

SULFUR-BEARING MOLECULES IN MASSIVE STAR-FORMING REGIONS: OBSERVATIONS OF OCS, CS, H₂S, AND SO

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ABSTRACT

We studied the sulfur chemistry of massive star-forming regions through single-dish submillimeter spectroscopy. OCS, O¹³CS, ¹³CS, H₂S, and SO transitions were observed toward a sample of massive star-forming regions with embedded UCH II or CH II regions. These sources could be divided into H II-hot core and H II-only sources based on their CH₃CN emission. Our results show that the OCS line of thirteen sources is optically thick, with optical depth ranging from 5 to 16. Column densities of these molecules were computed under LTE conditions. CS column densities were also derived using its optically thin isotopologue ¹³CS. H₂S is likely to be the most abundant gas-phase sulfuretted molecule in hot massive cores. Both the column density and abundance of sulfur-bearing molecules decrease significantly from H II-hot core to H II-only sources. Ages derived from hot core models appear to be consistent with star formation theories, suggesting that abundance ratios of [CS]/[SO], [SO]/[OCS], and [OCS]/[CS] could be used as chemical clocks in massive star-forming regions.

Key words: ISM: abundances – ISM: clouds – ISM: molecules – molecular processes – radio lines: ISM

1. INTRODUCTION

Sulfur-bearing molecules are not only excellent tracers of early protostellar evolution (van der Tak et al. 2003) but are also of specific interest because of their rapid evolution in warm gas (Charnley 1997; Hatchell et al. 1998a, 1998b). Unlike nitrogen chemistry or oxygen chemistry, most of the molecules involved are easy to detect at millimeter and submillimeter wavelengths and column densities can be calculated easily, thus sulfur is a potential good clock on timescales relevant to the embedded phase of star formation.

Chemical evolution of sulfur-bearing molecules have been well studied in low-mass star-forming regions, e.g., Buckle & Fuller (2003) studied the chemical evolution of H₂S, SO, and SO₂ in low-mass star-forming regions. They demonstrated that the chemical evolution of sulfur-bearing species was a potentially valuable probe of chemical timescales in low-mass star-forming regions. However, this issue has been less developed in the context of high-mass star-forming regions. Hatchell et al. (1998a, 1998b) studied the sulfur chemistry of eight sources associated with UCH II regions. The observed abundance ratios vary little from source to source. Van der Tak et al. (2003) studied the sulfur chemistry in nine high-mass protostars and found that high-energy transitions probing the inner parts of the protostars where the temperature exceeds 100 K were mandatory for using S-bearing molecules as chemical clocks. They proposed that OCS was the main carrier of sulfur on grains, based on the high excitation temperature of the molecule and its high abundance in the protostars.

Since hydrogenation is expected to be the most efficient process on grains, H₂S could be formed on interstellar grains (van Dishoeck & Blake 1998). It is now generally accepted that during the cold collapse phase, sulfur atoms freeze out onto grains and remain there in the form of H₂S until core heating begins. At this point, H₂S is evaporated from the grains and rapidly undergoes reactions that drive the production of SO and

SO₂ (e.g., Charnley 1997; Wakelam et al. 2004, 2011). The initial destruction of H₂S by H₃O⁺ is more efficient in the high-mass objects than in low-mass protostars, as water is more abundant there (van der Tak et al. 2006). Viti et al. (2004) extended chemical models of high-mass star-forming cores by including the experimental results on desorption from grains in the evaporation of icy mantles formed in star-forming regions and developed a time-dependent gas-grain model. They found that distinct chemical events occur at specific grain temperatures. They showed that, provided there is a monotonic increase in the temperature of the gas and dust surrounding the protostar, the changes in the chemical evolution of each species due to differential desorption are important. They proposed that the abundance ratio of sulfur-bearing species, like [CS]/[H₂S], [OCS]/[H₂S], [CS]/[SO], and [SO]/[OCS], can be used to investigate the evolution of the early phases of massive star formation, and that sulfur-bearing species are good “chemical clocks.” Esplugues et al. (2014) modeled the sulfur chemistry evolution of the hot core and the plateau in Orion KL using the chemical models in Viti et al. (2004), with two different phases in calculations. Phase I starts with the collapsing cloud and the depletion of atoms and molecules onto grain surfaces, while Phase II starts when a central protostar is formed and the evaporation from grains takes place. They reproduced the observations of sulfur-bearing molecules (SO, SO₂, CS, OCS, H₂S, and H₂CS) in the hot core and plateau.

Herpin et al. (2009) observed sulfur-bearing species of two mid-infrared-quiet and two brighter massive cores, with one bright massive core associated with hypercompact H II (HCH II) regions. They found that the SO and SO₂ relative abundances increased with time, while that of CS and OCS decreased. This trend seems to be in contradiction with the chemical model of Viti et al. (2004), and more observations are needed to investigate whether this is biased by the small observing sample.

Table 1
Observed Transitions

Transition	Frequency (GHz)	E_u (K)	Eins. A. (s^{-1})
H ₂ S 2 ₂₀ –2 ₁₁	216.710437	84.0	4.87×10^{-5}
SO 5 ₆ –4 ₅	219.949433	35.0	1.36×10^{-4}
OCS 19–18	231.060991	110.2	3.58×10^{-5}
O ¹³ CS 19–18	230.317500	109.9	3.54×10^{-5}
¹³ CS 5–4	231.220768	33.1	2.51×10^{-4}
C ¹⁸ O 2–1	219.560354	15.8	6.01×10^{-3}

Considering that the systematic understanding of the chemistry of sulfur-bearing molecules in massive star-forming regions is far from complete, we carried out observations of OCS, O¹³CS, ¹³CS, H₂S, and SO toward a sample of 36 massive star-forming regions with embedded H II regions and studied the evolution of molecular column densities and abundance ratios. These results are compared with chemical models of the hot core. In this paper we first introduce the observations and data reductions in Section 2. In Section 3, we present the observing results. The analysis and discussion of the results are presented in Section 4, followed by a summary in Section 5.

2. OBSERVATIONS AND DATA REDUCTIONS

The observations were conducted in 2013 July at the 10.4 m telescope of the Caltech Submillimeter Observatory (CSO) using the 230 receiver, which is a double-side band (DSB) receiver. The data were obtained in position switching mode with an offset of 30' in azimuth. The Fast Fourier Transform Spectrometer provides a 4 GHz window of DSB coverage with a velocity resolution of 0.35 km s⁻¹ at 230 GHz. The total integration times ranged from 3 to 20 minutes, while system temperatures ranged from 250 to 400 K. OCS 19–18 (231.060991 GHz), O¹³CS 19–18 (230.317500 GHz), ¹³CS 5–4 (231.220768 GHz), H₂S 2₂₀–2_{1,1} (216.710437 GHz), and SO 5₆–4₅ (219.949433 GHz) were observed simultaneously within the receiving band, with OCS, O¹³CS, and ¹³CS transitions observed in USB, and H₂S and SO transitions observed in LSB. The C¹⁸O 2–1 transition was obtained simultaneously, which was used to derive the column density of H₂ and molecular abundances. Table 1 gives details about the observed transitions.

Part of our sample is selected from Shirley et al.'s (2003) massive star-forming cores associated with water masers. The observing center was the water maser position from the catalog of Cesaroni et al. (1988). Other sources are chosen from Zhu et al. (2008), Wood & Churchwell (1989), Kurtz et al. (1994), and Beck et al. (1998). Table 1 lists detailed information of these 36 sources. The classifications of the associated H II region and references are also listed. Most of these sources are associated with UCH II regions, while W75OH is associated with the HCH II region (Li et al. 2012). G10.6–0.4, W33IRS1, W33IRS2, W33IRS3, and W51M are associated with CH II regions (Wood & Churchwell 1989; Beck et al. 1998). S76E is associated with H II regions (Li et al. 2012).

Data reduction was carried out using the GILDAS/CLASS package⁶ developed at IRAM. Linear baselines were subtracted and calibrated onto T_{MB} using a beam efficiency of 0.69. The

final spectra have rms noise levels (in T_{MB}) ranging from 0.03 to 0.14 K at the velocity resolution of 0.35 km s⁻¹.

3. RESULTS

The whole 8 GHz spectra data indicate that some sources show many high-excitation lines from complex molecules like CH₃CN, whereas others show mainly low-excitation lines like ¹³CO and C¹⁸O (J. Li et al. 2015, in preparation). High-excitation energy lines require high temperatures to excite, so the no-detections mean that there is no evidence for hot cores in line-poor sources. Since massive star formation takes place on rather short timescales and in clustered environments, the transitions from one stage into the next are smooth and not always clearly distinct. There might be overlaps among the hot massive core and UCH II region stages. Due to the large beam, regions with different evolutionary stages in their proximity can also be observed as one object showing characteristics of different stages. This situation was noted by Hatchell et al. (1998a), and they divided their UCH II sample into line-rich and line-poor sources. Thus we also divided our sources into two groups. Sources with detectable CH₃CN emission (220 GHz) above the 2 σ (~ 0.1 K) level are labeled as H II-hot core (H II-HC), while sources without detectable CH₃CN emission above the 2 σ (~ 0.1 K) level are labeled as H II-only. The brightness of CH₃CN emission does not correlate with distance, implying intrinsic differences in the source properties rather than simply beam-dilution effects (Hatchell et al. 1998a). For the H II-HC sources, though they are already surrounded by an ionized medium, they still show rich hot core chemistry. On the contrary, the H II-only sources without a hot core signature could better represent chemistry in the UCH II stage. Thus the H II-HC sources are regarded to be younger than the H II-only sources because they contain younger components (Gerner et al. 2014). There are 25 H II-HC sources and 11 H II-only sources in our sample. The classification information is also listed in Table 2.

