

University of Science and Technology of China

September 20, 2024

Resurgent large genus asymptotics of intersection numbers

j/w B. Eynard, E. Garcia-Failde, P. Gregori, D. Lewański

arXiv: [AG/2309.03143](#)

Alessandro Giacchetto

ETH Zurich

A case study: $m!$

Counting problem: $c_m = \# \left\{ \begin{array}{l} \text{arrangements of } m \text{ distinct objects} \\ \text{into } m \text{ distinct boxes} \end{array} \right\}$

Solution: $c_m = m! = \begin{cases} m \cdot c_{m-1} & m > 1 \\ 1 & m = 1 \end{cases}$

Pro: exact

Con: recursive

Asymptotics: $c_m = \sqrt{2\pi m} \left(\frac{m}{e}\right)^m \left(1 + O(m^{-1})\right)$

Con: asymptotically exact

Pro: closed-form

A case study: $m!$

Counting problem: $c_m = \# \left\{ \begin{array}{l} \text{arrangements of } m \text{ distinct objects} \\ \text{into } m \text{ distinct boxes} \end{array} \right\}$

Solution:
$$c_m = m! = \begin{cases} m \cdot c_{m-1} & m > 1 \\ 1 & m = 1 \end{cases}$$

Pro: exact

Con: recursive

Asymptotics:
$$c_m = \sqrt{2\pi m} \left(\frac{m}{e}\right)^m \left(1 + O(m^{-1})\right)$$

Con: asymptotically exact

Pro: closed-form

A case study: $m!$

Counting problem: $c_m = \# \left\{ \begin{array}{l} \text{arrangements of } m \text{ distinct objects} \\ \text{into } m \text{ distinct boxes} \end{array} \right\}$

Solution:
$$c_m = m! = \begin{cases} m \cdot c_{m-1} & m > 1 \\ 1 & m = 1 \end{cases}$$

Pro: exact

Con: recursive

Asymptotics:
$$c_m = \sqrt{2\pi m} \left(\frac{m}{e}\right)^m \left(1 + O(m^{-1})\right)$$

Con: asymptotically exact

Pro: closed-form

A case study: $m!$

Counting problem: $c_m = \# \left\{ \begin{array}{l} \text{arrangements of } m \text{ distinct objects} \\ \text{into } m \text{ distinct boxes} \end{array} \right\}$

Solution:
$$c_m = m! = \begin{cases} m \cdot c_{m-1} & m > 1 \\ 1 & m = 1 \end{cases}$$

Pro: exact

Con: recursive

Asymptotics:
$$c_m = \sqrt{2\pi m} \left(\frac{m}{e}\right)^m \left(1 + O(m^{-1})\right)$$

Con: asymptotically exact

Pro: closed-form

A case study: $m!$

Counting problem: $c_m = \# \left\{ \begin{array}{l} \text{arrangements of } m \text{ distinct objects} \\ \text{into } m \text{ distinct boxes} \end{array} \right\}$

Solution:
$$c_m = m! = \begin{cases} m \cdot c_{m-1} & m > 1 \\ 1 & m = 1 \end{cases}$$

Pro: exact

Con: recursive

Asymptotics:
$$c_m = \sqrt{2\pi m} \left(\frac{m}{e}\right)^m \left(1 + O(m^{-1})\right)$$

Con: asymptotically exact

Pro: closed-form

A case study: $m!$

Counting problem: $c_m = \# \left\{ \begin{array}{l} \text{arrangements of } m \text{ distinct objects} \\ \text{into } m \text{ distinct boxes} \end{array} \right\}$

Solution:
$$c_m = m! = \begin{cases} m \cdot c_{m-1} & m > 1 \\ 1 & m = 1 \end{cases}$$

Pro: exact

Con: recursive

Asymptotics:
$$c_m = \sqrt{2\pi m} \left(\frac{m}{e}\right)^m \left(1 + \frac{1}{12}m^{-1} + O(m^{-2})\right)$$

Con: asymptotically exact

Pro: closed-form

A case study: $m!$

Counting problem: $c_m = \# \left\{ \begin{array}{l} \text{arrangements of } m \text{ distinct objects} \\ \text{into } m \text{ distinct boxes} \end{array} \right\}$

Solution:
$$c_m = m! = \begin{cases} m \cdot c_{m-1} & m > 1 \\ 1 & m = 1 \end{cases}$$

Pro: exact

Con: recursive

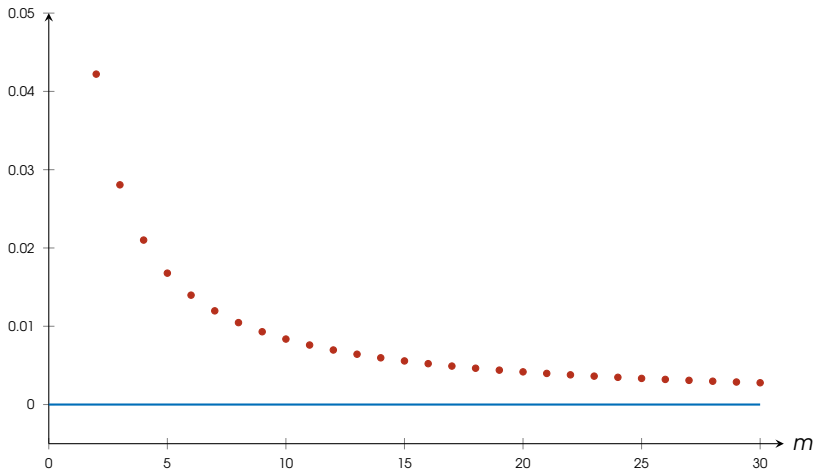
Asymptotics:
$$c_m = \sqrt{2\pi m} \left(\frac{m}{e}\right)^m \left(1 + \frac{1}{12}m^{-1} + \frac{1}{288}m^{-2} + O(m^{-3})\right)$$

Con: asymptotically exact

Pro: closed-form

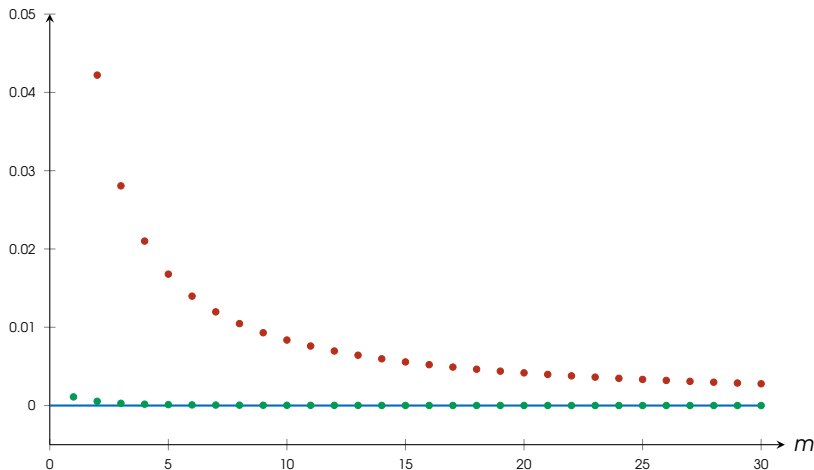
Visualising Stirling's formula

$$\frac{m!}{\sqrt{2\pi m} \left(\frac{m}{e}\right)^m} - 1 = O(m^{-1})$$



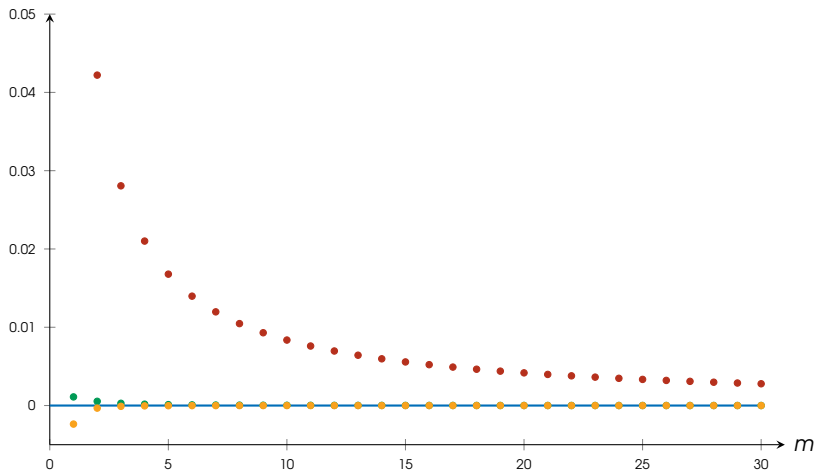
Visualising Stirling's formula

$$\frac{m!}{\sqrt{2\pi m} \left(\frac{m}{e}\right)^m} - \left(1 + \frac{1}{12}m^{-1}\right) = \mathcal{O}(m^{-2})$$



Visualising Stirling's formula

$$\frac{m!}{\sqrt{2\pi m} \left(\frac{m}{e}\right)^m} - \left(1 + \frac{1}{12}m^{-1} + \frac{1}{288}m^{-2}\right) = \mathcal{O}(m^{-3})$$



ψ -class intersection numbers

$$\langle \tau_{d_1} \cdots \tau_{d_n} \rangle = \int_{\overline{\mathcal{M}}_{g,n}} \psi_1^{d_1} \cdots \psi_n^{d_n} \quad d_i \geq 0, \quad d_1 + \cdots + d_n = 3g - 3 + n$$

- Volumes of moduli spaces of **metric ribbon graphs**
- Building block for all **tautological intersection numbers**:
 - Weil–Petersson volumes
 - Masur–Veech volumes
 - Hurwitz numbers
 - Gromov–Witten invariants for targets with s.s. quantum cohomology
 - ...
- Compute the perturbative expansion of **topological 2d gravity**

ψ -class intersection numbers

$$\langle \tau_{d_1} \cdots \tau_{d_n} \rangle = \int_{\overline{\mathcal{M}}_{g,n}} \psi_1^{d_1} \cdots \psi_n^{d_n} \quad d_i \geq 0, \quad d_1 + \cdots + d_n = 3g - 3 + n$$

