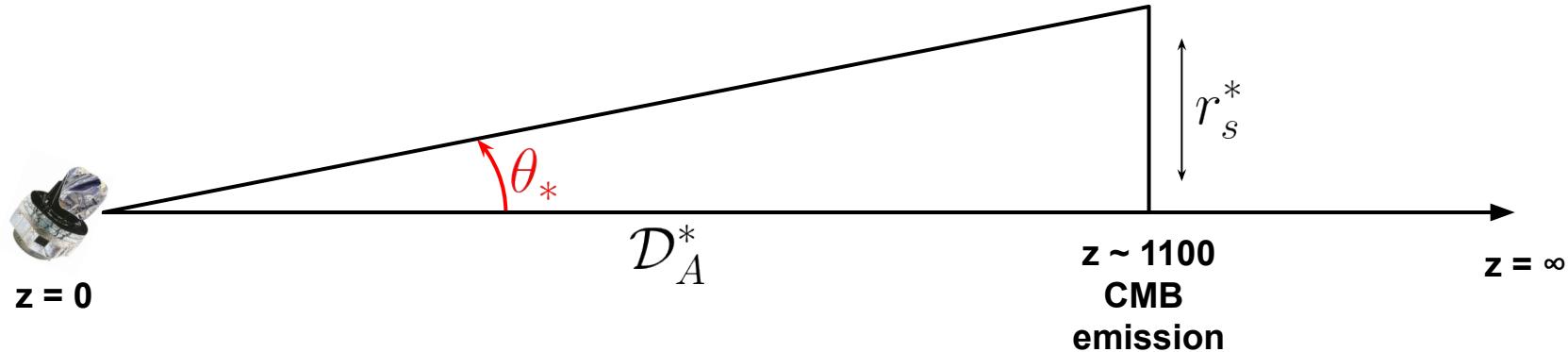
$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*}$$

Measuring H_0 from the CMB



$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*} \quad \xrightarrow{r_s^* = \int_{z^*}^{\infty} \frac{dz}{H(z)} c_s(z)}$$

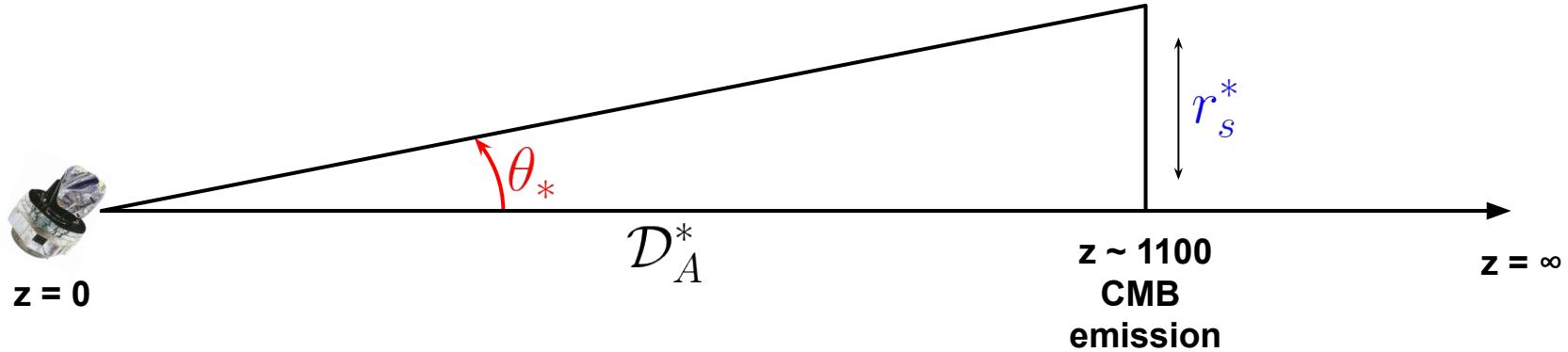
Measuring H_0 from the CMB



$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*} \longrightarrow r_s^* = \int_{z^*}^{\infty} \frac{dz}{H(z)} c_s(z) \longrightarrow c_s(z) = c \sqrt{\frac{1}{3 [1 + 3\rho_b^0/4\rho_\gamma^0(1+z)^{-1}]}}$$

$$H_{\text{early}}^2(z) = \frac{8\pi G}{3} [\rho_r^0(1+z)^4 + (\rho_b^0 + \rho_c^0)(1+z)^3]$$

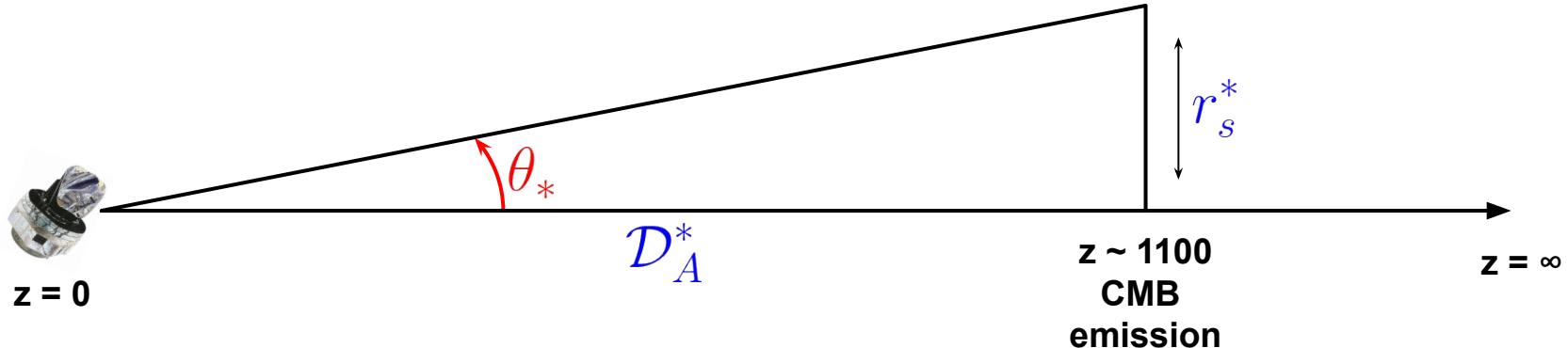
Measuring H_0 from the CMB



Now \mathcal{D}_A^* is known

$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*}$$

Measuring H_0 from the CMB

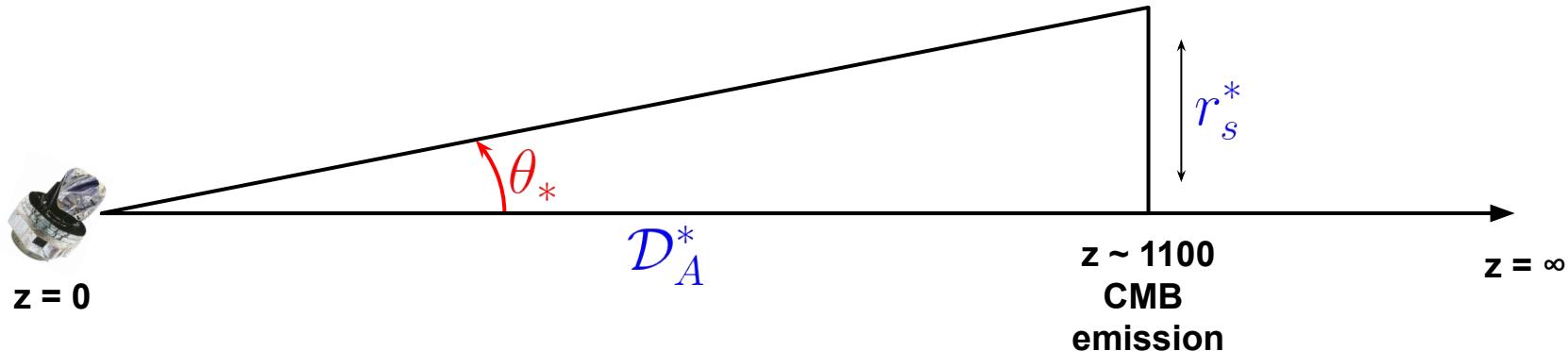


Now \mathcal{D}_A^* is known

$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*} \quad \xrightarrow{\mathcal{D}_A^* = c \int_0^{z^*} \frac{dz}{H(z)}}$$

$$H_{\text{late}}^2(z) = \frac{8\pi G}{3} [(\rho_b^0 + \rho_c^0)(1+z)^3 + \rho_\Lambda]$$

Measuring H_0 from the CMB



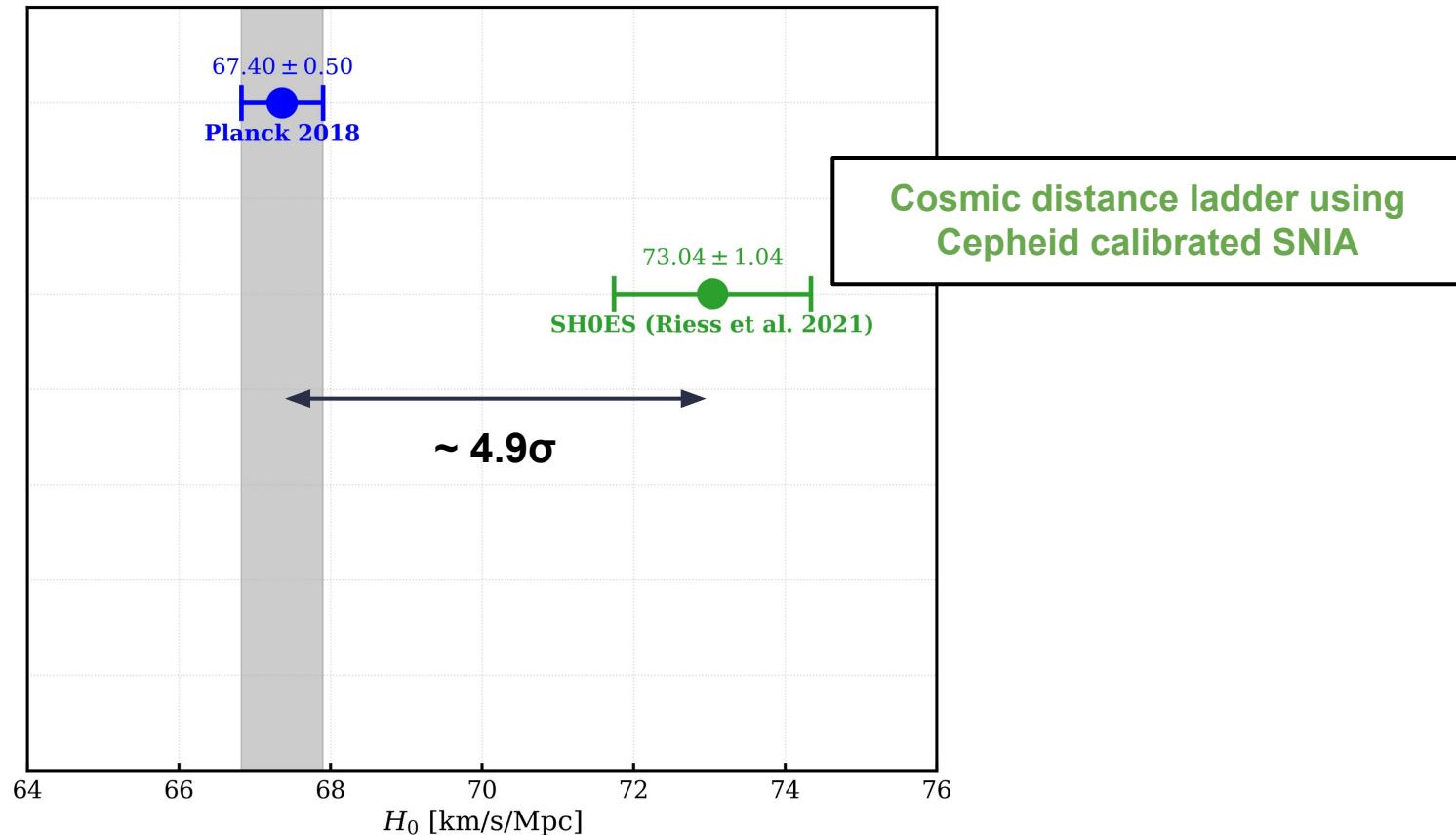
Now \mathcal{D}_A^* is known

$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*} \quad \mathcal{D}_A^* = c \int_0^{z^*} \frac{dz}{H(z)}$$

