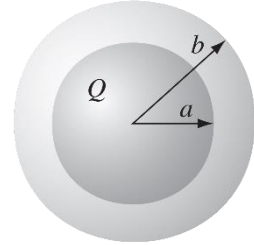


Department of Physics, Shiv Nadar Institution of Eminence
Spring 2025
PHY102: Introduction to Physics-II
Tutorial – 9

1. A spherical conductor of radius **a** carries a charge **Q** as shown in the figure. It is surrounded by linear dielectric material of susceptibility χ_e out to radius **b**. Find the energy of this configuration.



Q1)

$$\mathbf{D} = \begin{cases} 0, & (r < a) \\ \frac{Q}{4\pi r^2} \hat{\mathbf{r}}, & (r > a) \end{cases}, \quad \mathbf{E} = \begin{cases} 0, & (r < a) \\ \frac{Q}{4\pi\epsilon r^2} \hat{\mathbf{r}}, & (a < r < b) \\ \frac{Q}{4\pi\epsilon_0 r^2} \hat{\mathbf{r}}, & (r > b) \end{cases}.$$

$$\begin{aligned} W &= \frac{1}{2} \int \mathbf{D} \cdot \mathbf{E} \, d\tau = \frac{1}{2} \frac{Q^2}{(4\pi)^2} 4\pi \left\{ \frac{1}{\epsilon} \int_a^b \frac{1}{r^2} \frac{1}{r^2} r^2 dr + \frac{1}{\epsilon_0} \int_b^\infty \frac{1}{r^2} dr \right\} = \frac{Q^2}{8\pi} \left\{ \frac{1}{\epsilon} \left(\frac{-1}{r} \right) \Big|_a^b + \frac{1}{\epsilon_0} \left(\frac{-1}{r} \right) \Big|_b^\infty \right\} \\ &= \frac{Q^2}{8\pi\epsilon_0} \left\{ \frac{1}{(1 + \chi_e)} \left(\frac{1}{a} - \frac{1}{b} \right) + \frac{1}{b} \right\} = \boxed{\frac{Q^2}{8\pi\epsilon_0(1 + \chi_e)} \left(\frac{1}{a} + \frac{\chi_e}{b} \right)}. \end{aligned}$$

2. A thick spherical shell (inner radius a , outer radius b) is made of dielectric material with a “frozen-in” polarization :

$$\vec{P}(\vec{r}) = \frac{k}{r} \hat{r}$$

where k is a constant and r is the distance from the center (see **Fig. 1** below). (There is no *free* charge in the problem.) Find the electric field in all three regions by two different methods :

(a) Locate all the *bound charge*, and use the original Gauss’s law to calculate the field it produces.

(b) Find \vec{D} from the Gauss’s Law analogue for the Displacement, and then get \vec{E} from the defining equation for \vec{D} . [Notice that the *second method is much faster*, and it avoids any explicit reference to the bound charges.]

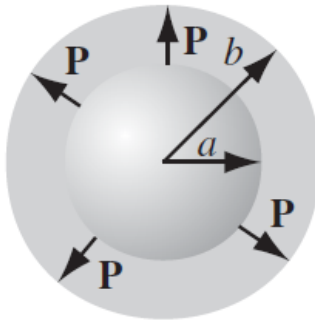


Fig. 1

Working in Spherical Polar Coordinates, we get :

$$(a) \rho_b = -\nabla \cdot \vec{P} = -\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{k}{r} \right) = -\frac{k}{r^2}; \quad \sigma_b = \vec{P} \cdot \hat{n} = \begin{cases} +\vec{P} \cdot \hat{r} = k/b & (\text{at } r = b), \\ -\vec{P} \cdot \hat{r} = -k/a & (\text{at } r = a). \end{cases}$$

Gauss’s law $\Rightarrow \vec{E} = \frac{1}{4\pi\epsilon_0} \frac{Q_{\text{enc}}}{r^2} \hat{r}$. For $r < a$, $Q_{\text{enc}} = 0$, so $\boxed{\vec{E} = 0}$. For $r > b$, $Q_{\text{enc}} = 0$ (Prob. 4.14), so $\boxed{\vec{E} = 0}$.

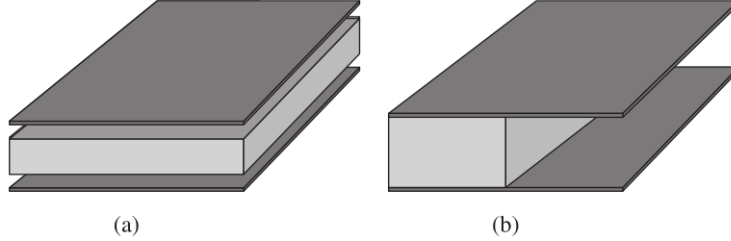
For $a < r < b$, $Q_{\text{enc}} = \left(\frac{-k}{a} \right) (4\pi a^2) + \int_a^r \left(\frac{-k}{\bar{r}^2} \right) 4\pi \bar{r}^2 d\bar{r} = -4\pi k a - 4\pi k (r - a) = -4\pi k r$; so $\boxed{\vec{E} = -(k/\epsilon_0 r) \hat{r}}$.

