
An Exploration of Multigrid Methods

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Abstract

We used both one and two dimensional poisson problems to study multigrid methods for solving partial differential equations. Using iterative solvers for linear systems we show how coarsening the discretization can lead to approximations which converge to the true solution of the PDE with fewer iterations of the solver.

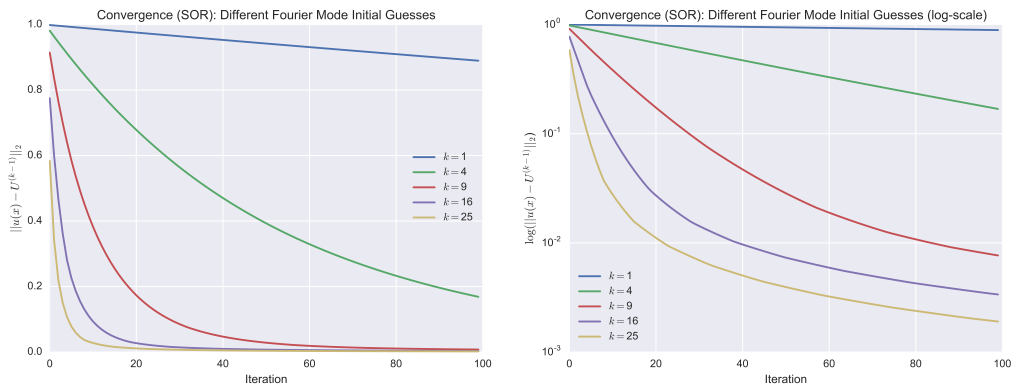
1 Motivation

The motivation for using coarser grids, while convergence analysis shows that finer grids should lead to more accurate approximations, comes from an observation about waves, discretizations, and aliasing. The iterative methods, often referred to as smoothers, smooth out error such that in the earlier iterations it is the high frequency components of the error that vanish first. As the algorithms sweep more times the error gets smoother, containing lower frequencies and tending toward zero. The trouble with the classical iterative methods like this is that low frequencies in the error can take many iterations to smooth. In general these iterative methods are $\mathcal{O}(n^2)$. Thus, some improvement in speed is desired.

The standard iterative methods for solving elliptic PDEs, often have a property in which the smoothing rate is not only small but also independent of the grid spacing h . Looking at a numerical example in which we apply weighted Jacobi iteration with $\omega = \frac{2}{3}$ to the problem $u''(x) = 0$ with zero boundary conditions gives us insight into this phenomena. Our initial guess is a fourier mode that looks like:

$$w_j = \sin\left(\frac{jk\pi}{N}\right) \quad 0 \leq j \leq N \text{ and } 1 \leq k \leq N-1$$

Here $N-1$ is the number of interior points. A convergence analysis gives us:



We see that error clearly decreases with each relaxation sweep but that the rate of decrease is much more faster for high frequency initial guesses than low frequency ones. Of course the first way to improve a relaxation is to increase the number of iterations, or improve the initial guess, however, relaxations on a coarser grid is less expensive. Moreover, smooth components on a fine grid projected onto a coarser grid look more

oscillatory. If we look at fourier modes associated with a mesh of discretization h , then we have:

$$w_{k,2j}^h = \sin\left(\frac{2jk\pi}{N}\right) = \sin\left(\frac{jk\pi}{N/2}\right) = w_{k,j}^{2h} \quad 1 \leq k < \frac{N}{2}$$

. This seemingly trivial manipulation reveals that in passing from a fine grid to coarse grid, the smooth wave becomes more sinusoidal and gains frequency.

From the Shannon Sampling Theorem we know that to retain all of the wave information, we need the discretization to have just over two points per wavelength.¹ The implication of this on our work is that the highest frequency in the error is determined by the mesh grid. Knowing this, we can coarsen the grid so that we have fewer sample points of the error function and then the low frequencies will be among the higher ones still contained in the coarse-grid error. By smoothing the error on this coarser grid we can eliminate more components of the error than we could without the coarsening. We leave the explanation and implementation of this for later sections.