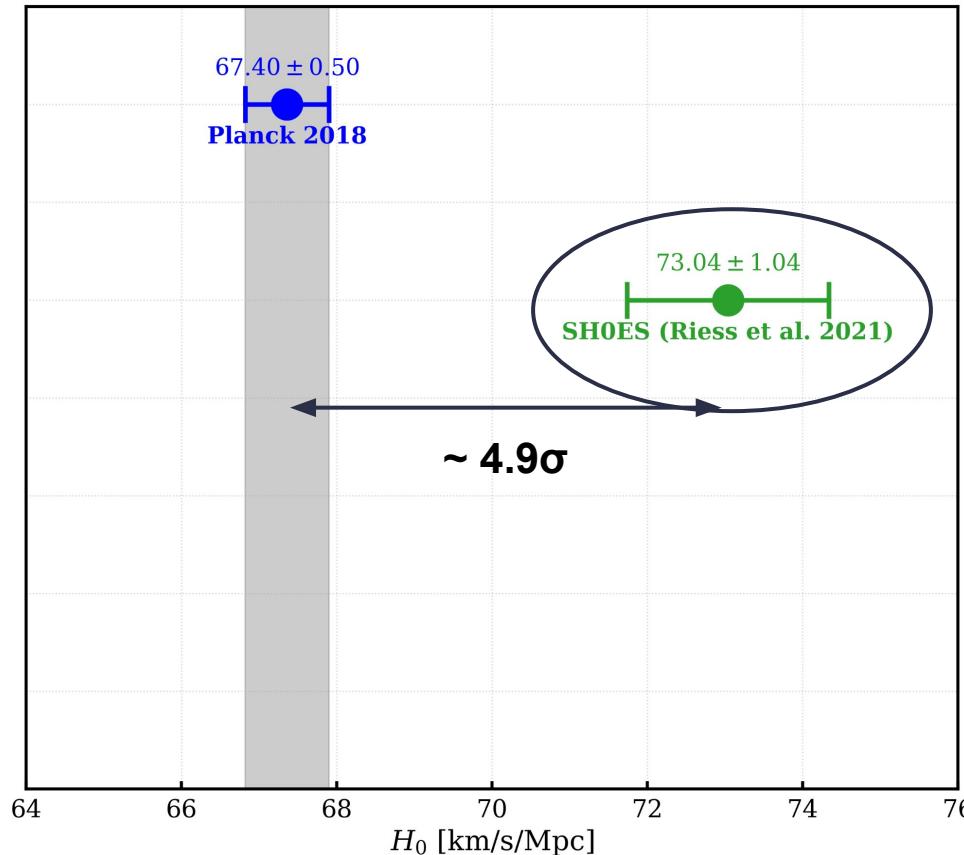
$$H_0^2 = \frac{8\pi G}{3} [\rho_b^0 + \rho_c^0 + \rho_\Lambda]$$

$$H_{\text{late}}^2(z) = \frac{8\pi G}{3} [(\rho_b^0 + \rho_c^0)(1+z)^3 + \rho_\Lambda]$$

The Hubble tension as of today



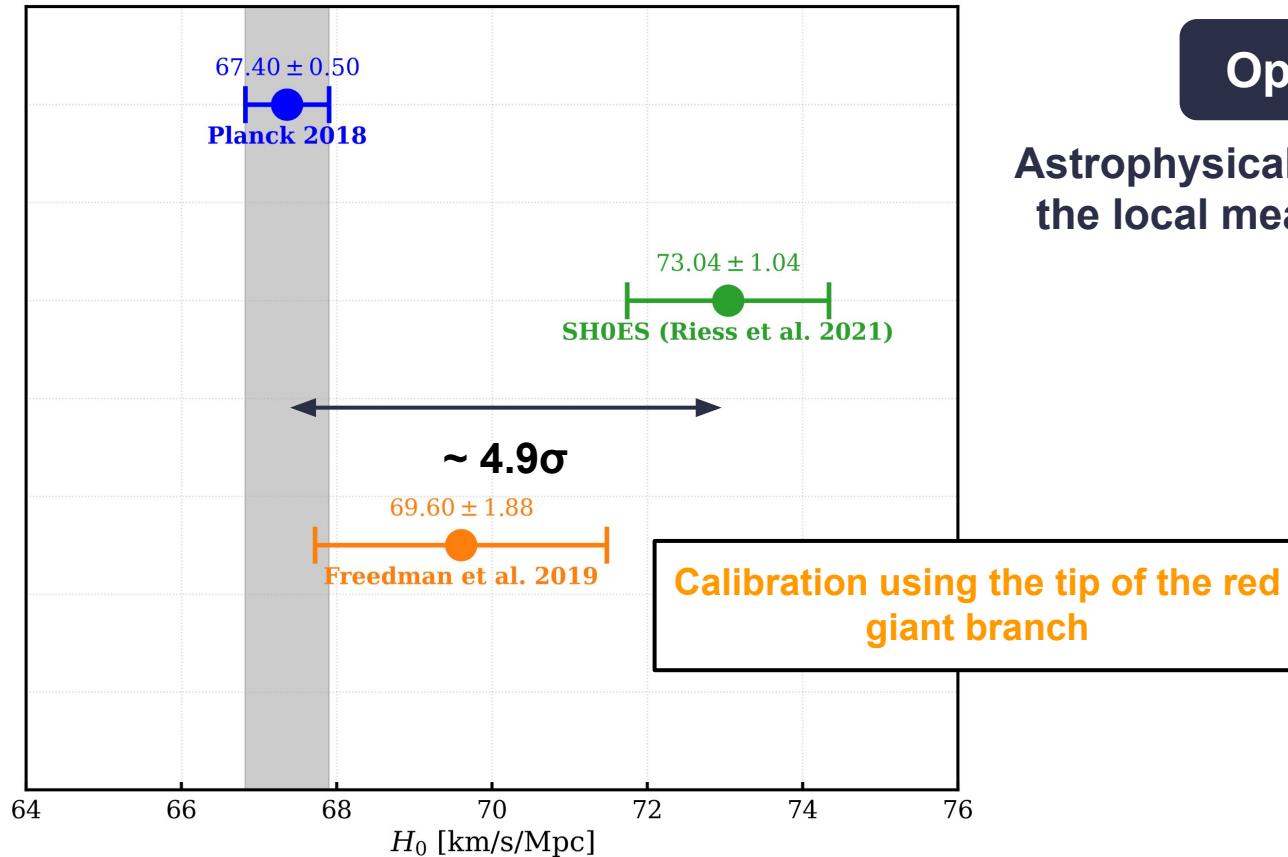
The Hubble tension as of today



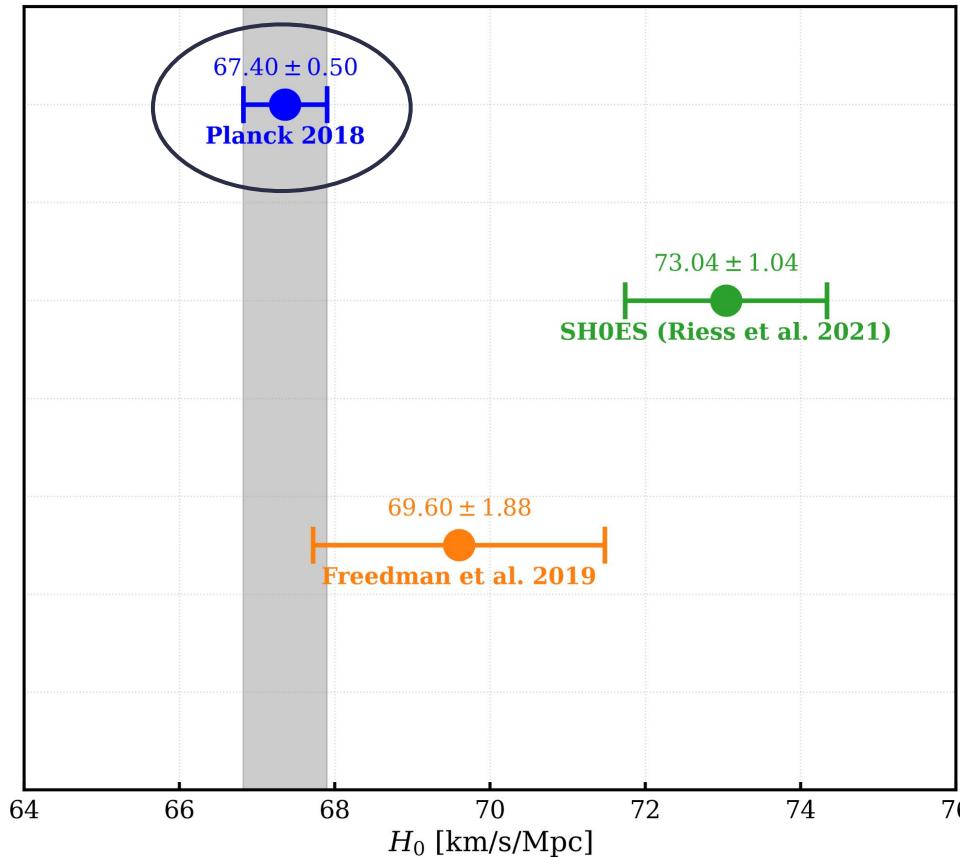
Option 1

Astrophysical biases affecting the local measurement of H_0

The Hubble tension as of today



The Hubble tension as of today



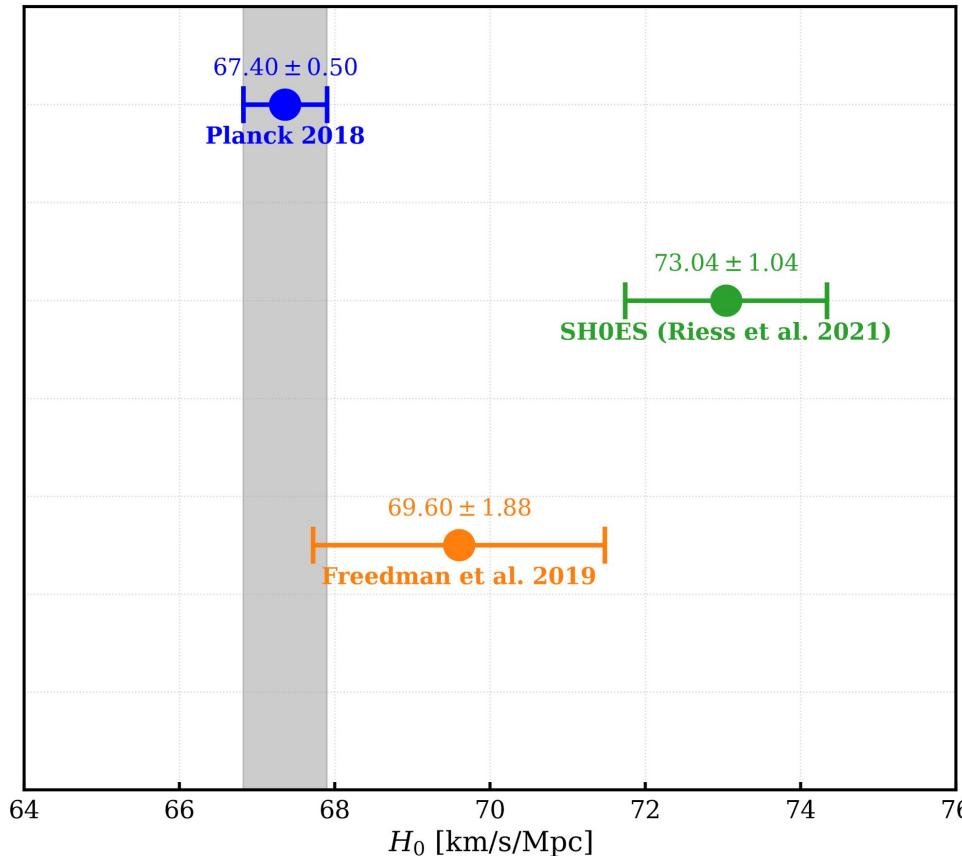
Option 1

Astrophysical biases affecting
the local measurement of H_0

Option 2

Instrumental systematic effect
biasing the value of H_0
inferred from the CMB

The Hubble tension as of today



Option 1

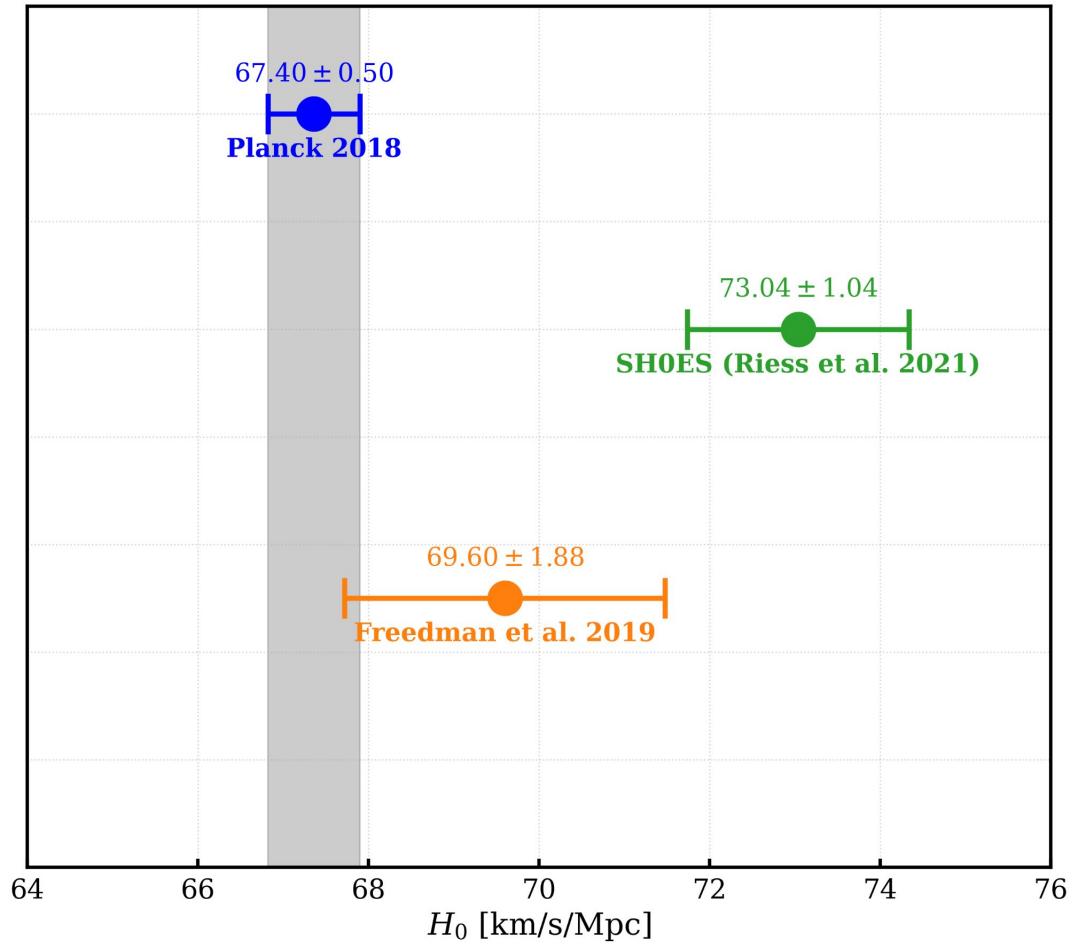
Astrophysical biases affecting
the local measurement of H_0