$$(b) \oint \vec{D} \cdot d\vec{a} = Q_{f,\text{enc}} = 0 \Rightarrow \vec{D} = 0 \text{ everywhere. } \vec{D} = \epsilon_0 \vec{E} + \vec{P} = 0 \Rightarrow \vec{E} = (-1/\epsilon_0) \vec{P}, \text{ so}$$

$$\boxed{\vec{E} = 0 \text{ (for } r < a \text{ and } r > b);}$$

$$\boxed{\vec{E} = -(k/\epsilon_0 r) \hat{r} \text{ (for } a < r < b).}$$

3. Suppose we half-fill a parallel-plate capacitor in two ways, as shown in the two figures. By what fraction is the capacitance increased when the material is distributed as shown in each case? For a given potential difference V between the plates, find E , D , and P , in each region, and the free and bound charge on all surfaces.



With no dielectric, $C_0 = A\epsilon_0/d$

In configuration (a), with $+\sigma$ on upper plate, $-\sigma$ on lower, $D = \sigma$ between the plates. $E = \sigma/\epsilon_0$ (in air) and $E = \sigma/\epsilon$ (in dielectric). So $V = \frac{\sigma}{\epsilon_0} \frac{d}{2} + \frac{\sigma}{\epsilon} \frac{d}{2} = \frac{Qd}{2\epsilon_0 A} \left(1 + \frac{\epsilon_0}{\epsilon}\right)$.

$$C_a = \frac{Q}{V} = \frac{\epsilon_0 A}{d} \left(\frac{2}{1 + \epsilon_0/\epsilon} \right) \Rightarrow \boxed{\frac{C_a}{C_0} = \frac{2\epsilon_r}{1 + \epsilon_r}}$$

In configuration (b), with potential difference V : $E = V/d$, so $\sigma = \epsilon_0 E = \epsilon_0 V/d$ (in air).

$P = \epsilon_0 \chi_e E = \epsilon_0 \chi_e V/d$ (in dielectric), so $\sigma_b = -\epsilon_0 \chi_e V/d$ (at top surface of dielectric).

$\sigma_{\text{tot}} = \epsilon_0 V/d = \sigma_f + \sigma_b = \sigma_f - \epsilon_0 \chi_e V/d$, so $\sigma_f = \epsilon_0 V(1 + \chi_e)/d = \epsilon_0 \epsilon_r V/d$ (on top plate above dielectric).

$$\Rightarrow C_b = \frac{Q}{V} = \frac{1}{V} \left(\sigma \frac{A}{2} + \sigma_f \frac{A}{2} \right) = \frac{A}{2V} \left(\epsilon_0 \frac{V}{d} + \epsilon_0 \frac{V}{d} \epsilon_r \right) = \frac{A\epsilon_0}{d} \left(\frac{1 + \epsilon_r}{2} \right). \quad \boxed{\frac{C_b}{C_0} = \frac{1 + \epsilon_r}{2}}$$

[Which is greater? $\frac{C_b}{C_0} - \frac{C_a}{C_0} = \frac{1 + \epsilon_r}{2} - \frac{2\epsilon_r}{1 + \epsilon_r} = \frac{(1 + \epsilon_r)^2 - 4\epsilon_r}{2(1 + \epsilon_r)} = \frac{1 + 2\epsilon_r + 4\epsilon_r^2 - 4\epsilon_r}{2(1 + \epsilon_r)} = \frac{(1 - \epsilon_r)^2}{2(1 + \epsilon_r)} > 0$. So $C_b > C_a$.]

4. A point charge q is placed in a medium whose permittivity ϵ changes with the distance r from q as $\epsilon = 1 + \frac{A}{r}$ where A is a constant. Show that the potential at any point is given by

$$\varphi(r) = \frac{q}{4\pi A} \ln \left(1 + \frac{A}{r} \right)$$

Q1 Considering a spherical Gaussian surface with the charge at its centre, we can write from Gauss's law

$$\oint \vec{D} \cdot d\vec{S} = q$$

Because of spherical symmetry \vec{D} is constant all over the Gaussian surface and is normal to it at every point. Therefore,

$$D \cdot 4\pi r^2 = q \Rightarrow D = \frac{q}{4\pi r^2}$$

$$\begin{aligned} \therefore E = \frac{D}{\epsilon} &= \frac{q}{4\pi \epsilon r^2} = \frac{q}{4\pi \left(1 + \frac{A}{r}\right) r^2} = \frac{q}{4\pi r(r+A)} \\ &= \frac{q}{4\pi A} \left[\frac{1}{r} - \frac{1}{r+A} \right] \end{aligned}$$

$$\begin{aligned} \therefore \phi(r) &= -\int \vec{E} \cdot d\vec{r} = -\int E dr \\ &= -\frac{q}{4\pi A} [\ln r - \ln(r+A)] + C \\ &= \frac{q}{4\pi A} \ln \left(1 + \frac{A}{r}\right) + C \end{aligned}$$

$C \rightarrow$ constant here

Since $\phi \rightarrow 0$ as $r \rightarrow \infty$, we get $C = 0$

$$\boxed{\therefore \phi(r) = \frac{q}{4\pi A} \ln \left(1 + \frac{A}{r}\right)}$$