- Volumes of moduli spaces of **metric ribbon graphs**
- Building block for all **tautological intersection numbers**:
 - Weil–Petersson volumes
 - Masur–Veech volumes
 - Hurwitz numbers
 - Gromov–Witten invariants for targets with s.s. quantum cohomology
 - ...
- Compute the perturbative expansion of **topological 2d gravity**

ψ -class intersection numbers

$$\langle \tau_{d_1} \cdots \tau_{d_n} \rangle = \int_{\overline{\mathcal{M}}_{g,n}} \psi_1^{d_1} \cdots \psi_n^{d_n} \quad d_i \geq 0, \quad d_1 + \cdots + d_n = 3g - 3 + n$$

- Volumes of moduli spaces of **metric ribbon graphs**
- Building block for all **tautological intersection numbers**:
 - Weil–Petersson volumes
 - Masur–Veech volumes
 - Hurwitz numbers
 - Gromov–Witten invariants for targets with s.s. quantum cohomology
 - ...
- Compute the perturbative expansion of **topological 2d gravity**

ψ -class intersection numbers

$$\langle \tau_{d_1} \cdots \tau_{d_n} \rangle = \int_{\overline{\mathcal{M}}_{g,n}} \psi_1^{d_1} \cdots \psi_n^{d_n} \quad d_i \geq 0, \quad d_1 + \cdots + d_n = 3g - 3 + n$$

- Volumes of moduli spaces of **metric ribbon graphs**
- Building block for all **tautological intersection numbers**:
 - Weil–Petersson volumes
 - Masur–Veech volumes
 - Hurwitz numbers
 - Gromov–Witten invariants for targets with s.s. quantum cohomology
 - ...
- Compute the perturbative expansion of **topological 2d gravity**

Solution

Normalisation: $\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = \langle \tau_{d_1} \cdots \tau_{d_n} \rangle \prod_{i=1}^n (2d_i + 1)!!$

Witten conjecture/Kontsevich theorem, early '90s:

$$\begin{aligned} \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle &= \sum_{m=2}^n (2d_m + 1) \langle\langle \tau_{d_1+d_m-1} \tau_{d_2} \cdots \widehat{\tau_{d_m}} \cdots \tau_{d_n} \rangle\rangle \\ &\quad + \frac{1}{2} \sum_{a+b=d_1-2} \left(\langle\langle \tau_a \tau_b \tau_{d_2} \cdots \tau_{d_n} \rangle\rangle + \sum_{\substack{g_1+g_2=g \\ l_1 \sqcup l_2 = \{d_2, \dots, d_n\}}} \langle\langle \tau_a \tau_{l_1} \rangle\rangle \langle\langle \tau_b \tau_{l_2} \rangle\rangle \right) \end{aligned}$$

with initial data $\langle\langle \tau_0^3 \rangle\rangle = 1$ and $\langle\langle \tau_1 \rangle\rangle = \frac{1}{8}$.

Solution

Normalisation: $\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = \langle \tau_{d_1} \cdots \tau_{d_n} \rangle \prod_{i=1}^n (2d_i + 1)!!$

Witten conjecture/Kontsevich theorem, early '90s:

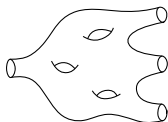
$$\begin{aligned} \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle &= \sum_{m=2}^n (2d_m + 1) \langle\langle \tau_{d_1+d_m-1} \tau_{d_2} \cdots \widehat{\tau_{d_m}} \cdots \tau_{d_n} \rangle\rangle \\ &\quad + \frac{1}{2} \sum_{a+b=d_1-2} \left(\langle\langle \tau_a \tau_b \tau_{d_2} \cdots \tau_{d_n} \rangle\rangle + \sum_{\substack{g_1+g_2=g \\ l_1 \sqcup l_2 = \{d_2, \dots, d_n\}}} \langle\langle \tau_a \tau_{l_1} \rangle\rangle \langle\langle \tau_b \tau_{l_2} \rangle\rangle \right) \end{aligned}$$

with initial data $\langle\langle \tau_0^3 \rangle\rangle = 1$ and $\langle\langle \tau_1 \rangle\rangle = \frac{1}{8}$.

Solution

Normalisation: $\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = \langle \tau_{d_1} \cdots \tau_{d_n} \rangle \prod_{i=1}^n (2d_i + 1)!!$

Witten conjecture/Kontsevich theorem, early '90s:



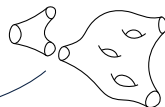
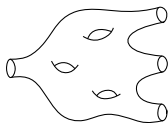
$$\begin{aligned}
 \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle &= \sum_{m=2}^n (2d_m + 1) \langle\langle \tau_{d_1+d_m-1} \tau_{d_2} \cdots \widehat{\tau_{d_m}} \cdots \tau_{d_n} \rangle\rangle \\
 &+ \frac{1}{2} \sum_{a+b=d_1-2} \left(\langle\langle \tau_a \tau_b \tau_{d_2} \cdots \tau_{d_n} \rangle\rangle + \sum_{\substack{g_1+g_2=g \\ l_1 \sqcup l_2 = \{d_2, \dots, d_n\}}} \langle\langle \tau_a \tau_{l_1} \rangle\rangle \langle\langle \tau_b \tau_{l_2} \rangle\rangle \right)
 \end{aligned}$$

with initial data $\langle\langle \tau_0^3 \rangle\rangle = 1$ and $\langle\langle \tau_1 \rangle\rangle = \frac{1}{8}$.

Solution

Normalisation: $\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = \langle \tau_{d_1} \cdots \tau_{d_n} \rangle \prod_{i=1}^n (2d_i + 1)!!$

Witten conjecture/Kontsevich theorem, early '90s:



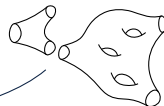
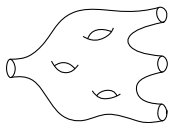
$$\begin{aligned}
 \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle &= \sum_{m=2}^n (2d_m + 1) \langle\langle \tau_{d_1+d_m-1} \tau_{d_2} \cdots \widehat{\tau_{d_m}} \cdots \tau_{d_n} \rangle\rangle \\
 &\quad + \frac{1}{2} \sum_{a+b=d_1-2} \left(\langle\langle \tau_a \tau_b \tau_{d_2} \cdots \tau_{d_n} \rangle\rangle + \sum_{\substack{g_1+g_2=g \\ l_1 \sqcup l_2 = \{d_2, \dots, d_n\}}} \langle\langle \tau_a \tau_{l_1} \rangle\rangle \langle\langle \tau_b \tau_{l_2} \rangle\rangle \right)
 \end{aligned}$$

with initial data $\langle\langle \tau_0^3 \rangle\rangle = 1$ and $\langle\langle \tau_1 \rangle\rangle = \frac{1}{8}$.

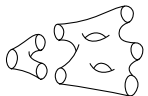
Solution

Normalisation: $\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = \langle \tau_{d_1} \cdots \tau_{d_n} \rangle \prod_{i=1}^n (2d_i + 1)!!$

Witten conjecture/Kontsevich theorem, early '90s:



$$\begin{aligned} \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle &= \sum_{m=2}^n (2d_m + 1) \langle\langle \tau_{d_1+d_m-1} \tau_{d_2} \cdots \widehat{\tau_{d_m}} \cdots \tau_{d_n} \rangle\rangle \\ &\quad + \frac{1}{2} \sum_{a+b=d_1-2} \left(\langle\langle \tau_a \tau_b \tau_{d_2} \cdots \tau_{d_n} \rangle\rangle + \sum_{\substack{g_1+g_2=g \\ l_1 \sqcup l_2 = \{d_2, \dots, d_n\}}} \langle\langle \tau_a \tau_{l_1} \rangle\rangle \langle\langle \tau_b \tau_{l_2} \rangle\rangle \right) \end{aligned}$$

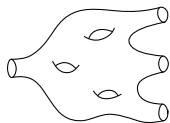
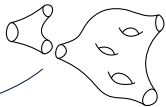


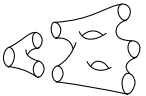
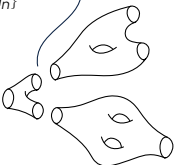
with initial data $\langle\langle \tau_0^3 \rangle\rangle = 1$ and $\langle\langle \tau_1 \rangle\rangle = \frac{1}{8}$.

Solution

Normalisation: $\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = \langle \tau_{d_1} \cdots \tau_{d_n} \rangle \prod_{i=1}^n (2d_i + 1)!!$

Witten conjecture/Kontsevich theorem, early '90s:

$$\begin{aligned}
 \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle &= \sum_{m=2}^n (2d_m + 1) \langle\langle \tau_{d_1+d_m-1} \tau_{d_2} \cdots \widehat{\tau_{d_m}} \cdots \tau_{d_n} \rangle\rangle \\
 &\quad + \frac{1}{2} \sum_{a+b=d_1-2} \left(\langle\langle \tau_a \tau_b \tau_{d_2} \cdots \tau_{d_n} \rangle\rangle + \sum_{\substack{g_1+g_2=g \\ l_1 \sqcup l_2 = \{d_2, \dots, d_n\}}} \langle\langle \tau_a \tau_{l_1} \rangle\rangle \langle\langle \tau_b \tau_{l_2} \rangle\rangle \right)
 \end{aligned}$$



with initial data $\langle\langle \tau_0^3 \rangle\rangle = 1$ and $\langle\langle \tau_1 \rangle\rangle = \frac{1}{8}$.

Large genus asymptotics

Uniformly in d_1, \dots, d_n as $g \rightarrow \infty$:

$$\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{(\frac{2}{3})^{2g-2+n}} (1 + o(g^{-1}))$$

- Conjectured by [Delecroix–Goujard–Zograf–Zorich](#), 2019
- Proved by [Aggarwal](#), 2020
(combinatorial analysis of Witten–Kontsevich topological recursion)
- Proved by [Guo–Yang](#), 2021
(combinatorial analysis of the determinantal formula)

Questions

- Universal strategy, adaptable to different problems?
- ‘Geometric’ meaning of the formula?
- Subleading corrections?