Spectra from each of the molecules toward all the sources are given in Figure 1. SO emission is stronger than other lines in all the sources except for W31, G10.6–0.4 and W42, in which it is weaker than H₂S emission. OCS emission was detected in 26 sources, most of which are H II-HC sources. O¹³CS emission was detected in 15 sources, all of which have been detected in OCS except for G35.20–1.74. ¹³CS was detected in 33 sources, while H₂S was detected in 30 sources, including all the H II-HC sources and some H II-only sources. SO was detected in all the sources, with integrated intensities ranging from 0.82 (G78.44 +2.66) to 61.97 K km s⁻¹ (W51M). We fitted the detected lines with Gaussian profiles. Most lines could be satisfactorily fitted with a single Gaussian except for SO, which is consistent with van der Tak et al. (2003). The Gaussian fitting results, including peak temperature (T_{mb}), integrated line intensity, LSR velocity, and FWHM linewidth of OCS, O¹³CS, ¹³CS, H₂S, and C¹⁸O are presented in Tables 3–5. For SO, which could not be satisfactorily fitted with a single Gaussian, we present the integrated intensity, LSR velocity, and linewidth in Table 4.

4. ANALYSIS AND DISCUSSIONS

4.1. OCS Optical Depth

The OCS/O¹³CS intensity ratio is significantly lower than the expected abundance ratio of about 45, suggesting that OCS is

⁶ <http://www.iram.fr/IRAMFR/GILDAS>.

Table 2
Source Sample

Source Name	R.A. (J2000) (h m s)	Decl. (deg)	Association	Reference	Chemically
W3(OH)	02:27:04.69	61:52:25.5	UCH II	Li2012	H II-HC
RCW142	17:50:15.13	-28:54:31.5	UCH II	Shirley2003	H II-HC
G5.89-0.39	18:00:30.39	-24:04:00.1	UCH II	Zhu2008	H II-HC
M8E	18:04:53.25	-24:26:42.4	UCH II	Shirley2003	H II-only
G8.67-0.36	18:06:19.03	-21:37:32.3	UCH II	Li2012	H II-HC
W31	18:08:38.32	-19:51:49.7	UCH II	Shirley2003	H II-HC
G10.6-0.4	18:10:28.70	-19:55:48.6	CH II	WC1989	H II-HC
G11.94-0.62	18:14:01.10	-18:53:24.2	UCH II	Zhu2008	H II-HC
W33cont	18:14:13.67	-17:55:25.2	UCH II	Shirley2003	H II-HC
W33irs1	18:14:15.00	-17:55:54.9	CH II	Beck1998	H II-HC
W33irs2	18:14:14.49	-17:55:14.9	CH II	Beck1998	H II-HC
W33irs3	18:14:13.30	-17:55:40.0	CH II	Beck1998	H II-HC
G14.33-0.64	18:18:54.71	-16:47:49.7	UCH II	Shirley2003	H II-HC
W42	18:38:12.46	-07:12:10.1	UCH II	Li2012	H II-HC
G29.96-0.02	18:46:03.90	-02:39:22.0	UCH II	Zhu2008	H II-HC
G30.54+0.02	18:46:59.5	-02:07:26	UCH II	Zhu2008	H II-only
G33.92+0.11	18:52:50.00	00:55:28.9	UCH II	Zhu2008	H II-only
G34.26+0.15	18:53:18.53	01:14:57.9	UCH II	WC1989	H II-HC
S76E	18:56:10.43	07:53:14.1	H II	Li2012	H II-HC
G35.20-0.74	18:58:12.73	01:40:36.5	UCH II	Shirley2003	H II-HC
G35.20-1.74	19:01:47.00	01:13:07.9	UCH II	Kurtz1994	H II-only
G37.543-0.11	19:00:15.70	04:03:12.9	UCH II	WC1989	H II-only
G43.89-0.78	19:14:26.2	09:22:34	UCH II	Zhu2008	H II-only
G45.45+0.06	19:14:21.30	11:09:12.9	UCH II	Zhu2008	H II-only
W51M	19:23:43.86	14:30:29.4	CH II	Li2012	H II-HC
W51D	19:23:39.90	14:31:06.0	UCH II	Zhu2008	H II-HC
G61.48+0.09	19:46:47.33	25:12:45.6	UCH II	Zhu2008	H II-only
K3-50 A	20:01:45.59	33:32:42.0	UCH II	Zhu2008	H II-only
G78.44+2.66	20:19:29.21	40:56:36.6	UCH II	Kurtz1994	H II-only
W75N	20:38:36.93	42:37:37.4	UCH II	Li2012	H II-HC
DR 21S	20:39:00.80	42:19:29.7	UCH II	Li2012	H II-HC
W75(OH)	20:39:01.00	42:22:49.8	HCH II	Li2012	H II-HC
G81.7+0.5	20:39:00.90	42:19:44.0	UCH II	Kurtz1994	H II-only
CepA	22:56:18.13	62:01:46.3	UCH II	Li2012	H II-HC
NGC 7538	23:13:44.85	61:26:50.6	UCH II	Shirley2003	H II-HC
NGC 7538A	23:13:45.59	61:28:18.0	UCH II	Zhu2008	H II-HC

References. Beck1998: Beck et al. (1998), Kurtz1994: Kurtz et al. (1994), Li2012: Li et al. (2012), Shirley2003: Shirley et al. (2003), WC1989: Wood & Churchwell (1989), Zhu2008: Zhu et al. (2008).

optically thick. This is consistent with SMA observations of G10.6-0.4 (Liu et al. 2010) and PdBI observations of the intermediate mass protostar NGC 7129 FIRS 2 (Fuente et al. 2014). Fuente et al. (2014) estimated the opacity of the OCS 19-18 line to be about 12 in NGC 7129 FIRS 2, while Schilke et al. (1997) calculated an optical depth of about 3.5 for Orion KL. We estimated the OCS optical depth by assuming that the excitation temperature is the same for both isotopologues. Tercero et al. (2010) found an average value of $^{12}\text{C}/^{13}\text{C} = 45$ in their study of sulfur carbon molecules of Orion KL, and thus we adopted $^{12}\text{C}/^{13}\text{C}$ of 45 for all these sources. With these assumptions we can write the following relation (Myers et al. 1983):

$$\frac{T_{\text{O}^{13}\text{CS}}}{T_{\text{OCS}}} = \frac{1 - \exp(-\tau_{\text{OCS}}/45)}{1 - \exp(-\tau_{\text{OCS}})}. \quad (1)$$

The computed τ_{OCS} is tabulated in Table 3. OCS 19-18 emissions are optically thick in all sources detected in O^{13}CS , with the optical depth ranging from 5 to 16. It appears that OCS lines are saturated in hot cores. We found that O^{13}CS is

stronger than OCS in DR 21S and G35.20-1.74, which is possibly caused by a large temperature gradient. Observations with higher angular resolution are needed to investigate this issue.

4.2. Column Density and Abundance

Molecular column densities have been estimated for all observed species. We assumed that the emission is extended and uniformly distributed within the beam. Therefore, the beam filling factor is ~ 1 . For all species, we assume local thermodynamic equilibrium (LTE) conditions and optically thin emission. We used the following formula to estimate molecular column densities (Vasyunina et al. 2014):

$$N_{\text{tot}} = \frac{8\pi}{\lambda^3 A} \frac{1}{J_{\mu}(T_{\text{ex}}) - J_{\mu}(T_{\text{bg}})} \frac{1}{1 - \exp(-h\nu/kT_{\text{ex}})} \times \frac{Q_{\text{rot}}}{g_u \exp(-E_l/kT_{\text{ex}})} \int T_{\text{mb}} dv, \quad (2)$$