Option 2

Instrumental systematic effect
biasing the value of H_0
inferred from the CMB

Option 3

Physics beyond Λ CDM



Option 1

Astrophysical biases affecting
the local measurement of H_0

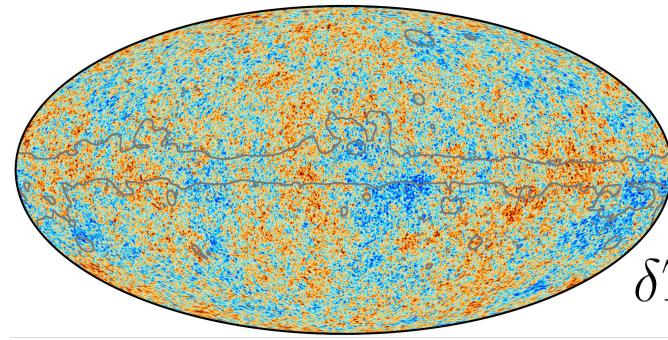
Option 2

Instrumental systematic effect
biasing the value of H_0
inferred from the CMB

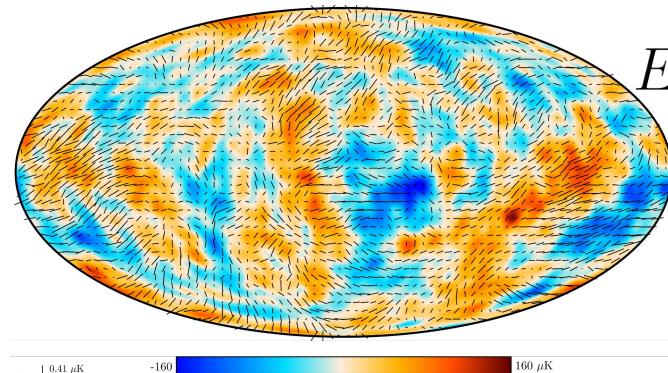
Option 3

Physics beyond Λ CDM

Systematics in the CMB



$$\delta T^{\text{sky}}(\hat{n})$$



$$E^{\text{sky}}(\hat{n})$$

$$\mathcal{I}_{T/E}$$



$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$
$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

Systematics in the CMB

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

$$\mathcal{I}_T = \mathcal{F}_T * c * B_T$$

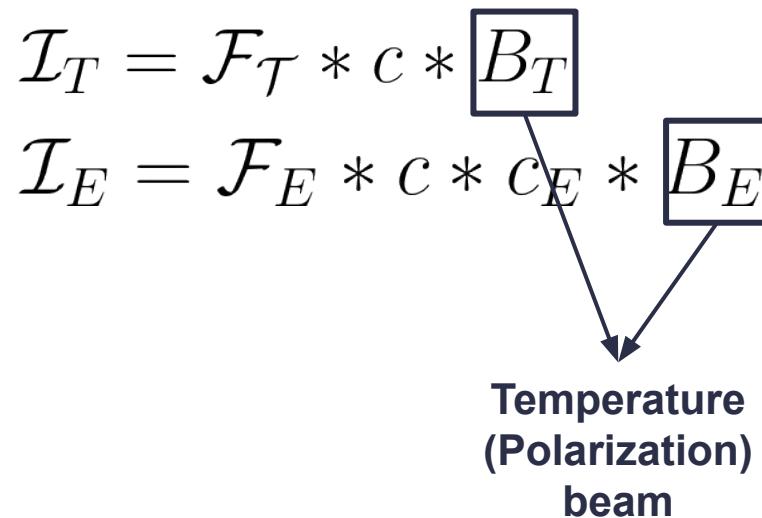
$$\mathcal{I}_E = \mathcal{F}_E * c * c_E * B_E$$

Systematics in the CMB

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)



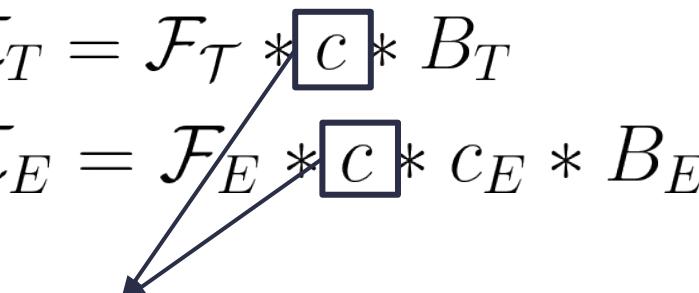
Systematics in the CMB

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)
- Calibration

$$\mathcal{I}_T = \mathcal{F}_T * \boxed{c} * B_T$$
$$\mathcal{I}_E = \mathcal{F}_E * \boxed{c} * c_E * B_E$$



Global
calibration

Systematics in the CMB

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$
$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)
- Calibration
- Polarization efficiency

$$\mathcal{I}_T = \mathcal{F}_T * c * B_T$$

$$\mathcal{I}_E = \mathcal{F}_E * c * \boxed{c_E} * B_E$$



Polarization
efficiency

Systematics in the CMB

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)
- Calibration
- Polarization efficiency
- Transfer functions (map-making)

$$\mathcal{I}_T = \boxed{\mathcal{F}_T} * c * B_T$$

$$\mathcal{I}_E = \boxed{\mathcal{F}_E} * c * c_E * B_E$$

Temperature
(polarization)
transfer function

Systematics in the CMB

$$\begin{aligned}\delta T^{\text{obs}}(\hat{n}) &= \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n}) \\ E^{\text{obs}}(\hat{n}) &= \mathcal{I}_E * E^{\text{sky}}(\hat{n})\end{aligned}$$

- Finite angular resolution (beams)
- Calibration
- Polarization efficiency
- Transfer functions (map-making)

These instrumental effects are multiplicative in harmonic space

$$C_\ell^{TT,\text{obs}} = (\mathcal{F}_\ell^T)^2 c^2 (B_\ell^T)^2 C_\ell^{TT}$$

$$C_\ell^{EE,\text{obs}} = (\mathcal{F}_\ell^E)^2 c^2 c_E^2 (B_\ell^E)^2 C_\ell^{EE}$$

$$C_\ell^{TE,\text{obs}} = \mathcal{F}_\ell^T \mathcal{F}_\ell^E c^2 c_E B_\ell^T B_\ell^E C_\ell^{EE}$$

Correlation coefficient of T and E modes

10

$$\mathcal{R}_\ell^{TE} = \frac{\langle a_{\ell m}^T a_{\ell m}^{E*} \rangle}{\sqrt{\langle a_{\ell m}^T a_{\ell m}^{T*} \rangle \langle a_{\ell m}^E a_{\ell m}^{E*} \rangle}} = \frac{C_\ell^{TE}}{\sqrt{C_\ell^{TT} C_\ell^{EE}}}$$

Correlation coefficient of T and E modes

$$\mathcal{R}_\ell^{TE} = \frac{\langle a_{\ell m}^T a_{\ell m}^{E*} \rangle}{\sqrt{\langle a_{\ell m}^T a_{\ell m}^{T*} \rangle \langle a_{\ell m}^E a_{\ell m}^{E*} \rangle}} = \frac{C_\ell^{TE}}{\sqrt{C_\ell^{TT} C_\ell^{EE}}}$$

$$\mathcal{R}_\ell^{TE,\text{obs}} = \frac{\mathcal{F}_\ell^T \mathcal{F}_\ell^E c^2 c_E B_\ell^T B_\ell^E C_\ell^{TE}}{\sqrt{(\mathcal{F}_\ell^T)^2 c^2 (B_\ell^T)^2 C_\ell^{TT} \times (\mathcal{F}_\ell^E)^2 c^2 c_E^2 (B_\ell^E)^2 C_\ell^{EE}}}$$

Correlation coefficient of T and E modes

$$\mathcal{R}_\ell^{TE} = \frac{\langle a_{\ell m}^T a_{\ell m}^{E*} \rangle}{\sqrt{\langle a_{\ell m}^T a_{\ell m}^{T*} \rangle \langle a_{\ell m}^E a_{\ell m}^{E*} \rangle}} = \frac{C_\ell^{TE}}{\sqrt{C_\ell^{TT} C_\ell^{EE}}}$$

$$\mathcal{R}_\ell^{TE,\text{obs}} = \frac{\cancel{\mathcal{F}_\ell^T \mathcal{F}_\ell^E c^2 c_E B_\ell^T B_\ell^E C_\ell^{TE}}}{\sqrt{(\cancel{\mathcal{F}_\ell^T})^2 c^2 (B_\ell^T)^2 C_\ell^{TT} \times (\cancel{\mathcal{F}_\ell^E})^2 c^2 c_E^2 (B_\ell^E)^2 C_\ell^{EE}}} = \mathcal{R}_\ell^{TE}$$

Correlation coefficient of T and E modes

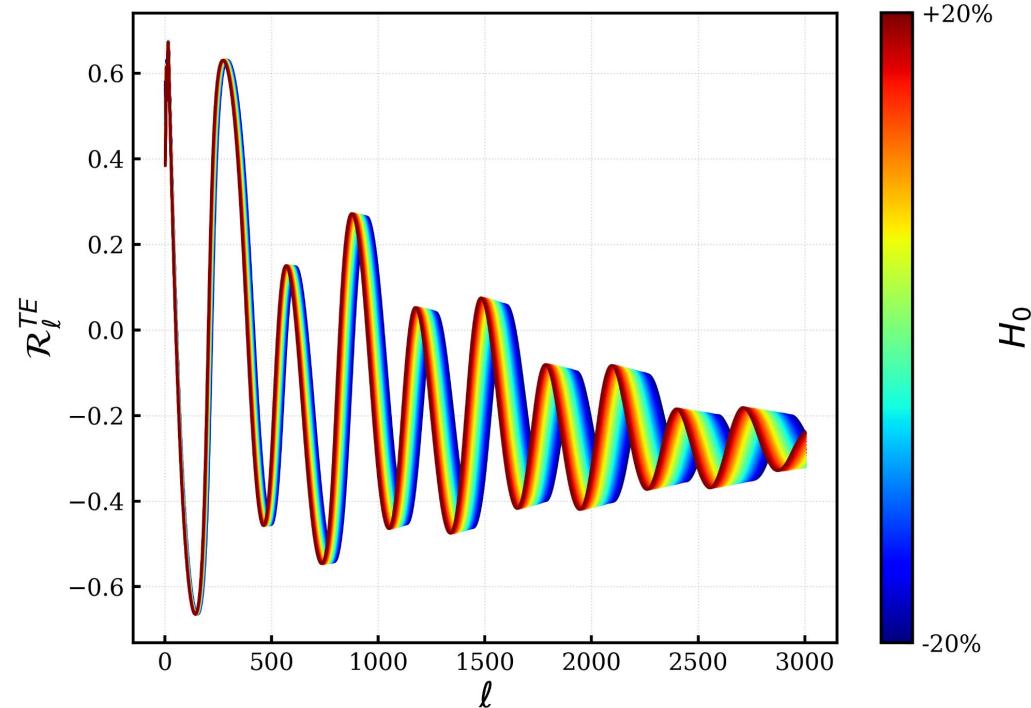
11

- The correlation coefficient is an observable insensitive to multiplicative biases
→ unbiased constraints on cosmological parameters

Correlation coefficient of T and E modes

11

- The correlation coefficient is an observable insensitive to multiplicative biases
 - unbiased constraints on cosmological parameters
- Particularly sensitive to H_0



Building a likelihood for \mathbf{R}^{TE}

Gaussian likelihood based on a Planck multifrequency likelihood (HiLLiPoP)

3 frequencies : **100, 143 and 217 GHz**

$$\ln \mathcal{L} \simeq -\frac{1}{2} (\Delta \mathcal{R}^{\text{vec}})^T \boldsymbol{\Xi}^{-1} (\Delta \mathcal{R}^{\text{vec}})$$

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$$\ln \mathcal{L} \simeq -\frac{1}{2} (\Delta \mathcal{R}^{\text{vec}})^T \boldsymbol{\Xi}^{-1} (\Delta \mathcal{R}^{\text{vec}})$$



$$\Delta \mathcal{R}_\ell^{TE, \nu_1 \times \nu_2} = \hat{\mathcal{R}}_\ell^{TE, \nu_1 \times \nu_2} - \mathcal{R}_\ell^{TE, \nu_1 \times \nu_2, \text{model}}$$

Building a likelihood for \mathbf{R}^{TE}

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**Unbiased
estimator (data)**

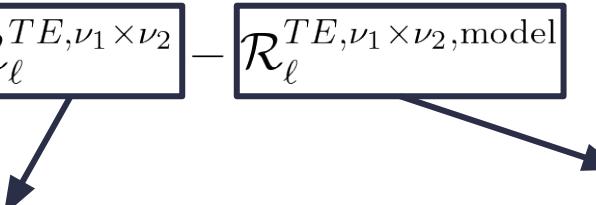
Building a likelihood for \mathbf{R}^{TE}

Gaussian likelihood based on a Planck multifrequency likelihood (HiLLiPoP)

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$$\ln \mathcal{L} \simeq -\frac{1}{2} (\Delta \mathcal{R}^{\text{vec}})^T \boldsymbol{\Xi}^{-1} (\Delta \mathcal{R}^{\text{vec}})$$

