



ATLAS Note

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2 WZ + Heavy Flavor Production in pp collisions 3 at $\sqrt{s} = 13$ TeV

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The ATLAS Collaboration

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A measurement of WZ produced with an associated heavy flavor jet is performed using 140
6 fb^{-1} of proton-proton collision data at $\sqrt{s} = 13$ TeV from the ATLAS experiment at the
7 LHC. The measurement is performed in the fully leptonic decay mode, $\text{WZ} \rightarrow \ell\bar{\nu}\ell\bar{\nu}$. The
8 cross-section of $\text{WZ} + b\text{-jets}$ is measured to be $X \pm X \pm X$, while the cross-section of $\text{WZ} +$
9 charm is measured as X , with a correlation of X between the two processes.

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¹¹ **Contents**

12 1 Changes and outstanding items**13 1.1 Changelog**

14 This is version 6

15 1.1.1 Changes relative to v5

- 16 • added list of DSIDs to an appendix

17 1.1.2 Changes relative to v4

- 18 • Fixed various typos, clarified wording
19 • Expanded info about JER uncertainties, electron systematics, theory uncertainties
20 • removed a table on lepton selection, included information in the text instead
21 • Plotted lepton Pt Z and W for Zjet CR, rather than lep 1 and 2
22 • fixed binning in kinematic plots
23 • Included prefit and postfit yield tables
24 • added signal modelling systematics
25 • included alternate fit studies with tZ included in signal

26 1.1.3 Changes relative to v3

- 27 • Merged introduction into executive summary, including unblinding details and list of
28 SRs/CRs used
29 • listed ptag used (p4133), and release (AB 21.2.127)
30 • Included table reftab:xsecUnc listing x-sec uncertainties used
31 • Removed selection criteria listed in table ?? (QMisID, AmbiguityType) that were removed
32 from the analysis
33 • specified variable used to make truth jet flavor determination (HadronConeExclTruthLa-
34 belID)
35 • fixed bug in MtLepMet calculation, updated selection/fits to account for this
36 • Included plots of MtLepMet and PtZ, swapped lep 1 and 2 p_T plots for lep W and lep Z
37 plots

- 38 • updated tZ BDT training to reduce overfitting, updated plots to include error bars, feature
39 importance
- 40 • updated table ?? to clarify selection, fix the tZ_BDT cut used
- 41 • replace a few broken ntuples which included large weight events
- 42 • include DL1r distribution for Z+jets and $t\bar{t}$ VRs
- 43 • Expanded section on fakes, included information on derived scale factors from VRs.
- 44 • Changed the kinematic plots to include $p_T(Z)$ and $m_T(W)$, list lepton p_T based on W and
45 Z candidates.

46 **1.1.4 Changes relative to v2**

- 47 • Added alternate VBS samples to include missing b-jet diagrams
- 48 • Included a section on tZ interference effects, ??.
- 49 • Updated to reflect changes for 2018, including the move to PFlow jets, DL1r, updated
50 trigger, and updated AnalysisBase version (now 21.2.127)
- 51 • Revised fit regions, using separate 1-jet and 2-jet fits, with all 2-j regions included
- 52 • updated plots for tZ BDT, added details about the model
- 53 • Included truth jet information

54 **1.1.5 Changes relative to v1**

- 55 • Added GRL list
- 56 • Fixed latex issue in line 92, typo in line 172
- 57 • Added tables ?? and ??, summarizing the event and object selection
- 58 • Added table ??, which includes the DSID of samples used
- 59 • Included reference to WZ inclusive paper in introduction

60 **1.2 Outstanding Items**

- 61 • Complete interference studies, apply any interference effects observed as a systematic
62 • Update results section with additional studies, possibly including:
63 – Truth jet migration studies
64 – Simultaneous fit over 1j and 2j
65 – Impact of allowing tZ to float
66 • Unblind, update plots and fits to include data
67 • Add cross-section, significance once unblinded

68 2 Executive Summary

69 The production of WZ in association with a heavy flavor jet represents an important background
70 for many major analyses. This includes any process with leptons and b-jets in the final state,
71 such as $t\bar{t}H$, $t\bar{t}W$, and $t\bar{t}Z$. While precise measurements have been made of WZ production
72 [1], $WZ +$ heavy flavor remains poorly understood. This is largely because the QCD processes
73 involved in the production of the b-jet make it difficult to simulate accurately. This introduces a
74 large uncertainty for analyses that include this process as a background.

75 Motivated by its relevance to the $t\bar{t}H$ multilepton analysis, we perform a study of the fully
76 leptonic decay mode of this channel; that is, events where both the W and Z decay leptonically.
77 Because WZ has no associated jets at leading order, while the major backgrounds for this channel
78 tend to have high jet multiplicity, events with more than two jets are rejected. This gives a final
79 state signature of three leptons and one or two jets.

80 Events that meet this selection criteria are sorted into pseudo-continuous b-tagging regions based
81 on the DL1r b-tag score of their associated jets. This is done to separate $WZ +$ b-jet events from
82 $WZ +$ charm and $WZ +$ light jets. These regions are fit to data in order make a more accurate
83 estimate of the contribution of $WZ +$ heavy-flavor, where heavy-flavor jets include b-jets and
84 charm jets. The full Run-2 dataset collected by the ATLAS detector, representing 139 fb^{-1} of
85 data from pp collisions at $\sqrt{s} = 13 \text{ TeV}$, is used for this study.

86 All backgrounds are accounted for using Monte Carlo, with the simulation of non-prompt lepton
87 backgrounds - $Z+jets$ and $t\bar{t}$ - validated using non-prompt Validation Regions.

88 Section ?? details the data and Monte Carlo (MC) samples used in the analysis. The reconstruction
89 of various physics objects is described in section ???. Section ?? describes the event selection
90 applied to these samples, along the definitions of the various regions used in the fit. The
91 multivariate analysis techniques used to separate the tZ background from $WZ +$ heavy flavor are
92 described in section ???. Section ?? describes the various sources of systematic uncertainties
93 considered in the fit. Finally, the results of the analysis are summarized in section ??, followed
94 by a brief conclusion in section ??.

95 The analysis aims to report a cross-section measurement of $WZ+b$ and $WZ+charm$, along with
96 their correlations, for both 1-jet and 2-jet exclusive regions. The proposed fiducial region for
97 these measurements includes events with three leptons, where the invariant mass of at least one
98 opposite charge, same flavor lepton pair falls within 10 GeV of 91.2 GeV , with 1 or 2 associated
99 jets. An alternate version of the measurement is included in the appendix, which considers tZ as
100 part of the $WZ+b$ signal.

101 The current state of thee analysis shows blinded results for thee full Run 2 dataset. Regions
102 containing $>5\%$ $WZ+b$ events are blinded, and results are from Asimov, MC only fits. Expected
103 significance and cross-section numbers are reported.

104 **3 Data and Monte Carlo Samples**

105 Both data and Monte Carlo samples used in this analysis were prepared in the `xAOD` format,
 106 which was used to produce a `DxAOD` sample in the `HIGG8D1` derivation framework. The `HIGG8D1`
 107 framework is designed for the $t\bar{t}H$ multi-lepton analysis, which targets events with multiple
 108 leptons as well as tau hadrons. This framework skims the dataset to remove unneeded variables
 109 as well as entire events. Events are removed from the derivations that do not meet the following
 110 selection:

- 111 • at least two light leptons within a range $|\eta| < 2.6$, with leading lepton $p_T > 15$ GeV and
 112 subleading lepton $p_T > 5$ GeV
- 113 • at least one light lepton with $p_T > 15$ GeV within a range $|\eta| < 2.6$, and at least two hadronic
 114 taus with $p_T > 15$ GeV.

115 Samples were then generated from these `HIGG8D1` derivations with p-tag of p4134 using Ana-
 116 lysisBase version 21.2.127 modified to include custom variables..

117 **3.1 Data Samples**

118 The study uses a sample of proton-proton collision data collected by the ATLAS detector from
 119 2015 through 2018 at an energy of $\sqrt{s} = 13$ TeV, which represents an integrated luminosity of
 120 139 fb^{-1} . This data set was collected with a bunch crossing rate of 25 ns. All data used in this
 121 analysis was verified by data quality checks, having been included in the following Good Run
 122 Lists:

- 123 • `data15_13TeV.periodAllYear_DetStatus-v79-repro20-02_DQDefects-00-02-02`
 124 `_PHYS_StandardGRL_All_Good_25ns.xml`
- 125 • `data16_13TeV.periodAllYear_DetStatus-v88-pro20-21_DQDefects-00-02-04`
 126 `_PHYS_StandardGRL_All_Good_25ns.xml`
- 127 • `data17_13TeV.periodAllYear_DetStatus-v97-pro21-13_Uncknown_PHYS_StandardGRL`
 128 `_All_Good_25ns_Triggerno17e33prim.xml`
- 129 • `data18_13TeV.periodAllYear_DetStatus-v102-pro22-04_Uncknown_PHYS_StandardGRL`
 130 `_All_Good_25ns_Triggerno17e33prim.xml`

131 Runs included from the AllYear period containers are included.

132 3.2 Monte Carlo Samples

133 Several different generators were used to produce Monte Carlo simulations of the signal and
 134 background processes. For all samples, the response of the ATLAS detector is simulated using
 135 Geant4. The WZ signal samples are simulated using Sherpa 2.2.2 [2]. Specific information about
 136 the Monte Carlo samples being used can be found in Table ???. A list of the specific samples
 137 used by data set ID is shown in Table ???.

Table 1: The configurations used for event generation of signal and background processes, including the event generator, matrix element (ME) order, parton shower algorithm, and parton distribution function (PDF).

Process	Event generator	ME order	Parton Shower	PDF
WZ, VV	SHERPA 2.2.2	MEPS NLO	SHERPA	CT10
tZ	MG5_AMC	LO	PYTHIA 6	CTEQ6L1
t̄tW	MG5_AMC (SHERPA 2.1.1)	NLO (LO multileg)	PYTHIA 8 (SHERPA)	NNPDF 3.0 NLO (NNPDF 3.0 NLO)
t̄t(Z/γ* → ll)	MG5_AMC	NLO	PYTHIA 8	NNPDF 3.0 NLO
t̄tH	MG5_AMC (MG5_AMC)	NLO (NLO)	PYTHIA 8 (HERWIG++)	NNPDF 3.0 NLO [Ball:2014uwa] (CT10 [ct10])
tHqb	MG5_AMC	LO	PYTHIA 8	CT10
tHW	MG5_AMC (SHERPA 2.1.1)	NLO (LO multileg)	HERWIG++ (SHERPA)	CT10 (NNPDF 3.0 NLO)
tWZ	MG5_AMC	NLO	PYTHIA 8	NNPDF 2.3 LO
t̄t, t̄t̄t̄	MG5_AMC	LO	PYTHIA 8	NNPDF 2.3 LO
t̄tW+W-	MG5_AMC	LO	PYTHIA 8	NNPDF 2.3 LO
t̄t	Powheg-BOX v2 [powhegtt]	NLO	PYTHIA 8	NNPDF 3.0 NLO
t̄tγ	MG5_AMC	LO	PYTHIA 8	NNPDF 2.3 LO
s-, t-channel, Wt single top	Powheg-BOX v1 [powhegstp]	NLO	PYTHIA 6	CT10
qqVV, VVV				
Z → l+l-	SHERPA 2.2.1	MEPS NLO	SHERPA	NNPDF 3.0 NLO

Sample	DSID
WZ	364253, 364739-42
VV	364250, 364254, 364255, 363355-60, 364890
t̄W	410155
t̄Z	410156, 410157, 410218-20
low mass t̄Z	410276-8
Rare Top	410397, 410398, 410399
single Top	410658-9, 410644-5
three Top	304014
four Top	410080
t̄WW	410081
Z + jets	364100-41
low mass Z + jets	364198-215
W + jets	364156-97
Vγ	364500-35
tZ	410560
tW	410013-4
WtZ	410408
VVV	364242-9
VH	342284-5
WtH	341998
t̄tγ	410389
t̄t	410470
t̄tH	345873-5, 346343-5

Table 2: List of Monte Carlo samples by data set ID used in the analysis.

138 4 Object Reconstruction

139 All regions defined in this analysis share a common lepton, jet, and overall event preselection.
 140 The selection applied to each physics object is detailed here; the event preselection, and the
 141 selection used to define the various fit regions, is described in Section ??.

142 4.1 Trigger

143 Events are required to be selected by dilepton triggers, as summarized in Table ??.

Dilepton triggers (2015)	
$\mu\mu$ (asymm.)	HLT_mu18_mu8noL1
ee (symm.)	HLT_2e12_lhloose_L12EM10VH
$e\mu, \mu e$ (\sim symm.)	HLT_e17_lhloose_mu14
Dilepton triggers (2016)	
$\mu\mu$ (asymm.)	HLT_mu22_mu8noL1
ee (symm.)	HLT_2e17_lhvloose_nod0
$e\mu, \mu e$ (\sim symm.)	HLT_e17_lhloose_nod0_mu14
Dilepton triggers (2017)	
$\mu\mu$ (asymm.)	HLT_mu22_mu8noL1
ee (symm.)	HLT_2e24_lhvloose_nod0
$e\mu, \mu e$ (\sim symm.)	HLT_e17_lhloose_nod0_mu14
Dilepton triggers (2018)	
$\mu\mu$ (asymm.)	HLT_mu22_mu8noL1
ee (symm.)	HLT_2e24_lhvloose_nod0
$e\mu, \mu e$ (\sim symm.)	HLT_e17_lhloose_nod0_mu14

Table 3: List of lowest p_T -threshold, un-prescaled dilepton triggers used for 2015–2018 data taking.

4.2 Light leptons

- 144 Electron candidates are reconstructed from energy clusters in the electromagnetic calorimeter
 145 that are associated with charged particle tracks reconstructed in the inner detector [3]. Electron
 146 candidates are required to have $p_T > 10$ GeV and $|\eta_{\text{cluster}}| < 2.47$. Muon candidates are
 147 reconstructed by combining inner detector tracks with track segments or full tracks in the muon
 148 spectrometer [4]. Muon candidates are required to have $p_T > 10$ GeV and $|\eta| < 2.5$. Candidates
 149 in the transition region between different electromagnetic calorimeter components, $1.37 <$
 150 $|\eta_{\text{cluster}}| < 1.52$, are rejected. A multivariate likelihood discriminant combining shower shape
 151 and track information is used to distinguish real electrons from hadronic showers (fake electrons).
 152 To further reduce the non-prompt electron contribution, the track is required to be consistent
 153 with originating from the primary vertex; requirements are imposed on the transverse impact
 154 parameter significance ($|d_0|/\sigma_{d_0} < 5$) and the longitudinal impact parameter ($|\Delta z_0 \sin \theta_\ell| < 0.5$
 155 mm).
 156
- 157 Muon candidates are reconstructed by combining inner detector tracks with track segments or
 158 full tracks in the muon spectrometer [4]. Muon candidates are required to have $p_T > 10$ GeV
 159 and $|\eta| < 2.5$. The longitudinal impact parameter is the same for both electrons and muons, while
 160 muons are required to pass a slightly tighter transverse impact parameter, $|d_0|/\sigma_{d_0} < 3$.

161 All leptons are required to be isolated, as defined through the standard `isolationFixedCutLoose`
 162 working point supported by combined performance groups. Leptons are additionally required to
 163 pass a non-prompt BDT selection described in detail in [5]. Optimized working points and scale
 164 factors for this BDT are taken from that analysis.

165 4.3 Jets

166 Jets are reconstructed from calibrated topological clusters built from energy deposits in the
 167 calorimeters [6], using the anti- k_t algorithm with a radius parameter $R = 0.4$. Jets with energy
 168 contributions likely arising from noise or detector effects are removed from consideration [7],
 169 and only jets satisfying $p_T > 25$ GeV and $|\eta| < 2.5$ are used in this analysis. For jets with
 170 $p_T < 60$ GeV and $|\eta| < 2.4$, a jet-track association algorithm is used to confirm that the jet
 171 originates from the selected primary vertex, in order to reject jets arising from pileup collisions
 172 [8].

173 4.4 B-tagged Jets

174 In order to make a measurement of $WZ +$ heavy flavor it is necessary to distinguish these events
 175 from $WZ +$ light jets. For this purpose, the DL1r b-tagging algorithm is used to distinguish
 176 heavy flavor jets from lighter ones. The DL1r algorithm uses jet kinematics, particularly jet
 177 vertex information, as input for a neural network which assigns each jet a score designed to
 178 reflect how likely that jet is to have originated from a b-quark.

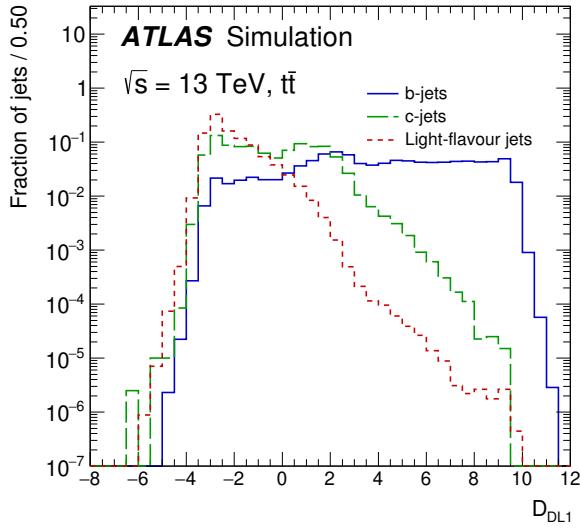


Figure 1: Output distribution of the DL1r algorithm for b-jets, charm jets, and light jets

179 From the output of the BDT, working points (WPs) are developed based on the efficiency of truth
 180 b-jets at particular values of the DL1r algorithm. The working points used in this analysis are
 181 summarized in Table ??.

WP	none	loose	medium	tight	tightest
b eff.	-	85%	77%	70%	60%

Table 4: B-tagging Working Points by tightness and b-jet efficiency

182 A tighter WP will accept fewer b-jets, but reject a higher fraction of charm and light jets.
 183 Generally, analyses that include b-jets will use a fixed working point, for example, requiring that
 184 a jet pass the 70% threshold. By instead treating these working point as bins, e.g. events with
 185 jets that fall between the 85% and 77% WPs fall into one bin, while events with jets passing the
 186 60% WP fall into another, and looking at the full psuedo-continuous DL1r spectrum of the jets,
 187 additional information can be gained. The psuedo-continuous b-tag spectrum is used in this case
 188 to separate out WZ + b, WZ + charm, and WZ + light.

189 **4.5 Missing transverse energy**

190 Missing transverse momentum (E_T^{miss}) is used as part of the event selection. The missing
 191 transverse momentum vector is defined as the inverse of the sum of the transverse momenta of
 192 all reconstructed physics objects as well as remaining unclustered energy, the latter of which is
 193 estimated from low- p_T tracks associated with the primary vertex but not assigned to a hard object,
 194 with object definitions taken from [9]. Light leptons considered in the E_T^{miss} reconstruction are
 195 required to have $p_T > 10 \text{ GeV}$, while jets are required to have $p_T > 20 \text{ GeV}$.

196 **5 Event Selection and Signal Region Definitions**

197 Event are required to pass a preselection described in Section ?? and summarized in Table ??.
 198 Those that pass this preselection are divided into various fit regions described in Section ??,
 199 based on the number of jets in the event, and the b-tag score of those jets.

200 **5.1 Event Preselection**

201 Events are required to include exactly three reconstructed light leptons passing the requirement
 202 described in ??, which have a total charge of ± 1 . As the opposite sign lepton is found to be
 203 prompt the vast majority of the time [5], it is required to be loose and isolated, as defined though
 204 the standard `isolationFixedCutLoose` working point supported by combined performance

205 groups. The same sign leptons are required to be very tightly isolated, as per the recommended
 206 `isolationFixedCutTight`.

207 The leptons are ordered in the analysis code as 0, 1, and 2. Lepton 0 is the lepton whose charge
 208 is opposite the other two. Lepton 1 is the lepton closest to the opposite charge lepton, i.e. the
 209 smallest ΔR , leaving lepton 2 as the lepton further from the opposite charge lepton. Lepton 0
 210 is required to have $p_T > 10$ GeV, while the same sign leptons, 1 and 2, are required to have
 211 $p_T > 20$ GeV to reduce the contribution of non-prompt leptons.

212 The invariant mass of at least one pair of opposite sign, same flavor leptons is required to fall
 213 within 10 GeV of the mass of the Z boson, 91.2 GeV. Events where one of the opposite sign pairs
 214 have an invariant mass less than 12 GeV are rejected in order to suppress low mass resonances.

215 An additional requirement is placed on the missing transverse energy, $E_T^{\text{miss}} > 20$ GeV, and the
 216 transverse mass of the W candidate, $m(E_T^{\text{miss}} + l_{\text{other}}) > 30$ GeV, where E_T^{miss} is the missing
 217 transverse energy, and l_{other} is the lepton not included in the Z-candidate.

218 Events are required to have one or two reconstructed jets passing the selection described in
 219 Section ???. Events with more than two jets are rejected in order to reduce the contribution of
 220 backgrounds such as $t\bar{t}Z$ and $t\bar{t}W$, which tend to have higher jet multiplicity.

Event Selection

Exactly three leptons with charge ± 1

Two tight Iso, tight ID same-charge leptons with $p_T > 20$ GeV

One loose Iso, medium ID opposite charge lepton with $p_T > 10$ GeV
 $m(l^+l^-)$ within 10 GeV of 91.2 GeV

Transverse mass of W-candidate, $m_T(E_T^{\text{miss}} + \text{lep}_{\text{other}}) > 30$ GeV

Missing transverse energy, $E_T^{\text{miss}} > 20$ GeV

One or two jets with $p_T > 25$ GeV

Table 5: Summary of the selection applied to events for inclusion in the fit

221 The event yields in the preselection region for both data and Monte Carlo are summarized in
 222 Table ??, which shows good agreement between data and Monte Carlo, and demonstrates that
 223 this region consists primarily of WZ events. The WZ events are split into WZ + b, WZ + c, and
 224 WZ + l based on the truth flavor of the associated jet in the event. Specifically, this determination
 225 is made based on the `HadronConeExclTruthLabelID` of the jet, as recommended by the b-
 226 tagging working group [BtagWG]. In this ordering b-jet supersedes charm, which supersedes
 227 light. That is, WZ + l events contain no charm and no b jets at truth level, WZ + c contain at
 228 least one truth charm and no b-jets, and WZ + b contains at least one truth b-jet.

Process	Events
WZ + b	167.6 ± 6.5
WZ + c	1080 ± 40
WZ + l	7220 ± 310
Other VV	850 ± 140
t̄tW	16.8 ± 2.3
t̄tZ	115 ± 17
rare Top	2.2 ± 0.1
Single top	0.10 ± 0.45
Three top	0.01 ± 0.01
Four top	0.02 ± 0.01
t̄tWW	0.23 ± 0.05
Z + jets	600 ± 260
V + γ	37 ± 54
tZ	190 ± 70
tW	5.5 ± 1.2
WtZ	25.8 ± 1.1
VVV	26.2 ± 0.9
VH	94 ± 7
t̄t	108.68 ± 8
t̄tH	4.3 ± 0.5
Total	10600 ± 530
Data	10574

Table 6: Event yields in the preselection region at 139.0 fb^{-1}

²²⁹ Here Other VV represents diboson processes other than WZ, and consists predominantly of
²³⁰ $ZZ \rightarrow ll\bar{l}\bar{l}$ events where one of the leptons is not reconstructed.

²³¹ Simulations are further validated by comparing the kinematic distributions of the Monte Carlo
²³² with data, which are shown in Figure ???. Here, bins with 5% or more WZ+b are blinded.