2 Multigrid Method

notes: derive residual equation, and the role of the residual equations

In this section, we will consider the theory and analyze why multigrid methods work so well. We will delve deep into gaining some understanding of the linear algebraic implications of mesh restriction and interpolation. More importantly, we will look at the spectral picture of the multigrid method. The standard cycle of multigrids is the standard V cycle which looks like:

[insert diagram of a V cycle]

Transitions from one discretization to another discretization require a mapping of some kind from the current approximation of the true solution or vector to the same vector on the next mesh. We will adopt the same notation that Briggs introduces. Ω^h represents a grid with spacing $h = \Delta x$, and naturally Ω^{2h} represents a coarser grid. \mathbf{v}^h is an approximation of the true solution \mathbf{u}^h , and therefore error is $\mathbf{e}^h = \mathbf{u}^h - \mathbf{v}^h$. The interpolation matrix maps a vector on a coarse mesh, say Ω^{2h} to Ω^h and is denoted T_{2h}^h , while the restriction matrix maps a vector from $\Omega^h \rightarrow \Omega^{2h}$ and is denoted by R_h^{2h} .

Assuming, we have $N+1$ total points in our boundary, and $N-1$ interior points, then the restriction matrix R_h^{2h} maps $\mathbb{R}^{N-1} \rightarrow \mathbb{R}^{N/2-1}$. Since we have $N-1$ unknown values and are mapping to a lower dimension subspace, then the rank of our restriction matrix is $\frac{N}{2} - 1$ and the null space of the restriction matrix is $\frac{N}{2}$. Let us explore how this matrix transforms the eigenvalues of our iterative matrix A^h on mesh Ω^h . The eigenvectors or modes of A^h are given by

$$\mathbf{w}_k^h = \left(\sin\left(\frac{k\pi}{N}\right), \dots, \sin\left(\frac{(N-1)k\pi}{N}\right) \right)$$

¹C. E. Shannon, "Communications in the presence of noise", Proc. IRE, vol. 37, pp. 10-21, Jan. 1949.

where the j th component of the vector is $w_{k,j}^h = \sin\left(\frac{jk\pi}{N}\right)$. When the restriction operator is applied directly to these vectors we get:

$$R_h^{2h} \mathbf{w}_k^h = \frac{1}{4} \begin{bmatrix} 1 & 2 & 1 & & \\ & 1 & 2 & 1 & \\ & & \ddots & \ddots & \ddots \\ & & & 1 & 2 & 1 \end{bmatrix} \mathbf{w}_k^h \quad (1)$$

The j th component of this is:

$$\Rightarrow (R_h^{2h} \mathbf{w}_k^h)_j = \frac{1}{4} \left(\sin\left(\frac{(j-1)k\pi}{N}\right) + 2 \sin\left(\frac{jk\pi}{N}\right) + \sin\left(\frac{(j+1)k\pi}{N}\right) \right) \quad (2)$$

Using the following two trigonometric identities, $\sin(\theta_1 + \theta_2) = \sin \theta_1 \cos \theta_2 + \cos \theta_1 \sin \theta_2$ and $\cos(2\theta) = 2 \cos^2 \theta - 1$, leads us to:

$$\frac{1}{4} \left[2 \sin\left(\frac{2jk\pi}{N}\right) \cos\left(\frac{k\pi}{N}\right) + 2 \sin\left(\frac{2jk\pi}{N}\right) \right] = \sin\left(\frac{jk\pi}{N/2}\right) \cos^2\left(\frac{k\pi}{2N}\right) \text{ for } 1 \leq k \leq \frac{N}{2}$$

What we find is that this result holds for the first $k/2$ eigenvectors or one half of the unknowns.

$$R_h^{2h} \mathbf{w}_k^h = \cos^2\left(\frac{k\pi}{2N}\right) \mathbf{w}_k^{2h}, \quad 1 \leq k \leq \frac{N}{2} \quad (3)$$

The restriction matrix produces a new k th mode of A^{2h} such that the old k th mode of A^h gets multiplied by a constant.

