

Université de Bordeaux

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# Resurgent large genus asymptotics of intersection numbers

j/w B. Eynard, E. Garcia-Failde, P. Gregori, D. Lewański

arXiv: [AG/2309.03143](#)

Alessandro Giacchetto

ETH Zürich

# A case study: $m!$

Counting problem:  $c_m = \# \left\{ \begin{array}{l} \text{arrangements of } m \text{ distinct objects} \\ \text{into } m \text{ distinct boxes} \end{array} \right\}$

Solution:  $c_m = m! = \begin{cases} m \cdot c_{m-1} & m > 1 \\ 1 & m = 1 \end{cases}$

Pro: exact

Con: recursive

Asymptotics:  $c_m = \sqrt{2\pi m} \left(\frac{m}{e}\right)^m \left(1 + O(m^{-1})\right)$

Con: asymptotically exact

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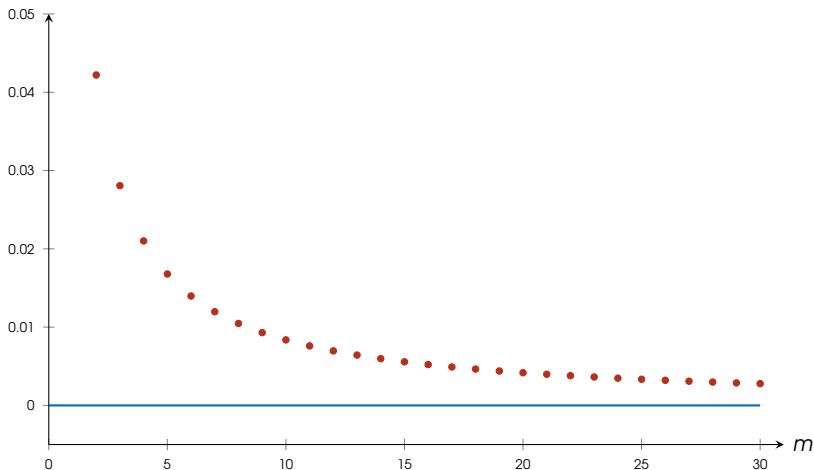
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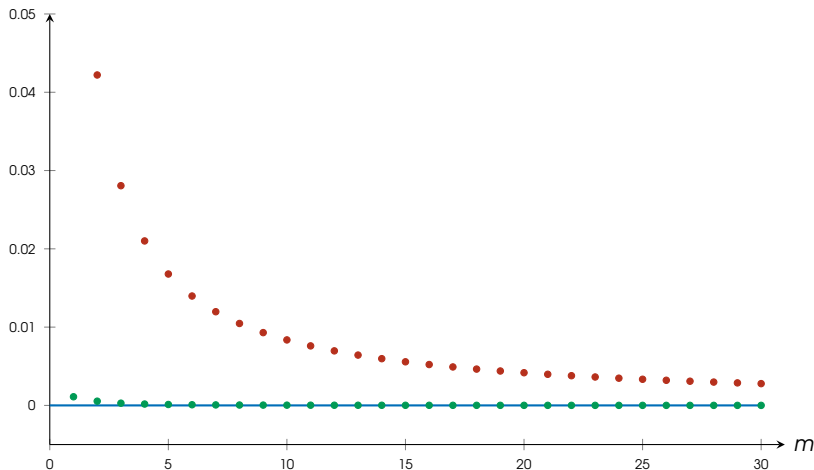
# Visualising Stirling's formula

$$\frac{m!}{\sqrt{2\pi m} \left(\frac{m}{e}\right)^m} - 1 = \mathcal{O}(m^{-1})$$



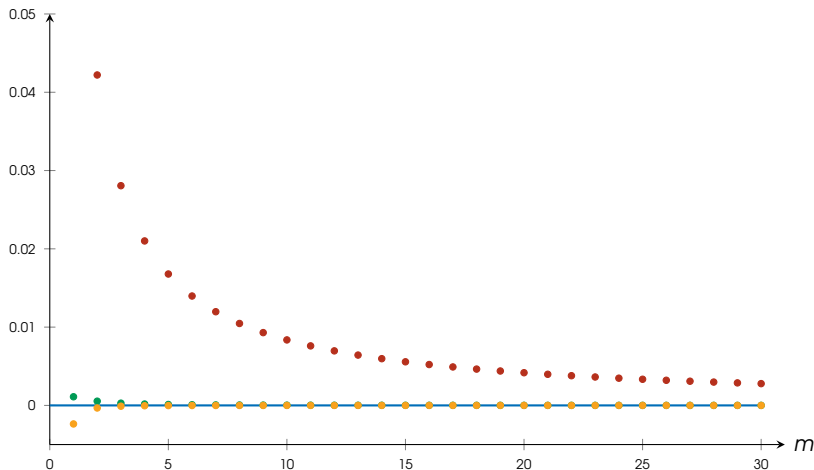
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# $\psi$ -class intersection numbers

