

Cosmology 16/9-14

Error in homework set

#7: minus signs missing

hyperbolic surfaces: ~~not~~ euclidean $\Rightarrow ds^2 = dx^2 + dy^2 - dz^2$
 $x^2 + y^2 - z^2 = -R^2$

Homework: we'll add phi-dimension

Today: space-time metrics

Similar to before: (4) ^{new}

$$ds^2 = \sum_{\mu, \nu=0}^3 g_{\mu\nu} dx^\mu dy^\nu$$

math is the same, just adding a time dimension.

$$\begin{aligned} dx^0 &= c dt \\ dx^1 &= dx \\ dx^2 &= dy \\ dx^3 &= dz \end{aligned}$$

Invariant

Minkowski: flat

$$g_{\mu\nu} = \eta_{\mu\nu} = \begin{bmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

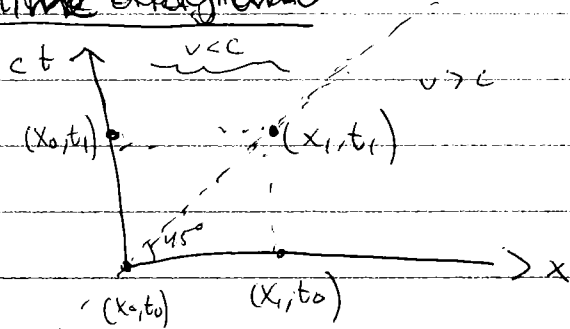
$$ds^2 = -c^2 dt^2 + dx^2 + dy^2 + dz^2$$

used in special relativity

$$\begin{aligned} ds^2 < 0 & \text{ timelike} \\ ds^2 > 0 & \text{ spacelike} \\ ds^2 = 0 & \text{ lightlike} \end{aligned}$$

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Space-time diagram



Flat static space, one dim in space

$$ds^2 = -c^2 dt^2 + dx^2$$

Integrate ds^2 twice

$$\Delta s^2 = -c^2 \Delta t^2 + \Delta x^2$$

Different examples:

Case #1: $\Delta s^2 = -c^2 (t_1 - t_0)^2 + (x_1 - x_0)^2 = \Delta x^2$ } (x_0, t_0) and (x_1, t_1)
 $\Delta s^2 > 0 \Rightarrow$ space like separation

Case #2: $\Delta s^2 = -c^2 (t_1 - t_0)^2 + (x_1 - x_0)^2 = 0$ } (x_0, t_0) and (x_1, t_1)
 $\Rightarrow \Delta s^2 = 0 \Rightarrow$ lightlike-separation

$$c = \frac{\Delta x}{\Delta t} \quad \left. \begin{array}{l} \text{slope of line} \\ \Rightarrow \text{speed of light} \end{array} \right\} \begin{array}{l} \text{sign from} \\ \text{unimodular} \\ \text{space} \end{array}$$

Null-geodesic

Case #3: $ds^2 = -c^2 (t_1 - t_0)^2 + (x_0 - x_1)^2 = -c^2 \Delta t^2$
 $(x_0, t_0) (x_1, t_1)$ $ds^2 < 0 \Rightarrow$ time-like separation

Transformations: Galilean transformations: $dx' = dx - v dt$
 $dt' = dt$

Lorentz transformations:

$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$\left. \begin{array}{l} dx' = \gamma(dx - v dt) \\ dt' = \gamma\left(dt - \frac{v}{c^2} dx\right) \end{array} \right\}$$

Minkowski-space is Lorentz invariant

$$(ds')^2 = -c^2 (dt')^2 + (dx')^2 = -c^2 \gamma^2 \left(dt - \frac{v}{c^2} dx\right)^2 + \gamma^2 (dx - v dt)^2$$

Space-time

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Minkowski: Lorentz invariant

$$\Rightarrow (ds')^2 = ds^2$$

} \Rightarrow inertial frame

start from simple spacetime metric \Rightarrow special relativity will hold.

Robinson-Walker not ~~true~~ that it is Lorentz invariant.

$$ds^2 = -c^2 dt^2 + a(t)^2 \left[\frac{dx^2}{1 - \frac{kx^2}{R^2}} + x^2 d\Omega^2 \right]$$

$$d\Omega^2 = d\theta^2 + \sin^2\theta d\phi^2$$

k : curvature constant

Null-geodesic

$ds=0$ } lightlike

$$\Rightarrow c^2 dt^2 = a(t)^2 \left[\frac{dx^2}{1 - \frac{kx^2}{R^2}} \right]$$

turning off θ, ϕ
only radial motion
for convenience: use R .