Large genus asymptotics

Uniformly in d_1, \dots, d_n as $g \rightarrow \infty$:

$$\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{(\frac{2}{3})^{2g-2+n}} (1 + o(g^{-1}))$$

- Conjectured by [Delecroix–Goujard–Zograf–Zorich](#), 2019
- Proved by [Aggarwal](#), 2020
(combinatorial analysis of Witten–Kontsevich topological recursion)
- Proved by [Guo–Yang](#), 2021
(combinatorial analysis of the determinantal formula)

Questions

- Universal strategy, adaptable to different problems?
- ‘Geometric’ meaning of the formula?
- Subleading corrections?

Large genus asymptotics

Uniformly in d_1, \dots, d_n as $g \rightarrow \infty$:

$$\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{(\frac{2}{3})^{2g-2+n}} \left(1 + O(g^{-1})\right)$$

- Conjectured by [Delecroix–Goujard–Zograf–Zorich](#), 2019
- Proved by [Aggarwal](#), 2020
(combinatorial analysis of Witten–Kontsevich topological recursion)
- Proved by [Guo–Yang](#), 2021
(combinatorial analysis of the determinantal formula)

Questions

- Universal strategy, adaptable to different problems?
- ‘Geometric’ meaning of the formula?
- Subleading corrections?

Large genus asymptotics: new perspective

Answers

- Universal strategy: resurgent analysis of the determinantal formula
- Geometric meaning: Airy functions

$$y^2 - x = 0 \quad \xrightarrow{\text{quantisation}} \quad \left(\hbar^2 \frac{d^2}{dx^2} - x \right) \psi(x, \hbar) = 0$$

- Subleading corrections: algorithm + properties

Uniformly in d_1, \dots, d_n as $g \rightarrow \infty$ (here $(x)_k = x(x-1) \cdots (x-k+1)$):

$$\begin{aligned} \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = S \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{A^{2g-2+n}} & \left(1 + \frac{A}{2g-3+n} \alpha_1 + \cdots \right. \\ & \left. + \frac{A^k}{(2g-3+n)_k} \alpha_k + O(g^{-k-1}) \right) \end{aligned}$$

Large genus asymptotics: new perspective

Answers

- Universal strategy: resurgent analysis of the determinantal formula
- Geometric meaning: Airy functions

$$y^2 - x = 0 \quad \xrightarrow{\text{quantisation}} \quad \left(\hbar^2 \frac{d^2}{dx^2} - x \right) \psi(x, \hbar) = 0$$

- Subleading corrections: algorithm + properties

Uniformly in d_1, \dots, d_n as $g \rightarrow \infty$ (here $(x)_k = x(x-1)\cdots(x-k+1)$):

$$\begin{aligned} \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = S \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{A^{2g-2+n}} & \left(1 + \frac{A}{2g-3+n} \alpha_1 + \cdots \right. \\ & \left. + \frac{A^k}{(2g-3+n)_k} \alpha_k + O(g^{-k-1}) \right) \end{aligned}$$

Large genus asymptotics: new perspective

Answers

- Universal strategy: resurgent analysis of the determinantal formula
- Geometric meaning: Airy functions

$$y^2 - x = 0 \quad \xrightarrow{\text{quantisation}} \quad \left(\hbar^2 \frac{d^2}{dx^2} - x \right) \psi(x, \hbar) = 0$$

- Subleading corrections: algorithm + properties

Uniformly in d_1, \dots, d_n as $g \rightarrow \infty$ (here $(x)_k = x(x-1)\cdots(x-k+1)$):

$$\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = S \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{A^{2g-2+n}} \left(1 + \frac{A}{2g-3+n} \alpha_1 + \cdots \right. \\ \left. + \frac{A^k}{(2g-3+n)_k} \alpha_k + O(g^{-k-1}) \right)$$

$S = 1$
Stokes constant

Large genus asymptotics: new perspective

Answers

- Universal strategy: resurgent analysis of the determinantal formula
- Geometric meaning: Airy functions

$$y^2 - x = 0 \quad \xrightarrow{\text{quantisation}} \quad \left(\hbar^2 \frac{d^2}{dx^2} - x \right) \psi(x, \hbar) = 0$$

- Subleading corrections: algorithm + properties

Uniformly in d_1, \dots, d_n as $g \rightarrow \infty$ (here $(x)_k = x(x-1)\cdots(x-k+1)$):

$$\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = S \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{A^{2g-2+n}} \left(1 + \frac{A}{2g-3+n} \alpha_1 + \cdots \right. \\ \left. + \frac{A^k}{(2g-3+n)_k} \alpha_k + O(g^{-k-1}) \right)$$

$$A = 2/3$$

$$\psi \sim \frac{1}{\sqrt{2x^{1/4}}} e^{\pm \frac{A}{\hbar} x^{-3/2}}$$

Large genus asymptotics: new perspective

Answers

- Universal strategy: resurgent analysis of the determinantal formula
- Geometric meaning: Airy functions

$$y^2 - x = 0 \quad \xrightarrow{\text{quantisation}} \quad \left(\hbar^2 \frac{d^2}{dx^2} - x \right) \psi(x, \hbar) = 0$$

- Subleading corrections: algorithm + properties

Uniformly in d_1, \dots, d_n as $g \rightarrow \infty$ (here $(x)_k = x(x-1)\cdots(x-k+1)$):

$$\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = S \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{A^{2g-2+n}} \left(1 + \frac{A}{2g-3+n} \alpha_1 + \cdots \right. \\ \left. + \frac{A^k}{(2g-3+n)_k} \alpha_k + O(g^{-k-1}) \right)$$

Computable; polynomial in n and multiplicities of d_i

Large genus asymptotics: new perspective

Answers

- Universal strategy: resurgent analysis of the determinantal formula
- Geometric meaning: Airy functions

$$y^2 - x = 0 \quad \xrightarrow{\text{quantisation}} \quad \left(\hbar^2 \frac{d^2}{dx^2} - x \right) \psi(x, \hbar) = 0$$

- Subleading corrections: algorithm + properties

Uniformly in d_1, \dots, d_n as $g \rightarrow \infty$ (here $(x)_k = x(x-1)\cdots(x-k+1)$):

$$\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = S \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{A^{2g-2+n}} \left(1 + \frac{A}{2g-3+n} \alpha_1 + \cdots \right. \\ \left. \frac{A^k}{(2g-3+n)_k} \alpha_k + O(g^{-k-1}) \right)$$

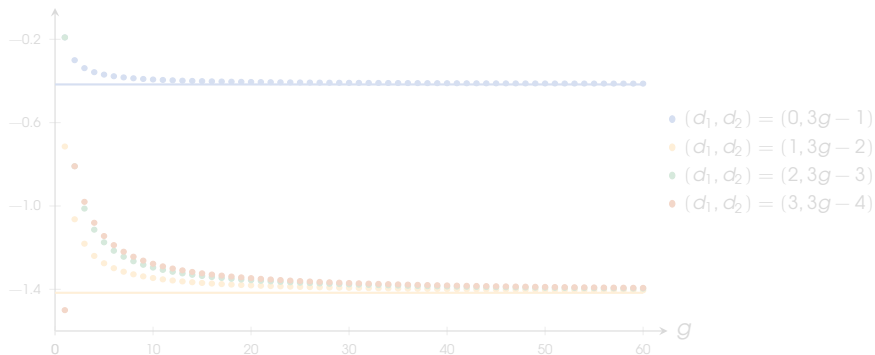
$$\alpha_1 = -\frac{17-15n+3n^2}{12} - \frac{(3-n)(n-p_0)}{2} - \frac{(n-p_0)_2}{4}$$

where $p_0 = \#\{d_i = 0\}$

Visualising the large genus asymptotics

$$\frac{2g-3+n}{2/3} \left(\frac{\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle}{\frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{(2/3)^{2g-2+n}}} - 1 \right) = \alpha_1(n, p_0) + O(g^{-1})$$

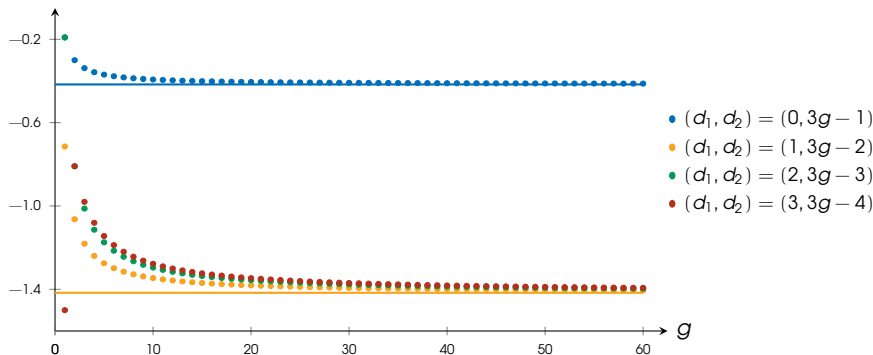
For $n = 2$:



Visualising the large genus asymptotics

$$\frac{2g-3+n}{2/3} \left(\frac{\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle}{\frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{(2/3)^{2g-2+n}}} - 1 \right) = \alpha_1(n, p_0) + O(g^{-1})$$

For $n = 2$:



Darboux meets Borel

Darboux's idea:

- Abs. convergent power series:

$$\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$$

- Get a holomorphic function around the origin, take analytic continuation
- The large m **asymptotics** of a_m is totally controlled by the behaviour of $\hat{\varphi}$ at its **singularities**

Borel's idea:

- Divergent power series:

$$\tilde{\varphi}(\hbar) = \sum_{m \geq 0} a_m \hbar^m, \quad |a_m| = O(A^{-m} m!)$$

- The **Borel transform**

$$\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$$

is now convergent

- Apply Darboux's idea

Darboux meets Borel

Darboux's idea:

- Abs. convergent power series:

$$\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$$

- Get a holomorphic function around the origin, take analytic continuation
- The large m **asymptotics** of a_m is totally controlled by the behaviour of $\hat{\varphi}$ at its **singularities**

Borel's idea:

- Divergent power series:

$$\tilde{\varphi}(\hbar) = \sum_{m \geq 0} a_m \hbar^m, \quad |a_m| = O(A^{-m} m!)$$

- The **Borel transform**

$$\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$$

is now convergent

- Apply Darboux's idea

Darboux's result: sketch of the proof

Take a convergent power series: $\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$

