



ATLAS Note

GROUP-2017-XX

15th March 2021



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2 WZ + Heavy Flavor Production in pp collisions 3 at $\sqrt{s} = 13 \text{ TeV}$

4

The ATLAS Collaboration

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A measurement of WZ produced with an associated heavy flavor jet is performed using 140
6 fb^{-1} of proton-proton collision data at $\sqrt{s} = 13 \text{ TeV}$ from the ATLAS experiment at the
7 LHC. The measurement is performed in the fully leptonic decay mode, $\text{WZ} \rightarrow \ell\bar{\nu}\ell\bar{\nu}$. The
8 cross-section of $\text{WZ} + b\text{-jets}$ is measured to be $X \pm X \pm X$, while the cross-section of $\text{WZ} +$
9 charm is measured as X , with a correlation of X between the two processes.

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48 **1 Changes and outstanding items**

49 **1.1 Changelog**

50 This is version 6

51 **1.1.1 Changes relative to v5**

- 52 • added list of DSIDs to an appendix
53 • included systematics on jet migrations
54 • Updated results section to include simultaneous fit over 1-jet and 2-jet bins, describe
55 unfolding procedure
56 • Updated other sections to account for this change

57 **1.1.2 Changes relative to v4**

- 58 • Fixed various typos, clarified wording
59 • Expanded info about JER uncertainties, electron systematics, theory uncertainties
60 • removed a table on lepton selection, included information in the text instead
61 • Plotted lepton Pt Z and W for Zjet CR, rather than lep 1 and 2
62 • fixed binning in kinematic plots
63 • Included prefit and postfit yield tables
64 • added signal modelling systematics
65 • included alternate fit studies with tZ included in signal

66 **1.1.3 Changes relative to v3**

- 67 • Merged introduction into executive summary, including unblinding details and list of
68 SRs/CRs used
69 • listed ptag used (p4133), and release (AB 21.2.127)
70 • Included table reftab:xsecUnc listing x-sec uncertainties used
71 • Removed selection criteria listed in table ?? (QMisID, AmbiguityType) that were removed
72 from the analysis

- 73 • specified variable used to make truth jet flavor determination (HadronConeExclTruthLa-
 74 belID)
- 75 • fixed bug in MtLepMet calculation, updated selection/fits to account for this
- 76 • Included plots of MtLepMet and PtZ, swapped lep 1 and 2 p_T plots for lep W and lep Z
 77 plots
- 78 • updated tZ BDT training to reduce overfitting, updated plots to include error bars, feature
 79 importance
- 80 • updated table 7 to clarify selection, fix the tZ_BDT cut used
- 81 • replace a few broken ntuples which included large weight events
- 82 • include DL1r distribution for Z+jets and $t\bar{t}$ VRs
- 83 • Expanded section on fakes, included information on derived scale factors from VRs.
- 84 • Changed the kinematic plots to include $p_T(Z)$ and $m_T(W)$, list lepton p_T based on W and
 85 Z candidates.

86 **1.1.4 Changes relative to v2**

- 87 • Added alternate VBS samples to include missing b-jet diagrams
- 88 • Included a section on tZ interference effects, ??.
- 89 • Updated to reflect changes for 2018, including the move to PFlow jets, DL1r, updated
 90 trigger, and updated AnalysisBase version (now 21.2.127)
- 91 • Revised fit regions, using separate 1-jet and 2-jet fits, with all 2-j regions included
- 92 • updated plots for tZ BDT, added details about the model
- 93 • Included truth jet information

94 **1.1.5 Changes relative to v1**

- 95 • Added GRL list
- 96 • Fixed latex issue in line 92, typo in line 172
- 97 • Added tables 5 and ??, summarizing the event and object selection
- 98 • Added table 2, which includes the DSID of samples used
- 99 • Included reference to WZ inclusive paper in introduction

100 **1.2 Outstanding Items**

- 101 • Complete interference studies, apply any interference effects observed as a systematic
102 • Update results section with additional studies, possibly including:
103 – Truth jet migration studies
104 – Simultaneous fit over 1j and 2j
105 – Impact of allowing tZ to float
106 • Unblind, update plots and fits to include data
107 • Add cross-section, significance once unblinded

108 2 Executive Summary

109 The production of WZ in association with a heavy flavor jet represents an important background
 110 for many major analyses. This includes any process with leptons and b-jets in the final state,
 111 such as $t\bar{t}H$, $t\bar{t}W$, and $t\bar{t}Z$. While precise measurements have been made of WZ production
 112 [1], $WZ +$ heavy flavor remains poorly understood. This is largely because the QCD processes
 113 involved in the production of the b-jet make it difficult to simulate accurately. This introduces a
 114 large uncertainty for analyses that include this process as a background.

115 Motivated by its relevance to the $t\bar{t}H$ multilepton analysis, we perform a study of the fully
 116 leptonic decay mode of this channel; that is, events where both the W and Z decay leptonically.
 117 Because WZ has no associated jets at leading order, while the major backgrounds for this channel
 118 tend to have high jet multiplicity, events with more than two jets are rejected. This gives a final
 119 state signature of three leptons and one or two jets.

120 Events that meet this selection criteria are sorted into pseudo-continuous b-tagging regions based
 121 on the DL1r b-tag score of their associated jets. This is done to separate $WZ +$ b-jet events from
 122 $WZ +$ charm and $WZ +$ light jets. These regions are fit to data in order make a more accurate
 123 estimate of the contribution of $WZ +$ heavy-flavor, where heavy-flavor jets include b-jets and
 124 charm jets. The full Run-2 dataset collected by the ATLAS detector, representing 139 fb^{-1} of
 125 data from pp collisions at $\sqrt{s} = 13 \text{ TeV}$, is used for this study.

126 All backgrounds are accounted for using Monte Carlo, with the simulation of non-prompt lepton
 127 backgrounds - $Z+jets$ and $t\bar{t}$ - validated using non-prompt Validation Regions.

128 Section 3 details the data and Monte Carlo (MC) samples used in the analysis. The reconstruction
 129 of various physics objects is described in section 4. Section 5 describes the event selection applied
 130 to these samples, along the definitions of the various regions used in the fit. The multivariate
 131 analysis techniques used to separate the tZ background from $WZ +$ heavy flavor are described in
 132 section 6. Section 7 describes the various sources of systematic uncertainties considered in the
 133 fit. Finally, the results of the analysis are summarized in section 8, followed by a brief conclusion
 134 in section 9.

135 The analysis aims to report a cross-section measurement of $WZ+b$ and $WZ+charm$, along with
 136 their correlations, for both 1-jet and 2-jet exclusive regions. The proposed fiducial region for
 137 these measurements includes events with three leptons, where the invariant mass of at least one
 138 opposite charge, same flavor lepton pair falls within 10 GeV of 91.2 GeV , with 1 or 2 associated
 139 jets. An alternate version of the measurement is included in the appendix, which considers tZ as
 140 part of the $WZ+b$ signal.

141 The current state of the analysis shows blinded results for the full Run 2 dataset. Regions
 142 containing $>5\%$ $WZ+b$ events are blinded, and results are from Asimov, MC only fits. Expected
 143 significance and cross-section numbers are reported.

¹⁴⁴ 3 Data and Monte Carlo Samples

¹⁴⁵ Both data and Monte Carlo samples used in this analysis were prepared in the `xAOD` format,
¹⁴⁶ which was used to produce a `DxAOD` sample in the `HIGG8D1` derivation framework. The `HIGG8D1`
¹⁴⁷ framework is designed for the $t\bar{t}H$ multi-lepton analysis, which targets events with multiple
¹⁴⁸ leptons as well as tau hadrons. This framework skims the dataset to remove unneeded variables
¹⁴⁹ as well as entire events. Events are removed from the derivations that do not meet the following
¹⁵⁰ selection:

- ¹⁵¹ • at least two light leptons within a range $|\eta| < 2.6$, with leading lepton $p_T > 15$ GeV and
¹⁵² subleading lepton $p_T > 5$ GeV
- ¹⁵³ • at least one light lepton with $p_T > 15$ GeV within a range $|\eta| < 2.6$, and at least two hadronic
¹⁵⁴ taus with $p_T > 15$ GeV.

¹⁵⁵ Samples were then generated from these `HIGG8D1` derivations with p-tag of p4134 using Ana-
¹⁵⁶ lysisBase version 21.2.127 modified to include custom variables..

¹⁵⁷ 3.1 Data Samples

¹⁵⁸ The study uses a sample of proton-proton collision data collected by the ATLAS detector from
¹⁵⁹ 2015 through 2018 at an energy of $\sqrt{s} = 13$ TeV, which represents an integrated luminosity of
¹⁶⁰ 139 fb^{-1} . This data set was collected with a bunch crossing rate of 25 ns. All data used in this
¹⁶¹ analysis was verified by data quality checks, having been included in the following Good Run
¹⁶² Lists:

- ¹⁶³ • `data15_13TeV.periodAllYear_DetStatus-v79-repro20-02_DQDefects-00-02-02`
¹⁶⁴ `_PHYS_StandardGRL_All_Good_25ns.xml`
- ¹⁶⁵ • `data16_13TeV.periodAllYear_DetStatus-v88-pro20-21_DQDefects-00-02-04`
¹⁶⁶ `_PHYS_StandardGRL_All_Good_25ns.xml`
- ¹⁶⁷ • `data17_13TeV.periodAllYear_DetStatus-v97-pro21-13_Unknown_PHYS_StandardGRL`
¹⁶⁸ `_All_Good_25ns_Triggerno17e33prim.xml`
- ¹⁶⁹ • `data18_13TeV.periodAllYear_DetStatus-v102-pro22-04_Unknown_PHYS_StandardGRL`
¹⁷⁰ `_All_Good_25ns_Triggerno17e33prim.xml`

¹⁷¹ Runs included from the AllYear period containers are included.

172 3.2 Monte Carlo Samples

173 Several different generators were used to produce Monte Carlo simulations of the signal and
 174 background processes. For all samples, the response of the ATLAS detector is simulated using
 175 Geant4. The WZ signal samples are simulated using Sherpa 2.2.2 [2]. Specific information
 176 about the Monte Carlo samples being used can be found in Table 1. A list of the specific samples
 177 used by data set ID is shown in Table 2.

Table 1: The configurations used for event generation of signal and background processes, including the event generator, matrix element (ME) order, parton shower algorithm, and parton distribution function (PDF).

Process	Event generator	ME order	Parton Shower	PDF
WZ, VV	SHERPA 2.2.2	MEPS NLO	SHERPA	CT10
tZ	MG5_AMC	LO	PYTHIA 6	CTEQ6L1
t̄tW	MG5_AMC (SHERPA 2.1.1)	NLO (LO multileg)	PYTHIA 8 (SHERPA)	NNPDF 3.0 NLO (NNPDF 3.0 NLO)
t̄t(Z/γ* → ll)	MG5_AMC	NLO	PYTHIA 8	NNPDF 3.0 NLO
t̄tH	MG5_AMC (MG5_AMC)	NLO (NLO)	PYTHIA 8 (HERWIG++)	NNPDF 3.0 NLO [Ball:2014uwa] (CT10 [ct10])
tHqb	MG5_AMC	LO	PYTHIA 8	CT10
tHW	MG5_AMC (SHERPA 2.1.1)	NLO (LO multileg)	HERWIG++ (SHERPA)	CT10 (NNPDF 3.0 NLO)
tWZ	MG5_AMC	NLO	PYTHIA 8	NNPDF 2.3 LO
t̄t, t̄t̄t̄	MG5_AMC	LO	PYTHIA 8	NNPDF 2.3 LO
t̄tW+W-	MG5_AMC	LO	PYTHIA 8	NNPDF 2.3 LO
t̄t	Powheg-BOX v2 [powhegtt]	NLO	PYTHIA 8	NNPDF 3.0 NLO
t̄tγ	MG5_AMC	LO	PYTHIA 8	NNPDF 2.3 LO
s-, t-channel, Wt single top	Powheg-BOX v1 [powhegstp]	NLO	PYTHIA 6	CT10
qqVV, VVV				
Z → l+l-	SHERPA 2.2.1	MEPS NLO	SHERPA	NNPDF 3.0 NLO

Sample	DSID
WZ	364253, 364739-42
VV	364250, 364254, 364255, 363355-60, 364890
t̄W	410155
t̄Z	410156, 410157, 410218-20
low mass t̄Z	410276-8
Rare Top	410397, 410398, 410399
single Top	410658-9, 410644-5
three Top	304014
four Top	410080
t̄WW	410081
Z + jets	364100-41
low mass Z + jets	364198-215
W + jets	364156-97
Vγ	364500-35
tZ	410560
tW	410013-4
WtZ	410408
VVV	364242-9
VH	342284-5
WtH	341998
t̄tγ	410389
t̄t	410470
t̄tH	345873-5, 346343-5

Table 2: List of Monte Carlo samples by data set ID used in the analysis.

178 4 Object Reconstruction

179 All regions defined in this analysis share a common lepton, jet, and overall event preselection.
 180 The selection applied to each physics object is detailed here; the event preselection, and the
 181 selection used to define the various fit regions, is described in Section 5.

182 4.1 Trigger

183 Events are required to be selected by dilepton triggers, as summarized in Table 3.

Dilepton triggers (2015)	
$\mu\mu$ (asymm.)	HLT_mu18_mu8noL1
ee (symm.)	HLT_2e12_lhloose_L12EM10VH
$e\mu, \mu e$ (\sim symm.)	HLT_e17_lhloose_mu14
Dilepton triggers (2016)	
$\mu\mu$ (asymm.)	HLT_mu22_mu8noL1
ee (symm.)	HLT_2e17_lhvloose_nod0
$e\mu, \mu e$ (\sim symm.)	HLT_e17_lhloose_nod0_mu14
Dilepton triggers (2017)	
$\mu\mu$ (asymm.)	HLT_mu22_mu8noL1
ee (symm.)	HLT_2e24_lhvloose_nod0
$e\mu, \mu e$ (\sim symm.)	HLT_e17_lhloose_nod0_mu14
Dilepton triggers (2018)	
$\mu\mu$ (asymm.)	HLT_mu22_mu8noL1
ee (symm.)	HLT_2e24_lhvloose_nod0
$e\mu, \mu e$ (\sim symm.)	HLT_e17_lhloose_nod0_mu14

Table 3: List of lowest p_T -threshold, un-prescaled dilepton triggers used for 2015-2018 data taking.

184 **4.2 Light leptons**

- 185 Electron candidates are reconstructed from energy clusters in the electromagnetic calorimeter
 186 that are associated with charged particle tracks reconstructed in the inner detector [3]. Electron
 187 candidates are required to have $p_T > 10$ GeV and $|\eta_{\text{cluster}}| < 2.47$. Muon candidates are
 188 reconstructed by combining inner detector tracks with track segments or full tracks in the muon
 189 spectrometer [4]. Muon candidates are required to have $p_T > 10$ GeV and $|\eta| < 2.5$. Candidates
 190 in the transition region between different electromagnetic calorimeter components, $1.37 <$
 191 $|\eta_{\text{cluster}}| < 1.52$, are rejected. A multivariate likelihood discriminant combining shower shape
 192 and track information is used to distinguish real electrons from hadronic showers (fake electrons).
 193 To further reduce the non-prompt electron contribution, the track is required to be consistent
 194 with originating from the primary vertex; requirements are imposed on the transverse impact
 195 parameter significance ($|d_0|/\sigma_{d_0} < 5$) and the longitudinal impact parameter ($|\Delta z_0 \sin \theta_\ell| < 0.5$
 196 mm).
- 197 Muon candidates are reconstructed by combining inner detector tracks with track segments or
 198 full tracks in the muon spectrometer [4]. Muon candidates are required to have $p_T > 10$ GeV
 199 and $|\eta| < 2.5$. The longitudinal impact parameter is the same for both electrons and muons, while
 200 muons are required to pass a slightly tighter transverse impact parameter, $|d_0|/\sigma_{d_0} < 3$.

201 All leptons are required to be isolated, as defined through the standard `isolationFixedCutLoose`
 202 working point supported by combined performance groups. Leptons are additionally required to
 203 pass a non-prompt BDT selection described in detail in [5]. Optimized working points and scale
 204 factors for this BDT are taken from that analysis.

205 **4.3 Jets**

206 Jets are reconstructed from calibrated topological clusters built from energy deposits in the
 207 calorimeters [6], using the anti- k_t algorithm with a radius parameter $R = 0.4$. Jets with energy
 208 contributions likely arising from noise or detector effects are removed from consideration [7],
 209 and only jets satisfying $p_T > 25$ GeV and $|\eta| < 2.5$ are used in this analysis. For jets with
 210 $p_T < 60$ GeV and $|\eta| < 2.4$, a jet-track association algorithm is used to confirm that the jet
 211 originates from the selected primary vertex, in order to reject jets arising from pileup collisions
 212 [8].

213 **4.4 B-tagged Jets**

214 In order to make a measurement of $WZ +$ heavy flavor it is necessary to distinguish these events
 215 from $WZ +$ light jets. For this purpose, the DL1r b-tagging algorithm is used to distinguish
 216 heavy flavor jets from lighter ones. The DL1r algorithm uses jet kinematics, particularly jet
 217 vertex information, as input for a neural network which assigns each jet a score designed to
 218 reflect how likely that jet is to have originated from a b-quark.

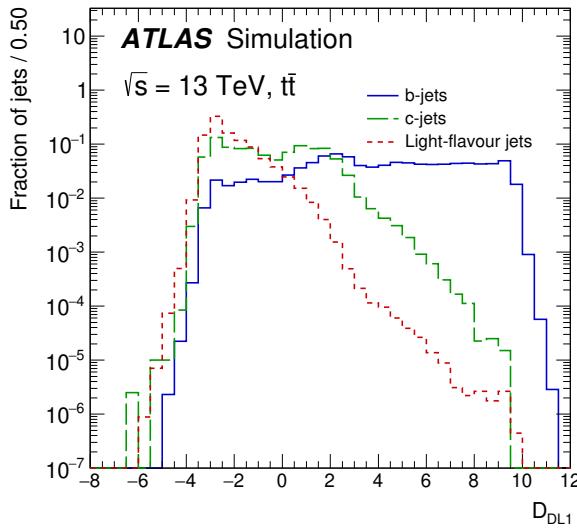


Figure 1: Output distribution of the DL1r algorithm for b-jets, charm jets, and light jets

219 From the output of the BDT, working points (WPs) are developed based on the efficiency of truth
 220 b-jets at particular values of the DL1r algorithm. The working points used in this analysis are
 221 summarized in Table 4.

WP	none	loose	medium	tight	tightest
b eff.	-	85%	77%	70%	60%

Table 4: B-tagging Working Points by tightness and b-jet efficiency

222 A tighter WP will accept fewer b-jets, but reject a higher fraction of charm and light jets.
 223 Generally, analyses that include b-jets will use a fixed working point, for example, requiring that
 224 a jet pass the 70% threshold. By instead treating these working point as bins, e.g. events with
 225 jets that fall between the 85% and 77% WPs fall into one bin, while events with jets passing the
 226 60% WP fall into another, and looking at the full psuedo-continuous DL1r spectrum of the jets,
 227 additional information can be gained. The psuedo-continuous b-tag spectrum is used in this case
 228 to separate out WZ + b, WZ + charm, and WZ + light.

229 4.5 Missing transverse energy

230 Missing transverse momentum (E_T^{miss}) is used as part of the event selection. The missing
 231 transverse momentum vector is defined as the inverse of the sum of the transverse momenta of
 232 all reconstructed physics objects as well as remaining unclustered energy, the latter of which is
 233 estimated from low- p_T tracks associated with the primary vertex but not assigned to a hard object,
 234 with object definitions taken from [9]. Light leptons considered in the E_T^{miss} reconstruction are
 235 required to have $p_T > 10 \text{ GeV}$, while jets are required to have $p_T > 20 \text{ GeV}$.

236 5 Event Selection and Signal Region Definitions

237 Event are required to pass a preselection described in Section 5.1 and summarized in Table 5.
 238 Those that pass this preselection are divided into various fit regions described in Section 5.2,
 239 based on the number of jets in the event, and the b-tag score of those jets.

240 5.1 Event Preselection

241 Events are required to include exactly three reconstructed light leptons passing the requirement
 242 described in 4.2, which have a total charge of ± 1 . As the opposite sign lepton is found to be
 243 prompt the vast majority of the time [5], it is required to be loose and isolated, as defined though
 244 the standard `isolationFixedCutLoose` working point supported by combined performance

245 groups. The same sign leptons are required to be very tightly isolated, as per the recommended
 246 `isolationFixedCutTight`.

247 The leptons are ordered in the analysis code as 0, 1, and 2. Lepton 0 is the lepton whose charge
 248 is opposite the other two. Lepton 1 is the lepton closest to the opposite charge lepton, i.e. the
 249 smallest ΔR , leaving lepton 2 as the lepton further from the opposite charge lepton. Lepton 0
 250 is required to have $p_T > 10$ GeV, while the same sign leptons, 1 and 2, are required to have
 251 $p_T > 20$ GeV to reduce the contribution of non-prompt leptons.

252 The invariant mass of at least one pair of opposite sign, same flavor leptons is required to fall
 253 within 10 GeV of the mass of the Z boson, 91.2 GeV. Events where one of the opposite sign pairs
 254 have an invariant mass less than 12 GeV are rejected in order to suppress low mass resonances.

255 An additional requirement is placed on the missing transverse energy, $E_T^{\text{miss}} > 20$ GeV, and the
 256 transverse mass of the W candidate, $m(E_T^{\text{miss}} + l_{\text{other}}) > 30$ GeV, where E_T^{miss} is the missing
 257 transverse energy, and l_{other} is the lepton not included in the Z-candidate.

258 Events are required to have one or two reconstructed jets passing the selection described in
 259 Section 4.3. Events with more than two jets are rejected in order to reduce the contribution of
 260 backgrounds such as $t\bar{t}Z$ and $t\bar{t}W$, which tend to have higher jet multiplicity.

Event Selection

- Exactly three leptons with charge ± 1
 - Two tight Iso, tight ID same-charge leptons with $p_T > 20$ GeV
 - One loose Iso, medium ID opposite charge lepton with $p_T > 10$ GeV
 - $m(l^+l^-)$ within 10 GeV of 91.2 GeV
 - Transverse mass of W-candidate, $m_T(E_T^{\text{miss}} + \text{lep}_{\text{other}}) > 30$ GeV
 - Missing transverse energy, $E_T^{\text{miss}} > 20$ GeV
 - One or two jets with $p_T > 25$ GeV
-

Table 5: Summary of the selection applied to events for inclusion in the fit

261 The event yields in the preselection region for both data and Monte Carlo are summarized in
 262 Table 5.1, which shows good agreement between data and Monte Carlo, and demonstrates that
 263 this region consists primarily of WZ events. The WZ events are split into WZ + b, WZ + c, and
 264 WZ + l based on the truth flavor of the associated jet in the event. Specifically, this determination
 265 is made based on the `HadronConeExclTruthLabelID` of the jet, as recommended by the b-
 266 tagging working group [BtagWG]. In this ordering b-jet supersedes charm, which supersedes
 267 light. That is, WZ + l events contain no charm and no b jets at truth level, WZ + c contain at
 268 least one truth charm and no b-jets, and WZ + b contains at least one truth b-jet.

Process	Events
WZ + b	167.6 ± 6.5
WZ + c	1080 ± 40
WZ + l	7220 ± 310
Other VV	850 ± 140
t̄tW	16.8 ± 2.3
t̄tZ	115 ± 17
rare Top	2.2 ± 0.1
Single top	0.10 ± 0.45
Three top	0.01 ± 0.01
Four top	0.02 ± 0.01
t̄tWW	0.23 ± 0.05
Z + jets	600 ± 260
V + γ	37 ± 54
tZ	190 ± 70
tW	5.5 ± 1.2
WtZ	25.8 ± 1.1
VVV	26.2 ± 0.9
VH	94 ± 7
t̄t	108.68 ± 8
t̄tH	4.3 ± 0.5
Total	10600 ± 530
Data	10574

Table 6: Event yields in the preselection region at 139.0 fb^{-1}

²⁶⁹ Here Other VV represents diboson processes other than WZ, and consists predominantly of
²⁷⁰ ZZ → ll̄ll events where one of the leptons is not reconstructed.

²⁷¹ Simulations are further validated by comparing the kinematic distributions of the Monte Carlo
²⁷² with data, which are shown in Figure 2. Here, bins with 5% or more WZ+b are blinded.