WZ Fit Region - Inclusive

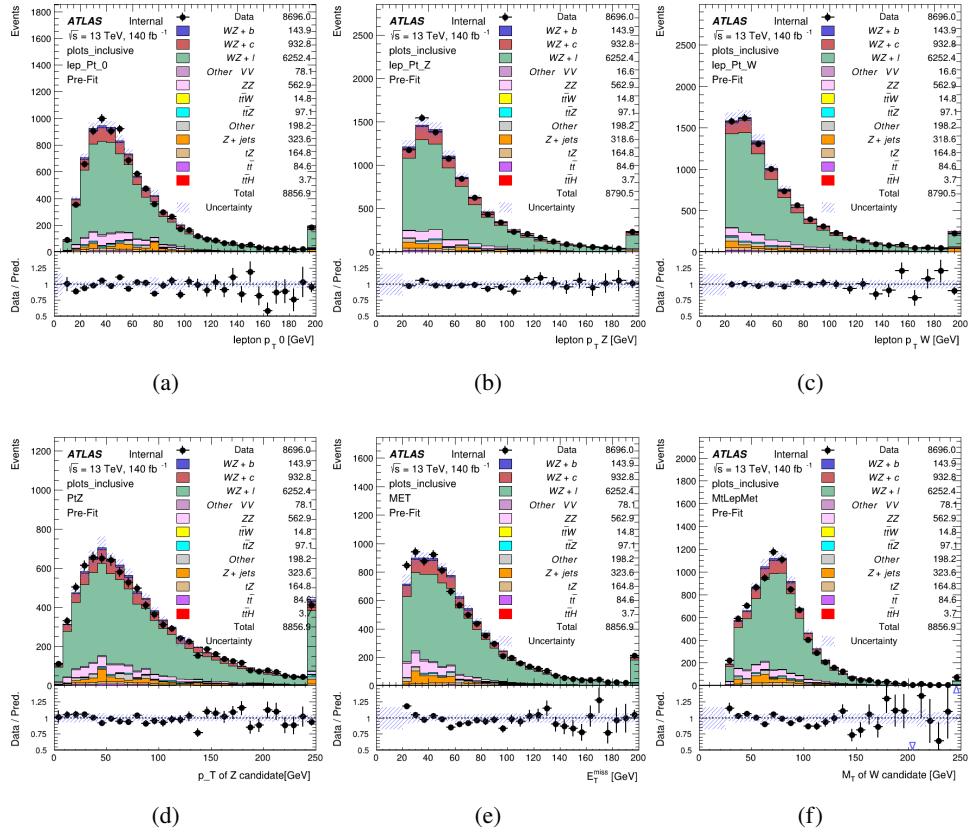


Figure 2: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the M_T of the W candidate.

233 5.2 Fit Regions

Once preselection has been applied, the remaining events are categorized into one of twelve orthogonal regions. The regions used in the fit are summarized in Table ??.

Table 7: A list of the regions used in the fit and the selection used for each.

Region	Selection
1j, <85%	$N_{\text{jets}} = 1, n\text{Jets_DL1r_85} = 0$
1j, 85%-77%	$N_{\text{jets}} = 1, n\text{Jets_DL1r_85} = 1, n\text{Jets_DL1r_77}=0$
1j, 77%-70%	$N_{\text{jets}} = 1, n\text{Jets_DL1r_77} = 1, n\text{Jets_DL1r_70}=0$
1j, 70%-60%	$N_{\text{jets}} = 1, n\text{Jets_DL1r_70} = 1, n\text{Jets_DL1r_60}=0$
1j, >60%	$N_{\text{jets}} = 1, n\text{Jets_DL1r_60} = 1, tZ \text{ BDT} > 0.725$
1j tZ CR	$N_{\text{jets}} = 1, n\text{Jets_DL1r_60} = 1, tZ \text{ BDT} < 0.725$
2j, <85%	$N_{\text{jets}} = 2, n\text{Jets_DL1r_85} = 0$
2j, 85%-77%	$N_{\text{jets}} = 2, n\text{Jets_DL1r_85} >= 1, n\text{Jets_DL1r_77}=0$
2j, 77%-70%	$N_{\text{jets}} = 2, n\text{Jets_DL1r_77} >= 1, n\text{Jets_DL1r_70}=0$
2j, 70%-60%	$N_{\text{jets}} = 2, n\text{Jets_DL1r_70} >= 1, n\text{Jets_DL1r_60}=0$
2j, >60%	$N_{\text{jets}} = 2, n\text{Jets_DL1r_60} >= 1, tZ \text{ BDT} > 0.725$
2j tZ CR	$N_{\text{jets}} = 2, n\text{Jets_DL1r_60} >= 1, tZ \text{ BDT} < 0.725$

236 The working points discussed in Section ?? are used to separate events into fit regions based on
 237 the highest working point reached by a jet in each event. Because the background composition
 238 differs significantly based on the number of b-jets, events are further subdivided into 1 jet and 2
 239 jet regions in order to minimize the impact of background uncertainties.

240 An additional tZ control region is created based on the BDT described in Section ?? . The region
 241 with 1-jet passing the 60% working point is split in two - a signal enriched region of events with
 242 a BDT score greater than 0.03, and a tZ control region including events with less than 0.03. This
 243 cutoff is arrived at by performing an Asimov fit with a variety of cutoffs, and selecting the value
 244 that produces the highest significance for the measurement of WZ + b.

245 The modeling in each region is validated by comparing data and MC predictions for various
 246 kinematic distributions. These plot are shown in Figures ??-??.

