



## ATLAS Note

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# 2 WZ + Heavy Flavor Production in pp collisions 3 at $\sqrt{s} = 13$ TeV

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The ATLAS Collaboration

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A measurement of WZ produced with an associated heavy flavor jet is performed using 140  
6  $\text{fb}^{-1}$  of proton-proton collision data at  $\sqrt{s} = 13$  TeV from the ATLAS experiment at the  
7 LHC. The measurement is performed in the fully leptonic decay mode,  $\text{WZ} \rightarrow \ell\bar{\nu}\ell\bar{\nu}$ . The  
8 cross-section of  $\text{WZ} + b\text{-jets}$  is measured to be  $X \pm X \pm X$ , while the cross-section of  $\text{WZ} +$   
9 charm is measured as  $X$ , with a correlation of  $X$  between the two processes.

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<sup>44</sup> **1 Changes and outstanding items**

<sup>45</sup> **1.1 Changelog**

<sup>46</sup> This is version 4

<sup>47</sup> **1.1.1 Changes relative to v3**

- <sup>48</sup> • Merged introduction into executive summary, including unblinding details and list of  
<sup>49</sup> SRs/CRs used
- <sup>50</sup> • listed ptag used (p4133), and release (AB 21.2.127)
- <sup>51</sup> • Included table reftab:xsecUnc listing x-sec uncertainties used
- <sup>52</sup> • Removed selection criteria listed in table 4 (QMisID, AmbiguityType) that were removed  
<sup>53</sup> from the analysis
- <sup>54</sup> • specified variable used to make truth jet flavor determination (HadronConeExclTruthLa-  
<sup>55</sup> belIID)
- <sup>56</sup> • fixed bug in MtLepMet calculation, updated selection/fits to account for this
- <sup>57</sup> • Included plots of MtLepMet and PtZ, swapped lep 1 and 2  $p_T$  plots for lep W and lep Z  
<sup>58</sup> plots
- <sup>59</sup> • updated tZ BDT training to reduce overfitting, updated plots to include error bars, feature  
<sup>60</sup> importance
- <sup>61</sup> • updated table 8 to clarify selection, fix the tZ\_BDT cut used
- <sup>62</sup> • replace a few broken ntuples which included large weight events
- <sup>63</sup> • include DL1r distribution for Z+jets and  $t\bar{t}$  VRs
- <sup>64</sup> • Expanded section on fakes, included information on derived scale factors from VRs.
- <sup>65</sup> • Changed the kinematic plots to include  $p_T(Z)$  and  $m_T(W)$ , list lepton  $p_T$  based on W and  
<sup>66</sup> Z candidates.

67 **1.1.2 Changes relative to v2**

- 68     • Added alternate VBS samples to include missing b-jet diagrams  
69     • Included a section on tZ interference effects, ??.  
70     • Updated to reflect changes for 2018, including the move to PFlow jets, DL1r, updated  
71       trigger, and updated AnalysisBase version (now 21.2.127)  
72     • Revised fit regions, using separate 1-jet and 2-jet fits, with all 2-j regions included  
73     • updated plots for tZ BDT, added details about the model  
74     • Included truth jet information

75 **1.1.3 Changes relative to v1**

- 76     • Added GRL list  
77     • Fixed latex issue in line 92, typo in line 172  
78     • Added tables 6 and 4, summarizing the event and object selection  
79     • Added table 2, which includes the DSID of samples used  
80     • Included reference to WZ inclusive paper in introduction

81 **1.2 Outstanding Items**

- 82     • Complete interference studies, apply any interference effects observed as a systematic  
83     • Update results section with additional studies, possibly including:  
84       – Truth jet migration studies  
85       – Simultaneous fit over 1j and 2j  
86       – Impact of allowing tZ to float  
87     • Unblind, update plots and fits to include data  
88     • Add cross-section, significance once unblinded

---

## 89 **2 Executive Summary**

90 The production of  $WZ$  in association with a heavy flavor jet represents an important background  
 91 for many major analyses. This includes any process with leptons and b-jets in the final state,  
 92 such as  $t\bar{t}H$ ,  $t\bar{t}W$ , and  $t\bar{t}Z$ . While precise measurements have been made of  $WZ$  production  
 93 [1],  $WZ +$  heavy flavor remains poorly understood. This is largely because the QCD processes  
 94 involved in the production of the b-jet make it difficult to simulate accurately. This introduces a  
 95 large uncertainty for analyses that include this process as a background.

96 Motivated by its relevance to the  $t\bar{t}H$  multilepton analysis, we perform a study of the fully  
 97 leptonic decay mode of this channel; that is, events where both the W and Z decay leptonically.  
 98 Because  $WZ$  has no associated jets at leading order, while the major backgrounds for this channel  
 99 tend to have high jet multiplicity, events with more than two jets are rejected. This gives a final  
 100 state signature of three leptons and one or two jets.

101 Events that meet this selection criteria are sorted into pseudo-continuous b-tagging regions based  
 102 on the DL1r b-tag score of their associated jets. This is done to separate  $WZ +$  b-jet events  
 103 from  $WZ +$  charm and  $WZ +$  light jets. These orthogonal categories, in addition to a tZ Control  
 104 Region formed using an MVA, are the Signal Regions for the analysis. These regions are fit to  
 105 data in order make a more accurate estimate of the contribution of  $WZ +$  heavy-flavor, where  
 106 heavy-flavor jets include b-jets and charm jets. Separate fits are performed for 1-jet and 2-jet  
 107 events. The full Run-2 dataset collected by the ATLAS detector, representing  $139 \text{ fb}^{-1}$  of data  
 108 from pp collisions at  $\sqrt{s} = 13 \text{ TeV}$ , is used for this study.

109 All backgrounds are accounted for using Monte Carlo, with the simulation of non-prompt lepton  
 110 backgrounds -  $Z+jets$  and  $t\bar{t}$  - validated using non-prompt Validation Regions.

111 Section 3 details the data and Monte Carlo (MC) samples used in the analysis. The reconstruction  
 112 of various physics objects is described in section 4. Section 5 describes the event selection applied  
 113 to these samples, along the definitions of the various regions used in the fit. The multivariate  
 114 analysis techniques used to separate the tZ background from  $WZ +$  heavy flavor are described in  
 115 section 6. Section 7 describes the various sources of systematic uncertainties considered in the  
 116 fit. Finally, the results of the analysis are summarized in section 8, followed by a brief conclusion  
 117 in section 9.

118 The current state of the analysis shows blinded results for the full 2018 dataset. Regions containing  
 119  $>5\%$   $WZ+b$  events are blinded, and results are from Asimov, MC only fits. In addition to adding  
 120 some additional information to this note, remaining tasks include performing  $WZ/tZ$  interference  
 121 studies, finalizing the presentation of results, and unblinding.

## 122 **3 Data and Monte Carlo Samples**

123 Both data and Monte Carlo samples used in this analysis were prepared in the xAOD format,  
 124 which was used to produce a Dx AOD sample in the HIGG8D1 derivation framework. The HIGG8D1

125 framework is designed for the  $t\bar{t}H$  multi-lepton analysis, which targets events with multiple  
 126 leptons as well as tau hadrons. This framework skims the dataset to remove unneeded variables  
 127 as well as entire events. Events are removed from the derivations that do not meet the following  
 128 selection:

- 129 • at least two light leptons within a range  $|\eta| < 2.6$ , with leading lepton  $p_T > 15$  GeV and  
 130 subleading lepton  $p_T > 5$  GeV
- 131 • at least one light lepton with  $p_T > 15$  GeV within a range  $|\eta| < 2.6$ , and at least two hadronic  
 132 taus with  $p_T > 15$  GeV.

133 Samples were then generated from these HIGG8D1 derivations with p-tag of p4134 using a  
 134 modified version of AnalysisBase version 21.2.127.

### 135 3.1 Data Samples

136 The study uses a sample of proton-proton collision data collected by the ATLAS detector from  
 137 2015 through 2018 at an energy of  $\sqrt{s} = 13$  TeV, which represents an integrated luminosity of  
 138  $139 \text{ fb}^{-1}$ . This data set was collected with a bunch crossing rate of 25 ns. All data used in this  
 139 analysis was verified by data quality checks, having been included in the following Good Run  
 140 Lists:

- 141 • data15\_13TeV.periodAllYear\_DetStatus-v79-repro20-02\_DQDefects-00-02-02  
 142 \_PHYS\_StandardGRL\_All\_Good\_25ns.xml
- 143 • data16\_13TeV.periodAllYear\_DetStatus-v88-pro20-21\_DQDefects-00-02-04  
 144 \_PHYS\_StandardGRL\_All\_Good\_25ns.xml
- 145 • data17\_13TeV.periodAllYear\_DetStatus-v97-pro21-13\_Unknown\_PHYS\_StandardGRL  
 146 \_All\_Good\_25ns\_Triggerno17e33prim.xml
- 147 • data18\_13TeV.periodAllYear\_DetStatus-v102-pro22-04\_Unknown\_PHYS\_StandardGRL  
 148 \_All\_Good\_25ns\_Triggerno17e33prim.xml

149 Runs included from the AllYear period containers are included.

### 150 3.2 Monte Carlo Samples

151 Several different generators were used to produce Monte Carlo simulations of the signal and  
 152 background processes. For all samples, the response of the ATLAS detector is simulated using  
 153 Geant4. The WZ signal samples are simulated using Sherpa 2.2.2 [2]. Specific information  
 154 about the Monte Carlo samples being used can be found in table 1. A list of the specific samples  
 155 used by data set ID is shown in table 2.

Table 1: The configurations used for event generation of signal and background processes, including the event generator, matrix element (ME) order, parton shower algorithm, and parton distribution function (PDF).

Process	Event generator	ME order	Parton Shower	PDF
WZ, VV	SHERPA 2.2.2	MEPS NLO	SHERPA	CT10
tZ	MG5_AMC	LO	PYTHIA 6	CTEQ6L1
t̄W	MG5_AMC (SHERPA 2.1.1)	NLO (LO multileg)	PYTHIA 8 (SHERPA)	NNPDF 3.0 NLO (NNPDF 3.0 NLO)
t̄(Z/γ* → ll)	MG5_AMC	NLO	PYTHIA 8	NNPDF 3.0 NLO
t̄H	MG5_AMC (MG5_AMC)	NLO (NLO)	PYTHIA 8 (HERWIG++)	NNPDF 3.0 NLO [Ball:2014uwa] (CT10 [ <b>ct10</b> ])
tHqb	MG5_AMC	LO	PYTHIA 8	CT10
tHW	MG5_AMC (SHERPA 2.1.1)	NLO (LO multileg)	HERWIG++ (SHERPA)	CT10 (NNPDF 3.0 NLO)
tWZ	MG5_AMC	NLO	PYTHIA 8	NNPDF 2.3 LO
t̄t, t̄t̄t̄	MG5_AMC	LO	PYTHIA 8	NNPDF 2.3 LO
t̄W+W-	MG5_AMC	LO	PYTHIA 8	NNPDF 2.3 LO
t̄t	Powheg-BOX v2 [ <b>powheggtt</b> ]	NLO	PYTHIA 8	NNPDF 3.0 NLO
t̄t̄γ	MG5_AMC	LO	PYTHIA 8	NNPDF 2.3 LO
s-, t-channel, Wt single top	Powheg-BOX v1 [ <b>powhegstp</b> ]	NLO	PYTHIA 6	CT10
qqVV, VVV Z → l+l-	SHERPA 2.2.1	MEPS NLO	SHERPA	NNPDF 3.0 NLO

## 4 Object Reconstruction

All regions defined in this analysis share a common lepton, jet, and overall event preselection. The selection applied to each physics object is detailed here; the event preselection, and the selection used to define the various fit regions, is described in section 5.

### 4.1 Trigger

Events are required to be selected by dilepton triggers, as summarized in table 3.

### 4.2 Light leptons

Electron candidates are reconstructed from energy clusters in the electromagnetic calorimeter that are associated with charged particle tracks reconstructed in the inner detector [3]. Electron candidates are required to have  $p_T > 10$  GeV and  $|\eta_{\text{cluster}}| < 2.47$ . Candidates in the transition region between different electromagnetic calorimeter components,  $1.37 < |\eta_{\text{cluster}}| < 1.52$ , are rejected. A multivariate likelihood discriminant combining shower shape and track information is used to distinguish real electrons from hadronic showers (fake electrons). To further reduce

Sample	DSID
WZ	364253, 364739-42
VV	364250, 364254, 364255, 363355-60, 364890
t̄W	410155
t̄Z	410156, 410157, 410218-20
low mass t̄Z	410276-8
Rare Top	410397, 410398, 410399
single Top	410658-9, 410644-5
three Top	304014
four Top	410080
t̄WW	410081
Z + jets	364100-41
low mass Z + jets	364198-215
W + jets	364156-97
Vγ	364500-35
tZ	410560
tW	410013-4
WtZ	410408
VVV	364242-9
VH	342284-5
WtH	341998
t̄tγ	410389
t̄t	410470
t̄tH	345873-5, 346343-5

Table 2: List of Monte Carlo samples by data set ID used in the analysis.

the non-prompt electron contribution, the track is required to be consistent with originating from the primary vertex; requirements are imposed on the transverse impact parameter significance ( $|d_0|/\sigma_{d_0}$ ) and the longitudinal impact parameter ( $|\Delta z_0 \sin \theta_\ell|$ ), as shown in table 4.

Muon candidates are reconstructed by combining inner detector tracks with track segments or full tracks in the muon spectrometer [4]. Muon candidates are required to have  $p_T > 10$  GeV and  $|\eta| < 2.5$ .

All leptons are required to be isolated, and pass a non-prompt BDT selection described in detail in [5].

### 4.3 Jets

Jets are reconstructed from calibrated topological clusters built from energy deposits in the calorimeters [6], using the anti- $k_t$  algorithm with a radius parameter  $R = 0.4$ . Jets with energy contributions likely arising from noise or detector effects are removed from consideration [7],

Dilepton triggers (2015)		
$\mu\mu$ (asymm.)	HLT_mu18_mu8noL1	
$ee$ (symm.)	HLT_2e12_lhloose_L12EM10VH	
$e\mu, \mu e$ ( $\sim$ symm.)	HLT_e17_lhloose_mu14	
Dilepton triggers (2016)		
$\mu\mu$ (asymm.)	HLT_mu22_mu8noL1	
$ee$ (symm.)	HLT_2e17_lhvloose_nod0	
$e\mu, \mu e$ ( $\sim$ symm.)	HLT_e17_lhloose_nod0_mu14	
Dilepton triggers (2017)		
$\mu\mu$ (asymm.)	HLT_mu22_mu8noL1	
$ee$ (symm.)	HLT_2e24_lhvloose_nod0	
$e\mu, \mu e$ ( $\sim$ symm.)	HLT_e17_lhloose_nod0_mu14	
Dilepton triggers (2018)		
$\mu\mu$ (asymm.)	HLT_mu22_mu8noL1	
$ee$ (symm.)	HLT_2e24_lhvloose_nod0	
$e\mu, \mu e$ ( $\sim$ symm.)	HLT_e17_lhloose_nod0_mu14	

Table 3: List of lowest  $p_T$ -threshold, un-prescaled dilepton triggers used for 2015-2018 data taking.

	$e$			$\mu$						
	L	$L^*$	T	L	$L^*$	T				
FixedCutLoose	No	Yes		No	Yes					
Non-prompt lepton BDT	No	Yes		No	Yes					
Identification	Loose	Tight		Loose	Medium					
Transverse impact parameter significance $ d_0 /\sigma_{d_0}$	$< 5$			$< 3$						
Longitudinal impact parameter $ z_0 \sin \theta $	$< 0.5$ mm									

Table 4: Loose (L), loose and minimally-isolated ( $L^*$ ), and tight (T) light lepton definitions.

and only jets satisfying  $p_T > 25$  GeV and  $|\eta| < 2.5$  are used in this analysis. For jets with  $p_T < 60$  GeV and  $|\eta| < 2.4$ , a jet-track association algorithm is used to confirm that the jet originates from the selected primary vertex, in order to reject jets arising from pileup collisions [8].

185 **4.4 B-tagged Jets**

186 In order to make a measurement of  $WZ +$  heavy flavor it is necessary to distinguish these events  
 187 from  $WZ +$  light jets. For this purpose, the DL1r b-tagging algorithm is used to distinguish  
 188 heavy flavor jets from lighter ones. The DL1r algorithm uses jet kinematics, particularly jet  
 189 vertex information, as input for a neural network which assigns each jet a score designed to  
 190 reflect how likely that jet is to have originated from a b-quark.

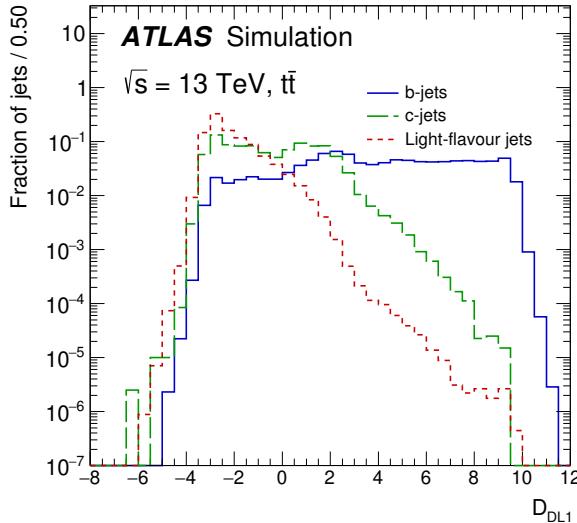


Figure 1: Output distribution of the DL1r algorithm for b-jets, charm jets, and light jets

191 From the output of the BDT, working points (WPs) are developed based on the efficiency of truth  
 192 b-jets at particular values of the DL1r algorithm. The working points used in this analysis are  
 193 summarized in table 5.

WP	none	loose	medium	tight	tightest
b eff.	-	85%	77%	70%	60%

Table 5: B-tagging Working Points by tightness and b-jet efficiency

194 A tighter WP will accept fewer b-jets, but reject a higher fraction of charm and light jets.  
 195 Generally, analyses that include b-jets will use a fixed working point, for example, requiring that  
 196 a jet pass the 70% threshold. By instead treating these working point as bins, e.g. events with  
 197 jets that fall between the 85% and 77% WPs fall into one bin, while events with jets passing the  
 198 60% WP fall into another, and looking at the full psuedo-continuous DL1r spectrum of the jets,  
 199 additional information can be gained. The psuedo-continuous b-tag spectrum is used in this case  
 200 to separate out  $WZ + b$ ,  $WZ + \text{charm}$ , and  $WZ + \text{light}$ .

---

201 **4.5 Missing transverse energy**

202 Missing transverse momentum ( $E_T^{\text{miss}}$ ) is used as part of the event selection. The missing  
 203 transverse momentum vector is defined as the inverse of the sum of the transverse momenta of  
 204 all reconstructed physics objects as well as remaining unclustered energy, the latter of which is  
 205 estimated from low- $p_T$  tracks associated with the primary vertex but not assigned to a hard object,  
 206 with object definitions taken from [9]. Light leptons considered in the  $E_T^{\text{miss}}$  reconstruction are  
 207 required to have  $p_T > 10$  GeV, while jets are required to have  $p_T > 20$  GeV.

208 **5 Event Selection and Signal Region Definitions**

209 Event are required to pass a preselection described in section 5.1 and summarized in table 6.  
 210 Those that pass this preselection are divided into various fit regions described in section 5.2,  
 211 based on the number of jets in the event, and the b-tag score of those jets.

212 **5.1 Event Preselection**

213 Events are required to include exactly three reconstructed light leptons passing the requirement  
 214 described in 4.2, which have a total charge of  $\pm 1$ . As the opposite sign lepton is found to  
 215 be prompt the vast majority of the time [5], it is required to be loose and isolated, as defined  
 216 though the standard `isolationFixedCutLoose` working point supported by combined per-  
 217 formance groups. The same sign leptons are required to be very tight, as per the recommended  
 218 `isolationFixedCutTight`.