WZ Fit Region - Inclusive

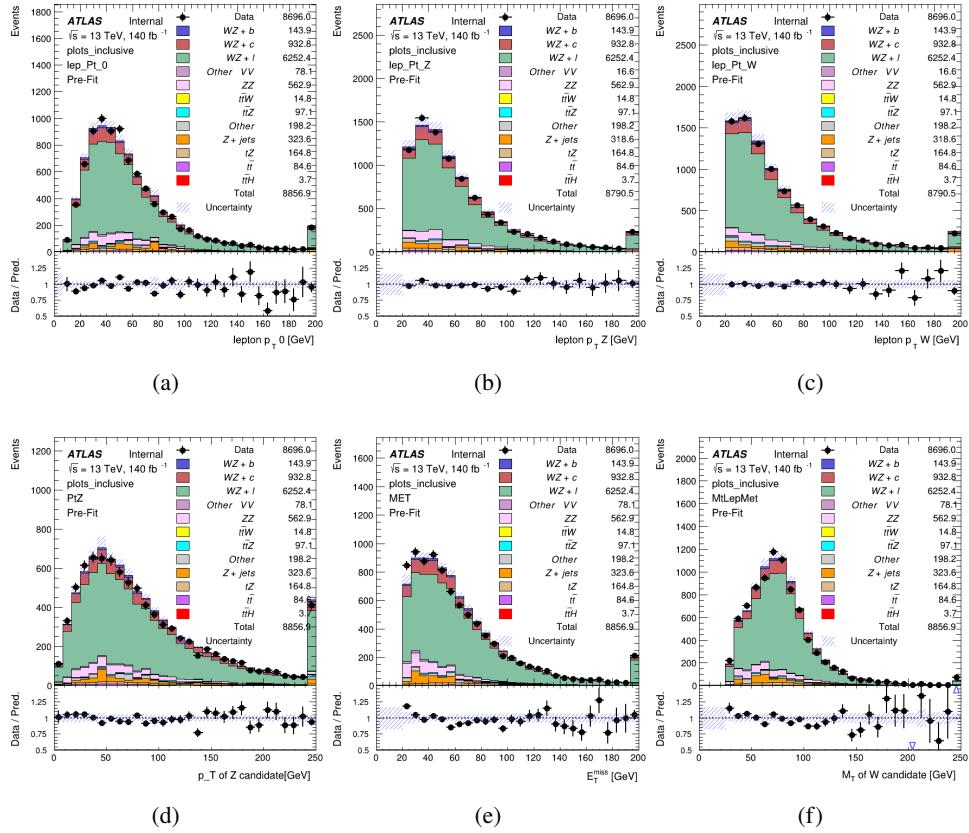


Figure 2: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

5.2 Fit Regions

Once preselection has been applied, the remaining events are categorized into one of twelve orthogonal regions. The regions used in the fit are summarized in Table 7.

Table 7: A list of the regions used in the fit and the selection used for each.

Region	Selection
1j, <85%	$N_{\text{jets}} = 1, n\text{Jets_DL1r_85} = 0$
1j, 85%-77%	$N_{\text{jets}} = 1, n\text{Jets_DL1r_85} = 1, n\text{Jets_DL1r_77}=0$
1j, 77%-70%	$N_{\text{jets}} = 1, n\text{Jets_DL1r_77} = 1, n\text{Jets_DL1r_70}=0$
1j, 70%-60%	$N_{\text{jets}} = 1, n\text{Jets_DL1r_70} = 1, n\text{Jets_DL1r_60}=0$
1j, >60%	$N_{\text{jets}} = 1, n\text{Jets_DL1r_60} = 1, tZ \text{ BDT} > 0.725$
1j tZ CR	$N_{\text{jets}} = 1, n\text{Jets_DL1r_60} = 1, tZ \text{ BDT} < 0.725$
2j, <85%	$N_{\text{jets}} = 2, n\text{Jets_DL1r_85} = 0$
2j, 85%-77%	$N_{\text{jets}} = 2, n\text{Jets_DL1r_85} >= 1, n\text{Jets_DL1r_77}=0$
2j, 77%-70%	$N_{\text{jets}} = 2, n\text{Jets_DL1r_77} >= 1, n\text{Jets_DL1r_70}=0$
2j, 70%-60%	$N_{\text{jets}} = 2, n\text{Jets_DL1r_70} >= 1, n\text{Jets_DL1r_60}=0$
2j, >60%	$N_{\text{jets}} = 2, n\text{Jets_DL1r_60} >= 1, tZ \text{ BDT} > 0.725$
2j tZ CR	$N_{\text{jets}} = 2, n\text{Jets_DL1r_60} >= 1, tZ \text{ BDT} < 0.725$

276 The working points discussed in Section 4.4 are used to separate events into fit regions based on
 277 the highest working point reached by a jet in each event. Because the background composition
 278 differs significantly based on the number of b-jets, events are further subdivided into 1 jet and 2
 279 jet regions in order to minimize the impact of background uncertainties.

280 An additional tZ control region is created based on the BDT described in Section 6. The region
 281 with 1-jet passing the 60% working point is split in two - a signal enriched region of events with
 282 a BDT score greater than 0.03, and a tZ control region including events with less than 0.03. This
 283 cutoff is arrived at by performing an Asimov fit with a variety of cutoffs, and selecting the value
 284 that produces the highest significance for the measurement of WZ + b.

285 The modeling in each region is validated by comparing data and MC predictions for various
 286 kinematic distributions. These plot are shown in Figures 3-16.