where λ is the rest wavelength of the transition, A is the

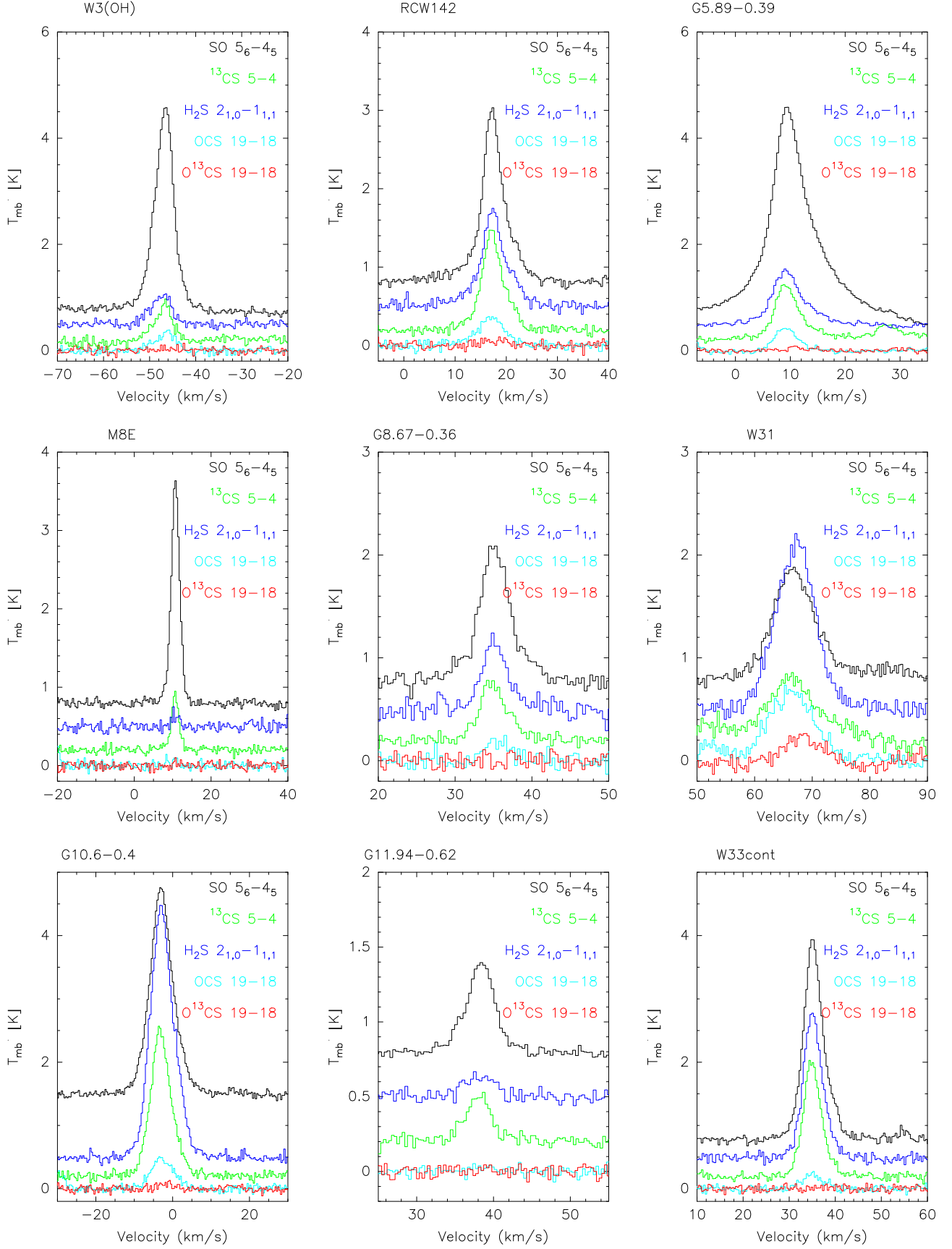


Figure 1. Spectra of ^{13}CS , OCS , O^{13}CS , H_2S , and SO for all observed sources. The line identifications are labeled in the upper right of each figure.

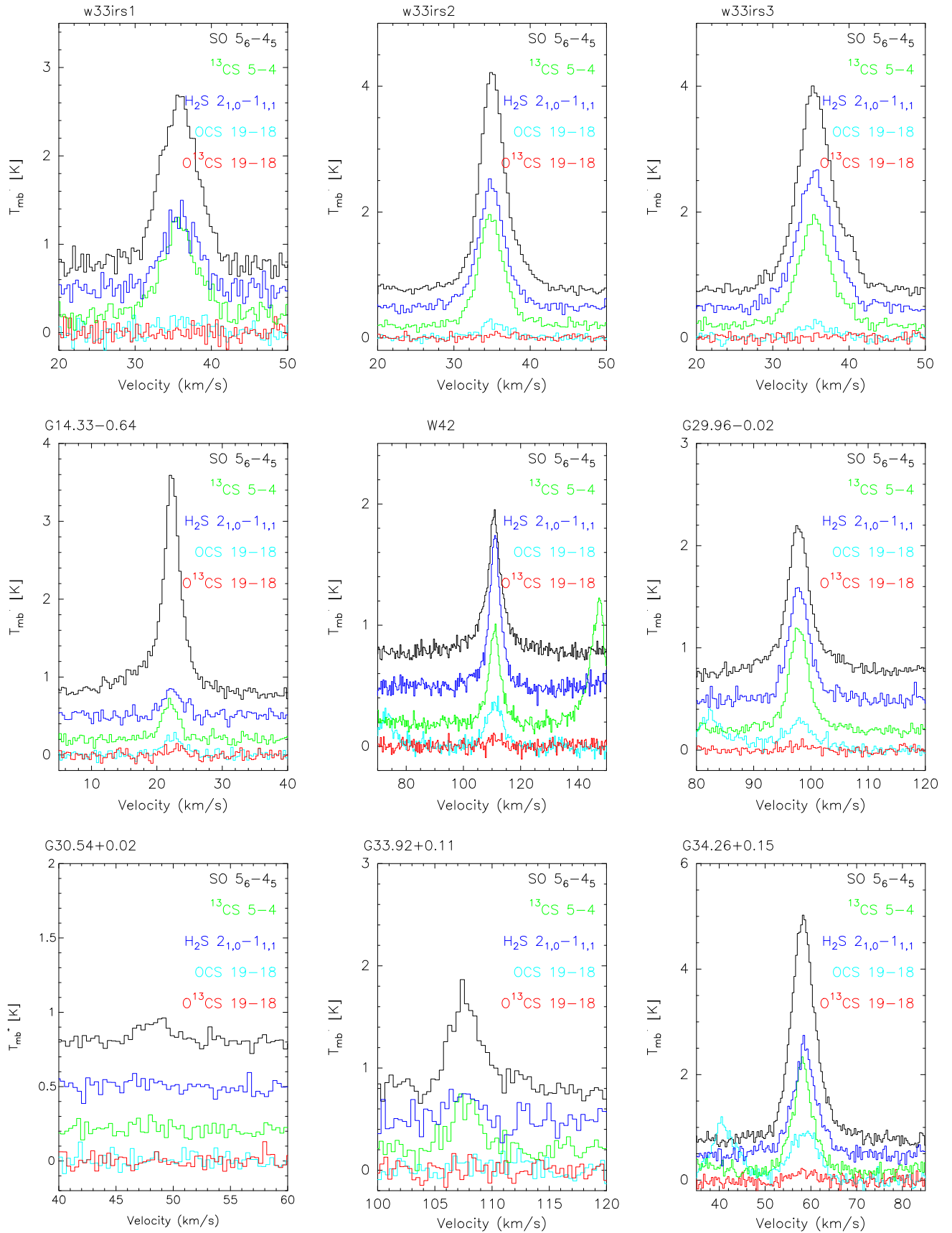


Figure 1. (Continued.)

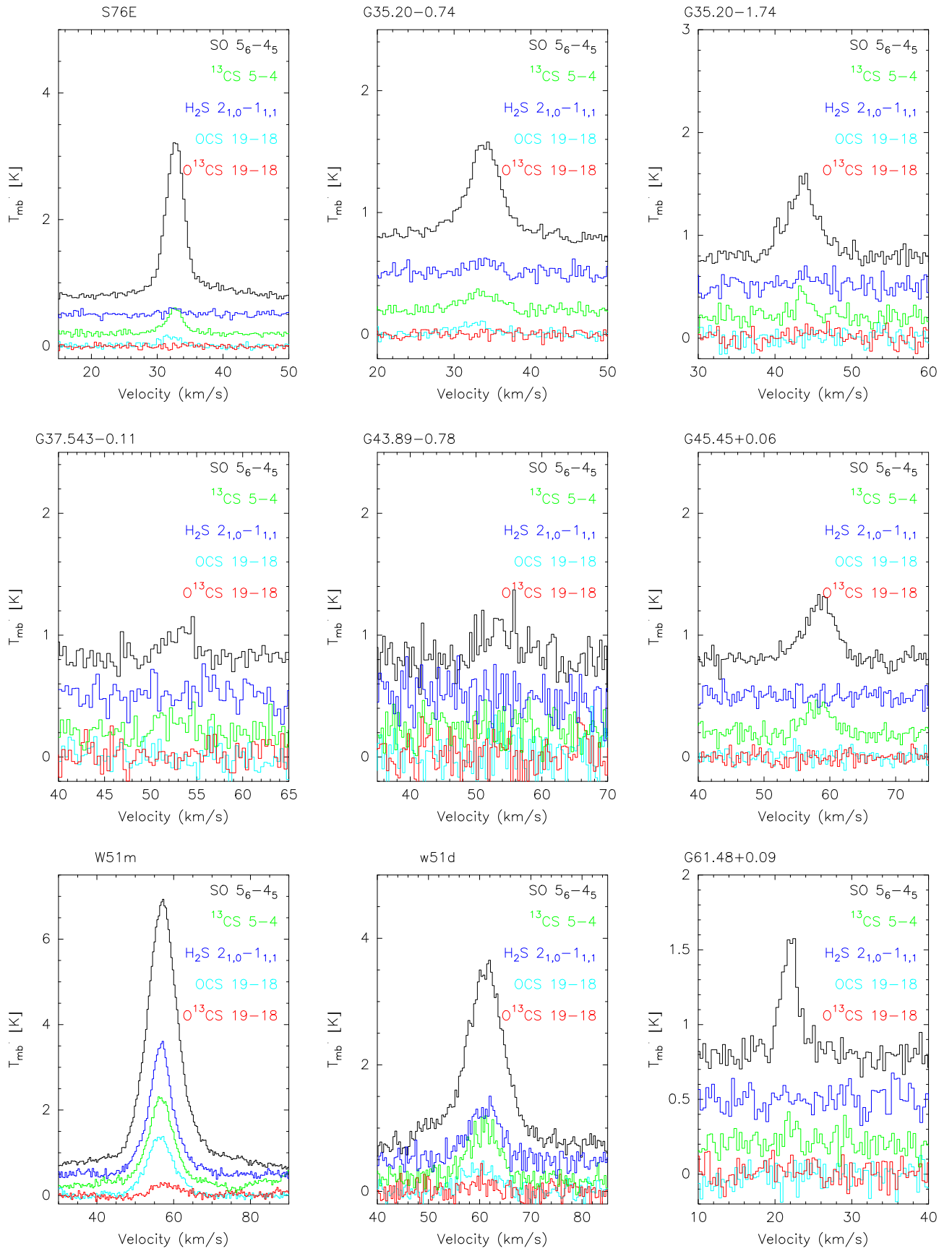


Figure 1. (Continued.)