WZ Fit Region - 1j Inclusive

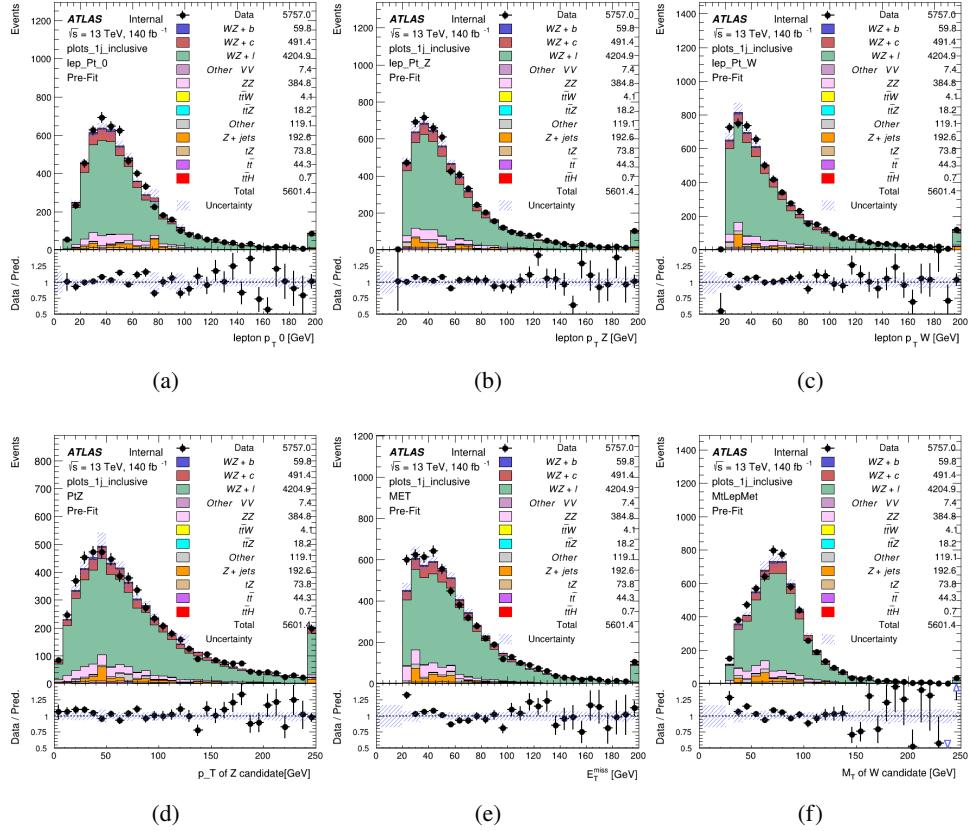


Figure 3: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 1j < 85% WP

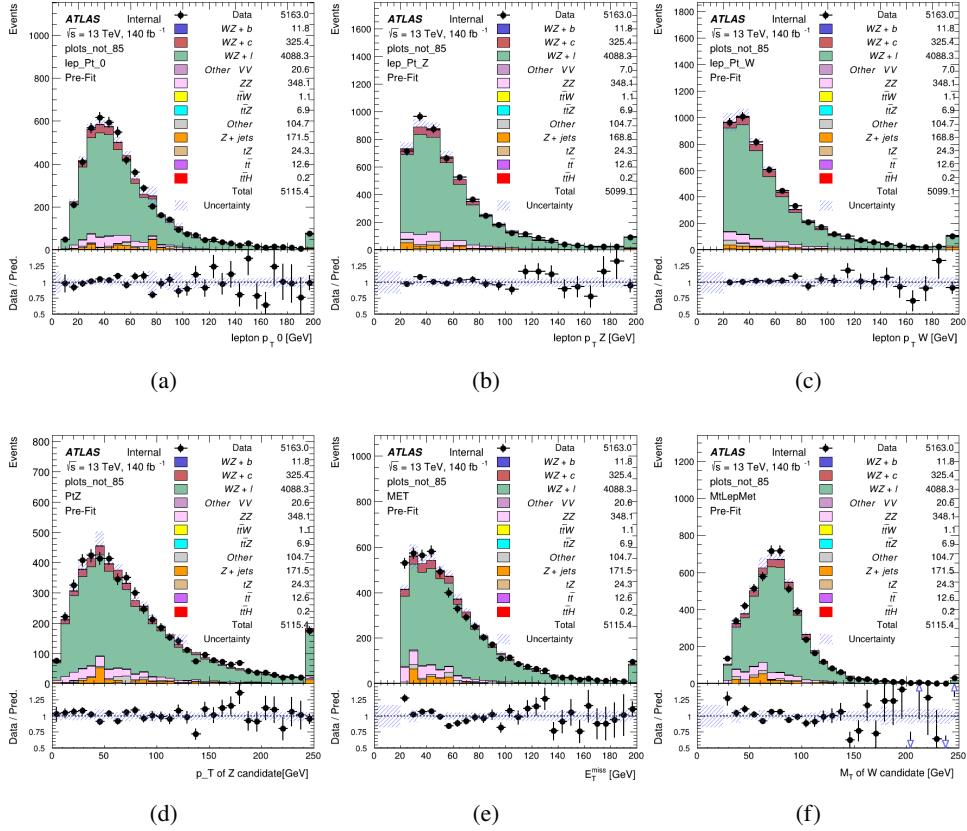


Figure 4: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 1j 77-85% WP

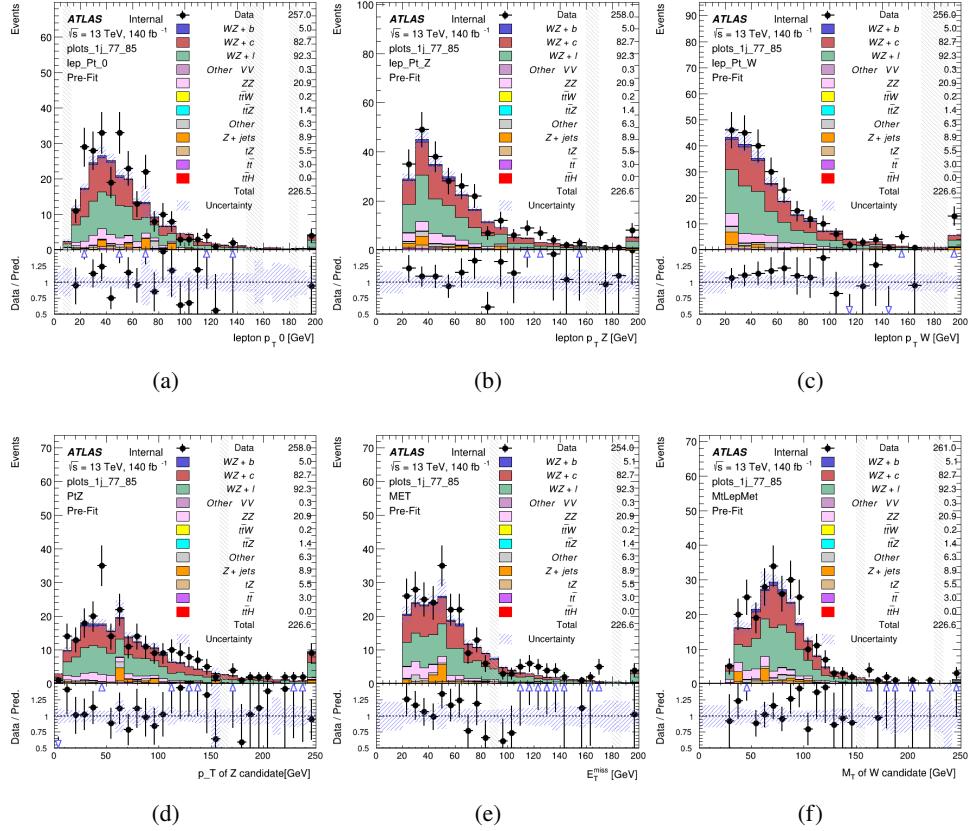


Figure 5: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 1j 70-77% WP

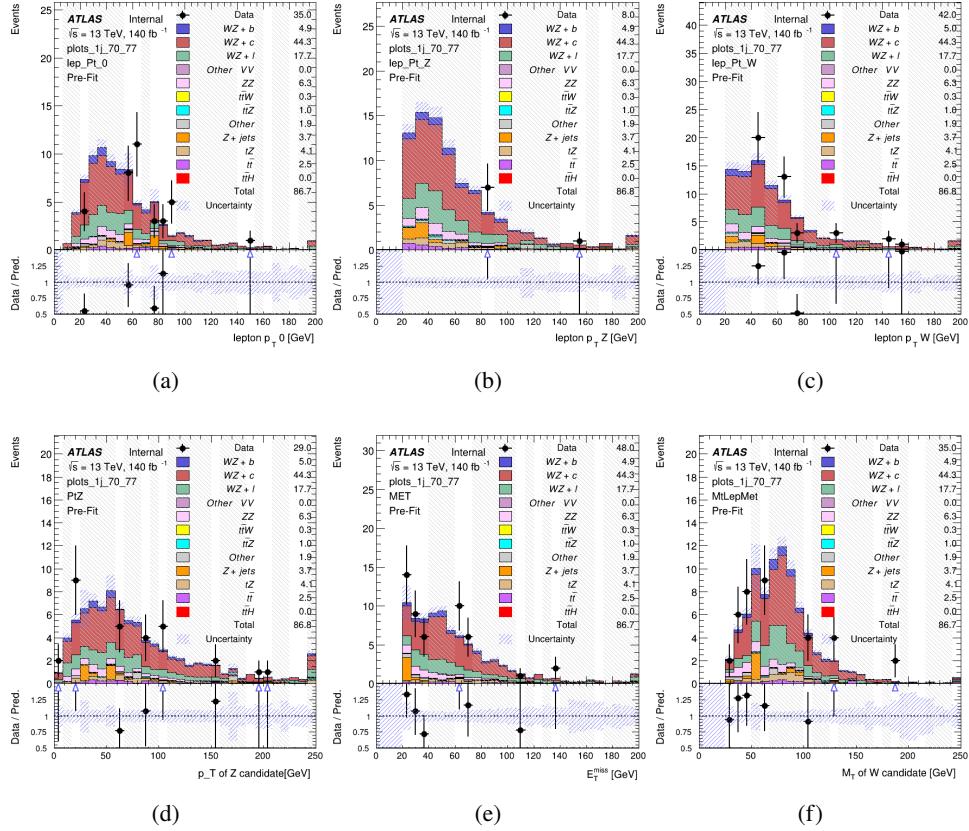


Figure 6: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 1j 60-70% WP

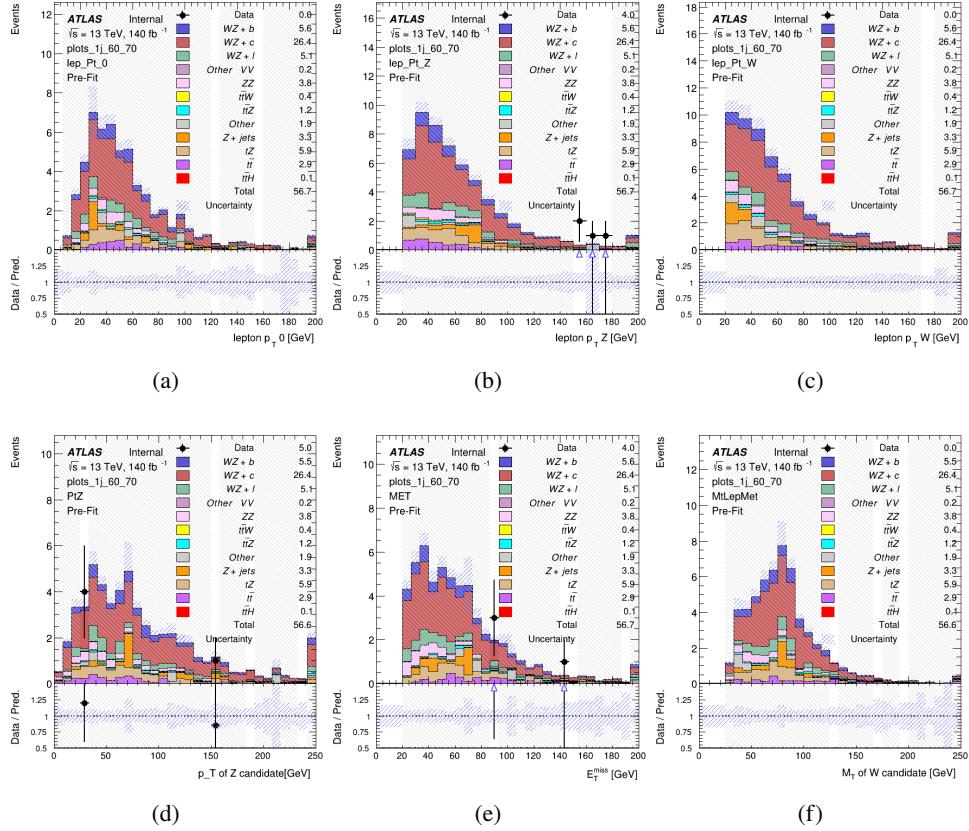


Figure 7: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 1j 60% WP

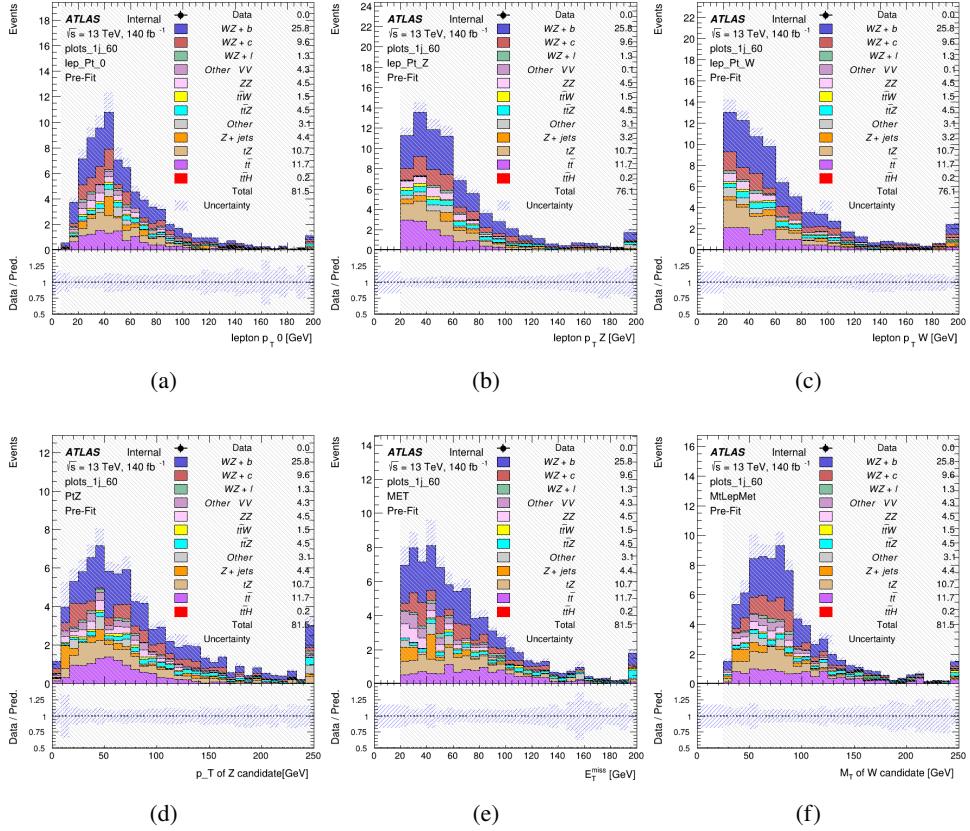


Figure 8: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - tZ-CR

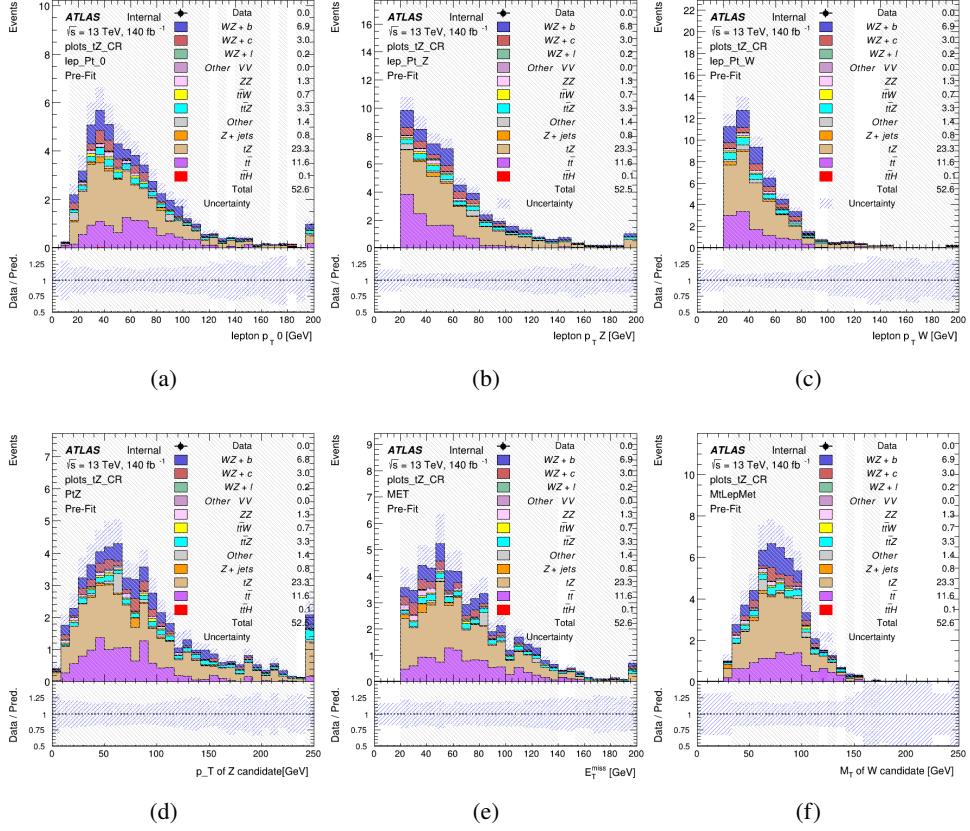


Figure 9: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 2j Inclusive

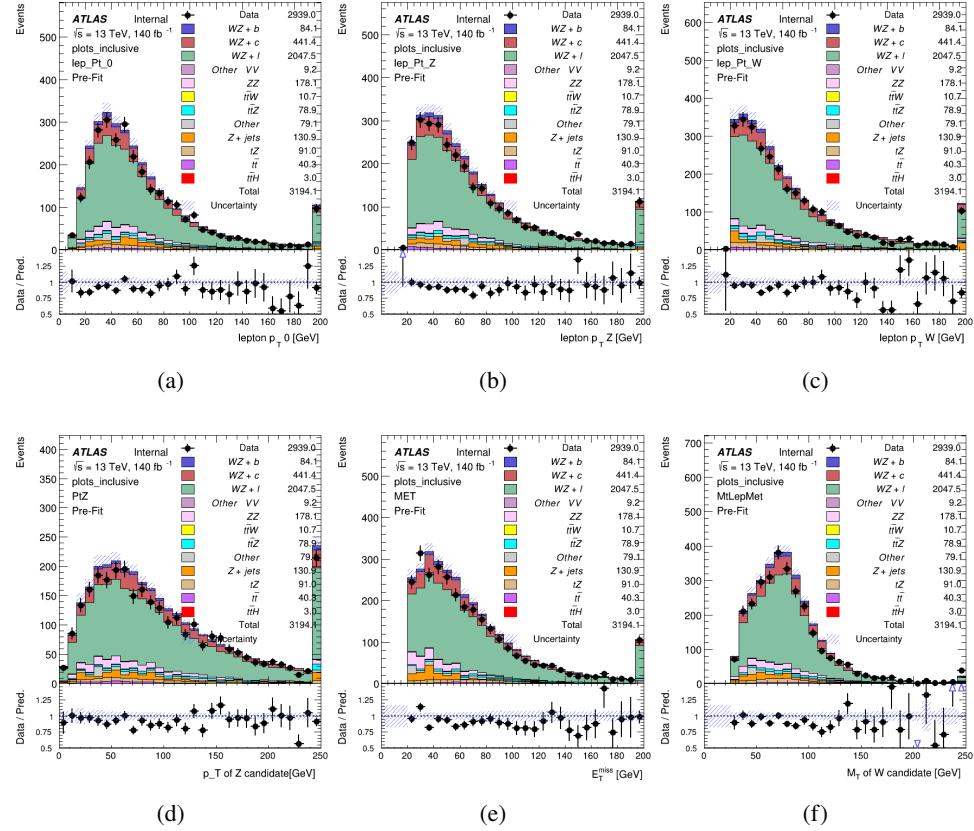


Figure 10: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 2j < 85% WP

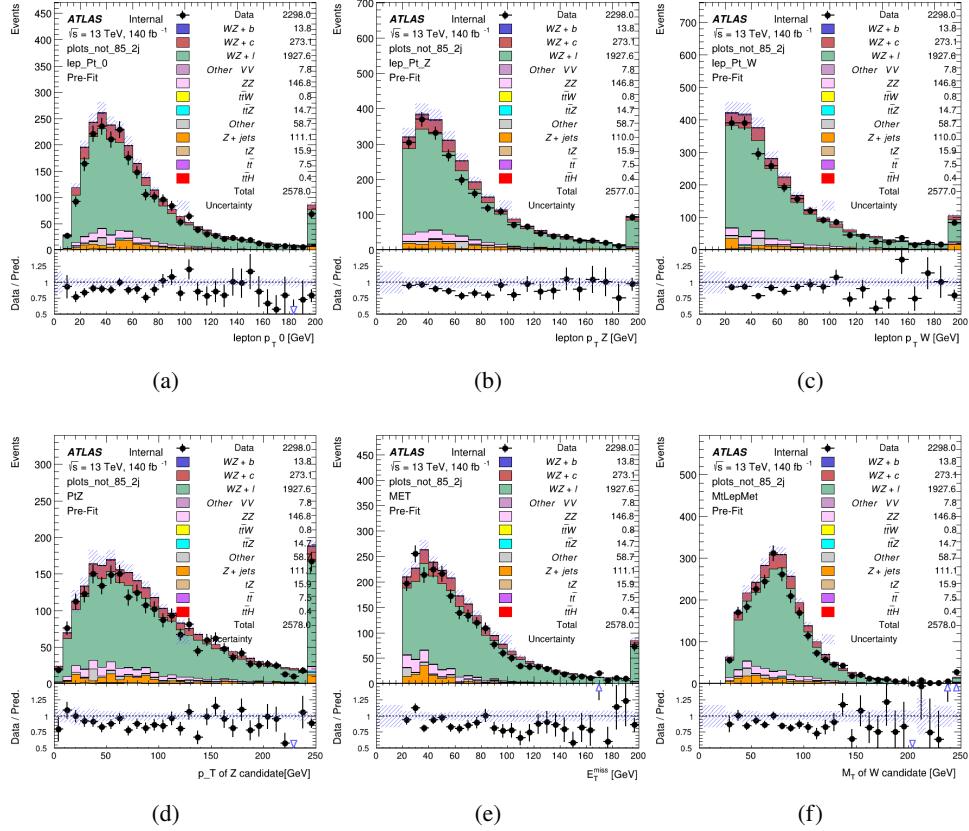


Figure 11: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 2j 77-85% WP

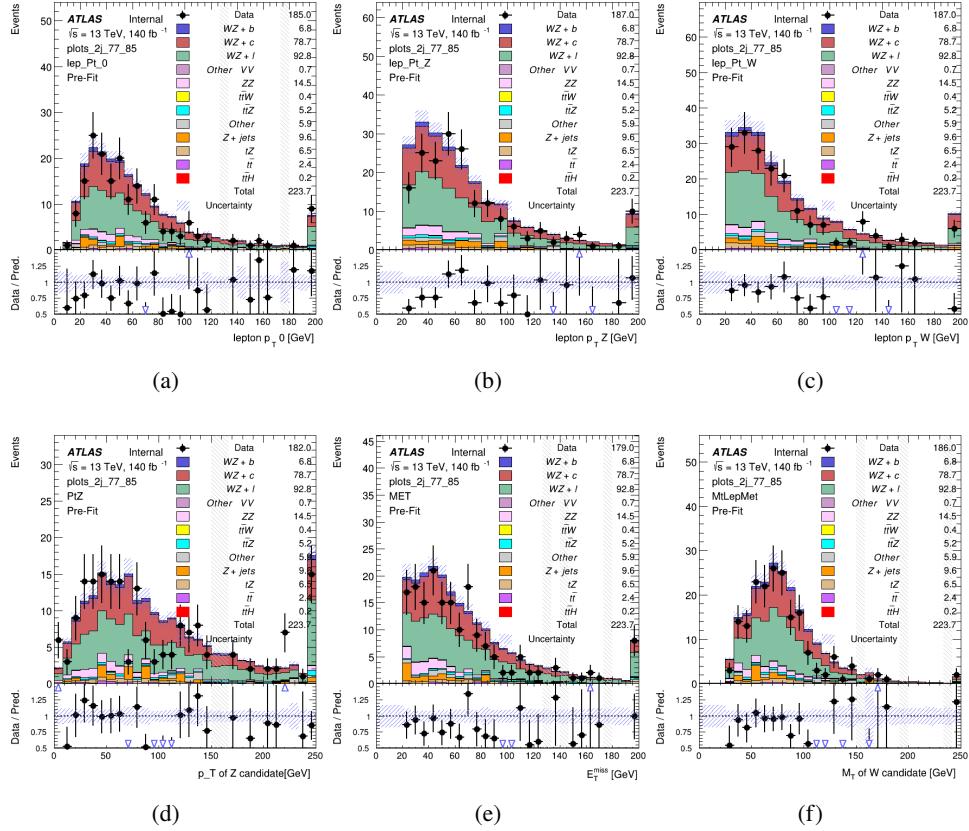


Figure 12: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 2j 70-77% WP

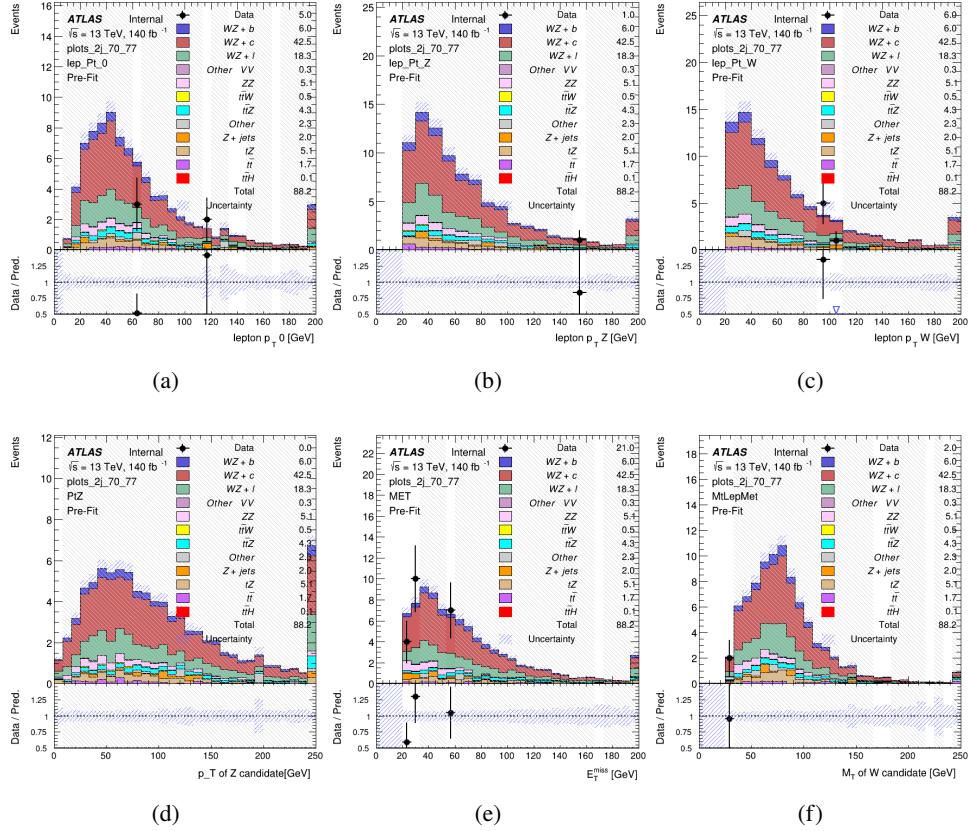


Figure 13: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - 2j 60-70% WP

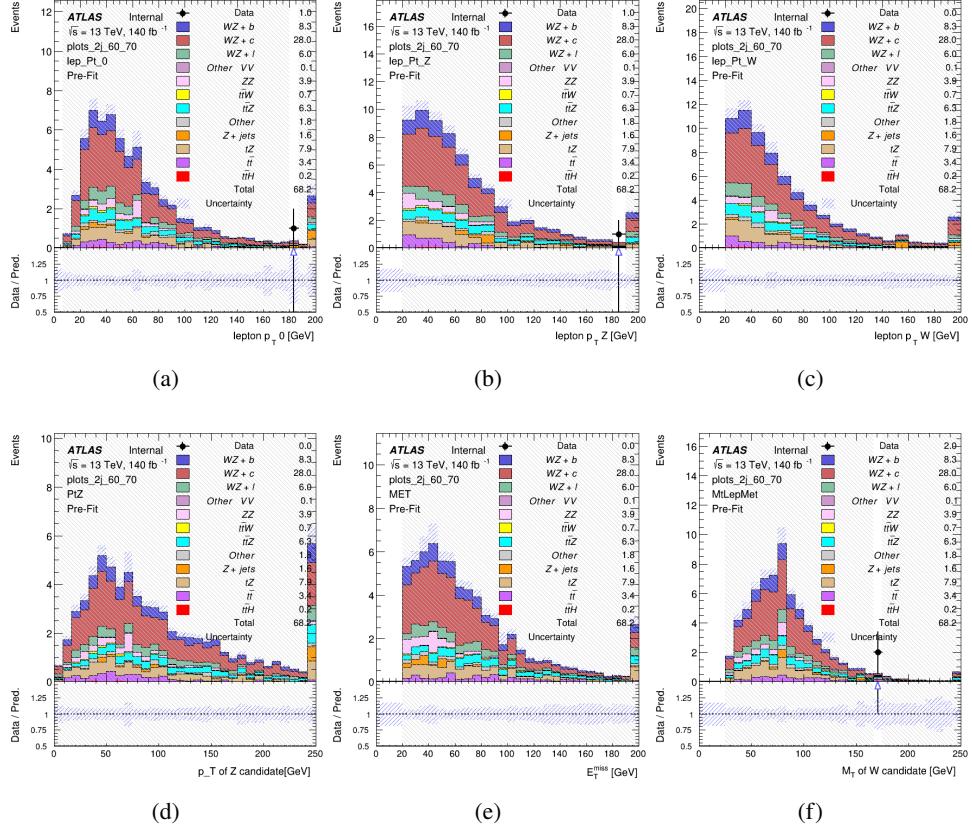


Figure 14: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

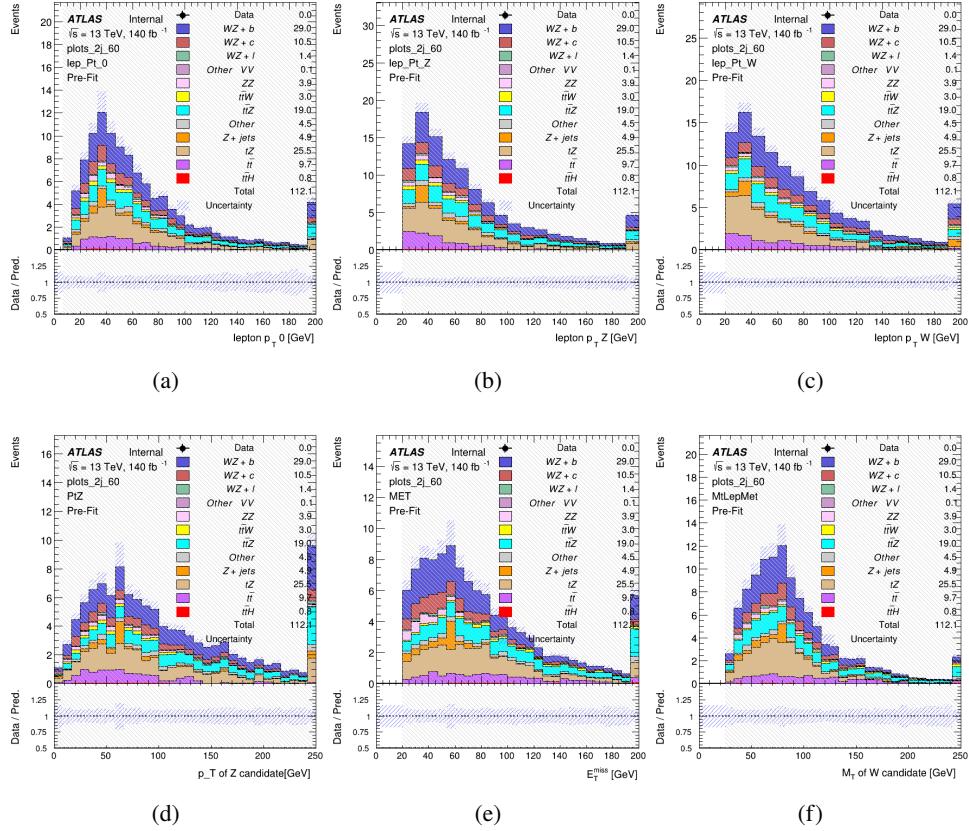
WZ Fit Region - 2j 60% WP

Figure 15: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

WZ Fit Region - tZ-CR-2j

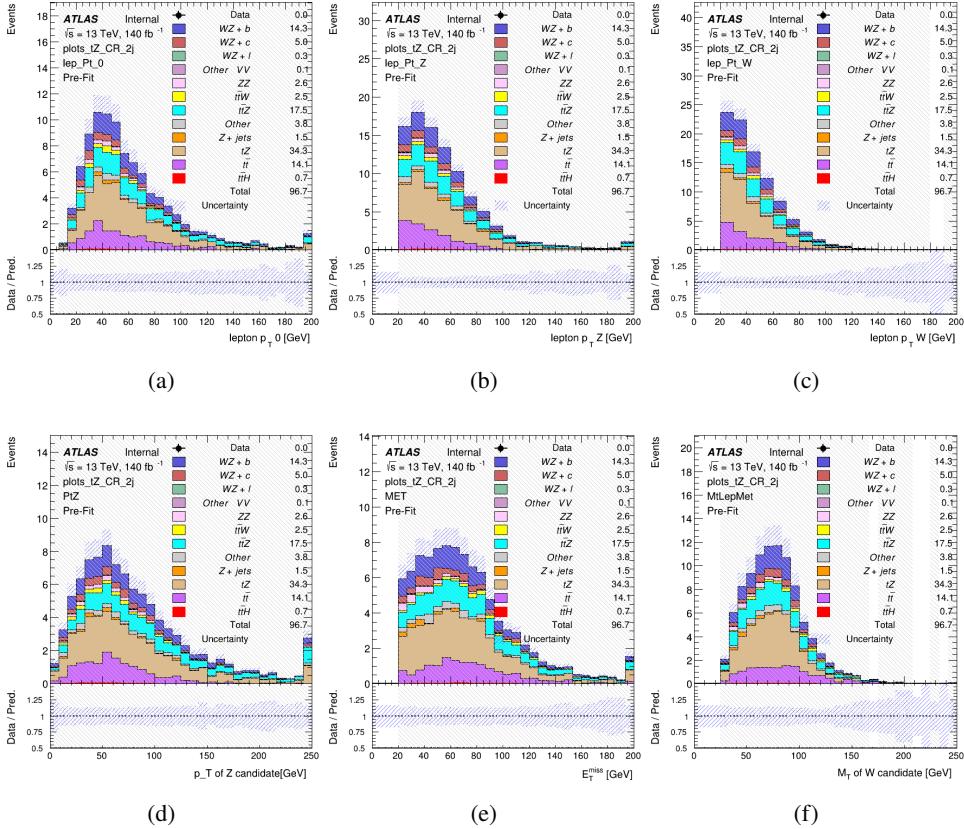


Figure 16: Comparisons between the data and MC distributions in the preselection region for (a) the p_T of the opposite sign lepton, (b) the p_T of the other lepton from the Z candidate, (c) the p_T of the lepton from the W candidate, (d) the p_T of the Z candidate, (e) the E_T^{miss} , and (f) the m_T of the W candidate.

5.3 Non-Prompt Lepton Estimation

Two processes act as sources of non-prompt leptons appear in the analysis: $t\bar{t}$ and $Z+jet$ production both produce two prompt leptons, and each contribute to the 3l region when an additional non-prompt lepton appears in the event. The contribution of these processes is estimated with Monte Carlo simulations, which are validated using enriched validation regions.

5.3.1 $t\bar{t}$ Validation

$t\bar{t}$ events can produce two prompt leptons from the decay of each of the tops. These top decays produce two b-quarks, the decay of which can produce additional non-prompt leptons, which occasionally pass the event preselection. In order to validate that the Monte Carlo accurately

256 simulates this process accurately, the MC prediction in a non-prompt $t\bar{t}$ enriched validation
257 region is compared to data.

258 The $t\bar{t}$ validation region is similar to the preselection region - three leptons meeting the criteria
259 described in Section ?? are required, and the requirements on E_T^{miss} remain the same. However,
260 the selection requiring a lepton pair form a Z-candidate are reversed. Events where the invariant
261 mass of any two opposite sign, same flavor leptons falls within 10 GeV of 91.2 GeV are rejected.
262 This ensures the $t\bar{t}$ validation region is orthogonal to the preselection region.

263 Further, because the jet multiplicity of $t\bar{t}$ events tends to be higher than WZ, the number of jets
264 in each event is required to be greater than 1. As b-jets are almost invariably produced from top
265 decays, at least one b-tagged jet passing the 70% DL1r WP in each event is required. Various
266 kinematic plots of this region are shown in Figure ??.

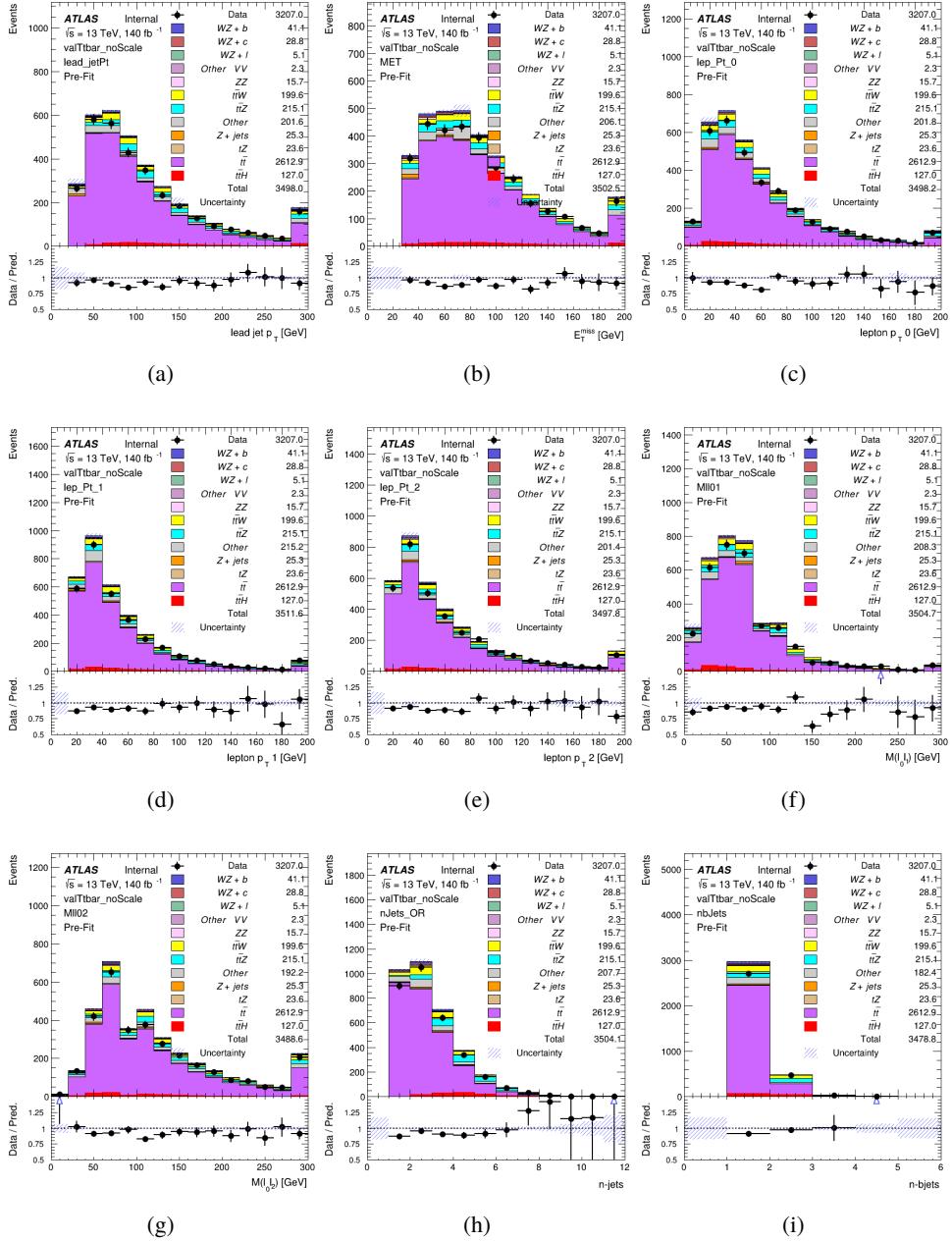


Figure 17: Comparisons between the data and MC distributions in the $t\bar{t}$ validation region for (a) the p_T of the leading jet, (b) the missing transverse energy, (c) the p_T of lepton 0, (d) p_T of lepton 1, (e) p_T of lepton 2, (f) the invariant mass of leptons 0 and 1, (g) the invariant mass of leptons 0 and 2, (h) the number of jets, (i) the number of b-tagged jets.

267 The shape of each distribution agrees quite well between data and MC, with a constant offset
268 between the two. This is accounted for by applying a constant correction factor of 0.883 to the $t\bar{t}$

269 MC prediction. Plots showing the kinematics of the $t\bar{t}$ VR after this correction factor has been
 270 applied are shown in Figure ??.

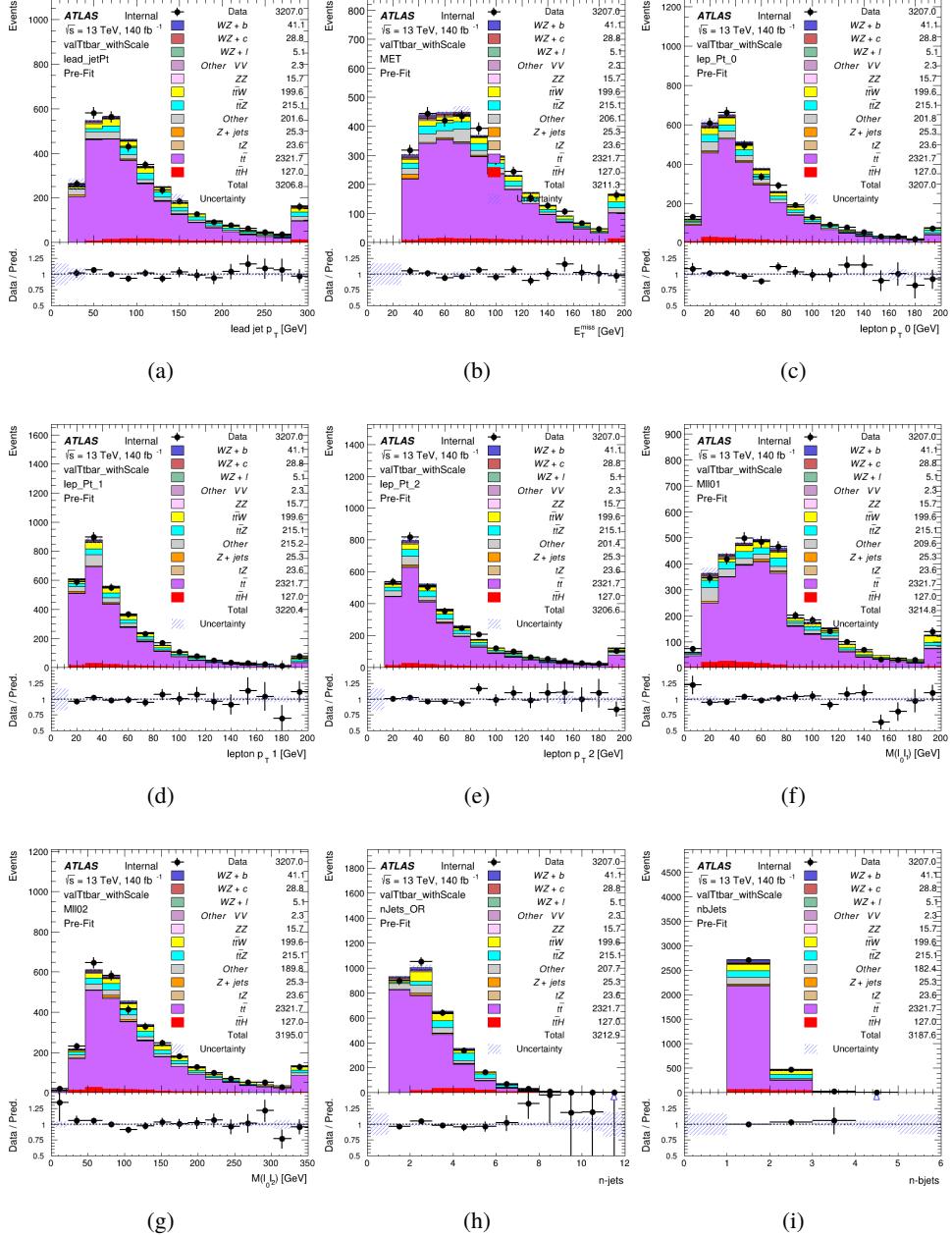


Figure 18: Comparisons between the data and MC distributions in the $t\bar{t}$ validation region after the correction factor has been applied for (a) the p_T of the leading jet, (b) the missing transverse energy, (c) the p_T of lepton 0, (d) p_T of lepton 1, (e) p_T of lepton 2, (f) the invariant mass of leptons 0 and 1, (g) the invariant mass of leptons 0 and 2, (h) the number of jets, (i) the number of b-tagged jets.