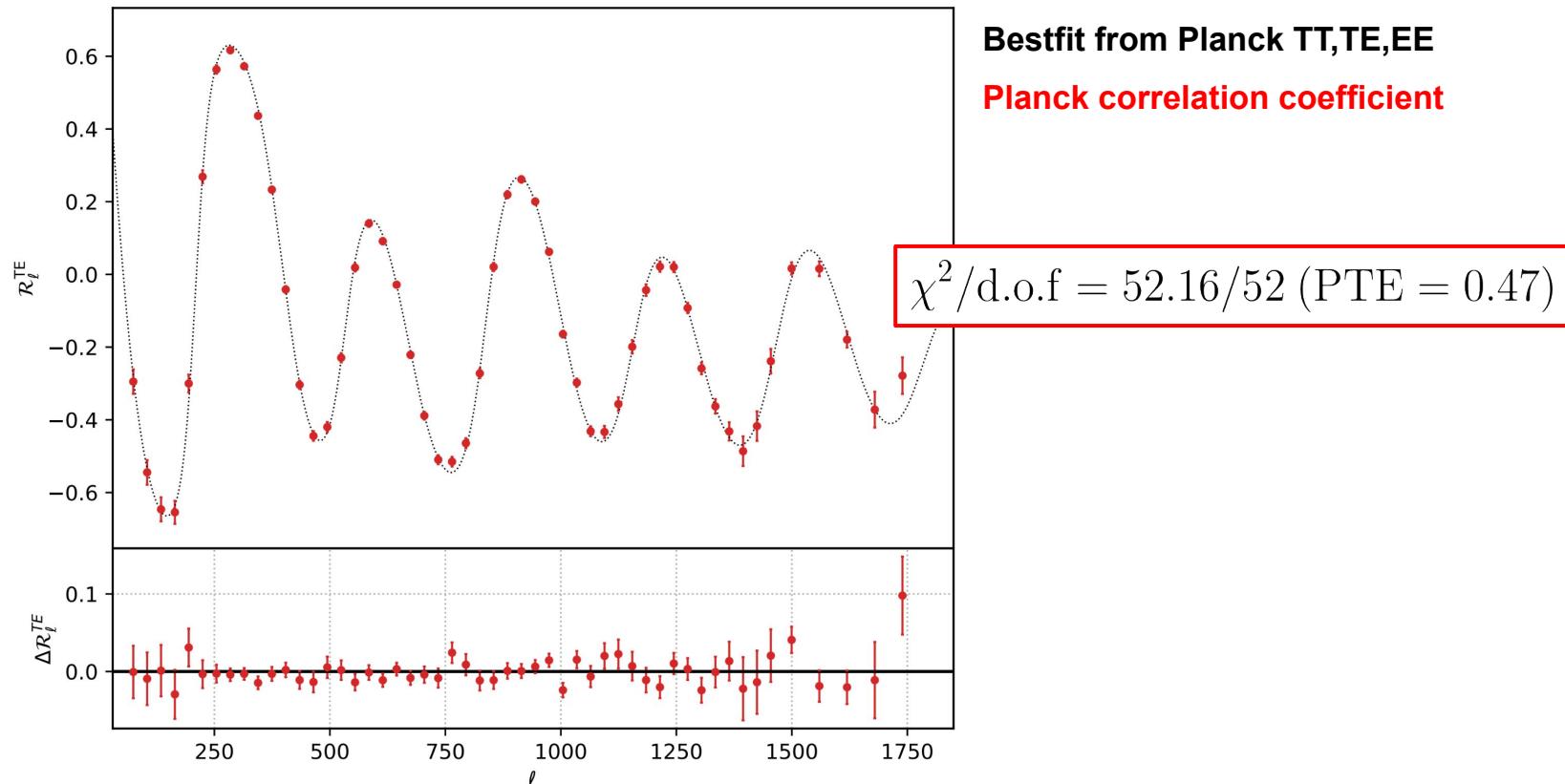
$$\Delta \mathcal{R}_\ell^{TE, \nu_1 \times \nu_2} = \hat{\mathcal{R}}_\ell^{TE, \nu_1 \times \nu_2} - \mathcal{R}_\ell^{TE, \nu_1 \times \nu_2, \text{model}}$$



$$\frac{C_\ell^{TE, \nu_1 \times \nu_2, \text{model}}(p_{\text{cosmo}}, p_{\text{fg}})}{\sqrt{C_\ell^{TT, \nu_1 \times \nu_2, \text{model}}(p_{\text{cosmo}}, p_{\text{fg}}) C_\ell^{EE, \nu_1 \times \nu_2, \text{model}}(p_{\text{cosmo}}, p_{\text{fg}})}}$$

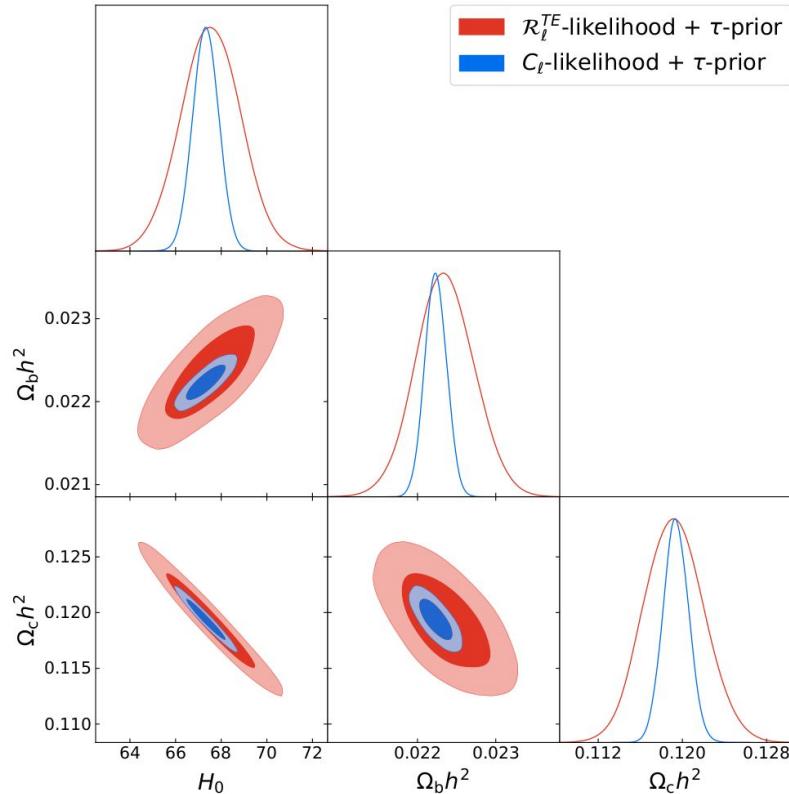
Unbiased estimator (data)

Planck correlation coefficient



Cosmological results from R^{TE}

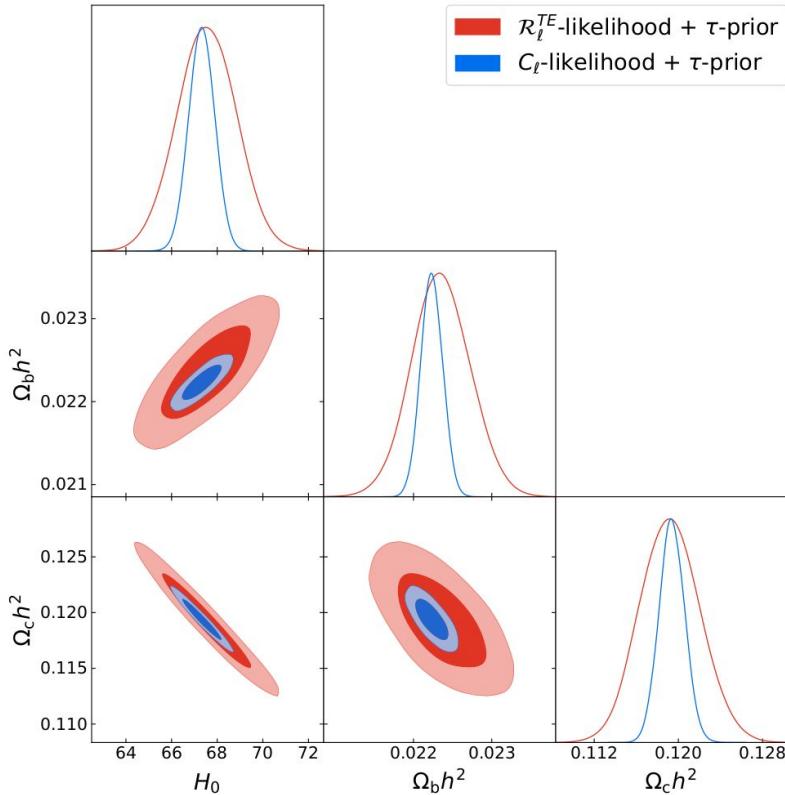
14



$$H_0 = 67.5 \pm 1.3 \text{ [km/s/Mpc]}$$

Cosmological results from R^{TE}

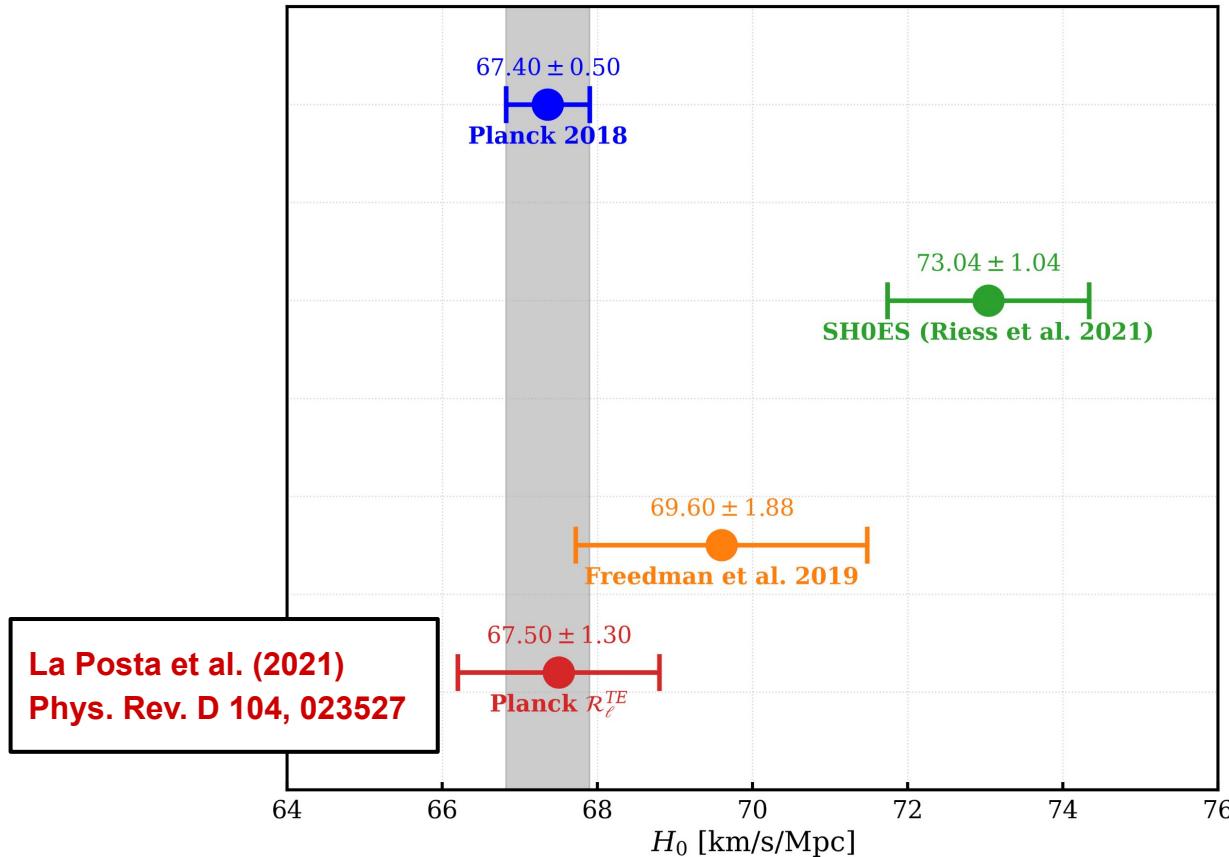
14



3.3 σ away from the latest
SH0ES measurement

$$H_0 = 67.5 \pm 1.3 \text{ [km/s/Mpc]}$$

Hubble tension



Independent measurements of H_0 from the ground

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Independent measurements of H_0 from the ground

17

Atacama Cosmology Telescope

6m telescope in the Atacama desert
(Chile ~5000m high)

ACT DR4 ([Choi+ 2020](#), [Aiola+ 2020](#))

data collected from 2013 to 2016

Cosmological analysis on ~5400 deg²

observed at 98 and 150 GHz

South Pole Telescope

10m primary mirror
(South Pole ~2800m high)

SPT-3G results ([Dutcher+ 2021](#))