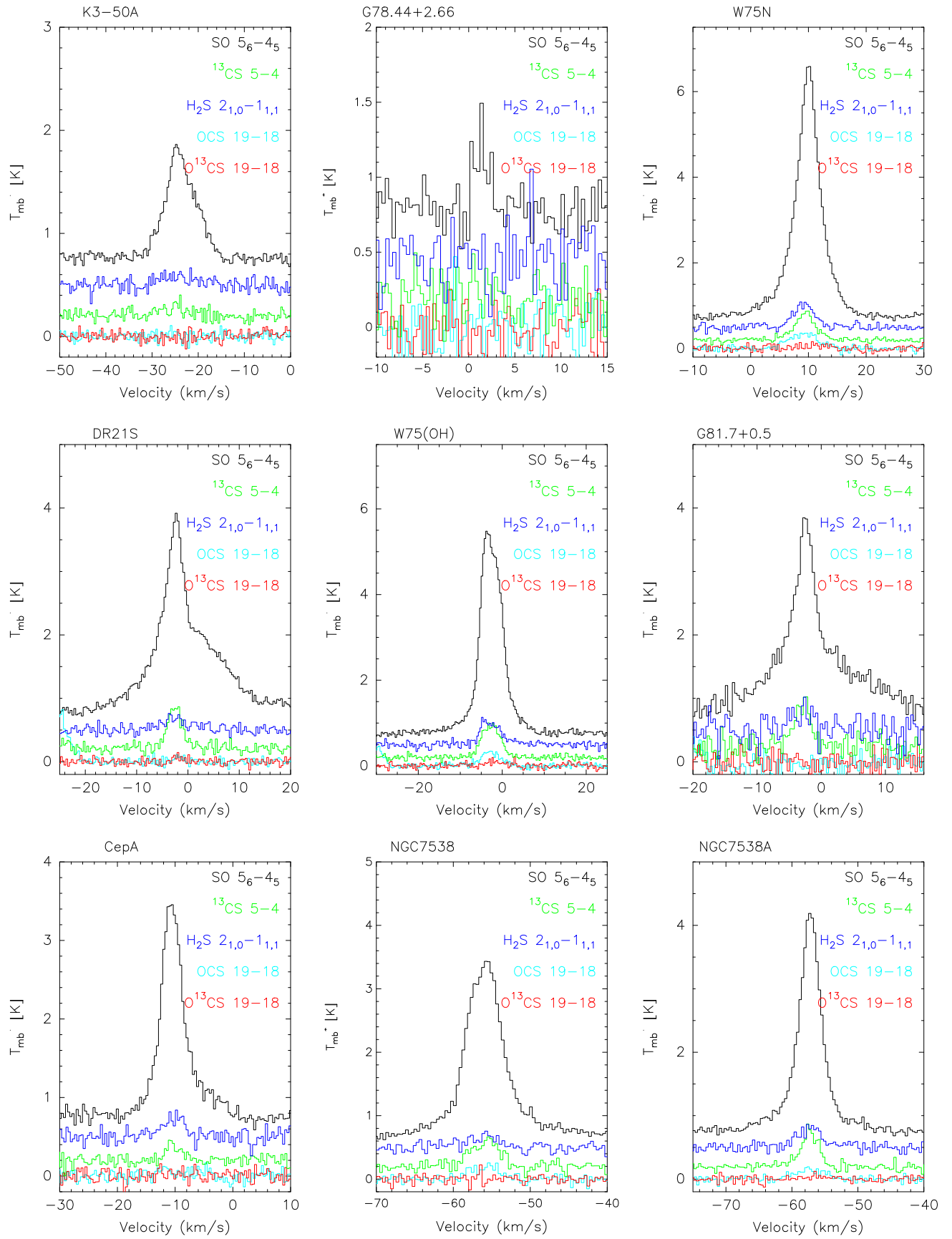


Figure 1. (Continued.)

Table 3
Results of OCS 19–18, O¹³CS 19–18 Transitions and Opacity of OCS

Source Name	OCS 19–18				O ¹³ CS 19–18				Opacity
	T_{mb} (K)	$\int T_{\text{mb}} d\nu$ (K km s ⁻¹)	V_{LSR} (km s ⁻¹)	FWHM (km s ⁻¹)	T_{mb} (K)	$\int T_{\text{mb}} d\nu$ (K km s ⁻¹)	V_{LSR} (km s ⁻¹)	FWHM (km s ⁻¹)	
W3(OH)	0.33(.04)	1.78(0.10)	-46.47(.14)	5.09(.33)	0.03(.04)	0.26(.13)	-45.43(1.97)	8.00(5.86)	7.2
RCW142	0.35(.04)	2.33(0.09)	17.02(.12)	6.34(.30)	0.06(.03)	0.41(0.09)	17.79(.64)	6.35(1.34)	8.7
G5.89–0.39	0.42(.06)	2.14(0.10)	9.23(.11)	4.76(.26)	0.04(.01)	0.29(0.04)	11.47(.42)	5.49(1.45)	6.6
M8E	0.06(.03)	0.14(0.06)	9.94(.50)	3.36(.87)	<0.04	<0.09
G8.67–0.36	0.19(.06)	0.78(0.10)	35.52(.24)	3.79(.53)	<0.06	<0.25
W31	0.67(.14)	6.68(0.42)	66.33(.28)	9.49(.71)	0.23(.06)	1.68(0.14)	68.32(.29)	6.75(.61)	13.1
G10.6–0.4	0.46(.09)	3.48(0.19)	-2.70(.19)	6.96(.45)	0.09(.03)	0.39(0.06)	-1.77(.36)	4.46(.79)	5.3
G11.94–0.62	0.03(.03)	0.10(0.04)	38.64(.56)	2.98(1.12)	<0.03	<0.10
W33cont	0.22(.04)	1.23(0.07)	35.02(.16)	5.32(.38)	<0.04	<0.22
W33irs1	0.13(.06)	0.91(0.17)	36.61(.62)	6.36(1.32)	<0.09	<0.63
W33irs2	0.23(.04)	1.03(0.07)	35.20(.14)	4.33(.39)	<0.03	<0.13
W33irs3	0.20(.04)	1.09(0.10)	35.70(.21)	4.98(.53)	<0.06	<0.33
G14.33–0.64	0.22(.04)	0.88(0.07)	22.52(.15)	3.78(.37)	0.10(.04)	0.12(0.07)	23.13(.40)	2.60(.99)	6.6
W42	0.35(.07)	2.49(0.17)	110.9(.23)	6.78(.56)	0.07(.03)	0.36(0.07)	11.3(.46)	4.80(1.03)	7
G29.96–0.02	0.23(.03)	2.10(0.09)	97.76(.15)	8.28(.45)	0.04(.03)	0.23(0.06)	98.17(.64)	4.81(1.42)	5.2
G30.54+0.02
G33.92+0.11	0.09(.07)	0.48(0.13)	110.6(.84)	5.57(1.25)	<0.07	<0.37
G34.26+0.15	0.90(.07)	7.03(0.75)	58.49(.39)	7.35(.94)	0.16(.38)	0.90(0.17)	58.53(.56)	5.50(1.06)	6.1
S76E	0.14(.03)	0.57(0.06)	32.01(.16)	3.67(.42)	<0.03	<0.43
G35.20–0.74	0.09(.03)	0.48(0.06)	32.55(.28)	5.33(.77)	<0.03	<0.16
G35.20–1.74	0.07(.06)	0.46(0.13)	45.49(.84)	5.61(.80)	...
G37.543–0.11
G43.89–0.78
G45.45+0.06
W51M	1.32(.06)	11.94(0.68)	56.54(.23)	8.49(.58)	0.26(.07)	2.22(0.17)	58.14(.32)	7.91(.68)	9.3
W51D	0.35(.09)	2.90(0.26)	60.07(.35)	7.71(.75)	<0.16(0.07)	<0.79(.19)	60.19(0.54)	4.60(1.26)	...
G61.48+0.09
K3–50 A
G78.44+2.66
W75N	0.35(.04)	2.19(0.10)	9.11(.14)	5.94(.32)	0.07(.06)	0.48(0.13)	10.98(.89)	5.97(2.08)	11.1
DR 21S	0.09(.04)	0.17(0.06)	-1.78(.29)	1.79(.58)	0.09(.04)	0.19(0.06)	-1.99(.37)	2.10(.87)	...
W75(OH)	0.32(.04)	1.86(0.10)	-2.82(.13)	5.43(.33)	0.10(.04)	0.57(0.14)	-1.61(.68)	5.39(1.83)	16.5
G81.7+0.5
CepA
NGC 7538	0.19(.04)	0.96(0.09)	-55.88(.24)	4.63(.49)	<0.04	<0.20
NGC 7538A	0.14(.04)	0.83(0.07)	-57.31(.25)	5.22(.52)	0.04(.03)	0.22(0.09)	-56.08(.92)	4.76(1.91)	13.9

Einstein coefficient, g_u is the upper state degeneracy, $J_u(T_{\text{ex}})$ and $J_u(T_{\text{bg}})$ are the values of the Planck function at excitation and background temperatures, respectively, Q_{rot} is the partition function, and E_l is the energy of the lower level (Lee et al. 2009). For g_u , A , and E_l , we used values from The Cologne Database for Molecular Spectroscopy (CDMS; Müller et al. 2001, 2005) for most of the species. We assumed a uniform excitation temperature T_{ex} of 50 K (van der Tak et al. 2003) and calculated the partition function Q_{rot} for every source by interpolating data from the CDMS for the assumed T_{ex} .