271 The modeling is further validated by looking at the yield in the $t\bar{t}$ VR for each DL1r WP, giving
 272 a clearer correspondence to the signal regions used in the fit. Each region shown in Figure ??
 273 requires one or more jets pass the listed WP, with no jets passing the next highest WP.

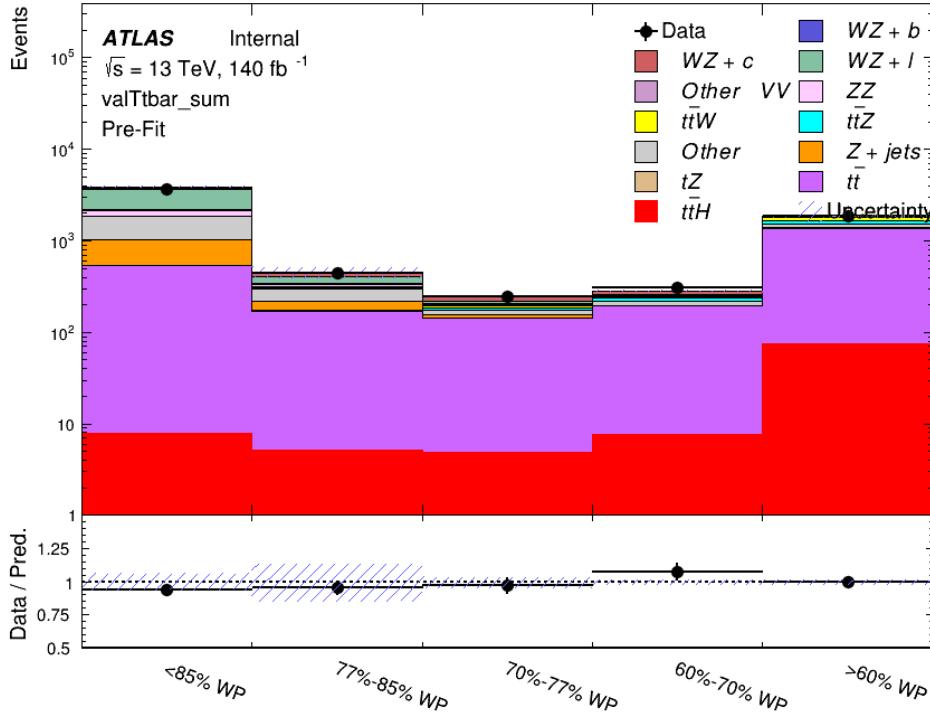


Figure 19: Data and MC comparisons for each DL1r WP for both 1-jet and 2-jet regions, after the $t\bar{t}$ VR selection and correction factor have been applied

274 As data and MC are found to agree within 10% for each of these working points, a 10% systematic
 275 uncertainty on the $t\bar{t}$ prediction is included for the analysis.

276 5.3.2 Z+jets Validation

277 Similar to $t\bar{t}$, a non-prompt Z+jets validation region is produced in order to validate the MC
 278 predictions. The lepton requirements remain the same as the preselection region. Because no
 279 neutrinos are present for this process, the E_T^{miss} cut is reversed, requiring $E_T^{\text{miss}} < 30 \text{ GeV}$. This
 280 also ensures this validation region is orthogonal to the preselection region. Further, the number
 281 of jets in each event is required to be greater than or equal to one. Various kinematic plots of this
 282 region are shown below. The general agreement between data and MC in each of these suggests
 283 that the non-prompt contribution of Z+jets is well modeled by Monte Carlo.

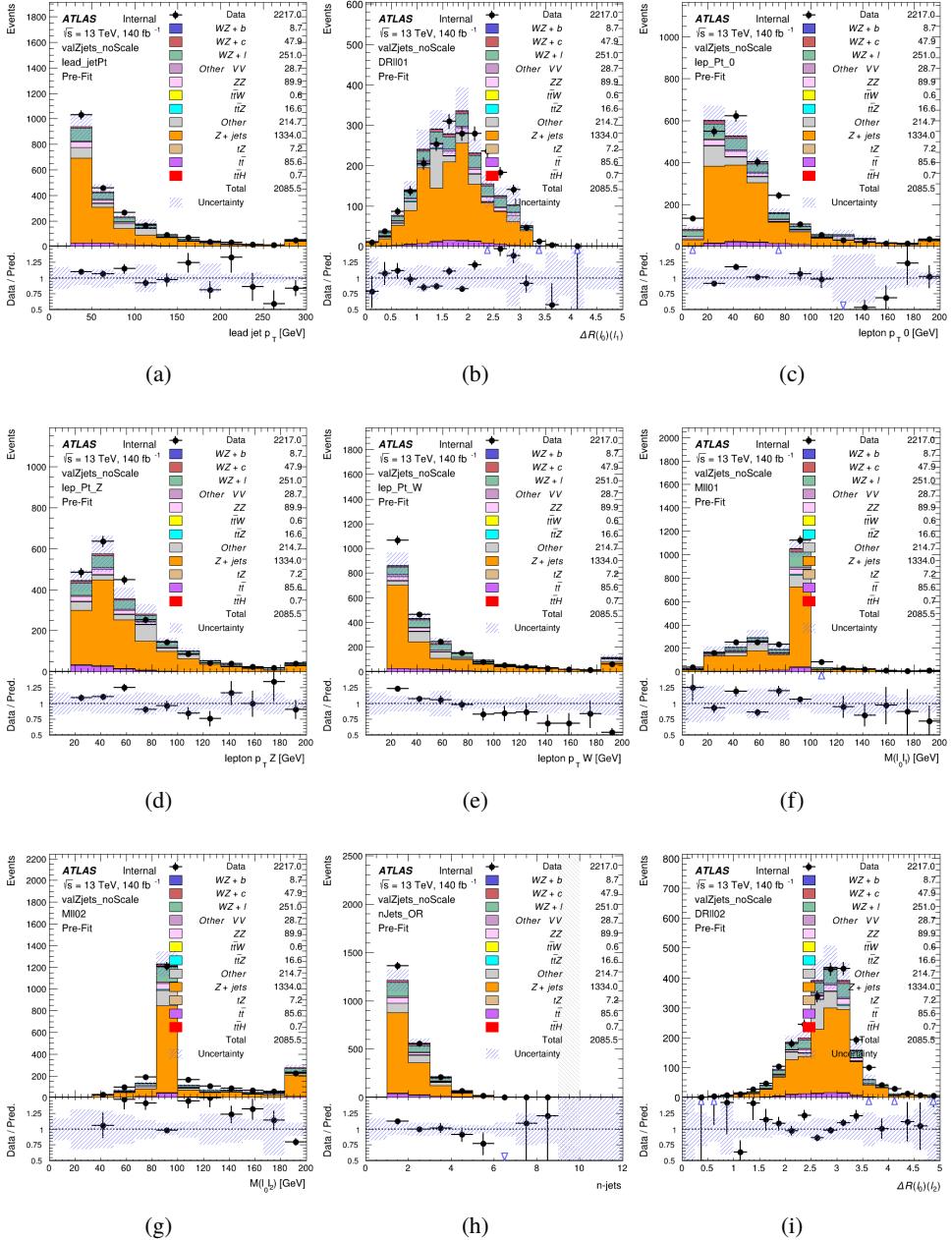


Figure 20: Comparisons between the data and MC distributions in the Z+jets validation region for (a) the p_T of the leading jet, (b) ΔR between leptons 0 and 1, (c) the p_T of lepton 0, (d) p_T of SS lepton from the Z candidate, (e) p_T of the SS lepton from the W candidate, (f) the invariant mass of leptons 0 and 1, (g) the invariant mass of leptons 0 and 2, (h) the number of jets, (i) ΔR between leptons 0 and 2. Includes only statistical uncertainties

284 While there is general agreement between data and MC within statistical uncertainty, the shape

285 of the p_T spectrum of the lepton from the W is found to differ. To account for this discrepancy,
286 a variable correction factor is applied to Z+jets. χ^2 minimization of the lepton 2 p_T spectrum is
287 performed to derive a correction factor of $1.53 - 6.6 * 10^{-6}(\text{lep}_\text{Pt}_W)$. Kinematic plots of the
288 Z + jets validation region after this correction factor has been applied are shown in Figure ??.

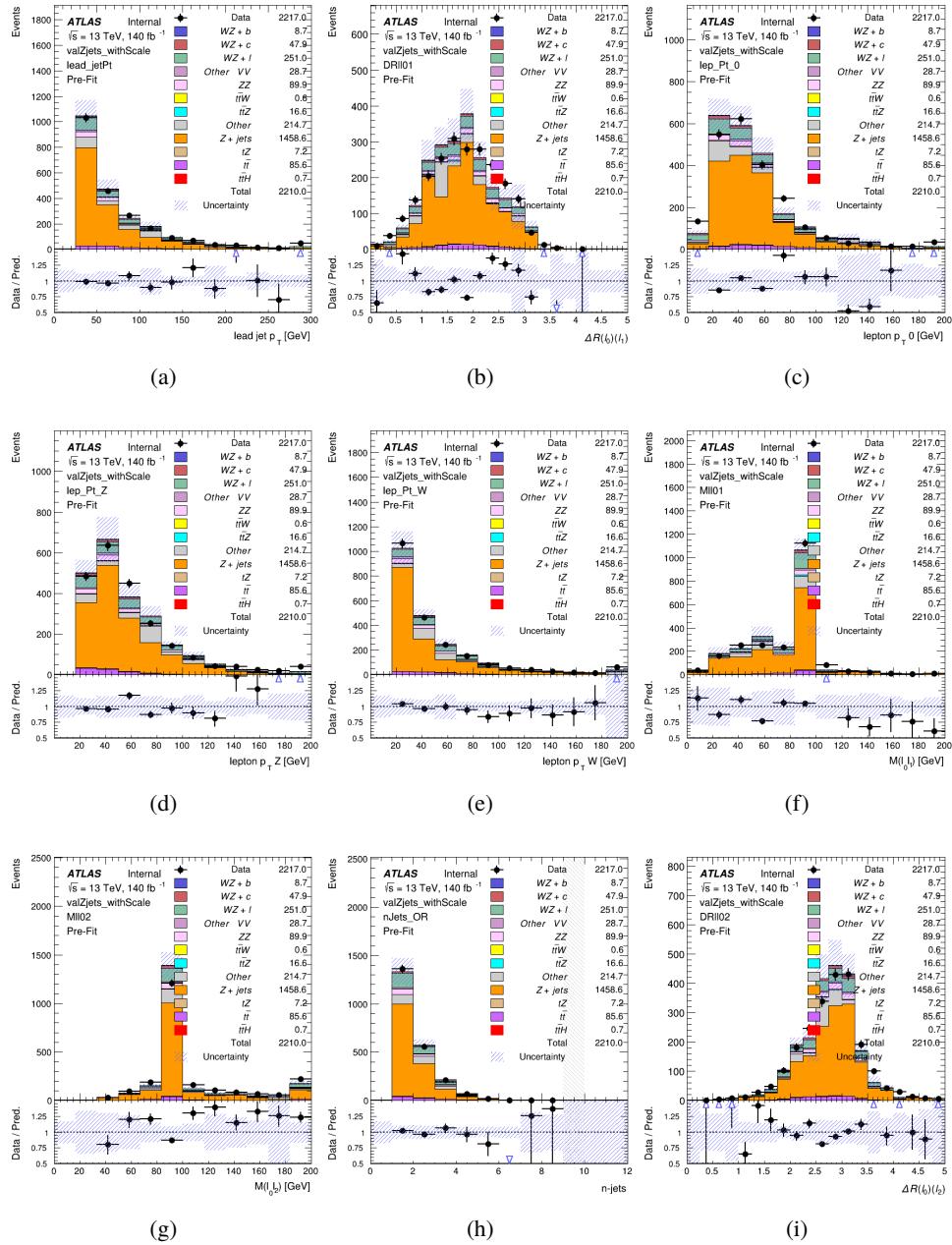


Figure 21: Comparisons between the data and MC distributions in the Z+jets validation region after the correction factor has been applied for (a) the p_T of the leading jet, (b) ΔR between leptons 0 and 1, (c) the p_T of lepton 0, (d) p_T of SS lepton from the Z candidate, (e) p_T of the SS lepton from the W candidate, (f) the invariant mass of leptons 0 and 1, (g) the invariant mass of leptons 0 and 2, (h) the number of jets, (i) ΔR between leptons 0 and 2

289 The modeling is further validated by looking at the yield in the Z+jets VR for each DL1r WP,

290 giving a clearer correspondence to the signal regions used in the fit. Each region shown in Figure
 291 ?? requires one or more jets pass the listed WP, with no jets passing the next highest WP.

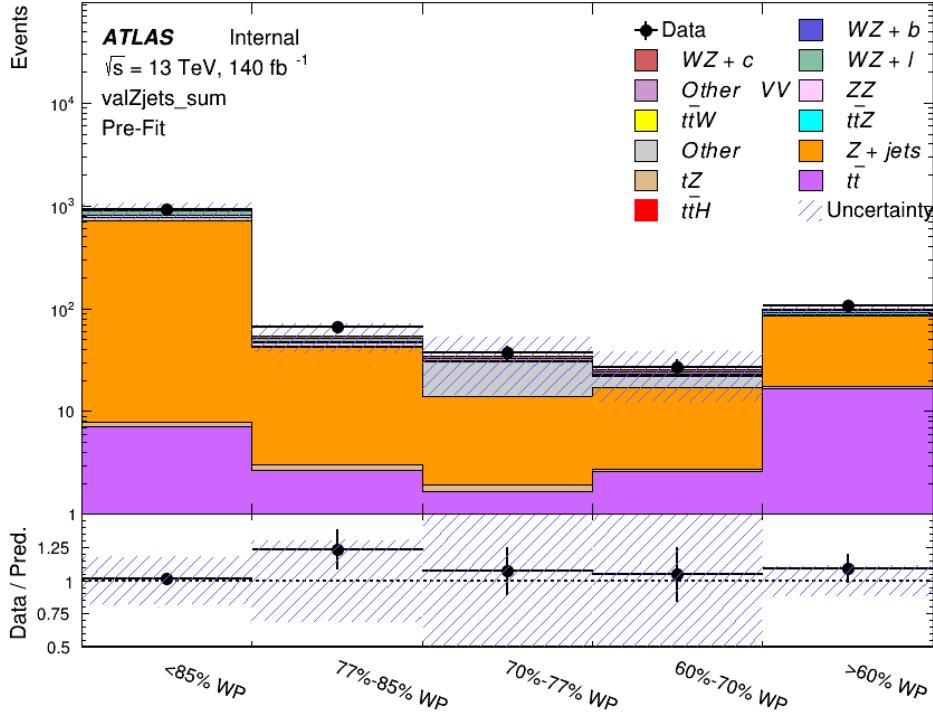


Figure 22: Data and MC comparisons for each DL1r WP for both 1-jet and 2-jet regions, after the Z+jets VR selection and correction factor have been applied

292 For each of the b-tagging working points considered, the data falls within 20% of the MC
 293 prediction once this correction factor has been applied. Therefore, a 20% systematic uncertainty
 294 is applied to Z + jets in the analysis.

295 6 tZ Interference Studies and Separation Multivariate Analysis

296 Because tZ produces a final state identical to signal, it represents a predominant background in
 297 the most signal enriched regions. That is, the region with one jet passing the 60% DL1r WP.
 298 Therefore, a boosted decision tree (BDT) algorithm is trained using TMVA [10] to separate WZ
 299 + heavy flavor from tZ.

300 Separation between tZ and WZ + heavy flavor is achieved in part by reconstructing the invariant
 301 mass of the top candidate, which clusters more closely to the top mass for tZ than WZ + heavy
 302 flavor.

303 The result of this BDT is used to create a tZ enriched region in the fit, reducing its impact on the
 304 measurement of WZ + heavy flavor.

305 **6.1 Top Mass Reconstruction**

306 The reconstruction of the top mass follows the procedure described in detail in section 6.1 of
 307 [11]. The mass of the top quark candidate is reconstructed from the jet, the lepton not included
 308 in the Z-candidate, and a reconstructed neutrino. In the case that there is one jet in the event,
 309 there is only possible b-jet candidate. For events with two jets, the jet with the highest DL1r
 310 score is used.

311 The neutrino from the W decay is expected to be the only source of E_T^{miss} . Therefore, the E_T
 312 and ϕ of the neutrino are taken from the E_T^{miss} measurement. This leaves the z-component of
 313 the neutrino momentum, $p_{\nu z}$ as the only unknown.

314 This unknown is solved for by taking the combined invariant mass of the lepton and neutrino to
 315 give the invariant mass of the W boson:

$$316 \quad (p_l + p_\nu)^2 = m_W^2$$

317 Expanding this out into components, this equation gives:

$$318 \quad \sqrt{p_{T\nu}^2 + p_{z\nu}^2} E_l = \frac{m_w^2 - m_l^2}{2} + p_{T\nu}(p_{lx}\cos\phi_\nu + p_{ly}\sin\phi_\nu) + p_{lz}p_{\nu z}$$

319 This equation gives two solutions for $p_{\nu z}$. For cases where only one of these solutions is real,
 320 that is taken as the value of $p_{\nu z}$. For instances with two real solutions, the one which is shown
 321 to be correct in the largest fraction of simulations is taken. For cases when no real solution is
 322 found, often because of detector effects, the value of E_T^{miss} is varied in decreasing increments of
 323 100 MeV until a real solution is found.

324 The reconstructed top mass distribution for tZ and WZ + b can be seen in figure ??.

325 **6.2 tZ BDT**

326 A Boosted Decision Tree (BDT), specifically XGBoost [xgboost_cite], is used to provide separa-
 327 tion between tZ and WZ+b. The following kinematic variables are used as inputs:

- 328 • The invariant mass of the reconstructed top candidate
- 329 • p_T of each of the leptons, jet
- 330 • The invariant mass of each combination of lepton pairs, $M(l\bar{l})$
- 331 • E_T^{miss}

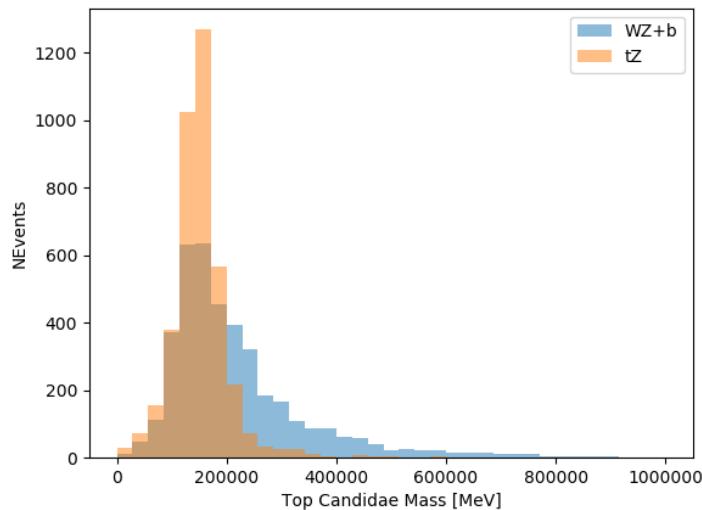


Figure 23: Reconstructed top mass distributions for tZ and WZ + b, measured in MeV.

- 332 • Distance between each combination of leptons, $\Delta R(l\bar{l})$
- 333 • Distance between each lepton and the jet, $\Delta R(lj)$

334 The training samples included only events meeting the requirements of the 1-jet, $>60\%$ region,
 335 i.e. passing all the selection described in section ?? and having exactly one jet which passes the
 336 tightest (60%) DL1r working point.

337 The distributions of a few of these features for both signal and background is shown in figure
 338 ??.

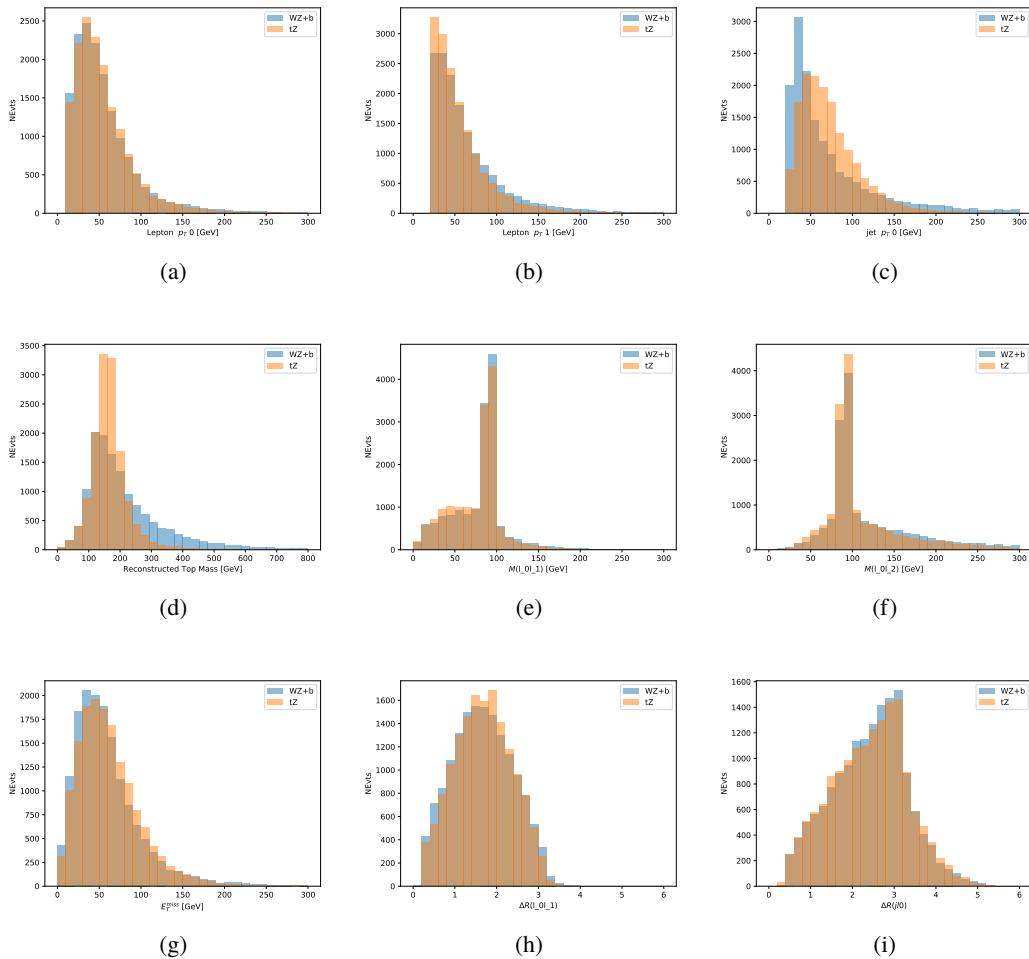


Figure 24: Distribution of input features of the BDT for signal (WZ) and background (tZ). Both are scaled to an equal number of events. (a), (b) and (c) show the p_T of lepton 0, lepton 1, and the jet, (d) show the reconstructed top mass, (e) and (f) show the invariant mass of leptons 0 and 1, leptons 0 and 2. (g) shows the E_T^{miss} of each event. (h) and (i) show the ΔR between lepton 0 and lepton 1, and the jet.

339 A sample of 20,000 background (tZ) and signal (WZ+b) Monte Carlo events are used to train
 340 the BDT. And additional 5,000 events are reserved for testing the model, in order to prevent
 341 over-fitting. A total of 750 decision trees with a maximum depth of 6 branches are used to build
 342 the model. These parameters are chosen empirically, by training several models with different
 343 parameters and selecting the one that gave the best separation for the test sample.