219 The leptons are ordered in the analysis code as 0, 1, and 2. Lepton 0 is the lepton whose charge  
 220 is opposite the other two. Lepton 1 is the lepton closest to the opposite charge lepton, i.e. the  
 221 smallest  $\Delta R$ , leaving lepton 2 as the lepton further from the opposite charge lepton. Lepton 0  
 222 is required to have  $p_T > 10$  GeV, while the same sign leptons, 1 and 2, are required to have  
 223  $p_T > 20$  GeV to reduce the contribution of non-prompt leptons.

224 The invariant mass of at least one pair of opposite sign, same flavor leptons is required to fall  
 225 within 10 GeV of the mass of the Z boson, 91.2 GeV. Events where one of the opposite sign pairs  
 226 have an invariant mass less than 12 GeV are rejected in order to suppress low mass resonances.

227 An additional requirement is placed on the missing transverse energy,  $E_T^{\text{miss}} > 20$  GeV, and the  
 228 transverse mass of the W candidate,  $m(E_T^{\text{miss}} + l_{\text{other}}) > 30$  GeV, where  $E_T^{\text{miss}}$  is the missing  
 229 transverse energy, and  $l_{\text{other}}$  is the lepton not included in the Z-candidate.

230 Events are required to have one or two reconstructed jets passing the selection described in  
 231 section 4.3. Events with more than two jets are rejected in order to reduce the contribution of  
 232 backgrounds such as  $t\bar{t}Z$  and  $t\bar{t}W$ , which tend to have higher jet multiplicity.

Event Selection
Exactly three leptons with charge $\pm 1$
Two same-charge leptons with $p_T > 20 \text{ GeV}$
One opposite charge lepton with $p_T > 10 \text{ GeV}$
$m(l^+l^-)$ within 10 GeV of 91.2 GeV
Transverse mass of W-candidate, $m_T(E_T^{\text{miss}} + \text{lep}_{\text{other}}) > 30 \text{ GeV}$
Missing transverse energy, $E_T^{\text{miss}} > 20 \text{ GeV}$
One or two jets with $p_T > 25 \text{ GeV}$

Table 6: Summary of the selection applied to events for inclusion in the fit

233 The event yields in the preselection region for both data and Monte Carlo are summarized in  
 234 table 5.1, which shows good agreement between data and Monte Carlo, and demonstrates that  
 235 this region consists primarily of WZ events. The WZ events are split into WZ + b, WZ + c, and  
 236 WZ + 1 based on the truth flavor of the heaviest associated jet in the event. Specifically, this  
 237 determination is made based on the HadronConeExclTruthLabelID of the jet. That is, WZ +  
 238 1 events contain no charm and b jets at truth level, WZ + c contain at least one truth charm and  
 239 no b-jets, and WZ + b contains at least one truth b-jet.

Process	Events
$WZ + b$	$167.64 \pm 6.45$
$WZ + c$	$1080.91 \pm 39.28$
$WZ + l$	$7223.37 \pm 312.53$
Other VV	$849.79 \pm 142.13$
$t\bar{t}W$	$16.81 \pm 2.31$
$t\bar{t}Z$	$114.68 \pm 17.40$
rare Top	$2.20 \pm 0.14$
Single top	$0.10 \pm 0.45$
Three top	$0.01 \pm 0.01$
Four top	$0.02 \pm 0.01$
$t\bar{t}WW$	$0.23 \pm 0.05$
$Z + \text{jets}$	$601.16 \pm 260.13$
$V + \gamma$	$36.51 \pm 54.34$
$tZ$	$194.64 \pm 65.74$
$tW$	$5.49 \pm 1.24$
$WtZ$	$25.80 \pm 1.07$
$VVV$	$26.21 \pm 0.87$
$VH$	$94.34 \pm 7.35$
$t\bar{t}$	$107.68 \pm 8.14$
$t\bar{t}H$	$4.28 \pm 0.46$
Total	$10556.8 \pm 533.4$
Data	10574

Table 7: Events yields in the preselection region at  $138.9 \text{ fb}^{-1}$ 

<sup>240</sup> Here Other VV represents diboson processes other than WZ, and consists predominantly of  
<sup>241</sup>  $ZZ \rightarrow ll\bar{l}\bar{l}$  events where one of the leptons is not reconstructed.

<sup>242</sup> Simulations are further validated by comparing the kinematic distributions of the Monte Carlo  
<sup>243</sup> with data, which are shown in figures 2. Here, bins with 5% or more WZ+b are blinded.

### WZ Fit Region - Inclusive

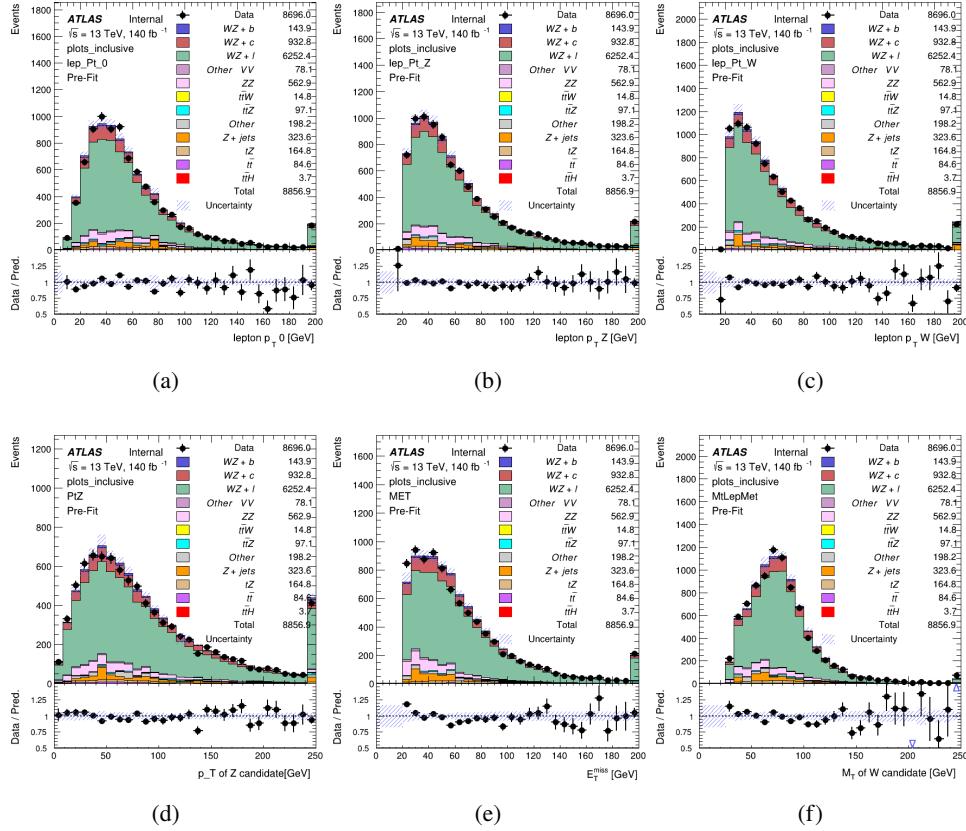


Figure 2: Comparisons between the data and MC distributions in the preselection region for (a) the  $p_T$  of the opposite sign lepton, (b) the  $p_T$  of the other lepton from the Z candidate, (c) the  $p_T$  of the lepton from the W candidate, (d) the  $p_T$  of the Z candidate, (e) the  $E_T^{\text{miss}}$ , and (f) the  $M_T$  of the W candidate.

## 5.2 Fit Regions

Once preselection has been applied, the remaining events are categorized into one of twelve orthogonal regions. The regions used in the fit are summarized in table 8.

The working points discussed in section 4.4 are used to separate events into fit regions based on the highest working point reached by a jet in each event. Because the background composition differs significantly based on the number of b-jets, events are further subdivided into 1 jet and 2 jet regions in order to minimize the impact of background uncertainties.

An additional tZ control region is created based on the BDT described in section 6. The region with 1-jet passing the 60% working point is split in two - a signal enriched region of events with a BDT score greater than 0.03, and a tZ control region including events with less than 0.03. This

Table 8: A list of the regions used in the fit and the selection used for each.

Region	Selection
1j, <85%	$N_{\text{jets}} = 1, n\text{Jets\_DL1r\_85} = 0$
1j, 85%-77%	$N_{\text{jets}} = 1, n\text{Jets\_DL1r\_85} = 1, n\text{Jets\_DL1r\_77}=0$
1j, 77%-70%	$N_{\text{jets}} = 1, n\text{Jets\_DL1r\_77} = 1, n\text{Jets\_DL1r\_70}=0$
1j, 70%-60%	$N_{\text{jets}} = 1, n\text{Jets\_DL1r\_70} = 1, n\text{Jets\_DL1r\_60}=0$
1j, >60%	$N_{\text{jets}} = 1, n\text{Jets\_DL1r\_60} = 1, tZ \text{ BDT} > 0.725$
1j tZ CR	$N_{\text{jets}} = 1, n\text{Jets\_DL1r\_60} = 1, tZ \text{ BDT} < 0.725$
2j, <85%	$N_{\text{jets}} = 2, n\text{Jets\_DL1r\_85} = 0$
2j, 85%-77%	$N_{\text{jets}} = 2, n\text{Jets\_DL1r\_85} >= 1, n\text{Jets\_DL1r\_77}=0$
2j, 77%-70%	$N_{\text{jets}} = 2, n\text{Jets\_DL1r\_77} >= 1, n\text{Jets\_DL1r\_70}=0$
2j, 70%-60%	$N_{\text{jets}} = 2, n\text{Jets\_DL1r\_70} >= 1, n\text{Jets\_DL1r\_60}=0$
2j, >60%	$N_{\text{jets}} = 2, n\text{Jets\_DL1r\_60} >= 1, tZ \text{ BDT} > 0.725$
2j tZ CR	$N_{\text{jets}} = 2, n\text{Jets\_DL1r\_60} >= 1, tZ \text{ BDT} < 0.725$

254 cutoff is arrived at by performing an Asimov fit with a variety of cutoffs, and selecting the value  
 255 that produces the highest significance for the measurement of  $WZ + b$ .

256 The modeling in each region is validated by comparing data and MC predictions for various  
 257 kinematic distributions. These plot are shown in figures 3-16.

## WZ Fit Region - 1j Inclusive

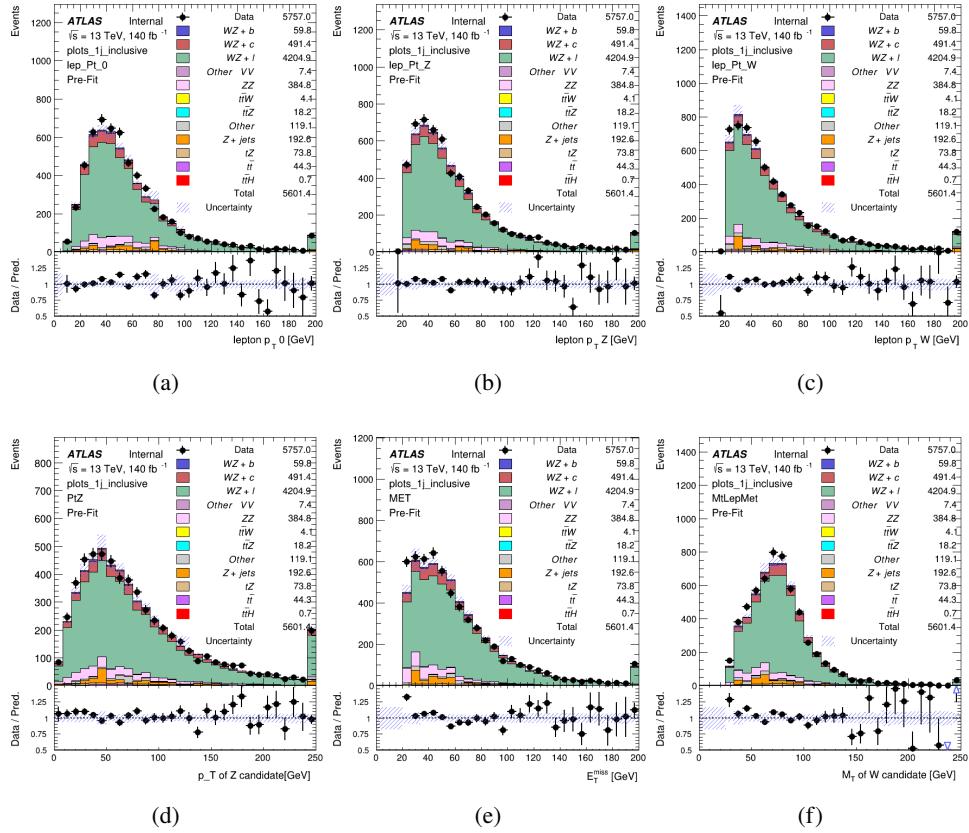


Figure 3: Comparisons between the data and MC distributions in the preselection region for (a) the  $p_T$  of the opposite sign lepton, (b) the  $p_T$  of the other lepton from the Z candidate, (c) the  $p_T$  of the lepton from the W candidate, (d) the  $p_T$  of the Z candidate, (e) the  $E_T^{\text{miss}}$ , and (f) the  $M_T$  of the W candidate.

## WZ Fit Region - 1j &lt; 85% WP

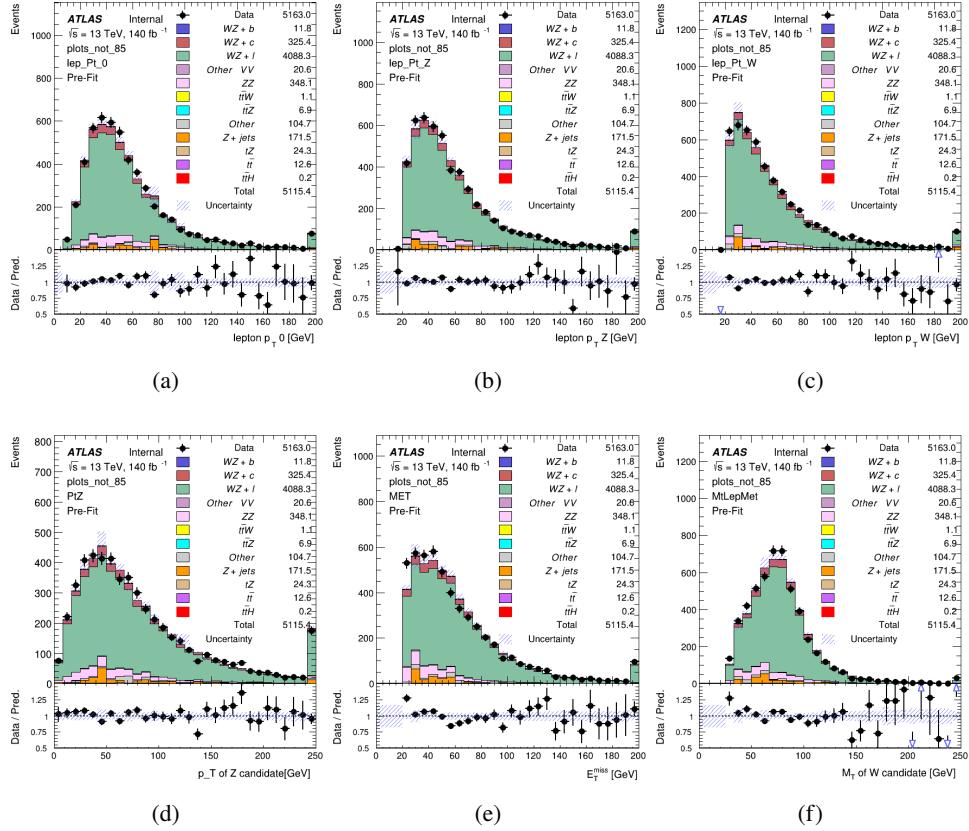


Figure 4: Comparisons between the data and MC distributions in the preselection region for (a) the  $p_T$  of the opposite sign lepton, (b) the  $p_T$  of the other lepton from the  $Z$  candidate, (c) the  $p_T$  of the lepton from the  $W$  candidate, (d) the  $p_T$  of the  $Z$  candidate, (e) the  $E_T^{\text{miss}}$ , and (f) the  $m_T$  of the  $W$  candidate.

## WZ Fit Region - 1j 77-85% WP

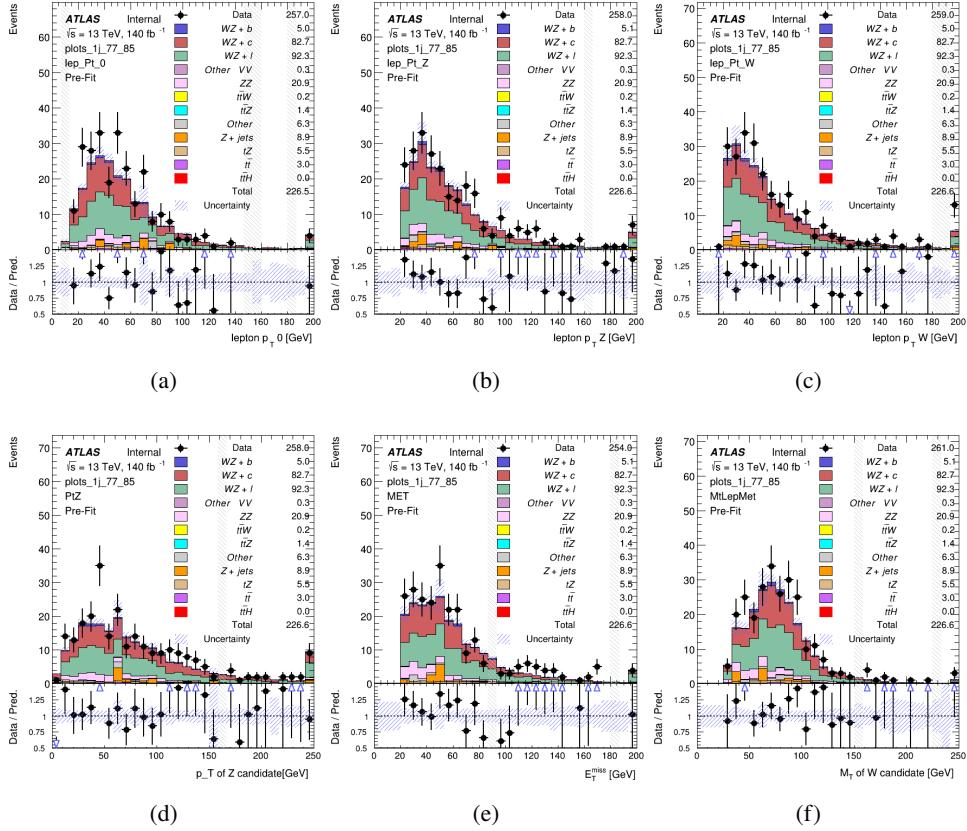


Figure 5: Comparisons between the data and MC distributions in the preselection region for (a) the  $p_T$  of the opposite sign lepton, (b) the  $p_T$  of the other lepton from the  $Z$  candidate, (c) the  $p_T$  of the lepton from the  $W$  candidate, (d) the  $p_T$  of the  $Z$  candidate, (e) the  $E_T^{\text{miss}}$ , and (f) the  $m_T$  of the  $W$  candidate.