Suppose its analytic continuation has a single log singularity at $s = A$:

$$\hat{\varphi}(s) = (\text{holomorphic @ } A) \log(s - A) + \text{holomorphic @ } A$$

$$a_m = \frac{m!}{2\pi i} \oint_c \frac{\hat{\varphi}(s)}{s^{m+1}} ds$$



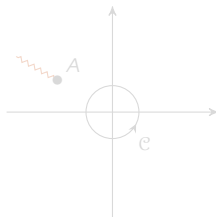
Darboux's result: sketch of the proof

Take a convergent power series: $\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$

Suppose its analytic continuation has a single log singularity at $s = A$:

$$\hat{\varphi}(s) = (\text{holomorphic @ } A) \log(s - A) + \text{holomorphic @ } A$$

$$a_m = \frac{m!}{2\pi i} \oint_c \frac{\hat{\varphi}(s)}{s^{m+1}} ds$$



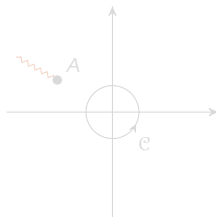
Darboux's result: sketch of the proof

Take a convergent power series: $\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$

Suppose its analytic continuation has a single log singularity at $s = A$:

$$\hat{\varphi}(s) = -\frac{s}{2\pi} \hat{\psi}(s - A) \log(s - A) + \text{holomorphic @ } A$$

$$a_m = \frac{m!}{2\pi i} \oint_c \frac{\hat{\varphi}(s)}{s^{m+1}} ds$$



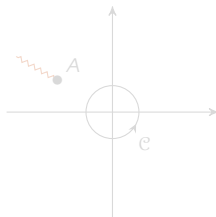
Darboux's result: sketch of the proof

Take a convergent power series: $\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$

Suppose its analytic continuation has a single log **singularity** at $s = A$:

$$\hat{\varphi}(s) = -\frac{s}{2\pi} \hat{\psi}(s-A) \log(s-A) + \text{holomorphic @ } A$$

$$a_m = \frac{m!}{2\pi i} \oint_c \frac{\hat{\varphi}(s)}{s^{m+1}} ds$$



Darboux's result: sketch of the proof

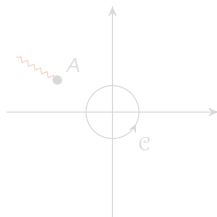
Take a convergent power series: $\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$

Suppose its analytic continuation has a single log **singularity** at $s = A$:

$$\hat{\varphi}(s) = -\frac{S}{2\pi} \hat{\psi}(s - A) \log(s - A) + \text{holomorphic @ } A$$

Stokes constant
 $S \in \mathbb{C}$

$$a_m = \frac{m!}{2\pi i} \oint_c \frac{\hat{\varphi}(s)}{s^{m+1}} ds$$



Darboux's result: sketch of the proof

Take a convergent power series: $\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$

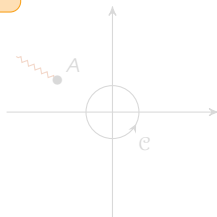
Suppose its analytic continuation has a single log singularity at $s = A$:

$$\hat{\varphi}(s) = -\frac{S}{2\pi} \hat{\psi}(s-A) \log(s-A) + \text{holomorphic @ } A$$

Stokes constant
 $S \in \mathbb{C}$

minor
holom @origin
 $\hat{\psi}(s) = \sum_{k \geq 0} \frac{b_k}{k!} s^k$

$$a_m = \frac{m!}{2\pi i} \oint_c \frac{\hat{\varphi}(s)}{s^{m+1}} ds$$



Darboux's result: sketch of the proof

Take a convergent power series: $\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$

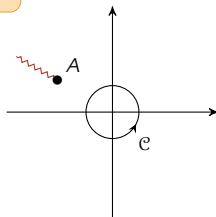
Suppose its analytic continuation has a single log **singularity** at $s = A$:

$$\hat{\varphi}(s) = -\frac{S}{2\pi} \hat{\psi}(s-A) \log(s-A) + \text{holomorphic @ } A$$

Stokes constant
 $S \in \mathbb{C}$

minor
holom @origin
 $\hat{\psi}(s) = \sum_{k \geq 0} \frac{b_k}{k!} s^k$

$$a_m = \frac{m!}{2\pi i} \oint_c \frac{\hat{\varphi}(s)}{s^{m+1}} ds$$



Darboux's result: sketch of the proof

Take a convergent power series: $\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$

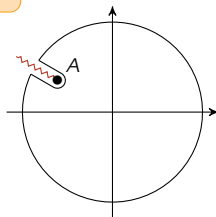
Suppose its analytic continuation has a single log **singularity** at $s = A$:

$$\hat{\varphi}(s) = -\frac{S}{2\pi} \hat{\psi}(s-A) \log(s-A) + \text{holomorphic @ } A$$

Stokes constant
 $S \in \mathbb{C}$

minor
holom @origin
 $\hat{\psi}(s) = \sum_{k \geq 0} \frac{b_k}{k!} s^k$

$$a_m = \frac{m!}{2\pi i} \oint_c \frac{\hat{\varphi}(s)}{s^{m+1}} ds$$



Darboux's result: sketch of the proof

Take a convergent power series: $\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$

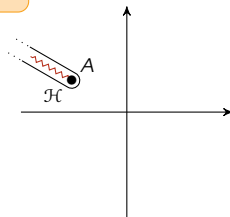
Suppose its analytic continuation has a single log singularity at $s = A$:

$$\hat{\varphi}(s) = -\frac{S}{2\pi} \hat{\psi}(s-A) \log(s-A) + \text{holomorphic @ } A$$

Stokes constant
 $S \in \mathbb{C}$

minor
holom @origin
 $\hat{\psi}(s) = \sum_{k \geq 0} \frac{b_k}{k!} s^k$

$$a_m = \frac{m!}{2\pi i} \int_{\mathcal{H}} \frac{\hat{\varphi}(s)}{s^{m+1}} ds$$



Darboux's result: sketch of the proof

Take a convergent power series: $\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$

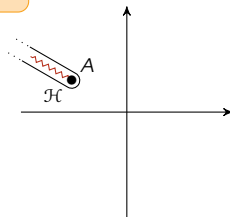
Suppose its analytic continuation has a single log singularity at $s = A$:

$$\hat{\varphi}(s) = -\frac{S}{2\pi} \hat{\psi}(s-A) \log(s-A) + \text{holomorphic @ } A$$

Stokes constant
 $S \in \mathbb{C}$

minor
holom @origin
 $\hat{\psi}(s) = \sum_{k \geq 0} \frac{b_k}{k!} s^k$

$$a_m = -\frac{m!}{2\pi i} \int_{\mathcal{H}} \frac{S}{2\pi} \frac{\hat{\psi}(s-A)}{s^{m+1}} \log(s-A) ds$$



Darboux's result: sketch of the proof

Take a convergent power series: $\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$

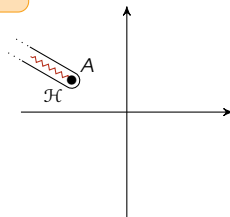
Suppose its analytic continuation has a single log singularity at $s = A$:

$$\hat{\varphi}(s) = -\frac{S}{2\pi} \hat{\psi}(s-A) \log(s-A) + \text{holomorphic @ } A$$

Stokes constant
 $S \in \mathbb{C}$

minor
holom @origin
 $\hat{\psi}(s) = \sum_{k \geq 0} \frac{b_k}{k!} s^k$

$$a_m = \frac{S}{2\pi} m! \int_A^\infty \frac{\hat{\psi}(s-A)}{s^{m+1}} ds$$



Darboux's result: sketch of the proof

Take a convergent power series: $\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$

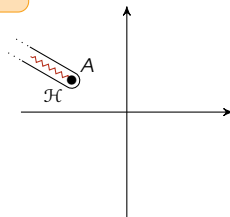
Suppose its analytic continuation has a single log **singularity** at $s = A$:

$$\hat{\varphi}(s) = -\frac{S}{2\pi} \hat{\psi}(s-A) \log(s-A) + \text{holomorphic @ } A$$

Stokes constant
 $S \in \mathbb{C}$

minor
holom @origin
 $\hat{\psi}(s) = \sum_{k \geq 0} \frac{b_k}{k!} s^k$

$$a_m = \frac{S}{2\pi} \frac{\Gamma(m)}{A^m} \left(b_0 + \frac{A}{m-1} b_1 + \dots \right. \\ \left. + \frac{A^k}{(m-1)_k} b_k + O(m^{-k-1}) \right)$$



Darboux meets Borel: summary

- Given: $\tilde{\varphi}(\hbar) = \sum_m a_m \hbar^m$ divergent
- Borel transform: $\hat{\varphi}(s) = \sum_m \frac{a_m}{m!} s^m$ convergent
- Suppose you can compute:
 - ① Log singularities: A_1, \dots, A_n
 - ② Stokes constants: S_1, \dots, S_n
 - ③ Minors: $\hat{\psi}_1, \dots, \hat{\psi}_n$
- Large m asymptotics:

$$a_m = \frac{\Gamma(m)}{2\pi} \left(\frac{S_1}{A_1^m} \left(b_{1,0} + \frac{A_1}{m-1} b_{1,1} + \frac{A_1^2}{(m-1)(m-2)} b_{1,2} + \dots \right) \right. \\ \left. + \dots \right. \\ \left. + \frac{S_n}{A_n^m} \left(b_{n,0} + \frac{A_n}{m-1} b_{n,1} + \frac{A_n^2}{(m-1)(m-2)} b_{n,2} + \dots \right) \right)$$

Darboux meets Borel: summary

- Given: $\tilde{\varphi}(\hbar) = \sum_m a_m \hbar^m$ divergent
- Borel transform: $\hat{\varphi}(s) = \sum_m \frac{a_m}{m!} s^m$ convergent
- Suppose you can compute:
 - 1 Log singularities: A_1, \dots, A_n
 - 2 Stokes constants: S_1, \dots, S_n
 - 3 Minors: $\hat{\psi}_1, \dots, \hat{\psi}_n$
- Large m asymptotics:

$$a_m = \frac{\Gamma(m)}{2\pi} \left(\frac{S_1}{A_1^m} \left(b_{1,0} + \frac{A_1}{m-1} b_{1,1} + \frac{A_1^2}{(m-1)(m-2)} b_{1,2} + \dots \right) \right. \\ \left. + \dots \right. \\ \left. + \frac{S_n}{A_n^m} \left(b_{n,0} + \frac{A_n}{m-1} b_{n,1} + \frac{A_n^2}{(m-1)(m-2)} b_{n,2} + \dots \right) \right)$$