Applying a similar analysis for when $\frac{N}{2} \leq k \leq N-1$, yields

$$R_h^{2h} \mathbf{w}_{k'}^h = -\sin^2\left(\frac{k\pi}{2N}\right) \mathbf{w}_k^{2h}, \quad 1 \leq k \leq \frac{N}{2} \text{ where } k' = N - k$$

Essentially, what is going on here is the phenomena of aliasing since the restriction matrix is acting on the $(N-k)$ th mode of A^h and multiplying it by a constant multiple of the . The interpretation behind this result is that the oscillatory modes on Ω^h cannot be represented in the coarser mesh Ω^{2h} and the restriction operator transforms these nodes into relatively smooth modes on the coarser mesh. The reason this matrix has no inverse is because there are two eigenvectors (eigenvector k and eigenvector $N-k$) on A^h that get mapped to the same k th mode on A^{2h} .

INCLUDE SOME GRAPHS HERE TO DESCRIBE THIS ALIASING

We define a set W_k^h that takes the span of these complementary nodes (the $(N-k)$ th eigenvector and the k th eigenvector), $W_k^h = \text{span}\{\mathbf{w}_k^h, \mathbf{w}_{N-k}^h\}$. We note that:

$$w_{N-k,j}^h = \sin\left(\frac{j(N-k)\pi}{N}\right) = \sin\left(j\pi - \frac{jk\pi}{N}\right)$$

$$\sin(j\pi) \cos\left(\frac{jk\pi}{N}\right) - \cos(j\pi) \sin\left(\frac{jk\pi}{N}\right) = (-1)^{j+1} \sin\left(\frac{jk\pi}{N}\right) = (-1)^{j+1} w_{k,j}^h \quad (4)$$

INCLUDE IMAGE OF COMPLEMENTARY PAIRS

Of course, this two element to one element mapping means that R_h^{2h} contains non-trivial elements in its null space. Let's take the unit vector \mathbf{e}_i^h on Ω^h , and see how it gets transformed by R_h^{2h} .

$$R_h^{2h} \mathbf{e}_i^h = \begin{bmatrix} 1 & 2 & 1 & & & \\ & & 1 & 2 & 1 & \\ & & & \ddots & \ddots & \ddots \\ & & & & 1 & 2 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ \vdots \\ 1 \\ \vdots \\ 0 \end{bmatrix} \quad (5)$$

We notice that there are two cases due to overlapping ones in the odd columns of R_h^{2h} . If i is odd in \mathbf{e}_i^h , then we have a contribution from two rows whereas if i is even, there is only one row contribution from only the component with 2. We have the following:

$$R_h^{2h} \mathbf{e}_i^h = \begin{cases} \frac{1}{4} \left(\mathbf{e}_{\frac{i-1}{2}}^{2h} + \mathbf{e}_{\frac{i+1}{2}}^{2h} \right) & i \text{ odd} \\ \frac{1}{4} \left(2\mathbf{e}_{\frac{i}{2}}^{2h} \right) & i \text{ even} \end{cases} \quad (6)$$

These unit vectors do not form the basis of the null space because this matrix product is not zero. How about we look at the transformation of these unit vectors under the iterative method first, namely $R_h^{2h} A^h \mathbf{e}_i^h$.

$$R_h^{2h} A^h \mathbf{e}_i^h = R_h^{2h} (-\mathbf{e}_{i-1}^h + 2\mathbf{e}_i^h - \mathbf{e}_{i+1}^h) \quad (7)$$

If the i is even then $i-1$ and $i+1$ are odd and we'll get extra terms that do not go to zero, whereas if i is odd, we get:

$$R_h^{2h} A^h \mathbf{e}_i^h = -\frac{1}{4} \left(2\mathbf{e}_{\frac{i-1}{2}}^{2h} \right) + \frac{1}{2} \left(\mathbf{e}_{\frac{i-1}{2}}^{2h} + \mathbf{e}_{\frac{i+1}{2}}^{2h} \right) - \frac{1}{2} \mathbf{e}_{\frac{i+1}{2}}^{2h} = 0 \quad (8)$$