We did not detect emission from ³³SO (217.831762 GHz) above 2σ (~ 0.1 K) in these sources, suggesting that ³²SO emission is almost optically thin. Previous observations also indicate that ³²SO was at most moderately optically thick (e.g., Buckle & Fuller 2003), thus we derived SO column densities with the assumption of optical thinness. Crockett et al. (2014) observed multi- J transitions of H₂S and its isotopologues at frequencies ranging from 480 to 1907 GHz. They found that H₂S is optically thick, with optical depth ranging from 2 to 40. However, given that we did not observe isotopologues of H₂S,

we assume that it is optically thin for simplicity. The column density and abundance of H₂S might be underestimated because of this, but the comparison between H II–HC and H II–only groups will not be affected too much. C¹⁸O column densities were also derived with the assumption optical thinness. Using the abundance of $N(\text{H}_2)/N(\text{C}^{18}\text{O}) = 7 \times 10^6$ (Castets & Langer 1995), the $N(\text{H}_2)$ based on C¹⁸O data are derived. To derive the CS column density, we used the optically thin emission of ¹³CS by assuming ¹²C/¹³C = 45 (Tercero et al. 2010). For sources detected in O¹³CS emission, the OCS column density was also derived from its less abundant isotopologue O¹³CS, again assuming a ¹²C/¹³C ratio of 45. For sources detected in OCS but not detected in O¹³CS, the OCS column densities were derived from OCS emission by assuming an optical depth of 5. Three of our sources, including G5.89–0.39, G29.96–0.02, and G34.26+0.15, were also observed by Hatchell et al. (1998a, 1998b). It is noted that the column densities derived here are in agreement with their results. Column density for H₂, CS, OCS, H₂S, and SO molecules, as well as abundance for H₂, CS, OCS, H₂S, and SO molecules are listed in Table 6. Median values of the

Table 4
Results of ^{13}CS 5–4, H_2S 2_{2,0}–2_{1,1}, and SO 5₆–4₅ Transitions

Source Name	^{13}CS 5–4				H_2S 2 _{2,0} –2 _{1,1}				SO 5 ₆ –4 ₅		
	T_{mb} (K)	$\int T_{\text{mb}} d\nu$ (K km s ^{–1})	V_{LSR} (km s ^{–1})	FWHM (km s ^{–1})	T_{mb} (K)	$\int T_{\text{mb}} d\nu$ (K km s ^{–1})	V_{LSR} (km s ^{–1})	FWHM (km s ^{–1})	$\int T_{\text{mb}} d\nu$ (K km s ^{–1})	V_{LSR} (km s ^{–1})	FWHM (km s ^{–1})
W3(OH)	0.64(.06)	3.00(.13)	–46.92(.09)	4.49(2.12)	0.54(.04)	2.94(0.10)	–47.45(.09)	5.18(.20)	20.08	–46.93	6.06
RCW142	1.16(.04)	6.41(.07)	17.29(.03)	5.18(.08)	1.09(.06)	7.09(0.13)	17.49(.05)	6.12(.51)	13.14	17.08	9.31
G5.89–0.39	0.97(.06)	5.84(.28)	9.40(.12)	5.62(0.34)	0.94(.03)	6.62(0.09)	9.43(.04)	6.61(.11)	47.63	10.21	19.44
M8E	0.70(.03)	1.70(.04)	10.65(.03)	2.27(.08)	0.14(.04)	0.36(0.06)	10.77(.18)	2.36(.34)	7.89	10.67	2.92
G8.67–0.36	0.54(.03)	2.45(.07)	34.96(.07)	4.33(0.17)	0.65(.04)	2.67(0.14)	35.14(.10)	3.87(.24)	7.18	35.11	6.89
W31	0.54(.06)	6.88(1.00)	66.35(.76)	11.964(2.41)	1.59(.06)	13.67(0.61)	67.13(.19)	8.03(.44)	10.46	67.35	10.38
G10.6–0.4	2.20(.04)	15.01(.12)	–3.09(.02)	6.39(.06)	3.80(.04)	28.94(0.13)	–2.77(.02)	7.16(.04)	24.32	–2.77	7.74
G11.94–0.62	0.30(.03)	1.26(.06)	38.03(.08)	3.96(.21)	0.13(.03)	0.68(0.06)	38.17(.19)	4.66(.41)	2.83	38.23	5.52
W33cont	1.75(.04)	8.70(.09)	35.02(.02)	4.66(.05)	2.20(.06)	11.51(0.12)	35.15(.02)	4.89(.06)	17.02	35.42	6.02
W33irs1	0.99(.07)	5.77(.16)	35.78(.07)	5.48(0.19)	0.81(.09)	4.55(0.19)	35.92(.11)	5.27(.25)	11.02	35.82	5.43
W33irs2	1.70(.04)	7.09(.09)	34.86(.02)	3.93(.06)	1.87(.06)	8.16(0.10)	34.91(.03)	4.09(.06)	14.97	35.22	5.20
W33irs3	1.62(.06)	8.84(.12)	35.46(.03)	5.11(.08)	2.06(.07)	12.10(0.14)	35.56(.03)	5.54(.08)	18.39	35.59	5.93
G14.33–0.64	0.39(.04)	1.83(.22)	22.70(.37)	4.40(2.77)	0.33(.04)	1.26(0.09)	22.42(.13)	3.61(.30)	11.95	21.46	7.60
W42	0.71(.04)	3.91(.51)	110.9(.31)	5.12(.84)	1.14(.06)	6.04(0.12)	110.0(.04)	4.94(.12)	7.48	110.53	8.52
G29.96–0.02	0.97(.03)	4.26(.61)	97.92(.28)	4.14(.71)	1.06(.06)	4.97(0.12)	97.92(.05)	4.41(.12)	8.64	97.74	8.90
G30.54+0.02	0.49	48.93	4.17
G33.92+0.11	0.46(.07)	1.41(.12)	107.7(.11)	2.84(.28)	0.26(.06)	0.75(0.14)	107.4(.29)	2.73(.54)	3.03	107.93	3.28
G34.26+0.15	1.90(.09)	10.41(.19)	58.20(.04)	5.14(.12)	1.90(.09)	12.35(0.22)	58.64(.05)	6.10(.13)	27.65	58.67	7.40
S76E	0.38(.03)	1.23(.06)	32.66(.07)	3.06(.18)	0.10(.03)	0.22(0.06)	32.68(.22)	2.00(.56)	10.07	33.26	6.68
G35.20–0.74	0.13(.03)	0.78(.06)	33.45(.22)	5.36(.48)	0.10(.04)	0.43(0.07)	33.74(.36)	3.93(.71)	4.34	33.55	6.34
G35.20–1.74	0.26(.06)	0.54(.09)	43.69(.14)	2.00(.42)	0.10(.04)	0.49(0.19)	44.22(.74)	4.61(2.63)	3.58	43.64	4.85
G37.543–0.11	0.11(.05)	0.45(.11)	52.62(.52)	3.86(.84)	0.86	53.20	3.57
G43.89–0.78	1.93	54.15	6.15
G45.45+0.06	0.22(.04)	1.07(.10)	58.29(.21)	4.59(.45)	2.76	58.12	5.57
W51M	1.99(.07)	17.83 (3.16)	56.71(.71)	8.42(1.87)	2.84(.07)	23.30(0.20)	56.70(.03)	7.70(.08)	61.97	57.54	11.24
W51D	0.80(.10)	5.94(.25)	60.74(.15)	7.01(.34)	0.81(.12)	6.26(0.29)	60.96(.16)	7.20(.38)	28.38	60.61	11.62
G61.48+0.09	0.15(.06)	0.28(.07)	22.06(.23)	1.17(.39)	2.11	22.10	2.72
K3–50 A	0.09(.04)	0.65(.12)	–25.25(.62)	7.53(1.58)	0.09(.04)	0.81(0.13)	–24.46(.77)	8.77(1.56)	6.89	–24.33	5.71
G78.44+2.66	0.82	1.37	1.63
W75N	0.61(.04)	2.72(.09)	9.43(.06)	4.17(.15)	0.55(.07)	2.35(0.12)	9.23(.10)	4.05(.23)	32.22	9.76	7.20
DR 21S	0.65(.07)	2.35(.12)	–2.56(.08)	3.36(.21)	0.16(.06)	1.29(0.16)	–2.42(.44)	7.35(1.18)	27.36	–1.25	15.18
W75OH	0.77(.04)	4.20(.09)	–2.70(0.06)	5.12(.13)	0.54(.06)	3.16(0.13)	–2.93(.11)	5.54(.26)	30.56	–2.65	6.72
G81.7+0.5	0.72(.12)	2.46(.22)	–2.91(.14)	3.22(.37)	0.29(.17)	1.25(0.30)	–3.09(.50)	4.02(1.10)	25.26	–1.66	14.83
CepA	0.20(.04)	0.67(.07)	–10.36(.17)	3.04(.41)	0.26(.06)	1.12(0.14)	–10.20(.28)	4.10(.67)	14.89	–9.90	7.11
NGC 7538	0.41(.06)	1.83(.12)	–55.81(.13)	4.25(.28)	0.17(.06)	0.87(0.14)	–56.12(.39)	4.75(.99)	16.38	–55.98	6.12
NGC 7538A	0.55(.04)	1.80(.07)	–57.27(.06)	3.04(.15)	0.30(.06)	1.39(0.12)	–57.31(.16)	4.23(.41)	14.67	–57.47	5.13