WZ Fit Region - 1j Inclusive

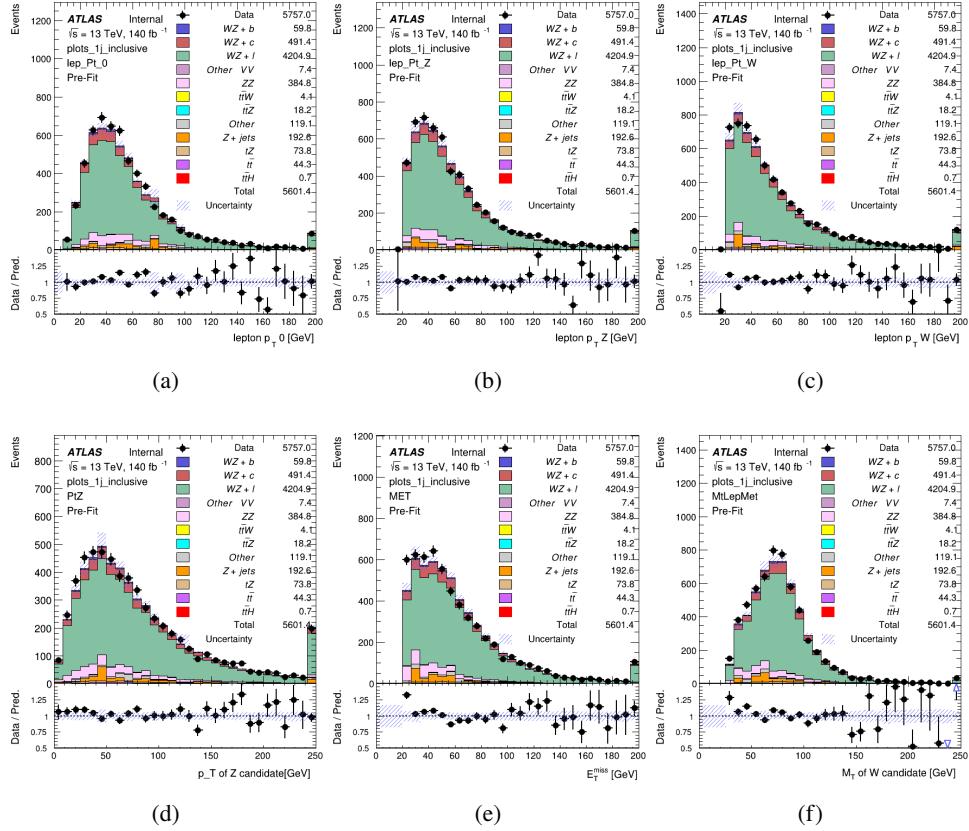


Figure 3: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 1j < 85% WP

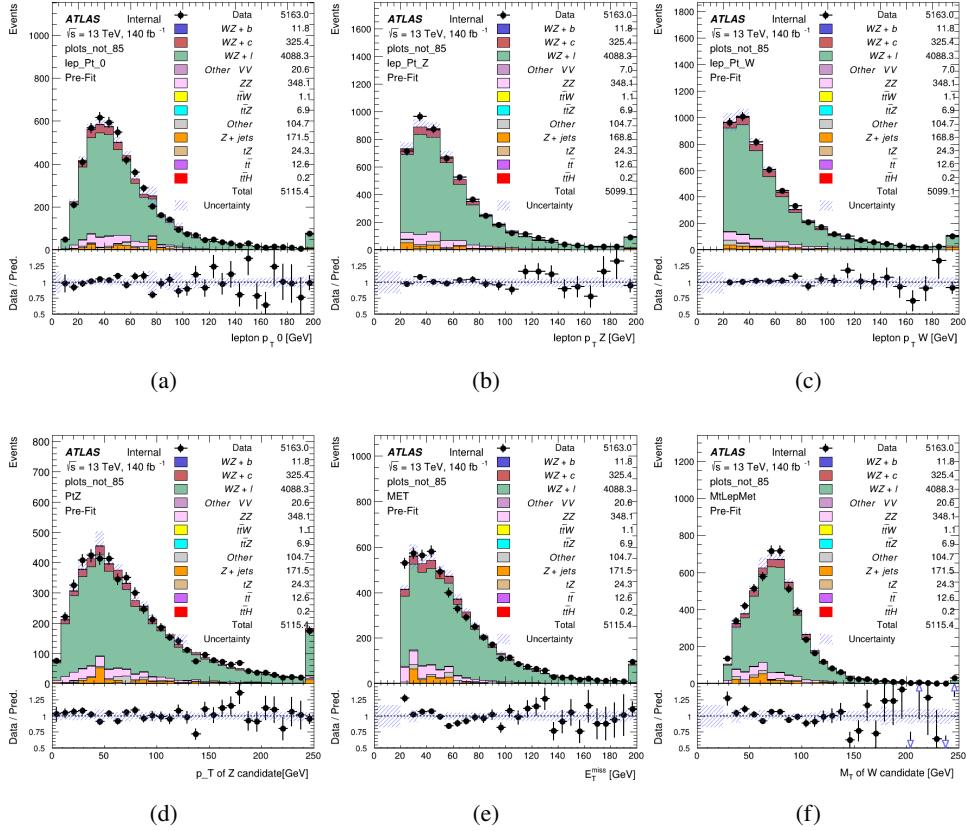


Figure 4: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 1j 77-85% WP

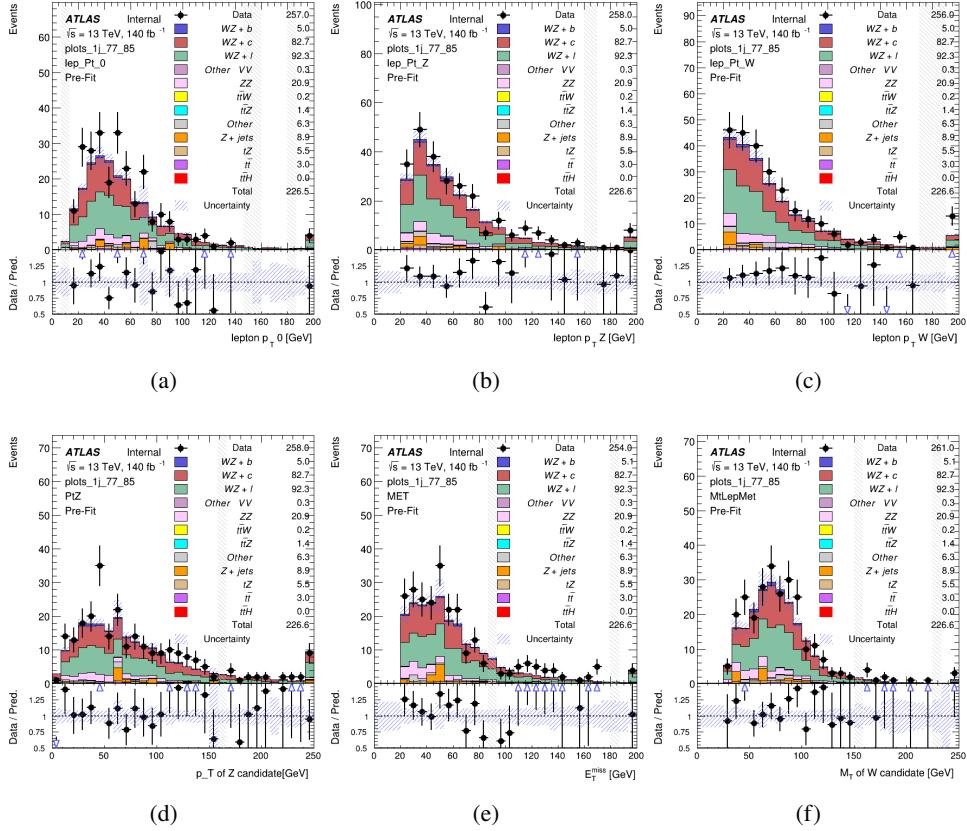


Figure 5: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 1j 70-77% WP

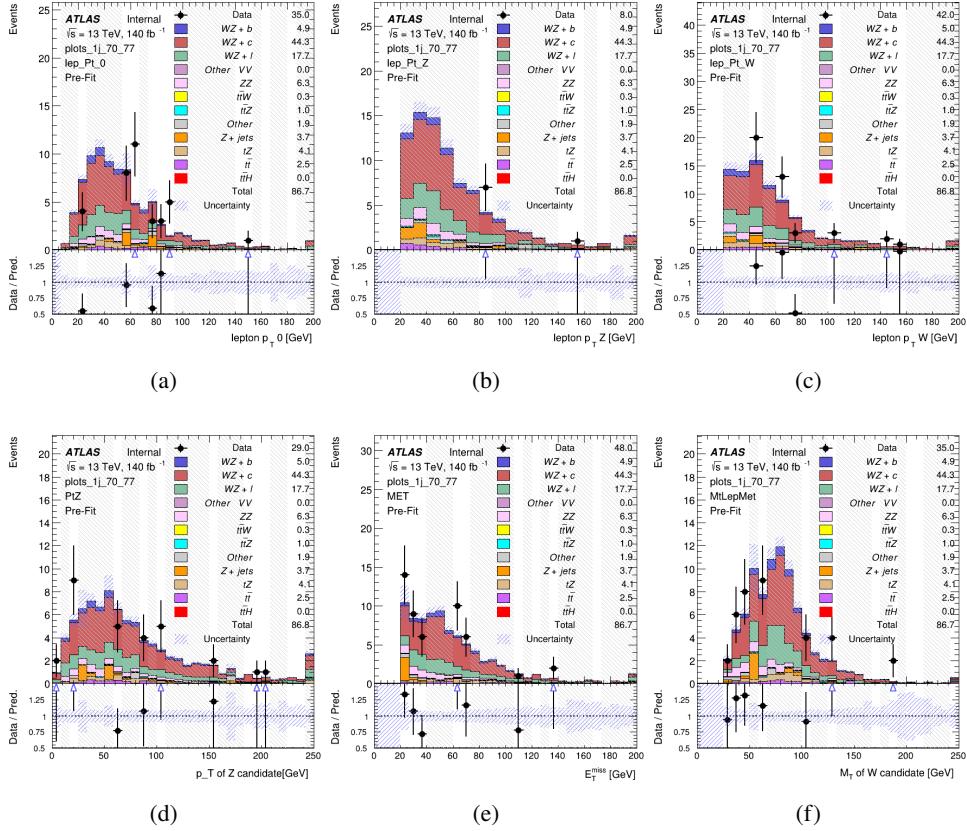


Figure 6: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 1j 60-70% WP

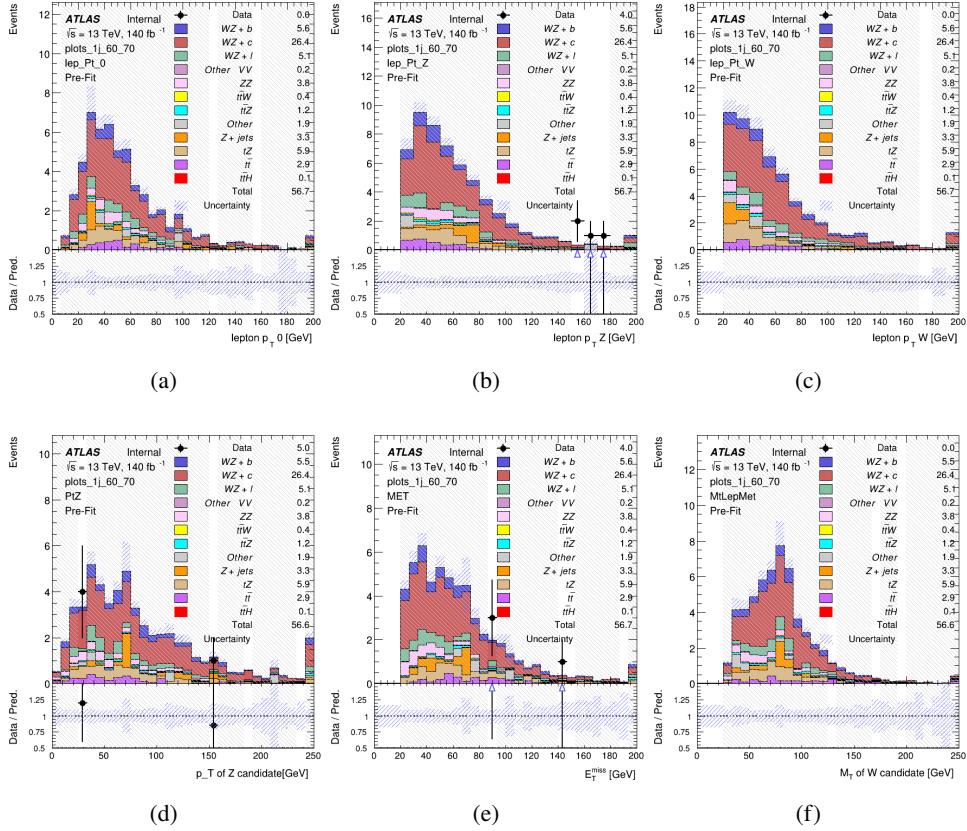


Figure 7: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 1j 60% WP

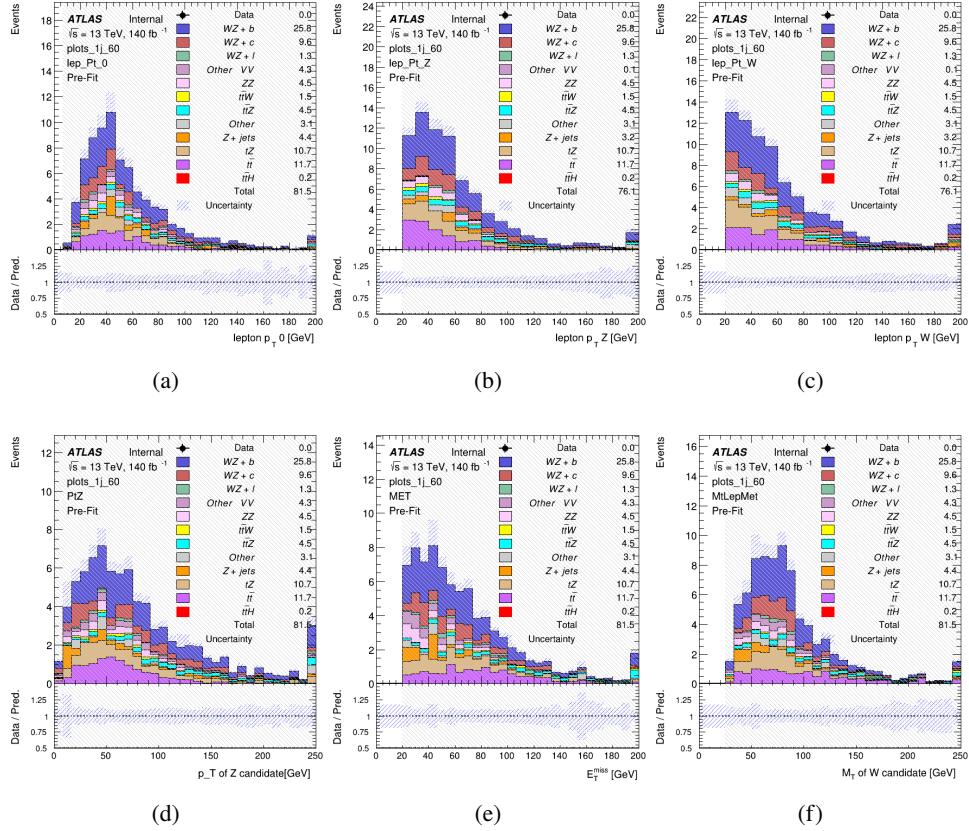


Figure 8: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - tZ-CR

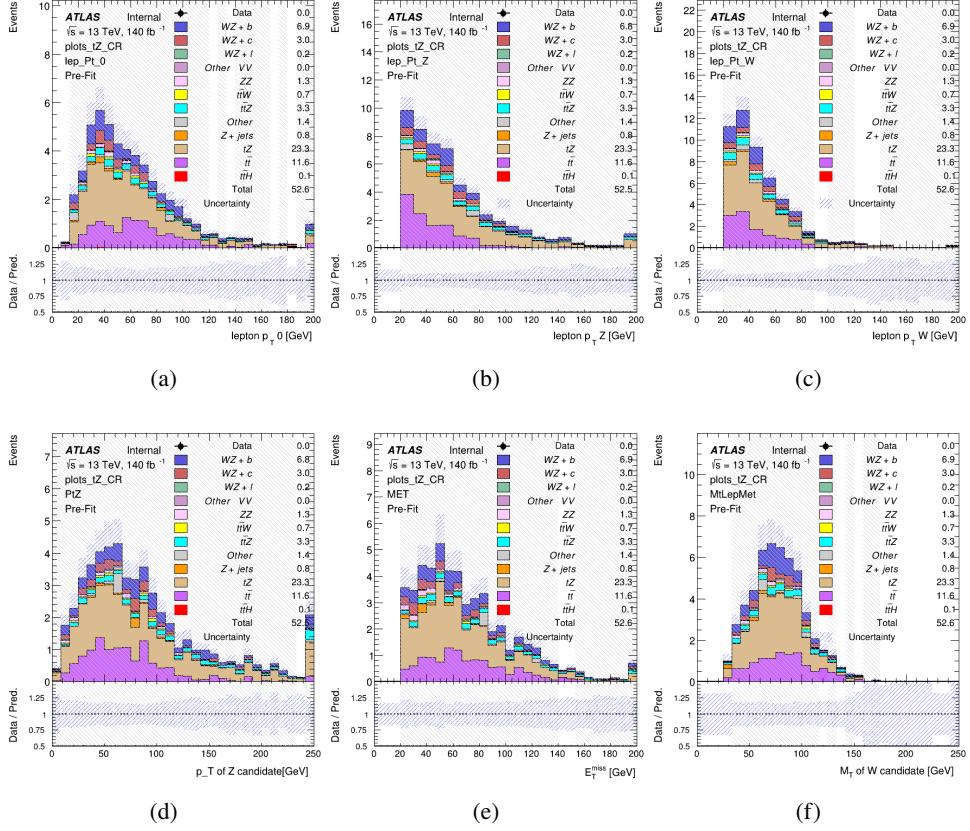


Figure 9: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 2j Inclusive

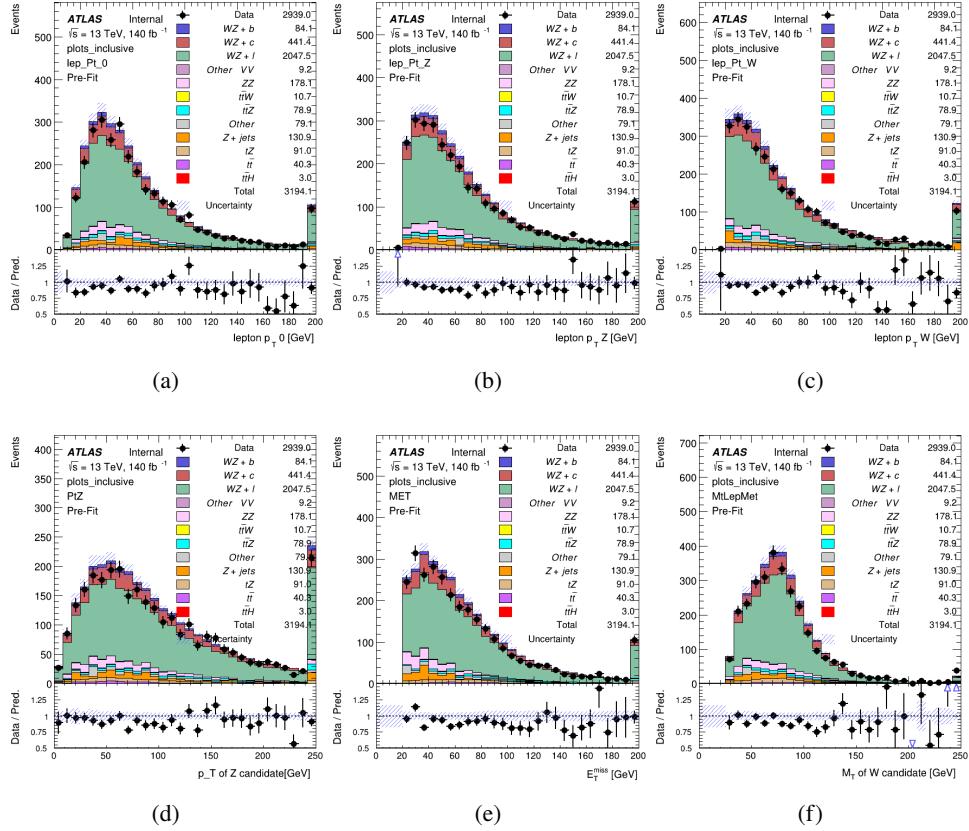


Figure 10: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 2j < 85% WP

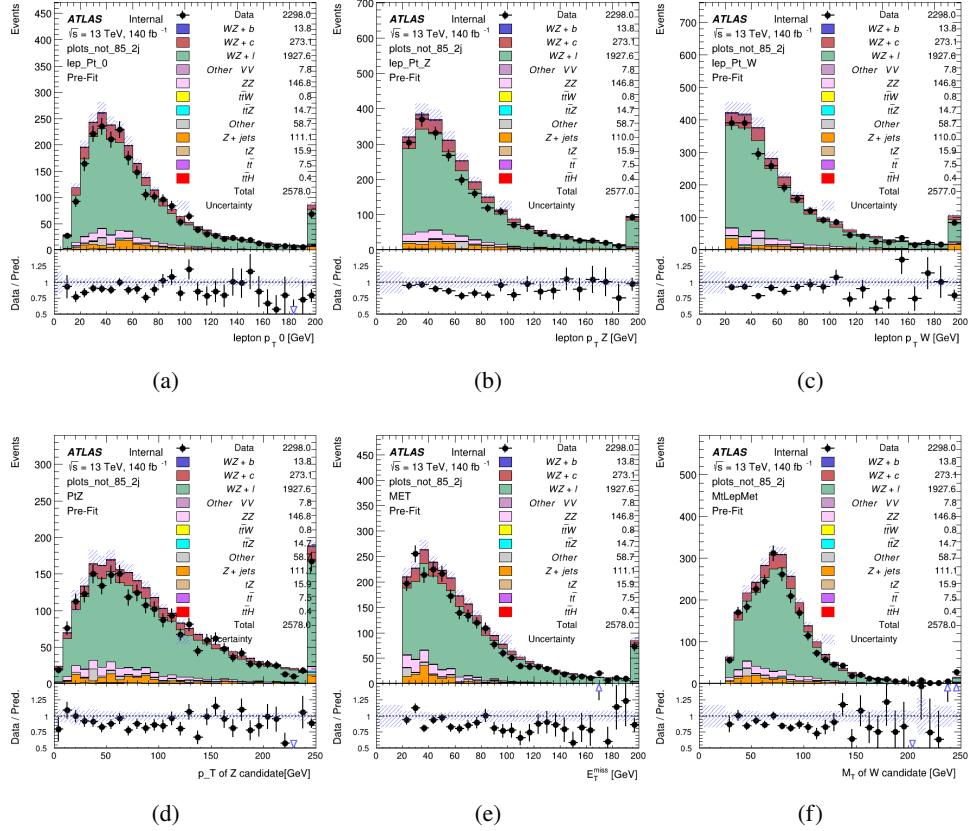


Figure 11: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the M_T of the W candidate.

WZ Fit Region - 2j 77-85% WP

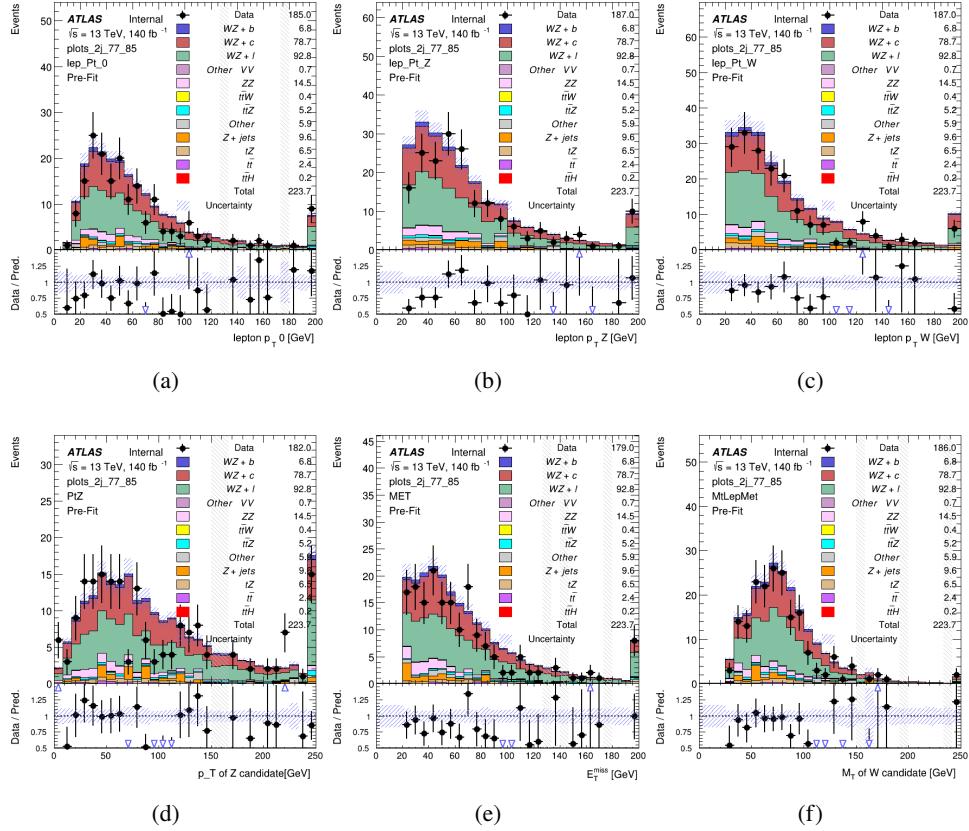


Figure 12: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 2j 70-77% WP

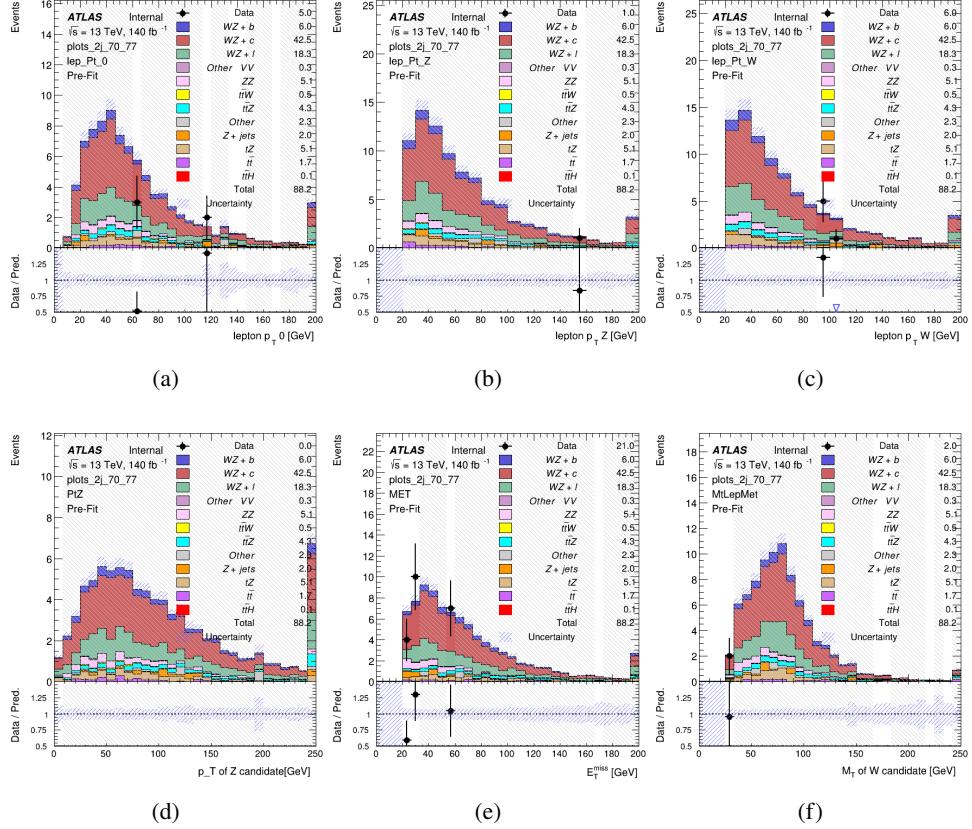


Figure 13: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 2j 60-70% WP

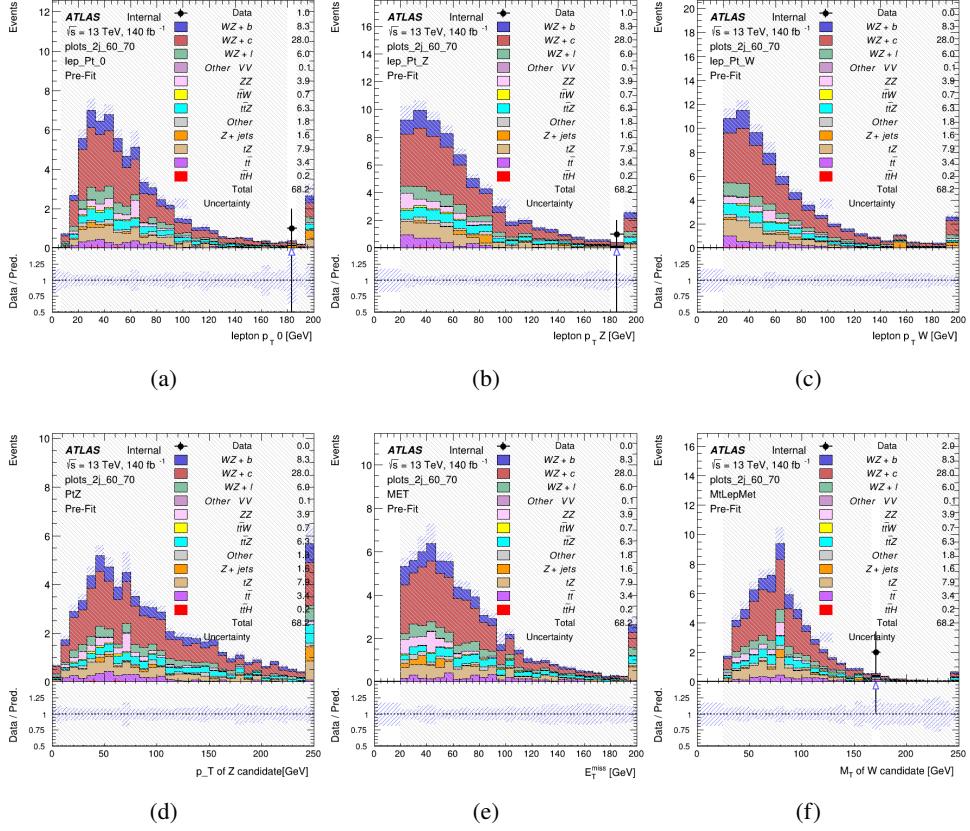


Figure 14: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

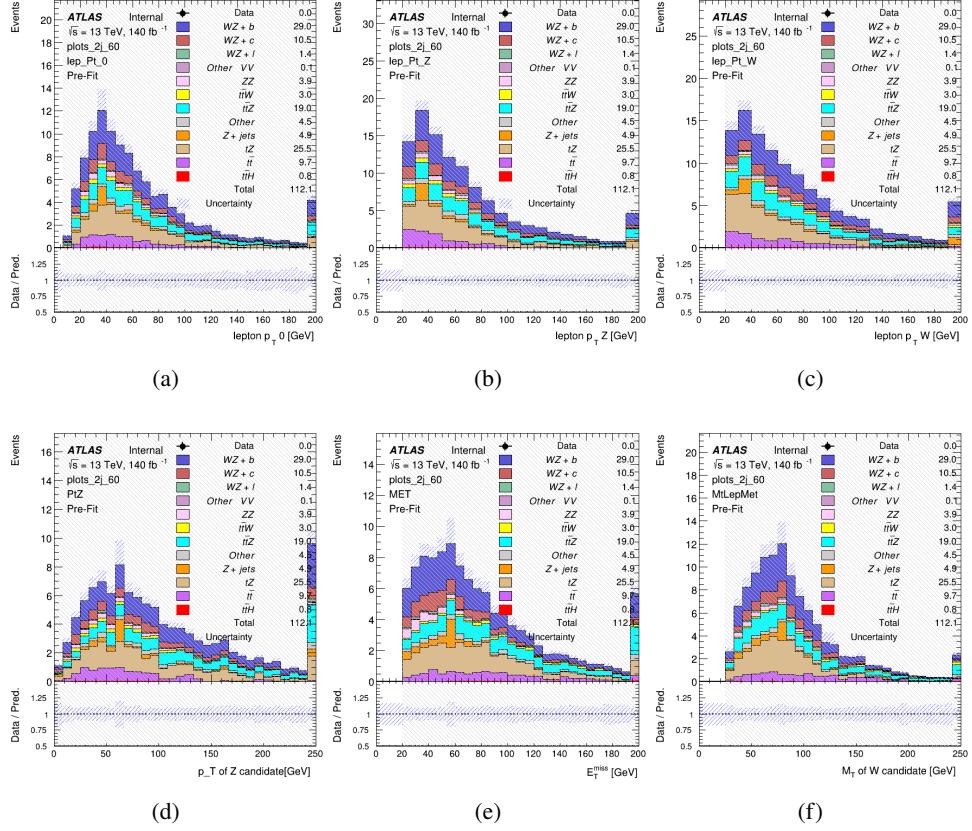
WZ Fit Region - 2j 60% WP

Figure 15: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - tZ-CR-2j

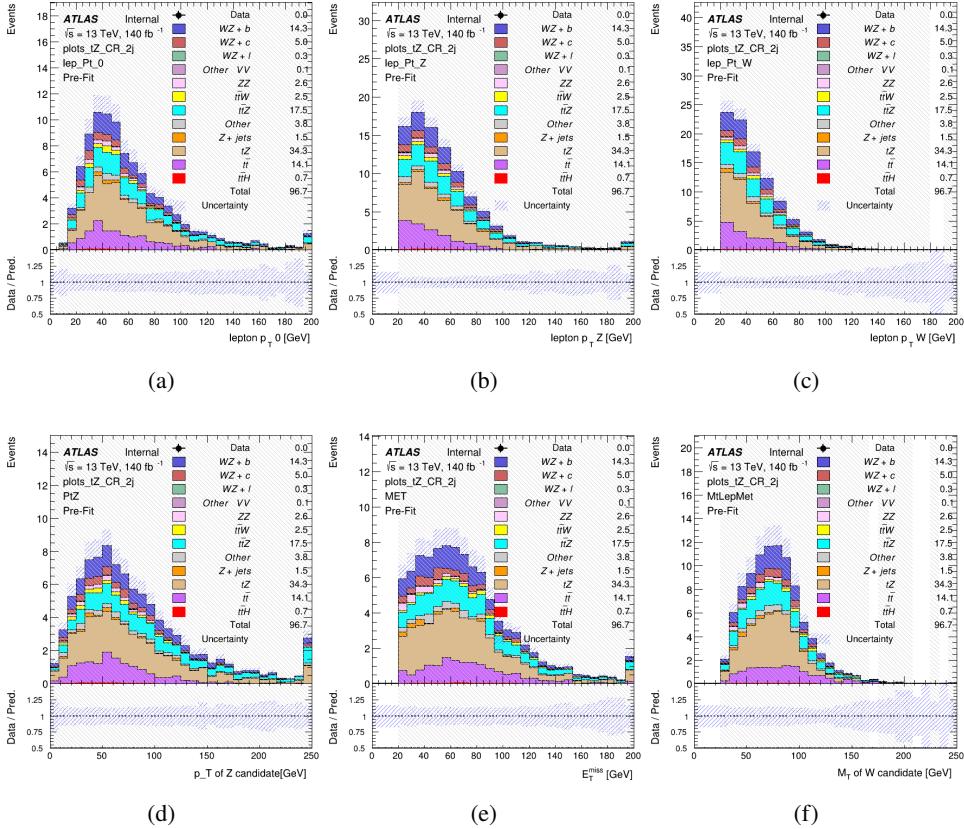


Figure 16: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

5.3 Non-Prompt Lepton Estimation

Two processes act as sources of non-prompt leptons appear in the analysis: $t\bar{t}$ and $Z+jet$ production both produce two prompt leptons, and each contribute to the 31 region when an additional non-prompt lepton appears in the event. The contribution of these processes is estimated with Monte Carlo simulations, which are validated using enriched validation regions.

5.3.1 $t\bar{t}$ Validation

$t\bar{t}$ events can produce two prompt leptons from the decay of each of the tops. These top decays produce two b-quarks, the decay of which can produce additional non-prompt leptons, which occasionally pass the event preselection. In order to validate that the Monte Carlo accurately

296 simulates this process accurately, the MC prediction in a non-prompt $t\bar{t}$ enriched validation
297 region is compared to data.

298 The $t\bar{t}$ validation region is similar to the preselection region - three leptons meeting the criteria
299 described in Section 5 are required, and the requirements on E_T^{miss} remain the same. However,
300 the selection requiring a lepton pair form a Z-candidate are reversed. Events where the invariant
301 mass of any two opposite sign, same flavor leptons falls within 10 GeV of 91.2 GeV are rejected.
302 This ensures the $t\bar{t}$ validation region is orthogonal to the preselection region.

303 Further, because the jet multiplicity of $t\bar{t}$ events tends to be higher than WZ, the number of jets
304 in each event is required to be greater than 1. As b-jets are almost invariably produced from top
305 decays, at least one b-tagged jet passing the 70% DL1r WP in each event is required. Various
306 kinematic plots of this region are shown in Figure 17.

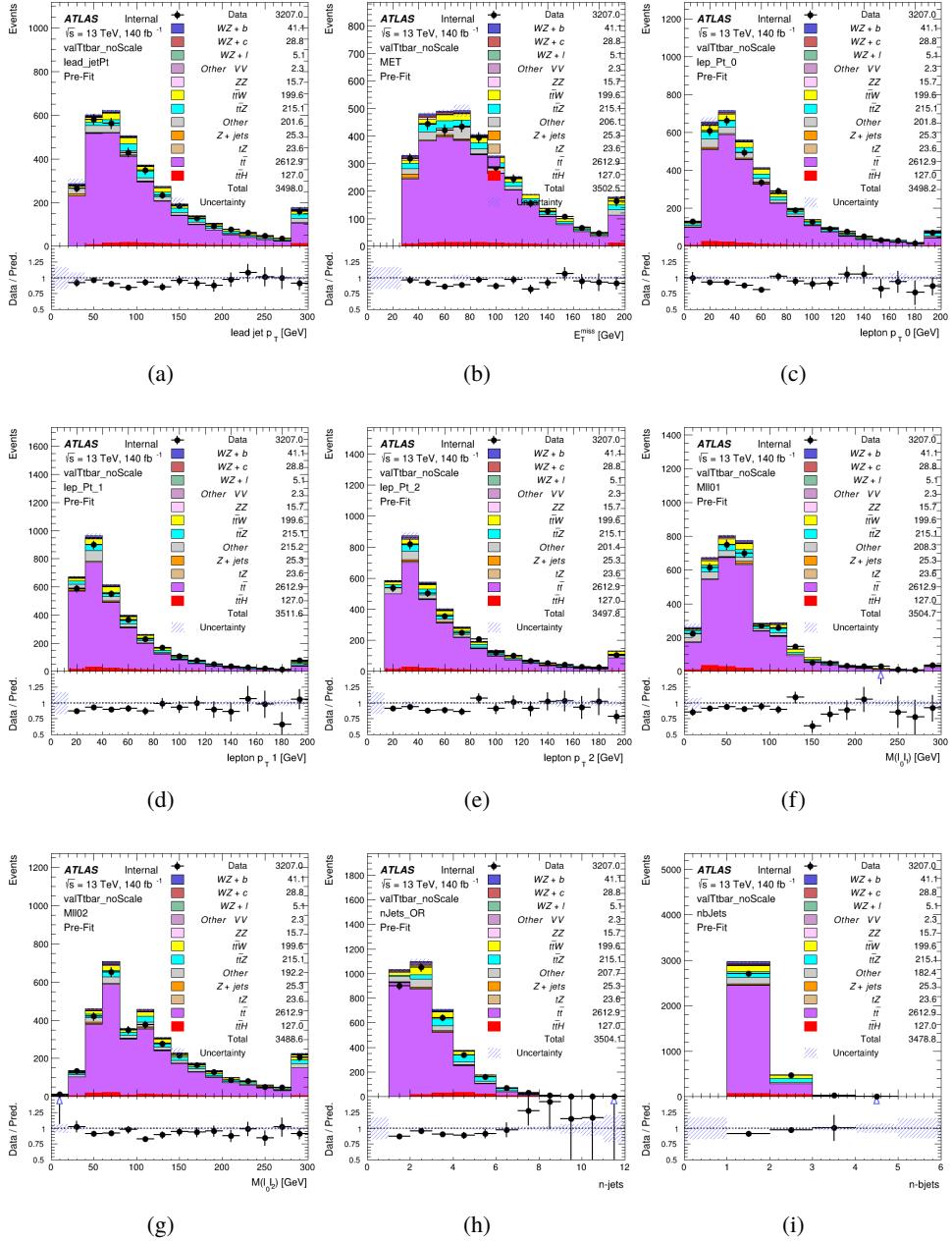


Figure 17: Comparisons between the data and MC distributions in the $t\bar{t}$ validation region for (a) the p_T of the leading jet, (b) the missing transverse energy, (c) the p_T of lepton 0, (d) p_T of lepton 1, (e) p_T of lepton 2, (f) the invariant mass of leptons 0 and 1, (g) the invariant mass of leptons 0 and 2, (h) the number of jets, (i) the number of b-tagged jets.

307 The shape of each distribution agrees quite well between data and MC, with a constant offset
 308 between the two. This is accounted for by applying a constant correction factor of 0.883 to the $t\bar{t}$

309 MC prediction. Plots showing the kinematics of the $t\bar{t}$ VR after this correction factor has been
310 applied are shown in Figure 18.