We see that the vectors $\mathbf{n}_i = A^h \mathbf{e}_i^h$ when i is odd form the basis of the null space of the full weighting operator. The null space contains both smooth and oscillatory modes of A^h because the vector $A^h \mathbf{e}_i^h$, which is a vector in the range of A^h , can be decomposed as a linear combination of the eigenvectors mentioned above.

Turning our attention to the interpolation matrix, T_{2h}^h , which has a reverse mapping from $\mathbb{R}^{N/2-1} \rightarrow \mathbb{R}^{N-1}$ and has full rank because there are fewer columns than rows. Like before, we want to ask how this matrix affects the eigenvectors or modes of matrix A^{2h} . Note that this time, we are considering matrix A^{2h} to be consistent with the analysis in restriction. We know that the j th component of the k th eigenvector of A^{2h} is:

$$w_{k,j}^{2h} = \sin \left(\frac{jk\pi}{N/2} \right) \text{ where } 1 \leq k \leq \frac{N}{2} \text{ and } 0 \leq j \leq \frac{N}{2}$$

The j th component of $T_{2h}^h \mathbf{w}_k^{2h}$ looks like:

$$T_{2h}^h \mathbf{w}_k^{2h} = \frac{1}{2} \begin{bmatrix} 1 & & & \\ 2 & & & \\ 1 & 1 & & \\ & 2 & & \\ & 1 & 1 & \\ & & \ddots & \ddots \end{bmatrix} \mathbf{w}_k^{2h}$$

$$(T_{2h}^h \mathbf{w}_k^{2h})_j = \begin{cases} \sin\left(\frac{k(j/2)\pi}{N/2}\right) & j \text{ even} \\ \frac{1}{2} \sin\left(\frac{k\frac{j-1}{2}\pi}{N/2}\right) + \frac{1}{2} \sin\left(\frac{k\frac{j+1}{2}\pi}{N/2}\right) & j \text{ odd} \end{cases}$$

For j odd we have:

$$(T_{2h}^h \mathbf{w}_k^{2h})_j = \frac{1}{2} \sin\left(\frac{k(j-1)\pi}{N}\right) + \frac{1}{2} \sin\left(\frac{k(j+1)\pi}{N}\right)$$

Using the trigonometric sum identity and double angle cosine identity $2\cos^2\theta - 1 = \cos 2\theta$, we get that for the j is odd case:

$$\begin{aligned} &\Rightarrow \sin\left(\frac{kj\pi}{N}\right) \cos\left(\frac{k\pi}{N}\right) = \sin\left(\frac{kj\pi}{N}\right) \left(2\cos^2\left(\frac{k\pi}{2N}\right) - 1\right) \\ &\Rightarrow \sin\left(\frac{kj\pi}{N}\right) \cos^2\left(\frac{k\pi}{2N}\right) + \sin\left(\frac{kj\pi}{N}\right) \left(\cos^2\left(\frac{k\pi}{2N}\right) - 1\right) \text{ and we have } k = N - k' \\ &\Rightarrow \sin\left(\frac{kj\pi}{N}\right) \cos^2\left(\frac{k\pi}{2N}\right) - \sin\left(\frac{(N-k')j\pi}{N}\right) \sin^2\left(\frac{k\pi}{2N}\right) \\ &\Rightarrow \sin\left(\frac{kj\pi}{N}\right) \cos^2\left(\frac{k\pi}{2N}\right) - \sin^2\left(\frac{k\pi}{2N}\right) \left[\sin(j\pi) \cos\left(\frac{k'j\pi}{N}\right) - \cos(j\pi) \sin\left(\frac{k'j\pi}{N}\right)\right] \\ &\Rightarrow \sin\left(\frac{kj\pi}{N}\right) \cos^2\left(\frac{k\pi}{2N}\right) - \sin^2\left(\frac{k\pi}{2N}\right) \left[(-1)^{j+1} \sin\left(\frac{k'j\pi}{N}\right)\right] \text{ and because } j \text{ is odd ...} \\ &\cos^2\left(\frac{k\pi}{2N}\right) \sin\left(\frac{kj\pi}{N}\right) - \sin^2\left(\frac{k\pi}{2N}\right) \sin\left(\frac{k'j\pi}{N}\right) \end{aligned} \tag{9}$$