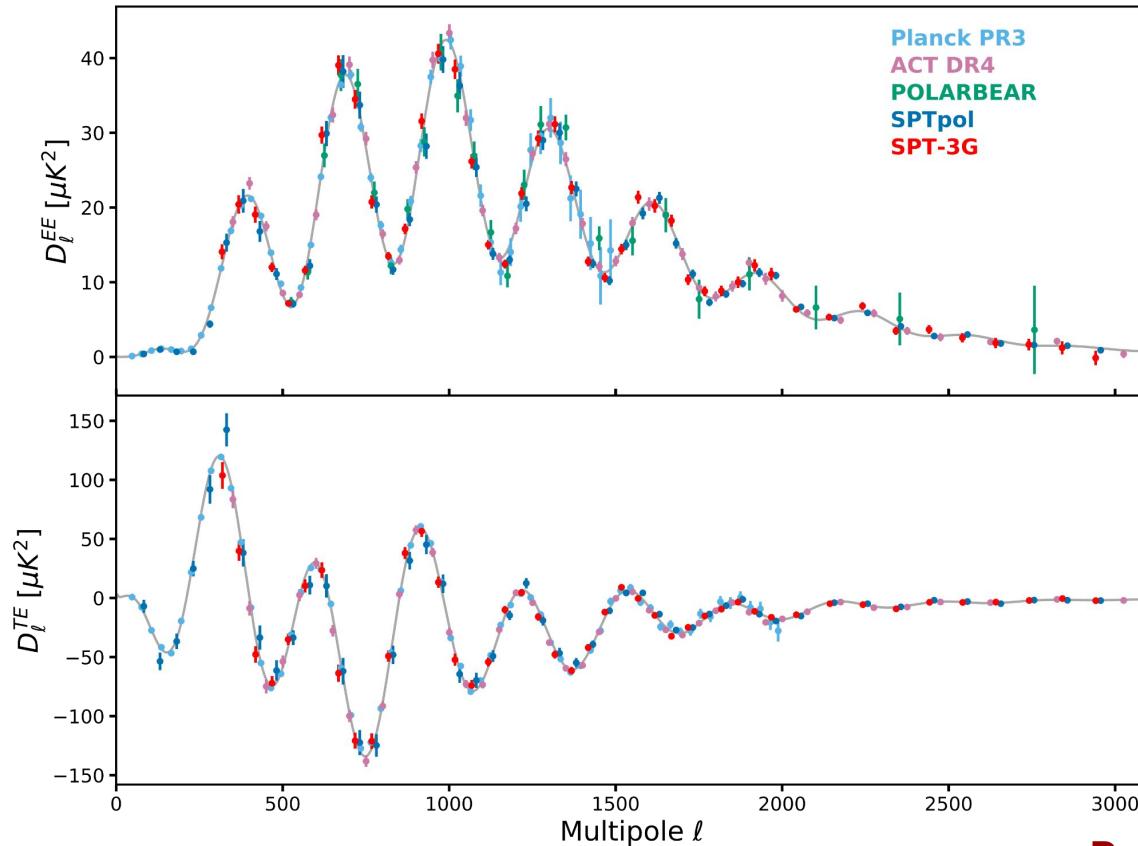
4 month period in 2018

Cosmological analysis on ~1500 deg²

observed at 95, 150 and 220 GHz

CMB Power Spectra

18



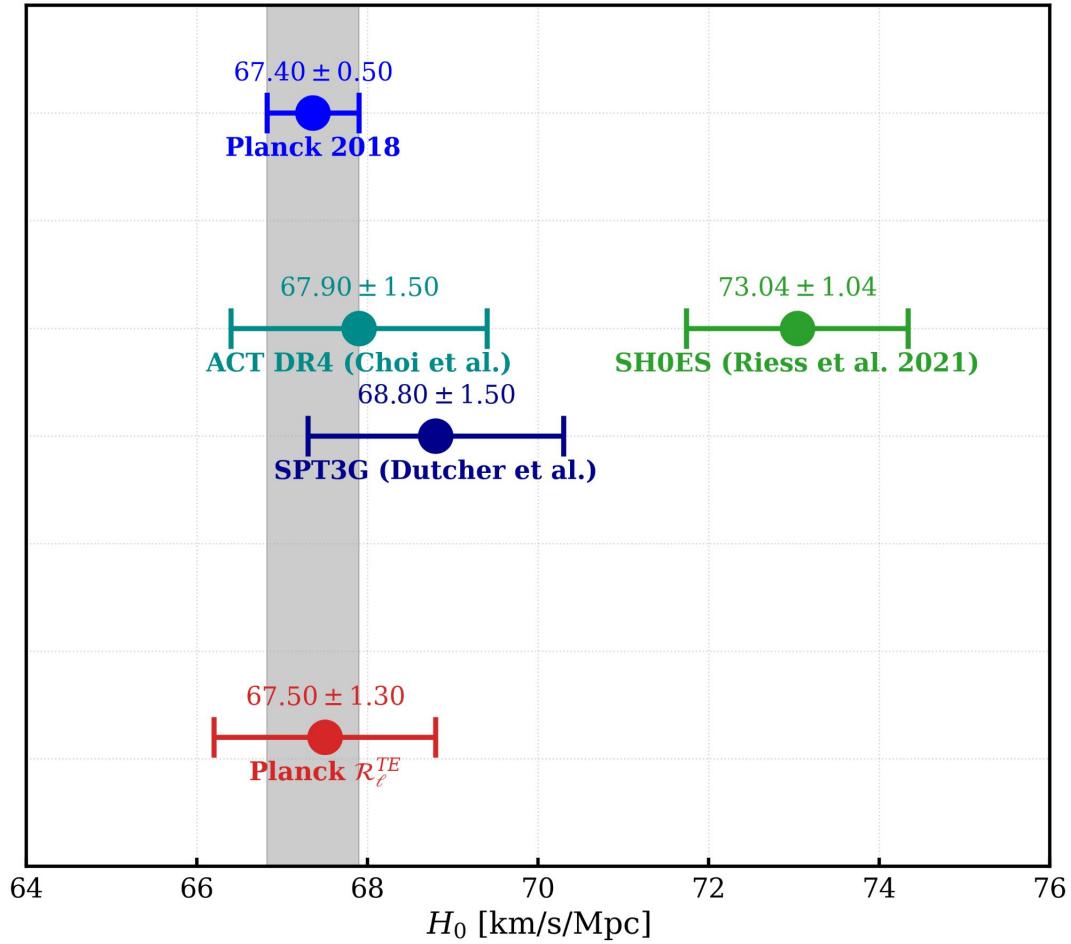
Option 2

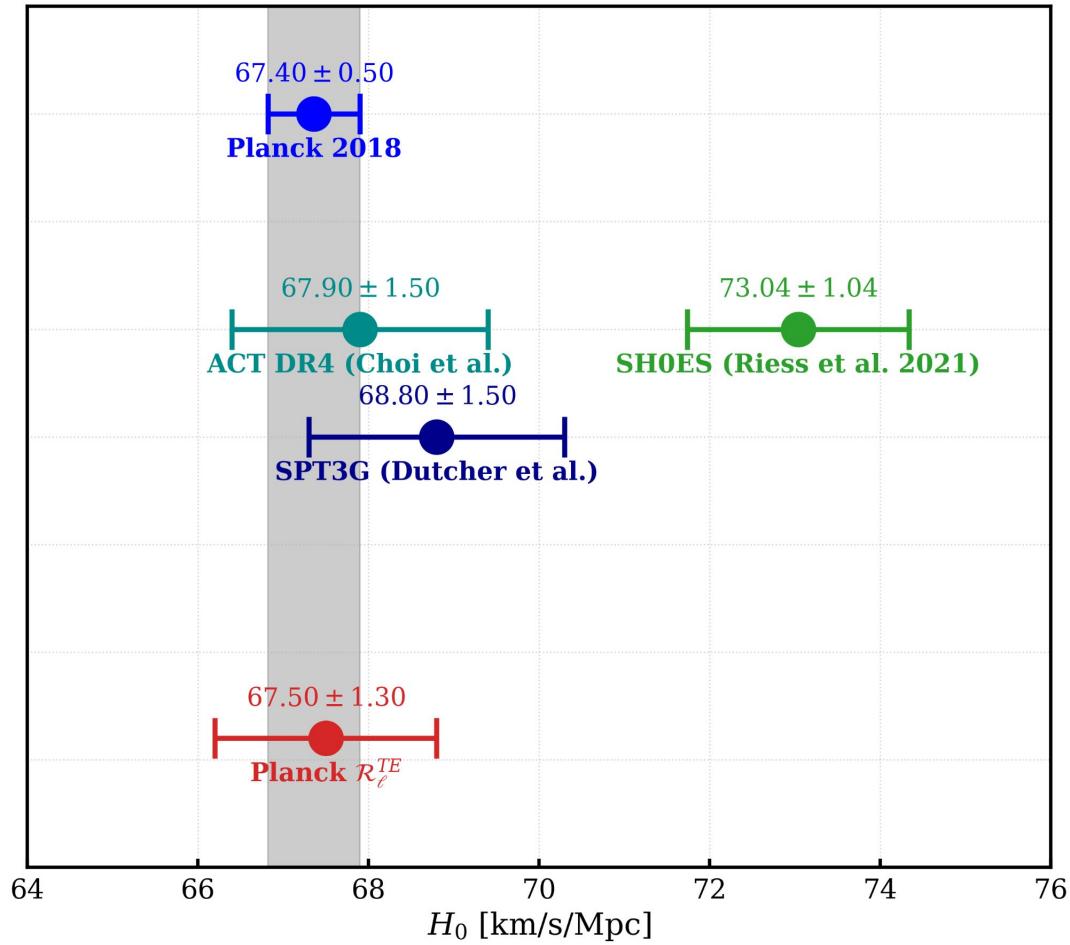
Instrumental systematic effect
biasing the value of H_0
inferred from the CMB



Hard to shift the CMB inferred H_0
with a systematic effect :

- Independent measurements from Planck, ACT and SPT
- Constraint from the correlation coefficient, robust against multiplicative systematics





Option 1

Astrophysical biases affecting
the local measurement of H_0

Option 2

Instrumental systematic effect
biasing the value of H_0
inferred from the CMB

Option 3

Physics beyond Λ CDM

Early-time modification to Λ CDM

Motivation : obtain a higher value of H_0 from the CMB

Early-time modification to Λ CDM

Motivation : obtain a higher value of H_0 from the CMB \longrightarrow lower \mathcal{D}_A^*

Early-time modification to Λ CDM

Motivation : obtain a higher value of H_0 from the CMB \longrightarrow lower \mathcal{D}_A^*

$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*} \longrightarrow \text{Decrease } r_s^* = \int_{z^*}^{\infty} \frac{dz}{H(z)} c_s(z)$$

\downarrow

$$\frac{3H_{\text{early}}^2(z)}{8\pi G} = \rho_r(z) + \rho_m(z)$$

Fixed by
observations

One proposed solution : Early Dark Energy

Motivation : obtain a higher value of H_0 from the CMB \longrightarrow lower \mathcal{D}_A^*

$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*} \longrightarrow \text{Decrease } r_s^* = \int_{z^*}^{\infty} \frac{dz}{H(z)} c_s(z)$$

Fixed by observations

$$\frac{3H_{\text{early}}^2(z)}{8\pi G} = \rho_r(z) + \rho_m(z) + \rho_{\text{EDE}}(z)$$

EDE phenomenology

The EDE component is described as a scalar field ϕ (Poulin+ 2019, Smith+ 2019)

Background evolution : $\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0$

$V(\phi) = m^2 f^2 \left[1 - \cos \left(\frac{\phi}{f} \right) \right]^3$

axion-like potential

EDE phenomenology

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M : mass parameter

EDE phenomenology

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m : mass parameter

f : decay constant

ϕ_i : initial field value

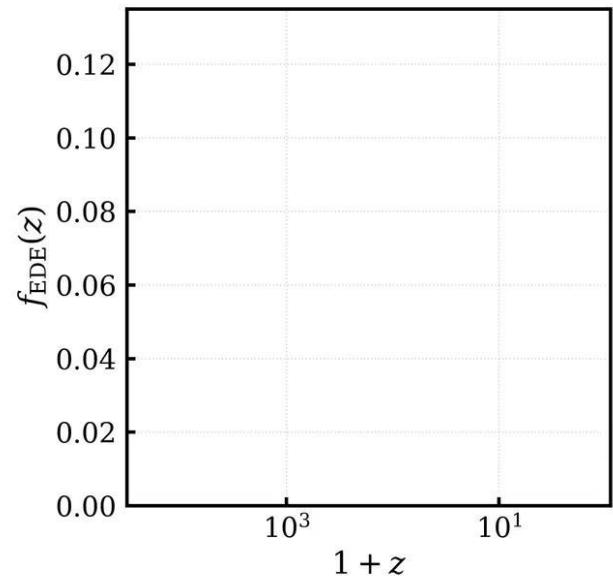
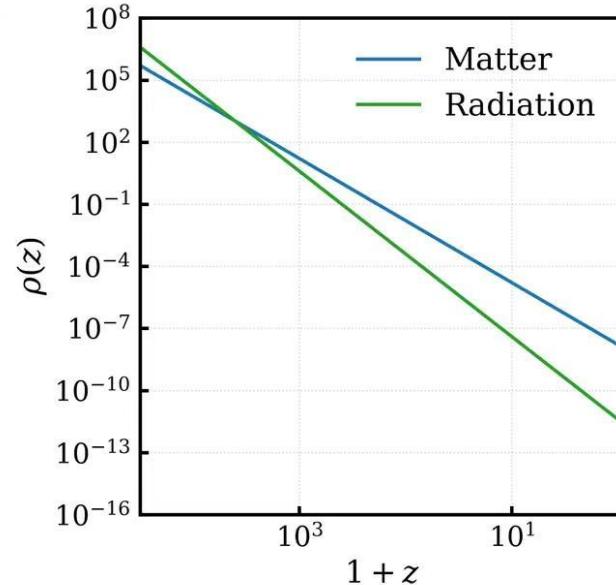
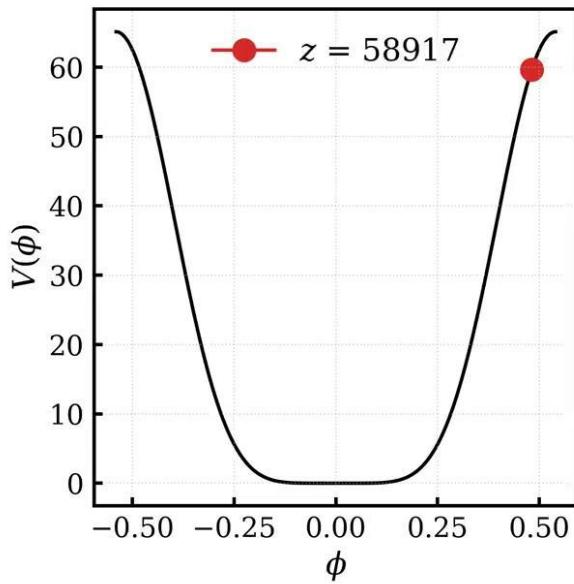
Early Dark Energy : frozen at early times