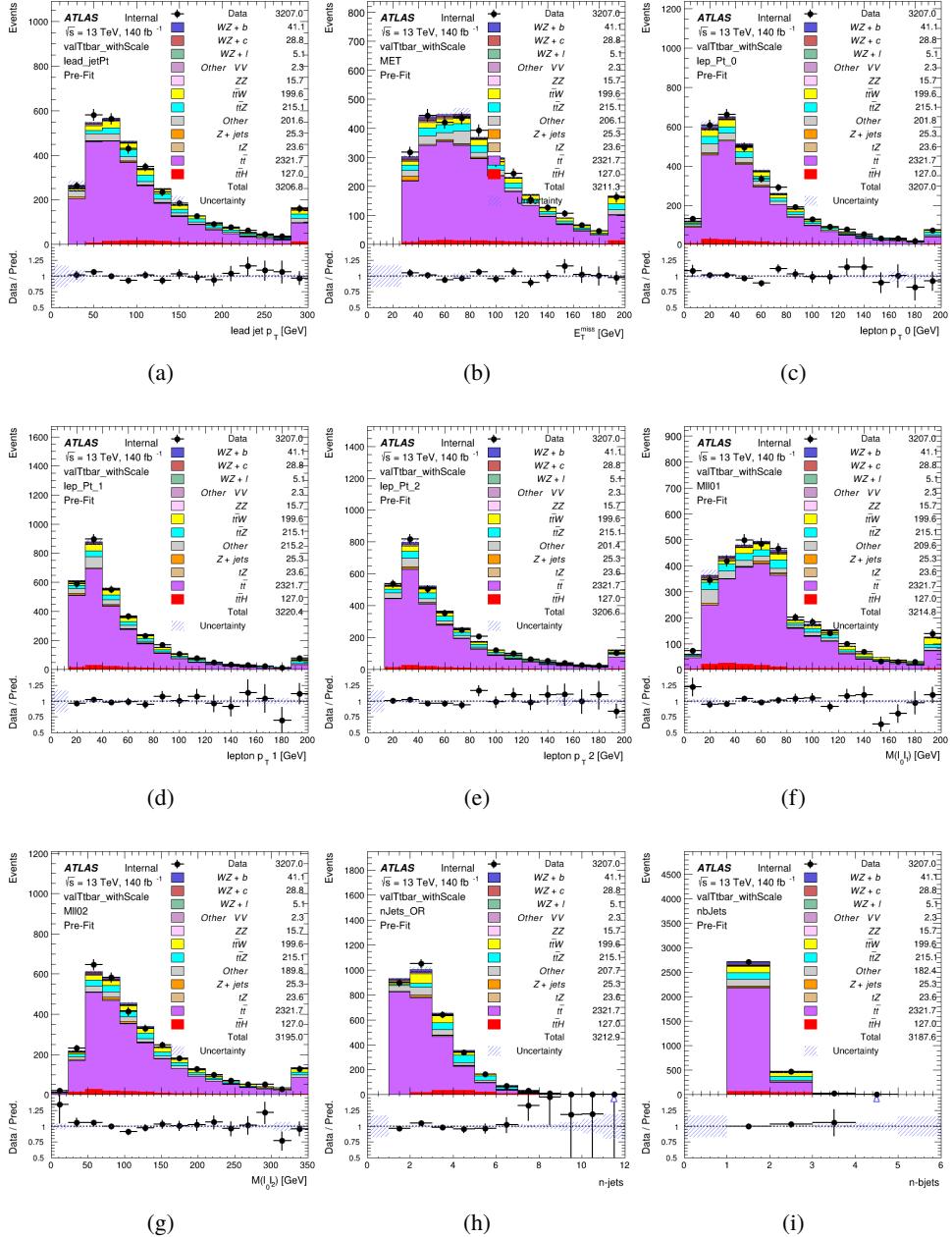


Figure 18: Comparisons between the data and MC distributions in the $t\bar{t}$ validation region after the correction factor has been applied for (a) the p_T of the leading jet, (b) the missing transverse energy, (c) the p_T of lepton 0, (d) p_T of lepton 1, (e) p_T of lepton 2, (f) the invariant mass of leptons 0 and 1, (g) the invariant mass of leptons 0 and 2, (h) the number of jets, (i) the number of b-tagged jets.

311 The modeling is further validated by looking at the yield in the $t\bar{t}$ VR for each DL1r WP, giving
 312 a clearer correspondence to the signal regions used in the fit. Each region shown in Figure 22
 313 requires one or more jets pass the listed WP, with no jets passing the next highest WP.

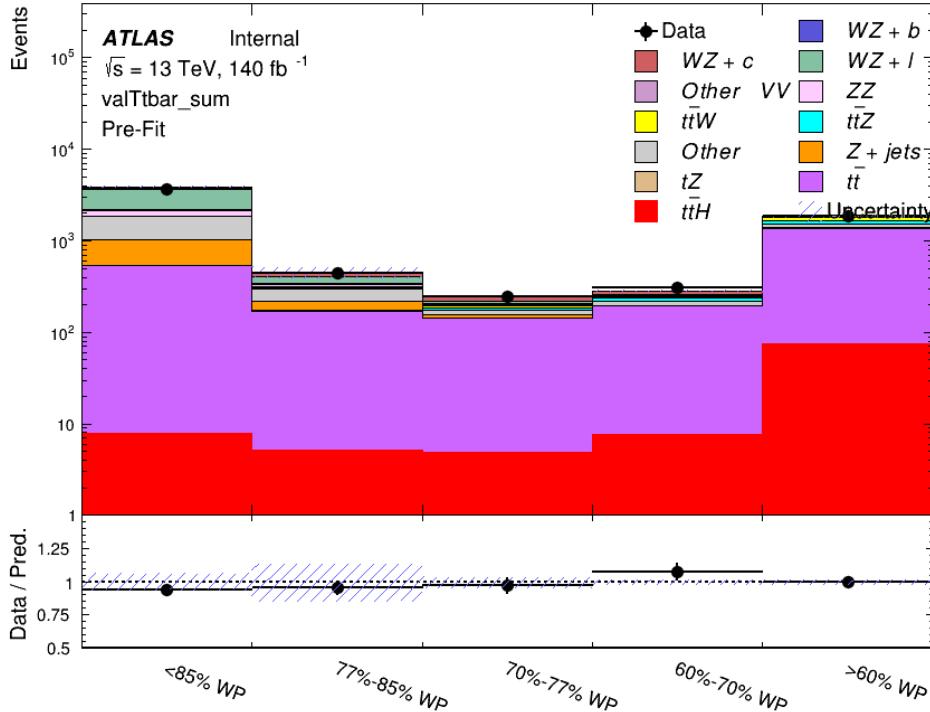


Figure 19: Data and MC comparisons for each DL1r WP for both 1-jet and 2-jet regions, after the $t\bar{t}$ VR selection and correction factor have been applied

314 As data and MC are found to agree within 10% for each of these working points, a 10% systematic
 315 uncertainty on the $t\bar{t}$ prediction is included for the analysis.

316 5.3.2 Z+jets Validation

317 Similar to $t\bar{t}$, a non-prompt Z+jets validation region is produced in order to validate the MC
 318 predictions. The lepton requirements remain the same as the preselection region. Because no
 319 neutrinos are present for this process, the E_T^{miss} cut is reversed, requiring $E_T^{\text{miss}} < 30 \text{ GeV}$. This
 320 also ensures this validation region is orthogonal to the preselection region. Further, the number
 321 of jets in each event is required to be greater than or equal to one. Various kinematic plots of this
 322 region are shown below. The general agreement between data and MC in each of these suggests
 323 that the non-prompt contribution of Z+jets is well modeled by Monte Carlo.

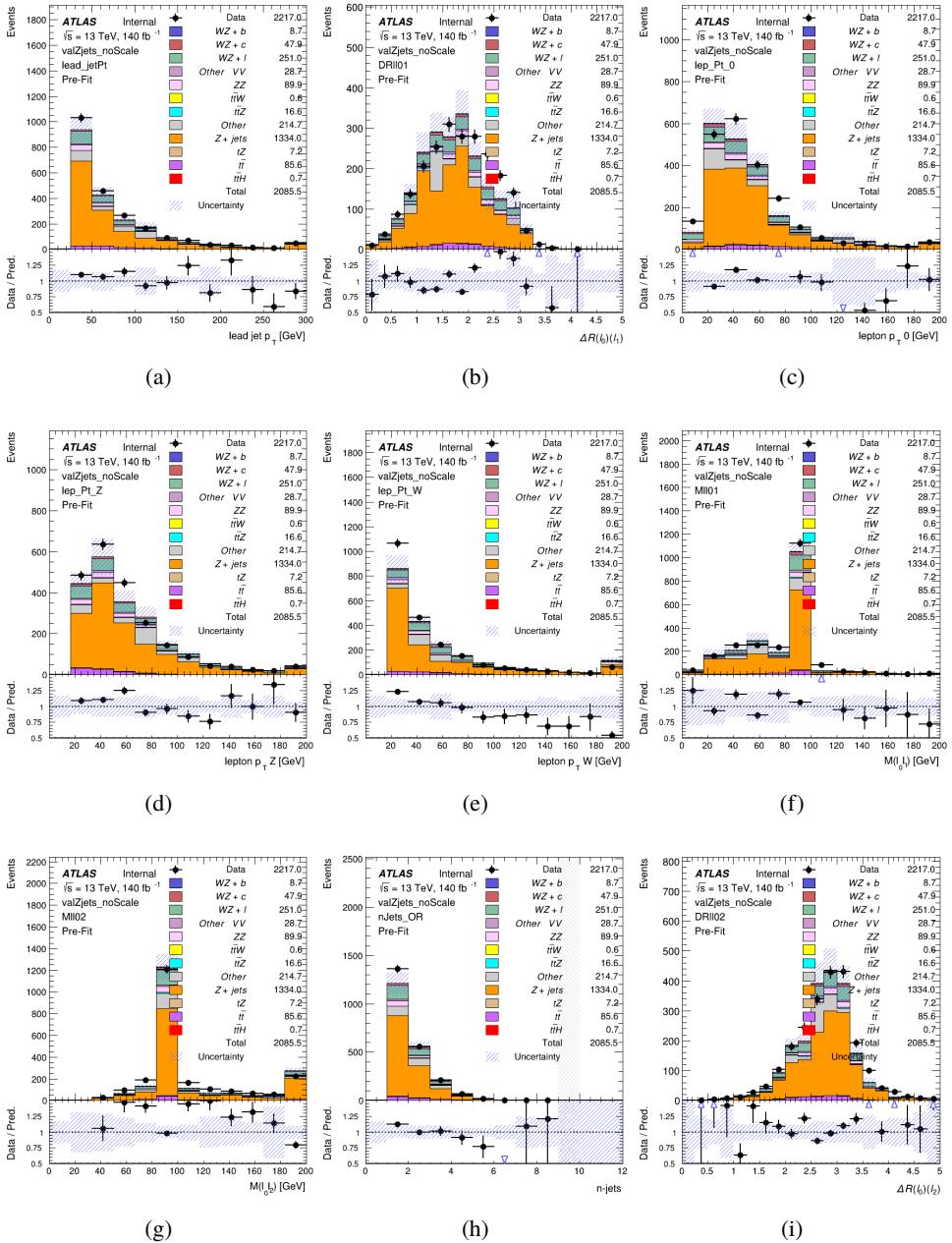


Figure 20: Comparisons between the data and MC distributions in the Z+jets validation region for (a) the p_T of the leading jet, (b) ΔR between leptons 0 and 1, (c) the p_T of lepton 0, (d) p_T of SS lepton from the Z candidate, (e) p_T of the SS lepton from the W candidate, (f) the invariant mass of leptons 0 and 1, (g) the invariant mass of leptons 0 and 2, (h) the number of jets, (i) ΔR between leptons 0 and 2. Includes only statistical uncertainties

324 While there is general agreement between data and MC within statistical uncertainty, the shape

325 of the p_T spectrum of the lepton from the W is found to differ. To account for this discrepancy,
326 a variable correction factor is applied to $Z+jets$. χ^2 minimization of the lepton 2 p_T spectrum is
327 performed to derive a correction factor of $1.53 - 6.6 * 10^{-6}(lep_Pt_W)$. Kinematic plots of the
328 $Z + jets$ validation region after this correction factor has been applied are shown in Figure 21.

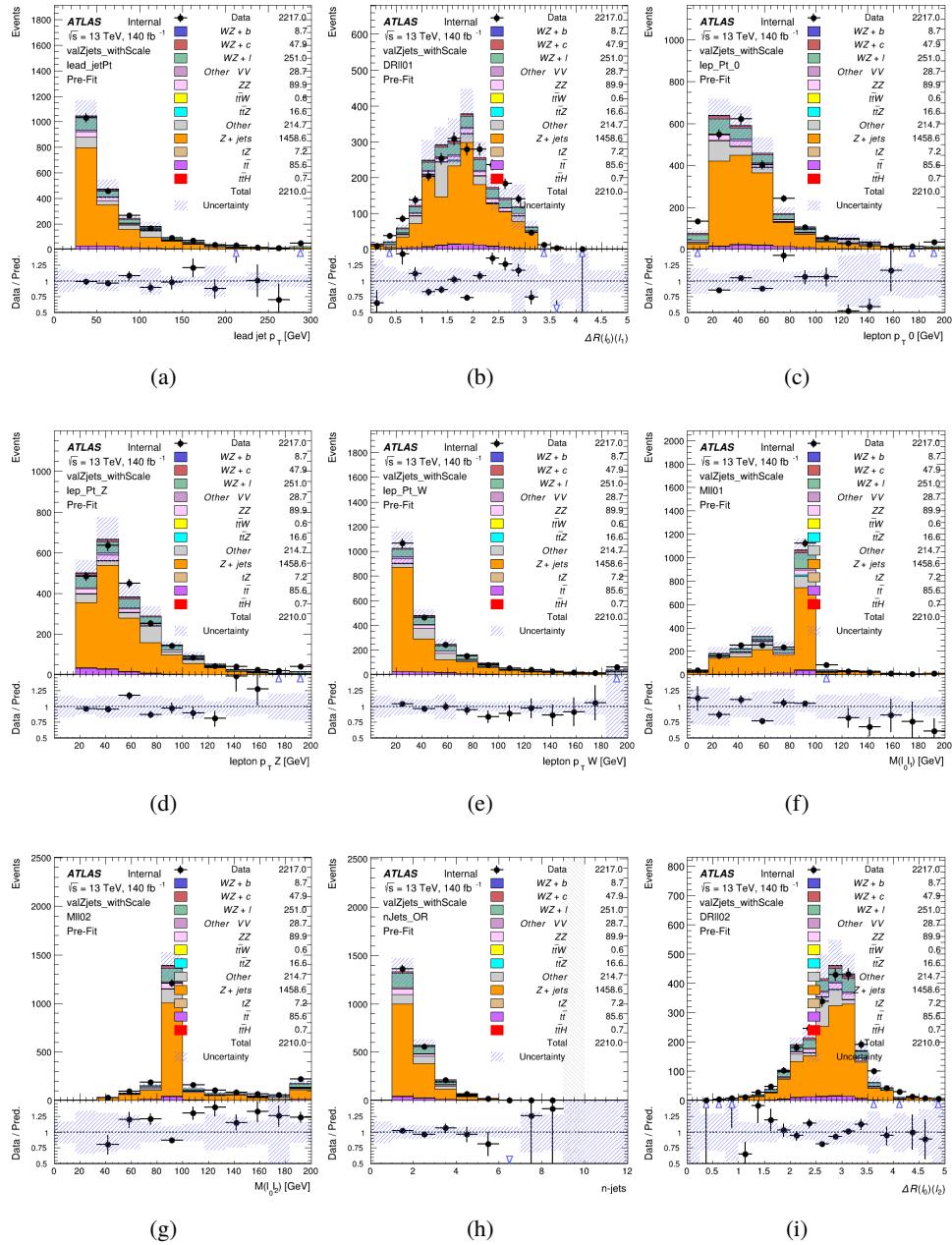


Figure 21: Comparisons between the data and MC distributions in the Z+jets validation region after the correction factor has been applied for (a) the p_T of the leading jet, (b) ΔR between leptons 0 and 1, (c) the p_T of lepton 0, (d) p_T of SS lepton from the Z candidate, (e) p_T of the SS lepton from the W candidate, (f) the invariant mass of leptons 0 and 1, (g) the invariant mass of leptons 0 and 2, (h) the number of jets, (i) ΔR between leptons 0 and 2

329 The modeling is further validated by looking at the yield in the Z+jets VR for each DL1r WP,

330 giving a clearer correspondence to the signal regions used in the fit. Each region shown in Figure
 331 22 requires one or more jets pass the listed WP, with no jets passing the next highest WP.

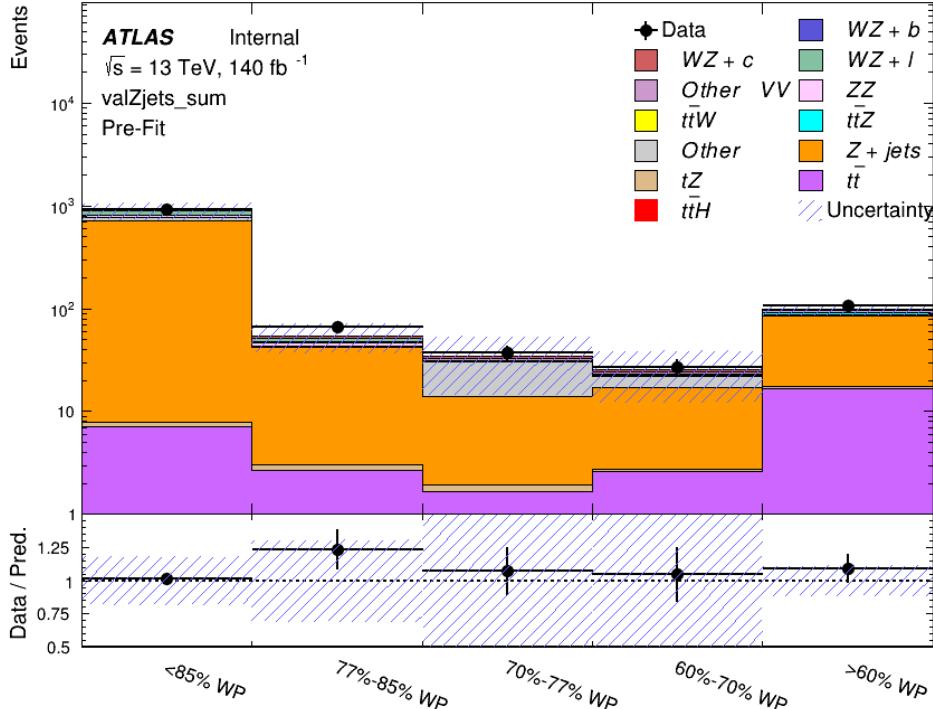


Figure 22: Data and MC comparisons for each DL1r WP for both 1-jet and 2-jet regions, after the Z+jets VR selection and correction factor have been applied

332 For each of the b-tagging working points considered, the data falls within 20% of the MC
 333 prediction once this correction factor has been applied. Therefore, a 20% systematic uncertainty
 334 is applied to Z + jets in the analysis.

335 6 tZ Separation Multivariate Analysis

336 Because tZ produces a final state identical to signal, it represents a predominant background in
 337 the most signal enriched regions. That is, the region with one jet passing the 60% DL1r WP.
 338 Therefore, a boosted decision tree (BDT) algorithm is trained using TMVA [10] to separate WZ
 339 + heavy flavor from tZ.

340 Separation between tZ and WZ + heavy flavor is achieved in part by reconstructing the invariant
 341 mass of the top candidate, which clusters more closely to the top mass for tZ than WZ + heavy
 342 flavor.

343 The result of this BDT is used to create a tZ enriched region in the fit, reducing its impact on the
 344 measurement of WZ + heavy flavor.

345 **6.1 Top Mass Reconstruction**

346 The reconstruction of the top mass follows the procedure described in detail in section 6.1 of
 347 [11]. The mass of the top quark candidate is reconstructed from the jet, the lepton not included
 348 in the Z-candidate, and a reconstructed neutrino. In the case that there is one jet in the event,
 349 there is only possible b-jet candidate. For events with two jets, the jet with the highest DL1r
 350 score is used.

351 The neutrino from the W decay is expected to be the only source of E_T^{miss} . Therefore, the E_T
 352 and ϕ of the neutrino are taken from the E_T^{miss} measurement. This leaves the z-component of
 353 the neutrino momentum, $p_{\nu z}$ as the only unknown.

354 This unknown is solved for by taking the combined invariant mass of the lepton and neutrino to
 355 give the invariant mass of the W boson:

$$356 \quad (p_l + p_\nu)^2 = m_W^2$$

357 Expanding this out into components, this equation gives:

$$358 \quad \sqrt{p_{T\nu}^2 + p_{z\nu}^2} E_l = \frac{m_w^2 - m_l^2}{2} + p_{T\nu}(p_{lx}\cos\phi_\nu + p_{ly}\sin\phi_\nu) + p_{lz}p_{\nu z}$$

359 This equation gives two solutions for $p_{\nu z}$. For cases where only one of these solutions is real,
 360 that is taken as the value of $p_{\nu z}$. For instances with two real solutions, the one which is shown
 361 to be correct in the largest fraction of simulations is taken. For cases when no real solution is
 362 found, often because of detector effects, the value of E_T^{miss} is varied in decreasing increments of
 363 100 MeV until a real solution is found.

364 The reconstructed top mass distribution for tZ and WZ + b can be seen in figure 23.

365 **6.2 tZ BDT**

366 A Boosted Decision Tree (BDT), specifically XGBoost [xgboost_cite], is used to provide separa-
 367 tion between tZ and WZ+b. The following kinematic variables are used as inputs:

- 368 • The invariant mass of the reconstructed top candidate
- 369 • p_T of each of the leptons, jet
- 370 • The invariant mass of each combination of lepton pairs, $M(l\bar{l})$
- 371 • E_T^{miss}

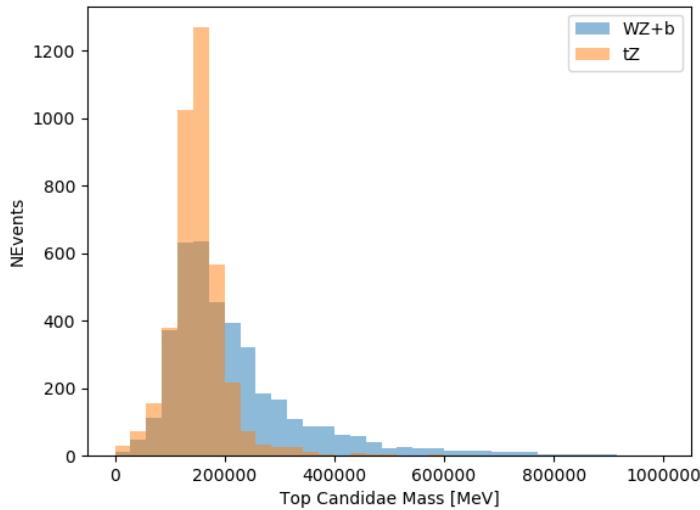


Figure 23: Reconstructed top mass distributions for tZ and WZ + b, measured in MeV.

- 372 • Distance between each combination of leptons, $\Delta R(ll)$
- 373 • Distance between each lepton and the jet, $\Delta R(lj)$

374 The training samples included only events meeting the requirements of the 1-jet, >60% region,
 375 i.e. passing all the selection described in section 5 and having exactly one jet which passes the
 376 tightest (60%) DL1r working point.

377 The distributions of a few of these features for both signal and background is shown in figure
 378 24.

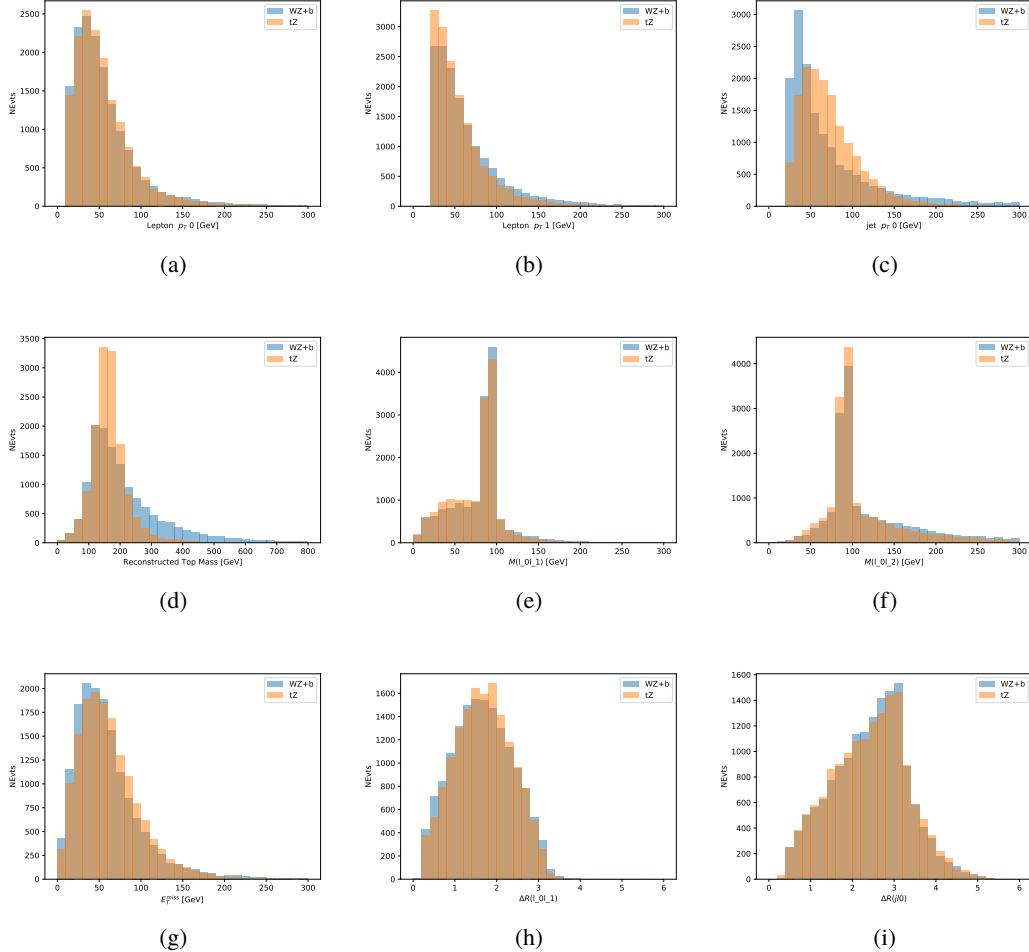


Figure 24: Distribution of input features of the BDT for signal (WZ) and background (tZ). Both are scaled to an equal number of events. (a), (b) and (c) show the p_T of lepton 0, lepton 1, and the jet, (d) show the reconstructed top mass, (e) and (f) show the invariant mass of leptons 0 and 1, leptons 0 and 2. (g) shows the E_T^{miss} of each event. (h) and (i) show the ΔR between lepton 0 and lepton 1, and the jet.