Darboux meets Borel: summary

- Given: $\tilde{\varphi}(\hbar) = \sum_m a_m \hbar^m$ divergent
- Borel transform: $\hat{\varphi}(s) = \sum_m \frac{a_m}{m!} s^m$ convergent
- Suppose you can compute:
 - 1 Log singularities: A_1, \dots, A_n
 - 2 Stokes constants: S_1, \dots, S_n
 - 3 Minors: $\hat{\psi}_1, \dots, \hat{\psi}_n$
- Large m asymptotics:

$$a_m = \frac{\Gamma(m)}{2\pi} \left(\frac{S_1}{A_1^m} \left(b_{1,0} + \frac{A_1}{m-1} b_{1,1} + \frac{A_1^2}{(m-1)(m-2)} b_{1,2} + \dots \right) \right. \\ \left. + \dots \right. \\ \left. + \frac{S_n}{A_n^m} \left(b_{n,0} + \frac{A_n}{m-1} b_{n,1} + \frac{A_n^2}{(m-1)(m-2)} b_{n,2} + \dots \right) \right)$$

Darboux meets Borel: summary

- Given: $\tilde{\varphi}(\hbar) = \sum_m a_m \hbar^m$ divergent
- Borel transform: $\hat{\varphi}(s) = \sum_m \frac{a_m}{m!} s^m$ convergent
- Suppose you can compute:
 - ① Log singularities: A_1, \dots, A_n
 - ② Stokes constants: S_1, \dots, S_n
 - ③ Minors: $\hat{\psi}_1, \dots, \hat{\psi}_n$
- Large m asymptotics:

$$\begin{aligned} a_m = \frac{\Gamma(m)}{2\pi} & \left(\frac{S_1}{A_1^m} \left(b_{1,0} + \frac{A_1}{m-1} b_{1,1} + \frac{A_1^2}{(m-1)(m-2)} b_{1,2} + \dots \right) \right. \\ & + \dots \\ & \left. + \frac{S_n}{A_n^m} \left(b_{n,0} + \frac{A_n}{m-1} b_{n,1} + \frac{A_n^2}{(m-1)(m-2)} b_{n,2} + \dots \right) \right) \end{aligned}$$

Main example: Airy function

The formal (WKB) solutions of the **Airy ODE**, $\left(\hbar^2 \frac{d^2}{dx^2} - x\right) \psi(x, \hbar) = 0$, are

$$\psi_{\text{Ai}}(x; \hbar) = \frac{e^{-\frac{2}{3\hbar}x^{3/2}}}{\sqrt{2}x^{1/4}} \left(1 - \frac{5}{48x^{3/2}}\hbar + \frac{385}{4608x^3}\hbar^2 + \dots\right)$$

$$\psi_{\text{Bi}}(x; \hbar) = \psi_{\text{Ai}}(x; -\hbar)$$

$\tilde{\psi}_{\text{Ai}}$ is a divergent series in \hbar . Its Borel transform has a single log singularity at $s = +\frac{4}{3}x^{3/2}$:

$$\hat{\psi}_{\text{Ai}}(x; s) = -\frac{1}{2\pi} \hat{\psi}_{\text{Bi}}(x; s - \frac{4}{3}x^{3/2}) \log(s - \frac{4}{3}x^{3/2}) + \text{holom}$$

For the Airy function: $(\text{sing}, \text{Stokes cnstrn}, \text{minor}) = (+\frac{4}{3}x^{3/2}, 1, \hat{\psi}_{\text{Bi}})$

Main example: Airy function

The formal (WKB) solutions of the **Airy ODE**, $\left(\hbar^2 \frac{d^2}{dx^2} - x\right) \psi(x, \hbar) = 0$, are

$$\psi_{\text{Ai}}(x; \hbar) = \frac{e^{-\frac{2}{3\hbar}x^{3/2}}}{\sqrt{2}x^{1/4}} \left(1 - \frac{5}{48x^{3/2}}\hbar + \frac{385}{4608x^3}\hbar^2 + \dots\right)$$

$$\psi_{\text{Bi}}(x; \hbar) = \psi_{\text{Ai}}(x; -\hbar)$$

$\tilde{\psi}_{\text{Ai}}$ is a divergent series in \hbar . Its Borel transform has a single log singularity at $s = +\frac{4}{3}x^{3/2}$:

$$\hat{\psi}_{\text{Ai}}(x; s) = -\frac{1}{2\pi} \hat{\psi}_{\text{Bi}}(x; s - \frac{4}{3}x^{3/2}) \log(s - \frac{4}{3}x^{3/2}) + \text{holom}$$

For the Airy function: $(\text{sing}, \text{Stokes cnstn}, \text{minor}) = (+\frac{4}{3}x^{3/2}, 1, \hat{\psi}_{\text{Bi}})$

Main example: Airy function

The formal (WKB) solutions of the **Airy ODE**, $(\hbar^2 \frac{d^2}{dx^2} - x) \psi(x, \hbar) = 0$, are

$$\psi_{\text{Ai}}(x; \hbar) = \frac{e^{-\frac{2}{3\hbar} x^{3/2}}}{\sqrt{2} x^{1/4}} \left(1 - \frac{5}{48 x^{3/2}} \hbar + \frac{385}{4608 x^3} \hbar^2 + \dots \right) = e^{-\frac{2}{3\hbar} x^{3/2}} \tilde{\psi}_{\text{Ai}}(x; \hbar)$$

$$\psi_{\text{Bi}}(x; \hbar) = \psi_{\text{Ai}}(x; -\hbar)$$

$\tilde{\psi}_{\text{Ai}}$ is a divergent series in \hbar . Its Borel transform has a single log singularity at $s = +\frac{4}{3} x^{3/2}$:

$$\hat{\psi}_{\text{Ai}}(x; s) = -\frac{1}{2\pi} \hat{\psi}_{\text{Bi}}(x; s - \frac{4}{3} x^{3/2}) \log(s - \frac{4}{3} x^{3/2}) + \text{holom}$$

For the Airy function: (sing, Stokes cnstrn, minor) = $(+\frac{4}{3} x^{3/2}, 1, \hat{\psi}_{\text{Bi}})$

Main example: Airy function

The formal (WKB) solutions of the **Airy ODE**, $(\hbar^2 \frac{d^2}{dx^2} - x) \psi(x, \hbar) = 0$, are

$$\psi_{\text{Ai}}(x; \hbar) = \frac{e^{-\frac{2}{3\hbar} x^{3/2}}}{\sqrt{2} x^{1/4}} \left(1 - \frac{5}{48 x^{3/2}} \hbar + \frac{385}{4608 x^3} \hbar^2 + \dots \right) = e^{-\frac{2}{3\hbar} x^{3/2}} \tilde{\psi}_{\text{Ai}}(x; \hbar)$$

$$\psi_{\text{Bi}}(x; \hbar) = \psi_{\text{Ai}}(x; -\hbar)$$

$\tilde{\psi}_{\text{Ai}}$ is a divergent series in \hbar . Its Borel transform has a single log singularity at $s = +\frac{4}{3} x^{3/2}$:

$$\hat{\psi}_{\text{Ai}}(x; s) = -\frac{1}{2\pi} \hat{\psi}_{\text{Bi}}(x; s - \frac{4}{3} x^{3/2}) \log(s - \frac{4}{3} x^{3/2}) + \text{holom}$$

For the Airy function: $(\text{sing}, \text{Stokes cnstrn}, \text{minor}) = (+\frac{4}{3} x^{3/2}, 1, \hat{\psi}_{\text{Bi}})$

Main example: Airy function

The formal (WKB) solutions of the **Airy ODE**, $(\hbar^2 \frac{d^2}{dx^2} - x) \psi(x, \hbar) = 0$, are

$$\psi_{\text{Ai}}(x; \hbar) = \frac{e^{-\frac{2}{3\hbar} x^{3/2}}}{\sqrt{2} x^{1/4}} \left(1 - \frac{5}{48 x^{3/2}} \hbar + \frac{385}{4608 x^3} \hbar^2 + \dots \right) = e^{-\frac{2}{3\hbar} x^{3/2}} \tilde{\psi}_{\text{Ai}}(x; \hbar)$$

$$\psi_{\text{Bi}}(x; \hbar) = \psi_{\text{Ai}}(x; -\hbar)$$

$\tilde{\psi}_{\text{Ai}}$ is a divergent series in \hbar . Its Borel transform has a single log singularity at $s = +\frac{4}{3} x^{3/2}$:

$$\hat{\psi}_{\text{Ai}}(x; s) = -\frac{1}{2\pi} \hat{\psi}_{\text{Bi}}(x; s - \frac{4}{3} x^{3/2}) \log(s - \frac{4}{3} x^{3/2}) + \text{holom}$$

Bairy fnct resurges
at the sing of Airy fnct

For the Airy function: (sing, Stokes cnstrn, minor) = $(+\frac{4}{3} x^{3/2}, 1, \hat{\psi}_{\text{Bi}})$

Main example: Airy function

The formal (WKB) solutions of the **Airy ODE**, $(\hbar^2 \frac{d^2}{dx^2} - x) \psi(x, \hbar) = 0$, are

$$\psi_{\text{Ai}}(x; \hbar) = \frac{e^{-\frac{2}{3\hbar} x^{3/2}}}{\sqrt{2} x^{1/4}} \left(1 - \frac{5}{48 x^{3/2}} \hbar + \frac{385}{4608 x^3} \hbar^2 + \dots \right) = e^{-\frac{2}{3\hbar} x^{3/2}} \tilde{\psi}_{\text{Ai}}(x; \hbar)$$

$$\psi_{\text{Bi}}(x; \hbar) = \psi_{\text{Ai}}(x; -\hbar)$$

$\tilde{\psi}_{\text{Ai}}$ is a divergent series in \hbar . Its Borel transform has a single log singularity at $s = +\frac{4}{3} x^{3/2}$:

$$\hat{\psi}_{\text{Ai}}(x; s) = -\frac{1}{2\pi} \hat{\psi}_{\text{Bi}}(x; s - \frac{4}{3} x^{3/2}) \log(s - \frac{4}{3} x^{3/2}) + \text{holom}$$

Bairy fnct resurges
at the sing of Airy fnct

For the Airy function: (sing, Stokes cnstrn, minor) = $(+\frac{4}{3} x^{3/2}, 1, \hat{\psi}_{\text{Bi}})$

Strategy towards large genus asymptotics

Goal

Compute the large genus asymptotics of $\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle$

- 1 Take the generating series

$$W_n(x_1, \dots, x_n; \hbar) = \sum_{g \geq 0} \hbar^{2g-2+n} W_{g,n}(x_1, \dots, x_n)$$

- W_n is a divergent series in \hbar . Take its Borel transform and study its singularity structure
- Get the large genus asymptotics (with subleading contributions!)