344 The results of the BDT training are shown in figure ???. The output scores for both signal and
 345 background events is shown on the left. The right shows the receiving operating characteristic
 346 (ROC) curve that results from the MVA. The ROC curve represents the background rejection
 347 as a function of signal efficiency, where each point on the curve represents a different response

348 score. The ROC curve of the BDT is compared to the performance of using an optimal set of flat
 349 selections on the same set of input variables.

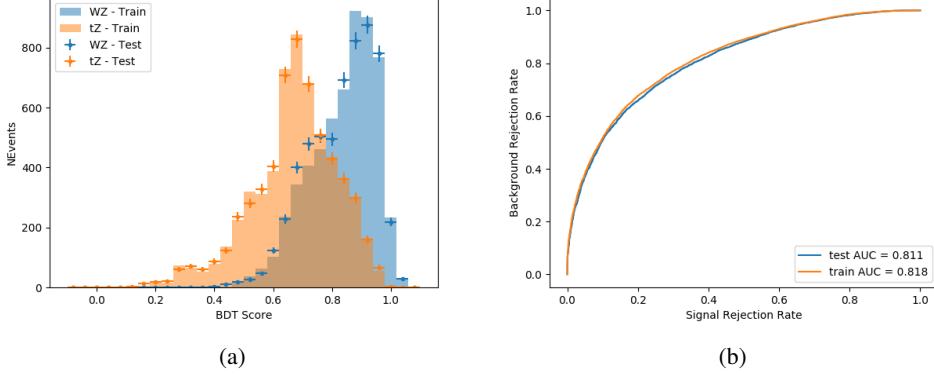


Figure 25: Distribution of the BDT response for signal and background events on the left, the ROC curve for the BDT on the right.

350 The relative important of each input feature in the model, measured by how often they appeared
 351 in the decision trees, is shown in figure ??.

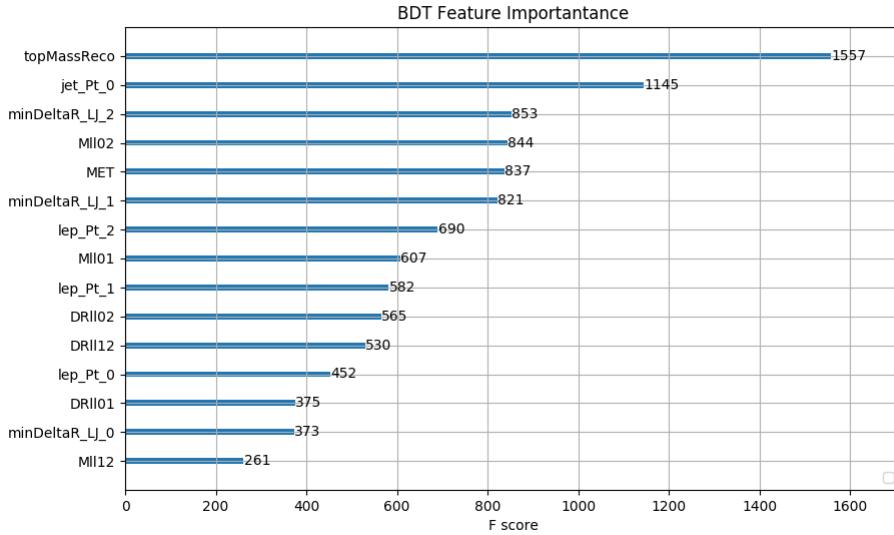


Figure 26: Relative importance of each input feature in the model.

352 These results suggest that some amount of separation can be achieved between these two pro-
 353 cesses, with a high BDT score selecting a set of events that is pure in $WZ + b$. A BDT score

354 of 0.725 is selected as a cutoff, where events with scores higher than this form a signal enriched
 355 region, and events with scores lower than this form a tZ control region. This cutoff is selected by
 356 varying the value of this cutoff in stat-only Asimov fits, and selecting the value that minimizes
 357 the statistical uncertainty on WZ + b.

358 7 Systematic Uncertainties

359 The systematic uncertainties that are considered are summarized in Table ???. These are imple-
 360 mented in the fit either as a normalization factors or as a shape variation or both in the signal
 361 and background estimations. The numerical impact of each of these uncertainties is outlined in
 362 Section ???.

Table 8: Sources of systematic uncertainty considered in the analysis.

Systematic uncertainty	Components
Luminosity	1
Pileup reweighting	1
Physics Objects	
Electron	6
Muon	15
Jet energy scale	28
Jet energy resolution	8
Jet vertex fraction	1
Jet flavor tagging	131
E_T^{miss}	3
Total (Experimental)	194
Signal Modeling	
Shape modelling	3
Renormalization and factorization scales	5
Background Modeling	
Cross section	15
Renormalization and factorization scales	12
Total (Signal and background modeling)	35
Total (Overall)	229

363 The uncertainty in the combined integrated luminosity is derived from a calibration of the
 364 luminosity scale performed for 13 TeV proton-proton collisions [12], [LUCID2].

365 The experimental uncertainties are related to the reconstruction and identification of light leptons
 366 and b-tagging of jets, and to the reconstruction of E_T^{miss} . The TOTAL electron ID correlation

367 model is used, corresponding to 1 electron ID systematic. Electron ID is found to be a subleading
368 systematic that is unconstrained by the fit, making it an appropriate choice for this analysis.

369 The sources which contribute to the uncertainty in the jet energy scale (JES) [13] are decom-
370 posed into uncorrelated components and treated as independent sources in the analysis. The
371 CategoryReduction model is used to account for JES uncertainties, which decomposes the uncer-
372 tainties into 30 nuisance parameters included in the fit. The SimpleJER model is used to account
373 for jet energy resolution (JER) uncertainties, and 8 JER uncertainty components unclued as
374 NPs in the fit.

375 The uncertainties in the b-tagging efficiencies measured in dedicated calibration analyses [14] are
376 also decomposed into uncorrelated components. The large number of components for b-tagging
377 is due to the calibration of the distribution of the MVA discriminant.

378 The full list of systematic uncertainties considered in the analysis is summarized in Tables ??,
379 ?? and ??.

380

Experimental Systematics on Leptons and E_T^{miss}			
Type	Description	Systematics Name	Application
Trigger			
Scale Factors	Trigger Efficiency	lepSFTrigTight_MU(EL)_SF_Trigger_STAT(SYST)	Event Weight
Muons			
Efficiencies	Reconstruction and Identification	lepSFObjTight_MU_SF_ID_STAT(SYST)	Event Weight
	Isolation	lepSFObjTight_MU_SF_Isol_STAT(SYST)	Event Weight
	Track To Vertex Association	lepSFObjTight_MU_SF_TTVA_STAT(SYST)	Event Weight
p_T Scale	p_T Scale	MUONS_SCALE	p_T Correction
Resolution	Inner Detector Energy Resolution	MUONS_ID	p_T Correction
	Muon Spectrometer Energy Resolution	MUONS_MS	p_T Correction
Electrons			
Efficiencies	Reconstruction	lepSFObjTight_EL_SF_ID	Event Weight
	Identification	lepSFObjTight_EL_SF_Reco	Event Weight
	Isolation	lepSFObjTight_EL_SF_Isol	Event Weight
Scale Factor	Energy Scale	EG_SCALE_ALL	Energy Correction
Resolution	Energy Resolution	EG_RESOLUTION_ALL	Energy Correction
E_T^{miss}			
Soft Tracks Terms	Resolution	MET_SoftTrk_ResoPerp	p_T Correction
	Resolution	MET_SoftTrk_ResoPara	p_T Correction
	Scale	MET_SoftTrk_ScaleUp	p_T Correction
	Scale	MET_SoftTrk_ScaleDown	p_T Correction

Table 9: Summary of experimental systematics considered for leptons and E_T^{miss} . Includes type, description, name of systematic as used in the fit, and mode of application. The mode of application indicates the systematic evaluation, e.g. as an overall event re-weighting (Event Weight) or rescaling (p_T Correction).

Experimental Systematics on Jets			
Type	Origin	Systematics Name	Application
Jet Vertex Tagger		JVT	Event Weight
Energy Scale	Calibration Method	JET_21NP_ JET_EffectiveNP_1-19	p_T Correction p_T Correction
	η inter-calibration	JET_EtaIntercalibration_Modelling JET_EtaIntercalibration_NonClosure JET_EtaIntercalibration_TotalStat	p_T Correction p_T Correction p_T Correction
	High p_T jets	JET_SingleParticle_HighPt	p_T Correction
	Pile-Up	JET_Pileup_OffsetNPV JET_Pileup_OffsetMu JET_Pileup_PtTerm JET_Pileup_RhoTopology	p_T Correction p_T Correction p_T Correction p_T Correction
	Non Closure	JET_PunchThrough_MC15	p_T Correction
	Flavour	JET_Flavor_Response JET_BJES_Response JET_Flavor_Composition	p_T Correction p_T Correction p_T Correction
Resolution		JET_JER_SINGLE_NP	Event Weight

Table 10: Jet systematics take into account effects of jets calibration method, η inter-calibration, high p_T jets, pile-up, and flavor response. They are all diagonalised into effective parameters.

Experimental Systematics on b-tagging		
Type	Origin	Systematic Name
Scale Factors	DL1r b-tagger efficiency on b originated jets in bins of η	DL1r_Continuous_EventWeight_B0-29
	DL1r b-tagger efficiency on c originated jets in bins of η	DL1r_Continuous_EventWeight_C0-19
	DL1r b-tagger efficiency on light flavoured originated jets in bins of η and p_T	DL1r_Continuous_EventWeight_Light0-79
	DL1r b-tagger extrapolation efficiency	DL1r_Continuous_EventWeight_extrapolation DL1r_Continuous_EventWeight_extrapolation_from_charm

Table 11: Summary of experimental systematics to be included for b-tagging of jets in the analysis, using the continuous DL1r tagging algorithm. All of the b-tagging related systematics are applied as event weights. From left: type, description, and the name of systematic used in the fit.

- ³⁸¹ Theoretical uncertainties applied to MC predictions, including cross section, PDF, and scale
³⁸² uncertainties are taken from theory calculations, with the exception of non-prompt and diboson
³⁸³ backgrounds. The cross-section uncertainty on tZ is taken from [tZ_paper]. Derivation of the
³⁸⁴ non-prompt background uncertainties, Z+jets and tt>, are explained in detail in Section ??.
- ³⁸⁵ The other VV + heavy flavor processes (namely VV+b and VV+charm, which primarily consist
³⁸⁶ of ZZ events) are also poorly understood, because these processes involve the same physics as
³⁸⁷ WZ + heavy flavor, and have also not been measured. Therefore, a conservative 50% uncertainty
³⁸⁸ is applied to those samples. While this uncertainty is large, it is found to have little impact on
³⁸⁹ the significance of the final result.
- ³⁹⁰ The theory uncertainties applied to the predominate background estimates are summarized in
³⁹¹ Table ??.

Process	X-section [%]
tZ	X-sec: ± 15.2 QCD Scale: $^{+5.2}_{-1.3}$ PDF($+\alpha_S$): ± 1.2
t <bar>t} H (aMC@NLO+Pythia8)</bar>	QCD Scale: $^{+5.8}_{-9.2}$ PDF($+\alpha_S$): ± 3.6
t <bar>t} Z (aMC@NLO+Pythia8)</bar>	QCD Scale: $^{+9.6}_{-11.3}$ PDF($+\alpha_S$): ± 4
t <bar>t} W (aMC@NLO+Pythia8)</bar>	QCD Scale: $^{+12.9}_{-11.5}$ PDF($+\alpha_S$): ± 3.4
VV + b/charm (Sherpa 2.2.1)	± 50
VV + light (Sherpa 2.2.1)	± 6
t <bar>t}</bar>	± 10
Z + jets	± 20
Others	± 50

Table 12: Summary of theoretical uncertainties for MC predictions in the analysis.

392 Additional signal uncertainties are estimated by comparing estimates from the nominal Sherpa
 393 WZ samples with alternate WZ samples generated with Powheg+PYTHIA8 (DSID 361601).
 394 The shape of the templates used in the fit are compared between these two samples for WZ + b,
 395 WZ + charm and WZ + light, as shown in Figures ?? and ?? . Each of these plots are normalized
 396 to unity in order to capture differences in shape.

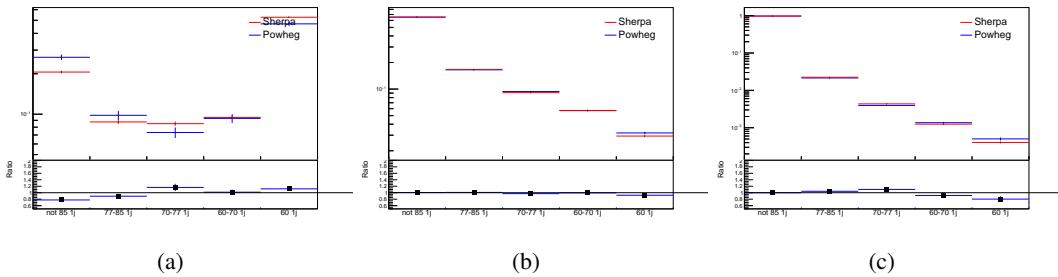


Figure 27: Comparison between Sherpa and Powheg predictions of the distribution of (a) WZ + b, (b) WZ + charm, and (c) WZ + light among the various b-tag WPs used in the 1-jet fit.

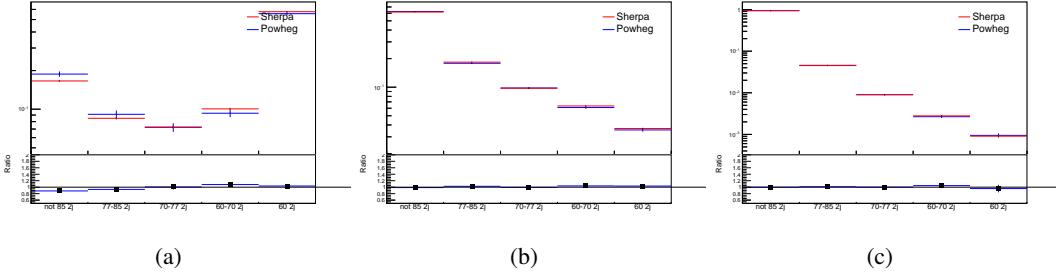


Figure 28: Comparison between Sherpa and Powheg predictions of the distribution of (a) WZ + b, (b) WZ + charm, and (c) WZ + light among the various b-tag WPs used in the 2-jet fit.

397 Separate systematics are included in the fit for WZ + b, WZ + charm and WZ + light, where
 398 the distribution among each of the fit regions is varied based on the prediction of the Powheg
 399 sample.

400 8 Results

401 A separate maximum-likelihood fit is performed over the 1-jet and 2-jet fit regions in order to
 402 extract the best-fit value of the WZ + b-jet and WZ + charm jet contributions. The WZ + b,
 403 WZ + charm and WZ + light contributions are allowed to float, with the remaining background
 404 contributions are held fixed. **The current fit strategy treats the WZ + b-jet contribution as**
 405 **the parameter of interest, with the normalization of the WZ + charm and the WZ + light**
 406 **contributions taken as systematic uncertainties. This could however be adjusted, depending**
 407 **on whether it is decided the goal of the analysis should be to measure WZ+b specifically or**
 408 **WZ + heavy flavor overall.** The result of the fit is used to extract the cross-section of WZ +
 409 heavy-flavor production.

410 A maximum likelihood fit to data is performed simultaneously in the regions described in Section
 411 ???. The parameters μ_{WZ+b} , $\mu_{WZ+charm}$, $\mu_{WZ+light}$, where $\mu = \sigma_{\text{observed}}/\sigma_{\text{SM}}$, are extracted
 412 from the fit.

413 The Asimov fit for 1-jet events gives an expected μ value of $1.00^{+0.37}_{-0.35}(\text{stat})^{+0.24}_{-0.22}(\text{sys})$ for WZ
 414 + b. The fitted cross-section modifiers for WZ + charm and WZ + light are $1.00 \pm 0.17 \pm 0.15$
 415 and $1.00 \pm 0.04 \pm 0.07$, respectively.

416 The expected cross-section of WZ+b with 1 associated jet obtained from the fit is $1.74^{+0.65}_{-0.61}(\text{stat})^{+0.42}_{-0.37}(\text{sys})$
 417 fb with an expected significance of 2.2σ . The expected cross-section of WZ + charm is measured
 418 to be $14.6 \pm 2.5(\text{stat}) \pm 2.1(\text{sys})$ fb, with a correlation of -0.23.

419 For 2-jet events, the fit gives an expected μ value of $1.00^{+0.34}_{-0.33}(\text{stat})^{+0.29}_{-0.28}(\text{sys})$ for WZ + b.
 420 The fitted cross-section modifiers for WZ + charm and WZ + light are $1.00 \pm 0.17 \pm 0.15$ and
 421 $1.00 \pm 0.04 \pm 0.08$, respectively.

⁴²² The expected $WZ + b$ cross-section in the 2-jet region is $2.46^{+0.83}_{-0.81}(\text{stat})^{+0.073}_{-0.68}(\text{sys}) \text{ fb}$ with
⁴²³ an expected significance of 2.6σ . The 2-jet expected cross-section of $WZ + \text{charm}$ is $12.7 \pm 2.2(\text{stat}) \pm 2.0(\text{sys}) \text{ fb}$, and the correlation between $WZ + \text{charm}$ and $WZ + b$ is -0.26.

⁴²⁵ 8.1 1-jet Fit Results

⁴²⁶ **The results of the fit are currently blinded.**

⁴²⁷ The pre-fit yields in each of the regions used in the fit are shown in Table ??, and summarized in
⁴²⁸ Figure ??.

⁴²⁹

Sample	1j, <85% WP	1j, 77%-85% WP	1j, 70%-77% WP	1j, 60%-70% WP	1j, >60% WP	tZ CR
$WZ + b$	11 ± 2	4.7 ± 0.5	4.6 ± 0.4	5.1 ± 0.4	24 ± 2	6.0 ± 0.5
$WZ + c$	318 ± 22	81 ± 6	43.1 ± 3.6	25.8 ± 2.6	9.4 ± 1.8	2.9 ± 0.6
$WZ + l$	4020 ± 250	91 ± 13	17 ± 3	4.9 ± 1.6	1.3 ± 0.4	0.2 ± 0.1
Other VV	6.2 ± 0.6	0.2 ± 0.4	0.2 ± 0.04	0.07 ± 0.1	0.1 ± 0.1	0.1 ± 0.2
ZZ	336 ± 26	17.8 ± 2.1	4.3 ± 0.6	1.7 ± 0.5	0.36 ± 0.08	0.10 ± 0.03
$t\bar{t}W$	1.1 ± 0.2	0.2 ± 0.1	0.3 ± 0.1	0.4 ± 0.1	1.5 ± 0.3	0.7 ± 0.2
$t\bar{t}Z$	6.8 ± 1.2	1.4 ± 0.3	1.0 ± 0.2	1.2 ± 0.2	4.4 ± 0.8	3.2 ± 0.6
$Z + \text{jets}$	169 ± 38	8.9 ± 1.9	3.7 ± 0.8	3.3 ± 0.7	3.2 ± 0.7	0.8 ± 0.17
$V + \gamma$	45 ± 28	1.9 ± 2.4	0.1 ± 0.1	0.02 ± 0.01	1.0 ± 0.9	0.02 ± 0.03
tZ	24.3 ± 4.3	5.5 ± 1.1	4.1 ± 0.8	5.9 ± 1.1	10.7 ± 2.0	23 ± 4
tW	1.4 ± 0.8	0.2 ± 0.5	0.0 ± 0.2	0.7 ± 0.6	0.26 ± 0.42	0.39 ± 0.41
WtZ	2.3 ± 1.2	0.6 ± 0.3	0.3 ± 0.21	0.27 ± 0.2	1.1 ± 0.7	0.6 ± 0.5
VVV	12.4 ± 0.5	0.93 ± 0.06	0.35 ± 0.03	0.13 ± 0.02	0.14 ± 0.03	0.02 ± 0.01
VH	40 ± 6	2.6 ± 1.4	0.9 ± 0.8	0.7 ± 0.8	0.5 ± 0.6	0.0 ± 0.0
$t\bar{t}$	12.1 ± 1.6	2.9 ± 0.6	2.5 ± 0.5	2.8 ± 0.5	11.2 ± 1.4	10.9 ± 1.5
$t\bar{t}H$	0.24 ± 0.03	0.05 ± 0.01	0.04 ± 0.01	0.06 ± 0.01	0.20 ± 0.03	0.13 ± 0.02
Total	5010 ± 260	227 ± 24	88 ± 12	57 ± 8	76 ± 16	53 ± 8

Table 13: Pre-fit yields in each of the 1-jet fit regions.

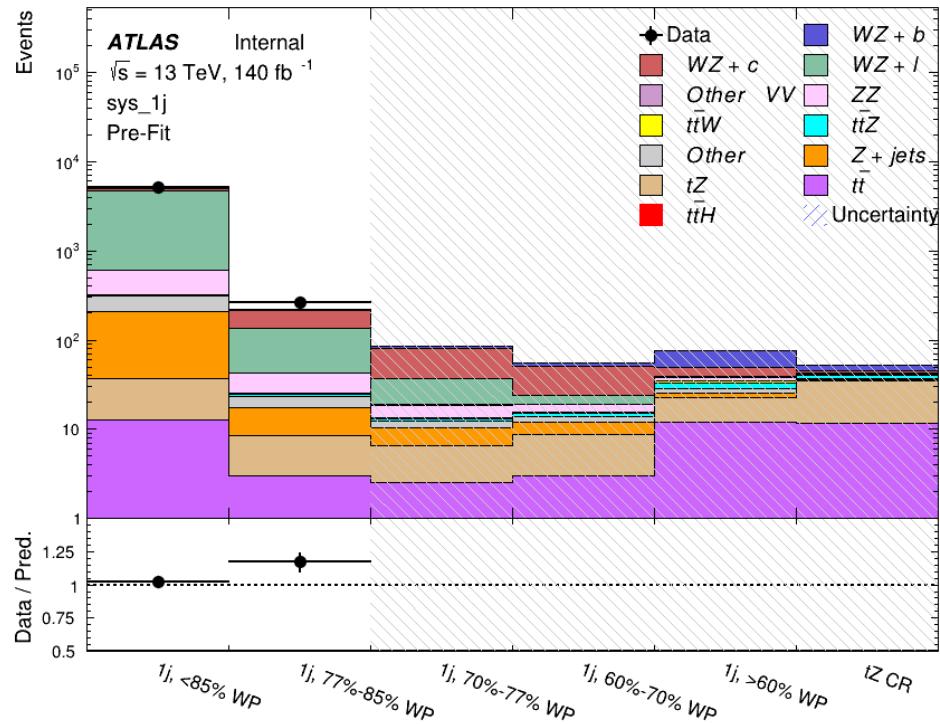


Figure 29: Pre-fit summary of the 1-jet fit regions.

430 The post-fit yields in each region are summarized in Figure ??.

431

	1j, <85% WP	1j, 77%-85% WP	1j, 70%-77% WP	1j, 60%-70% WP	1j, >60% WP	tZ CR
WZ + b	11 ± 5	4.7 ± 2.0	4.6 ± 2.0	5.1 ± 2.1	24 ± 10	6.0 ± 2.50
WZ + c	320 ± 60	80 ± 14	43 ± 7	26 ± 5	9.4 ± 2.3	2.9 ± 0.7
WZ + l	4020 ± 130	90 ± 11	17.3 ± 2.8	4.9 ± 1.6	1.3 ± 0.4	0.23 ± 0.13
Other VV	6.2 ± 0.6	0.92 ± 0.07	0.02 ± 0.01	0.01 ± 0.01	0.01 ± 0.01	0.01 ± 0.01
ZZ	346 ± 57	19 ± 5	4.3 ± 0.8	2.7 ± 0.5	2.4 ± 0.1	2.1 ± 0.6
t̄tW	1.09 ± 0.21	0.2 ± 0.1	0.1 ± 0.1	0.4 ± 0.1	1.5 ± 0.3	0.1 ± 0.2
t̄tZ	6.8 ± 1.2	1.4 ± 0.3	1.0 ± 0.2	1.2 ± 0.2	4.4 ± 0.7	3.2 ± 0.5
rare Top	0.14 ± 0.04	0.04 ± 0.02	0.04 ± 0.0	0.1 ± 0.03	0.14 ± 0.04	0.15 ± 0.05
t̄tWW	0.04 ± 0.03	0.01 ± 0.02	0.01 ± 0.01	0.01 ± 0.01	0.01 ± 0.02	0.01 ± 0.01
Z + jets	169 ± 37	8.9 ± 1.9	3.7 ± 0.8	3.3 ± 0.7	3.2 ± 0.7	0.8 ± 0.2
W + jets	0.01 ± 0.01					
V + γ	46 ± 28	1.9 ± 2.4	0.1 ± 0.1	0.0 ± 0.2	1.0 ± 0.9	0.0 ± 0.0
tZ	24 ± 4	5.5 ± 1.0	4.1 ± 0.8	5.9 ± 1.1	10.7 ± 1.8	23.3 ± 3.7
tW	1.37 ± 0.82	0.18 ± 0.26	0.01 ± 0.12	0.67 ± 0.64	0.26 ± 0.42	0.39 ± 0.41
WtZ	2.3 ± 1.2	0.6 ± 0.3	0.3 ± 0.2	0.3 ± 0.2	1.1 ± 0.6	0.6 ± 0.3
VVV	12.4 ± 0.4	0.9 ± 0.1	0.4 ± 0.1	0.13 ± 0.02	0.14 ± 0.03	0.02 ± 0.01
VH	40 ± 6	2.6 ± 1.4	0.9 ± 0.8	0.7 ± 0.8	0.4 ± 0.6	0.01 ± 0.01
t̄t	12.1 ± 1.6	2.9 ± 0.6	2.5 ± 0.5	2.8 ± 0.5	11.2 ± 1.5	10.9 ± 1.4
t̄tH	0.24 ± 0.03	0.05 ± 0.01	0.04 ± 0.01	0.06 ± 0.01	0.20 ± 0.03	0.13 ± 0.02
Total	5100 ± 110	227 ± 12	87 ± 6	56.7 ± 4.4	76 ± 9	52.5 ± 4.2

Table 14: Post-fit yields in each of the 1-jet fit regions.