## WZ Fit Region - 1j 70-77% WP

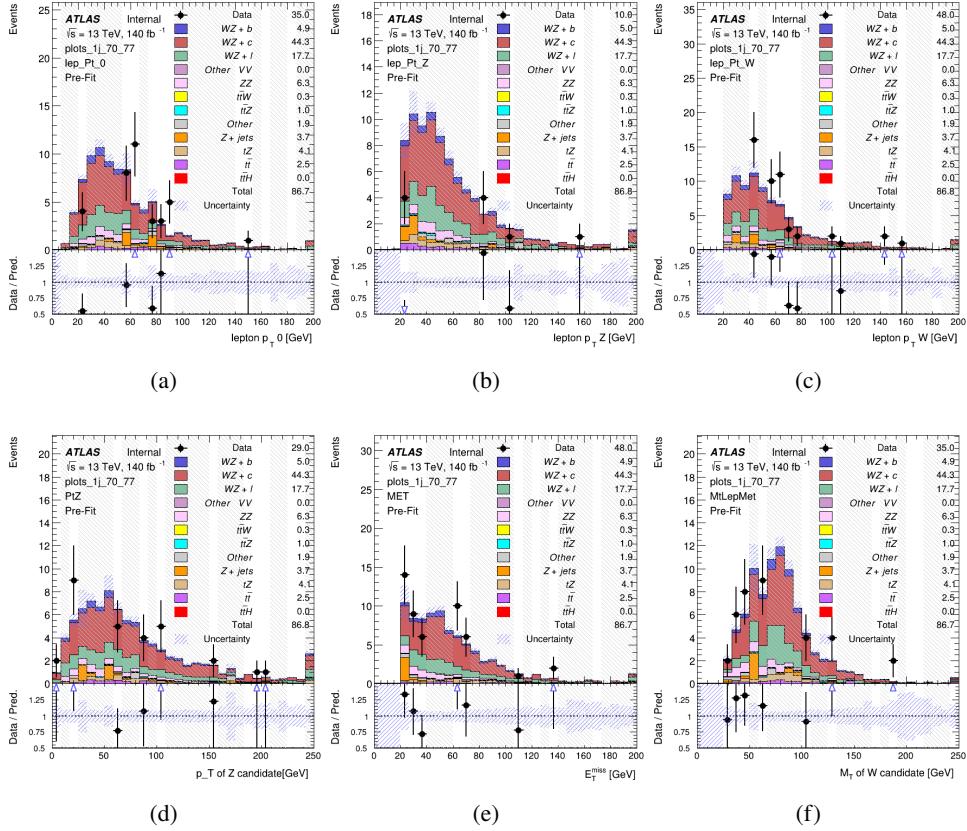


Figure 6: Comparisons between the data and MC distributions in the preselection region for (a) the  $p_T$  of the opposite sign lepton, (b) the  $p_T$  of the other lepton from the  $Z$  candidate, (c) the  $p_T$  of the lepton from the  $W$  candidate, (d) the  $p_T$  of the  $Z$  candidate, (e) the  $E_T^{\text{miss}}$ , and (f) the  $m_T$  of the  $W$  candidate.

## WZ Fit Region - 1j 60-70% WP

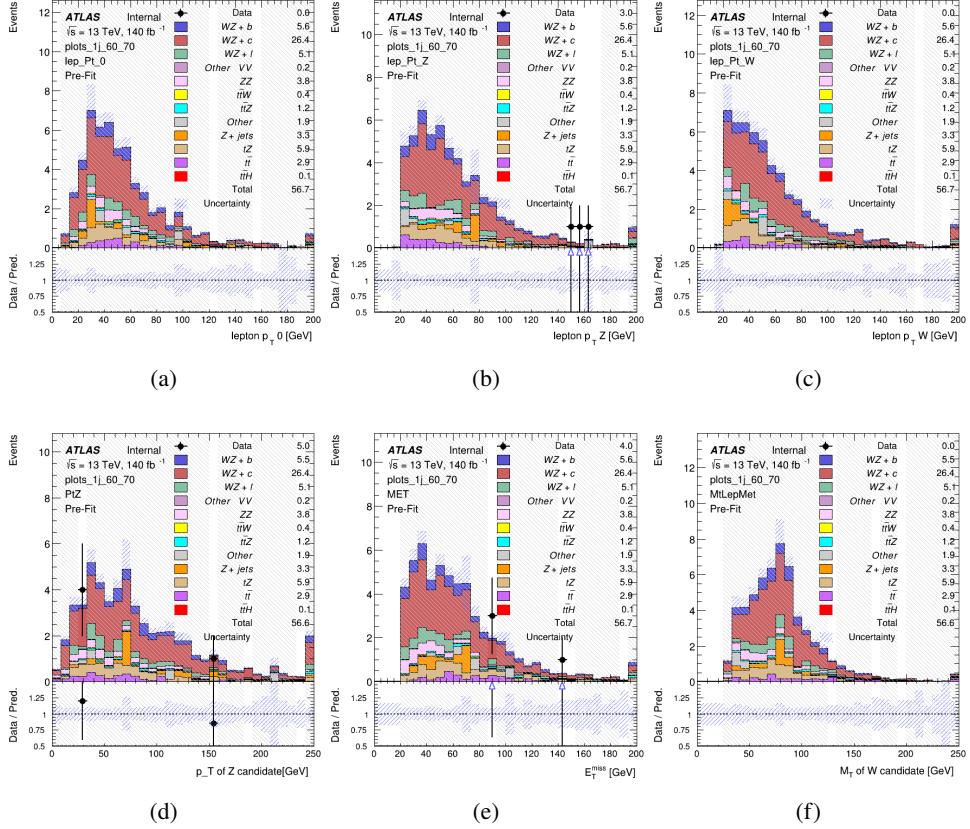


Figure 7: Comparisons between the data and MC distributions in the preselection region for (a) the  $p_T$  of the opposite sign lepton, (b) the  $p_T$  of the other lepton from the  $Z$  candidate, (c) the  $p_T$  of the lepton from the  $W$  candidate, (d) the  $p_T$  of the  $Z$  candidate, (e) the  $E_T^{\text{miss}}$ , and (f) the  $m_T$  of the  $W$  candidate.

## WZ Fit Region - 1j 60% WP

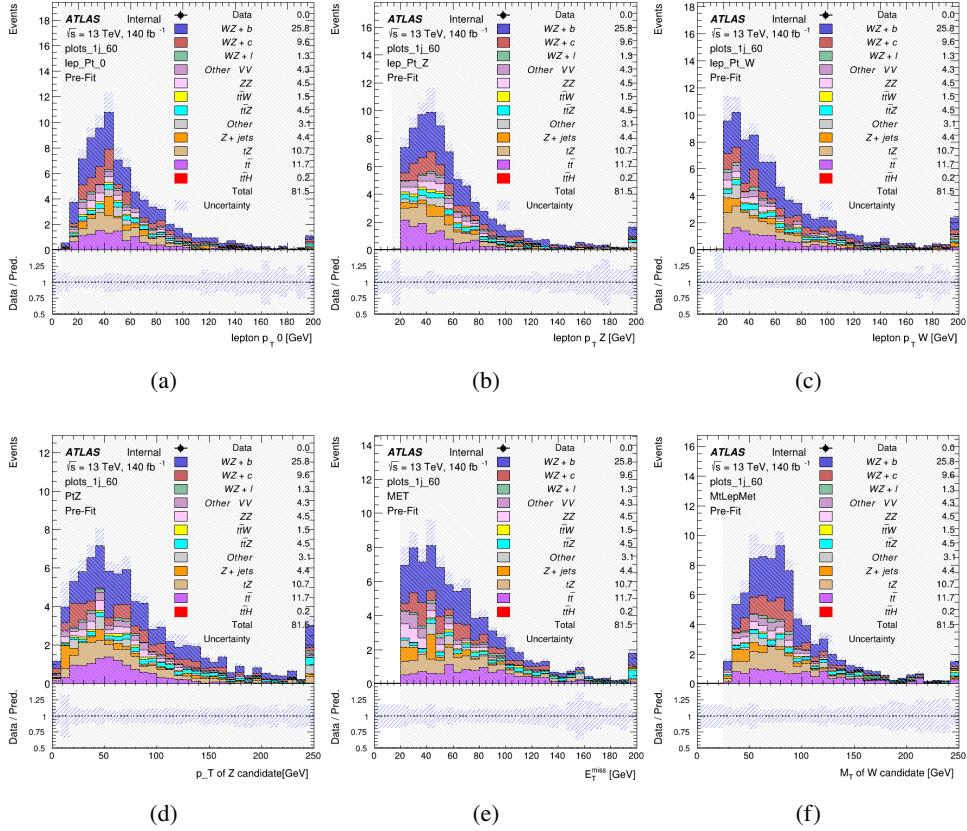


Figure 8: Comparisons between the data and MC distributions in the preselection region for (a) the  $p_T$  of the opposite sign lepton, (b) the  $p_T$  of the other lepton from the Z candidate, (c) the  $p_T$  of the lepton from the W candidate, (d) the  $p_T$  of the Z candidate, (e) the  $E_T^{\text{miss}}$ , and (f) the  $m_T$  of the W candidate.

## WZ Fit Region - tZ-CR

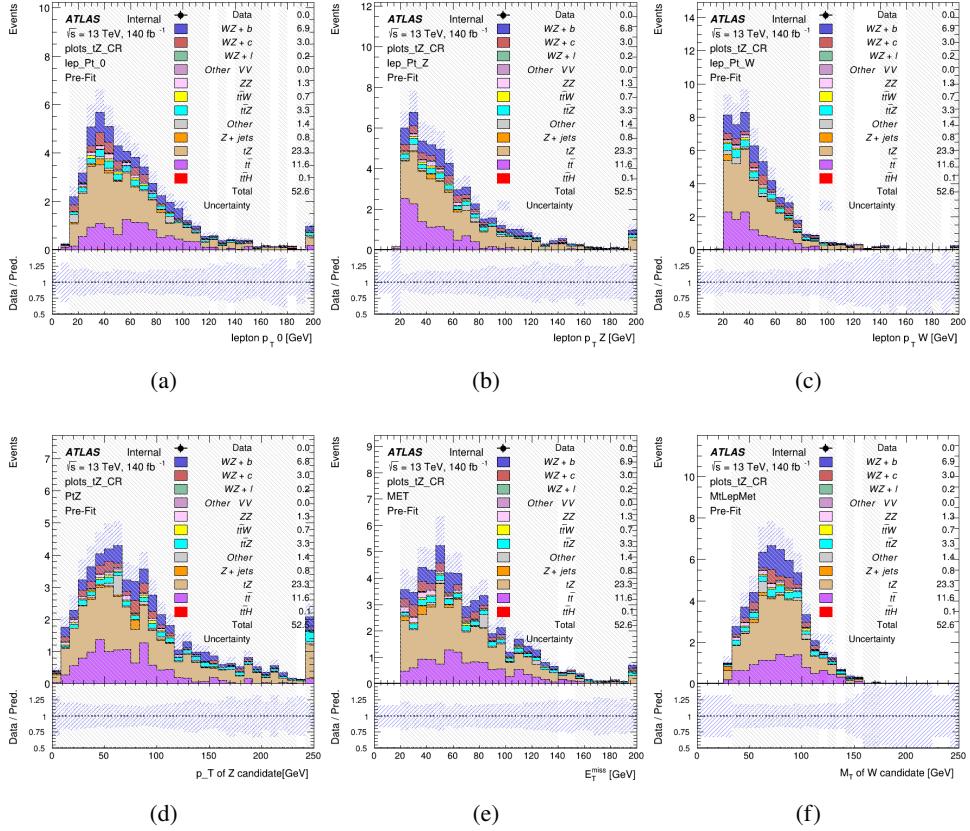


Figure 9: Comparisons between the data and MC distributions in the preselection region for (a) the  $p_T$  of the opposite sign lepton, (b) the  $p_T$  of the other lepton from the Z candidate, (c) the  $p_T$  of the lepton from the W candidate, (d) the  $p_T$  of the Z candidate, (e) the  $E_T^{\text{miss}}$ , and (f) the  $m_T$  of the W candidate.

### WZ Fit Region - 2j Inclusive

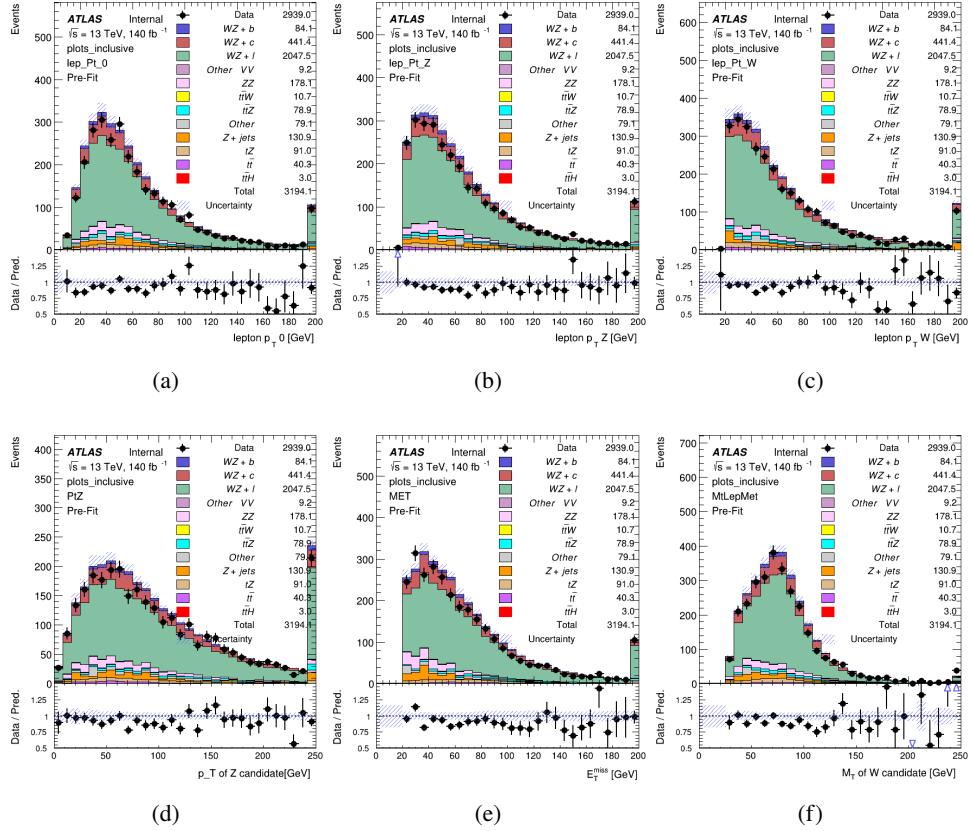


Figure 10: Comparisons between the data and MC distributions in the preselection region for (a) the  $p_T$  of the opposite sign lepton, (b) the  $p_T$  of the other lepton from the Z candidate, (c) the  $p_T$  of the lepton from the W candidate, (d) the  $p_T$  of the Z candidate, (e) the  $E_T^{\text{miss}}$ , and (f) the  $m_T$  of the W candidate.

## WZ Fit Region - 2j &lt; 85% WP

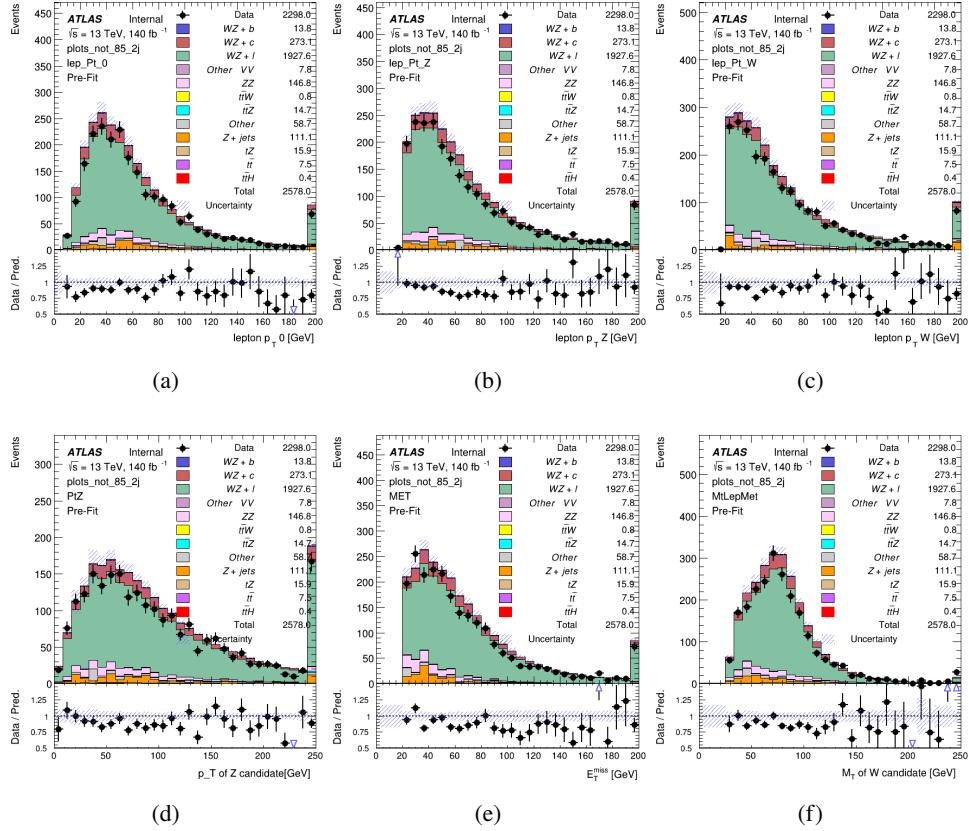


Figure 11: Comparisons between the data and MC distributions in the preselection region for (a) the  $p_T$  of the opposite sign lepton, (b) the  $p_T$  of the other lepton from the Z candidate, (c) the  $p_T$  of the lepton from the W candidate, (d) the  $p_T$  of the Z candidate, (e) the  $E_T^{\text{miss}}$ , and (f) the  $m_T$  of the W candidate.

## WZ Fit Region - 2j 77-85% WP

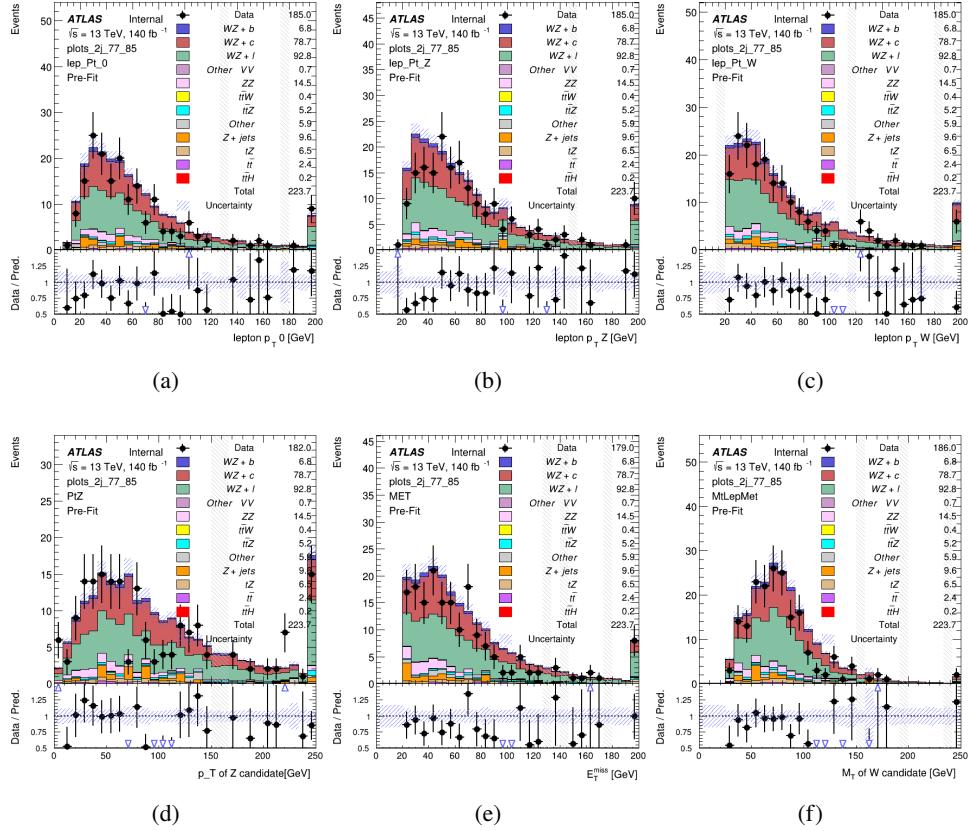


Figure 12: Comparisons between the data and MC distributions in the preselection region for (a) the  $p_T$  of the opposite sign lepton, (b) the  $p_T$  of the other lepton from the Z candidate, (c) the  $p_T$  of the lepton from the W candidate, (d) the  $p_T$  of the Z candidate, (e) the  $E_T^{\text{miss}}$ , and (f) the  $m_T$  of the W candidate.