Strategy towards large genus asymptotics

Goal

Compute the large genus asymptotics of $\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle$

- 1 Take the generating series

$$W_n(x_1, \dots, x_n; \hbar) = \sum_{g \geq 0} \hbar^{2g-2+n} W_{g,n}(x_1, \dots, x_n)$$

- 2 W_n is a divergent series in \hbar . Take its Borel transform and study its singularity structure
- 3 Get the large genus asymptotics (with subleading contributions!)

Strategy towards large genus asymptotics

Goal

Compute the large genus asymptotics of $\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle$

- 1 Take the generating series

$$\begin{aligned} W_n(x_1, \dots, x_n; \hbar) &= \sum_{g \geq 0} \hbar^{2g-2+n} W_{g,n}(x_1, \dots, x_n) \\ &= (-2)^{-(2g-2+n)} \sum_{d_1, \dots, d_n} \frac{\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle}{2x_1^{d_1+3/2} \cdots 2x_n^{d_n+3/2}} \end{aligned}$$

- 2 W_n is a divergent series in \hbar . Take its Borel transform and study its singularity structure
- 3 Get the large genus asymptotics (with subleading contributions!)

Strategy towards large genus asymptotics

Goal

Compute the large genus asymptotics of $\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle$

- 1 Take the generating series

$$W_n(x_1, \dots, x_n; \hbar) = \sum_{g \geq 0} \hbar^{2g-2+n} W_{g,n}(x_1, \dots, x_n)$$

$$= (-2)^{-(2g-2+n)} \sum_{d_1, \dots, d_n} \frac{\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle}{2x_1^{d_1+3/2} \cdots 2x_n^{d_n+3/2}}$$

n fixed
 $\hbar \rightarrow$ genus
 $x_i \rightarrow d_i$

- 2 W_n is a divergent series in \hbar . Take its Borel transform and study its singularity structure
- 3 Get the large genus asymptotics (with subleading contributions!)

Strategy towards large genus asymptotics

Goal

Compute the large genus asymptotics of $\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle$

- 1 Take the generating series

$$W_n(x_1, \dots, x_n; \hbar) = \sum_{g \geq 0} \hbar^{2g-2+n} W_{g,n}(x_1, \dots, x_n)$$

$$= (-2)^{-(2g-2+n)} \sum_{d_1, \dots, d_n} \frac{\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle}{2x_1^{d_1+3/2} \cdots 2x_n^{d_n+3/2}}$$

n fixed
 $\hbar \rightarrow$ genus
 $x_i \rightarrow d_i$

- 2 W_n is a divergent series in \hbar . Take its Borel transform and study its singularity structure
- 3 Get the large genus asymptotics (with subleading contributions!)

Strategy towards large genus asymptotics

Goal

Compute the large genus asymptotics of $\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle$

- 1 Take the generating series

$$W_n(x_1, \dots, x_n; \hbar) = \sum_{g \geq 0} \hbar^{2g-2+n} W_{g,n}(x_1, \dots, x_n)$$

$$= (-2)^{-(2g-2+n)} \sum_{d_1, \dots, d_n} \frac{\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle}{2x_1^{d_1+3/2} \cdots 2x_n^{d_n+3/2}}$$

n fixed
 $\hbar \rightarrow$ genus
 $x_i \rightarrow d_i$

- 2 W_n is a divergent series in \hbar . Take its Borel transform and study its singularity structure
- 3 Get the large genus asymptotics (with subleading contributions!)

Determinantal formula: setup

Arrange the formal Airy functions as

$$\Psi(x, \hbar) = \begin{pmatrix} \psi_{Ai} & \psi_{Bi} \\ \psi'_{Ai} & \psi'_{Bi} \end{pmatrix} \in \mathbf{SL}(2, \mathbb{C}).$$

It solves the system $(\hbar \frac{d}{dx} - \begin{pmatrix} 0 & 1 \\ x & 0 \end{pmatrix})\Psi = 0$.

Define the matrix $M = \Psi^{-1} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \Psi$:

$$M(x, \hbar) = \begin{pmatrix} \frac{1}{2}(\psi'_{Ai}\psi_{Bi} + \psi_{Ai}\psi'_{Bi}) & \psi_{Ai}\psi_{Bi} \\ \psi'_{Ai}\psi'_{Bi} & -\frac{1}{2}(\psi_{Ai}\psi'_{Bi} + \psi'_{Ai}\psi_{Bi}) \end{pmatrix} \in \mathfrak{sl}(2, \mathbb{C}).$$

Crucial facts:

- M contains only quadratic products Ai-Bi
- The exponential terms cancel out $\implies M$ is a formal series in \hbar .

Determinantal formula: setup

Arrange the formal Airy functions as

$$\Psi(x, \hbar) = \begin{pmatrix} \psi_{Ai} & \psi_{Bi} \\ \psi'_{Ai} & \psi'_{Bi} \end{pmatrix} \in \mathbf{SL}(2, \mathbb{C}).$$

It solves the system $(\hbar \frac{d}{dx} - \begin{pmatrix} 0 & 1 \\ x & 0 \end{pmatrix})\Psi = 0$.

Define the matrix $M = \Psi \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \Psi^{-1}$:

$$M(x, \hbar) = \begin{pmatrix} \frac{1}{2}(\psi'_{Ai}\psi_{Bi} + \psi_{Ai}\psi'_{Bi}) & \psi_{Ai}\psi_{Bi} \\ \psi'_{Ai}\psi'_{Bi} & -\frac{1}{2}(\psi_{Ai}\psi'_{Bi} + \psi'_{Ai}\psi_{Bi}) \end{pmatrix} \in \mathfrak{sl}(2, \mathbb{C}).$$

Crucial facts:

- M contains only quadratic products Ai-Bi
- The exponential terms cancel out $\implies M$ is a formal series in \hbar .

Determinantal formula: setup

Arrange the formal Airy functions as

$$\Psi(x, \hbar) = \begin{pmatrix} \psi_{Ai} & \psi_{Bi} \\ \psi'_{Ai} & \psi'_{Bi} \end{pmatrix} \in \mathbf{SL}(2, \mathbb{C}).$$

It solves the system $(\hbar \frac{d}{dx} - \begin{pmatrix} 0 & 1 \\ x & 0 \end{pmatrix})\Psi = 0$.

Define the matrix $M = \Psi \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \Psi^{-1}$:

$$M(x, \hbar) = \begin{pmatrix} \frac{1}{2}(\psi'_{Ai}\psi_{Bi} + \psi_{Ai}\psi'_{Bi}) & \psi_{Ai}\psi_{Bi} \\ \psi'_{Ai}\psi'_{Bi} & -\frac{1}{2}(\psi_{Ai}\psi'_{Bi} + \psi'_{Ai}\psi_{Bi}) \end{pmatrix} \in \mathfrak{sl}(2, \mathbb{C}).$$

Crucial facts:

- M contains only quadratic products Ai-Bi
- The exponential terms cancel out $\implies M$ is a formal series in \hbar .

Determinantal formula: setup

Arrange the formal Airy functions as

$$\Psi(x, \hbar) = \begin{pmatrix} \psi_{Ai} & \psi_{Bi} \\ \psi'_{Ai} & \psi'_{Bi} \end{pmatrix} \in \mathbf{SL}(2, \mathbb{C}).$$

It solves the system $(\hbar \frac{d}{dx} - \begin{pmatrix} 0 & 1 \\ x & 0 \end{pmatrix})\Psi = 0$.

Define the matrix $M = \Psi \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \Psi^{-1}$:

$$M(x, \hbar) = \begin{pmatrix} \frac{1}{2}(\tilde{\psi}'_{Ai}\tilde{\psi}_{Bi} + \tilde{\psi}_{Ai}\tilde{\psi}'_{Bi}) & \tilde{\psi}_{Ai}\tilde{\psi}_{Bi} \\ \tilde{\psi}'_{Ai}\tilde{\psi}'_{Bi} & -\frac{1}{2}(\tilde{\psi}_{Ai}\tilde{\psi}'_{Bi} + \tilde{\psi}'_{Ai}\tilde{\psi}_{Bi}) \end{pmatrix} \in \mathfrak{sl}(2, \mathbb{C}).$$

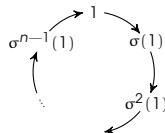
Crucial facts:

- M contains only quadratic products Ai-Bi
- The exponential terms cancel out $\implies M$ is a formal series in \hbar .