⁴³² A post-fit summary plot of the 1-jet fitted regions is shown in Figure ??:

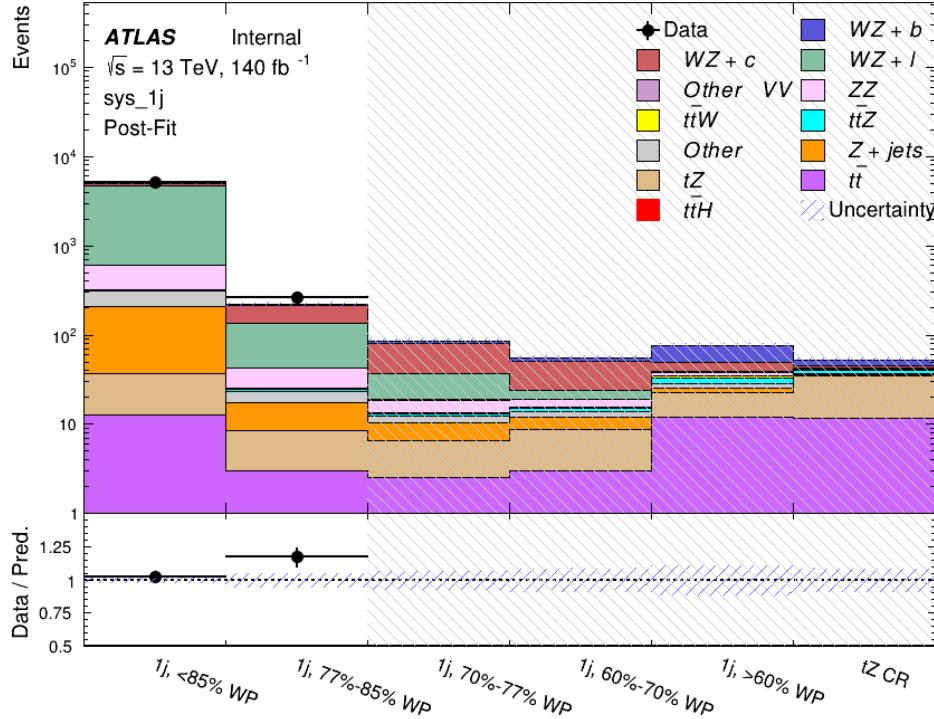


Figure 30: Post-fit summary of the 1-jet fit regions.

433 As described in Section ??, there are 226 systematic uncertainties that are considered as NPs in
 434 the fit. These NPs are constrained by Gaussian or log-normal probability density functions. The
 435 latter are used for normalisation factors to ensure that they are always positive. The expected
 436 number of signal and background events are functions of the likelihood. The prior for each NP is
 437 added as a penalty term, decreasing the likelihood as it is shifted away from its nominal value.

438 The impact of each NP is calculated by performing the fit with the parameter of interest held
 439 fixed, varied from its fitted value by its uncertainty, and calculating $\Delta\mu$ relative to the baseline
 440 fit. The impact of the most significant sources of systematic uncertainties is summarized in Table
 441 ??.

Uncertainty Source	$\Delta\mu$	
WZ + light cross-section	0.13	-0.12
WZ + charm cross-section	-0.10	0.12
Jet Energy Scale	0.08	0.13
tZ cross-section	-0.10	0.10
Jet Energy Resolution	-0.10	0.10
Luminosity	-0.08	0.09
Other Diboson + b cross-section	-0.07	0.07
Flavor tagging	0.05	0.05
t <bar>t} cross-section</bar>	-0.05	0.05
WZ cross-section - QCD scale	-0.04	0.03
Total Systematic Uncertainty	0.28	0.32

Table 15: Summary of the most significant sources of systematic uncertainty on the measurement of WZ + b with exactly one associated jet.

⁴⁴² The ranking and impact of those nuisance parameters with the largest contribution to the overall uncertainty is shown in Figure ??.

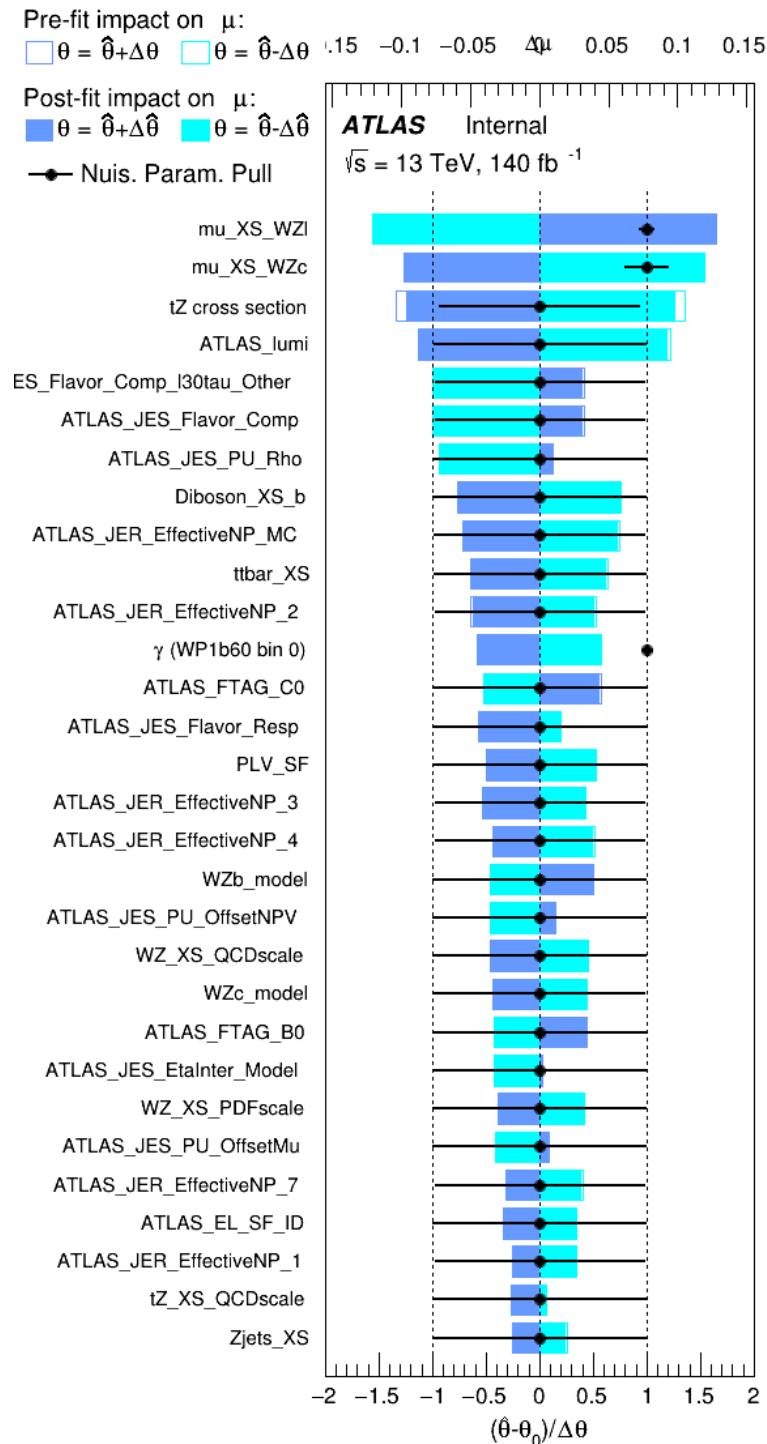


Figure 31: Impact of systematic uncertainties on the signal-strength of $WZ + b$ for events with exactly one jet

444 The large impact of the Jet Energy Scale and Jet Flavor Tagging is unsurprising, as the shape of
 445 the fit regions depends heavily on the modeling of the jets. The other major sources of uncertainty
 446 come from background modelling and cross-section uncertainty. The pie charts in Figure ??
 447 show that for the modelling uncertainties that contribute most correspond to the most significant
 448 backgrounds.

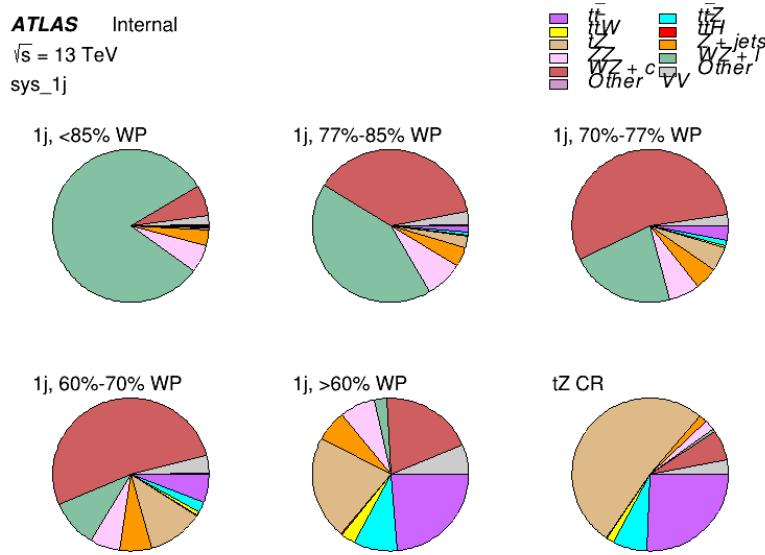


Figure 32: Post-fit background composition of the fit regions.

449 The correlations between these nuisance parameters are summarized in Figure ??.

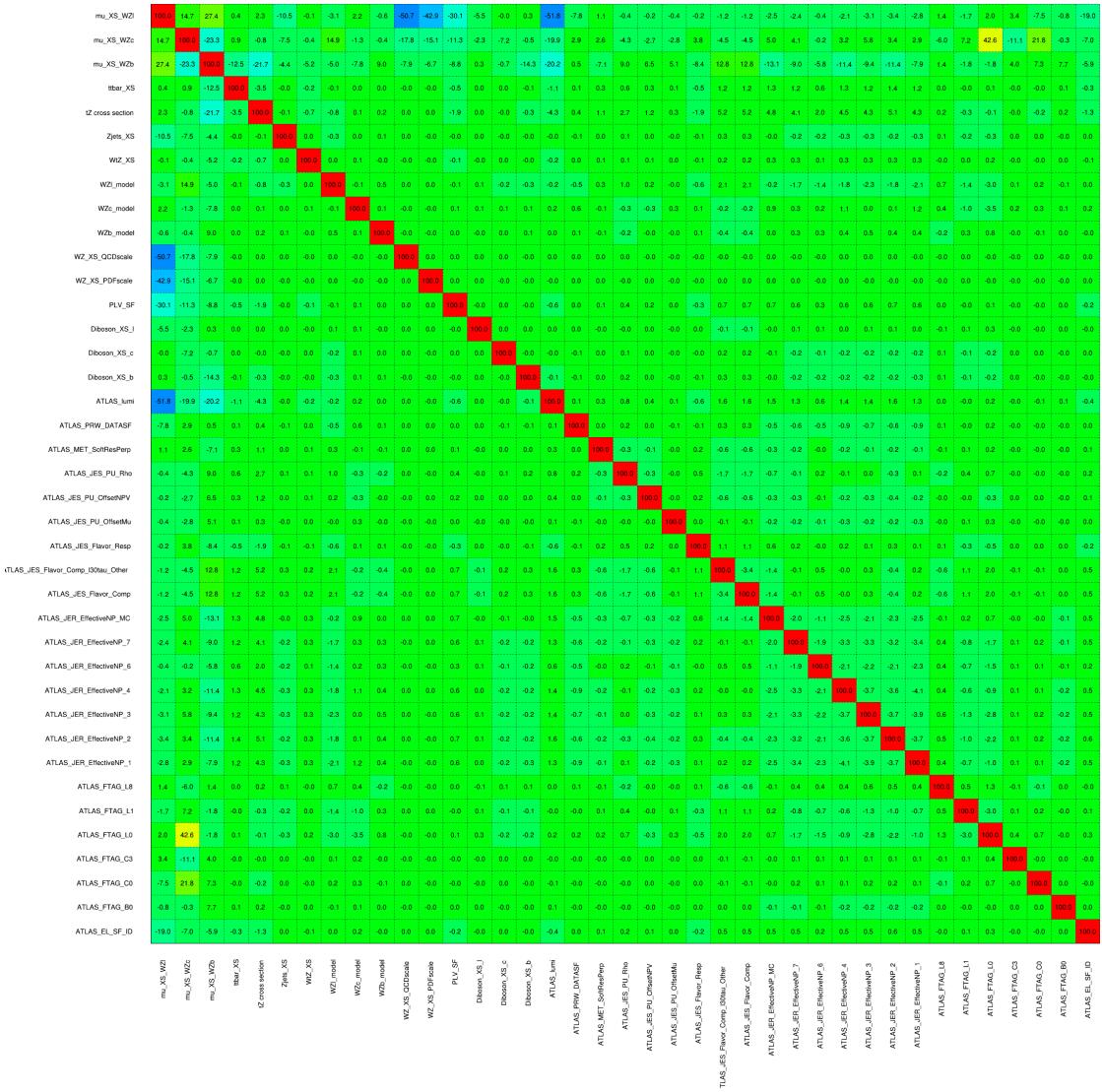


Figure 33: Correlations between nuisance parameters

450 The negative correlations between $\mu_{WZ+charm}$ and μ_{WZ+b} and $\mu_{WZ+light}$ are expected: WZ +
 451 charm is present in both the WZ + b and WZ + light enriched regions, therefore increasing the
 452 fraction of charm requires increasing the fraction of WZ + b and WZ + light. This reasoning
 453 also explains the positive correlation between μ_{WZ+b} and $\mu_{WZ+light}$.

454 Two of the major backgrounds in the region with the highest purity of WZ + b are tZ and Other
 455 VV + b, explaining the negative correlations between μ_{WZ+b} and the tZ cross section, and the
 456 VV + b cross section.

457 The high correlation between the luminosity and $\mu_{WZ+light}$ arises from the fact that the uncer-
 458 tainty on $\mu_{WZ+light}$ is very low (around 4%). Small changes in luminosity cause a change in
 459 the yield of $WZ + light$ that is large compared to its uncertainty, producing a large correlation
 460 between these two parameters.

461 8.2 2-jet Fit Results

462 **The results of the fit are currently blinded.**

463 Pre-fit yields in each of the 2-jet fit regions are shown in Figure ??.

	2j, <85% WP	2j, 77%-85% WP	2j, 70%-77% WP	2j, 60%-70% WP	2j, >60% WP	tZ CR 2j
WZ + b	13 ± 2	6.7 ± 0.5	5.6 ± 0.4	8.0 ± 0.6	31 ± 2	14 ± 1
WZ + c	260 ± 20	77 ± 6	41 ± 3	27 ± 3	11 ± 2	4.8 ± 0.6
WZ + l	1860 ± 150	90 ± 14	18 ± 3	5.8 ± 1.4	1.4 ± 0.4	0.25 ± 0.15
Other VV	7.63 ± 0.63	0.6 ± 0.5	0.16 ± 0.03	0.01 ± 0.01	0.1 ± 0.1	0.1 ± 0.1
ZZ	135 ± 20	14.1 ± 3.2	4.7 ± 0.8	4.0 ± 0.6	4.1 ± 0.7	3.1 ± 0.5
t <bar>t>W</bar>	0.8 ± 0.2	0.4 ± 0.1	0.5 ± 0.1	0.7 ± 0.2	4.3 ± 0.6	3.9 ± 0.6
t <bar>t>Z</bar>	14.7 ± 2.2	5.6 ± 0.8	4.5 ± 0.7	6.5 ± 1.1	25.4 ± 4.0	21.9 ± 3.4
rare Top	0.14 ± 0.04	0.07 ± 0.03	0.03 ± 0.02	0.09 ± 0.03	0.37 ± 0.07	0.6 ± 0.1
t <bar>t>WW</bar>	0.04 ± 0.03	0.02 ± 0.02	0.01 ± 0.01	0.01 ± 0.00	0.03 ± 0.03	0.01 ± 0.01
Z + jets	110.0 ± 22.9	9.6 ± 2.0	2.1 ± 0.50	1.6 ± 0.4	5.1 ± 1.1	1.5 ± 0.3
W + jets	0.0 ± 0.0	0.0 ± 0.0	0.0 ± 0.0	0.0 ± 0.0	0.0 ± 0.0	0.0 ± 0.0
V + γ	25 ± 18	0.5 ± 0.2	0.1 ± 0.1	0.1 ± 0.1	0.0 ± 0.02	0.05 ± 0.07
tZ	15.9 ± 2.9	6.9 ± 1.3	5.1 ± 1.0	8.0 ± 1.5	18.7 ± 3.2	36.4 ± 6.1
tW	0.9 ± 0.7	0.2 ± 0.3	0.0 ± 0.1	0.0 ± 0.0	0.8 ± 0.6	0.2 ± 0.2
WtZ	4.9 ± 2.5	1.5 ± 0.8	1.1 ± 0.6	1.3 ± 0.7	4.6 ± 2.4	3.3 ± 1.7
VVV	7.4 ± 0.3	1.0 ± 0.1	0.4 ± 0.1	0.2 ± 0.1	0.13 ± 0.03	0.04 ± 0.01
VH	19.5 ± 4.2	2.8 ± 1.6	0.7 ± 0.7	0.1 ± 0.2	0.0 ± 0.0	0.0 ± 0.0
t <bar>t></bar>	0.7 ± 0.4	0.1 ± 0.1	0.05 ± 0.06	0.15 ± 0.13	0.8 ± 0.5	2.3 ± 1.2
t <bar>t></bar>	6.8 ± 1.0	2.4 ± 0.5	1.8 ± 0.4	3.3 ± 0.6	8.4 ± 1.2	13.6 ± 1.7
t <bar>t>H</bar>	0.4 ± 0.1	0.2 ± 0.1	0.16 ± 0.02	0.23 ± 0.03	0.94 ± 0.11	1.03 ± 0.12
Total	2580 ± 160	229 ± 24	89 ± 13	69 ± 11	120 ± 15	108 ± 11

Table 16: Pre-fit yields in each of the 2-jet fit regions.

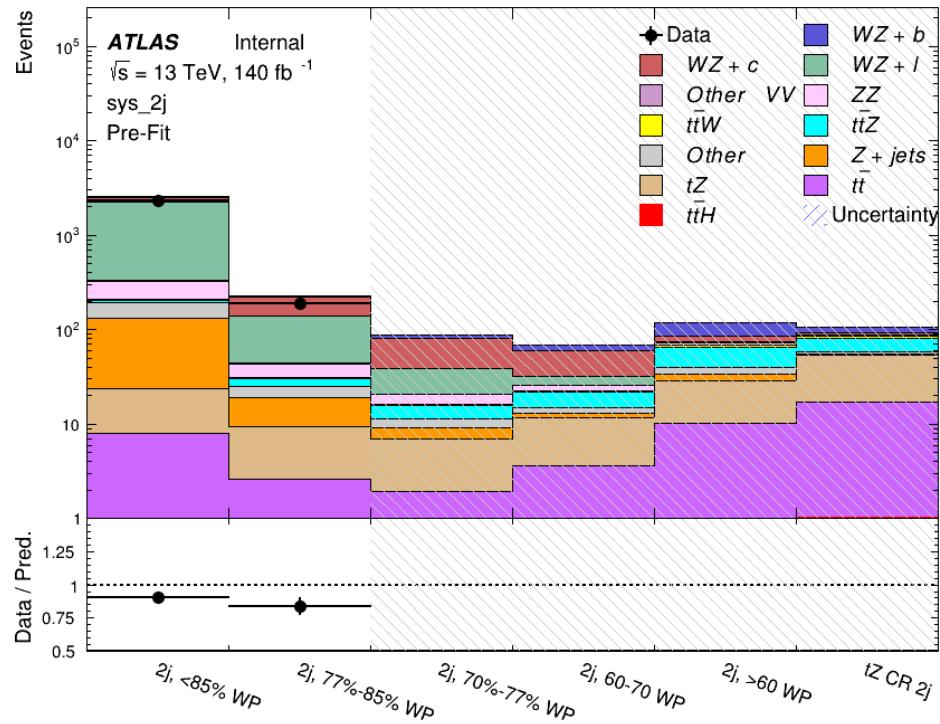


Figure 34: Pre-fit summary of the 2-jet fit regions.

⁴⁶⁴ The post-fit yields in each region are summarized in Figure ??.

	2j, <85% WP	2j, 77%-85% WP	2j, 70%-77% WP	2j, 60%-70% WP	2j, >60% WP	tZ CR 2j
WZ + b	13 ± 6	6.7 ± 2.9	5.8 ± 2.5	8.0 ± 3.5	31 ± 13	14 ± 5
WZ + c	260 ± 60	77 ± 15	41 ± 8	26 ± 5	10.9 ± 2.4	4.8 ± 1.1
WZ + l	1860 ± 90	90 ± 12	17.6 ± 2.8	5.8 ± 1.3	1.4 ± 0.4	0.3 ± 0.2
Other VV	7.6 ± 0.6	0.3 ± 0.3	0.3 ± 0.1	0.1 ± 0.06	0.03 ± 0.02	0.1 ± 0.1
ZZ	145 ± 30	11.3 ± 4.4	2.7 ± 1.6	1.0 ± 0.3	4.0 ± 0.1	2.4 ± 0.1
t̄tW	0.8 ± 0.2	0.4 ± 0.1	0.54 ± 0.12	0.74 ± 0.15	4.3 ± 0.6	3.9 ± 0.6
t̄tZ	14.7 ± 2.2	5.6 ± 0.8	4.5 ± 0.7	6.5 ± 1.0	25.4 ± 3.9	21.9 ± 3.3
rare Top	0.14 ± 0.04	0.07 ± 0.03	0.03 ± 0.02	0.09 ± 0.03	0.4 ± 0.1	0.6 ± 0.1
t̄tWW	0.04 ± 0.03	0.02 ± 0.02	0.01 ± 0.01	0.01 ± 0.01	0.03 ± 0.03	0.01 ± 0.01
Z + jets	110 ± 23	9.6 ± 2.0	2.1 ± 0.5	1.6 ± 0.4	5.1 ± 1.1	1.5 ± 0.3
W + jets	0.0 ± 0.0					
V + γ	25 ± 19	0.5 ± 0.2	0.1 ± 0.1	0.13 ± 0.14	0.0 ± 0.02	0.05 ± 0.07
tZ	15.9 ± 2.7	6.9 ± 1.2	5.1 ± 0.9	8.0 ± 1.4	18.7 ± 3.0	36 ± 6
tW	0.1 ± 0.7	0.2 ± 0.3	0.0 ± 0.1	0.0 ± 0.0	0.8 ± 0.6	0.2 ± 0.2
WtZ	4.9 ± 2.5	1.5 ± 0.8	1.1 ± 0.6	1.3 ± 0.7	4.6 ± 2.3	3.3 ± 1.7
VVV	7.4 ± 0.3	1.0 ± 0.1	0.36 ± 0.03	0.19 ± 0.03	0.13 ± 0.03	0.04 ± 0.01
VH	19 ± 4	2.8 ± 1.6	0.7 ± 0.7	0.1 ± 0.2	0.0 ± 0.0	0.0 ± 0.0
t̄t	6.8 ± 1.0	2.4 ± 0.5	1.8 ± 0.4	3.3 ± 0.6	8.4 ± 1.2	13.6 ± 1.7
t̄tH	0.40 ± 0.05	0.19 ± 0.03	0.16 ± 0.02	0.23 ± 0.03	0.94 ± 0.11	1.03 ± 0.11
Total	2580 ± 60	229 ± 11	89 ± 6	69.1 ± 4.1	120 ± 10	108 ± 6

Table 17: Post-fit yields in each of the 2-jet fit regions.