For the even case, we can do the following:

$$(T_{2h}^h \mathbf{w}_k^{2h})_j = \sin\left(\frac{kj\pi}{N}\right) = \cos^2\left(\frac{k\pi}{2N}\right) \sin\left(\frac{kj\pi}{N}\right) + \sin^2\left(\frac{k\pi}{2N}\right) \sin\left(\frac{kj\pi}{N}\right)$$

In the previous page we showed that complementary nodes are related by just a sign. Using this we can turn the above expression to take on a similar form. Just like Briggs, we'll define $c_k = \cos^2\left(\frac{k\pi}{2N}\right)$ and $s_k = \sin^2\left(\frac{k\pi}{2N}\right)$. Therefore, for both the even and odd case we can express the new eigenvectors as:

$$[T_{2h}^h \mathbf{w}_k^{2h}]_j = \cos^2 \left(\frac{k\pi}{2N} \right) \sin \left(\frac{kj\pi}{N} \right) - \sin^2 \left(\frac{k\pi}{2N} \right) \sin \left(\frac{kj\pi}{N} \right) \Rightarrow T_{2h}^h \mathbf{w}_k^{2h} = c_k \mathbf{w}_k^h - s_k \mathbf{w}_k^h \quad (10)$$

The takeaway of all this tedious, trigonometric algebra is that the interpolation matrix in some way excites the k th mode in Ω^{2h} to produce two modes on the finer mesh Ω^h .

A basis for the range of interpolation is of course given by the columns of T_{2h}^h because it is full rank (all of its columns are linearly independent). Just like the null space of the restriction matrix, these basis vectors can be decomposed as a linear combination of the modes of A^h and hence reveal that the column space of the interpolation matrix contains both smooth and oscillatory nodes of A^h (see Briggs).

The purpose of all this analysis on these two matrices subspaces and their affect on the modes of A^h is because the crux of multigrid methods is the *coarse grid correction scheme*. This scheme is a residual correction scheme that maps the residual of an iterative method on some mesh Ω^h to a coarser mesh to solve for error. This error is then interpolated upwards to obtain a better approximation. Essentially, coarse grid correction (CGC) is a series of matrix products.

1. We start with an approximation, \mathbf{v}^h to our true solution, \mathbf{u}^h and relax ν times.
 $\mathbf{v}^h \leftarrow P^\nu \mathbf{v}^h$, where P is a relaxation scheme
2. Obtain a residual \mathbf{r}^h and restrict it: $\mathbf{r}^{2h} = R_h^{2h}(\mathbf{f}^h - A^h \mathbf{v}^h)$
3. Solve the residual equation: $\mathbf{v}^{2h} = (A^{2h})^{-1} \mathbf{r}^{2h}$
4. Correct the approximation on the original mesh Ω^h : $\mathbf{v}^h \leftarrow \mathbf{v}^h + T_{2h}^h \mathbf{v}^{2h}$

These 4 steps can be summarized as a series of matrix products:

$$\mathbf{v}^h \leftarrow P^\nu \mathbf{v}^h + T_{2h}^h (A^{2h})^{-1} R_h^{2h} (\mathbf{f}^h - A^h P^\nu \mathbf{v}^h) \quad (11)$$

Naturally, this must satisfy:

$$\mathbf{u}^h = P^\nu \mathbf{u}^h + T_{2h}^h (A^{2h})^{-1} R_h^{2h} (\mathbf{f}^h - A^h P^\nu \mathbf{u}^h) \quad (12)$$

Taking their difference and considering this difference as the coarse grid correction operator (as Briggs does, page 74), we get a matrix operator:

$$[I - T_{2h}^h (A^{2h})^{-1} R_h^{2h} A^h] P^\nu \quad (13)$$

Considering this whole series of products and sums as one operation, Briggs states that in the scenario in which we do not relax, $\nu = 0$, then the subspaces $W_k^h = \text{span}\{w_k^h, w_{k'}^h\}$ is an invariant subspace.