Table 5
Observing Results of C¹⁸O

Source Name	T_{mb} (K)	$\int T_{\text{mb}} d\nu$ (K km s ⁻¹)	V_{LSR} (km s ⁻¹)	FWHM (km s ⁻¹)
W3(OH)	4.48(.05)	19.70(.09)	-46.77(.01)	4.13(.02)
RCW142	9.63(.06)	49.74(.13)	17.28(.01)	4.85(.02)
G5.89-0.39	9.39(.19)	43.91(.37)	9.23(.02)	4.39(.04)
M8E	7.30(.04)	17.76(.06)	10.71(.01)	2.29(.01)
G8.67-0.36	8.04(.06)	45.65(.13)	35.34(.01)	5.34(.02)
W31	5.35(.05)	34.73(.10)	66.96(.01)	6.09(.02)
G10.6-0.4	13.80(.08)	88.45(.17)	-3.06(.01)	6.01(.01)
G11.94-0.62	5.26(.04)	25.49(.08)	38.27(.01)	4.55(.02)
W33cont	16.0(.05)	92.27(.11)	35.16(.01)	5.40(.01)
W33irs1	13.20(.19)	83.39(.41)	35.51(.01)	5.93(.03)
W33irs2	18.80(.32)	87.85(.62)	34.96(.01)	4.40(.04)
W33irs3	15.60(.12)	100.6(.26)	35.59(.01)	6.06(.02)
G14.33-0.64	5.16(.04)	19.55(.08)	22.29(.01)	3.56(.02)
W42	6.63(.08)	31.01	110.1(.01)	4.40(.03)
G29.96-0.02	7.94(.16)	37.21(.32)	97.75(.02)	4.40(.05)
G30.54+0.02	2.02(.09)	7.37(.15)	47.81(.03)	3.43(.08)
G33.92+0.11	8.29(.10)	28.18(.17)	107.6(.01)	3.19(.02)
G34.26+0.15	14.50(.09)	78.75(.20)	57.58(.01)	5.11(.02)
S76E	3.40(.06)	11.82(.10)	32.92(.01)	3.27(.03)
G35.20-0.74	2.97(.10)	10.64(.01)	33.81(.03)	3.37(.05)
G35.20-1.74	3.34(.05)	18.71(.11)	43.10(.02)	5.27(.04)
G37.543-0.11	2.01(.08)	7.86(.15)	52.74(.03)	3.68(.08)
G43.89-0.78	4.41(.20)	17.40(.35)	54.02(.04)	3.70(.09)
G45.45+0.06	3.40(.06)	15.50(.12)	58.63(.02)	4.28(.04)
W51M	9.20(.14)	90.07(.39)	57.53(.02)	9.20(.05)
W51D	8.73(1.07)	63.87(2.73)	61.47(.13)	6.87(.38)
G61.48+0.09	4.07(.07)	12.48(.12)	21.80(.01)	2.88(.02)
K3-50 A	2.96(.04)	21.29(.03)	-23.54(.01)	6.77(.02)
G78.44+2.66	3.53(.13)	7.91(.17)	0.96(.02)	2.11(.05)
W75N	9.34(.11)	38.59(.21)	9.67(.01)	3.88(.03)
DR 21S	9.02(0.16)	36.67(0.30)	-2.49(.02)	3.82(.04)
W75(OH)	10.7(.05)	48.31(.10)	-2.83(.01)	4.25(.01)
G81.7+0.5	9.95(.21)	36.45(.36)	-2.51(.02)	3.44(.04)
CepA	8.24(.09)	32.18(.15)	-10.50(.01)	3.67(.02)
NGC 7538	3.64(.10)	21.53(.21)	-56.42(.03)	5.56(.06)
NGC 7538A	6.06(.12)	26.96(.24)	-56.96(.02)	4.18(.05)

column density and abundance for different evolutionary sequences are also shown.

The number distribution of the CS, OCS, H₂S, and SO column densities for H II-HC and H II-only sources are shown in Figure 2, while the number distribution of CS, OCS, H₂S, and SO abundances are shown in Figure 3. Since we took into account sources undetected while calculating the median values, some median values do not fall in the middle of the sources shown. For example, for H II-only sources, OCS was detected in less than half of the sources, and only an upper limit for the median value could be obtained. We could see a trend of decreasing column density and abundance with the evolution of the clumps from these figures, which is confirmed by the median values. The median value of the H₂ column densities are 2.2×10^{23} and 9.9×10^{22} cm⁻² for H II-HC sources and H II-only sources, respectively. The decrease of H₂ column densities might be caused by the expansion of H II regions. The decrease in CS, OCS, H₂S, and SO column density and abundance strongly suggests that these molecules are probably destroyed by the UV radiation from the forming stars. For H II-HC sources, median values of CS, OCS, H₂S, and SO are 6.4×10^{14} , 1.2×10^{15} , 7.9×10^{14} , and 2.3×10^{14} cm⁻²,

respectively. Since H₂S is assumed to be optically thin in the calculation, the column density and abundance of H₂S might be underestimated, and the column density and abundance of H₂S is likely to be higher than OCS. All together, these results suggest that H₂S is probably the main sulfur carrier molecule in massive hot cores, while the column density and abundance of OCS are generally higher than CS and SO. This trend is not consistent with the hot core model of Viti et al. (2004), in which CS abundance is always higher than OCS. OCS was found to be the most abundant sulfur-bearing species in the starburst galaxy NGC 253 (Martin et al. 2005). Maity & Kaiser (2013) investigated the formation of sulfur species in interstellar ices during the irradiation of CS₂ and O₂ ices with energetic electrons at 12 K, and the sulfuretted molecules produced during the irradiation were SO₂, SO₃, and OCS. The large overabundance of OCS compared with the time-dependent sulfur chemistry model seems to support the idea that OCS is injected into the gas phase from grain mantles by low-velocity shocks (Martin et al. 2005). It would be helpful to investigate the relationship between OCS abundance and low-velocity shock tracers such as HNCO (Rodríguez-Fernández et al. 2010; Li et al. 2013).

Esplagues et al. (2014) modeled the sulfur chemistry evolution of the hot core and the plateau in Orion KL using the chemical models in Viti et al. (2004), with two different phases in calculations. Phase I starts with the collapsing cloud and the depletion of atoms and molecules onto grain surfaces. In this phase, the material collapses and atoms and molecules are depleted onto grain surfaces. The density increases with time according to the so-called modified collapse. In Phase II, the density is constant, a central protostar is formed, and the sublimation from grains takes place due to the warming up of the region. A time-dependent evaporation model was employed, as instantaneous sublimation is a more appropriate approximation only if the mass of the central stars is very high. We found that their hot core model of Orion KL could reproduce the general trend (H₂S > OCS > CS > SO) observed in this paper. According to their predictions, such trend are reproduced at about 7.5×10^4 yr after the central star switches on (see Figure 12 of Esplagues et al. 2014). The models that reproduce the the median column densities and abundances of sulfur-bearing species are those with an initial sulfur abundance of 0.1 times the sulfur solar abundance ($0.1 S_{\odot}$) and a density of about 10^7 cm⁻³. The column density of SO is higher than CS by an order of magnitude (see Table 2 in Esplagues et al. 2014) in the Orion hot core, whereas we found CS > SO, suggesting that our sources are on average older than the Orion hot core.

4.3. Abundance Ratios

Abundance ratios of chemically related molecules and with similar excitation conditions are thought to be more accurate than absolute molecular abundances (Ginard et al. 2012). Previous chemical models and observations all suggest that abundance ratios of sulfur-bearing molecules can be used to investigate the evolution of the early phases of massive star formation, and that sulfur-bearing species are good “chemical clocks” (e.g., Charnley 1997; Viti et al. 2004; Wakelam et al. 2004, 2011; Herpin et al. 2009).

According to the hot core model of Orion KL (Esplagues et al. 2014), for a $10 M_{\odot}$ star with an initial sulfur abundance of $0.1 S_{\odot}$ and a density of at least 10^7 cm⁻³, (i) [SO]/[OCS] > 1 up to 70,000 yr and < 1 after that; (ii) [CS]/[SO] < 1 up to 70,000 yr