## WZ Fit Region - 2j 70-77% WP

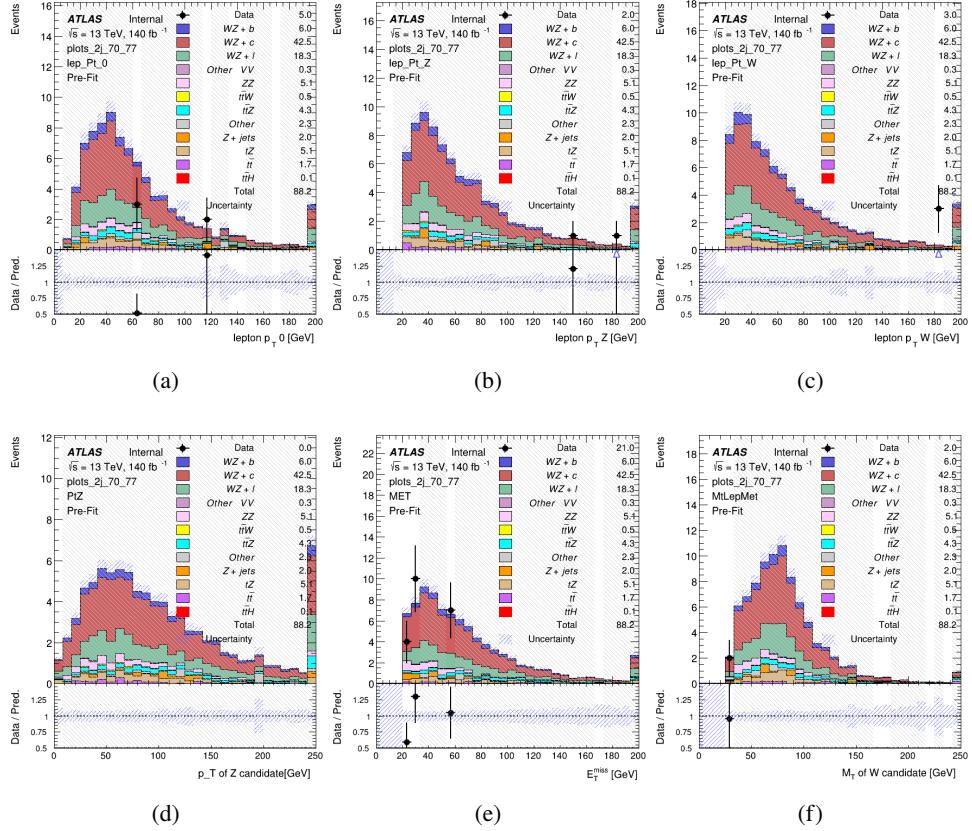


Figure 13: Comparisons between the data and MC distributions in the preselection region for (a) the  $p_T$  of the opposite sign lepton, (b) the  $p_T$  of the other lepton from the Z candidate, (c) the  $p_T$  of the lepton from the W candidate, (d) the  $p_T$  of the Z candidate, (e) the  $E_T^{\text{miss}}$ , and (f) the  $m_T$  of the W candidate.

## WZ Fit Region - 2j 60-70% WP

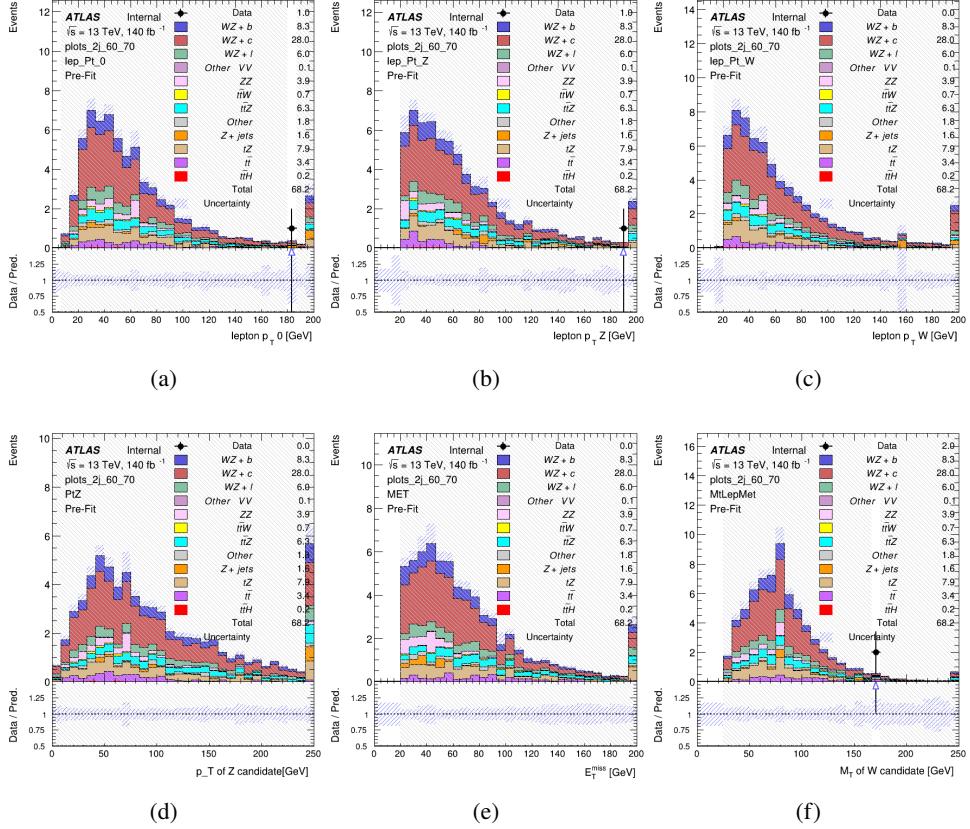


Figure 14: Comparisons between the data and MC distributions in the preselection region for (a) the  $p_T$  of the opposite sign lepton, (b) the  $p_T$  of the other lepton from the Z candidate, (c) the  $p_T$  of the lepton from the W candidate, (d) the  $p_T$  of the Z candidate, (e) the  $E_T^{\text{miss}}$ , and (f) the  $m_T$  of the W candidate.

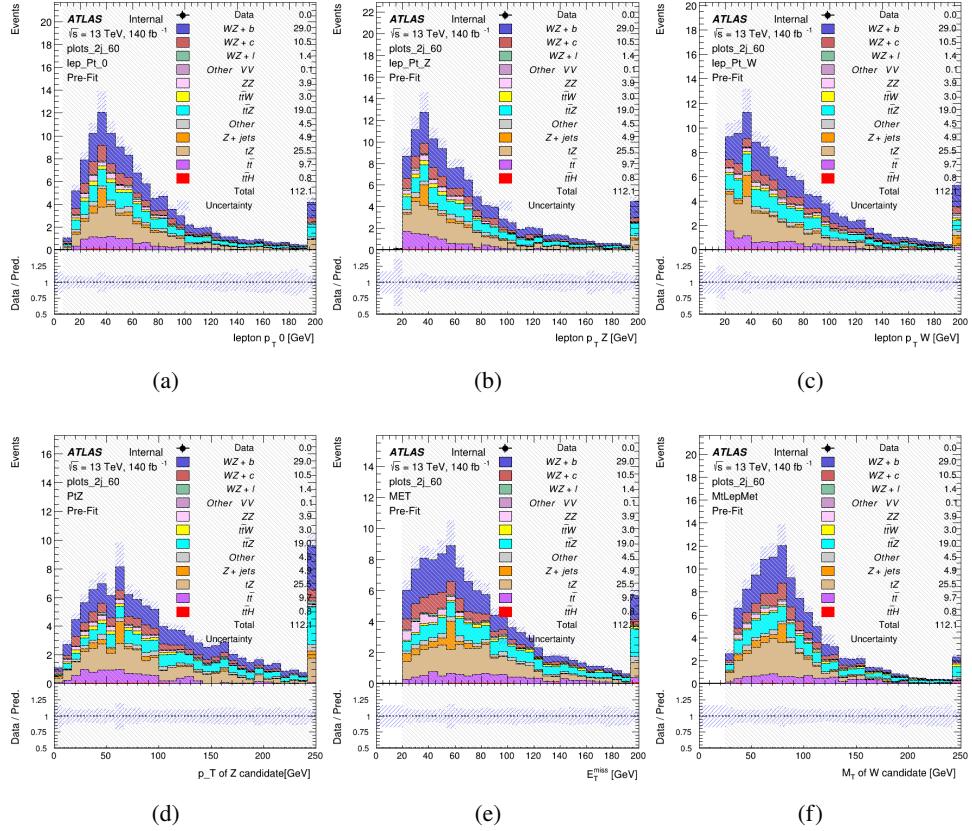
**WZ Fit Region - 2j 60% WP**

Figure 15: Comparisons between the data and MC distributions in the preselection region for (a) the  $p_T$  of the opposite sign lepton, (b) the  $p_T$  of the other lepton from the Z candidate, (c) the  $p_T$  of the lepton from the W candidate, (d) the  $p_T$  of the Z candidate, (e) the  $E_T^{\text{miss}}$ , and (f) the  $m_T$  of the W candidate.

## WZ Fit Region - tZ-CR-2j

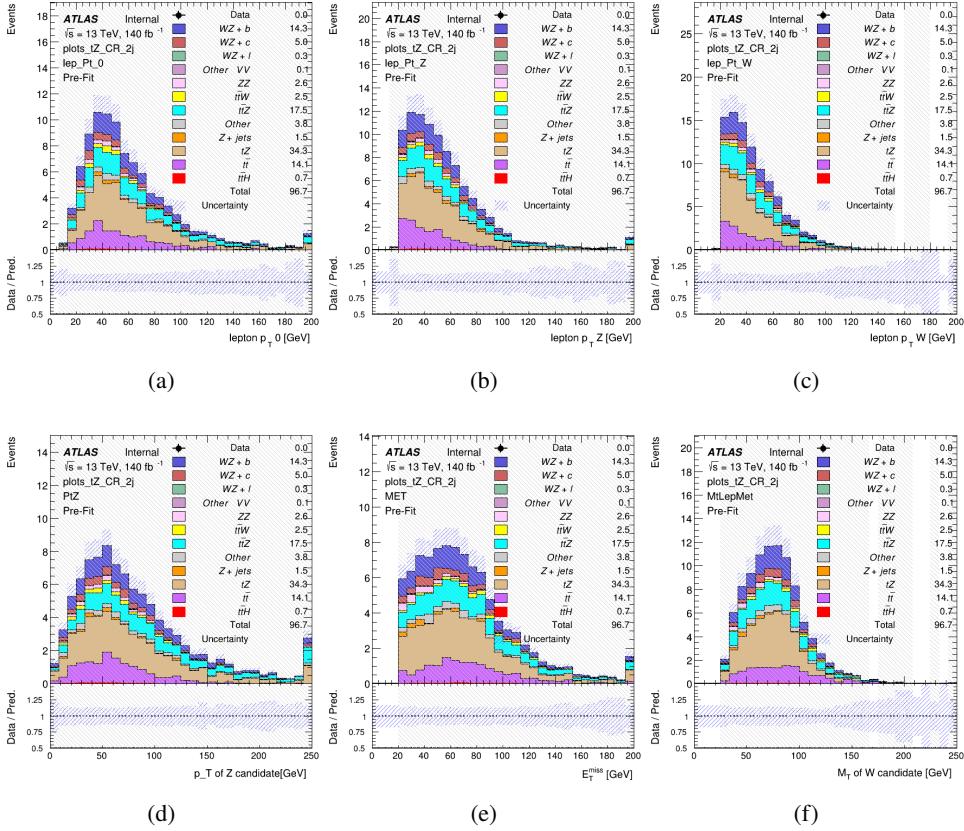


Figure 16: Comparisons between the data and MC distributions in the preselection region for (a) the  $p_T$  of the opposite sign lepton, (b) the  $p_T$  of the other lepton from the  $Z$  candidate, (c) the  $p_T$  of the lepton from the  $W$  candidate, (d) the  $p_T$  of the  $Z$  candidate, (e) the  $E_T^{\text{miss}}$ , and (f) the  $m_T$  of the  $W$  candidate.

### 5.3 Non-Prompt Lepton Estimation

Two processes act as sources of non-prompt leptons appear in the analysis:  $t\bar{t}$  and  $Z+\text{jet}$  production both produce two prompt leptons, and each contribute to the 3l region when an additional non-prompt lepton appears in the event. The contribution of these processes is estimated with Monte Carlo simulations, which are validated using enriched validation regions.

#### 5.3.1 $t\bar{t}$ Validation

$t\bar{t}$  events can produce two prompt leptons from the decay of each of the tops. These top decays produce two b-quarks, the decay of which can produce additional non-prompt leptons, which occasionally pass the event preselection. In order to validate that the Monte Carlo accurately

267 simulates this process accurately, the MC prediction in a non-prompt  $t\bar{t}$  enriched validation  
268 region is compared to data.

269 The  $t\bar{t}$  validation region is similar to the preselection region - three leptons meeting the criteria  
270 described in section 5 are required, and the requirements on  $E_T^{\text{miss}}$  remain the same. However,  
271 the selection requiring a lepton pair form a Z-candidate are reversed. Events where the invariant  
272 mass of any two opposite sign, same flavor leptons falls within 10 GeV of 91.2 GeV are rejected.  
273 This ensures the  $t\bar{t}$  validation region is orthogonal to the preselection region.

274 Further, because the jet multiplicity of  $t\bar{t}$  events tends to be higher than WZ, the number of jets  
275 in each event is required to be greater than 1. As b-jets are almost invariably produced from top  
276 decays, at least one b-tagged jet passing the 70% DL1r WP in each event is required. Various  
277 kinematic plots of this region are shown in figure 17.

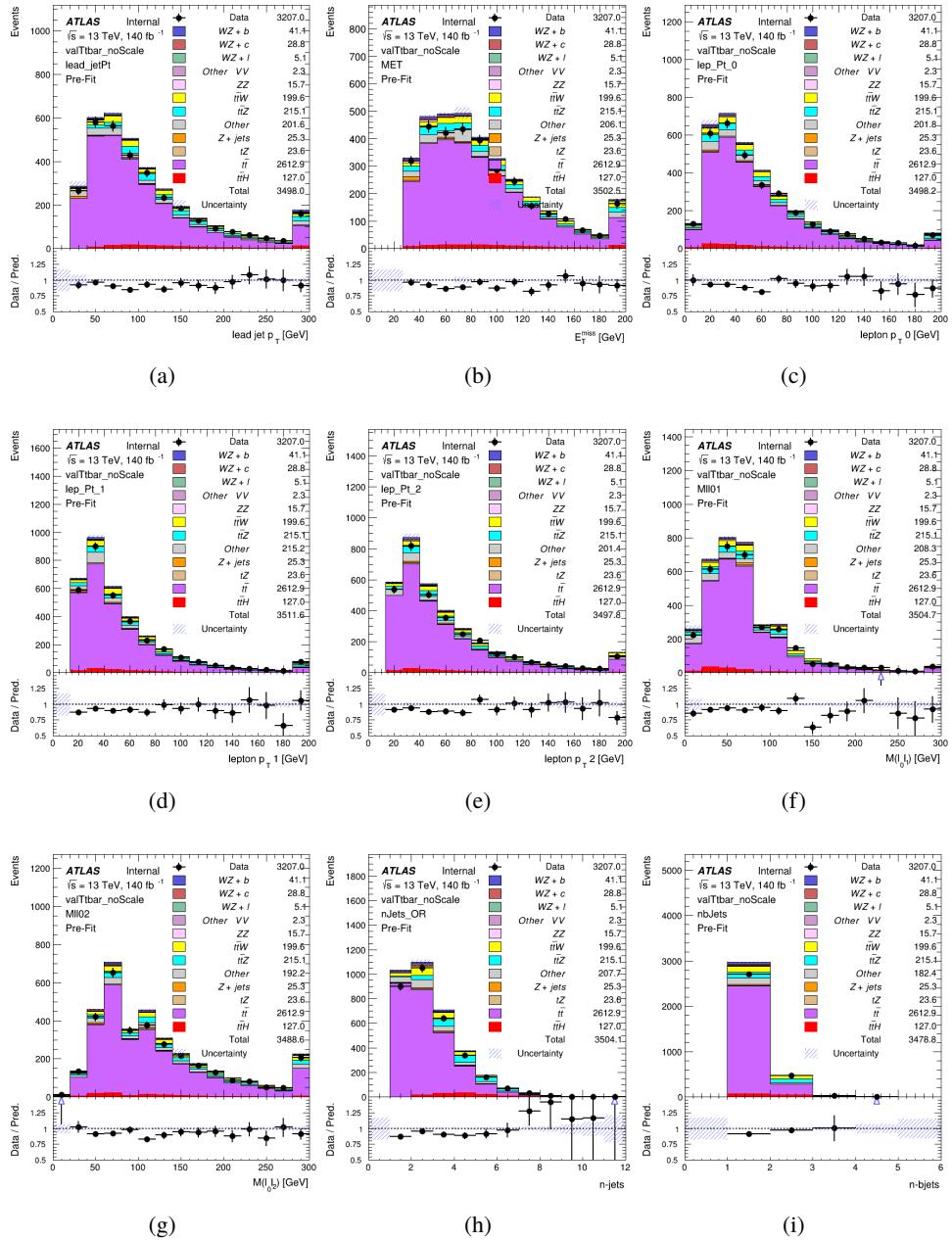


Figure 17: Comparisons between the data and MC distributions in the  $t\bar{t}$  validation region for (a) the  $p_T$  of the leading jet, (b) the missing transverse energy, (c) the  $p_T$  of lepton 0, (d)  $p_T$  of lepton 1, (e)  $p_T$  of lepton 2, (f) the invariant mass of leptons 0 and 1, (g) the invariant mass of leptons 0 and 2, (h) the number of jets, (i) the number of b-tagged jets.

278 The shape of each distribution agrees quite well between data and MC, with a constant offset  
 279 between the two. This is accounted for by applying a constant correction factor of 0.883 to the  $t\bar{t}$

280 MC prediction. Plots showing the kinematics of the  $t\bar{t}$  VR after this correction factor has been  
 281 applied are shown in figure 18.