"proper time" "cosmic time"

$a(t)$ is scale factor

$a(t=0) = 1$ by convention $t=0 \Rightarrow$ today!

scale factor $\hat{=}$ very small number in early universe

Lorentz inv.?

$$(ds')^2 = -c^2 (dt')^2 + a^2 (dx')^2$$

} we now have scale factor $a(t)$

$$(ds')^2 = -c^2 t^2 \left(dt - \frac{v}{c} dx \right)^2 + a^2 t^2 (dx - v dt)^2$$

if $a=1$ today it is Lorentz inv.

$$\hat{=} t^2 (a^2 v^2 - c^2) dt + t^2 \left(a^2 - \frac{v^2}{c^2} \right) dx^2 + 2t^2 v (1 - a^2) dx dt \neq ds^2$$

- concepts from special-rel are not necessarily valid in cosmology.

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We can now calculate distances

- proper distance: eg. how far away is galaxy which is moving away from us.

$\int dr, dr, dt = 0$ assume r is fixed line
and $dt = 0$.

Robinson-Walker:

$$ds = a(t) dr$$

$$d_p(t) = a(t) \int_0^r dr = a(t) \cdot r$$

Robinson Walker $ds^2 = a(t)^2 \cdot \frac{dx^2}{1 - \frac{kx^2}{R^2}}$

Now we need curvature.

$$d_p(t) = a(t) \cdot \int_0^x \frac{dx}{\sqrt{1 - \frac{k}{R^2} x^2}}$$

proper distance

Flat $k=0 \Rightarrow d_p(r) = a(t) x$

closed $k=1 \Rightarrow d_p(r) = a(t) \int_0^x \frac{dx}{\sqrt{1 - \frac{x^2}{R^2}}} = a(t) \int_0^y \frac{R dy}{\sqrt{1 - y^2}} = a(t) R \sinh^{-1}\left(\frac{x}{R}\right)$

open $k=-1 \Rightarrow d_p(r) = a(t) R \sinh^{-1}\left(\frac{x}{R}\right)$

$$d_p(t) = a(t) \int_0^x \frac{R dy}{\sqrt{1 - y^2}} = a(t) R \sinh^{-1}\left(\frac{x}{R}\right)$$

imp. for memorization

We really need to know scale factor.

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We have universe expanding, and we know distance to something: we can find change in distance

depends on curvature

$$\frac{d(p(t))}{dt} = \frac{d}{dt} [a(t) f(x)] = \dot{a} f(x)$$

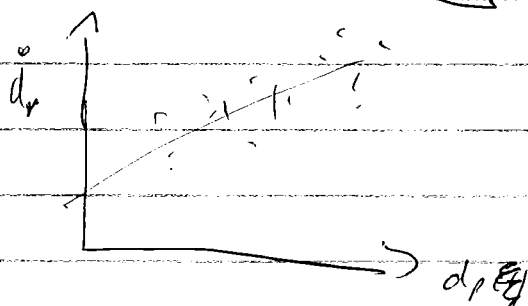
$$\left(R \sinh^{-1}\left(\frac{x}{a}\right) \right) \times R \sinh^{-1}\left(\frac{x}{a}\right) \{f(x)\}$$

$$f(x) = \frac{dp(t)}{a(t)}$$

$$\Rightarrow \boxed{\frac{dp(t)}{dt} = \frac{\dot{a}}{a} dp}$$

Expansion velocity of Universe.

We don't know what \dot{a} is yet, but we know proportionality!



\Rightarrow Hubble's plot:

$$\boxed{H_0 = \frac{\dot{a}}{a}} \quad \text{if } t = t_0 \text{ today}$$

Local Universe H_0 is a constant.

- slope of Hubble flow

Recession speed \uparrow as $d_p \uparrow$

\hookrightarrow we'll get to point where

recession speed is more than speed of light

$$dp^o = \frac{\dot{a}}{a} dp = v_p$$

if

$$v_p(t_0) = c$$

then

$$\boxed{dp(t_0) = \frac{c}{H_0}}$$

it's okay that points were away faster

Hubble distance

than speed of light. Robinson-Walker... not current event.

$$H_0 = 70 \pm 7 \text{ km/s/Mpc}$$

$$d_H(t_0) = \frac{3 \cdot 10^8 \text{ km/s}}{70 \cdot 1000 \frac{\text{km/s}}{\text{Mpc}}} = 4300 \text{ Mpc} = 14 \text{ Gly}$$