23

$$\ddot{\phi} + \boxed{3H\dot{\phi}} + V'(\phi) = 0$$

The field is initially frozen due to
Hubble friction ($H \gg m$)

acts as dark energy ($w = -1$)

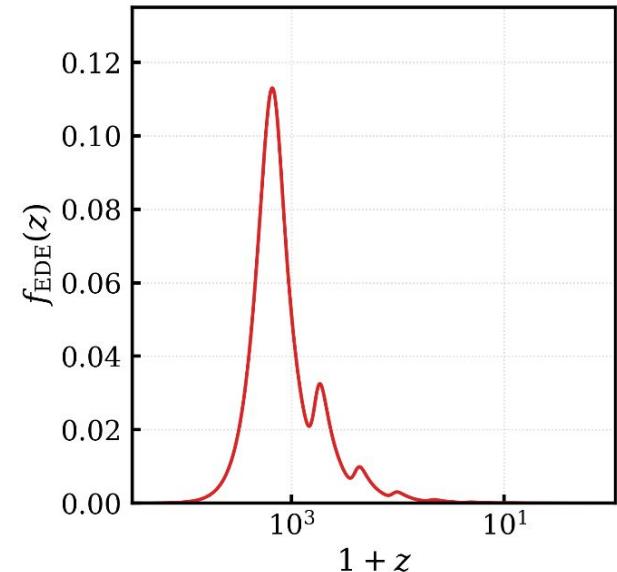
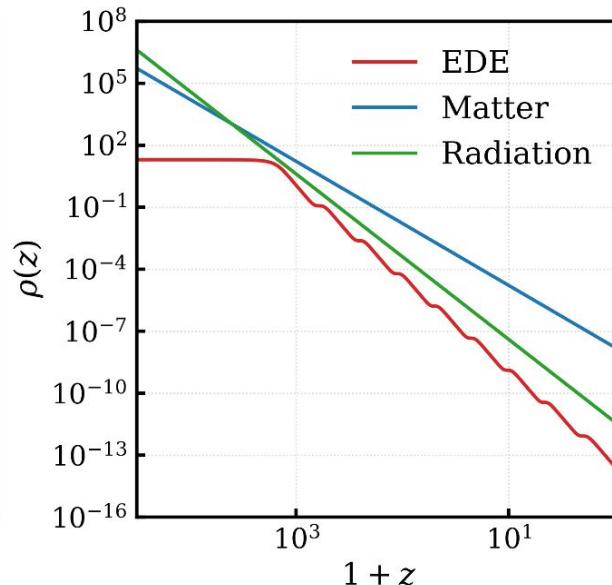
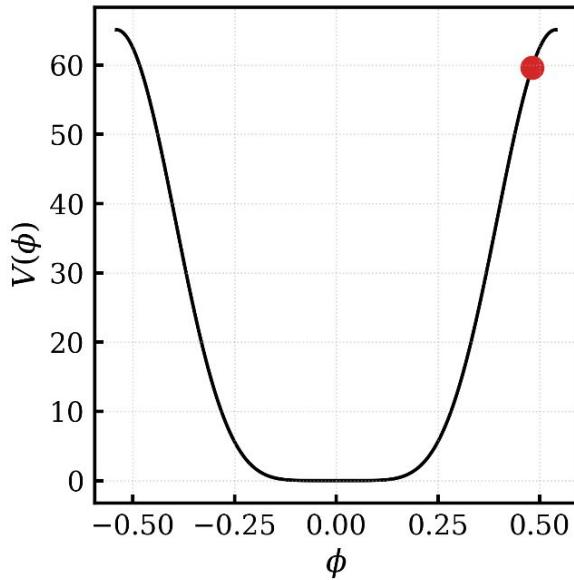


Early Dark Energy : frozen at early times

$$\ddot{\phi} + \boxed{3H\dot{\phi}} + V'(\phi) = 0$$

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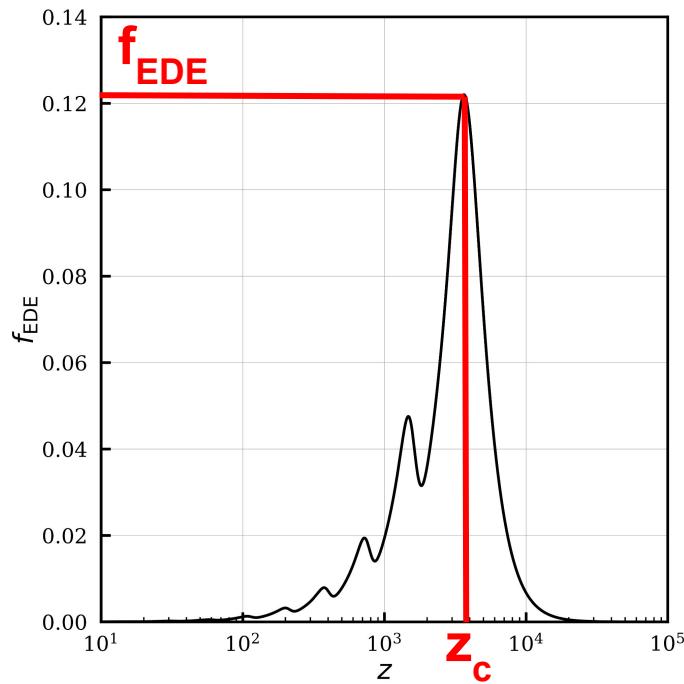


Early Dark Energy : phenomenological parametrization ²⁴

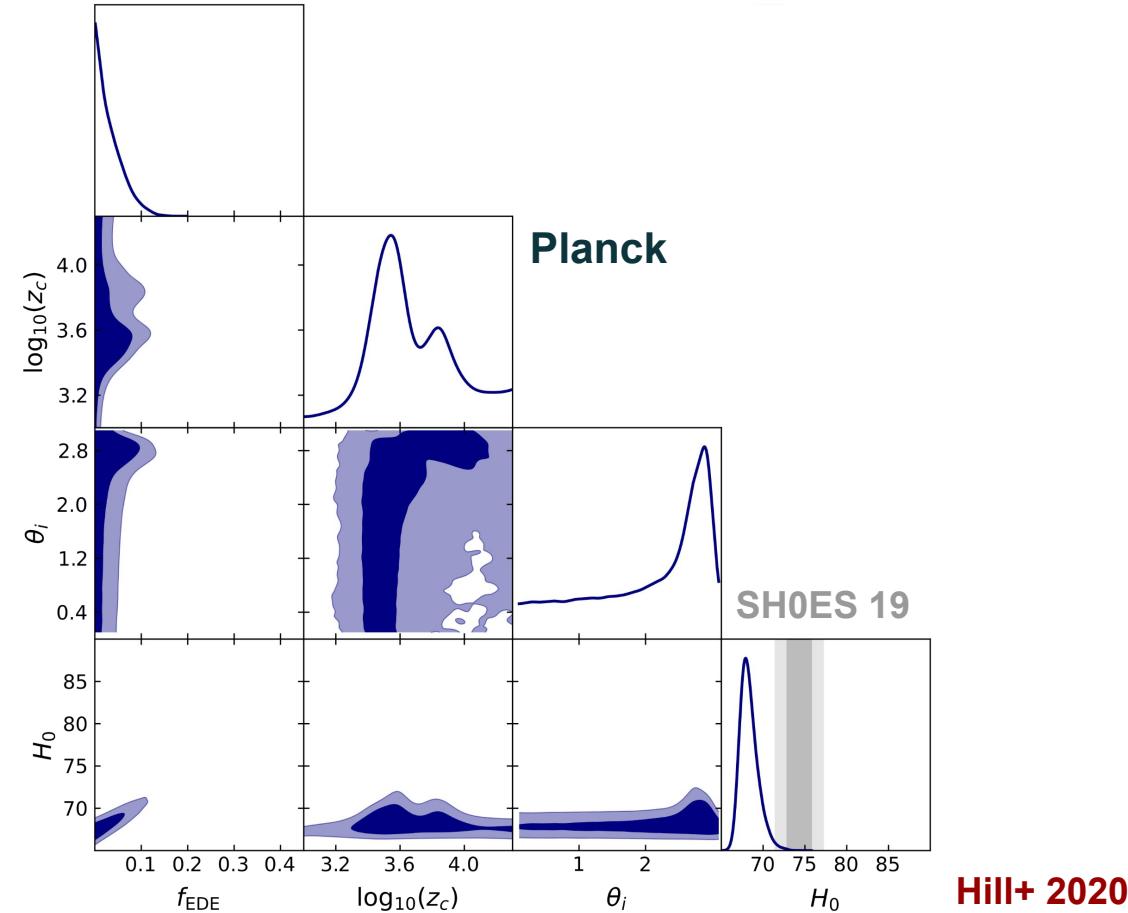
$$(m, f, \phi_i) \longrightarrow (f_{\text{EDE}}, z_c, \phi_i)$$