379 A sample of 20,000 background (tZ) and signal (WZ+b) Monte Carlo events are used to train
 380 the BDT. And additional 5,000 events are reserved for testing the model, in order to prevent
 381 over-fitting. A total of 750 decision trees with a maximum depth of 6 branches are used to build
 382 the model. These parameters are chosen empirically, by training several models with different
 383 parameters and selecting the one that gave the best separation for the test sample.

384 The results of the BDT training are shown in figure 25. The output scores for both signal and
 385 background events is shown on the left. The right shows the receiving operating characteristic
 386 (ROC) curve that results from the MVA. The ROC curve represents the background rejection
 387 as a function of signal efficiency, where each point on the curve represents a different response

388 score. The ROC curve of the BDT is compared to the performance of using an optimal set of flat
 389 selections on the same set of input variables.

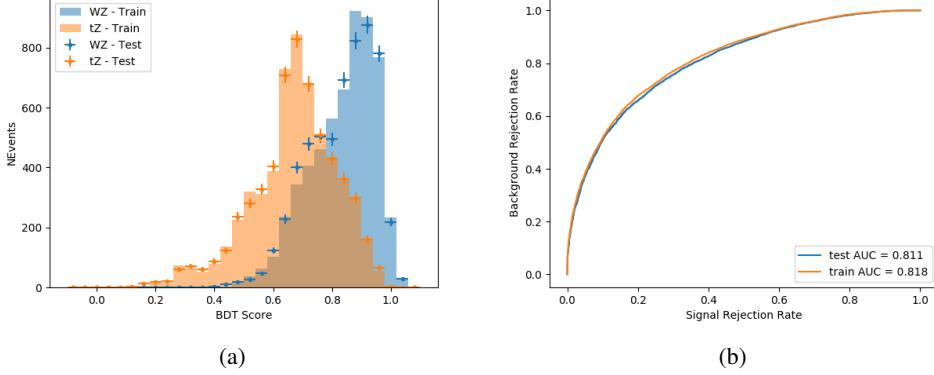


Figure 25: Distribution of the BDT response for signal and background events on the left, the ROC curve for the BDT on the right.

390 The relative important of each input feature in the model, measured by how often they appeared
 391 in the decision trees, is shown in figure 26.

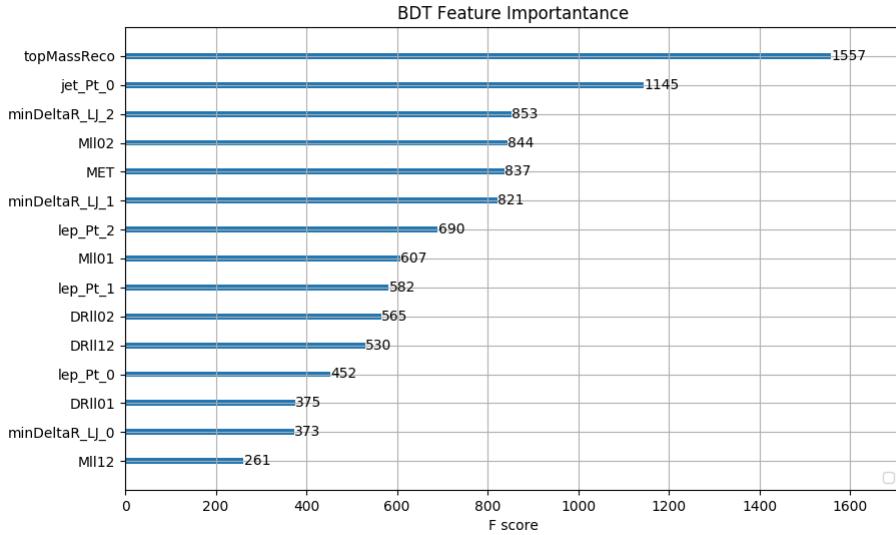


Figure 26: Relative importance of each input feature in the model.

392 These results suggest that some amount of separation can be achieved between these two pro-
 393 cesses, with a high BDT score selecting a set of events that is pure in $WZ + b$. A BDT score

394 of 0.725 is selected as a cutoff, where events with scores higher than this form a signal enriched
 395 region, and events with scores lower than this form a tZ control region. This cutoff is selected by
 396 varying the value of this cutoff in stat-only Asimov fits, and selecting the value that minimizes
 397 the statistical uncertainty on WZ + b.

398 7 Systematic Uncertainties

399 The systematic uncertainties that are considered are summarized in Table 8. These are imple-
 400 mented in the fit either as a normalization factors or as a shape variation or both in the signal
 401 and background estimations. The numerical impact of each of these uncertainties is outlined in
 402 Section 8.

Table 8: Sources of systematic uncertainty considered in the analysis.

Systematic uncertainty	Components
Luminosity	1
Pileup reweighting	1
Physics Objects	
Electron	6
Muon	15
Jet energy scale	28
Jet energy resolution	8
Jet vertex fraction	1
Jet flavor tagging	131
E_T^{miss}	3
Total (Experimental)	194
Signal Modeling	
Shape modelling	3
Renormalization and factorization scales	5
nJet Migration	4
Background Modeling	
Cross section	15
Renormalization and factorization scales	12
Total (Signal and background modeling)	35
Total (Overall)	229

403 The uncertainty in the combined integrated luminosity is derived from a calibration of the
 404 luminosity scale performed for 13 TeV proton-proton collisions [12], [LUCID2].

405 The experimental uncertainties are related to the reconstruction and identification of light leptons
 406 and b-tagging of jets, and to the reconstruction of E_T^{miss} . The TOTAL electron ID correlation

407 model is used, corresponding to 1 electron ID systematic. Electron ID is found to be a subleading
408 systematic that is unconstrained by the fit, making it an appropriate choice for this analysis.

409 The sources which contribute to the uncertainty in the jet energy scale (JES) [13] are decom-
410 posed into uncorrelated components and treated as independent sources in the analysis. The
411 CategoryReduction model is used to account for JES uncertainties, which decomposes the uncer-
412 tainties into 30 nuisance parameters included in the fit. The SimpleJER model is used to account
413 for jet energy resolution (JER) uncertainties, and 8 JER uncertainty components unclued as
414 NPs in the fit.

415 The uncertainties in the b-tagging efficiencies measured in dedicated calibration analyses [14] are
416 also decomposed into uncorrelated components. The large number of components for b-tagging
417 is due to the calibration of the distribution of the MVA discriminant.

418 The full list of systematic uncertainties considered in the analysis is summarized in Tables 9, 10
419 and 11.

420

Experimental Systematics on Leptons and E_T^{miss}			
Type	Description	Systematics Name	Application
Trigger			
Scale Factors	Trigger Efficiency	lepSFTrigTight_MU(EL)_SF_Trigger_STAT(SYST)	Event Weight
Muons			
Efficiencies	Reconstruction and Identification	lepSFObjTight_MU_SF_ID_STAT(SYST)	Event Weight
	Isolation	lepSFObjTight_MU_SF_Isol_STAT(SYST)	Event Weight
	Track To Vertex Association	lepSFObjTight_MU_SF_TTVA_STAT(SYST)	Event Weight
p_T Scale	p_T Scale	MUONS_SCALE	p_T Correction
Resolution	Inner Detector Energy Resolution	MUONS_ID	p_T Correction
	Muon Spectrometer Energy Resolution	MUONS_MS	p_T Correction
Electrons			
Efficiencies	Reconstruction	lepSFObjTight_EL_SF_ID	Event Weight
	Identification	lepSFObjTight_EL_SF_Reco	Event Weight
	Isolation	lepSFObjTight_EL_SF_Isol	Event Weight
Scale Factor	Energy Scale	EG_SCALE_ALL	Energy Correction
Resolution	Energy Resolution	EG_RESOLUTION_ALL	Energy Correction
E_T^{miss}			
Soft Tracks Terms	Resolution	MET_SoftTrk_ResoPerp	p_T Correction
	Resolution	MET_SoftTrk_ResoPara	p_T Correction
	Scale	MET_SoftTrk_ScaleUp	p_T Correction
	Scale	MET_SoftTrk_ScaleDown	p_T Correction

Table 9: Summary of experimental systematics considered for leptons and E_T^{miss} . Includes type, description, name of systematic as used in the fit, and mode of application. The mode of application indicates the systematic evaluation, e.g. as an overall event re-weighting (Event Weight) or rescaling (p_T Correction).

Experimental Systematics on Jets			
Type	Origin	Systematics Name	Application
Jet Vertex Tagger		JVT	Event Weight
Energy Scale	Calibration Method	JET_21NP_ JET_EffectiveNP_1-19	p_T Correction p_T Correction
	η inter-calibration	JET_EtaIntercalibration_Modelling JET_EtaIntercalibration_NonClosure JET_EtaIntercalibration_TotalStat	p_T Correction p_T Correction p_T Correction
	High p_T jets	JET_SingleParticle_HighPt	p_T Correction
	Pile-Up	JET_Pileup_OffsetNPV JET_Pileup_OffsetMu JET_Pileup_PtTerm JET_Pileup_RhoTopology	p_T Correction p_T Correction p_T Correction p_T Correction
	Non Closure	JET_PunchThrough_MC15	p_T Correction
	Flavour	JET_Flavor_Response JET_BJES_Response JET_Flavor_Composition	p_T Correction p_T Correction p_T Correction
Resolution		JET_JER_SINGLE_NP	Event Weight

Table 10: Jet systematics take into account effects of jets calibration method, η inter-calibration, high p_T jets, pile-up, and flavor response. They are all diagonalised into effective parameters.

Experimental Systematics on b-tagging		
Type	Origin	Systematic Name
Scale Factors	DL1r b-tagger efficiency on b originated jets in bins of η	DL1r_Continuous_EventWeight_B0-29
	DL1r b-tagger efficiency on c originated jets in bins of η	DL1r_Continuous_EventWeight_C0-19
	DL1r b-tagger efficiency on light flavoured originated jets in bins of η and p_T	DL1r_Continuous_EventWeight_Light0-79
	DL1r b-tagger extrapolation efficiency	DL1r_Continuous_EventWeight_extrapolation DL1r_Continuous_EventWeight_extrapolation_from_charm

Table 11: Summary of experimental systematics to be included for b-tagging of jets in the analysis, using the continuous DL1r tagging algorithm. All of the b-tagging related systematics are applied as event weights. From left: type, description, and the name of systematic used in the fit.

- 421 Theoretical uncertainties applied to MC predictions, including cross section, PDF, and scale
 422 uncertainties are taken from theory calculations, with the exception of non-prompt and diboson
 423 backgrounds. The cross-section uncertainty on tZ is taken from [tZ_paper]. Derivation of the
 424 non-prompt background uncertainties, Z+jets and tt>, are explained in detail in Section 5.3.
 425 The other VV + heavy flavor processes (namely VV+b and VV+charm, which primarily consist
 426 of ZZ events) are also poorly understood, because these processes involve the same physics as
 427 WZ + heavy flavor, and have also not been measured. Therefore, a conservative 50% uncertainty
 428 is applied to those samples. While this uncertainty is large, it is found to have little impact on
 429 the significance of the final result.
 430 The theory uncertainties applied to the predominate background estimates are summarized in
 431 Table 12.

Process	X-section [%]
tZ	X-sec: ± 15.2 QCD Scale: $^{+5.2}_{-1.3}$ PDF($+\alpha_S$): ± 1.2
t <bar>t} H (aMC@NLO+Pythia8)</bar>	QCD Scale: $^{+5.8}_{-9.2}$ PDF($+\alpha_S$): ± 3.6
t <bar>t} Z (aMC@NLO+Pythia8)</bar>	QCD Scale: $^{+9.6}_{-11.3}$ PDF($+\alpha_S$): ± 4
t <bar>t} W (aMC@NLO+Pythia8)</bar>	QCD Scale: $^{+12.9}_{-11.5}$ PDF($+\alpha_S$): ± 3.4
VV + b/charm (Sherpa 2.2.1)	± 50
VV + light (Sherpa 2.2.1)	± 6
t <bar>t}</bar>	± 10
Z + jets	± 20
Others	± 50

Table 12: Summary of theoretical uncertainties for MC predictions in the analysis.

432 Additional signal uncertainties are estimated by comparing estimates from the nominal Sherpa
 433 WZ samples with alternate WZ samples generated with Powheg+PYTHIA8 (DSID 361601).
 434 The shape of the templates used in the fit are compared between these two samples for WZ + b,
 435 WZ + charm and WZ + light, as shown in Figures 27 and 28. Each of these plots are normalized
 436 to unity in order to capture differences in shape.

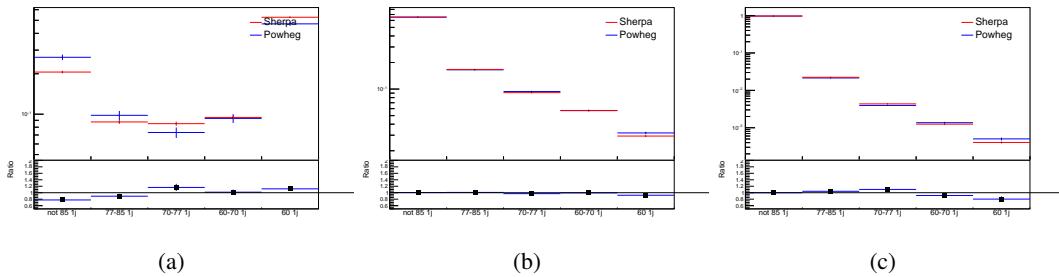


Figure 27: Comparison between Sherpa and Powheg predictions of the distribution of (a) WZ + b, (b) WZ + charm, and (c) WZ + light among the various b-tag WPs used in the 1-jet fit.

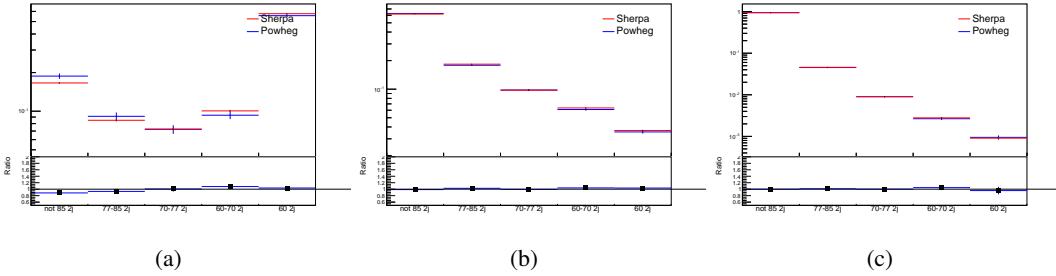


Figure 28: Comparison between Sherpa and Powheg predictions of the distribution of (a) WZ + b, (b) WZ + charm, and (c) WZ + light among the various b-tag WPs used in the 2-jet fit.

⁴³⁷ Separate systematics are included in the fit for WZ + b, WZ + charm and WZ + light, where
⁴³⁸ the distribution among each of the fit regions is varied based on the prediction of the Powheg
⁴³⁹ sample.

440 A similar approach is taken to account for uncertainties in migrations between the number of
 441 reco and truth jets. The fraction of events with 1 truth jet which fall into the 1 jet bin versus the
 442 2 jet bin at reco level is compared for Sherpa and Powheg. The same is done for events with 2
 443 truth jets. This comparison is shown in figure 29.

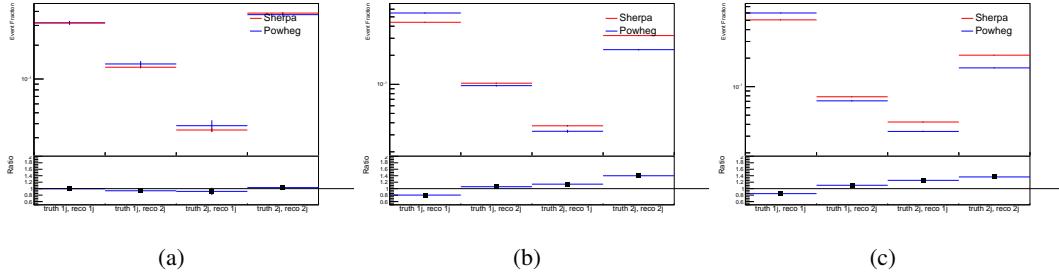


Figure 29: Comparison between Sherpa and Powheg predictions for truth jet migrations between the 1 and 2 jet reco bins for (a) WZ + b, (b) WZ + charm, and (c) WZ + light

444 A systematic is included where events are shifted between the 1-jet and 2-jet regions based on
445 the differences between these two shapes. This is done independently for each of the WZ + b,
446 WZ + charm, and WZ + light templates.

Additional systematics are included to account for the uncertainty in the contamination of 0 jet and 3 or more jet events (at truth level) in the 1 and 2 reco jet bins. Because these events fall outside the scope of this measurement, these events are included as a background. As such, a normalization, rather than a shape, uncertainty is applied for this background.

451 The number of WZ events with 0-jets and ≥ 3 -jets in the reconstructed 1-jet and 2-jet regions
 452 are compared for Sherpa and Powheg, as seen in figure 30. These differences are taken as separate
 453 normalization systematics on the yield of WZ+0-jet and WZ+ ≥ 3 -jet events.

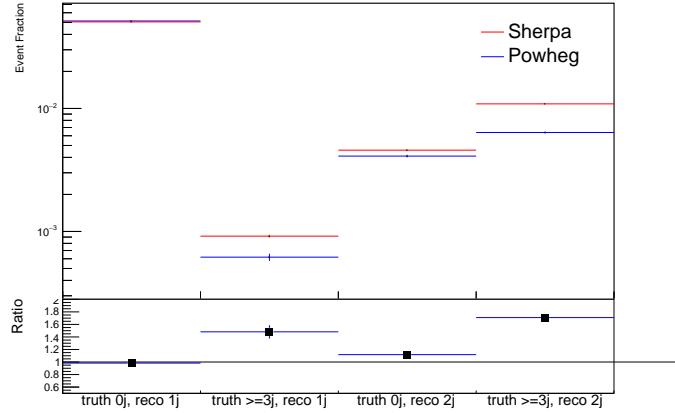


Figure 30: Comparison between Sherpa and Powheg predictions for truth jet migrations between the 1 and 2 jet reco bins

454 8 Results

455 A separate maximum-likelihood fit is performed over the various fit regions in order to extract
 456 the best-fit value of the WZ + b-jet and WZ + charm jet contributions. Separate measurements
 457 are made for WZ + 1-jet and WZ + 2-jet events. The WZ + b, WZ + charm and WZ + light
 458 contributions, separated into 1-jet and 2-jet samples at truth level, are allowed to float, with the
 459 remaining background contributions are held fixed. The result of the fit is used to extract the
 460 cross-section of WZ + heavy-flavor production.

461 A maximum likelihood fit to data is performed simultaneously in the regions described in Section
 462 5, summarized in figure 31. The parameters $\mu_{WZ+b-1-jet}$, $\mu_{WZ+charm1-jet}$, $\mu_{WZ+light-1-jet}$,
 463 $\mu_{WZ+b-2-jet}$, $\mu_{WZ+charm2-jet}$, $\mu_{WZ+light-2-jet}$, where $\mu = \sigma_{\text{observed}}/\sigma_{\text{SM}}$, are extracted from
 464 the fit. A simultaneous fit is performed over all 1-jet and 2-jet regions.

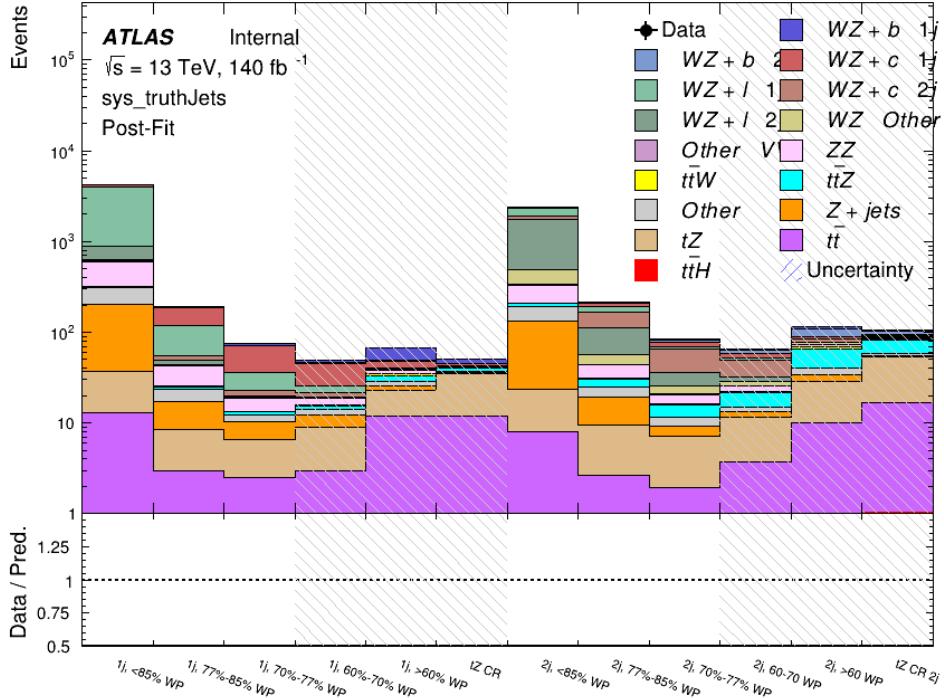


Figure 31: Post-fit summary of the fit regions.

465 **These numbers are preliminary, and need to be updated to account for the latest results.**

466 The Asimov fit for 1-jet events gives an expected μ value of $1.00^{+0.37}_{-0.35}(\text{stat})^{+0.24}_{-0.22}(\text{sys})$ for WZ
 467 + b. The fitted cross-section modifiers for WZ + charm and WZ + light are $1.00 \pm 0.17 \pm 0.15$
 468 and $1.00 \pm 0.04 \pm 0.07$, respectively.

469 The expected cross-section of WZ+b with 1 associated jet obtained from the fit is $1.74^{+0.75}_{-0.72}(\text{stat})^{+0.62}_{-0.57}(\text{sys})$
 470 fb with an expected significance of 1.8σ . The expected cross-section of WZ + charm is measured
 471 to be $14.6 \pm 2.5(\text{stat}) \pm 2.1(\text{sys})$ fb, with a correlation of -0.23.

472 For 2-jet events, the fit gives an expected μ value of $1.00^{+0.34}_{-0.33}(\text{stat})^{+0.29}_{-0.28}(\text{sys})$ for WZ + b.
 473 The fitted cross-section modifiers for WZ + charm and WZ + light are $1.00 \pm 0.17 \pm 0.15$ and
 474 $1.00 \pm 0.04 \pm 0.08$, respectively.

475 The expected WZ + b cross-section in the 2-jet region is $2.46^{+0.83}_{-0.81}(\text{stat})^{+0.73}_{-0.68}(\text{sys})$ fb with
 476 an expected significance of 2.1σ . The 2-jet expected cross-section of WZ + charm is $12.7 \pm$
 477 $2.2(\text{stat}) \pm 2.0(\text{sys})$ fb, and the correlation between WZ + charm and WZ + b is -0.26.

478 While the fit is performed for WZ+1-jet and WZ+2-jet simultaneously, the detailed results
 479 below

⁴⁸⁰ **8.1 1-jet Results**

⁴⁸¹ **The results of the fit are currently blinded.**

⁴⁸² The pre-fit yields in each of the regions used in the fit are shown in Table 8.1.