Table 6
Column Densities and Abundances

Source Name	Column Density (cm^{-2})					Abundance				
	H ₂	CS	OCS	H ₂ S	SO	CS	OCS	H ₂ S	SO	
W3(OH)	1.12×10^{23}	4.61×10^{14}	1.40×10^{15}	5.09×10^{14}	3.05×10^{14}	4.10×10^{-9}	1.21×10^{-8}	4.53×10^{-9}	2.71×10^{-9}	
RCW142	2.84×10^{23}	9.84×10^{14}	2.16×10^{15}	1.23×10^{15}	2.00×10^{14}	3.47×10^{-9}	7.61×10^{-9}	4.33×10^{-9}	7.04×10^{-10}	
G5.89–0.39	2.50×10^{23}	8.97×10^{14}	1.53×10^{15}	1.15×10^{15}	7.24×10^{14}	3.58×10^{-9}	6.10×10^{-9}	4.58×10^{-9}	2.89×10^{-9}	
M8E	1.01×10^{23}	2.61×10^{14}	$^a 8.20 \times 10^{13}$	6.24×10^{13}	1.20×10^{14}	2.57×10^{-9}	8.20×10^{-10}	6.15×10^{-10}	1.18×10^{-9}	
G8.67–0.36	2.60×10^{23}	3.76×10^{14}	$^a 4.57 \times 10^{14}$	4.62×10^{14}	1.09×10^{14}	1.44×10^{-9}	1.76×10^{-9}	1.77×10^{-9}	4.19×10^{-10}	
W31	1.98×10^{23}	1.06×10^{15}	8.86×10^{15}	2.37×10^{15}	1.59×10^{14}	5.33×10^{-9}	4.47×10^{-8}	1.19×10^{-8}	8.02×10^{-10}	
G10.6–0.4	5.04×10^{23}	2.30×10^{15}	2.06×10^{15}	5.01×10^{15}	3.70×10^{14}	4.56×10^{-9}	4.07×10^{-9}	9.94×10^{-9}	7.33×10^{-10}	
G11.94–0.62	1.45×10^{23}	1.93×10^{14}	$^a 5.85 \times 10^{13}$	1.18×10^{14}	4.30×10^{13}	1.33×10^{-9}	4.03×10^{-10}	8.10×10^{-10}	2.95×10^{-10}	
W33cont	5.26×10^{23}	1.34×10^{15}	$^a 7.20 \times 10^{14}$	1.99×10^{15}	2.59×10^{14}	2.53×10^{-9}	1.37×10^{-9}	3.79×10^{-9}	4.91×10^{-10}	
W33irs1	4.75×10^{23}	8.86×10^{14}	$^a 5.35 \times 10^{14}$	7.88×10^{14}	1.67×10^{14}	1.86×10^{-9}	1.13×10^{-9}	1.65×10^{-9}	3.52×10^{-10}	
W33irs2	5.01×10^{23}	1.09×10^{15}	$^a 6.05 \times 10^{14}$	1.41×10^{15}	2.27×10^{14}	2.17×10^{-9}	1.20×10^{-9}	2.82×10^{-9}	4.54×10^{-10}	
W33irs3	5.74×10^{23}	1.36×10^{15}	$^a 6.40 \times 10^{14}$	2.10×10^{15}	2.80×10^{14}	2.36×10^{-9}	1.12×10^{-9}	3.65×10^{-9}	4.87×10^{-10}	
G14.33–0.64	1.11×10^{23}	2.81×10^{14}	6.3×10^{14}	2.18×10^{14}	1.82×10^{14}	2.52×10^{-9}	5.67×10^{-9}	1.95×10^{-9}	1.63×10^{-9}	
W42	1.77×10^{23}	6.00×10^{14}	1.89×10^{15}	1.05×10^{15}	1.14×10^{14}	3.39×10^{-9}	1.07×10^{-8}	5.91×10^{-9}	6.43×10^{-10}	
G29.96–0.02	2.12×10^{23}	6.54×10^{14}	1.21×10^{15}	8.61×10^{14}	1.31×10^{14}	3.08×10^{-9}	5.71×10^{-9}	4.05×10^{-9}	6.19×10^{-10}	
G30.54+0.02	4.20×10^{22}	7.44×10^{12}	1.77×10^{-10}	
G33.92+0.11	1.61×10^{23}	2.17×10^{14}	$^a 2.82 \times 10^{14}$	1.30×10^{14}	4.61×10^{13}	1.34×10^{-9}	1.80×10^{-9}	8.08×10^{-10}	2.86×10^{-10}	
G34.26+0.15	4.49×10^{23}	1.60×10^{15}	4.74×10^{15}	2.14×10^{15}	4.20×10^{14}	3.55×10^{-9}	1.05×10^{-8}	4.76×10^{-9}	9.36×10^{-10}	
S76E	6.74×10^{22}	1.89×10^{14}	$^a 3.34 \times 10^{14}$	3.81×10^{13}	1.53×10^{14}	2.80×10^{-9}	5.00×10^{-9}	5.65×10^{-10}	2.27×10^{-9}	
G35.20–0.74	6.07×10^{22}	1.20×10^{14}	$^a 2.82 \times 10^{14}$	7.45×10^{13}	6.60×10^{13}	1.97×10^{-9}	4.65×10^{-9}	1.22×10^{-9}	1.08×10^{-9}	
G35.20–1.74	1.07×10^{23}	8.29×10^{13}	2.43×10^{15}	8.49×10^{13}	5.44×10^{13}	7.77×10^{-10}	2.27×10^{-8}	7.95×10^{-10}	5.09×10^{-10}	
G37.543–0.11	4.48×10^{22}	6.91×10^{13}	1.30×10^{13}	1.54×10^{-9}	2.91×10^{-10}	
G43.89–0.78	9.91×10^{22}	2.93×10^{13}	2.95×10^{-10}	
G45.45+0.06	8.84×10^{22}	1.64×10^{14}	4.19×10^{13}	1.85×10^{-9}	...	4.70×10^{-10}	4.74×10^{-10}	
W51M	5.13×10^{23}	2.73×10^{15}	1.17×10^{16}	4.04×10^{15}	9.42×10^{14}	5.32×10^{-9}	2.27×10^{-8}	7.85×10^{-9}	1.83×10^{-9}	
W51D	3.64×10^{23}	9.12×10^{14}	$^a 1.70 \times 10^{15}$	1.08×10^{15}	4.31×10^{14}	2.50×10^{-9}	4.67×10^{-9}	2.97×10^{-9}	1.18×10^{-9}	
G61.48+0.09	7.11×10^{22}	4.28×10^{13}	3.20×10^{13}	6.03×10^{-10}	...	4.18×10^{-8}	4.50×10^{-10}	
K3–50 A	1.20×10^{23}	9.98×10^{13}	...	1.40×10^{14}	1.05×10^{14}	8.22×10^{-10}	...	1.15×10^{-9}	8.62×10^{-10}	
G78.44+2.66	4.50×10^{22}	1.25×10^{13}	2.76×10^{-10}	
W75N	2.20×10^{23}	4.18×10^{14}	2.53×10^{15}	4.07×10^{14}	4.90×10^{14}	1.89×10^{-9}	1.14×10^{-8}	1.85×10^{-9}	2.22×10^{-9}	
DR 21S	2.09×10^{23}	3.61×10^{14}	1.0×10^{15}	2.23×10^{14}	4.16×10^{14}	1.72×10^{-9}	4.78×10^{-9}	1.06×10^{-9}	1.98×10^{-9}	
W75(OH)	2.75×10^{23}	6.45×10^{14}	3.00×10^{15}	5.47×10^{14}	4.65×10^{14}	2.34×10^{-9}	1.09×10^{-8}	1.98×10^{-9}	1.68×10^{-9}	
G81.7+0.5	2.08×10^{23}	3.78×10^{14}	...	2.17×10^{14}	3.84×10^{14}	1.81×10^{-9}	...	1.04×10^{-9}	1.84×10^{-9}	
CepA	1.83×10^{23}	1.03×10^{14}	...	1.94×10^{14}	2.26×10^{14}	5.60×10^{-10}	...	1.05×10^{-9}	1.23×10^{-9}	
NGC 7538	1.23×10^{23}	2.81×10^{14}	$^a 5.65 \times 10^{14}$	1.51×10^{14}	2.49×10^{14}	2.28×10^{-9}	4.60×10^{-9}	1.22×10^{-9}	2.02×10^{-9}	
NGC 7538A	1.54×10^{23}	2.76×10^{14}	1.16×10^{15}	2.41×10^{14}	2.23×10^{14}	1.79×10^{-9}	7.54×10^{-9}	1.56×10^{-9}	1.45×10^{-9}	
median _{H II – HC}	2.2×10^{23}	6.4×10^{14}	1.2×10^{15}	7.9×10^{14}	2.3×10^{14}	2.5×10^{-9}	5.0×10^{-9}	2.8×10^{-9}	1.1×10^{-9}	
median _{H II – only}	9.9×10^{22}	8.3×10^{13}	$< 8.2 \times 10^{13}$	4.2×10^{13}	4.2×10^{13}	8.2×10^{-10}	$< 8.1 \times 10^{-10}$	4.7×10^{-10}	4.5×10^{-10}	

Notes. The column density of CS is derived using ^{13}CS . The column density of OCS is derived using O^{13}CS except for those labeled with a , which is derived using OCS by assuming an optical depth of 5.

and >1 after that; and (iii) $[\text{OCS}]/[\text{CS}] >1$ up to 85,000 yr and <1 after that. Both $[\text{OCS}]/[\text{H}_2\text{S}]$ and $[\text{CS}]/[\text{H}_2\text{S}]$ are not chemical clock candidate, since they are always smaller than 1. Herpin et al. (2009) investigated possible correlations of the molecular abundance ratios with the order of evolution of four objects, and also proposed that molecular ratios like $[\text{OCS}]/[\text{H}_2\text{S}]$, $[\text{CS}]/[\text{H}_2\text{S}]$, $[\text{SO}]/[\text{OCS}]$, and $[\text{CS}]/[\text{SO}]$ might be a good indicator of evolution. However, their trend is different from that of Esplugues et al. (2014). We calculated abundance ratios and present them in Table 7. Individual abundance ratios for H II–HC sources in which the chemistry is dominated by hot core are discussed below.

$[\text{OCS}]/[\text{H}_2\text{S}]$. The median value for a H II–HC source is 1.8, which is higher than predicted in Esplugues et al. (2014). This inconsistency is likely to be caused by the assumption of optical thinness of H_2S .

$[\text{CS}]/[\text{H}_2\text{S}]$. The median values for H II–HC is 0.8, which is not in contradiction with the model prediction.

$[\text{SO}]/[\text{OCS}]$. The median values for H II–HC is 0.2, suggesting ages are older than 70,000 yr. $[\text{SO}]/[\text{OCS}]$ is smaller than 1 for all the sources except for M8E, which belongs to the H II–only group.

$[\text{CS}]/[\text{SO}]$. The median values for H II–HC and H II–only regions are 2.9 and 1.3, respectively. Both of them are larger than 1, suggesting ages older than 70,000 yr (Esplugues et al. 2014).

$[\text{OCS}]/[\text{CS}]$. The median values for H II–HC is 2.0, suggesting ages younger than 85,000 yr (Esplugues et al. 2014).