Field parameters

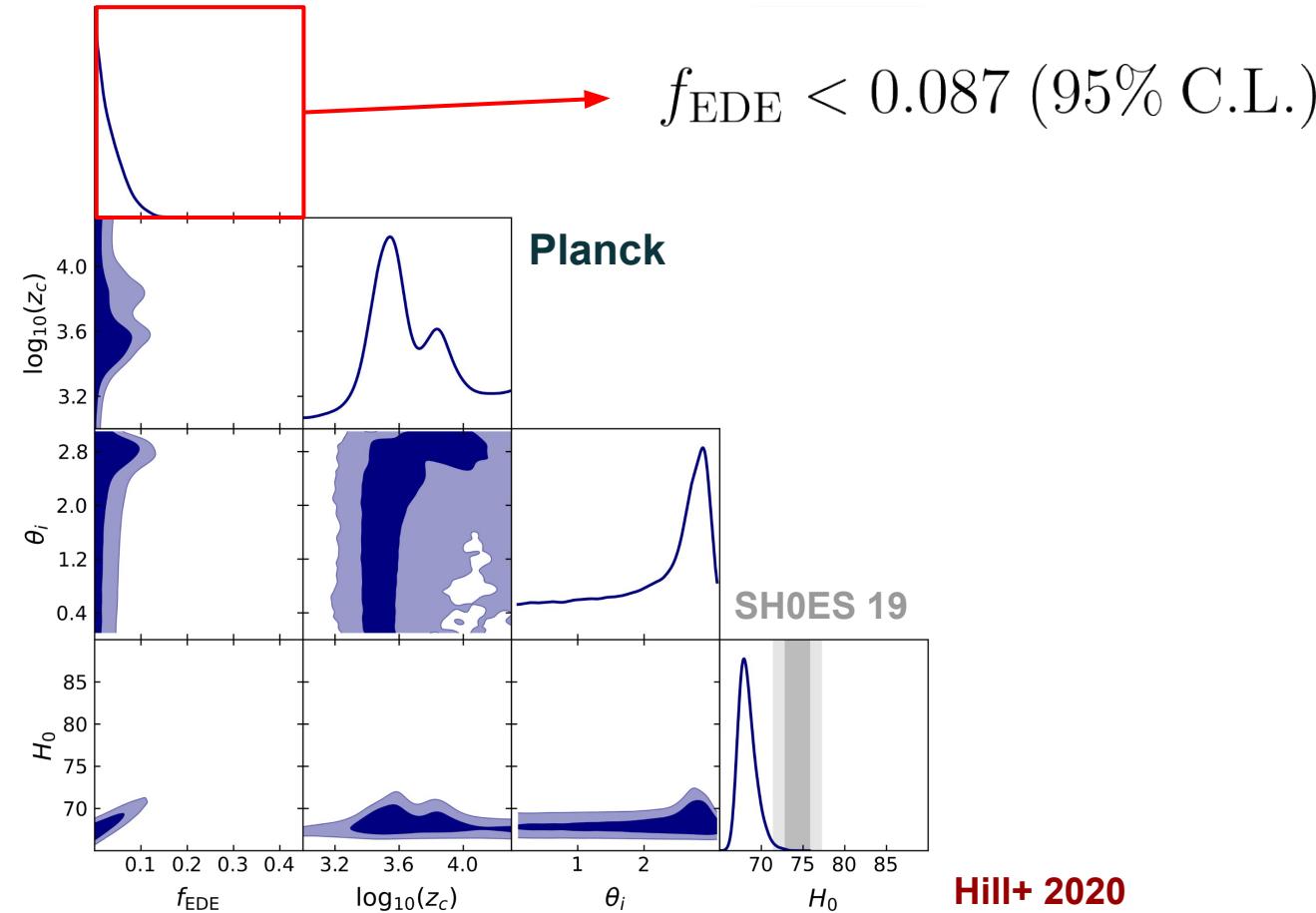
Phenomenological
parametrization



Constraints on EDE from Planck data

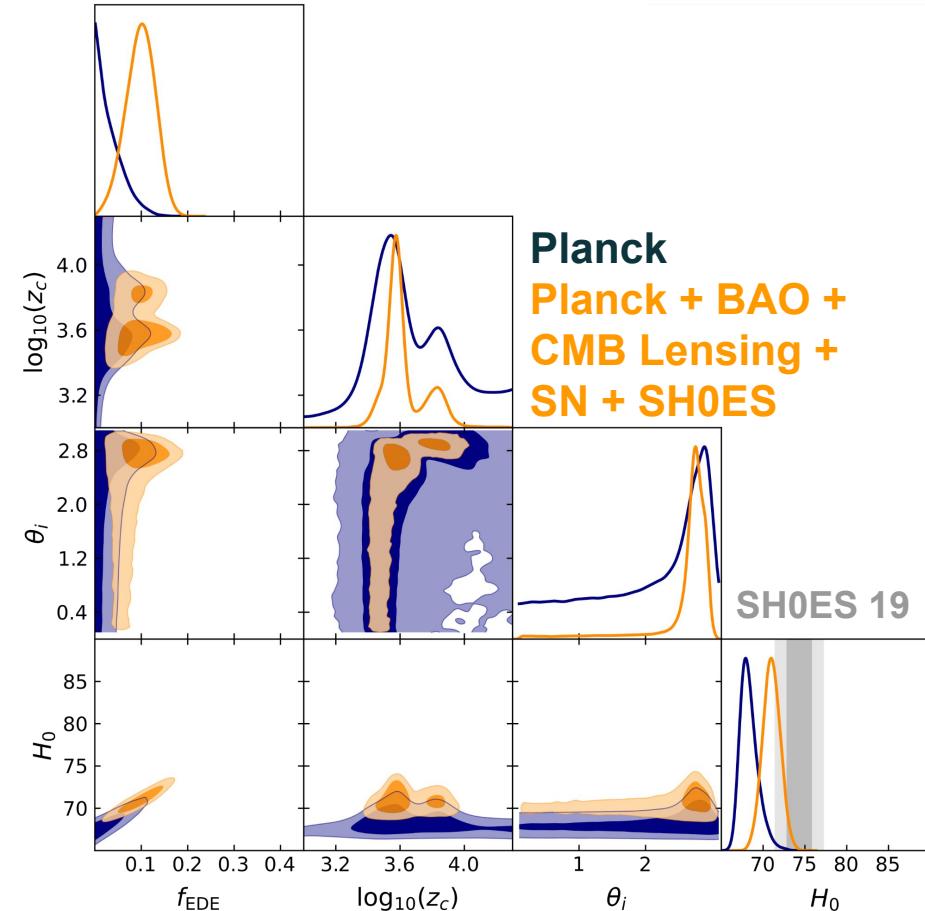


Constraints on EDE from Planck data



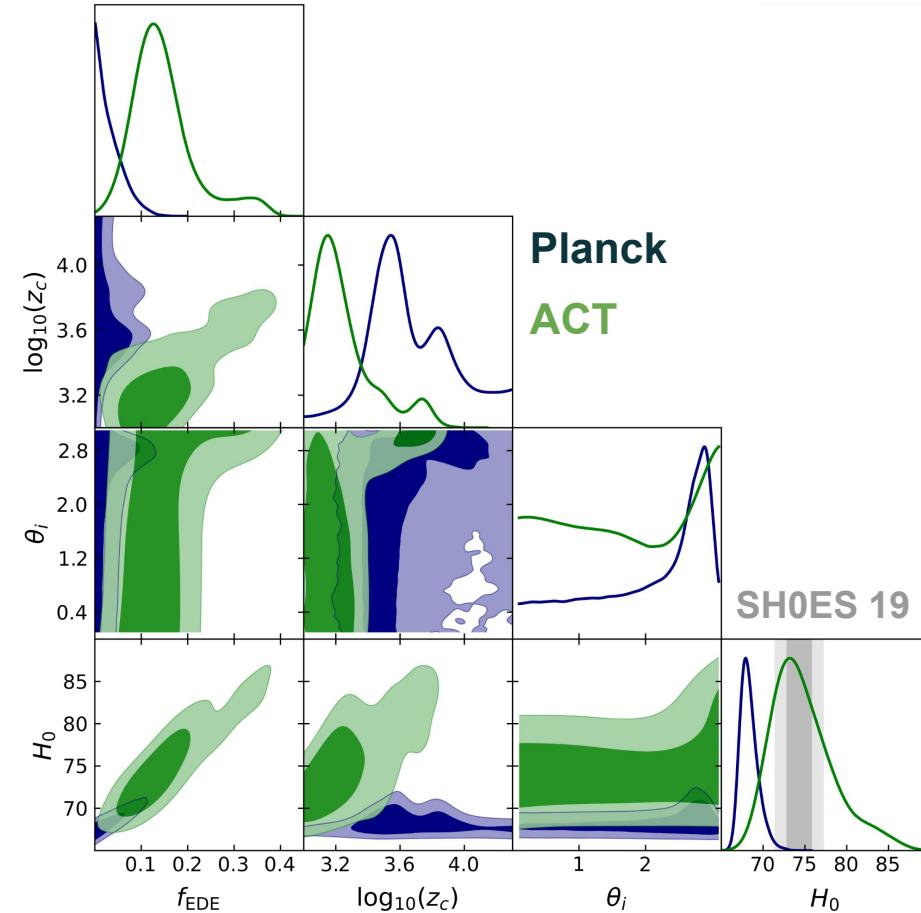
Results for a combination of Planck and SH0ES data

26



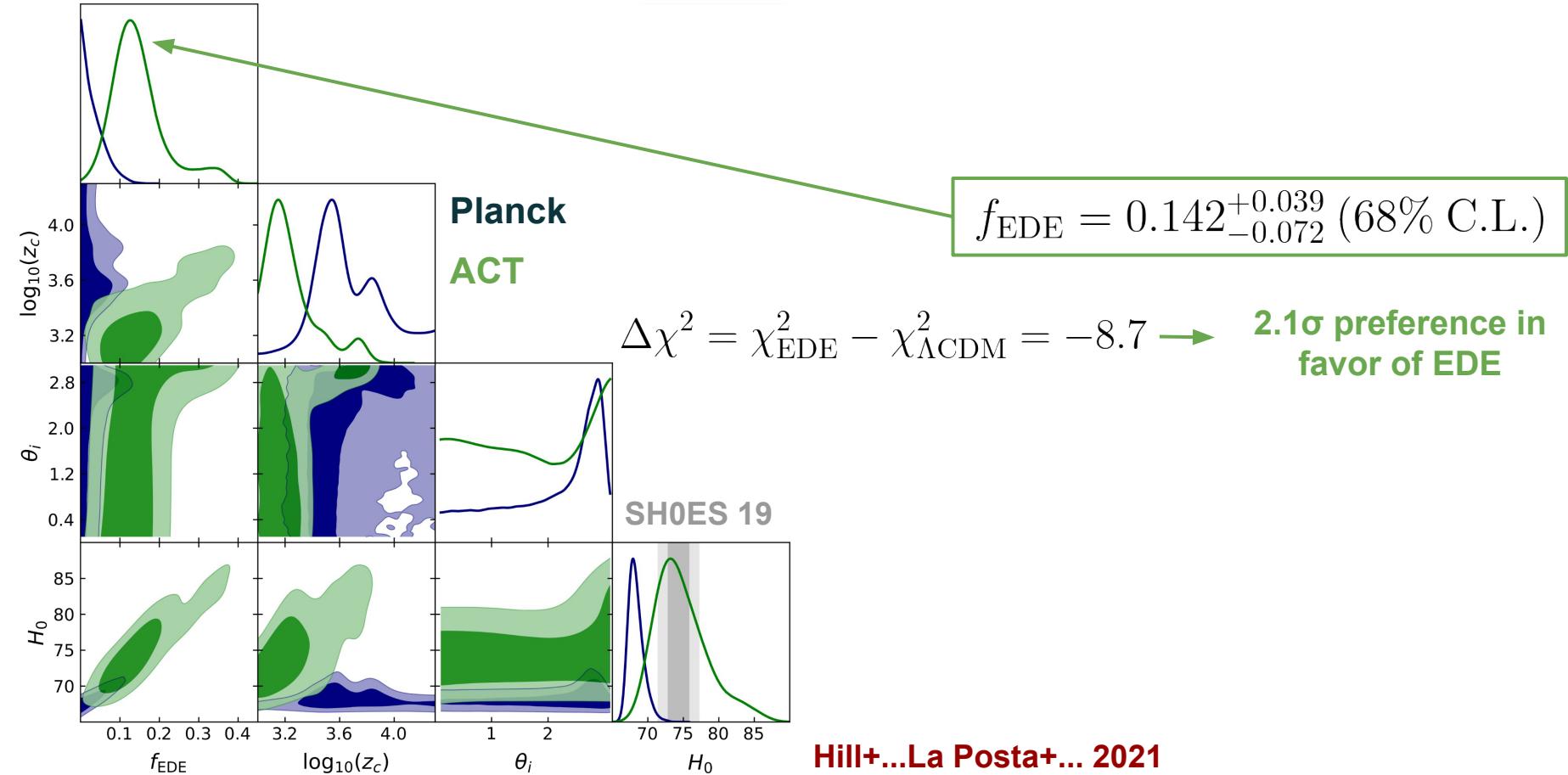
Poulin+ 2019, Smith+ 2019, Hill+ 2020

Additional constraints from ACTPol



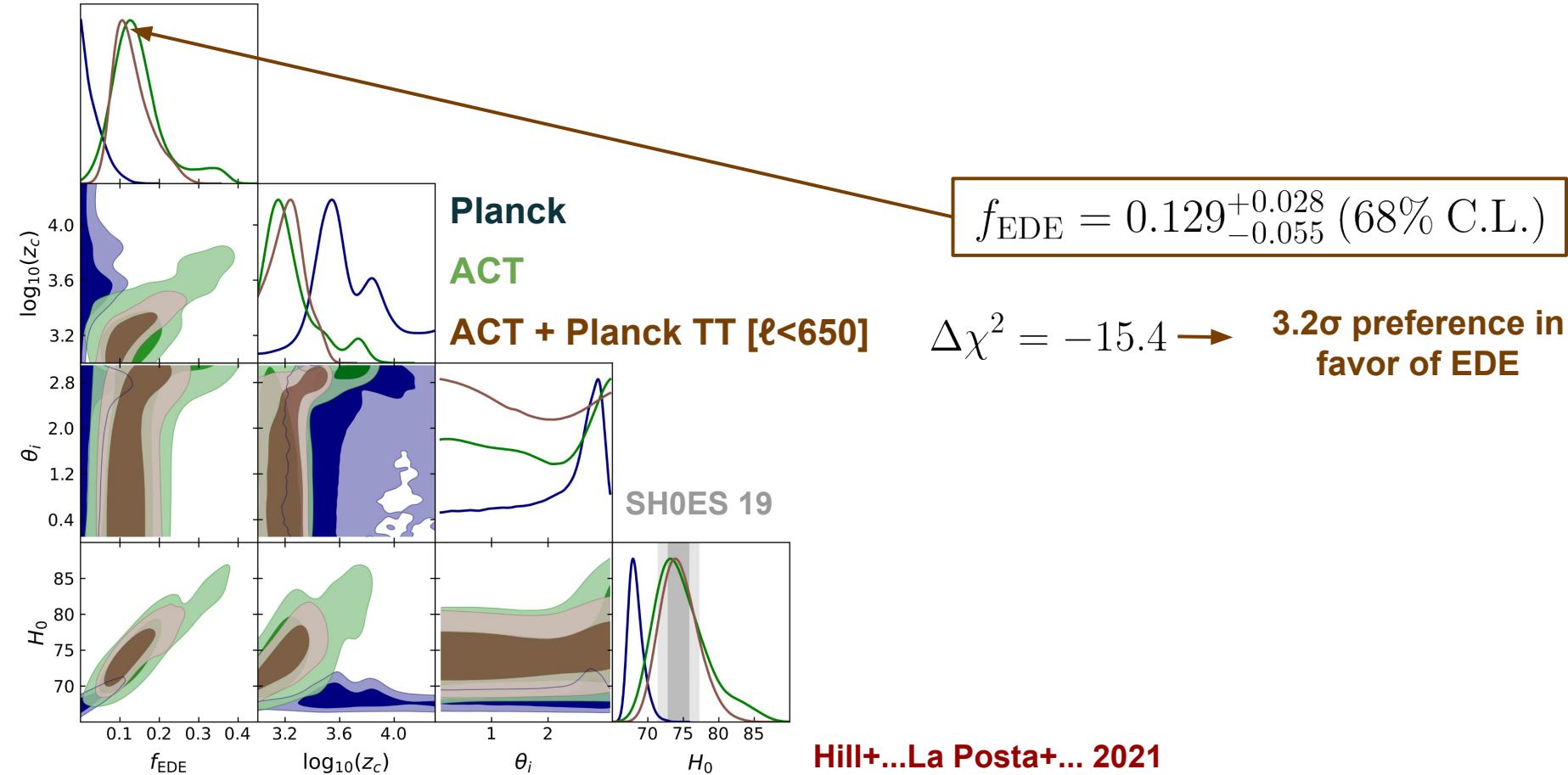
Additional constraints from ACTPol

27



Additional constraints from ACTPol

28



Summary of EDE results

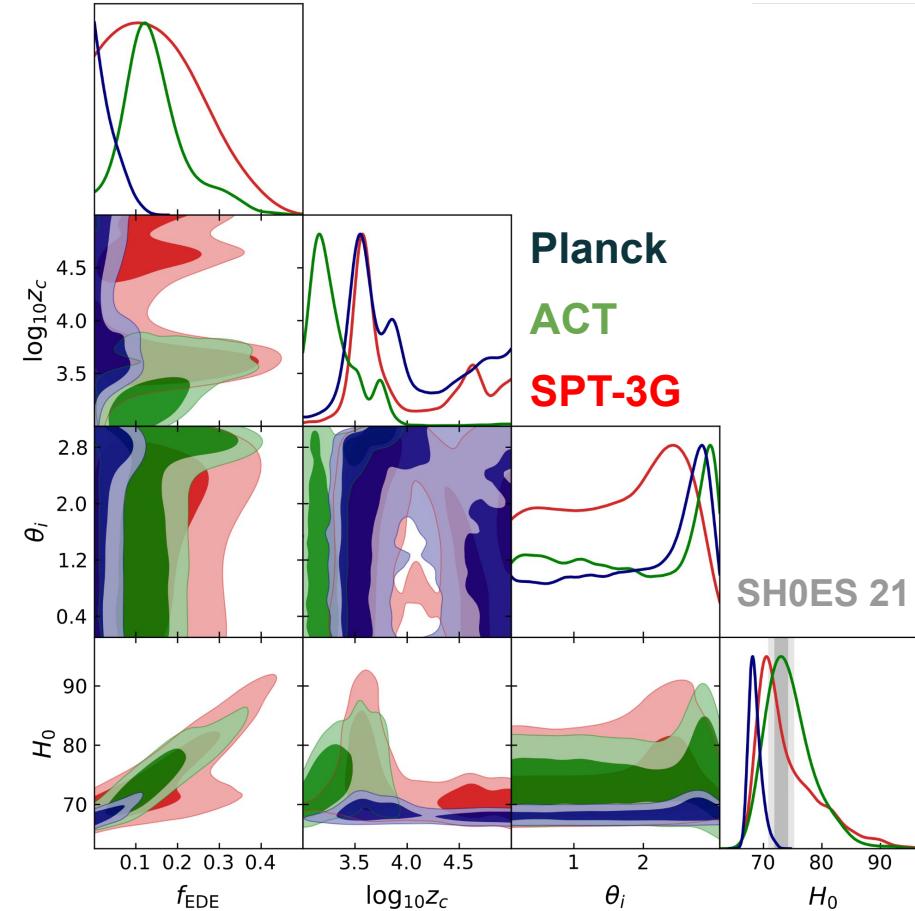
- Planck data alone don't favor high f_{EDE} values (Hill+ 2020)
- Planck data in combination with SH0ES show a preference for non-zero f_{EDE} (Poulin+ 2019, Smith+ 2019)
- ACT data alone favors EDE over Λ CDM (Hill+...La Posta+... 2021)