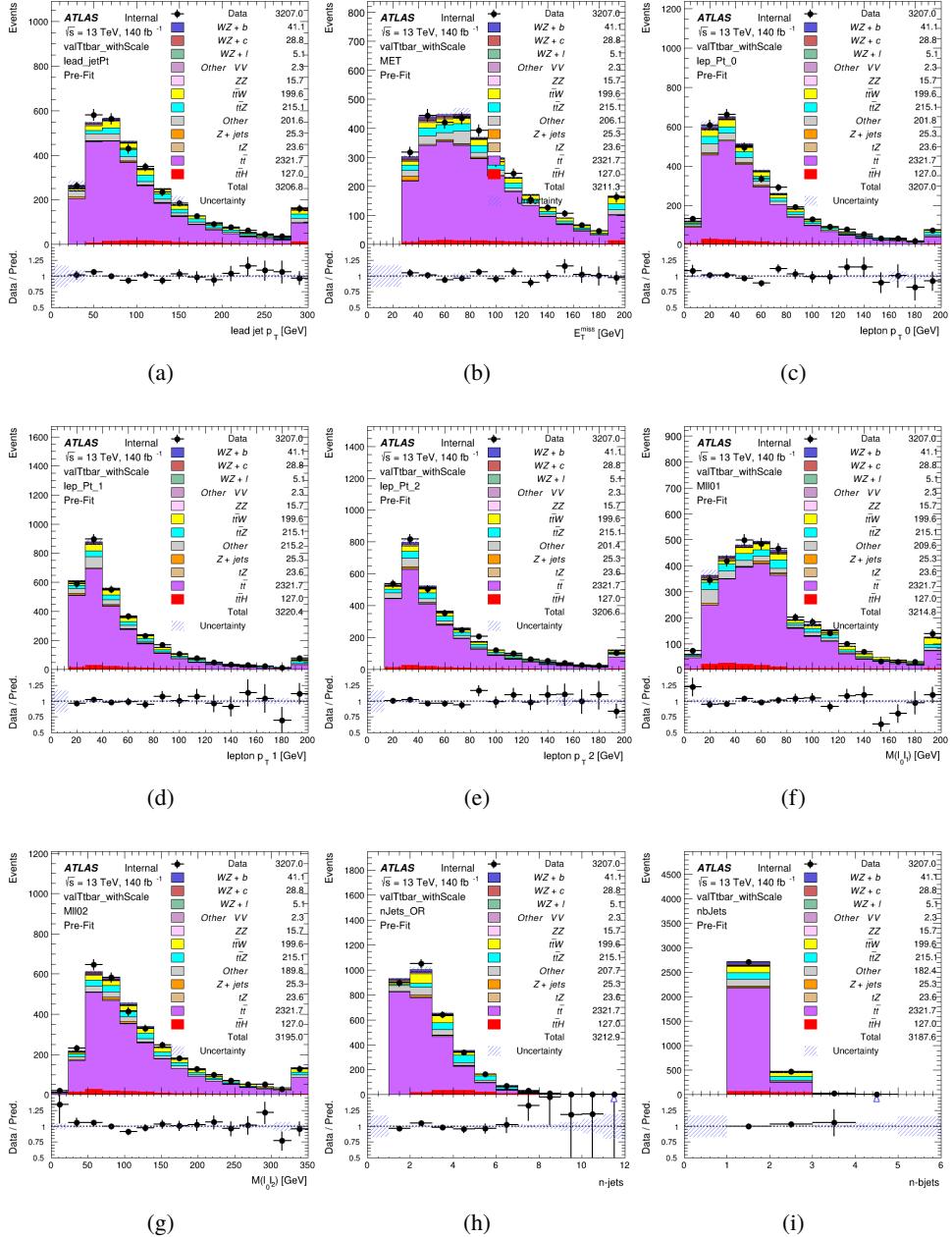


Figure 18: Comparisons between the data and MC distributions in the  $t\bar{t}$  validation region after the correction factor has been applied for (a) the  $p_T$  of the leading jet, (b) the missing transverse energy, (c) the  $p_T$  of lepton 0, (d)  $p_T$  of lepton 1, (e)  $p_T$  of lepton 2, (f) the invariant mass of leptons 0 and 1, (g) the invariant mass of leptons 0 and 2, (h) the number of jets, (i) the number of b-tagged jets.

282 The modeling is further validated by looking at the yield in the  $t\bar{t}$  VR for each DL1r WP, giving  
 283 a clearer correspondence to the signal regions used in the fit. Each region shown in figure 22  
 284 requires one or more jets pass the listed WP, with no jets passing the next highest WP.

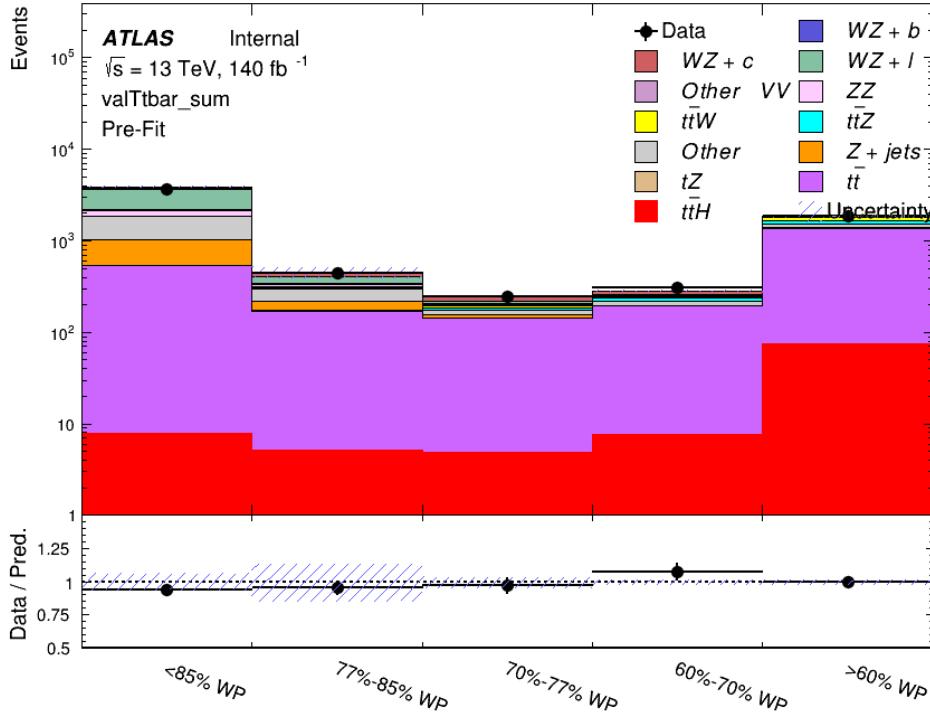


Figure 19: Data and MC comparisons for each DL1r WP for both 1-jet and 2-jet regions, after the  $t\bar{t}$  VR selection and correction factor have been applied

285 As data and MC are found to agree within 10% for each of these working points, a 10% systematic  
 286 uncertainty on the  $t\bar{t}$  prediction is included for the analysis.

### 287 5.3.2 Z+jets Validation

288 Similar to  $t\bar{t}$ , a non-prompt Z+jets validation region is produced in order to validate the MC  
 289 predictions. The lepton requirements remain the same as the preselection region. Because no  
 290 neutrinos are present for this process, the  $E_T^{\text{miss}}$  cut is reversed, requiring  $E_T^{\text{miss}} < 30 \text{ GeV}$ . This  
 291 also ensures this validation region is orthogonal to the preselection region. Further, the number  
 292 of jets in each event is required to be greater than or equal to one. Various kinematic plots of this  
 293 region are shown below. The general agreement between data and MC in each of these suggests  
 294 that the non-prompt contribution of Z+jets is well modeled by Monte Carlo.

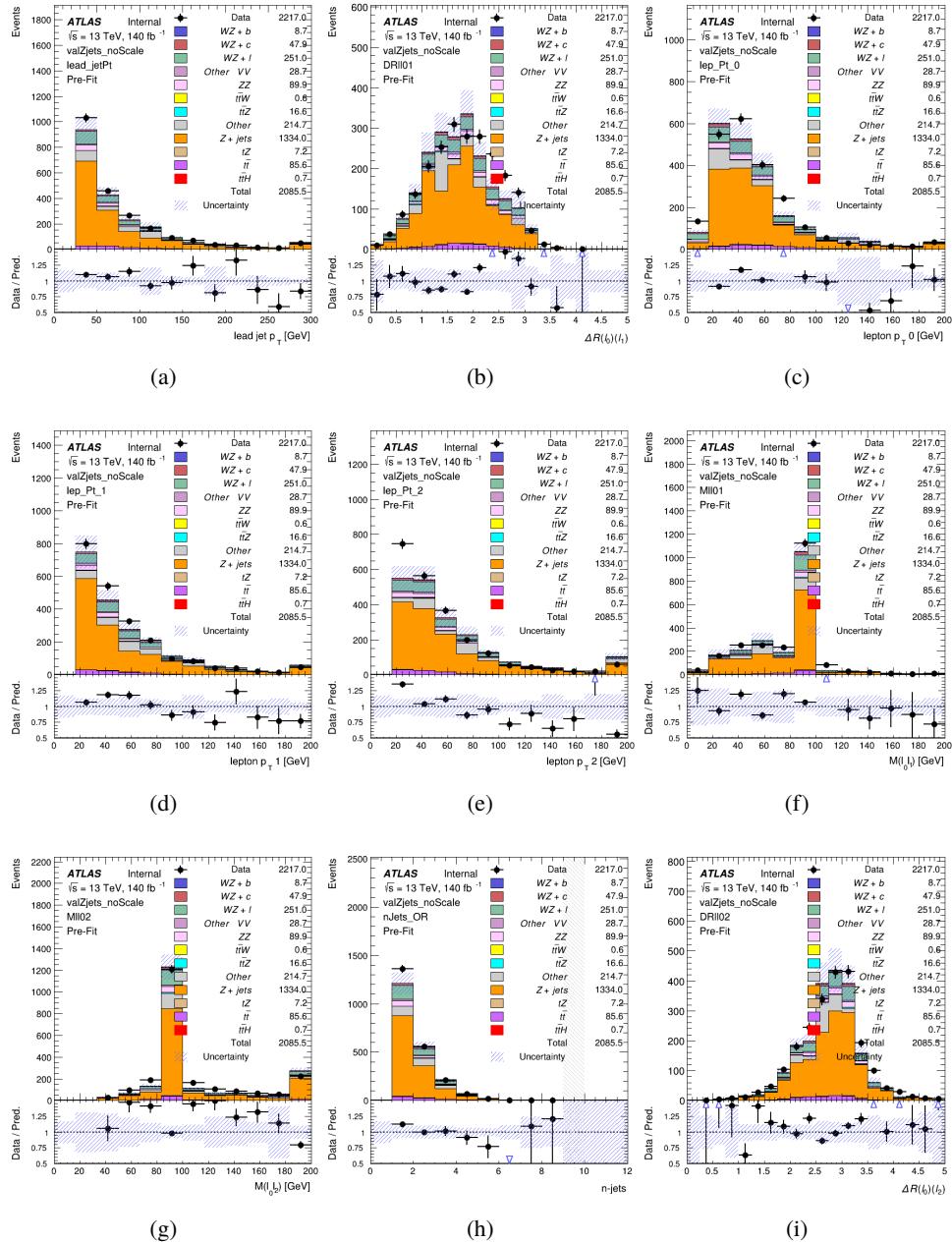


Figure 20: Comparisons between the data and MC distributions in the Z+jets validation region for (a) the  $p_T$  of the leading jet, (b)  $\Delta R$  between leptons 0 and 1, (c) the  $p_T$  of lepton 0, (d)  $p_T$  of lepton 1, (e)  $p_T$  of lepton 2, (f) the invariant mass of leptons 0 and 1, (g) the invariant mass of leptons 0 and 2, (h) the number of jets, (i)  $\Delta R$  between leptons 0 and 2. Includes only statistical uncertainties

295 While there is general agreement between data and MC within statistical uncertainty, the shape

296 of the  $p_T$  spectrum of lepton 2 is found to differ. To account for this discrepancy, a variable  
297 correction factor is applied to Z+jets.  $\chi^2$  minimization of the lepton 2  $p_T$  spectrum is performed  
298 to derive a correction factor of  $1.53 - 6.6 * 10^{-6}(\text{lep\_Pt\_2})$ . Kinematic plots of the Z + jets  
299 validation region after this correction factor has been aplied are shown in figure 21.

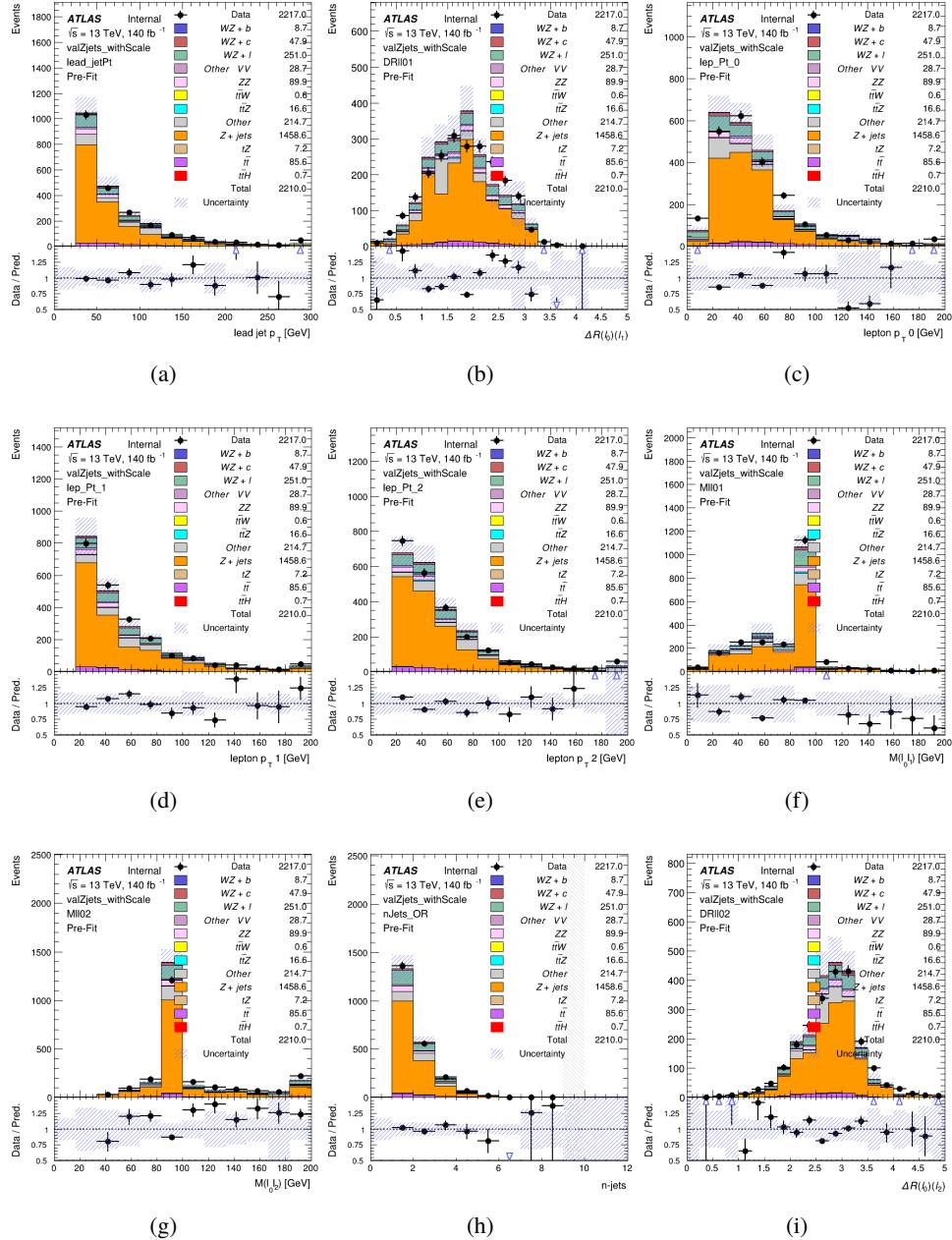


Figure 21: Comparisons between the data and MC distributions in the Z+jets validation region after the correction factor has been applied for (a) the  $p_T$  of the leading jet, (b)  $\Delta R$  between leptons 0 and 1, (c) the  $p_T$  of lepton 0, (d)  $p_T$  of lepton 1, (e)  $p_T$  of lepton 2, (f) the invariant mass of leptons 0 and 1, (g) the invariant mass of leptons 0 and 2, (h) the number of jets, (i)  $\Delta R$  between leptons 0 and 2

300 The modeling is further validated by looking at the yield in the Z+jets VR for each DL1r WP,

301 giving a clearer correspondence to the signal regions used in the fit. Each region shown in figure  
 302 22 requires one or more jets pass the listed WP, with no jets passing the next highest WP.

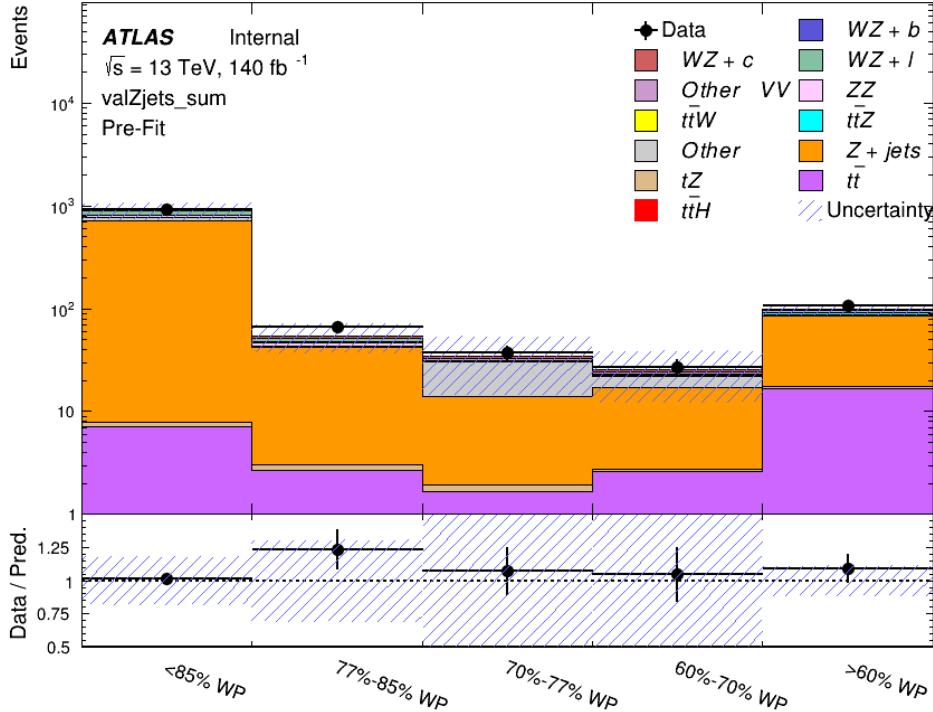


Figure 22: Data and MC comparisons for each DL1r WP for both 1-jet and 2-jet regions, after the Z+jets VR selection and correction factor have been applied

303 For each of the working points considered, the data falls within 20% of the MC prediction once  
 304 this correction factor has been applied. Therefore, a 20% systematic uncertainty is applied to Z  
 305 + jets in the analysis.

## 306 6 tZ Interference Studies and Separation Multivariate Analysis

307 Because tZ produces a final state identical to signal, it represents a predominant background in  
 308 the most signal enriched regions. That is, the region with one jet passing the 60% DL1r WP.  
 309 Therefore, a boosted decision tree (BDT) algorithm is trained using TMVA [10] to separate WZ  
 310 + heavy flavor from tZ.

311 Separation between tZ and WZ + heavy flavor is achieved in part by reconstructing the invariant  
 312 mass of the top candidate, which clusters more closely to the top mass for tZ than WZ + heavy  
 313 flavor.

314 The result of this BDT is used to create a tZ enriched region in the fit, reducing its impact on the  
 315 measurement of WZ + heavy flavor.

316 **6.1 Top Mass Reconstruction**

317 The reconstruction of the top mass follows the procedure described in detail in section 6.1 of  
 318 [11]. The mass of the top quark candidate is reconstructed from the jet, the lepton not included  
 319 in the Z-candidate, and a reconstructed neutrino. In the case that there is one jet in the event,  
 320 there is only possible b-jet candidate. For events with two jets, the jet with the highest DL1r  
 321 score is used.

322 The neutrino from the W decay is expected to be the only source of  $E_T^{\text{miss}}$ . Therefore, the  $E_T$   
 323 and  $\phi$  of the neutrino are taken from the  $E_T^{\text{miss}}$  measurement. This leaves the z-component of  
 324 the neutrino momentum,  $p_{\nu z}$  as the only unknown.