⁴⁶⁵ A post-fit summary of the fitted regions is shown in Figure ??:

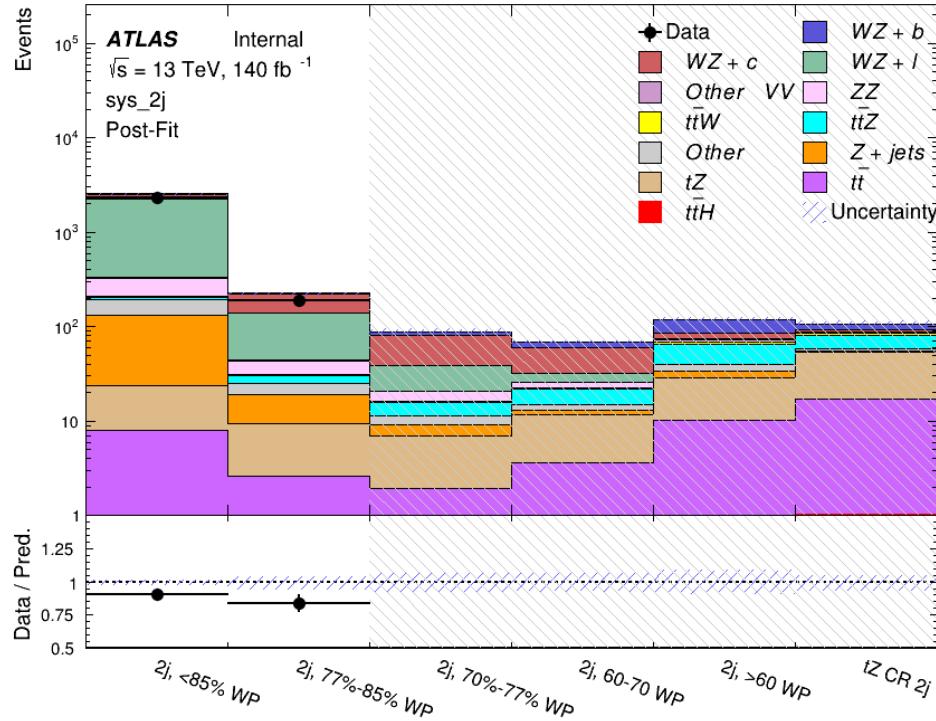


Figure 35: Post-fit summary of the fit over 2-jet regions.

466 The same set of systematic uncertainties consider for the 1-jet fit are included in the 2-jet fit as
 467 well. The impact of the most significant systematic uncertainties is summarized in Table ??.

Uncertainty Source	$\Delta\mu$	
WZ + c cross-section	-0.1	0.14
Luminosity	-0.11	0.12
tZ cross-section	-0.11	0.11
Jet Energy Scale	-0.11	0.11
ttZ cross-section - QCD scale	-0.08	0.09
WZ + l cross-section	0.08	-0.07
WtZ cross-section	-0.07	0.07
Flavor tagging	0.05	0.05
Other VV + b cross-section	-0.05	0.05
Jet Energy Resolution	-0.04	0.04
Total	0.29	0.31

Table 18: Summary of the most significant sources of systematic uncertainty on the measurement of WZ + b 2-jet events.

⁴⁶⁸ The ranking and impact of those nuisance parameters with the largest contribution to the overall
⁴⁶⁹ uncertainty is shown in Figure ??.

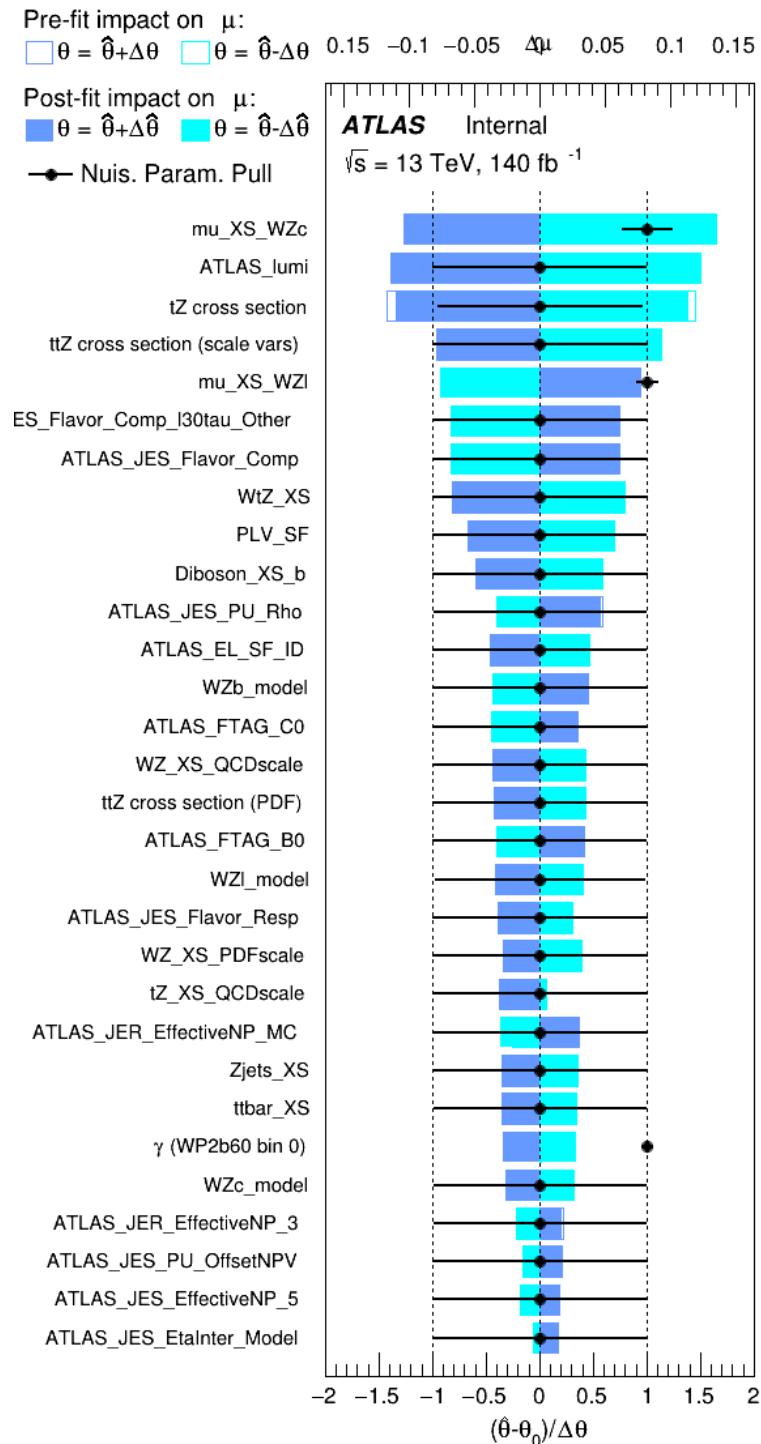


Figure 36: Impact of systematic uncertainties on the signal-strength of WZ + b in 2-jet events.

470 The large impact of the Jet Energy Scale and Jet Flavor Tagging is unsurprising, as the shape of
471 the fit regions depends heavily on the modeling of the jets. The other major sources of uncertainty
472 come from background modelling and cross-section uncertainty. The pie charts in Figure ??
473 show that for the modelling uncertainties that contribute most correspond to the most significant
474 backgrounds.

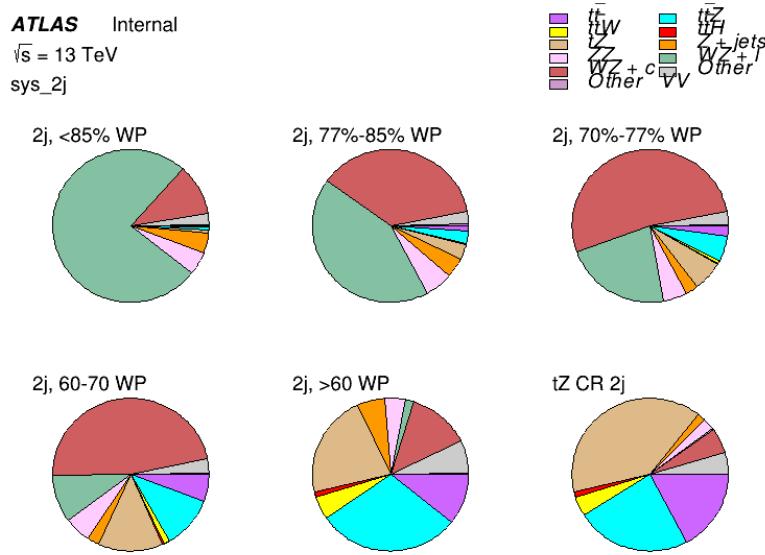


Figure 37: Post-fit background composition of the 2-jet fit regions.

475 The correlations between these nuisance parameters are summarized in Figure ??.



Figure 38: Correlations between nuisance parameters in the 2-jet fit

476 As in the 1-jet case, no significant, unexpected correlations are found between nuisance para-
 477 meters.

9 Conclusion

479 A measurement of $WZ + \text{heavy flavor}$ is performed using 140 fb^{-1} of $\sqrt{s} = 13 \text{ TeV}$ proton-
 480 proton collision data collected by the ATLAS detector at the LHC. The expected cross-section

481 of WZ+b with 1-jet is $1.74^{+0.65}_{-0.61}(\text{stat})^{+0.42}_{-0.37}(\text{sys}) \text{ fb}$, and $14.6 \pm 2.5(\text{stat}) \pm 2.1(\text{sys}) \text{ fb}$ for WZ
482 + charm, with a correlation of -0.23 between them. An expected significance of 2.2 is observed
483 for WZ + b in this region.

484 For the 2-jet regions, an expected significance of 2.6 is observed for WZ + b, with an expected
485 cross-section of $2.46^{+0.83}_{-0.81}(\text{stat})^{+0.073}_{-0.68}(\text{sys}) \text{ fb}$. For WZ + charm, a cross-section of $12.7 \pm$
486 $2.2(\text{stat}) \pm 2.0(\text{sys}) \text{ fb}$ is expected for 2-jet events. A correlation of -0.26 is observed for WZ+b
487 and WZ + charm.

488 **This section will be include final results once unblinded.**

489 **Appendices**

490 **tZ Interference Studies**

491 Because it includes an on-shell Z boson as well as a b-jet and W from the top decay, tZ
492 production represents an identical final state to WZ + b-jet. This implies the possibility of matrix
493 level interference between these two processes not accounted for in the Monte Carlo simulations,
494 which consider the two processes independently. Truth level studies are performed in order to
495 estimate the impact of these interference effects.

496 In order to estimate the matrix level interference effects between tZ and WZ + b-jet, two different
497 sets of simulations are produced using MadGraph 5 [**Madgraph**] - one which simulates these
498 two processes independently, and another where they are produced simultaneously, such that
499 interference effects are present. These two sets of samples are then compared, and the difference
500 between them can be taken to represent any interference effects.

501 MadGraph simulations of 10,000 tZ and 10,000 WZ + b-jet events are produced, along with
502 20,000 events where both are present, in the fiducial region where three leptons and at least one
503 jet are produced.

504 A selection mimicking the preselection used in the main analysis is applied to the samples: The
505 SS leptons are required to have $p_T > 20$ GeV, and > 10 GeV is required for the OS lepton. The
506 associated b-jet is required to have $p_T > 25$ GeV, and all physics objects are required to fall in a
507 range of $|\eta| < 2.5$.

508 The kinematics of these samples after the selection has been applied are shown below:

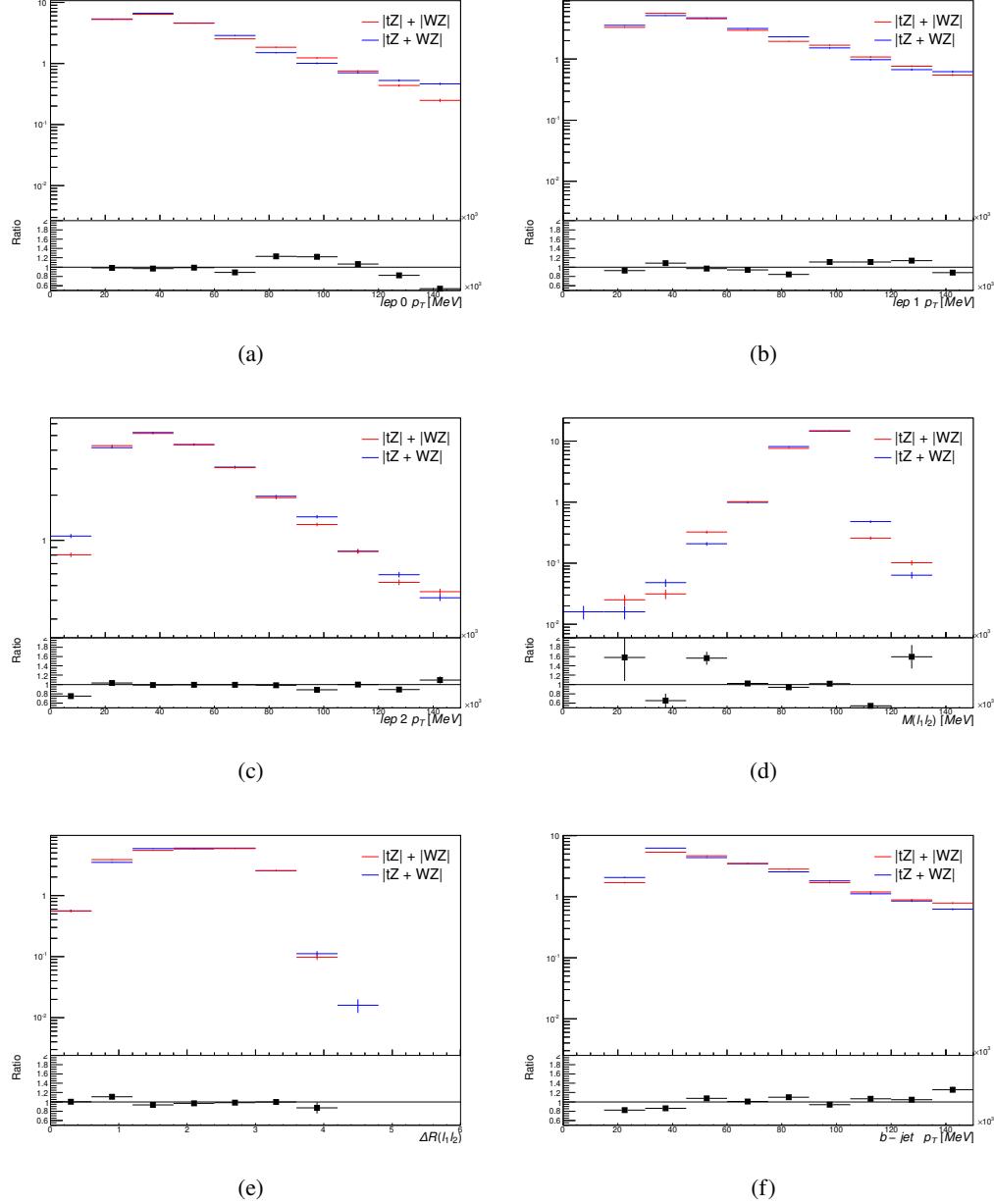


Figure 39: Comparisons between (a) the p_T of the lepton from the W, (b) and (c) show the p_T of the lepton from the Z, (d) the invariant mass of the Z-candidate, (e) ΔR of the leptons from the Z, and (f) the p_T of the b-jet, for WZ and tZ events generated with interference effects (blue) and without interference effects (red).

509 The overall cross-section of the two methods agree within error, and no significant differences
 510 in the kinematic distributions are seen. It is therefore concluded that interference effects do not

⁵¹¹ significantly impact the results.

512 **Alternate tZ Inclusive Fit**

- 513 While tZ is often considered as a distinct process from WZ + b, this could also be considered part
 514 of the signal. Alternate studies are performed where, using the same framework as the nominal
 515 analysis, a measurement of WZ + b is performed that includes tZ as part of WZ+b.
- 516 Because of this change, the tZ CR is no longer necessary, and only the five pseudo-continuous
 517 b-tag regions are used in the fit. Further, systematics related to the tZ cross-section are removed
 518 from the fit, as they are now encompassed by the normalization measurement of WZ + b. All
 519 other systematic uncertainties are carried over from the nominal analysis.
- 520 A post-fit summary of the 1-jet regions used in the fit are shown in Figure ??.

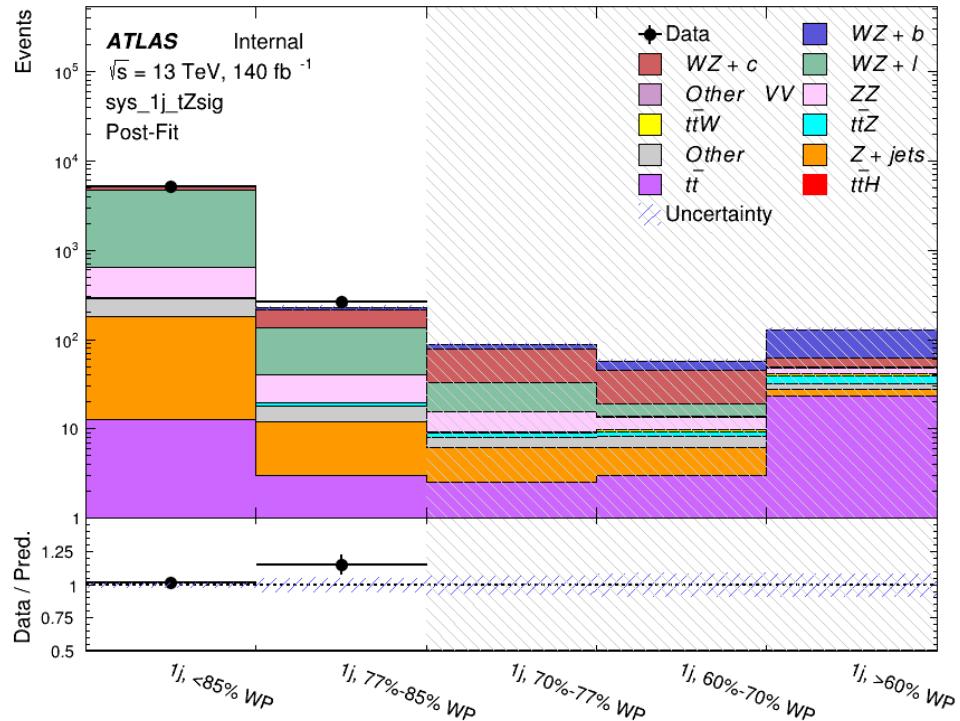


Figure 40: Post-fit summary of the 1-jet fit regions.

- 521 An expected WZ + b cross-section of $4.1^{+0.78}_{-0.74}(\text{stat})^{+0.53}_{-0.52}(\text{sys}) \text{ fb}$ is extracted from the fit, with
 522 an expected significance of 4.0σ .
- 523 The impact of the predominate systematics are summarized in Table ??.

Uncertainty Source	$\Delta\mu$	
WZ + light cross-section	0.08	-0.08
Jet Energy Scale	-0.06	0.08
Luminosity	-0.05	0.06
WZ + charm cross-section	-0.04	0.05
Other Diboson + b cross-section	-0.04	0.04
WZ cross-section - QCD scale	-0.04	0.03
t <bar>t</bar>	-0.03	0.03
Jet Energy Resolution	-0.03	0.03
Flavor tagging	-0.03	0.03
Z+jets cross section	-0.02	0.02
Total Systematic Uncertainty	-0.15	0.16

Table 19: Summary of the most significant sources of systematic uncertainty on the measurement of WZ + b with exactly one associated jet.

524 A post-fit summary of the 2-jet regions used in the fit are shown in Figure ??.

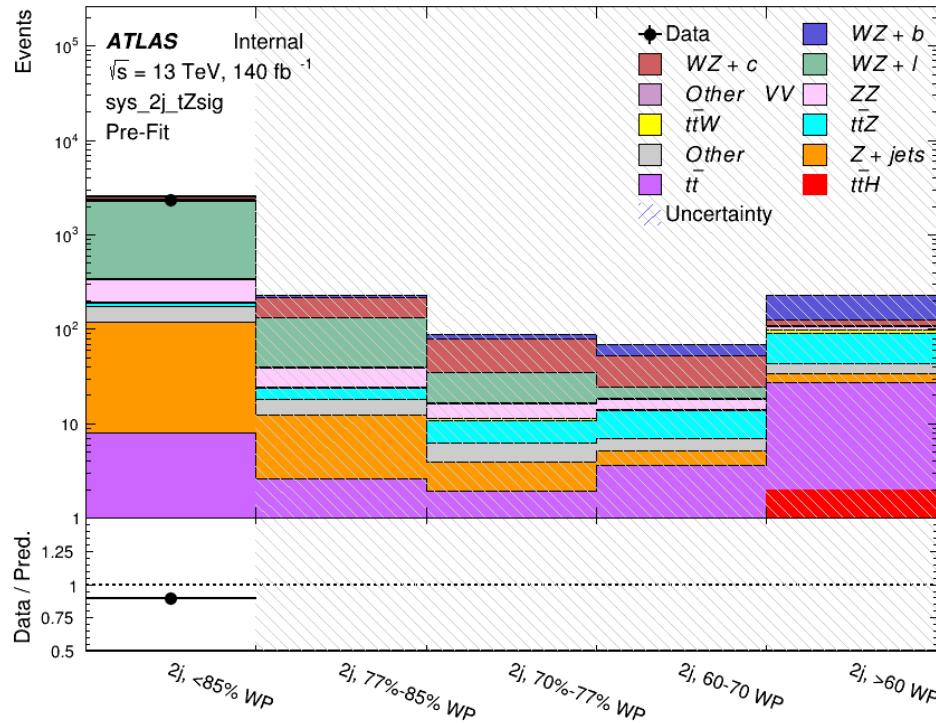


Figure 41: Post-fit summary of the 2-jet fit regions.

525 An expected WZ + b cross-section of $5.9^{+0.9}_{-0.9}(\text{stat})^{+0.9}_{-0.8}(\text{sys}) \text{ fb}$ is extracted from the fit, with an

⁵²⁶ expected significance of 5.3σ .

⁵²⁷ The impact of the predominate systematics are summarized in Table ??.

Uncertainty Source	$\Delta\mu$	
Luminosity	-0.07	0.07
Jet Energy Scale	0.06	0.05
ttZ cross-section - QCD scale	-0.05	0.05
WZ+l cross-section	0.05	-0.05
WZ+c cross-section	-0.03	0.05
WtZ cross-section	-0.03	0.03
WZ cross-section QCDscale	-0.03	0.03
Diboson cross-section b	-0.03	0.03
WZ cross-section - PDF	-0.03	0.03
Flavor Tagging	0.03	0.02
Total Systematic Uncertainty	-0.14	0.16

Table 20: Summary of the most significant sources of systematic uncertainty on the measurement of WZ + b with exactly one associated jet.