Determinantal formula: result

Cyclic permutations:

$$C_n = \{ \sigma \in S_n \mid \sigma \text{ has cycle type of length } n \}$$



Determinantal formula (Eynard et al., Bertola–Dubrovin–Yang):

$$W_n(x_1, \dots, x_n; \hbar) = \sum_{\sigma \in C_n} \text{sgn}(\sigma) \frac{\text{tr}(M(x_1, \hbar) M(x_{\sigma(1)}, \hbar) \cdots M(x_{\sigma^{n-1}(1)}, \hbar))}{(x_1 - x_{\sigma(1)})(x_2 - x_{\sigma(2)}) \cdots (x_n - x_{\sigma(n)})}$$

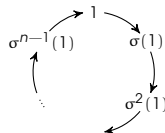
Example: $n = 2$

$$\begin{aligned} W_2 &= - \frac{\text{tr}(M(x_1, \hbar) M(x_2, \hbar))}{(x_1 - x_2)(x_2 - x_1)} \\ &= \frac{\tilde{\psi}_{\text{Ai},1} \tilde{\psi}_{\text{Bi},1} \tilde{\psi}'_{\text{Ai},2} \tilde{\psi}'_{\text{Bi},2} + \frac{1}{2} \tilde{\psi}_{\text{Ai},1} \tilde{\psi}'_{\text{Bi},1} \tilde{\psi}_{\text{Ai},2} \tilde{\psi}'_{\text{Bi},2} + \frac{1}{2} \tilde{\psi}_{\text{Ai},1} \tilde{\psi}'_{\text{Bi},1} \tilde{\psi}_{\text{Bi},2} \tilde{\psi}'_{\text{Ai},2}}{(x_1 - x_2)^2} + (x_1 \leftrightarrow x_2) \end{aligned}$$

Determinantal formula: result

Cyclic permutations:

$$C_n = \{ \sigma \in S_n \mid \sigma \text{ has cycle type of length } n \}$$



Determinantal formula (Eynard et al., Bertola–Dubrovin–Yang):

$$W_n(x_1, \dots, x_n; \hbar) = \sum_{\sigma \in C_n} \text{sgn}(\sigma) \frac{\text{tr}(M(x_1, \hbar) M(x_{\sigma(1)}, \hbar) \cdots M(x_{\sigma^{n-1}(1)}, \hbar))}{(x_1 - x_{\sigma(1)})(x_2 - x_{\sigma(2)}) \cdots (x_n - x_{\sigma(n)})}$$

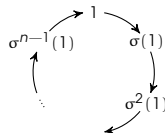
Example: $n = 2$

$$\begin{aligned} W_2 &= - \frac{\text{tr}(M(x_1, \hbar) M(x_2, \hbar))}{(x_1 - x_2)(x_2 - x_1)} \\ &= \frac{\tilde{\psi}_{Ai,1} \tilde{\psi}_{Bi,1} \tilde{\psi}'_{Ai,2} \tilde{\psi}'_{Bi,2} + \frac{1}{2} \tilde{\psi}_{Ai,1} \tilde{\psi}'_{Bi,1} \tilde{\psi}_{Ai,2} \tilde{\psi}'_{Bi,2} + \frac{1}{2} \tilde{\psi}_{Ai,1} \tilde{\psi}'_{Bi,1} \tilde{\psi}_{Bi,2} \tilde{\psi}'_{Ai,2}}{(x_1 - x_2)^2} + (x_1 \leftrightarrow x_2) \end{aligned}$$

Determinantal formula: result

Cyclic permutations:

$$C_n = \{ \sigma \in S_n \mid \sigma \text{ has cycle type of length } n \}$$



Determinantal formula (Eynard et al., Bertola–Dubrovin–Yang):

$$W_n(x_1, \dots, x_n; \hbar) = \sum_{\sigma \in C_n} \text{sgn}(\sigma) \frac{\text{tr}(M(x_1, \hbar) M(x_{\sigma(1)}, \hbar) \cdots M(x_{\sigma^{n-1}(1)}, \hbar))}{(x_1 - x_{\sigma(1)})(x_2 - x_{\sigma(2)}) \cdots (x_n - x_{\sigma(n)})}$$

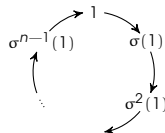
Example: $n = 2$

$$\begin{aligned} W_2 &= - \frac{\text{tr}(M(x_1, \hbar) M(x_2, \hbar))}{(x_1 - x_2)(x_2 - x_1)} \\ &= \frac{\tilde{\psi}_{Ai,1} \tilde{\psi}_{Bi,1} \tilde{\psi}'_{Ai,2} \tilde{\psi}'_{Bi,2} + \frac{1}{2} \tilde{\psi}_{Ai,1} \tilde{\psi}'_{Bi,1} \tilde{\psi}_{Ai,2} \tilde{\psi}'_{Bi,2} + \frac{1}{2} \tilde{\psi}_{Ai,1} \tilde{\psi}'_{Bi,1} \tilde{\psi}_{Bi,2} \tilde{\psi}'_{Ai,2}}{(x_1 - x_2)^2} + (x_1 \leftrightarrow x_2) \end{aligned}$$

Determinantal formula: result

Cyclic permutations:

$$C_n = \{ \sigma \in S_n \mid \sigma \text{ has cycle type of length } n \}$$



Determinantal formula (Eynard et al., Bertola–Dubrovin–Yang):

$$W_n(x_1, \dots, x_n; \hbar) = \sum_{\sigma \in C_n} \text{sgn}(\sigma) \frac{\text{tr}(M(x_1, \hbar) M(x_{\sigma(1)}, \hbar) \cdots M(x_{\sigma^{n-1}(1)}, \hbar))}{(x_1 - x_{\sigma(1)})(x_2 - x_{\sigma(2)}) \cdots (x_n - x_{\sigma(n)})}$$

Example: $n = 2$

$$\begin{aligned} W_2 &= - \frac{\text{tr}(M(x_1, \hbar) M(x_2, \hbar))}{(x_1 - x_2)(x_2 - x_1)} \\ &= \frac{\tilde{\psi}_{Ai,1} \tilde{\psi}_{Bi,1} \tilde{\psi}'_{Ai,2} \tilde{\psi}'_{Bi,2} + \frac{1}{2} \tilde{\psi}_{Ai,1} \tilde{\psi}'_{Bi,1} \tilde{\psi}_{Ai,2} \tilde{\psi}'_{Bi,2} + \frac{1}{2} \tilde{\psi}_{Ai,1} \tilde{\psi}'_{Bi,1} \tilde{\psi}_{Bi,2} \tilde{\psi}'_{Ai,2}}{(x_1 - x_2)^2} + (x_1 \leftrightarrow x_2) \end{aligned}$$

Singularity structure of \widehat{W}_n Singularity strct
of $\widehat{\psi}_{Ai}, \widehat{\psi}_{Bi}$ Singularity strct
of \widehat{W}_n

- $2n$ log singularities of \widehat{W}_n , located at

$$+ \frac{4}{3}x_i^{3/2} \quad \text{and} \quad - \frac{4}{3}x_i^{3/2}, \quad i = 1, \dots, n$$

- Stokes constants: $S = 1$

- Minors:

Ⓐ at $+\frac{4}{3}x_i^{3/2}$: replace each $\widehat{\psi}_{Ai}(x_i; s)$ with $\widehat{\psi}_{Bi}(x_i; s)$

Ⓑ at $-\frac{4}{3}x_i^{3/2}$: replace each $\widehat{\psi}_{Bi}(x_i; s)$ with $\widehat{\psi}_{Ai}(x_i; s)$

Singularity structure of \widehat{W}_n Singularity strct
of $\widehat{\psi}_{Ai}, \widehat{\psi}_{Bi}$ Singularity strct
of \widehat{W}_n

- $2n \log$ **singularities** of \widehat{W}_n , located at

$$+ \frac{4}{3}x_i^{3/2} \quad \text{and} \quad - \frac{4}{3}x_i^{3/2}, \quad i = 1, \dots, n$$

- **Stokes constants:** $S = 1$

- **Minors:**

Ⓐ at $+\frac{4}{3}x_i^{3/2}$: replace each $\widehat{\psi}_{Ai}(x_i; s)$ with $\widehat{\psi}_{Bi}(x_i; s)$

Ⓑ at $-\frac{4}{3}x_i^{3/2}$: replace each $\widehat{\psi}_{Bi}(x_i; s)$ with $\widehat{\psi}_{Ai}(x_i; s)$

Singularity structure of \widehat{W}_n Singularity strct
of $\widehat{\psi}_{Ai}, \widehat{\psi}_{Bi}$ Singularity strct
of \widehat{W}_n

- $2n \log$ **singularities** of \widehat{W}_n , located at

$$+ \frac{4}{3}x_i^{3/2} \quad \text{and} \quad - \frac{4}{3}x_i^{3/2}, \quad i = 1, \dots, n$$

- **Stokes constants:** $S = 1$

- **Minors:**

Ⓐ at $+\frac{4}{3}x_i^{3/2}$: replace each $\widehat{\psi}_{Ai}(x_i; s)$ with $\widehat{\psi}_{Bi}(x_i; s)$

Ⓑ at $-\frac{4}{3}x_i^{3/2}$: replace each $\widehat{\psi}_{Bi}(x_i; s)$ with $\widehat{\psi}_{Ai}(x_i; s)$

Singularity structure of \widehat{W}_n Singularity strct
of $\widehat{\psi}_{Ai}, \widehat{\psi}_{Bi}$ Singularity strct
of \widehat{W}_n

- $2n \log$ **singularities** of \widehat{W}_n , located at

$$+ \frac{4}{3}x_i^{3/2} \quad \text{and} \quad - \frac{4}{3}x_i^{3/2}, \quad i = 1, \dots, n$$

- **Stokes constants:** $S = 1$

- **Minors:**

Ⓐ at $+\frac{4}{3}x_i^{3/2}$: replace each $\widehat{\psi}_{Ai}(x_i; s)$ with $\widehat{\psi}_{Bi}(x_i; s)$

Ⓑ at $-\frac{4}{3}x_i^{3/2}$: replace each $\widehat{\psi}_{Bi}(x_i; s)$ with $\widehat{\psi}_{Ai}(x_i; s)$

Summary

Uniformly in d_1, \dots, d_n as $g \rightarrow \infty$:

$$\begin{aligned} \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = S \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{A^{2g-2+n}} & \left(1 + \frac{A}{2g-3+n} \alpha_1 + \cdots \right. \\ & \left. + \frac{A^k}{(2g-3+n)_k} \alpha_k + O(g^{-k-1}) \right) \end{aligned}$$

where:

- $S = 1$
Stokes constants of the **Airy** ODE
- $A = 2/3$
leading exp behaviour of Ai
- α_k polynomials in n and multiplicities of d_i (conj by Guo–Yang)
are computable from the asymptotic expansion coeffs of Ai