325 This unknown is solved for by taking the combined invariant mass of the lepton and neutrino to  
 326 give the invariant mass of the W boson:

$$327 \quad (p_l + p_\nu)^2 = m_W^2$$

328 Expanding this out into components, this equation gives:

$$329 \quad \sqrt{p_{T\nu}^2 + p_{z\nu}^2} E_l = \frac{m_w^2 - m_l^2}{2} + p_{T\nu}(p_{lx}\cos\phi_\nu + p_{ly}\sin\phi_\nu) + p_{lz}p_{\nu z}$$

330 This equation gives two solutions for  $p_{\nu z}$ . For cases where only one of these solutions is real,  
 331 that is taken as the value of  $p_{\nu z}$ . For instances with two real solutions, the one which is shown  
 332 to be correct in the largest fraction of simulations is taken. For cases when no real solution is  
 333 found, often because of detector effects, the value of  $E_T^{\text{miss}}$  is varied in decreasing increments of  
 334 100 MeV until a real solution is found.

335 The reconstructed top mass distribution for tZ and WZ + b can be seen in figure 23.

336 **6.2 tZ BDT**

337 A Boosted Decision Tree (BDT), specifically XGBoost [xgboost\_cite], is used to provide separa-  
 338 tion between tZ and WZ+b. The following kinematic variables are used as inputs:

- 339 • The invariant mass of the reconstructed top candidate
- 340 •  $p_T$  of each of the leptons, jet
- 341 • The invariant mass of each combination of lepton pairs,  $M(l\bar{l})$
- 342 •  $E_T^{\text{miss}}$

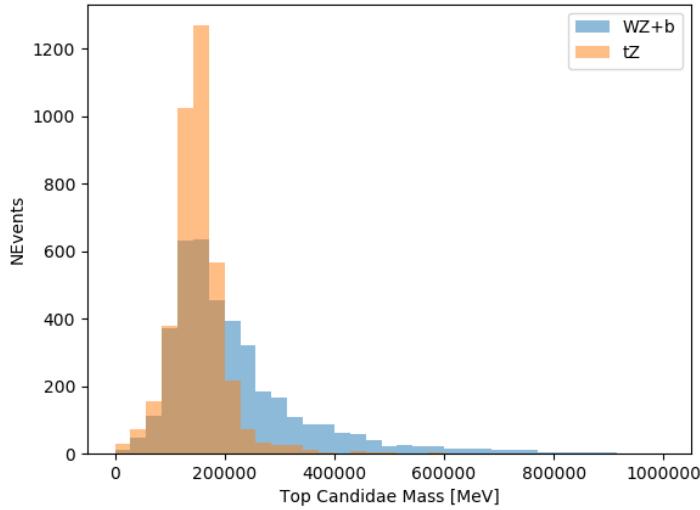


Figure 23: Reconstructed top mass distributions for tZ and WZ + b, measured in MeV.

- 343     • Distance between each combination of leptons,  $\Delta R(ll)$
- 344     • Distance between each lepton and the jet,  $\Delta R(lj)$

345     The training samples included only events meeting the requirements of the 1-jet, >60% region,  
 346     i.e. passing all the selection described in section 5 and having exactly one jet which passes the  
 347     tightest (60%) DL1r working point.

348     The distributions of a few of these features for both signal and background is shown in figure  
 349     24.

350     A sample of 20,000 background (tZ) and signal (WZ+b) Monte Carlo events are used to train  
 351     the BDT. And additional 5,000 events are reserved for testing the model, in order to prevent  
 352     over-fitting. A total of 750 decision trees with a maximum depth of 6 branches are used to build  
 353     the model. These parameters are chosen empirically, by training several models with different  
 354     parameters and selecting the one that gave the best separation for the test sample.

355     The results of the BDT training are shown in figure 25. The output scores for both signal and  
 356     background events is shown on the left. The right shows the receiving operating characteristic  
 357     (ROC) curve that results from the MVA. The ROC curve represents the background rejection  
 358     as a function of signal efficiency, where each point on the curve represents a different response  
 359     score. The ROC curve of the BDT is compared to the performance of using an optimal set of flat  
 360     selections on the same set of input variables.

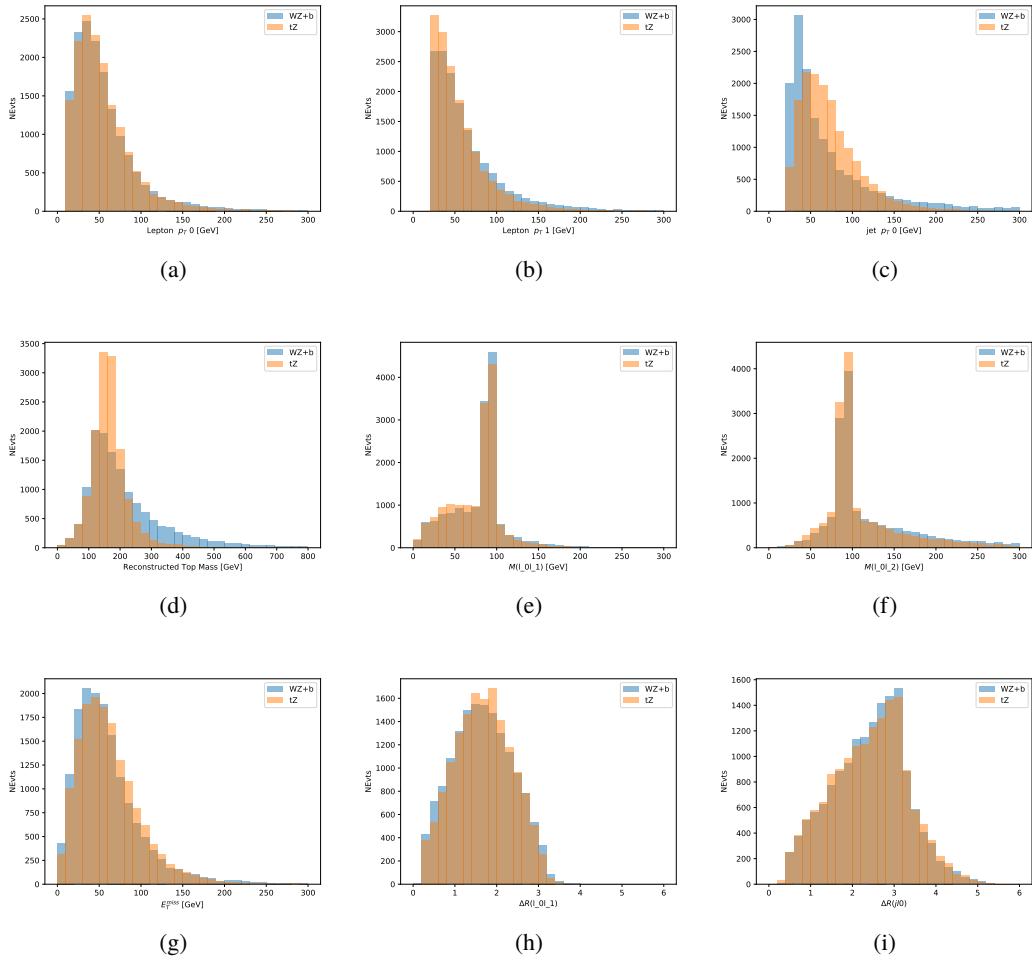
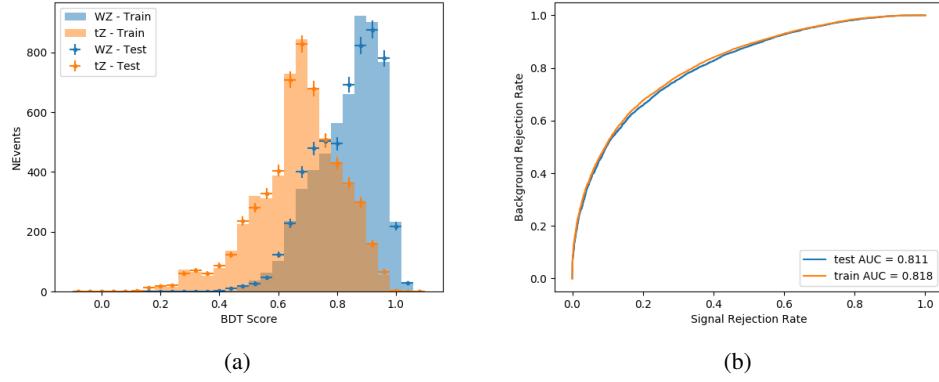


Figure 24: Distribution of input features of the BDT for signal (WZ) and background (tZ). Both are scaled to an equal number of events. (a), (b) and (c) show the  $p_T$  of lepton 0, lepton 1, and the jet, (d) show the reconstructed top mass, (e) and (f) show the invariant mass of leptons 0 and 1, leptons 0 and 2. (g) shows the  $E_T^{\text{miss}}$  of each event. (h) and (i) show the  $\Delta R$  between lepton 0 and lepton 1, and the jet.



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Figure 25: Distribution of the BDT response for signal and background events on the left, the ROC curve for the BDT on the right.  
40

361 The relative important of each input feature in the model, measured by how often they appeared  
 362 in the decision trees, is shown in figure 26.

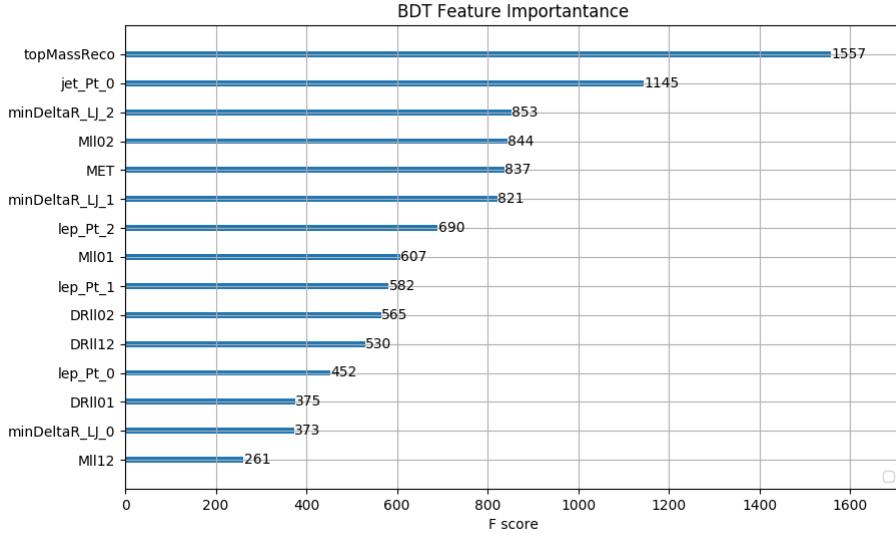


Figure 26: Relative importance of each input feature in the model.

363 These results suggest that some amount of separation can be achieved between these two pro-  
 364 cesses, with a high BDT score selecting a set of events that is pure in  $WZ + b$ .

## 365 7 Systematic Uncertainties

366 The systematic uncertainties that are considered are summarized in table 9. These are imple-  
 367 mented in the fit either as a normalization factors or as a shape variation or both in the signal  
 368 and background estimations. The numerical impact of each of these uncertainties is outlined in  
 369 section 8.

370 The uncertainty in the combined integrated luminosity is derived from a calibration of the  
 371 luminosity scale performed for 13 TeV proton-proton collisions [12], [LUCID2].

372 The experimental uncertainties are related to the reconstruction and identification of light leptons  
 373 and  $b$ -tagging of jets, and to the reconstruction of  $E_T^{\text{miss}}$ . The sources which contribute to the  
 374 uncertainty in the jet energy scale (JES) [13] are decomposed into uncorrelated components and  
 375 treated as independent sources in the analysis. The CategoryReduction model is used to account  
 376 for JES uncertainties, which decomposes

Table 9: Sources of systematic uncertainty considered in the analysis.

Systematic uncertainty	Components
Luminosity	1
Pileup reweighting	1
<b>Physics Objects</b>	
Electron	6
Muon	15
Jet energy scale and resolution	28
Jet vertex fraction	1
Jet flavor tagging	131
$E_T^{\text{miss}}$	3
Total (Experimental)	186
<b>Background Modeling</b>	
Cross section	24
Renormalization and factorization scales	10
Parton shower and hadronization model	2
Shower tune	4
Total (Signal and background modeling)	40
Total (Overall)	226

377 The uncertainties in the b-tagging efficiencies measured in dedicated calibration analyses [14] are  
 378 also decomposed into uncorrelated components. The large number of components for b-tagging  
 379 is due to the calibration of the distribution of the MVA discriminant.

380 The systematic uncertainties associated with the signal and background processes are accounted  
 381 for by varying the cross-section of each process within its uncertainty.

382 The full list of systematic uncertainties considered in the analysis is summarized in tables 10, 11  
 383 and 12.

384

Experimental Systematics on Leptons and $E_T^{\text{miss}}$			
Type	Description	Systematics Name	Application
<b>Trigger</b>			
Scale Factors	Trigger Efficiency	lepSFTrigTight_MU(EL)_SF_Trigger_STAT(SYST)	Event Weight
<b>Muons</b>			
Efficiencies	Reconstruction and Identification	lepSFObjTight_MU_SF_ID_STAT(SYST)	Event Weight
	Isolation	lepSFObjTight_MU_SF_Isol_STAT(SYST)	Event Weight
	Track To Vertex Association	lepSFObjTight_MU_SF_TTVA_STAT(SYST )	Event Weight
$p_T$ Scale	$p_T$ Scale	MUONS_SCALE	$p_T$ Correction
Resolution	Inner Detector Energy Resolution	MUONS_ID	$p_T$ Correction
	Muon Spectrometer Energy Resolution	MUONS_MS	$p_T$ Correction
<b>Electrons</b>			
Efficiencies	Reconstruction	lepSFObjTight_EL_SF_ID	Event Weight
	Identification	lepSFObjTight_EL_SF_Reco	Event Weight
	Isolation	lepSFObjTight_EL_SF_Isol	Event Weight
Scale Factor	Energy Scale	EG_SCALE_ALL	Energy Correction
Resolution	Energy Resolution	EG_RESOLUTION_ALL	Energy Correction
<b><math>E_T^{\text{miss}}</math></b>			
Soft Tracks Terms	Resolution	MET_SoftTrk_ResoPerp	$p_T$ Correction
	Resolution	MET_SoftTrk_ResoPara	$p_T$ Correction
	Scale	MET_SoftTrk_ScaleUp	$p_T$ Correction
	Scale	MET_SoftTrk_ScaleDown	$p_T$ Correction

Table 10: Summary of experimental systematics considered for leptons and  $E_T^{\text{miss}}$ . Includes type, description, name of systematic as used in the fit, and mode of application. The mode of application indicates the systematic evaluation, e.g. as an overall event re-weighting (Event Weight) or rescaling ( $p_T$  Correction).

Experimental Systematics on b-tagging		
Type	Origin	Systematic Name
Scale Factors	DL1r b-tagger efficiency on b originated jets in bins of $\eta$	DL1r_Continuous_EventWeight_B0-29
	DL1r b-tagger efficiency on c originated jets in bins of $\eta$	DL1r_Continuous_EventWeight_C0-19
	DL1r b-tagger efficiency on light flavoured originated jets in bins of $\eta$ and $p_T$	DL1r_Continuous_EventWeight_Light0-79
9th February 2014 2016ger extrapolation efficiency		DL1r_Continuous_EventWeight_extrap040on DL1r_Continuous_EventWeight_extrapolation_from_charm

Table 12: Summary of experimental systematics to be included for b-tagging of jets in the analysis, using the continuous DL1r tagging algorithm. All of the b-tagging related systematics are applied as event weights. From left: type, description, and the name of systematic used in the fit.

Experimental Systematics on Jets			
Type	Origin	Systematics Name	Application
Jet Vertex Tagger		JVT	Event Weight
Energy Scale	Calibration Method	JET_21NP_ JET_EffectiveNP_1-19	$p_T$ Correction $p_T$ Correction
	$\eta$ inter-calibration	JET_EtaIntercalibration_Modelling JET_EtaIntercalibration_NonClosure JET_EtaIntercalibration_TotalStat	$p_T$ Correction $p_T$ Correction $p_T$ Correction
	High $p_T$ jets	JET_SingleParticle_HighPt	$p_T$ Correction
	Pile-Up	JET_Pileup_OffsetNPV JET_Pileup_OffsetMu JET_Pileup_PtTerm JET_Pileup_RhoTopology	$p_T$ Correction $p_T$ Correction $p_T$ Correction $p_T$ Correction
	Non Closure	JET_PunchThrough_MC15	$p_T$ Correction
	Flavour	JET_Flavor_Response JET_BJES_Response JET_Flavor_Composition	$p_T$ Correction $p_T$ Correction $p_T$ Correction
Resolution		JET_JER_SINGLE_NP	Event Weight

Table 11: Jet systematics take into account effects of jets calibration method,  $\eta$  inter-calibration, high  $p_T$  jets, pile-up, and flavor response. They are all diagonalised into effective parameters.

385 Theoretical uncertainties applied to backgrounds, including cross section, PDF, and scale un-  
 386 certainties are taken from theory calculations, with the exception of non-prompt and diboson  
 387 backgrounds. Derivation of the non-prompt background uncertainties, Z+jets and  $t\bar{t}$ , are ex-  
 388 plained in detail in section 5.3. Because the other VV + heavy flavor processes are also poorly  
 389 understood, a conservative 50% uncertainty is applied to those samples. The cross-section  
 390 uncertainties applied to the predominate background estimates are summarized in table 13.

## 391 8 Results

392 A separate maximum-likelihood fit is performed over the 1-jet and 2-jet fit regions in order to  
 393 extract the best-fit value of the WZ + b-jet and WZ + charm jet contributions. The WZ + b,  
 394 WZ + charm and WZ + light contributions are allowed to float, with the remaining background  
 395 contributions are held fixed. **The current fit strategy treats the WZ + b-jet contribution as**  
 396 **the parameter of interest, with the normalization of the WZ + charm and the WZ + light**

Process	X-section [%]
WZ	QCD Scale: $^{+3.7}_{-3.4}$ PDF( $+\alpha_S$ ): $\pm 3.1$
tZ	X-sec: $\pm 15.2$ QCD Scale: $^{+5.2}_{-1.3}$ PDF( $+\alpha_S$ ): $\pm 1.2$
t̄H (aMC@NLO+Pythia8)	QCD Scale: $^{+5.8}_{-9.2}$ PDF( $+\alpha_S$ ): $\pm 3.6$
t̄Z (aMC@NLO+Pythia8)	QCD Scale: $^{+9.6}_{-11.3}$ PDF( $+\alpha_S$ ): $\pm 4$
t̄W (aMC@NLO+Pythia8)	QCD Scale: $^{+12.9}_{-11.5}$ PDF( $+\alpha_S$ ): $\pm 3.4$
VV + b/charm (Sherpa 2.2.1)	$\pm 50$
VV + light (Sherpa 2.2.1)	$\pm 6$
t̄t	$\pm 20$
Z + jets	$\pm 25$

Table 13: Summary of theoretical uncertainties for MC prediction of backgrounds in the analysis.

397 **contributions taken as systematic uncertainties. This could however be adjusted, depending**  
 398 **on whether it is decided the goal of the analysis should be to measure WZ+b specifically or**  
 399 **WZ + heavy flavor overall.** The result of the fit is used to extract the cross-section of WZ +  
 400 heavy-flavor production.

401 A maximum likelihood fit to data is performed simultaneously in the regions described in section  
 402 5. The parameters  $\mu_{WZ+b}$ ,  $\mu_{WZ+charm}$ ,  $\mu_{WZ+light}$ , where  $\mu = \sigma_{\text{observed}}/\sigma_{\text{SM}}$ , are extracted  
 403 from the fit.

## 404 8.1 1-jet Fit Results

405 **The results of the fit are currently blinded.** The post-fit yields in each region are summarized  
 406 in figure 27.