According to the analysis above, most of these H II–HC sources should be older than 70,000 yr, and younger than 85,000 yr. Both McKee & Tan (2003) and Mac Low et al. (2007) derived typical timescales for the formation of a massive star of about 10^5 yr. Gerner et al. (2014) also found that the total chemical timescale for the high-mass star formation process is on the order of 10^5 yr. It appears that our result is in agreement with theoretical estimates and chemical

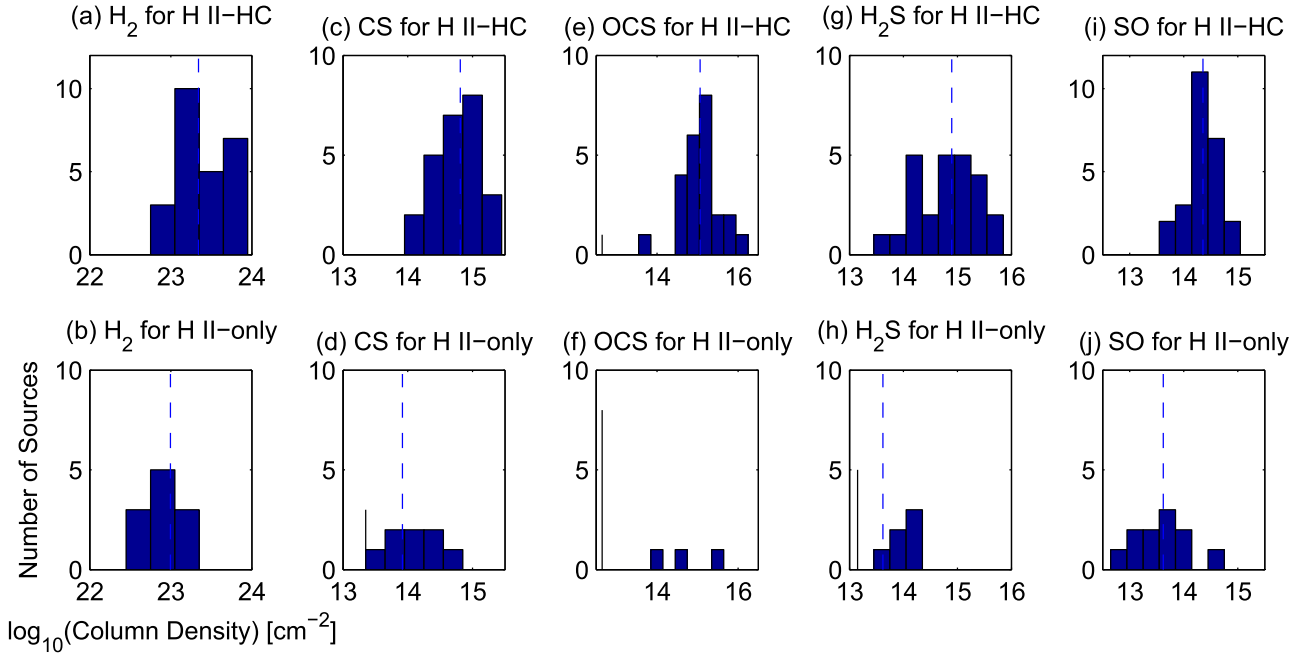


Figure 2. Histograms of the number distributions of H_2 , CS, OCS, H_2S , and SO column densities for H II-HC and H II-only sources. The vertical dashed lines indicate the median values of the column densities for each distribution. Median values of the column densities are given in Table 3. The vertical bars stand for sources undetected in a particular molecular line emission.

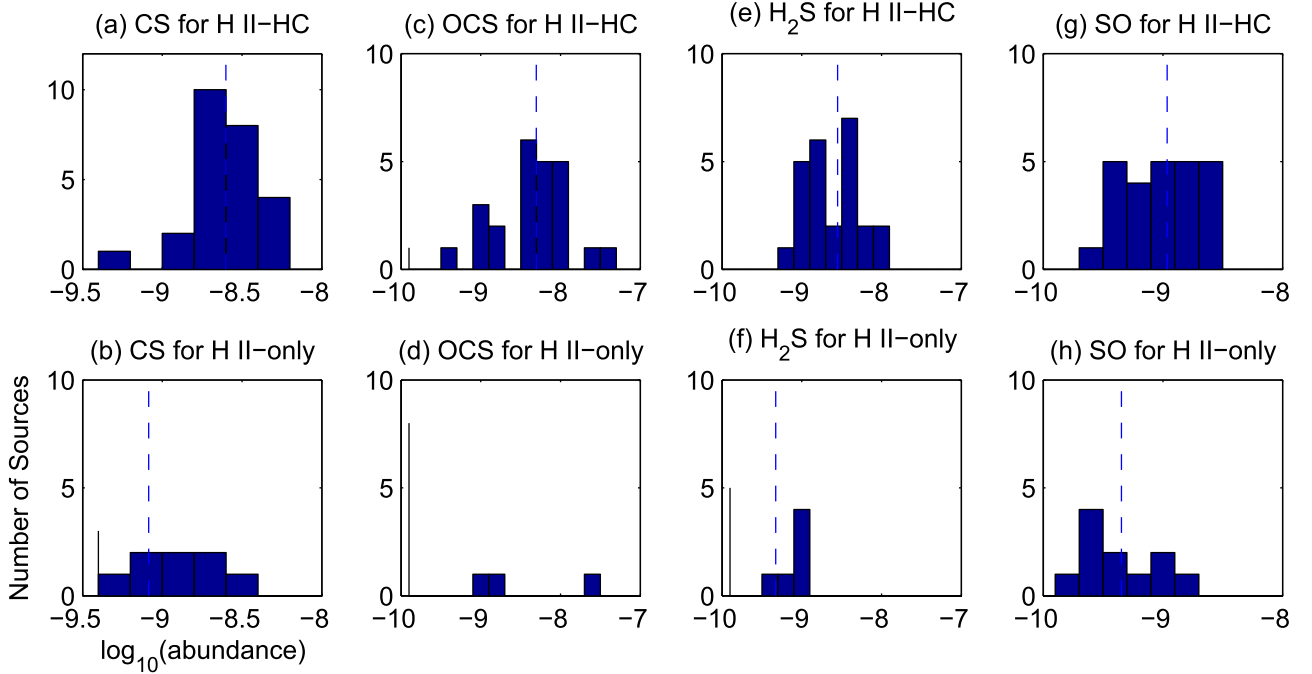


Figure 3. Histograms of the number distributions of CS, OCS, H_2S , and SO abundances for H II-HC and H II-only sources. The vertical dashed lines indicate the median values of the column densities for each distribution. Median values of the column densities are given in Table 3. Similar to Figure 2, the vertical bars also stand for sources undetected in a particular molecular line emission.

evolution model, suggesting that $[\text{SO}]/[\text{OCS}]$, $[\text{CS}]/[\text{SO}]$, and $[\text{OCS}]/[\text{CS}]$ could be used as a “chemical clock” in massive star-forming regions.

For H II-only sources, abundance ratios were obtained for only a few sources. Comparing with H II-HC sources, the abundance ratios have not changed significantly, suggesting that the sulfur species that are left are still from relatively young components. Hatchell et al. (1998a) proposed that the non-

detection of hot gas may be caused by the low column density of hot gas or the small source size of these sources. However, the obvious decline of abundance imply that models that take into account effects of UV radiation from forming stars are needed to explain results of H II-only sources.

As mentioned above, massive star formation always take place on rather short timescales and in clustered environments. Due to the large beam of our observations, some regions with

Table 7
Abundance Ratio

Source Name	OCS/H ₂ S	CS/H ₂ S	SO/OCS	CS/SO	OCS/CS
W3(OH)	2.69	0.90	0.22	1.51	2.97
RCW142	1.76	0.80	0.09	4.93	2.19
G5.89–0.39	1.33	0.78	0.47	1.24	1.70
M8E	1.32	4.19	1.46	2.18	0.31
G8.67–0.36	0.99	0.81	0.24	3.44	1.21
W31	3.74	0.45	0.02	6.65	8.38
G10.6–0.4	0.41	0.46	0.18	6.23	0.89
G11.94–0.62	0.50	1.64	0.73	4.50	0.30
W33cont	0.36	0.67	0.36	5.16	0.54
W33irs1	0.68	1.12	0.31	5.29	0.60
W33irs2	0.43	0.77	0.38	4.79	0.56
W33irs3	0.31	0.65	0.44	4.86	0.47
G14.33–0.64	2.90	1.29	0.29	1.55	2.25
W42	1.81	0.57	0.06	5.28	3.16
G29.96–0.02	1.41	0.76	0.11	4.98	1.85
G30.54+0.02
G33.92+0.11	2.17	1.67	0.16	4.70	1.30
G34.26+0.15	2.22	0.75	0.09	3.80	2.96
S76E	8.76	4.96	0.45	1.23	1.77
G35.20–0.74	3.78	1.61	0.24	1.82	2.35
G35.20–1.74	28.55	0.98	0.02	1.52	29.23
G37.543–0.11
G43.89–0.78
G45.45+0.06	3.92	...
W51M	2.90	0.68	0.08	2.90	4.28
W51D	1.57	0.84	0.25	2.11	1.86
G61.48+0.09	...	0.17	...	16.05	...
K3–50 A	...	0.71	...	0.95	...
G78.44+2.66
W75N	6.21	1.03	0.19	0.85	6.05
DR 21S	4.48	1.61	0.42	0.87	2.77
W75(OH)	5.49	1.18	0.15	1.39	4.66
G81.7+0.5	...	1.74	...	0.98	...
CepA	...	0.53	...	0.45	...
NGC 7538	3.75	1.86	0.44	1.13	2.01
NGC 7538A	4.81	1.15	0.19	1.24	4.19
median _{H II – HC}	1.8	0.8	0.2	2.9	2.0
median _{H II – only}	1.3	...

embedded H II regions also show characteristics of a hot core, and thus observations with higher resolution could better constrain the age of these sources.

5. SUMMARY

We have conducted a systematic study of sulfur-bearing molecules toward 36 massive star-forming regions with embedded H II regions using the CSO telescope. We divided these sources into H II-HC and H II-only to characterize their evolutionary stages. OCS emission is optically thick for thirteen sources, with optical depth ranging from 5 to 16. We calculated the column density and abundance of these molecules under LTE conditions, and H₂S seems to be the most abundant sulfuretted molecule. The hot core model of Orion KL could reproduce the general trend (H₂S > OCS > CS > SO) observed in this paper. The models that reproduce the the median column densities and abundances of sulfur-bearing species are those with an initial sulfur abundance of 0.1 times the sulfur solar abundance (0.1S_⊙) and a density of at least $5 \times 10^6 \text{ cm}^{-3}$. We found that the

column density and abundance of CS, OCS, H₂S, and SO decrease significantly from H II-HC to H II-only sources. We derived abundance ratios of [OCS]/[H₂S], [CS]/[H₂S], [SO]/[OCS], [CS]/[SO], and [OCS]/[CS], and made comparisons between the ages predicted by hot core models with massive star formation models. We conclude that these abundance ratios are potential good chemical clocks in massive star-forming regions.

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