528 **DSID list**

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 532 data17_13TeV.AllYear.physics_Main.PhysCont.DAOD_HIGG8D1.grp17_v01_p4134
 533 data18_13TeV.AllYear.physics_Main.PhysCont.DAOD_HIGG8D1.grp18_v01_p4134
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 741 mc16_13TeV.346343.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_allhad.deriv.DAOD_HIGG8D1.e7148_e5984_s3126_r9364_r9315_p4133
 742 mc16_13TeV.346344.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_semilep.deriv.DAOD_HIGG8D1.e7148_e5984_s3126_r9364_r9315_p4133
 743 mc16_13TeV.346345.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_dilep.deriv.DAOD_HIGG8D1.e7148_e5984_s3126_r9364_r9315_p4133
 744
 745 mc16d:
 746 mc16_13TeV.364253.Sherpa_222_NNPDF30NNLO_lllv.deriv.DAOD_HIGG8D1.e5916_e5984_s3126_r10201_r10210_p4133
 747 mc16_13TeV.364250.Sherpa_222_NNPDF30NNLO_llll.deriv.DAOD_HIGG8D1.e5894_e5984_s3126_r10201_r10210_p4133
 748 mc16_13TeV.364254.Sherpa_222_NNPDF30NNLO_llvv.deriv.DAOD_HIGG8D1.e5916_e5984_s3126_r10201_r10210_p4133
 749 mc16_13TeV.364255.Sherpa_222_NNPDF30NNLO_lvvv.deriv.DAOD_HIGG8D1.e5916_e5984_s3126_r10201_r10210_p4133
 750 mc16_13TeV.410155.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttW.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10201_r10210_p4133
 751 mc16_13TeV.410156.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttZnunu.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10201_r10210_p4133
 752 mc16_13TeV.410157.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttZqq.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10201_r10210_p4133
 753 mc16_13TeV.410218.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_tteee.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10201_r10210_p4133
 754 mc16_13TeV.410219.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttmmumu.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10201_r10210_p4133
 755 mc16_13TeV.410220.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_tttautau.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10201_r10210_p4133
 756 mc16_13TeV.410276.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttee_mll_1_5.deriv.DAOD_HIGG8D1.e6087_e5984_s3126_r10201_r10210_p4133
 757 mc16_13TeV.410277.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttmmumu_mll_1_5.deriv.DAOD_HIGG8D1.e6087_e5984_s3126_r10201_r10210_p4133
 758 mc16_13TeV.410278.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_tttautau_mll_1_5.deriv.DAOD_HIGG8D1.e6087_e5984_s3126_r10201_r10210_p4133
 759 mc16_13TeV.410397.MadGraphPythia8EvtGen_ttbar_wbee_MEN30LO_A14N23LO.deriv.DAOD_HIGG8D1.e6086_e5984_s3126_r10201_r10210_p4133
 760 mc16_13TeV.410398.MadGraphPythia8EvtGen_ttbar_wbmumu_MEN30LO_A14N23LO.deriv.DAOD_HIGG8D1.e6086_e5984_s3126_r10201_r10210_p4133
 761 mc16_13TeV.410399.MadGraphPythia8EvtGen_ttbar_wbtautau_MEN30LO_A14N23LO.deriv.DAOD_HIGG8D1.e6086_e5984_s3126_r10201_r10210_p4133
 762 mc16_13TeV.410658.PhPy8EG_A14_tchan_BW50_lept_top.deriv.DAOD_HIGG8D1.e6671_e5984_s3126_r10201_r10210_p4133
 763 mc16_13TeV.410659.PhPy8EG_A14_tchan_BW50_lept_antitop.deriv.DAOD_HIGG8D1.e6671_e5984_s3126_r10201_r10210_p4133
 764 mc16_13TeV.410644.PowhegPythia8EvtGen_A14_singletop_schan_lept_top.deriv.DAOD_HIGG8D1.e6527_e5984_s3126_r10201_r10210_p4133
 765 mc16_13TeV.410645.PowhegPythia8EvtGen_A14_singletop_schan_lept_antitop.deriv.DAOD_HIGG8D1.e6527_s3126_r10201_r10210_p4133
 766 mc16_13TeV.304014.MadGraphPythia8EvtGen_A14NNPDF23_3top_SM.deriv.DAOD_HIGG8D1.e4324_e5984_s3126_r10201_r10210_p4133
 767 mc16_13TeV.410080.MadGraphPythia8EvtGen_A14NNPDF23_4topSM.deriv.DAOD_HIGG8D1.e4111_e5984_s3126_r10201_r10210_p4133
 768 mc16_13TeV.410081.MadGraphPythia8EvtGen_A14NNPDF23_ttbar_WW.deriv.DAOD_HIGG8D1.e4111_e5984_s3126_r10201_r10210_p4133
 769 mc16_13TeV.364100.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV0_70_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_s3126_r10201_r10210_p4133
 770 mc16_13TeV.364101.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV0_70_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_s3126_r10201_r10210_p4133
 771 mc16_13TeV.364101.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV0_70_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10201_r10210_p4133
 772 mc16_13TeV.364102.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV0_70_BFilter.deriv.DAOD_HIGG8D1.e5271_s3126_r10201_r10210_p4133
 773 mc16_13TeV.364103.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV70_140_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_s3126_r10201_r10210_p4133
 774 mc16_13TeV.364104.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV70_140_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_s3126_r10201_r10210_p4133
 775 mc16_13TeV.364106.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV140_280_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_s3126_r10201_r10210_p4133
 776 mc16_13TeV.364107.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV140_280_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_s3126_r10201_r10210_p4133

901 mc16_13TeV.364189.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV70_140_BFilter.deriv.DAOD_HIGG8D1.e5340_s3126_r10201_r10210_p4133
 902 mc16_13TeV.364189.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV70_140_BFilter.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10201_r10210_p4133
 903 mc16_13TeV.364190.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV140_280_CVetoBVeto.deriv.DAOD_HIGG8D1.e5340_s3126_r10201_r10210_p4133
 904 mc16_13TeV.364190.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV140_280_CVetoBVeto.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_s3136_r10201_r10210_p4133
 905 mc16_13TeV.364191.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV140_280_CFilterBVeto.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_s3136_r10201_r10210_p4133
 906 mc16_13TeV.364191.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV140_280_CFilterBVeto.deriv.DAOD_HIGG8D1.e5340_s3126_r10201_r10210_p4133
 907 mc16_13TeV.364192.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV140_280_BFilter.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10201_r10210_p4133
 908 mc16_13TeV.364193.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV280_500_CVetoBVeto.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10201_r10210_p4133
 909 mc16_13TeV.364193.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV280_500_CVetoBVeto.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_s3136_r10201_r10210_p4133
 910 mc16_13TeV.364194.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV280_500_CFilterBVeto.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10201_r10210_p4133
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 913 mc16_13TeV.364195.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV280_500_BFilter.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10201_r10210_p4133
 914 mc16_13TeV.364195.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV280_500_BFilter.deriv.DAOD_HIGG8D1.e5340_s3126_r10201_r10210_p4133
 915 mc16_13TeV.364196.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV500_1000.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_s3136_r10201_r10210_p4133
 916 mc16_13TeV.364196.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV500_1000.deriv.DAOD_HIGG8D1.e5340_s3126_r10201_r10210_p4133
 917 mc16_13TeV.364197.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV1000_E_CMS.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10201_r10210_p4133
 918 mc16_13TeV.364500.Sherpa_222_NNPDF30NNLO_eegamma_pty_7_15.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 919 mc16_13TeV.364501.Sherpa_222_NNPDF30NNLO_eegamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 920 mc16_13TeV.364502.Sherpa_222_NNPDF30NNLO_eegamma_pty_35_70.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 921 mc16_13TeV.364503.Sherpa_222_NNPDF30NNLO_eegamma_pty_70_140.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
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 924 mc16_13TeV.364506.Sherpa_222_NNPDF30NNLO_mumugamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
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 927 mc16_13TeV.364509.Sherpa_222_NNPDF30NNLO_mumugamma_pty_140_E_CMS.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 928 mc16_13TeV.364510.Sherpa_222_NNPDF30NNLO_tautaugamma_pty_7_15.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 929 mc16_13TeV.364511.Sherpa_222_NNPDF30NNLO_tautaugamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 930 mc16_13TeV.364512.Sherpa_222_NNPDF30NNLO_tautaugamma_pty_35_70.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 931 mc16_13TeV.364513.Sherpa_222_NNPDF30NNLO_tautaugamma_pty_70_140.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 932 mc16_13TeV.364513.Sherpa_222_NNPDF30NNLO_tautaugamma_pty_70_140.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 933 mc16_13TeV.364514.Sherpa_222_NNPDF30NNLO_tautaugamma_pty_140_E_CMS.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 934 mc16_13TeV.364521.Sherpa_222_NNPDF30NNLO_enugamma_pty_7_15.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 935 mc16_13TeV.364522.Sherpa_222_NNPDF30NNLO_enugamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 936 mc16_13TeV.364523.Sherpa_222_NNPDF30NNLO_enugamma_pty_35_70.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 937 mc16_13TeV.364524.Sherpa_222_NNPDF30NNLO_enugamma_pty_70_140.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 938 mc16_13TeV.364525.Sherpa_222_NNPDF30NNLO_enugamma_pty_140_E_CMS.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 939 mc16_13TeV.364526.Sherpa_222_NNPDF30NNLO_munugamma_pty_7_15.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 940 mc16_13TeV.364527.Sherpa_222_NNPDF30NNLO_munugamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 941 mc16_13TeV.364528.Sherpa_222_NNPDF30NNLO_munugamma_pty_35_70.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 942 mc16_13TeV.364529.Sherpa_222_NNPDF30NNLO_munugamma_pty_70_140.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 943 mc16_13TeV.364530.Sherpa_222_NNPDF30NNLO_munugamma_pty_140_E_CMS.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 944 mc16_13TeV.364531.Sherpa_222_NNPDF30NNLO_taunugamma_pty_7_15.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 945 mc16_13TeV.364532.Sherpa_222_NNPDF30NNLO_taunugamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
 946 mc16_13TeV.364533.Sherpa_222_NNPDF30NNLO_taunugamma_pty_35_70.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10201_r10210_p4133
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 949 mc16_13TeV.410560.MadGraphPythia8EvtGen_A14_tZ_4fl_tchan_noAllHad.deriv.DAOD_HIGG8D1.e5803_e5984_s3126_r10201_r10210_p4133
 950 mc16_13TeV.410013.PowhegPythiaEvtGen_P2012_Wt_inclusive_top.deriv.DAOD_HIGG8D1.e3753_s3126_r10201_r10210_p4133
 951 mc16_13TeV.410014.PowhegPythiaEvtGen_P2012_Wt_inclusive_antitop.deriv.DAOD_HIGG8D1.e3753_s3126_r10201_r10210_p4133
 952 mc16_13TeV.410408.aMcAtNloPythia8EvtGen_tWZ_Ztoll_minDR1.deriv.DAOD_HIGG8D1.e6423_e5984_s3126_r10201_r10210_p4133
 953 mc16_13TeV.364242.Sherpa_222_NNPDF30NNLO_WWW_313v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10201_r10210_p4133
 954 mc16_13TeV.364243.Sherpa_222_NNPDF30NNLO_WWZ_412v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10201_r10210_p4133
 955 mc16_13TeV.364244.Sherpa_222_NNPDF30NNLO_WWZ_214v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10201_r10210_p4133
 956 mc16_13TeV.364245.Sherpa_222_NNPDF30NNLO_WZZ_511v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10201_r10210_p4133
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 960 mc16_13TeV.364249.Sherpa_222_NNPDF30NNLO_ZZZ_214v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10201_r10210_p4133
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 962 mc16_13TeV.342285.Pythia8EvtGen_A14NNPDF23LO_ZH125_inc.deriv.DAOD_HIGG8D1.e4246_e5984_s3126_r10201_r10210_p4133

963 mc16_13TeV.341998.aMcAtNloHppEG_UEEE5_CTEQ6L1_CT10ME_tWH125_gamgam_yt_plus1.deriv.DAOD_HIGG8D1.e4394_e5984_s3126_r10201_r10201_r1
 964 mc16_13TeV.410470.PhPy8EG_A14_ttbar_hdamp258p75_nonallhad.deriv.DAOD_HIGG8D1.e6337_e5984_s3126_r10201_r10210_p4133
 965 mc16_13TeV.410472.PhPy8EG_A14_ttbar_hdamp258p75_dil.deriv.DAOD_HIGG8D1.e6348_e5984_s3126_r10201_r10210_p4133
 966 mc16_13TeV.346343.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_allhad.deriv.DAOD_HIGG8D1.e7148_e5984_s3126_r10201_r10210_p4133
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 975 mc16_13TeV.361603.PowhegPy8EG_CT10nloME_AZNLOCTEQ6L1_ZZllll_mll4.deriv.DAOD_HIGG8D1.e4475_e5984_s3126_r10724_r10726_p4133
 976 mc16_13TeV.361604.PowhegPy8EG_CT10nloME_AZNLOCTEQ6L1_ZZvlll_mll4.deriv.DAOD_HIGG8D1.e4475_e5984_s3126_r10724_r10726_p4133
 977 mc16_13TeV.364288.Sherpa_222_NNPDF30NNLO_llll_lowMllPtComplement.deriv.DAOD_HIGG8D1.e6096_s3126_r10724_p3983
 978 mc16_13TeV.364287.Sherpa_222_NNPDF30NNLO_llvj_ss_EW6.deriv.DAOD_HIGG8D1.e6055_e5984_s3126_r10724_r10726_p3983
 979 mc16_13TeV.364289.Sherpa_222_NNPDF30NNLO_lllv_lowMllPtComplement.deriv.DAOD_HIGG8D1.e6133_e5984_s3126_r10724_r10726_p3983
 980 mc16_13TeV.364285.Sherpa_222_NNPDF30NNLO_llvj_EW6.deriv.DAOD_HIGG8D1.e6055_e5984_s3126_r10724_r10726_p3983
 981 mc16_13TeV.364284.Sherpa_222_NNPDF30NNLO_llljj_EW6.deriv.DAOD_HIGG8D1.e6055_s3126_r10724_p3983
 982 mc16_13TeV.364283.Sherpa_222_NNPDF30NNLO_llljjj_EW6.deriv.DAOD_HIGG8D1.e6055_s3126_r10724_p3983
 983 mc16_13TeV.364739.MGPY8EG_NNPDF30NLO_A14NNPDF23LO_lvlljjEW6_OFMinus.deriv.DAOD_HIGG8D1.e7421_e5984_s3126_r10724_r10726_p4133
 984 mc16_13TeV.364742.MGPY8EG_NNPDF30NLO_A14NNPDF23LO_lvlljjEW6_SFPlus.deriv.DAOD_HIGG8D1.e7421_e5984_s3126_r10724_r10726_p4133
 985 mc16_13TeV.364740.MGPY8EG_NNPDF30NLO_A14NNPDF23LO_lvlljjEW6_OFPlus.deriv.DAOD_HIGG8D1.e7421_e5984_s3126_r10724_r10726_p4133
 986 mc16_13TeV.364741.MGPY8EG_NNPDF30NLO_A14NNPDF23LO_lvlljjEW6_SFMinus.deriv.DAOD_HIGG8D1.e7421_e5984_s3126_r10724_r10726_p4133
 987 mc16_13TeV.364253.Sherpa_222_NNPDF30NNLO_llvv.deriv.DAOD_HIGG8D1.e5916_e5984_s3126_r10724_r10726_p4133
 988 mc16_13TeV.364250.Sherpa_222_NNPDF30NNLO_llll.deriv.DAOD_HIGG8D1.e5894_e5984_s3126_r10724_r10726_p4133
 989 mc16_13TeV.364254.Sherpa_222_NNPDF30NNLO_llvv.deriv.DAOD_HIGG8D1.e5916_e5984_s3126_r10724_r10726_p4133
 990 mc16_13TeV.364255.Sherpa_222_NNPDF30NNLO_lvvv.deriv.DAOD_HIGG8D1.e5916_e5984_s3126_r10724_r10726_p4133
 991 mc16_13TeV.410155.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttW.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10724_r10726_p4133
 992 mc16_13TeV.410156.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttZnunu.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10724_r10726_p4133
 993 mc16_13TeV.410157.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttZqq.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10724_r10726_p4133
 994 mc16_13TeV.410218.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttee.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10724_r10726_p4133
 995 mc16_13TeV.410219.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_ttmmumu.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10724_r10726_p4133
 996 mc16_13TeV.410220.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_tttautau.deriv.DAOD_HIGG8D1.e5070_e5984_s3126_r10724_r10726_p4133
 997 mc16_13TeV.410276.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_tteee_mll_1_5.deriv.DAOD_HIGG8D1.e6087_e5984_s3126_r10724_r10726_p4133
 998 mc16_13TeV.410277.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_tttautau_mll_1_5.deriv.DAOD_HIGG8D1.e6087_e5984_s3126_r10724_r10726_p4133
 999 mc16_13TeV.410278.aMcAtNloPythia8EvtGen_MEN30NLO_A14N23LO_tttautau_mll_1_5.deriv.DAOD_HIGG8D1.e6087_e5984_s3126_r10724_r10726_p4133
 1000 mc16_13TeV.410397.MadGraphPythia8EvtGen_ttbar_wbee_MEN30LO_A14N23LO.deriv.DAOD_HIGG8D1.e6086_e5984_s3126_r10724_r10726_p4133
 1001 mc16_13TeV.410398.MadGraphPythia8EvtGen_ttbar_wbmumu_MEN30LO_A14N23LO.deriv.DAOD_HIGG8D1.e6086_e5984_s3126_r10724_r10726_p4133
 1002 mc16_13TeV.410399.MadGraphPythia8EvtGen_ttbar_wbtautau_MEN30LO_A14N23LO.deriv.DAOD_HIGG8D1.e6086_e5984_s3126_r10724_r10726_p4133
 1003 mc16_13TeV.410658.PhPy8EG_A14_tchan_BW50_lept_top.deriv.DAOD_HIGG8D1.e6671_e5984_s3126_s3136_r10724_r10726_p4133
 1004 mc16_13TeV.410659.PhPy8EG_A14_tchan_BW50_lept_antitop.deriv.DAOD_HIGG8D1.e6671_e5984_s3126_s3136_r10724_r10726_p4133
 1005 mc16_13TeV.410644.PowhegPythia8EvtGen_A14_singletop_schan_lept_top.deriv.DAOD_HIGG8D1.e6527_e5984_s3126_r10724_r10726_p4133
 1006 mc16_13TeV.410645.PowhegPythia8EvtGen_A14_singletop_schan_lept_antitop.deriv.DAOD_HIGG8D1.e6527_e5984_s3126_r10724_r10726_p4133
 1007 mc16_13TeV.304014.MadGraphPythia8EvtGen_A14NNPDF23_3top_SM.deriv.DAOD_HIGG8D1.e4324_e5984_s3126_r10724_r10726_p4133
 1008 mc16_13TeV.410080.MadGraphPythia8EvtGen_A14NNPDF23_4topSM.deriv.DAOD_HIGG8D1.e4111_e5984_s3126_r10724_r10726_p4133
 1009 mc16_13TeV.410081.MadGraphPythia8EvtGen_A14NNPDF23_ttbarWW.deriv.DAOD_HIGG8D1.e4111_e5984_s3126_r10724_r10726_p4133
 1010 mc16_13TeV.364100.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV0_70_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_s3136_r10724
 1011 mc16_13TeV.364100.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV0_70_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 1012 mc16_13TeV.364101.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV0_70_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 1013 mc16_13TeV.364102.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV0_70_BFilter.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_s3136_r10724_r10726_p4133
 1014 mc16_13TeV.364102.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV0_70_BFilter.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 1015 mc16_13TeV.364103.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV70_140_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 1016 mc16_13TeV.364103.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV70_140_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_s3136_r10724_r10726_p4133
 1017 mc16_13TeV.364104.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV70_140_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 1018 mc16_13TeV.364104.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV70_140_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_s3136_r10724_r10726_p4133
 1019 mc16_13TeV.364106.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV140_280_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_s3136_r10724_r10726_p4133
 1020 mc16_13TeV.364106.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV140_280_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 1021 mc16_13TeV.364107.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV140_280_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 1022 mc16_13TeV.364107.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV140_280_CFilterBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_s3136_r10724_r10726_p4133
 1023 mc16_13TeV.364108.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV140_280_BFilter.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133
 1024 mc16_13TeV.364109.Sherpa_221_NNPDF30NNLO_Zmumu_MAXHTPTV280_500_CVetoBVeto.deriv.DAOD_HIGG8D1.e5271_e5984_s3126_r10724_r10726_p4133

1149 mc16_13TeV.364192.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV140_280_BFilter.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_s3136_r10724_r1
 1150 mc16_13TeV.364192.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV140_280_BFilter.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10724_r10726_p1
 1151 mc16_13TeV.364193.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV280_500_CVetoBVeto.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10724_r10726_p1
 1152 mc16_13TeV.364193.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV280_500_CVetoBVeto.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_s3136_r10724_r10726_p1
 1153 mc16_13TeV.364194.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV280_500_CFilterBVeto.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_s3136_r10724_r10726_p1
 1154 mc16_13TeV.364194.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV280_500_CFilterBVeto.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10724_r10726_p1
 1155 mc16_13TeV.364195.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV280_500_BFilter.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_s3136_r10724_r10726_p1
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 1157 mc16_13TeV.364196.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV500_1000.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_s3136_r10724_r10726_p1
 1158 mc16_13TeV.364196.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV500_1000.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10724_r10726_p4133
 1159 mc16_13TeV.364197.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV1000_E_CMS.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_s3136_r10724_r10726_p1
 1160 mc16_13TeV.364197.Sherpa_221_NNPDF30NNLO_Wtaunu_MAXHTPTV1000_E_CMS.deriv.DAOD_HIGG8D1.e5340_e5984_s3126_r10724_r10726_p4133
 1161 mc16_13TeV.364500.Sherpa_222_NNPDF30NNLO_eegamma_pty_7_15.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 1162 mc16_13TeV.364501.Sherpa_222_NNPDF30NNLO_eegamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 1163 mc16_13TeV.364502.Sherpa_222_NNPDF30NNLO_eegamma_pty_35_70.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 1164 mc16_13TeV.364503.Sherpa_222_NNPDF30NNLO_eegamma_pty_70_140.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
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 1167 mc16_13TeV.364506.Sherpa_222_NNPDF30NNLO_mumugamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 1168 mc16_13TeV.364507.Sherpa_222_NNPDF30NNLO_mumugamma_pty_35_70.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 1169 mc16_13TeV.364508.Sherpa_222_NNPDF30NNLO_mumugamma_pty_70_140.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 1170 mc16_13TeV.364509.Sherpa_222_NNPDF30NNLO_mumugamma_pty_140_E_CMS.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 1171 mc16_13TeV.364511.Sherpa_222_NNPDF30NNLO_tautaugamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 1172 mc16_13TeV.364512.Sherpa_222_NNPDF30NNLO_tautaugamma_pty_35_70.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 1173 mc16_13TeV.364513.Sherpa_222_NNPDF30NNLO_tautaugamma_pty_70_140.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 1174 mc16_13TeV.364514.Sherpa_222_NNPDF30NNLO_tautaugamma_pty_140_E_CMS.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 1175 mc16_13TeV.364522.Sherpa_222_NNPDF30NNLO_enugamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
 1176 mc16_13TeV.364523.Sherpa_222_NNPDF30NNLO_enugamma_pty_35_70.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
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 1183 mc16_13TeV.364532.Sherpa_222_NNPDF30NNLO_taunugamma_pty_15_35.deriv.DAOD_HIGG8D1.e5928_e5984_s3126_r10724_r10726_p4133
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 1188 mc16_13TeV.410408.aMcAtNloPythia8EvtGen_tWZ_Ztoll_minDR1.deriv.DAOD_HIGG8D1.e6423_e5984_s3126_r10724_r10726_p4133
 1189 mc16_13TeV.364242.Sherpa_222_NNPDF30NNLO_WWW_3l3v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10724_r10726_p4133
 1190 mc16_13TeV.364243.Sherpa_222_NNPDF30NNLO_WWZ_4l2v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10724_r10726_p4133
 1191 mc16_13TeV.364244.Sherpa_222_NNPDF30NNLO_WWZ_2l4v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10724_r10726_p4133
 1192 mc16_13TeV.364245.Sherpa_222_NNPDF30NNLO_WZZ_5l1v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10724_r10726_p4133
 1193 mc16_13TeV.364246.Sherpa_222_NNPDF30NNLO_WZZ_3l3v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10724_r10726_p4133
 1194 mc16_13TeV.364247.Sherpa_222_NNPDF30NNLO_ZZZ_6l0v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10724_r10726_p4133
 1195 mc16_13TeV.364248.Sherpa_222_NNPDF30NNLO_ZZZ_4l2v_EW6.deriv.DAOD_HIGG8D1.e5887_e5984_s3126_r10724_r10726_p4133
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 1197 mc16_13TeV.342284.Pythia8EvtGen_A14NNPDF23LO_WH125_inc.deriv.DAOD_HIGG8D1.e4246_e5984_s3126_r10724_r10726_p4133
 1198 mc16_13TeV.342285.Pythia8EvtGen_A14NNPDF23LO_ZH125_inc.deriv.DAOD_HIGG8D1.e4246_e5984_s3126_r10724_r10726_p4133
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 1201 mc16_13TeV.410470.PhPy8EG_A14_ttbar_hdamp258p75_nonallhad.deriv.DAOD_HIGG8D1.e6337_e5984_s3126_r10724_r10726_p4133
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 1206 mc16_13TeV.346344.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_semilep.deriv.DAOD_HIGG8D1.e7148_e5984_s3126_r10724_r10726_p4133
 1207 mc16_13TeV.346345.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_dilep.deriv.DAOD_HIGG8D1.e7148_e5984_a875_r10724_r10726_p4133
 1208 mc16_13TeV.346345.PhPy8EG_A14NNPDF23_NNPDF30ME_ttH125_dilep.deriv.DAOD_HIGG8D1.e7148_e5984_s3126_r10724_r10726_p4133
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