We will endeavour to prove this. Without any relaxation, our operator is:

$$I - T_{2h}^h (A^{2h})^{-1} R_h^{2h} A^h$$

In a homework problem, we found that k th eigenvalue of the matrix A^h is $\lambda_k^h = 4 \sin^2 \left(\frac{k\pi}{2N} \right)$. We therefore have:

1. $A^h \mathbf{w}_k^h = \lambda_k^h \mathbf{w}_k^h$
2. $R_j^{2h} A^h \mathbf{w}_k^h = \lambda_k^h \cos^2 \left(\frac{k\pi}{2N} \right) \mathbf{w}_k^h$
3. $(A^{2h})^{-1} R_j^{2h} A^h \mathbf{w}_k^h = \lambda_k^h \cos^2 \left(\frac{k\pi}{2N} \right) (A^{2h})^{-1} \mathbf{w}_k^{2h} = \frac{\lambda_k^h}{\lambda_k^{2h}} \cos^2 \left(\frac{k\pi}{2N} \right) \mathbf{w}_k^{2h}$
4. $T_{2h}^h (A^{2h})^{-1} R_j^{2h} A^h \mathbf{w}_k^h = \frac{\lambda_k^h}{\lambda_k^{2h}} \cos^2 \left(\frac{k\pi}{2N} \right) T_{2h}^h \mathbf{w}_k^{2h} = \frac{\lambda_k^h}{\lambda_k^{2h}} \cos^2 \left(\frac{k\pi}{2N} \right) [c_k \mathbf{w}_k^h - s_k \mathbf{w}_{k'}^h]$
5. Our final expression for the change on \mathbf{w}_k^h by this coarse grid correction operator is:

$$\mathbf{w}_k^h - \frac{\lambda_k^h}{\lambda_k^{2h}} \cos^2 \left(\frac{k\pi}{2N} \right) [c_k \mathbf{w}_k^h - s_k \mathbf{w}_{k'}^h]$$

Plugging in our eigenvalues into this, we obtain:

$$\mathbf{w}_k^h - \frac{4 \sin^2 \left(\frac{k\pi}{2N} \right)}{4 \sin^2 \left(\frac{k\pi}{N} \right)} \cos^2 \left(\frac{k\pi}{2N} \right) [c_k \mathbf{w}_k^h - s_k \mathbf{w}_{k'}^h] \quad (14)$$

$$\left(1 - \frac{\sin^2 \left(\frac{k\pi}{2N} \right) \cos^2 \left(\frac{k\pi}{2N} \right)}{\sin^2 \left(\frac{k\pi}{N} \right)} \right) \mathbf{w}_k^h + \frac{\sin^4 \left(\frac{k\pi}{2N} \right) \cos^2 \left(\frac{k\pi}{2N} \right)}{\sin^2 \left(\frac{k\pi}{N} \right)} \mathbf{w}_{k'}^h \quad (15)$$

Using the double angle formula $\sin 2\theta = 2 \sin \theta \cos \theta$ and $\cos^2 \theta + \sin^2 \theta = 1$, we can show (a lot of algebra):

$$\begin{aligned} & \left(1 - \frac{\cos^2 \left(\frac{k\pi}{2N} \right)}{4} \right) \mathbf{w}_k^h + \frac{1}{4} \sin^2 \left(\frac{k\pi}{2N} \right) \mathbf{w}_{k'}^h \\ \Rightarrow & \left(\frac{3}{4} + \frac{1}{4} \sin^2 \left(\frac{k\pi}{2N} \right) \right) \mathbf{w}_k^h + \frac{1}{4} \sin^2 \left(\frac{k\pi}{2N} \right) \mathbf{w}_{k'}^h \end{aligned}$$