Summary

Uniformly in d_1, \dots, d_n as $g \rightarrow \infty$:

$$\begin{aligned} \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = S \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{A^{2g-2+n}} & \left(1 + \frac{A}{2g-3+n} \alpha_1 + \cdots \right. \\ & \left. + \frac{A^k}{(2g-3+n)_k} \alpha_k + O(g^{-k-1}) \right) \end{aligned}$$

where:

- $S = 1$

Stokes constants of the **Airy ODE**

- $A = 2/3$

leading exp behaviour of Ai

- α_k polynomials in n and multiplicities of d_i (conj by Guo–Yang)

are computable from the asymptotic expansion coeffs of Ai

Summary

Uniformly in d_1, \dots, d_n as $g \rightarrow \infty$:

$$\begin{aligned} \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = S \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{A^{2g-2+n}} & \left(1 + \frac{A}{2g-3+n} \alpha_1 + \cdots \right. \\ & \left. + \frac{A^k}{(2g-3+n)_k} \alpha_k + O(g^{-k-1}) \right) \end{aligned}$$

where:

- $S = 1$

Stokes constants of the **Airy ODE**

- $A = 2/3$

leading exp behaviour of Ai

- α_k polynomials in n and multiplicities of d_i (conj by Guo–Yang)

are computable from the asymptotic expansion coeffs of Ai

Summary

Uniformly in d_1, \dots, d_n as $g \rightarrow \infty$:

$$\begin{aligned} \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle = S \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{A^{2g-2+n}} & \left(1 + \frac{A}{2g-3+n} \alpha_1 + \cdots \right. \\ & \left. + \frac{A^k}{(2g-3+n)_k} \alpha_k + O(g^{-k-1}) \right) \end{aligned}$$

where:

- $S = 1$

Stokes constants of the **Airy ODE**

- $A = 2/3$

leading exp behaviour of Ai

- α_k polynomials in n and multiplicities of d_i (conj by Guo–Yang)

are computable from the asymptotic expansion coeffs of Ai

Bessel

Norbury's intersection numbers (super WP volumes, BGW tau function):

$$\begin{aligned}
 \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle^\Theta &= \int_{\overline{\mathcal{M}}_{g,n}} \Theta_{g,n} \prod_{i=1}^n \psi_i^{d_i} (2d_i + 1)!! \\
 &= S \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{A^{2g-2+n}} \left(1 + \frac{A}{2g-3+n} \alpha_1 + \cdots \right. \\
 &\quad \left. + \frac{A^k}{(2g-3+n)_k} \alpha_k + O(g^{-k-1}) \right)
 \end{aligned}$$

where:

- $S = 2$
Stokes constants of the Bessel ODE
- $A = 2$
leading exp behaviour of K_0
- α_k polynomials in n and multiplicities of d_i
are computable from the asymptotic expansion coeffs of K_0

Bessel

Norbury's intersection numbers (super WP volumes, BGW tau function):

$$\begin{aligned}
 \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle^\Theta &= \int_{\overline{\mathcal{M}}_{g,n}} \Theta_{g,n} \prod_{i=1}^n \psi_i^{d_i} (2d_i + 1)!! \\
 &= S \frac{2^n}{4\pi} \frac{\Gamma(2g - 2 + n)}{A^{2g-2+n}} \left(1 + \frac{A}{2g - 3 + n} \alpha_1 + \cdots \right. \\
 &\quad \left. + \frac{A^k}{(2g - 3 + n)_k} \alpha_k + O(g^{-k-1}) \right)
 \end{aligned}$$

where:

- $S = 2$

Stokes constants of the Bessel ODE

- $A = 2$

leading exp behaviour of K_0

- α_k polynomials in n and multiplicities of d_i

are computable from the asymptotic expansion coeffs of K_0

Bessel

Norbury's intersection numbers (super WP volumes, BGW tau function):

$$\begin{aligned}
 \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle^\Theta &= \int_{\overline{\mathcal{M}}_{g,n}} \Theta_{g,n} \prod_{i=1}^n \psi_i^{d_i} (2d_i + 1)!! \\
 &= S \frac{2^n}{4\pi} \frac{\Gamma(2g - 2 + n)}{A^{2g-2+n}} \left(1 + \frac{A}{2g - 3 + n} \alpha_1 + \cdots \right. \\
 &\quad \left. + \frac{A^k}{(2g - 3 + n)_k} \alpha_k + O(g^{-k-1}) \right)
 \end{aligned}$$

where:

- $S = 2$

Stokes constants of the **Bessel ODE**

- $A = 2$

leading exp behaviour of K_0

- α_k polynomials in n and multiplicities of d_i

are computable from the asymptotic expansion coeffs of K_0

r-Airy

Witten's *r*-spin intersection numbers (*r*-KdV tau function):

$$\begin{aligned}
\langle\langle \tau_{a_1, d_1} \cdots \tau_{a_n, d_n} \rangle\rangle^{r\text{-spin}} &= \int_{\overline{\mathcal{M}}_{g,n}} c_w(a_1, \dots, a_n) \prod_{i=1}^n \psi_i^{d_i} (rd_i + a_i)!_{(r)} \\
&= \frac{2^n}{2\pi} \frac{\Gamma(2g-2+n)}{rg^{-1}-|d|} \left[\frac{S_{r,1}}{|A_{r,1}|^{2g-2+n}} \left(\alpha_0^{(r,1)} + \frac{|A_{r,1}|}{2g-3+n} \alpha_1^{(r,1)} + \cdots \right) \right. \\
&\quad + \cdots \\
&\quad + \frac{S_{r, \lfloor \frac{r-1}{2} \rfloor}}{|A_{r, \lfloor \frac{r-1}{2} \rfloor}|^{2g-2+n}} \left(\alpha_0^{(r, \lfloor \frac{r-1}{2} \rfloor)} + \frac{|A_{r, \lfloor \frac{r-1}{2} \rfloor}|^K}{2g-3+n} \alpha_1^{(r, \lfloor \frac{r-1}{2} \rfloor)} + \cdots \right) \\
&\quad \left. + \frac{\delta_r^{\text{even}}}{2} \frac{S_{r, \frac{r}{2}}}{|A_{r, \frac{r}{2}}|^{2g-2+n}} \left(\alpha_0^{(r, \frac{r}{2})} + \frac{|A_{r, \frac{r}{2}}|^K}{2g-3+n} \alpha_1^{(r, \frac{r}{2})} + \cdots \right) \right]
\end{aligned}$$

where $S_{r,i}$, $A_{r,i}$, $\alpha_k^{(r,i)}$ are obtained the *r*-Airy ODE.

r -Airy

Witten's r -spin intersection numbers (r -KdV tau function):

$$\begin{aligned}
\langle\langle \tau_{a_1, d_1} \cdots \tau_{a_n, d_n} \rangle\rangle^{r\text{-spin}} &= \int_{\overline{\mathcal{M}}_{g,n}} c_w(a_1, \dots, a_n) \prod_{i=1}^n \psi_i^{d_i} (rd_i + a_i)!_{(r)} \\
&= \frac{2^n}{2\pi} \frac{\Gamma(2g-2+n)}{rg^{-1}-|d|} \left[\frac{S_{r,1}}{|A_{r,1}|^{2g-2+n}} \left(\alpha_0^{(r,1)} + \frac{|A_{r,1}|}{2g-3+n} \alpha_1^{(r,1)} + \cdots \right) \right. \\
&\quad + \cdots \\
&\quad + \frac{S_{r, \lfloor \frac{r-1}{2} \rfloor}}{|A_{r, \lfloor \frac{r-1}{2} \rfloor}|^{2g-2+n}} \left(\alpha_0^{(r, \lfloor \frac{r-1}{2} \rfloor)} + \frac{|A_{r, \lfloor \frac{r-1}{2} \rfloor}|^K}{2g-3+n} \alpha_1^{(r, \lfloor \frac{r-1}{2} \rfloor)} + \cdots \right) \\
&\quad \left. + \frac{\delta_r^{\text{even}}}{2} \frac{S_{r, \frac{r}{2}}}{|A_{r, \frac{r}{2}}|^{2g-2+n}} \left(\alpha_0^{(r, \frac{r}{2})} + \frac{|A_{r, \frac{r}{2}}|^K}{2g-3+n} \alpha_1^{(r, \frac{r}{2})} + \cdots \right) \right]
\end{aligned}$$

where $S_{r,i}$, $A_{r,i}$, $\alpha_k^{(r,i)}$ are obtained the r -Airy ODE.

r -Airy

Witten's r -spin intersection numbers (r -KdV tau function):

$$\begin{aligned}
\langle\langle \tau_{a_1, d_1} \cdots \tau_{a_n, d_n} \rangle\rangle^{r\text{-spin}} &= \int_{\overline{\mathcal{M}}_{g,n}} c_w(a_1, \dots, a_n) \prod_{i=1}^n \psi_i^{d_i} (rd_i + a_i)!_{(r)} \\
&= \frac{2^n}{2\pi} \frac{\Gamma(2g-2+n)}{r^{g-1-|d|}} \left[\frac{S_{r,1}}{|A_{r,1}|^{2g-2+n}} \left(\alpha_0^{(r,1)} + \frac{|A_{r,1}|}{2g-3+n} \alpha_1^{(r,1)} + \cdots \right) \right. \\
&\quad + \cdots \\
&\quad + \frac{S_{r, \lfloor \frac{r-1}{2} \rfloor}}{|A_{r, \lfloor \frac{r-1}{2} \rfloor}|^{2g-2+n}} \left(\alpha_0^{(r, \lfloor \frac{r-1}{2} \rfloor)} + \frac{|A_{r, \lfloor \frac{r-1}{2} \rfloor}|^k}{2g-3+n} \alpha_1^{(r, \lfloor \frac{r-1}{2} \rfloor)} + \cdots \right) \\
&\quad \left. + \frac{\delta_r^{\text{even}}}{2} \frac{S_{r, \frac{r}{2}}}{|A_{r, \frac{r}{2}}|^{2g-2+n}} \left(\alpha_0^{(r, \frac{r}{2})} + \frac{|A_{r, \frac{r}{2}}|^k}{2g-3+n} \alpha_1^{(r, \frac{r}{2})} + \cdots \right) \right]
\end{aligned}$$

where $S_{r,i}$, $A_{r,i}$, $\alpha_k^{(r,i)}$ are obtained the r -Airy ODE.

谢谢你的注意力!