407 A post-fit summary plot of the 1-jet fitted regions is shown in figure 28:

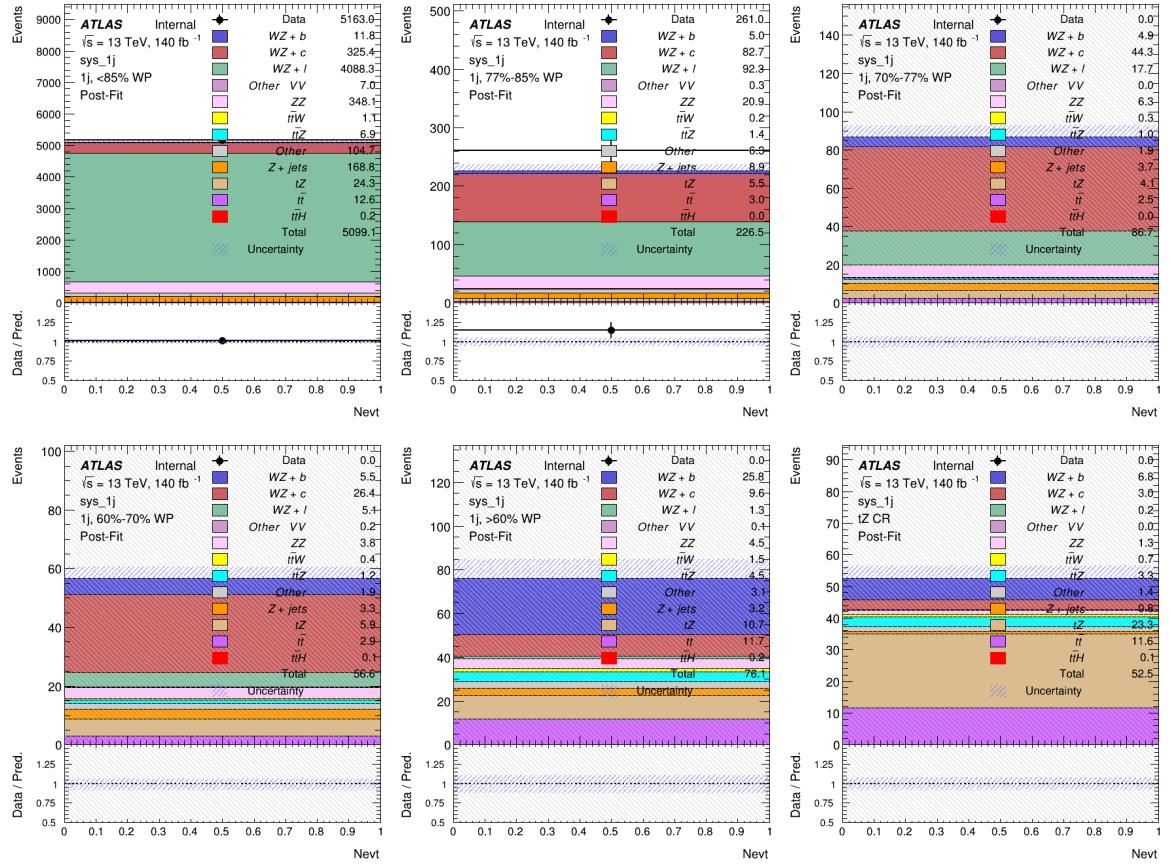


Figure 27: Data/MC results in each of the 1-jet regions after the fit has been performed.

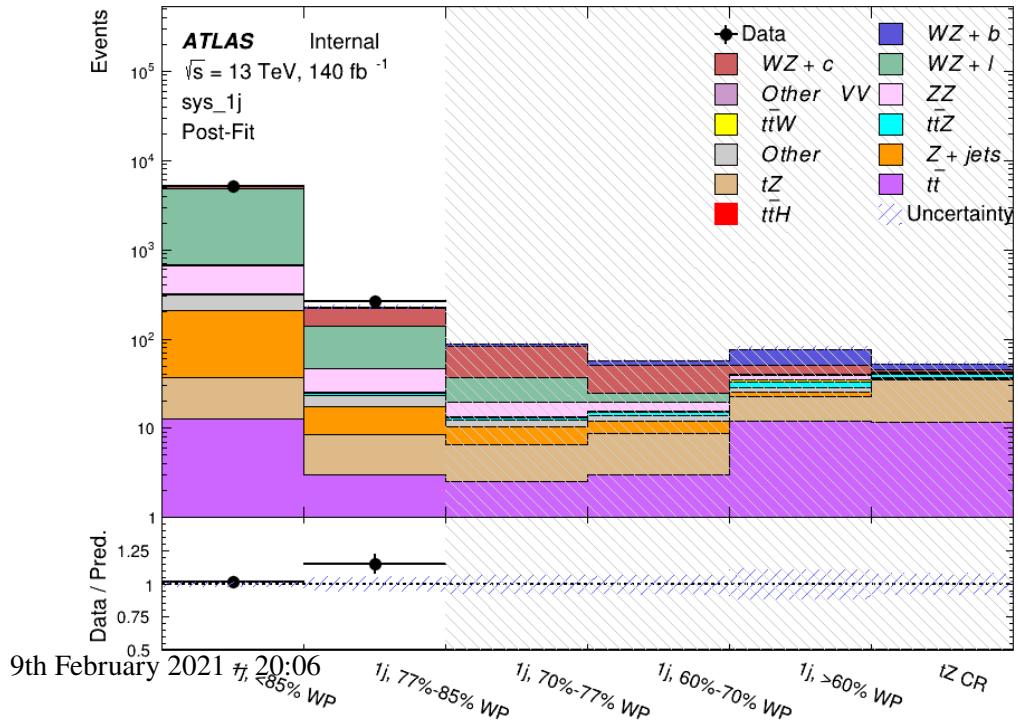


Figure 28: Post-fit summary of the 1-jet fit regions.

408 As described in section 7, there are 226 systematic uncertainties that are considered as NPs in  
 409 the fit. These NPs are constrained by Gaussian or log-normal probability density functions. The  
 410 latter are used for normalisation factors to ensure that they are always positive. The expected  
 411 numbers of signal and background events are functions of the likelihood. The prior for each  
 412 NP is added as a penalty term, decreasing the likelihood as it is shifted away from its nominal  
 413 value.

414 The impact of each NP is calculated by performing the fit with the parameter of interest held  
 415 fixed, varied from its fitted value by its uncertainty, and calculating  $\Delta\mu$  relative to the baseline  
 416 fit. The impact of the most significant sources of systematic uncertainties is summarized in table  
 417 14.

Uncertainty Source	$\Delta\mu$	
WZ + charm cross-section	-0.1966	0.2171
tZ cross-section	-0.1521	0.1518
WZ + light cross-section	0.1485	-0.1411
Other VV + b cross-section	-0.1115	0.1163
Flavor Tagging	0.0955	0.0957
Jet Energy Scale	0.0613	0.081
t <bar>t</bar>	-0.0662	0.0654
Luminosity	-0.0609	0.0655
Z + jets cross-section	-0.0284	0.0284
Other VV + charm cross-section	0.0207	-0.0202
Muon Trigger Scale Factor	0.019	0.0209
Total Systematic Uncertainty	0.3511	0.3679

Table 14: Summary of the most significant sources of systematic uncertainty on the measurement of WZ + b with exactly one associated jet.

418 The ranking and impact of those nuisance parameters with the largest contribution to the overall  
 419 uncertainty is shown in figure 29.

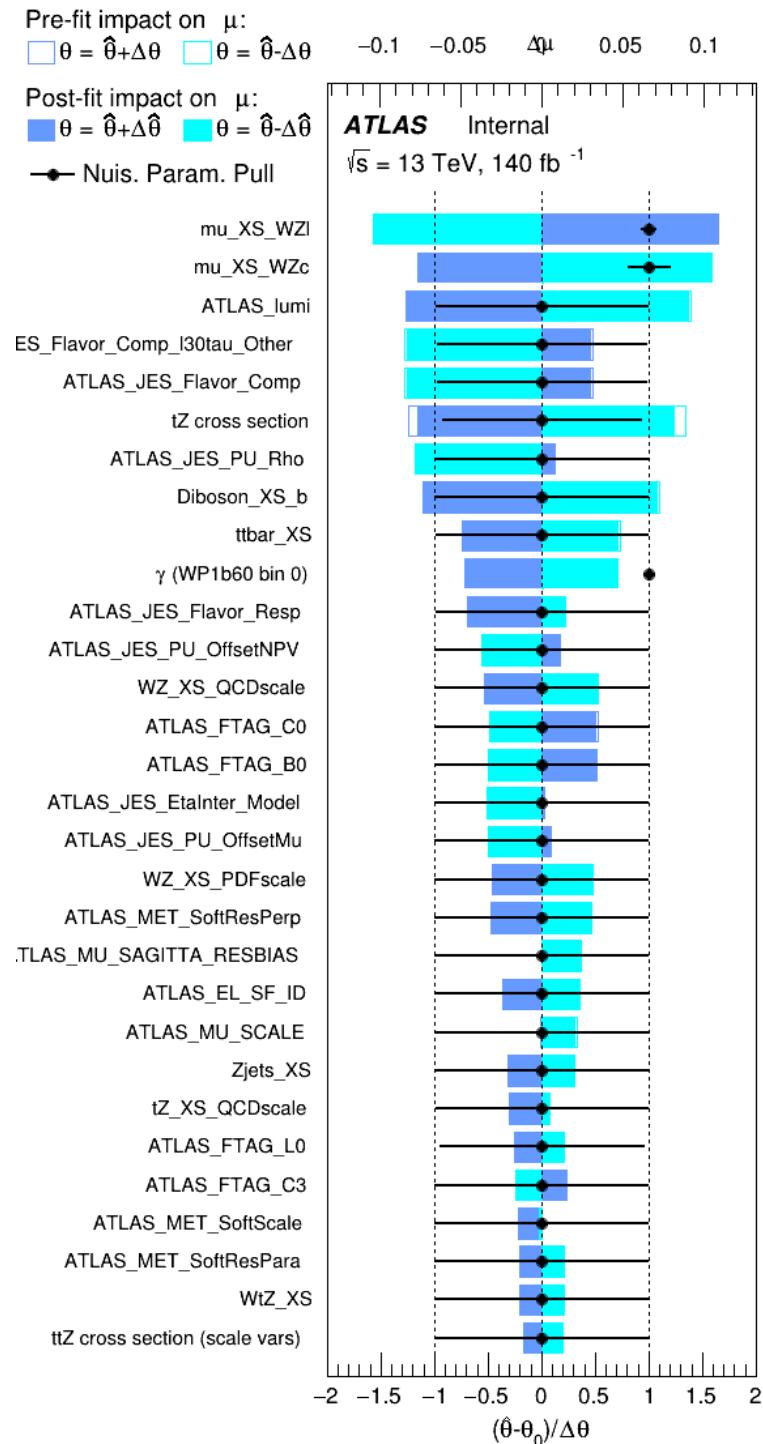


Figure 29: Impact of systematic uncertainties on the signal-strength of  $WZ + b$  for events with exactly one jet

420 The large impact of the Jet Energy Scale and Jet Flavor Tagging is unsurprising, as the shape  
 421 of the fit regions depends heavily on the modeling of the jets. The other major sources of  
 422 uncertainty come from background modelling and cross-section uncertainty. The pie charts in  
 423 figure 30 show that for the modelling uncertainties that contribute most correspond to the most  
 424 significant backgrounds.

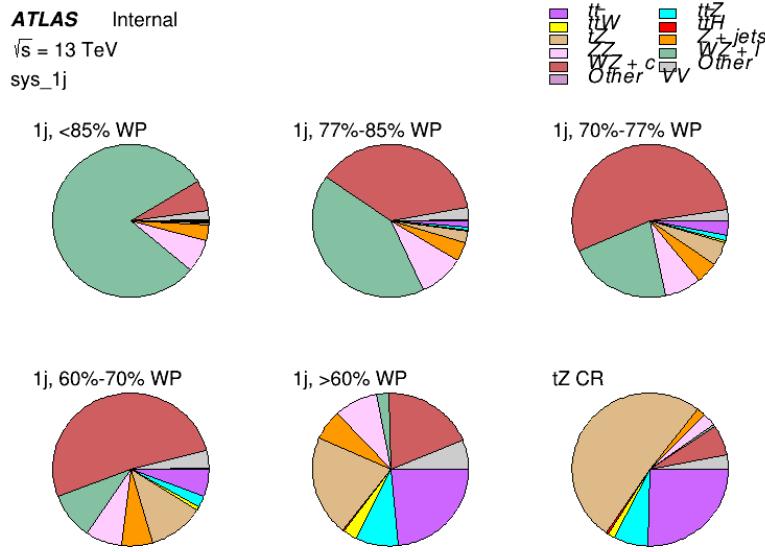


Figure 30: Post-fit background composition of the fit regions.

425 The correlations between these nuisance parameters are summarized in figure 31.

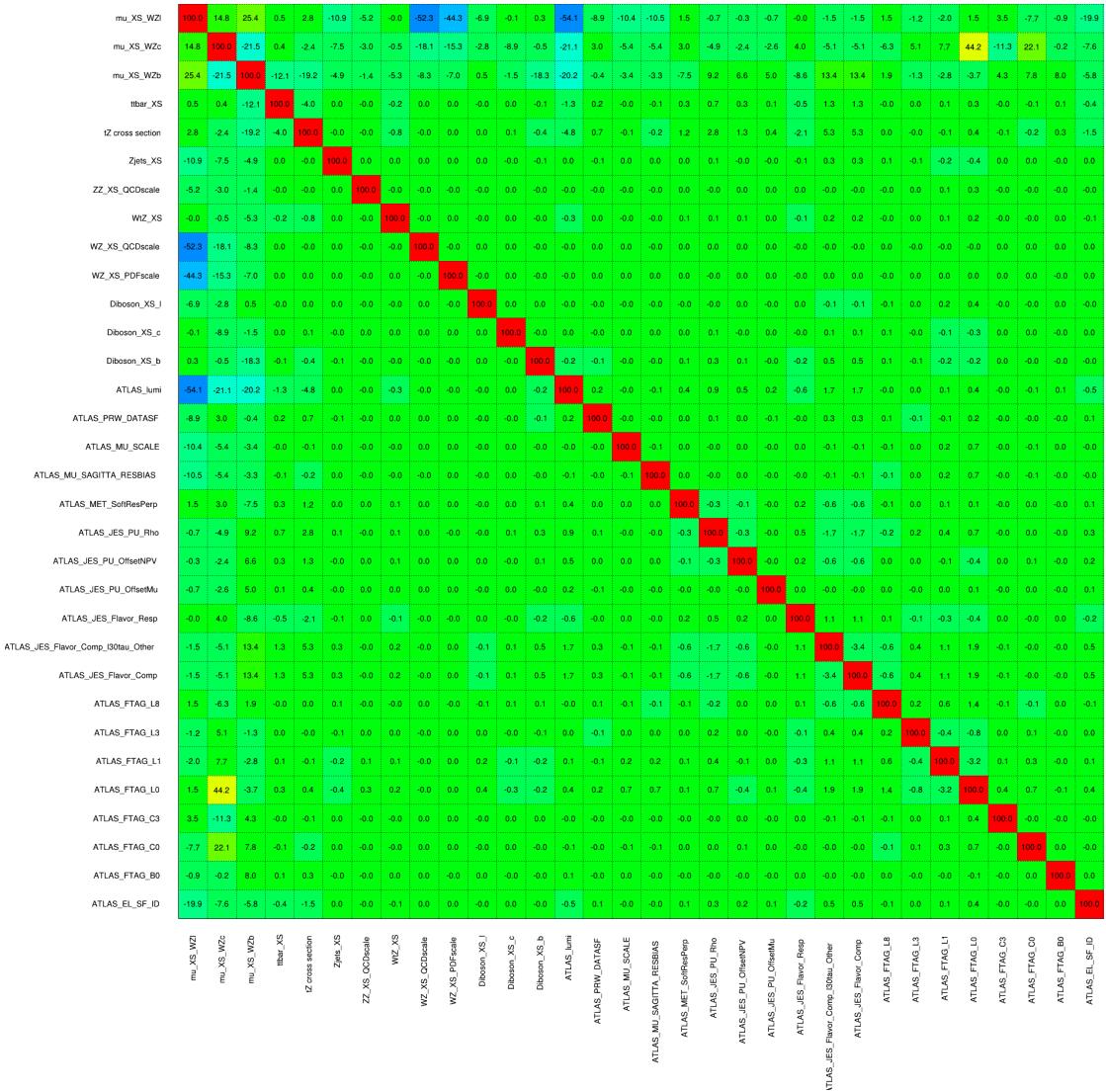


Figure 31: Correlations between nuisance parameters

426 The negative correlations between  $\mu_{WZ+charm}$  and  $\mu_{WZ+b}$  and  $\mu_{WZ+light}$  are expected: WZ +  
 427 charm is present in both the WZ + b and WZ + light enriched regions, therefore increasing the  
 428 fraction of charm requires increasing the fraction of WZ + b and WZ + light. This reasoning  
 429 also explains the positive correlation between  $\mu_{WZ+b}$  and  $\mu_{WZ+light}$ .

430 Two of the major backgrounds in the region with the highest purity of WZ + b are tZ and Other  
 431 VV + b, explaining the negative correlations between  $\mu_{WZ+b}$  and the tZ cross section, and the  
 432 VV + b cross section.

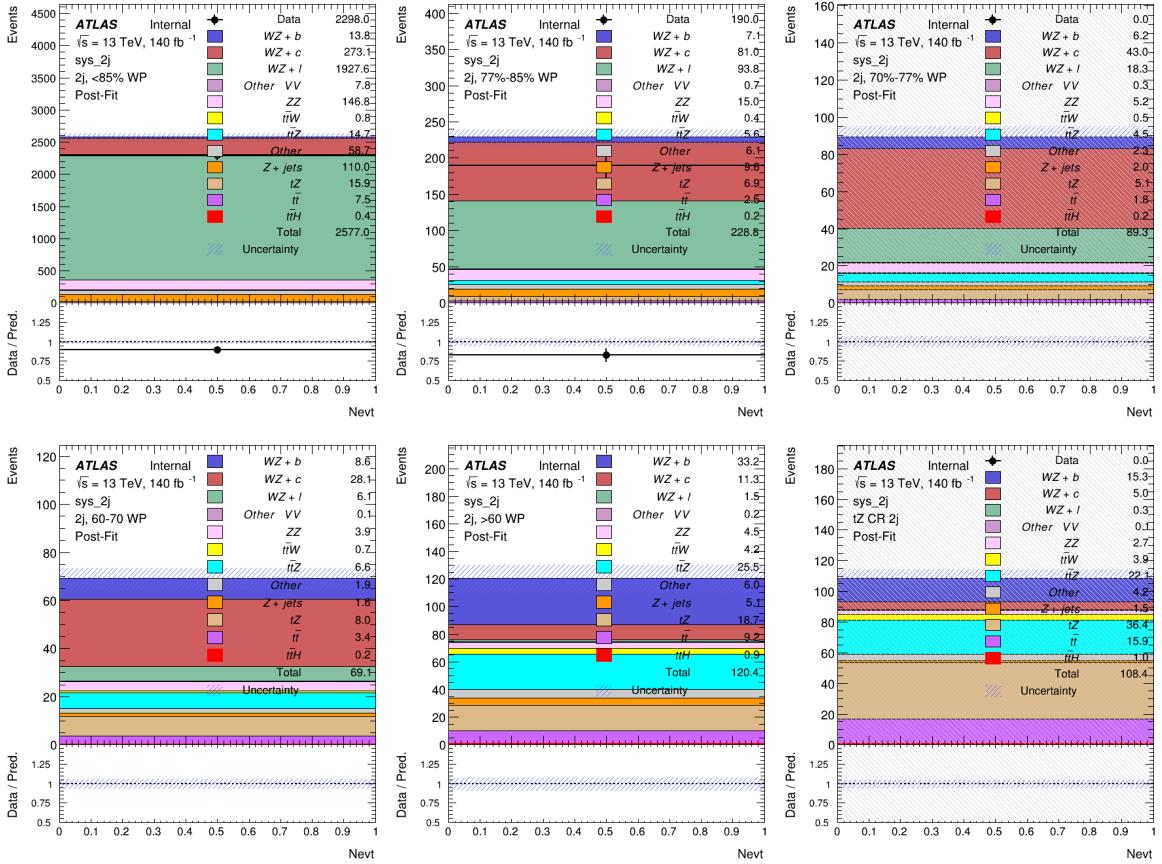


Figure 32: Data/MC results in each of the regions in the 2-jet fit after the fit has been performed.