⁴⁸³

Sample	1j, <85% WP	1j, 77%-85% WP	1j, 70%-77% WP	1j, 60%-70% WP	1j, >60% WP	tZ CR
WZ + b - 1j	8.1 ± 1.6	4.7 ± 0.5	4.6 ± 0.4	5.1 ± 0.4	18.2 ± 2.4	4.8 ± 0.6
WZ + c - 1j	260 ± 22	81 ± 6	43.1 ± 3.6	25.8 ± 2.6	9.4 ± 1.8	2.9 ± 0.6
WZ + l - 1j	3090 ± 250	91 ± 13	17 ± 3	4.9 ± 1.6	1.3 ± 0.4	0.2 ± 0.1
WZ + b - 2j	1.10 ± 0.37	0.44 ± 0.11	0.39 ± 0.06	0.62 ± 0.14	2.1 ± 0.5	0.59 ± 0.14
WZ + c - 2j	21 ± 5	5.6 ± 1.2	3.0 ± 0.7	2.0 ± 0.5	0.70 ± 0.20	0.30 ± 0.08
WZ + l - 2j	250 ± 60	5.7 ± 1.6	0.73 ± 0.53	0.31 ± 0.15	0.07 ± 0.06	0.01 ± 0.01
WZ - Other	13 ± 5	1.4 ± 0.4	0.42 ± 0.08	0.2 ± 0.01	0.30 ± 0.05	0.67 ± 0.15
Other VV	6.2 ± 0.6	0.2 ± 0.4	0.2 ± 0.04	0.07 ± 0.1	0.1 ± 0.1	0.1 ± 0.2
ZZ	336 ± 26	17.8 ± 2.1	4.3 ± 0.6	1.7 ± 0.5	0.36 ± 0.08	0.10 ± 0.03
t̄tW	1.1 ± 0.2	0.2 ± 0.1	0.3 ± 0.1	0.4 ± 0.1	1.5 ± 0.3	0.7 ± 0.2
t̄tZ	6.8 ± 1.2	1.4 ± 0.3	1.0 ± 0.2	1.2 ± 0.2	4.4 ± 0.8	3.2 ± 0.6
Z + jets	169 ± 38	8.9 ± 1.9	3.7 ± 0.8	3.3 ± 0.7	3.2 ± 0.7	0.8 ± 0.17
V + γ	45 ± 28	1.9 ± 2.4	0.1 ± 0.1	0.02 ± 0.01	1.0 ± 0.9	0.02 ± 0.03
tZ	24.3 ± 4.3	5.5 ± 1.1	4.1 ± 0.8	5.9 ± 1.1	10.7 ± 2.0	23 ± 4
tW	1.4 ± 0.8	0.2 ± 0.5	0.0 ± 0.2	0.7 ± 0.6	0.26 ± 0.42	0.39 ± 0.41
WtZ	2.3 ± 1.2	0.6 ± 0.3	0.3 ± 0.21	0.27 ± 0.2	1.1 ± 0.7	0.6 ± 0.5
VVV	12.4 ± 0.5	0.93 ± 0.06	0.35 ± 0.03	0.13 ± 0.02	0.14 ± 0.03	0.02 ± 0.01
VH	40 ± 6	2.6 ± 1.4	0.9 ± 0.8	0.7 ± 0.8	0.5 ± 0.6	0.0 ± 0.0
t̄t	12.1 ± 1.6	2.9 ± 0.6	2.5 ± 0.5	2.8 ± 0.5	11.2 ± 1.4	10.9 ± 1.5
t̄tH	0.24 ± 0.03	0.05 ± 0.01	0.04 ± 0.01	0.06 ± 0.01	0.20 ± 0.03	0.13 ± 0.02
Total	5010 ± 260	227 ± 24	88 ± 12	57 ± 8	76 ± 16	53 ± 8

Table 13: Pre-fit yields in each of the 1-jet fit regions.

⁴⁸⁴ The post-fit yields in each region are summarized in Figure 8.1.

⁴⁸⁵

	1j, <85% WP	1j, 77%-85% WP	1j, 70%-77% WP	1j, 60%-70% WP	1j, >60% WP	tZ CR
WZ + b	8.1 ± 4.9	4.7 ± 2.0	4.6 ± 2.0	5.1 ± 2.1	18 ± 10	5.0 ± 2.5
WZ + c	260 ± 60	80 ± 14	43 ± 7	26 ± 5	9.4 ± 2.3	2.9 ± 0.7
WZ + l	3090 ± 130	90 ± 11	17.3 ± 2.8	4.9 ± 1.6	1.3 ± 0.4	0.23 ± 0.13
WZ + b - 2j	1.10 ± 0.37	0.44 ± 0.11	0.39 ± 0.06	0.62 ± 0.14	2.1 ± 0.5	0.59 ± 0.14
WZ + c - 2j	21 ± 5	5.6 ± 1.2	3.0 ± 0.7	2.0 ± 0.5	0.70 ± 0.20	0.30 ± 0.08
WZ + l - 2j	250 ± 60	5.7 ± 1.6	0.73 ± 0.53	0.31 ± 0.15	0.07 ± 0.06	0.01 ± 0.01
WZ - Other	13 ± 5	1.4 ± 0.4	0.42 ± 0.08	0.2 ± 0.01	0.30 ± 0.05	0.67 ± 0.15
Other VV	6.2 ± 0.6	0.92 ± 0.07	0.02 ± 0.01	0.01 ± 0.01	0.01 ± 0.01	0.01 ± 0.01
ZZ	346 ± 57	19 ± 5	4.3 ± 0.8	2.7 ± 0.5	2.4 ± 0.1	2.1 ± 0.6
t̄tW	1.09 ± 0.21	0.2 ± 0.1	0.1 ± 0.1	0.4 ± 0.1	1.5 ± 0.3	0.1 ± 0.2
t̄tZ	6.8 ± 1.2	1.4 ± 0.3	1.0 ± 0.2	1.2 ± 0.2	4.4 ± 0.7	3.2 ± 0.5
rare Top	0.14 ± 0.04	0.04 ± 0.02	0.04 ± 0.0	0.1 ± 0.03	0.14 ± 0.04	0.15 ± 0.05
t̄tWW	0.04 ± 0.03	0.01 ± 0.02	0.01 ± 0.01	0.01 ± 0.01	0.01 ± 0.02	0.01 ± 0.01
Z + jets	169 ± 37	8.9 ± 1.9	3.7 ± 0.8	3.3 ± 0.7	3.2 ± 0.7	0.8 ± 0.2
W + jets	0.01 ± 0.01					
V + γ	46 ± 28	1.9 ± 2.4	0.1 ± 0.1	0.0 ± 0.2	1.0 ± 0.9	0.0 ± 0.0
tZ	24 ± 4	5.5 ± 1.0	4.1 ± 0.8	5.9 ± 1.1	10.7 ± 1.8	23.3 ± 3.7
tW	1.37 ± 0.82	0.18 ± 0.26	0.01 ± 0.12	0.67 ± 0.64	0.26 ± 0.42	0.39 ± 0.41
WtZ	2.3 ± 1.2	0.6 ± 0.3	0.3 ± 0.2	0.3 ± 0.2	1.1 ± 0.6	0.6 ± 0.3
VVV	12.4 ± 0.4	0.9 ± 0.1	0.4 ± 0.1	0.13 ± 0.02	0.14 ± 0.03	0.02 ± 0.01
VH	40 ± 6	2.6 ± 1.4	0.9 ± 0.8	0.7 ± 0.8	0.4 ± 0.6	0.01 ± 0.01
t̄t	12.1 ± 1.6	2.9 ± 0.6	2.5 ± 0.5	2.8 ± 0.5	11.2 ± 1.5	10.9 ± 1.4
t̄tH	0.24 ± 0.03	0.05 ± 0.01	0.04 ± 0.01	0.06 ± 0.01	0.20 ± 0.03	0.13 ± 0.02
Total	5100 ± 110	227 ± 12	87 ± 6	56.7 ± 4.4	76 ± 9	52.5 ± 4.2

Table 14: Post-fit yields in each of the 1-jet fit regions.

486 As described in Section 7, there are 226 systematic uncertainties that are considered as NPs in
 487 the fit. These NPs are constrained by Gaussian or log-normal probability density functions. The
 488 latter are used for normalisation factors to ensure that they are always positive. The expected
 489 number of signal and background events are functions of the likelihood. The prior for each NP is
 490 added as a penalty term, decreasing the likelihood as it is shifted away from its nominal value.

491 The impact of each NP is calculated by performing the fit with the parameter of interest held
 492 fixed, varied from its fitted value by its uncertainty, and calculating $\Delta\mu$ relative to the baseline
 493 fit. The impact of the most significant sources of systematic uncertainties is summarized in Table
 494 15.

Uncertainty Source	$\Delta\mu$	
WZ + 1-jet light XS	0.13	-0.18
WZ + charm cross-section	-0.10	0.12
Jet Energy Scale	0.1	-0.13
tZ cross-section	-0.10	0.10
WZ 1-jet/2-jet Migration	0.08	-0.07
Jet Energy Resolution	-0.10	0.10
Luminosity	-0.08	0.09
Other Diboson + b cross-section	-0.07	0.07
Flavor tagging	0.05	0.05
t <bar>t} cross-section</bar>	-0.05	0.05
Total Systematic Uncertainty	0.29	0.34

Table 15: Summary of the most significant sources of systematic uncertainty on the measurement of WZ + b with exactly one associated jet.

⁴⁹⁵ The ranking and impact of those nuisance parameters with the largest contribution to the overall
⁴⁹⁶ uncertainty is shown in Figure 32.

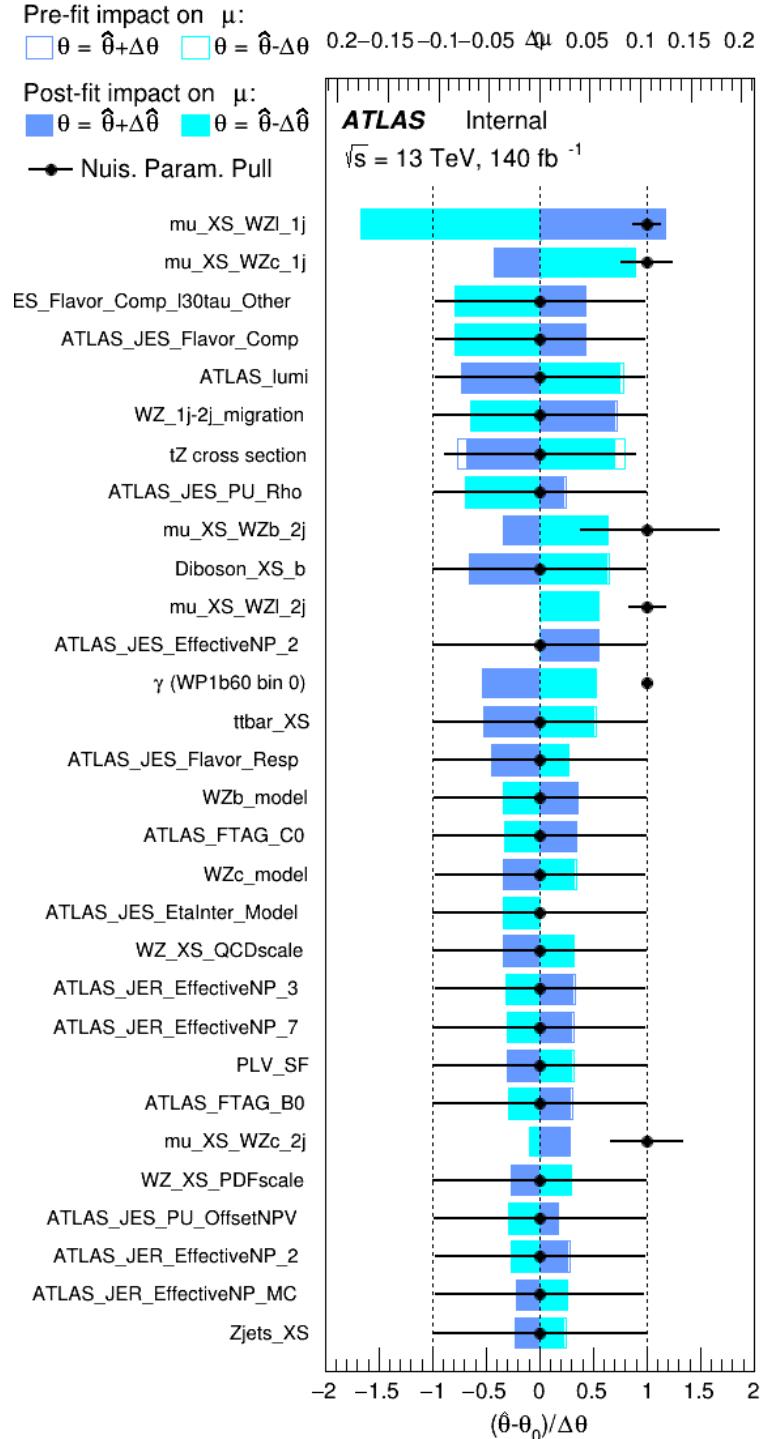


Figure 32: Impact of systematic uncertainties on the signal-strength of $WZ + b$ for events with exactly one jet

497 The large impact of the Jet Energy Scale and Jet Flavor Tagging is unsurprising, as the shape of
498 the fit regions depends heavily on the modeling of the jets. The other major sources of uncertainty
499 come from background modelling and cross-section uncertainty.

The correlations between these nuisance parameters are summarized in Figure 33.

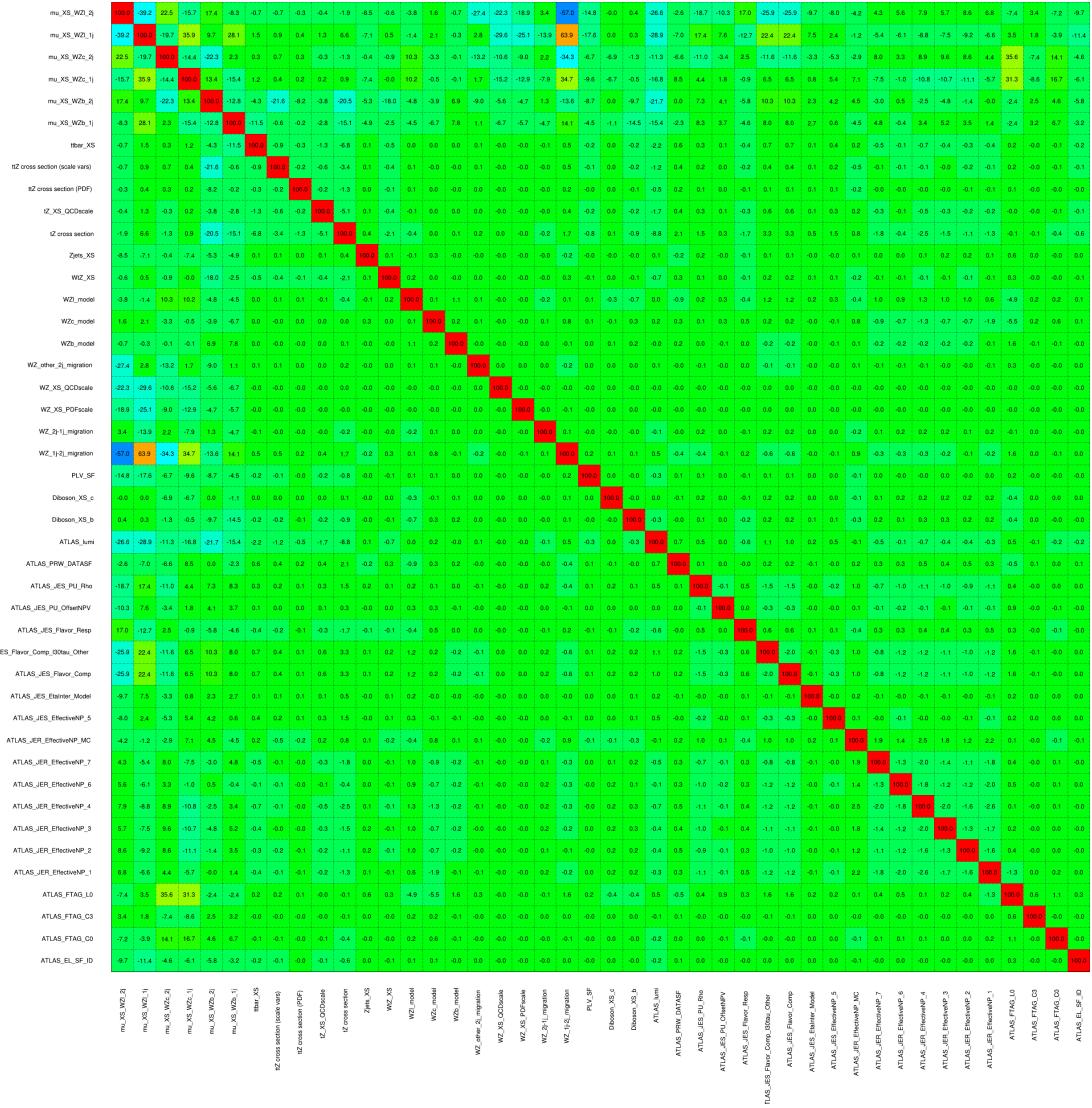


Figure 33: Correlations between nuisance parameters

501 The negative correlations between $\mu_{WZ+charm}$ and μ_{WZ+b} and $\mu_{WZ+light}$ are expected: WZ +
 502 charm is present in both the WZ + b and WZ + light enriched regions, therefore increasing the

503 fraction of charm requires increasing the fraction of $WZ + b$ and $WZ + \text{light}$. This reasoning
 504 also explains the positive correlation between μ_{WZ+b} and $\mu_{WZ+\text{light}}$.

505 Two of the major backgrounds in the region with the highest purity of $WZ + b$ are tZ and Other
 506 VV + b, explaining the negative correlations between μ_{WZ+b} and the tZ cross section, and the
 507 VV + b cross section.

508 The high correlation between the luminosity and $\mu_{WZ+\text{light}}$ arises from the fact that the uncer-
 509 tainty on $\mu_{WZ+\text{light}}$ is very low (around 4%). Small changes in luminosity cause a change in
 510 the yield of $WZ + \text{light}$ that is large compared to its uncertainty, producing a large correlation
 511 between these two parameters.

512 8.2 2-jet Results

513 **The results of the fit are currently blinded.**

514 Pre-fit yields in each of the 2-jet fit regions are shown in Figure 8.2.

	2j, <85% WP	2j, 77%-85% WP	2j, 70%-77% WP	2j, 60%-70% WP	2j, >60% WP	tZ CR 2j
$WZ + b - 2j$	3.1 ± 1.6	6.7 ± 0.5	5.6 ± 0.4	8.0 ± 0.6	31 ± 2	14 ± 1
$WZ + c - 2j$	180 ± 20	54 ± 6	41 ± 3	24 ± 3	11 ± 2	4.8 ± 0.6
$WZ + l - 2j$	1250 ± 150	90 ± 14	18 ± 3	5.8 ± 1.4	1.4 ± 0.4	0.25 ± 0.15
$WZ + b - 1j$	3.4 ± 0.6	1.52 ± 0.35	1.58 ± 0.23	1.95 ± 0.39	6.7 ± 1.1	1.9 ± 0.6
$WZ + c - 1j$	56 ± 14	17.6 ± 4.0	8.6 ± 2.2	6.3 ± 1.8	3.0 ± 0.9	0.7 ± 0.2
$WZ + l - 1j$	427 ± 120	24 ± 7	4.7 ± 2.3	1.6 ± 0.7	0.3 ± 0.2	0.01 ± 0.01
$WZ - \text{Other}$	129 ± 29	6.1 ± 4.6	1.2 ± 0.3	0.3 ± 0.2	2.9 ± 0.5	3.6 ± 0.6
Other VV	7.63 ± 0.63	0.6 ± 0.5	0.16 ± 0.03	0.01 ± 0.01	0.1 ± 0.1	0.1 ± 0.1
ZZ	135 ± 20	14.1 ± 3.2	4.7 ± 0.8	4.0 ± 0.6	4.1 ± 0.7	3.1 ± 0.5
$t\bar{t}W$	0.8 ± 0.2	0.4 ± 0.1	0.5 ± 0.1	0.7 ± 0.2	4.3 ± 0.6	3.9 ± 0.6
$t\bar{t}Z$	14.7 ± 2.2	5.6 ± 0.8	4.5 ± 0.7	6.5 ± 1.1	25.4 ± 4.0	21.9 ± 3.4
rare Top	0.14 ± 0.04	0.07 ± 0.03	0.03 ± 0.02	0.09 ± 0.03	0.37 ± 0.07	0.6 ± 0.1
$t\bar{t}WW$	0.04 ± 0.03	0.02 ± 0.02	0.01 ± 0.01	0.01 ± 0.00	0.03 ± 0.03	0.01 ± 0.01
$Z + \text{jets}$	110.0 ± 22.9	9.6 ± 2.0	2.1 ± 0.50	1.6 ± 0.4	5.1 ± 1.1	1.5 ± 0.3
$W + \text{jets}$	0.0 ± 0.0	0.0 ± 0.0	0.0 ± 0.0	0.0 ± 0.0	0.0 ± 0.0	0.0 ± 0.0
$V + \gamma$	25 ± 18	0.5 ± 0.2	0.1 ± 0.1	0.1 ± 0.1	0.0 ± 0.02	0.05 ± 0.07
tZ	15.9 ± 2.9	6.9 ± 1.3	5.1 ± 1.0	8.0 ± 1.5	18.7 ± 3.2	36.4 ± 6.1
tW	0.9 ± 0.7	0.2 ± 0.3	0.0 ± 0.1	0.0 ± 0.0	0.8 ± 0.6	0.2 ± 0.2
WtZ	4.9 ± 2.5	1.5 ± 0.8	1.1 ± 0.6	1.3 ± 0.7	4.6 ± 2.4	3.3 ± 1.7
VVV	7.4 ± 0.3	1.0 ± 0.1	0.4 ± 0.1	0.2 ± 0.1	0.13 ± 0.03	0.04 ± 0.01
VH	19.5 ± 4.2	2.8 ± 1.6	0.7 ± 0.7	0.1 ± 0.2	0.0 ± 0.0	0.0 ± 0.0
$t\bar{t}$	0.7 ± 0.4	0.1 ± 0.1	0.05 ± 0.06	0.15 ± 0.13	0.8 ± 0.5	2.3 ± 1.2
$t\bar{t}$	6.8 ± 1.0	2.4 ± 0.5	1.8 ± 0.4	3.3 ± 0.6	8.4 ± 1.2	13.6 ± 1.7
$t\bar{t}H$	0.4 ± 0.1	0.2 ± 0.1	0.16 ± 0.02	0.23 ± 0.03	0.94 ± 0.11	1.03 ± 0.12
Total	2580 ± 160	229 ± 24	89 ± 13	69 ± 11	120 ± 15	108 ± 11

Table 16: Pre-fit yields in each of the 2-jet fit regions.

⁵¹⁵ The post-fit yields in each region are summarized in Figure 8.2.

	2j, <85% WP	2j, 77%-85% WP	2j, 70%-77% WP	2j, 60%-70% WP	2j, >60% WP	tZ CR 2j
WZ + b	13 ± 6	6.7 ± 2.9	5.8 ± 2.5	8.0 ± 3.5	31 ± 13	14 ± 5
WZ + c	260 ± 60	77 ± 15	41 ± 8	26 ± 5	10.9 ± 2.4	4.8 ± 1.1
WZ + l	1860 ± 90	90 ± 12	17.6 ± 2.8	5.8 ± 1.3	1.4 ± 0.4	0.3 ± 0.2
WZ + b - 1j	3.4 ± 0.6	1.52 ± 0.35	1.58 ± 0.23	1.95 ± 0.39	6.7 ± 1.1	1.9 ± 0.6
WZ + c - 1j	56 ± 14	17.6 ± 4.0	8.6 ± 2.2	6.3 ± 1.8	3.0 ± 0.9	0.7 ± 0.2
WZ + l - 1j	427 ± 120	24 ± 7	4.7 ± 2.3	1.6 ± 0.7	0.3 ± 0.2	0.01 ± 0.01
WZ - Other	129 ± 29	6.1 ± 4.6	1.2 ± 0.3	0.3 ± 0.2	2.9 ± 0.5	3.6 ± 0.6
Other VV	7.6 ± 0.6	0.3 ± 0.3	0.3 ± 0.1	0.1 ± 0.06	0.03 ± 0.02	0.1 ± 0.1
ZZ	145 ± 30	11.3 ± 4.4	2.7 ± 1.6	1.0 ± 0.3	4.0 ± 0.1	2.4 ± 0.1
t̄tW	0.8 ± 0.2	0.4 ± 0.1	0.54 ± 0.12	0.74 ± 0.15	4.3 ± 0.6	3.9 ± 0.6
t̄tZ	14.7 ± 2.2	5.6 ± 0.8	4.5 ± 0.7	6.5 ± 1.0	25.4 ± 3.9	21.9 ± 3.3
rare Top	0.14 ± 0.04	0.07 ± 0.03	0.03 ± 0.02	0.09 ± 0.03	0.4 ± 0.1	0.6 ± 0.1
t̄tWW	0.04 ± 0.03	0.02 ± 0.02	0.01 ± 0.01	0.01 ± 0.01	0.03 ± 0.03	0.01 ± 0.01
Z + jets	110 ± 23	9.6 ± 2.0	2.1 ± 0.5	1.6 ± 0.4	5.1 ± 1.1	1.5 ± 0.3
W + jets	0.0 ± 0.0					
V + γ	25 ± 19	0.5 ± 0.2	0.1 ± 0.1	0.13 ± 0.14	0.0 ± 0.02	0.05 ± 0.07
tZ	15.9 ± 2.7	6.9 ± 1.2	5.1 ± 0.9	8.0 ± 1.4	18.7 ± 3.0	36 ± 6
tW	0.1 ± 0.7	0.2 ± 0.3	0.0 ± 0.1	0.0 ± 0.0	0.8 ± 0.6	0.2 ± 0.2
WtZ	4.9 ± 2.5	1.5 ± 0.8	1.1 ± 0.6	1.3 ± 0.7	4.6 ± 2.3	3.3 ± 1.7
VVV	7.4 ± 0.3	1.0 ± 0.1	0.36 ± 0.03	0.19 ± 0.03	0.13 ± 0.03	0.04 ± 0.01
VH	19 ± 4	2.8 ± 1.6	0.7 ± 0.7	0.1 ± 0.2	0.0 ± 0.0	0.0 ± 0.0
t̄t	6.8 ± 1.0	2.4 ± 0.5	1.8 ± 0.4	3.3 ± 0.6	8.4 ± 1.2	13.6 ± 1.7
t̄tH	0.40 ± 0.05	0.19 ± 0.03	0.16 ± 0.02	0.23 ± 0.03	0.94 ± 0.11	1.03 ± 0.11
Total	2580 ± 60	229 ± 11	89 ± 6	69.1 ± 4.1	120 ± 10	108 ± 6

Table 17: Post-fit yields in each of the 2-jet fit regions.