Summary of EDE results

- Planck data alone don't favor high f_{EDE} values (Hill+ 2020)
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- Motivates an analysis of EDE with public SPT-3G data

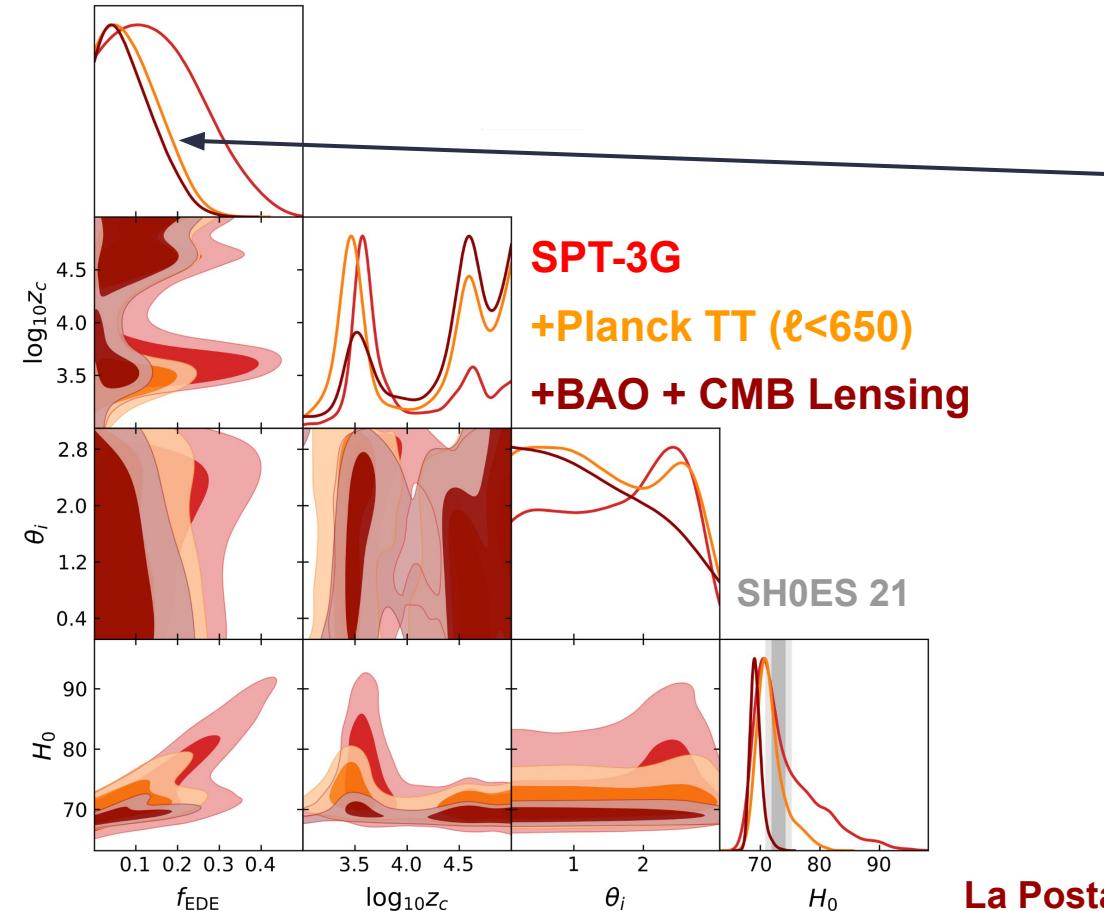
New results from SPT-3G public data

30



New results from SPT-3G public data

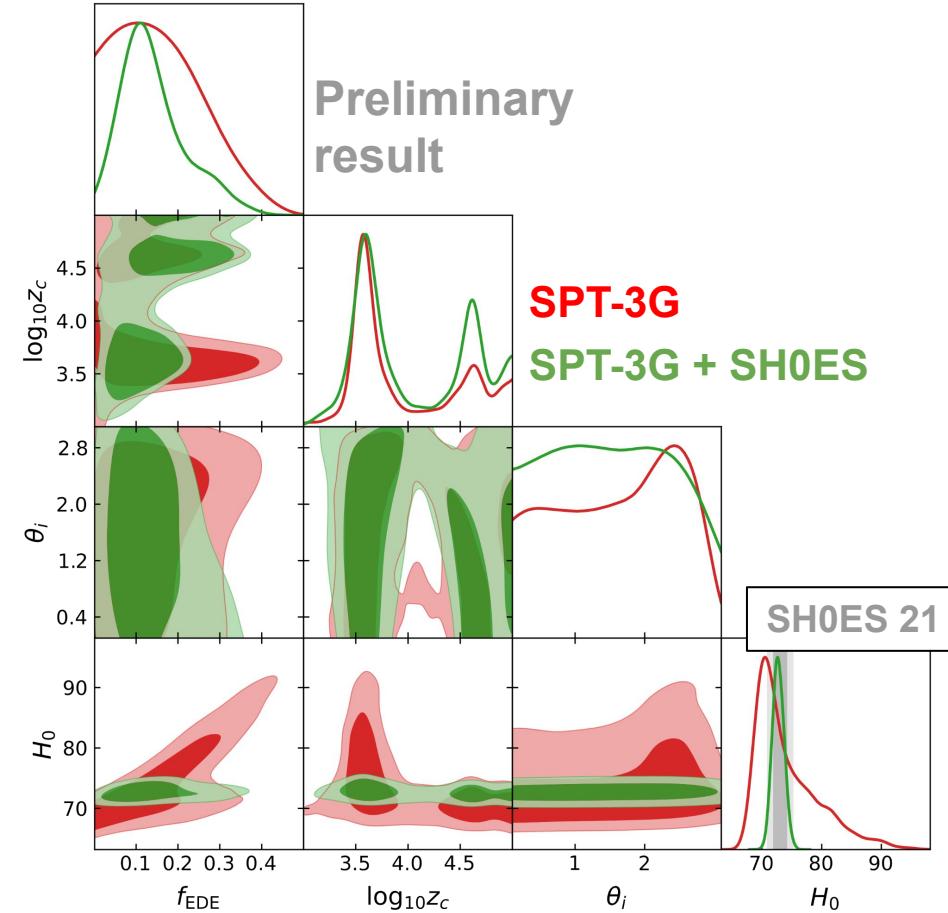
31



We tighten the constraint on f_{EDE} when we combine SPT3G and Planck TT ($\ell < 650$) or when we add LSS probes

Combining with SH0ES constraint

32



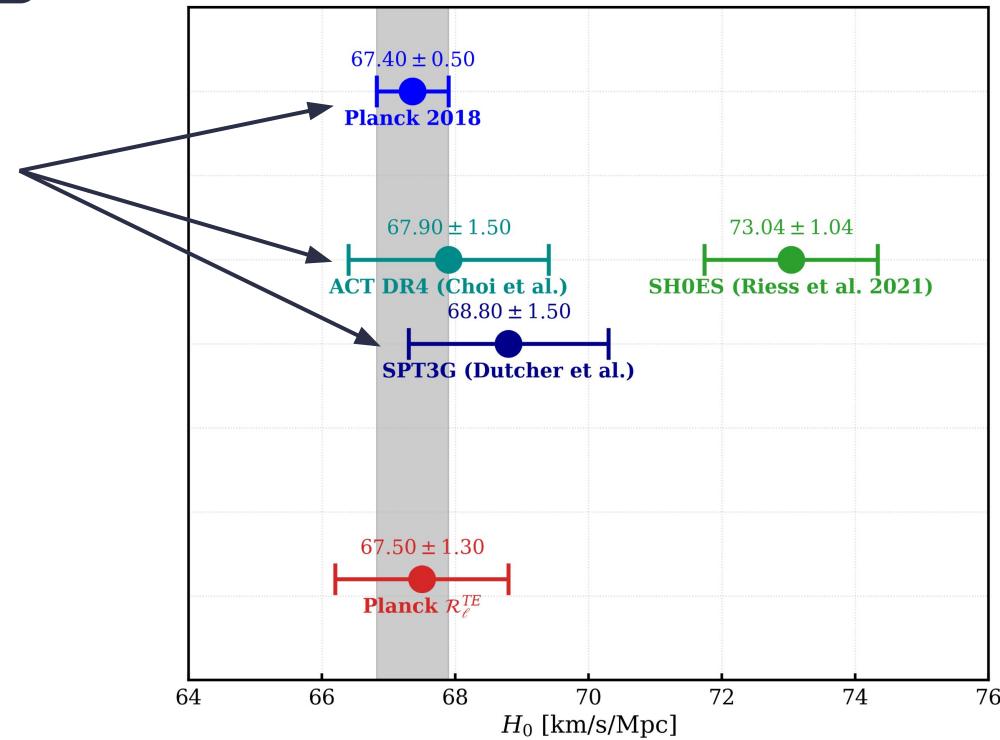
$$\Delta\chi^2_{\text{SPT-3G}} = -6.3$$

improvement of the fit to
SPT-3G data (with
respect to Λ CDM)

Conclusions

Option 2 : Systematics in CMB data

- 3 independent measurements of H_0

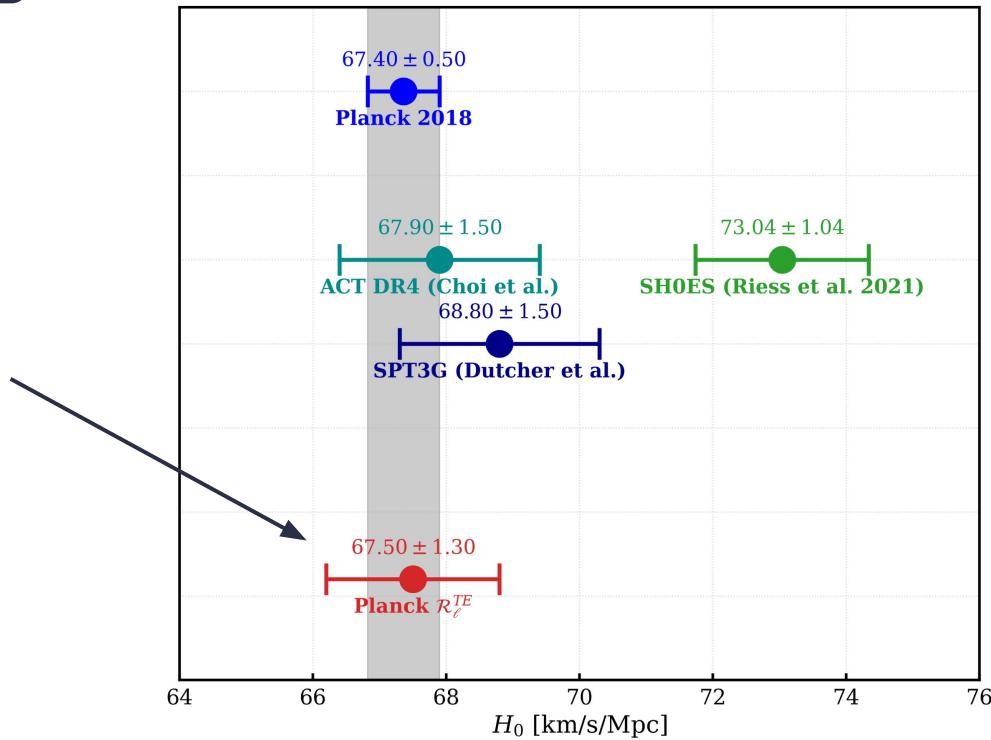


Conclusions

Option 2 : Systematics in CMB data

- 3 independent measurements of H_0
- Constraint from the **correlation coefficient** : insensitive to multiplicative systematic effects

It's hard to solve the Hubble tension with systematics in CMB data



Conclusions

Option 3 : Beyond Λ CDM physics - Early Dark Energy

- Planck data alone do not favor high f_{EDE} values
- Planck + SH0ES show a preference for $f_{\text{EDE}} \sim 10\%$
- ACT DR4 data favors EDE over Λ CDM (with $f_{\text{EDE}} \sim 10\%$)
- SPT-3G is not as constraining as ACT and Planck : but sees some degree of EDE when combined with SH0ES