433 The high correlation between the luminosity and  $\mu_{WZ+\text{light}}$  arises from the fact that the uncer-  
 434 tainty on  $\mu_{WZ+\text{light}}$  is very low (around 4%). Small changes in luminosity cause a change in  
 435 the yield of  $WZ + \text{light}$  that is large compared to its uncertainty, producing a large correlation  
 436 between these two parameters.

## 437 8.2 2-jet Fit Results

438 **The results of the fit are currently blinded.** The post-fit yields in each region are summarized  
 439 in figure 32.

440 A post-fit summary plot of the fitted regions is shown in figure 33:

441 The same set of systematic uncertainties consider for the 1-jet fit are included in the 2-jet fit as  
 442 well. The impact of the most significant systematic uncertainties is summarized in table 15.

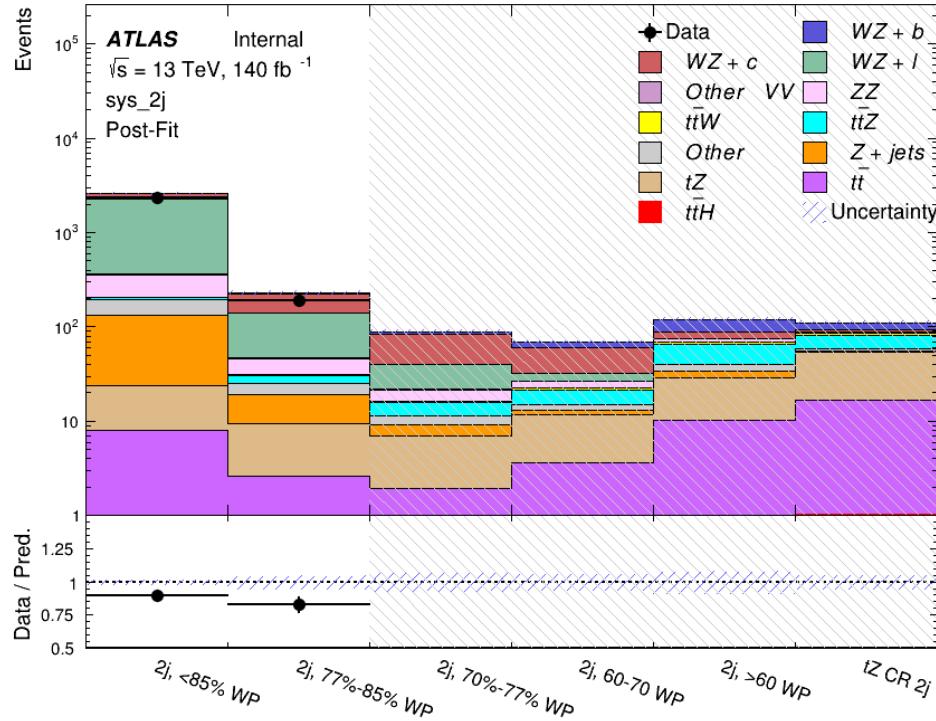


Figure 33: Post-fit summary of the fit over 2-jet regions.

Uncertainty Source	$\Delta\mu$	
WZ + charm cross-section	-0.1966	0.2171
tZ cross-section	-0.1521	0.1518
WZ + light cross-section	0.1485	-0.1411
Other VV + b cross-section	-0.1115	0.1163
Flavor Tagging	0.0955	0.0957
Jet Energy Scale	0.0613	0.081
t̄t cross-section	-0.0662	0.0654
Luminosity	-0.0609	0.0655
Z + jets cross-section	-0.0284	0.0284
Other VV + charm cross-section	0.0207	-0.0202
Muon Trigger Scale Factor	0.019	0.0209
Total Systematic Uncertainty	0.3511	0.3679

Table 15: Summary of the most significant sources of systematic uncertainty on the measurement of WZ + b 2-jet events.

<sup>443</sup> The ranking and impact of those nuisance parameters with the largest contribution to the overall  
<sup>444</sup> uncertainty is shown in figure 34.

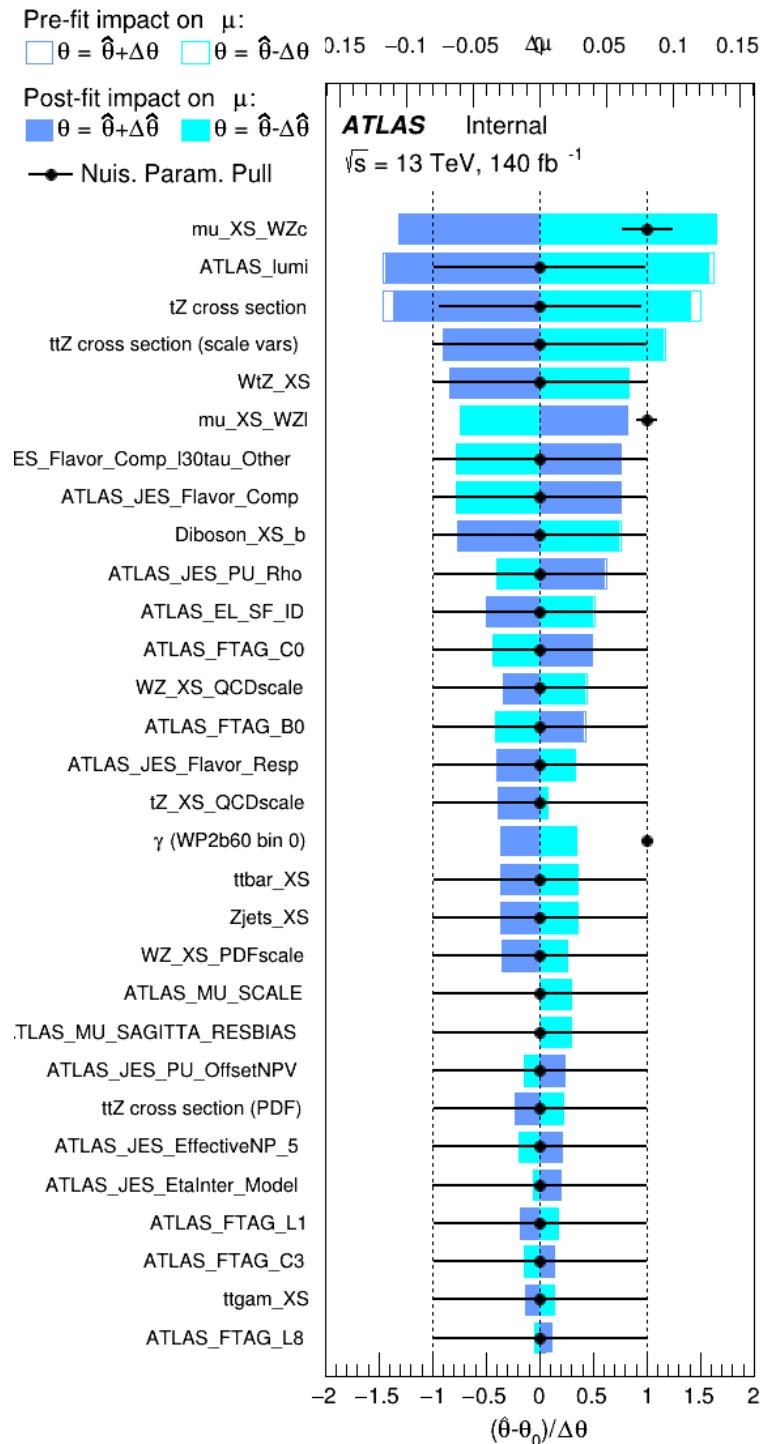


Figure 34: Impact of systematic uncertainties on the signal-strength of  $WZ + b$  in 2-jet events.

445 The large impact of the Jet Energy Scale and Jet Flavor Tagging is unsurprising, as the shape  
446 of the fit regions depends heavily on the modeling of the jets. The other major sources of  
447 uncertainty come from background modelling and cross-section uncertainty. The pie charts in  
448 figure 35 show that for the modelling uncertainties that contribute most correspond to the most  
449 significant backgrounds.

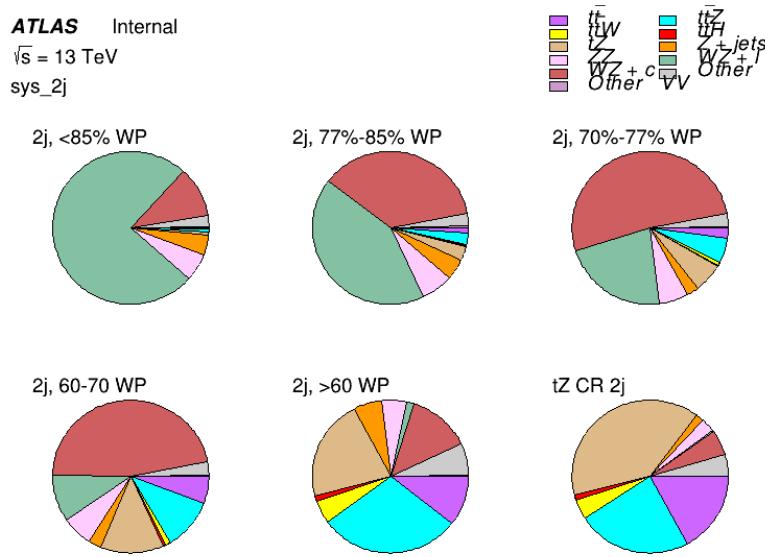


Figure 35: Post-fit background composition of the 2-jet fit regions.

450 The correlations between these nuisance parameters are summarized in figure 36.

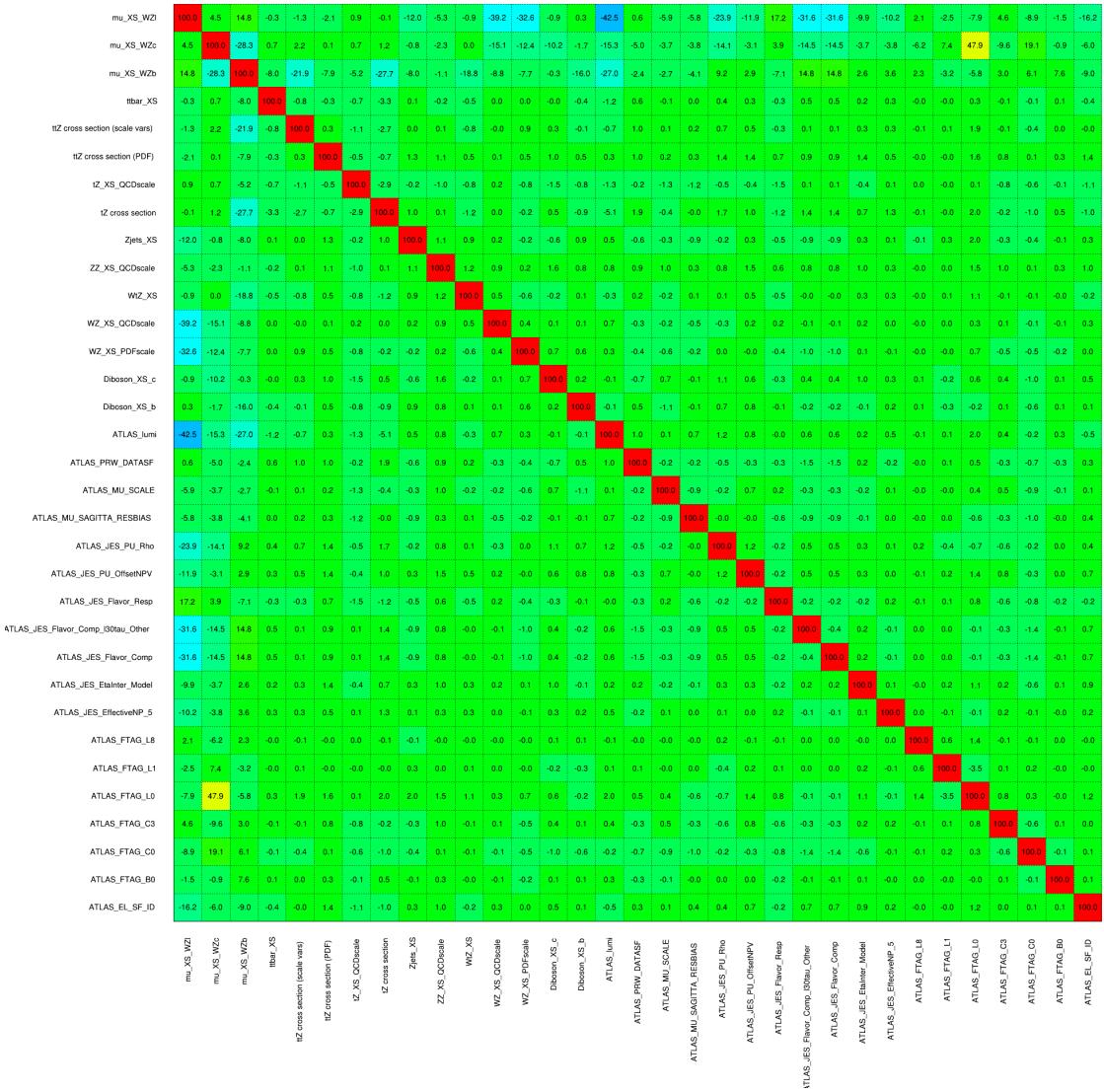


Figure 36: Correlations between nuisance parameters in the 2-jet fit

451 As in the 1-jet case, no significant, unexpected correlations are found between nuisance para-  
452 meters.

## 453 9 Conclusion

454 A measurement of  $WZ + \text{heavy flavor}$  is performed using  $140 \text{ fb}^{-1}$  of  $\sqrt{s} = 13 \text{ TeV}$  proton-  
455 proton collision data collected by the ATLAS detector at the LHC. This section will be include

456 **final results once unblinded.**

457 **Appendices**

458 **.1 tZ Interference Studies**

459 Because it includes an on-shell Z boson as well as a b-jet and W from the top decay, tZ production  
460 represents an identical final state to WZ + b-jet. This implies the possibility of matrix level  
461 interference between these two processes not accounted for in the Monte Carlo simulations,  
462 which consider the two processes independently. Truth level studies are performed in order to  
463 estimate the impact of these interference effects.

464 In order to estimate the matrix level interference effects between tZ and WZ + b-jet, two different  
465 sets of simulations are produced using MadGraph 5 [**Madgraph**] - one which simulates these  
466 two processes independently, and another where they are produced simultaneously, such that  
467 interference effects are present. These two sets of samples are then compared, and the difference  
468 between them can be taken to represent any interference effects. The flavor flavor scheme is used  
469 in both cases.

470 MadGraph simulations of 10,000 tZ and 10,000 WZ + b-jet events (excluding) are produced,  
471 along with 20,000 events where both are present, in the fiducial region where three leptons and  
472 at least one jet are produced. Both are generated in MadGraph

473 A selection mimicking the preselection used in the main analysis is applied to the samples: The  
474 SS leptons are required to have  $p_T > 20$  GeV, and  $> 10$  GeV is required for the OS lepton. The  
475 associated b-jet is required to have  $p_T > 25$  GeV, and all physics objects are required to fall in a  
476 range of  $|\eta| < 2.5$ .

477 The kinematics of these samples after the selection has been applied are shown below:

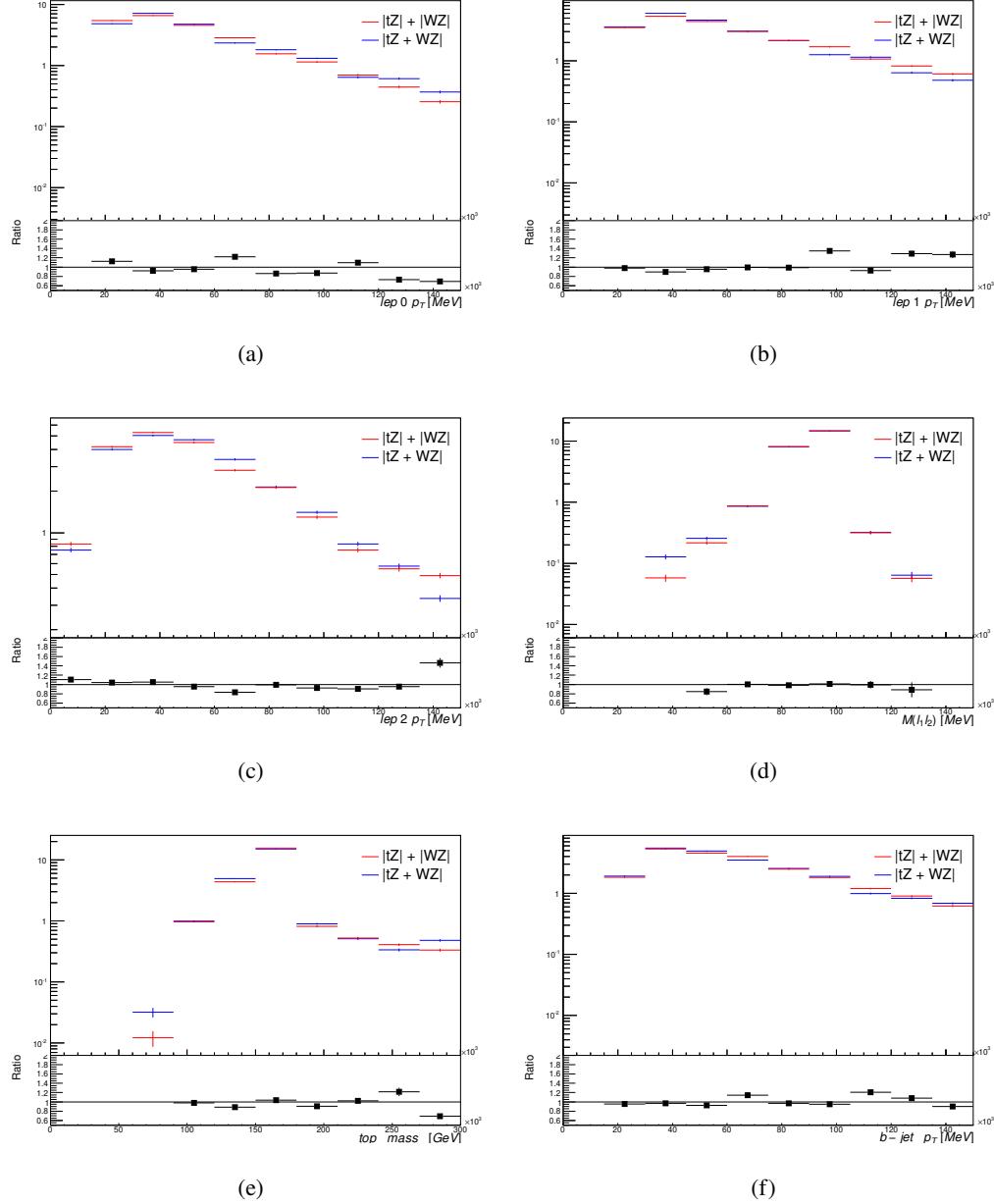


Figure 37: Comparisons between (a) the  $p_T$  of the lepton from the W, (b) and (c) show the  $p_T$  of the lepton from the Z, (d) the invariant mass of the Z-candidate, (e) the mass of the top candidate, and (f) the  $p_T$  of the b-jet, for WZ and tZ events generated with interference effects (blue) and without interference effects (red).

478 The overall cross-section of the two methods agree within error, and no significant differences  
479 in the kinematic distributions are seen. It is therefore concluded that interference effects do not

480 significantly impact the results.

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