⁵¹⁶ The same set of systematic uncertainties consider for the 1-jet fit are included in the 2-jet fit as
⁵¹⁷ well. The impact of the most significant systematic uncertainties is summarized in Table 18.

Uncertainty Source	$\Delta\mu$	
WZ + c 2-jet cross-section	0.19	0.15
WZ 0-jet, $>=3$ -jet cross-section	0.14	-0.14
Luminosity	-0.11	0.12
tZ cross-section	-0.11	0.11
Jet Energy Scale	-0.11	0.11
ttZ cross-section - QCD scale	-0.08	0.09
WZ + l cross-section	0.08	-0.07
WtZ cross-section	-0.07	0.07
Flavor tagging	0.05	0.05
Other VV + b cross-section	-0.05	0.05
Total	0.32	0.36

Table 18: Summary of the most significant sources of systematic uncertainty on the measurement of WZ + b 2-jet events.

⁵¹⁸ The ranking and impact of those nuisance parameters with the largest contribution to the overall
⁵¹⁹ uncertainty is shown in Figure 34.

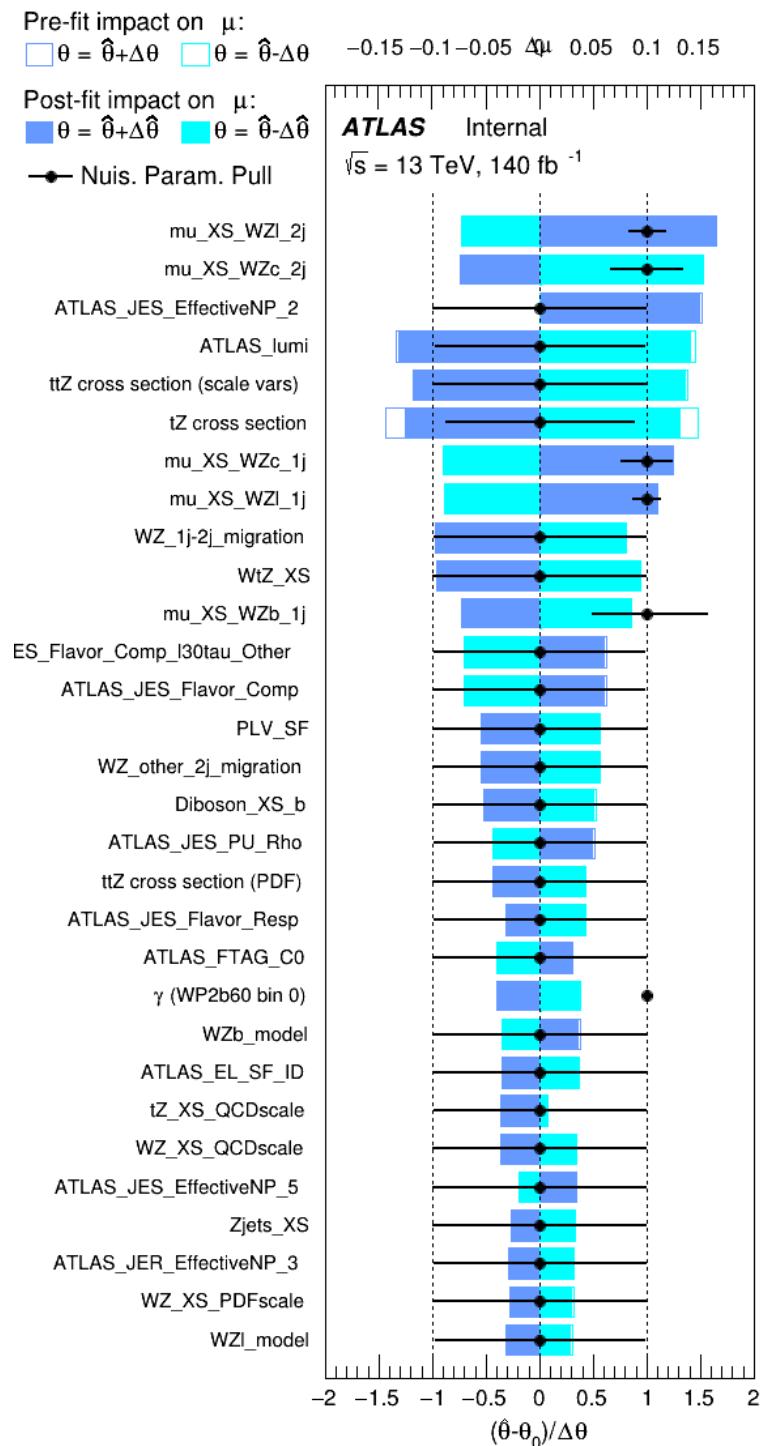


Figure 34: Impact of systematic uncertainties on the signal-strength of $WZ + b$ in 2-jet events.

520 The large impact of the Jet Energy Scale and Jet Flavor Tagging is unsurprising, as the shape of
521 the fit regions depends heavily on the modeling of the jets. The other major sources of uncertainty
522 come from background modelling and cross-section uncertainty.

523 The correlations between these nuisance parameters are summarized in Figure 35.

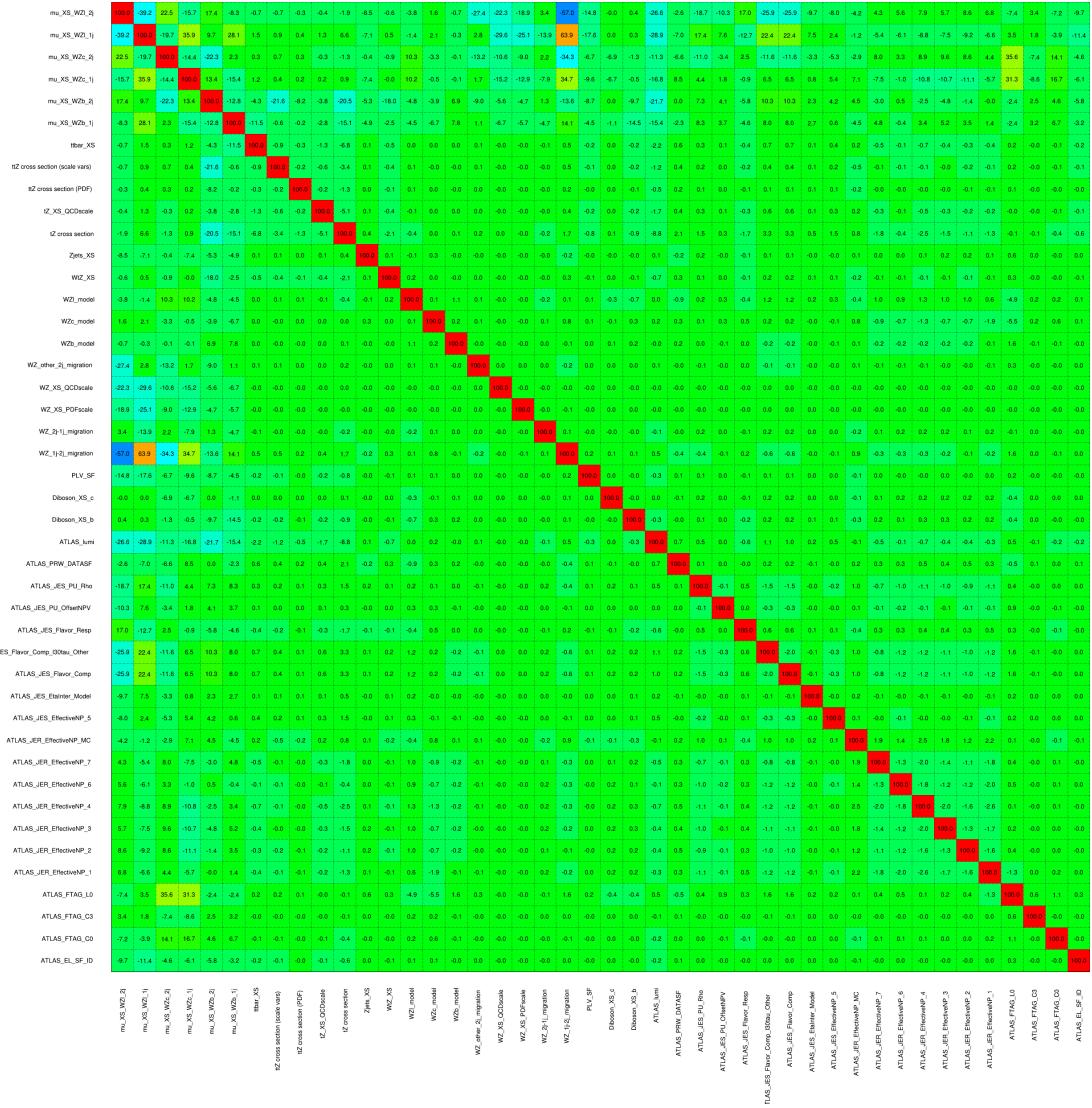


Figure 35: Correlations between nuisance parameters in the 2-jet fit

As in the 1-jet case, no significant, unexpected correlations are found between nuisance parameters.

526 9 Conclusion

527 A measurement of $WZ +$ heavy flavor is performed using 140 fb^{-1} of $\sqrt{s} = 13 \text{ TeV}$ proton-
528 proton collision data collected by the ATLAS detector at the LHC. The expected cross-section
529 of $WZ+b$ with 1-jet is $1.74^{+0.65}_{-0.61}(\text{stat})^{+0.42}_{-0.37}(\text{sys}) \text{ fb}$, and $14.6 \pm 2.5(\text{stat}) \pm 2.1(\text{sys}) \text{ fb}$ for WZ
530 + charm, with a correlation of -0.23 between them. An expected significance of 2.2 is observed
531 for $WZ + b$ in this region.

532 For the 2-jet regions, an expected significance of 2.6 is observed for $WZ + b$, with an expected
533 cross-section of $2.46^{+0.83}_{-0.81}(\text{stat})^{+0.73}_{-0.68}(\text{sys}) \text{ fb}$. For $WZ +$ charm, a cross-section of $12.7 \pm$
534 $2.2(\text{stat}) \pm 2.0(\text{sys}) \text{ fb}$ is expected for 2-jet events. A correlation of -0.26 is observed for $WZ+b$
535 and $WZ +$ charm.

536 **This section will be include final results once unblinded.**

537 **Appendices**

538 **.1 tZ Interference Studies**

539 Because it includes an on-shell Z boson as well as a b-jet and W from the top decay, tZ
540 production represents an identical final state to WZ + b-jet. This implies the possibility of matrix
541 level interference between these two processes not accounted for in the Monte Carlo simulations,
542 which consider the two processes independently. Truth level studies are performed in order to
543 estimate the impact of these interference effects.

544 In order to estimate the matrix level interference effects between tZ and WZ + b-jet, two different
545 sets of simulations are produced using MadGraph 5 [**Madgraph**] - one which simulates these
546 two processes independently, and another where they are produced simultaneously, such that
547 interference effects are present. These two sets of samples are then compared, and the difference
548 between them can be taken to represent any interference effects.

549 MadGraph simulations of 10,000 tZ and 10,000 WZ + b-jet events are produced, along with
550 20,000 events where both are present, in the fiducial region where three leptons and at least one
551 jet are produced.

552 A selection mimicking the preselection used in the main analysis is applied to the samples: The
553 SS leptons are required to have $p_T > 20$ GeV, and > 10 GeV is required for the OS lepton. The
554 associated b-jet is required to have $p_T > 25$ GeV, and all physics objects are required to fall in a
555 range of $|\eta| < 2.5$.

556 The kinematics of these samples after the selection has been applied are shown below:

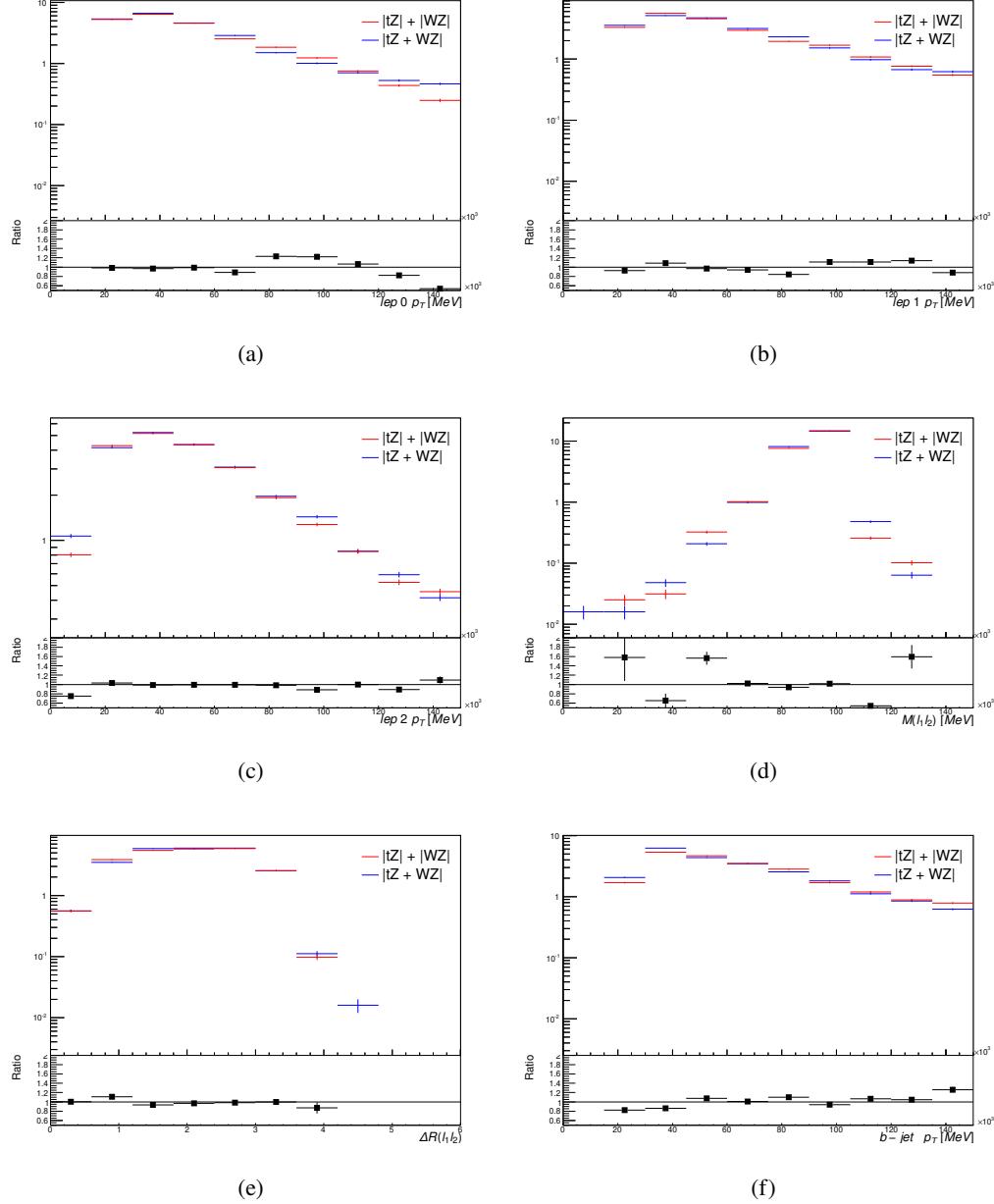


Figure 36: Comparisons between (a) the p_T of the lepton from the W, (b) and (c) show the p_T of the lepton from the Z, (d) the invariant mass of the Z-candidate, (e) ΔR of the leptons from the Z, and (f) the p_T of the b-jet, for WZ and tZ events generated with interference effects (blue) and without interference effects (red).

557 The overall cross-section of the two methods agree within error, and no significant differences
558 in the kinematic distributions are seen. It is therefore concluded that interference effects do not

559 significantly impact the results.

560 **.2 Alternate tZ Inclusive Fit**

561 While tZ is often considered as a distinct process from WZ + b, this could also be considered part
 562 of the signal. Alternate studies are performed where, using the same framework as the nominal
 563 analysis, a measurement of WZ + b is performed that includes tZ as part of WZ+b.

564 Because of this change, the tZ CR is no longer necessary, and only the five pseudo-continuous
 565 b-tag regions are used in the fit. Further, systematics related to the tZ cross-section are removed
 566 from the fit, as they are now encompassed by the normalization measurement of WZ + b. All
 567 other systematic uncertainties are carried over from the nominal analysis.

568 A post-fit summary of the 1-jet regions used in the fit are shown in Figure 37.

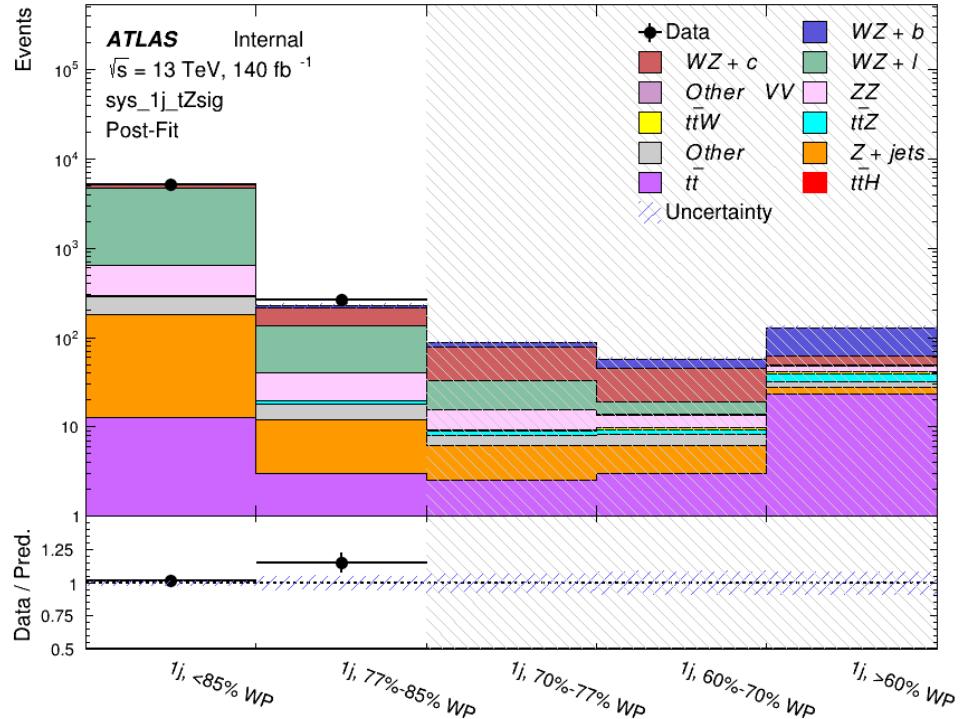


Figure 37: Post-fit summary of the 1-jet fit regions.

569 An expected WZ + b cross-section of $4.1^{+0.78}_{-0.74}(\text{stat})^{+0.53}_{-0.52}(\text{sys}) \text{ fb}$ is extracted from the fit, with
 570 an expected significance of 4.0σ .

571 The impact of the predominate systematics are summarized in Table 19.

Uncertainty Source	$\Delta\mu$	
WZ + light cross-section	0.08	-0.08
Jet Energy Scale	-0.06	0.08
Luminosity	-0.05	0.06
WZ + charm cross-section	-0.04	0.05
Other Diboson + b cross-section	-0.04	0.04
WZ cross-section - QCD scale	-0.04	0.03
t <bar>t</bar>	-0.03	0.03
Jet Energy Resolution	-0.03	0.03
Flavor tagging	-0.03	0.03
Z+jets cross section	-0.02	0.02
Total Systematic Uncertainty	-0.15	0.16

Table 19: Summary of the most significant sources of systematic uncertainty on the measurement of WZ + b with exactly one associated jet.

572 A post-fit summary of the 2-jet regions used in the fit are shown in Figure 38.

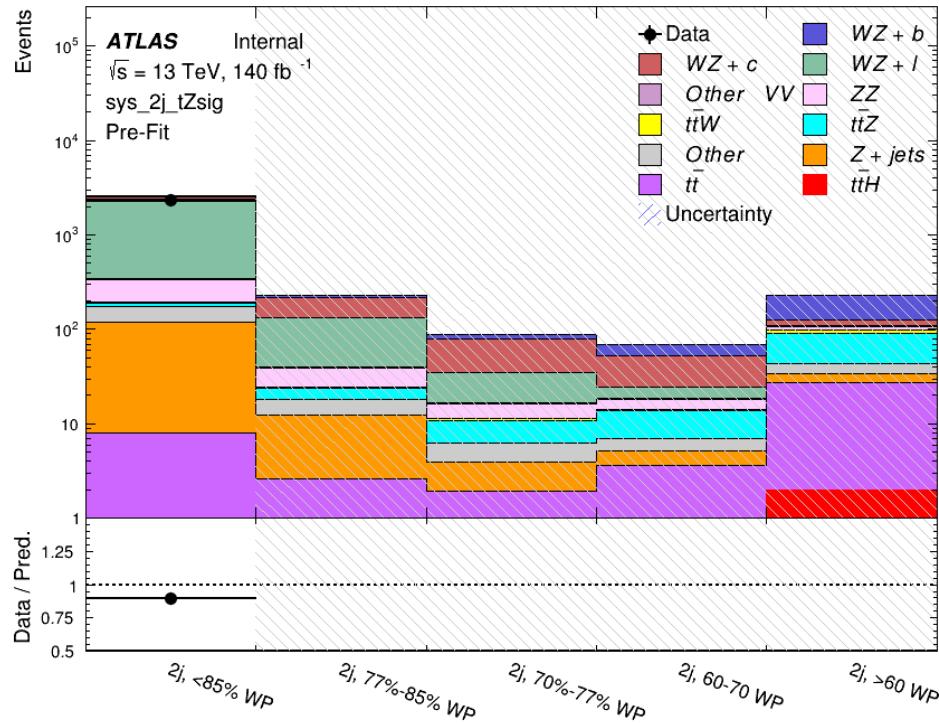


Figure 38: Post-fit summary of the 2-jet fit regions.

573 An expected WZ + b cross-section of $5.9^{+0.9}_{-0.9}(\text{stat})^{+0.9}_{-0.8}(\text{sys}) \text{ fb}$ is extracted from the fit, with an

⁵⁷⁴ expected significance of 5.3σ .

⁵⁷⁵ The impact of the predominate systematics are summarized in Table 20.

Uncertainty Source	$\Delta\mu$	
Luminosity	-0.07	0.07
Jet Energy Scale	0.06	0.05
ttZ cross-section - QCD scale	-0.05	0.05
WZ+l cross-section	0.05	-0.05
WZ+c cross-section	-0.03	0.05
WtZ cross-section	-0.03	0.03
WZ cross-section QCDscale	-0.03	0.03
Diboson cross-section b	-0.03	0.03
WZ cross-section - PDF	-0.03	0.03
Flavor Tagging	0.03	0.02
Total Systematic Uncertainty	-0.14	0.16

Table 20: Summary of the most significant sources of systematic uncertainty on the measurement of WZ + b with exactly one associated jet.