If we continue and expand things out, we can show this equals to, $s_k \mathbf{w}_k + s_k \mathbf{w}_{k'}$ and hence we obtain invariance of the element \mathbf{w}_k . We need to do this for the element $\mathbf{w}_{k'}$. To avoid the same type of analysis, we'll rely on Briggs to assert that we get $c_k \mathbf{w}_k + c_k \mathbf{w}_{k'}$. What this means is that we get smooth and oscillatory results when this coarse grid correction scheme is applied to these complementary modes.

The power of this method comes from the spectral properties of iterative solvers. When we introduce ν steps of relaxation of some iterative method, like Gauss-Seidel or SOR, then we get that our new operator becomes:

$$(I - T_{2h}^h (A^{2h})^{-1} R_h^{2h} A^h) P^\nu \quad (16)$$

Note, Following the same 5 steps above, we get that (letting λ_k be the eigenvalue of P associated with the k th eigenvector \mathbf{w}_k)

$$(I - T_{2h}^h (A^{2h})^{-1} R_h^{2h} A^h) P^\nu \mathbf{w}_k^h = (I - T_{2h}^h (A^{2h})^{-1} R_h^{2h} A^h) \lambda_k^\nu \mathbf{w}_k^h$$

What is new this time is a constant term - the eigenvalue! So we get similar results! The coarse grid correction scheme when operating on \mathbf{w}_k and $\mathbf{w}_{k'}$ returns:

$$\lambda_k^\nu s_k \mathbf{w}_k + \lambda_k^\nu s_k \mathbf{w}_{k'} \text{ and } \lambda_k^\nu c_k \mathbf{w}_k + \lambda_k^\nu c_k \mathbf{w}_{k'} \quad \text{where } 1 \leq k \leq \frac{N}{2}, \quad k' = N - k$$

The smoothing is the dominant term here! It is what makes the oscillatory modes small, while at the same time eliminating smooth modes, reflected in the s_k terms. When all terms are small particularly for small modes k relative to $N/2$ or when the number of relaxation increases, the end result is a complete process where both types of eigenvectors (low frequency and high frequency modes) are well damped. This is the spectral picture of the multigrid method

3 One Dimensional Problem

For the one dimensional problem we chose to demonstrate the Poisson problem with Dirichlet boundary conditions as follows:

$$\begin{aligned} \frac{d^2 u}{dx^2} &= f(x), \quad x \in (a, b), \\ u(a) &= \alpha, \\ u(b) &= \beta. \end{aligned} \tag{17}$$

The values of $f(x)$, a , b , α , and β are specified below for two different examples.

- Example #1:

$$f(x) = -\sin(x), \quad a = 0, \quad b = 2\pi, \quad \alpha = \beta = 0$$

- Example #2:

$$f(x) = e^x, \quad a = 0, \quad b = 1, \quad \alpha = 1, \quad \beta = e$$

The results of our method are presented in the Results section, however we will describe the implementation in the next section first.

3.1 Implementation

Regarding implementation, choosing an appropriate data structure is essential to achieving a high level efficiency otherwise, the benefits of the multigrid method is lost. Briggs, in *A Multigrid Tutorial*, adopts arrays as his primary data structure for one-dimensional problems. We adopted the same strategy and in general, any arbitrary grid will contain 2^{n-1} interior grid points plus two boundary points for a total of $2^n + 1$. The size of the arrays are transformed across different mesh grids.

3.2 Results

4 Two Dimensional Problem

4.1 Implementation

4.2 Results

5 Error Analysis

5.1 Multigrid

5.2 Gauss-Seidel

5.3 Successive Over Relaxation

6 Conclusion

7 Other