576 .3 DSID list

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 mc16_13TeV.410014.PowhegPythiaEvtGen_P2012_Wt_inclusive_antitop.deriv.DAOD_HIGG8D1.e3753_s3126_r9364_r9315_p4133
 mc16_13TeV.410408.aMcAtNloPythia8EvtGen_tWZ_Ztoll_minDR1.deriv.DAOD_HIGG8D1.e6423_e5984_s3126_r9364_r9315_p4133
 mc16_13TeV.364242.Sherpa_222_NNPDF30NNLO_WWW_3l3v_EW6.deriv.DAOD_HIGG8D1.e5887_s3126_r9364_r9315_p4133
 mc16_13TeV.364243.Sherpa_222_NNPDF30NNLO_WWZ_4l2v_EW6.deriv.DAOD_HIGG8D1.e5887_s3126_r9364_r9315_p4133
 mc16_13TeV.364244.Sherpa_222_NNPDF30NNLO_WWZ_2l4v_EW6.deriv.DAOD_HIGG8D1.e5887_s3126_r9364_r9315_p4133
 mc16_13TeV.364245.Sherpa_222_NNPDF30NNLO_WZZ_5l1v_EW6.deriv.DAOD_HIGG8D1.e5887_s3126_r9364_r9315_p4133
 mc16_13TeV.364246.Sherpa_222_NNPDF30NNLO_WZZ_3l3v_EW6.deriv.DAOD_HIGG8D1.e5887_s3126_r9364_r9315_p4133
 mc16_13TeV.364247.Sherpa_222_NNPDF30NNLO_ZZZ_6l0v_EW6.deriv.DAOD_HIGG8D1.e5887_s3126_r9364_r9315_p4133
 mc16_13TeV.364248.Sherpa_222_NNPDF30NNLO_ZZZ_4l2v_EW6.deriv.DAOD_HIGG8D1.e5887_s3126_r9364_r9315_p4133
 mc16_13TeV.364249.Sherpa_222_NNPDF30NNLO_ZZZ_2l4v_EW6.deriv.DAOD_HIGG8D1.e5887_s3126_r9364_r9315_p4133
 mc16_13TeV.342284.Pythia8EvtGen_A14NNPDF23LO_WH125_inc.deriv.DAOD_HIGG8D1.e4246_s3126_r9364_r9315_p4133
 mc16_13TeV.342285.Pythia8EvtGen_A14NNPDF23LO_ZH125_inc.deriv.DAOD_HIGG8D1.e4246_s3126_r9364_r9315_p4133
 mc16_13TeV.341998.aMcAtNloHppEG_UEEE5_CTEQ6L1_CT10ME_tWH125_gamgam_yt_plus1.deriv.DAOD_HIGG8D1.e4394_e5984_s3126_r9364_r9315_p4133
 mc16_13TeV.410389.MadGraphPythia8EvtGen_A14NNPDF23_tgamma_nonallhadronic.deriv.DAOD_HIGG8D1.e6155_e5984_s3126_r9364_r9315_p4133
 mc16_13TeV.410470.PhPy8EG_A14_ttbar_hdamp258p75_nonallhad.deriv.DAOD_HIGG8D1.e6337_e5984_s3126_r9364_r9315_p4133
 mc16_13TeV.410472.PhPy8EG_A14_ttbar_hdamp258p75_dil.deriv.DAOD_HIGG8D1.e6348_e5984_s3126_r9364_r9315_p4133
 mc16_13TeV.346343.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_allhad.deriv.DAOD_HIGG8D1.e7148_e5984_s3126_r9364_r9315_p4133
 mc16_13TeV.346344.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_semilep.deriv.DAOD_HIGG8D1.e7148_e5984_s3126_r9364_r9315_p4133
 mc16_13TeV.346345.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_dilep.deriv.DAOD_HIGG8D1.e7148_e5984_s3126_r9364_r9315_p4133

 mc16d:
 mc16_13TeV.364253.Sherpa_222_NNPDF30NNLO_lllv.deriv.DAOD_HIGG8D1.e5916_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.364250.Sherpa_222_NNPDF30NNLO_llll.deriv.DAOD_HIGG8D1.e5894_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.364254.Sherpa_222_NNPDF30NNLO_llvv.deriv.DAOD_HIGG8D1.e5916_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.364255.Sherpa_222_NNPDF30NNLO_lvvv.deriv.DAOD_HIGG8D1.e5916_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410155.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttW.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410156.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttZnunu.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410157.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttZqq.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410218.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttee.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410219.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttmumu.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410220.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttautau.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410276.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttee_mll_1_5.deriv.DAOD_HIGG8D1.e6087_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410277.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttmumu_mll_1_5.deriv.DAOD_HIGG8D1.e6087_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410278.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttautau_mll_1_5.deriv.DAOD_HIGG8D1.e6087_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410397.MadGraphPythia8EvtGen_ttbar_wbee_MEN30LO_A14N23LO.deriv.DAOD_HIGG8D1.e6086_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410398.MadGraphPythia8EvtGen_ttbar_wbmumu_MEN30LO_A14N23LO.deriv.DAOD_HIGG8D1.e6086_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410399.MadGraphPythia8EvtGen_ttbar_wbtautau_MEN30LO_A14N23LO.deriv.DAOD_HIGG8D1.e6086_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410658.PhPy8EG_A14_ttchan_BW50_lept_top.deriv.DAOD_HIGG8D1.e6671_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410659.PhPy8EG_A14_ttchan_BW50_lept_antitop.deriv.DAOD_HIGG8D1.e6671_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410644.PowhegPythia8EvtGen_A14_singletop_schan_lept_top.deriv.DAOD_HIGG8D1.e6527_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410645.PowhegPythia8EvtGen_A14_singletop_schan_lept_antitop.deriv.DAOD_HIGG8D1.e6527_s3126_r10201_r10210_p4133
 mc16_13TeV.304014.MadGraphPythia8EvtGen_A14NNPDF23_3top_SM.deriv.DAOD_HIGG8D1.e4324_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410080.MadGraphPythia8EvtGen_A14NNPDF23_4topSM.deriv.DAOD_HIGG8D1.e4111_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410081.MadGraphPythia8EvtGen_A14NNPDF23_ttbarWW.deriv.DAOD_HIGG8D1.e4111_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.364100.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV0_70_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_s3126_r10201_r10210_p4133
 mc16_13TeV.364101.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV0_70_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_s3126_r10201_r10210_p4133
 mc16_13TeV.364101.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV0_70_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.364102.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV0_70_BFilter.deriv.DAOD_HIGG8D1.e5271_s3126_r10201_r10210_p4133
 mc16_13TeV.364103.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV0_140_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_s3126_r10201_r10210_p4133
 mc16_13TeV.364104.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV0_140_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_s3126_r10201_r10210_p4133
 mc16_13TeV.364106.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV140_280_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_s3126_r10201_r10210_p4133
 mc16_13TeV.364107.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV140_280_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_s3126_r10201_r10210_p4133
 mc16_13TeV.364108.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV140_280_BFilter.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.364108.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV140_280_BFilter.deriv.DAOD_HIGG8D1.e5271_s3126_r10201_r10210_p4133
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 mc16_13TeV.364111.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV280_500_BFilter.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10201_r10210_p4133

mc16_13TeV.364192.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV140_280_BFilter.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.364193.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV280_500_CVetoBVeto.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.364193.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV280_500_CVetoBVeto.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_s3136_r10201_r10210_p4133
 mc16_13TeV.364194.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV280_500_CFilterBVeto.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10201_r10210_p4133
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 mc16_13TeV.364194.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV280_500_CFilterBVeto.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_s3136_r10201_r10210_p4133
 mc16_13TeV.364195.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV280_500_BFilter.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10201_r10210_p4133
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 mc16_13TeV.364196.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV500_1000.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_s3136_r10201_r10210_p4133
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 mc16_13TeV.364522.Sherpa_222_NNPDF30NNLO_enugamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
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 mc16_13TeV.342284.Pythia8EvtGen_A14NNPDF23LO_WH125_inc.deriv.DAOD_HIGG8D1.e4246_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.342285.Pythia8EvtGen_A14NNPDF23LO_ZH125_inc.deriv.DAOD_HIGG8D1.e4246_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.341998.aMcAtNloHppEG_UEEE5_CTEQ6L1_CT10ME_tWH125_gamgam_yt_plus1.deriv.DAOD_HIGG8D1.e4394_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410470.PhPy8EG_A14_ttbar_hdamp258p75_nonallhad.deriv.DAOD_HIGG8D1.e6337_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.410472.PhPy8EG_A14_ttbar_hdamp258p75_dil.deriv.DAOD_HIGG8D1.e6348_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.346343.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_allhad.deriv.DAOD_HIGG8D1.e7148_e5984_s3126_r10201_r10210_p4133
 mc16_13TeV.346344.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_semilep.deriv.DAOD_HIGG8D1.e7148_e5984_a875_r10201_r10210_p4133
 mc16_13TeV.346345.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_dilep.deriv.DAOD_HIGG8D1.e7148_e5984_a875_r10201_r10210_p4133

mc16e:

mc16_13TeV.361605.PowhegPy8EG_CT10nloME_AZNLOCTEQ6L1_ZZvvvv_mll4.deriv.DAOD_HIGG8D1.e4054_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.361602.PowhegPy8EG_CT10nloME_AZNLOCTEQ6L1_WZlvvv_mll4.deriv.DAOD_HIGG8D1.e4054_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.361601.PowhegPy8EG_CT10nloME_AZNLOCTEQ6L1_WZlvvv_mll4.deriv.DAOD_HIGG8D1.e4475_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.361600.PowhegPy8EG_CT10nloME_AZNLOCTEQ6L1_WWlvlv.deriv.DAOD_HIGG8D1.e4616_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.361603.PowhegPy8EG_CT10nloME_AZNLOCTEQ6L1_ZZlIII_mll4.deriv.DAOD_HIGG8D1.e4475_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.361604.PowhegPy8EG_CT10nloME_AZNLOCTEQ6L1_ZZvvll_mll4.deriv.DAOD_HIGG8D1.e4475_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364288.Sherpa_222_NNPDF30NNLO_llll_lowMIPtComplement.deriv.DAOD_HIGG8D1.e6096_s3126_r10724_p3983 mc16_13TeV.364287.Sherpa_222_NNPDF30NNLO_llvvjj_EW6.deriv.DAOD_HIGG8D1.e6055_e5984_s3126_r10724_r10726_p3983
 mc16_13TeV.364285.Sherpa_222_NNPDF30NNLO_llvvjj_EW6.deriv.DAOD_HIGG8D1.e6055_s3126_r10724_p3983
 mc16_13TeV.364284.Sherpa_222_NNPDF30NNLO_llvvjj_EW6.deriv.DAOD_HIGG8D1.e6055_s3126_r10724_p3983
 mc16_13TeV.364283.Sherpa_222_NNPDF30NNLO_lllljj_EW6.deriv.DAOD_HIGG8D1.e6055_s3126_r10724_p3983
 mc16_13TeV.364739.MGPy8EG_NNPDF30NLO_A14NNPDF23LO_lvlljjEW6_OFMinus.deriv.DAOD_HIGG8D1.e7421_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364742.MGPy8EG_NNPDF30NLO_A14NNPDF23LO_lvlljjEW6_SFPlus.deriv.DAOD_HIGG8D1.e7421_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364740.MGPy8EG_NNPDF30NLO_A14NNPDF23LO_lvlljjEW6_OFPlus.deriv.DAOD_HIGG8D1.e7421_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364741.MGPy8EG_NNPDF30NLO_A14NNPDF23LO_lvlljjEW6_SFMinus.deriv.DAOD_HIGG8D1.e7421_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364253.Sherpa_222_NNPDF30NNLO_llll.deriv.DAOD_HIGG8D1.e5916_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364250.Sherpa_222_NNPDF30NNLO_llll.deriv.DAOD_HIGG8D1.e5894_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364254.Sherpa_222_NNPDF30NNLO_llvv.deriv.DAOD_HIGG8D1.e5916_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364255.Sherpa_222_NNPDF30NNLO_llvv.deriv.DAOD_HIGG8D1.e5916_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410155.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttW.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410156.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttZnunu.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410157.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttZqq.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410218.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttee.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410219.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttmumu.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410220.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttautau.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410276.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttee_mll_1_5.deriv.DAOD_HIGG8D1.e6087_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410277.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttmumu_mll_1_5.deriv.DAOD_HIGG8D1.e6087_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410278.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttautau_mll_1_5.deriv.DAOD_HIGG8D1.e6087_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410397.MadGraphPythia8EvtGen_ttbar_wbee_MEN30LO_A14N23LO.deriv.DAOD_HIGG8D1.e6086_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410398.MadGraphPythia8EvtGen_ttbar_wbmumu_MEN30LO_A14N23LO.deriv.DAOD_HIGG8D1.e6086_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410399.MadGraphPythia8EvtGen_ttbar_wbtautau_MEN30LO_A14N23LO.deriv.DAOD_HIGG8D1.e6086_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410658.PhPy8EG_A14_tchan_BW50_lept_top.deriv.DAOD_HIGG8D1.e6671_e5984_s3126_s3136_r10724_r10726_p4133
 mc16_13TeV.410659.PhPy8EG_A14_tchan_BW50_lept_antitop.deriv.DAOD_HIGG8D1.e6671_e5984_s3126_s3136_r10724_r10726_p4133
 mc16_13TeV.410644.PowhegPythia8EvtGen_A14_singletop_schan_lept_top.deriv.DAOD_HIGG8D1.e6527_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410645.PowhegPythia8EvtGen_A14_singletop_schan_lept_antitop.deriv.DAOD_HIGG8D1.e6527_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.304014.MadGraphPythia8EvtGen_A14NNPDF23_3top_SM.deriv.DAOD_HIGG8D1.e4324_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410080.MadGraphPythia8EvtGen_A14NNPDF23_4topSM.deriv.DAOD_HIGG8D1.e4111_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410081.MadGraphPythia8EvtGen_A14NNPDF23_ttbarWW.deriv.DAOD_HIGG8D1.e4111_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364100.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV_70_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_s3136_r10724_r10726_p4133
 mc16_13TeV.364100.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV_70_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364101.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV_70_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364102.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV_70_BFilter.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_s3136_r10724_r10726_p4133
 mc16_13TeV.364102.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV_70_BFilter.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364103.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV70_140_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364103.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV70_140_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_s3136_r10724_r10726_p4133
 mc16_13TeV.364104.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV70_140_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364104.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV70_140_CFiltBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_s3136_r10724_r10726_p4133
 mc16_13TeV.364104.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV70_140_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364106.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV140_280_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_s3136_r10724_r10726_p4133
 mc16_13TeV.364106.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV140_280_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364107.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV140_280_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364107.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV140_280_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_s3136_r10724_r10726_p4133
 mc16_13TeV.364108.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV140_280_BFilter.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364109.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV280_500_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364110.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV280_500_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364110.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV280_500_BFilter.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_s3136_r10724_r10726_p4133
 mc16_13TeV.364111.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV280_500_BFilter.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_s3136_r10724_r10726_p4133
 mc16_13TeV.364111.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV280_500_BFilter.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364112.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV500_1000.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364113.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV1000_E_CMS.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364113.Sherpa_221_NNPDF30NNLO_Zmmumu_MAXHTPTV1000_E_CMS.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_s3136_r10724_r10726_p4133
 mc16_13TeV.364114.Sherpa_221_NNPDF30NNLO_Zee_MAXHTPTV0_70_CVetoBVeto.deriv.DAOD_HIGG8D1.e5299_e5984_s3126_s3136_r10724_r10726_p4133

mc16_13TeV.364196.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV500_1000.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_s3136_r10724_r10726_p4133
 mc16_13TeV.364196.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV500_1000.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364197.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV1000_E_CMS.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_s3136_r10724_r10726_p4133
 mc16_13TeV.364197.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV1000_E_CMS.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364500.Sherpa_222_NNPDF30NNLO_eegamma_pty_7_15.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364501.Sherpa_222_NNPDF30NNLO_eegamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364502.Sherpa_222_NNPDF30NNLO_eegamma_pty_35_70.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
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 mc16_13TeV.364504.Sherpa_222_NNPDF30NNLO_eegamma_pty_140_E_CMS.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
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 mc16_13TeV.364509.Sherpa_222_NNPDF30NNLO_mumugamma_pty_140_E_CMS.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364511.Sherpa_222_NNPDF30NNLO_tautaugamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364512.Sherpa_222_NNPDF30NNLO_tautaugamma_pty_35_70.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364513.Sherpa_222_NNPDF30NNLO_tautaugamma_pty_70_140.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364514.Sherpa_222_NNPDF30NNLO_tautaugamma_pty_140_E_CMS.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364522.Sherpa_222_NNPDF30NNLO_enugamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364523.Sherpa_222_NNPDF30NNLO_enugamma_pty_35_70.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364524.Sherpa_222_NNPDF30NNLO_enugamma_pty_70_140.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364525.Sherpa_222_NNPDF30NNLO_enugamma_pty_140_E_CMS.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364527.Sherpa_222_NNPDF30NNLO_munugamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364528.Sherpa_222_NNPDF30NNLO_munugamma_pty_35_70.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364529.Sherpa_222_NNPDF30NNLO_munugamma_pty_70_140.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364530.Sherpa_222_NNPDF30NNLO_munugamma_pty_140_E_CMS.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364532.Sherpa_222_NNPDF30NNLO_taunugamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364533.Sherpa_222_NNPDF30NNLO_taunugamma_pty_35_70.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364534.Sherpa_222_NNPDF30NNLO_taunugamma_pty_70_140.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364535.Sherpa_222_NNPDF30NNLO_taunugamma_pty_140_E_CMS.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410560.MadGraphPythia8EvtGen_A14_iZ_4fl_tchan_noAllHad.deriv.DAOD_HIGG8D1.e5803_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410408.aMcAtNloPythia8EvtGen_tWZ_Ztoll_minDR1.deriv.DAOD_HIGG8D1.e6423_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364242.Sherpa_222_NNPDF30NNLO_WWW_3l3v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364243.Sherpa_222_NNPDF30NNLO_WWZ_4l2v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364244.Sherpa_222_NNPDF30NNLO_WWZ_2l4v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364245.Sherpa_222_NNPDF30NNLO_WZZ_5l1v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364246.Sherpa_222_NNPDF30NNLO_WZZ_3l3v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364247.Sherpa_222_NNPDF30NNLO_ZZZ_6l0v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364248.Sherpa_222_NNPDF30NNLO_ZZZ_4l2v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.364249.Sherpa_222_NNPDF30NNLO_ZZZ_2l4v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.342284.Pythia8EvtGen_A14NNPDF23LO_WH125_inc.deriv.DAOD_HIGG8D1.e4246_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.342285.Pythia8EvtGen_A14NNPDF23LO_ZH125_inc.deriv.DAOD_HIGG8D1.e4246_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.341998.aMcAtNloHppEG_UEEEES_CTEQ6L1_CT10ME_iWH125_gamgam_yt_plus1.deriv.DAOD_HIGG8D1.e4394_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410389.MadGraphPythia8EvtGen_A14NNPDF23_ttgamma_nonallhadronic.deriv.DAOD_HIGG8D1.e6155_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410470.PhPy8EG_A14_ttbar_hdamp258p75_nonallhad.deriv.DAOD_HIGG8D1.e6337_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.410472.PhPy8EG_A14_ttbar_hdamp258p75_dil.deriv.DAOD_HIGG8D1.e6348_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.346343.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_allhad.deriv.DAOD_HIGG8D1.e7148_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.346343.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_allhad.deriv.DAOD_HIGG8D1.e7148_e5984_a875_r10724_r10726_p4133
 mc16_13TeV.346344.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_semilep.deriv.DAOD_HIGG8D1.e7148_e5984_a875_r10724_r10726_p4133
 mc16_13TeV.346344.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_semilep.deriv.DAOD_HIGG8D1.e7148_e5984_s3126_r10724_r10726_p4133
 mc16_13TeV.346345.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_dilep.deriv.DAOD_HIGG8D1.e7148_e5984_a875_r10724_r10726_p4133
 mc16_13TeV.346345.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_dilep.deriv.DAOD_HIGG8D1.e7148_e5984_s3126_r10724_r10726_p4133

577 **References**

- 578 [1] M. Aaboud et al. ‘Observation of electroweak $W^\pm Z$ boson pair production in association
579 with two jets in pp collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector’. In: *Phys.*
580 *Lett.* B793 (2019), pp. 469–492. doi: [10.1016/j.physletb.2019.05.012](https://doi.org/10.1016/j.physletb.2019.05.012). arXiv:
581 [1812.09740 \[hep-ex\]](https://arxiv.org/abs/1812.09740).
- 582 [2] T. Gleisberg et al. ‘Event generation with SHERPA 1.1’. In: *JHEP* 02 (2009), p. 007. doi:
583 [10.1088/1126-6708/2009/02/007](https://doi.org/10.1088/1126-6708/2009/02/007). arXiv: [0811.4622 \[hep-ph\]](https://arxiv.org/abs/0811.4622).
- 584 [3] ATLAS Collaboration. *Electron efficiency measurements with the ATLAS detector using*
585 *the 2015 LHC proton–proton collision data*. ATLAS-CONF-2016-024. 2016. URL: <https://cds.cern.ch/record/2157687>.
- 587 [4] ATLAS Collaboration. ‘Measurement of the muon reconstruction performance of the
588 ATLAS detector using 2011 and 2012 LHC proton–proton collision data’. In: *Eur. Phys.*
589 *J. C* 74 (2014), p. 3130. doi: [10.1140/epjc/s10052-014-3130-x](https://doi.org/10.1140/epjc/s10052-014-3130-x). arXiv: [1407.3935 \[hep-ex\]](https://arxiv.org/abs/1407.3935).
- 591 [5] *Evidence for the associated production of the Higgs boson and a top quark pair with the*
592 *ATLAS detector*. Tech. rep. ATLAS-CONF-2017-077. Geneva: CERN, Nov. 2017. URL:
593 <https://cds.cern.ch/record/2291405>.
- 594 [6] ATLAS Collaboration. *Jet Calibration and Systematic Uncertainties for Jets Reconstructed*
595 *in the ATLAS Detector at $\sqrt{s} = 13$ TeV*. ATL-PHYS-PUB-2015-015. 2015. URL:
596 <https://cds.cern.ch/record/2037613>.
- 597 [7] ATLAS Collaboration. *Selection of jets produced in 13 TeV proton–proton collisions with*
598 *the ATLAS detector*. ATLAS-CONF-2015-029. 2015. URL: <https://cds.cern.ch/record/2037702>.
- 600 [8] ATLAS Collaboration. ‘Performance of pile-up mitigation techniques for jets in pp col-
601 lisions at $\sqrt{s} = 8$ TeV using the ATLAS detector’. In: *Eur. Phys. J. C* 76 (2016), p. 581.
602 doi: [10.1140/epjc/s10052-016-4395-z](https://doi.org/10.1140/epjc/s10052-016-4395-z). arXiv: [1510.03823 \[hep-ex\]](https://arxiv.org/abs/1510.03823).
- 603 [9] ATLAS Collaboration. *Performance of missing transverse momentum reconstruction with*
604 *the ATLAS detector in the first proton–proton collisions at $\sqrt{s} = 13$ TeV*. ATL-PHYS-
605 PUB-2015-027. 2015. URL: <https://cds.cern.ch/record/2037904>.
- 606 [10] P. S. A. Hoecker. ‘TMVA 4 Toolkit for Multivariate Data Analysis with ROOT’. In:
607 *arXiv:physics/0703039* (2013).
- 608 [11] F. Cardillo et al. ‘Measurement of the fiducial and differential cross-section of a top quark
609 pair in association with a Z boson at 13 TeV with the ATLAS detector’. In: ATL-COM-
610 PHYS-2019-334 (Apr. 2019). URL: <https://cds.cern.ch/record/2672207>.
- 611 [12] ATLAS Collaboration. ‘Luminosity determination in pp collisions at $\sqrt{s} = 7$ TeV using
612 the ATLAS detector at the LHC’. In: *Eur. Phys. J. C* 71 (2011), p. 1630. doi: [10.1140/epjc/s10052-011-1630-5](https://doi.org/10.1140/epjc/s10052-011-1630-5). arXiv: [1101.2185 \[hep-ex\]](https://arxiv.org/abs/1101.2185).

- 614 [13] G. Aad et al. ‘Jet energy resolution in proton-proton collisions at $\sqrt{s} = 7$ TeV recorded
615 in 2010 with the ATLAS detector’. In: *The European Physical Journal C* 73.3 (Mar.
616 2013), p. 2306. ISSN: 1434-6052. DOI: [10.1140/epjc/s10052-013-2306-0](https://doi.org/10.1140/epjc/s10052-013-2306-0). URL:
617 <https://doi.org/10.1140/epjc/s10052-013-2306-0>.
- 618 [14] A. Collaboration. ‘Performance of b -jet identification in the ATLAS experiment’. In:
619 *Journal of Instrumentation* 11.04 (2016), P04008. URL: <http://stacks.iop.org/1748-0221/11/i=04/a=P04008>.
620