$$\langle \tau_{d_1} \cdots \tau_{d_n} \rangle_g = \int_{\overline{\mathcal{M}}_{g,n}} \psi_1^{d_1} \cdots \psi_n^{d_n} \quad d_i \geq 0, \quad d_1 + \cdots + d_n = 3g - 3 + n$$

- Volumes of moduli spaces of **metric ribbon graphs** (maps)

$$V_{g,n}(L_1, \dots, L_n) = \sum_{d_1 + \cdots + d_n = 3g - 3 + n} \langle \tau_{d_1} \cdots \tau_{d_n} \rangle_g \prod_{i=1}^n \frac{L_i^{2d_i}}{2^{d_i} d_i!}$$

- Building block for all **tautological intersection numbers**:
  - Weil–Petersson volumes
  - Masur–Veech volumes
  - Hurwitz numbers
  - ...
- Compute the perturbative expansion of **topological 2d gravity**

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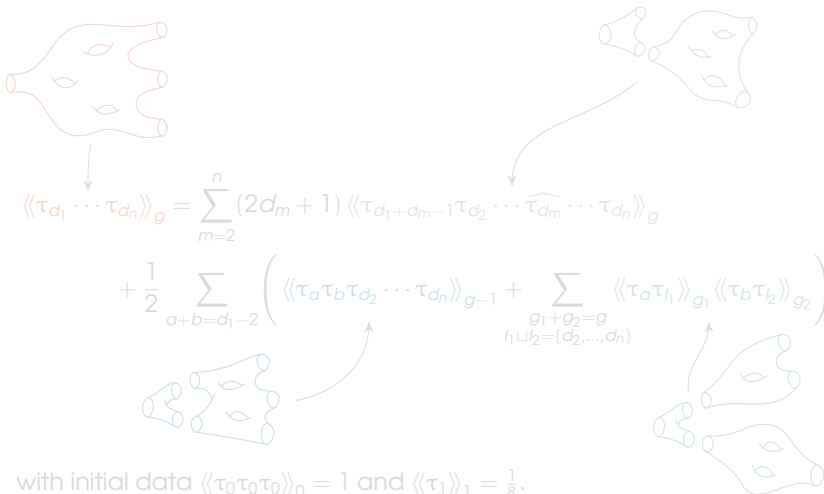
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## Solution

Normalisation:  $\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle_g = \langle \tau_{d_1} \cdots \tau_{d_n} \rangle_g \prod_{i=1}^n (2d_i + 1)!!$

Witten conjecture/Kontsevich theorem, early '90s:



$$\begin{aligned} \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle_g &= \sum_{m=2}^n (2d_m + 1) \langle\langle \tau_{d_1+d_m-1} \tau_{d_2} \cdots \widehat{\tau_{d_m}} \cdots \tau_{d_n} \rangle\rangle_g \\ &+ \frac{1}{2} \sum_{a+b=d_1-2} \left( \langle\langle \tau_a \tau_b \tau_{d_2} \cdots \tau_{d_n} \rangle\rangle_{g-1} + \sum_{\substack{g_1+g_2=g \\ I_1 \sqcup I_2 = \{d_2, \dots, d_n\}}} \langle\langle \tau_a \tau_{I_1} \rangle\rangle_{g_1} \langle\langle \tau_b \tau_{I_2} \rangle\rangle_{g_2} \right) \end{aligned}$$

with initial data  $\langle\langle \tau_0 \tau_0 \tau_0 \rangle\rangle_0 = 1$  and  $\langle\langle \tau_1 \rangle\rangle_1 = \frac{1}{8}$ .



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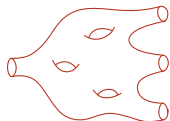

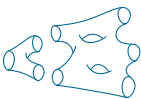
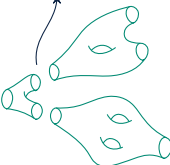



Diagram illustrating the Witten-Kontsevich theorem, showing a genus-2 surface with four boundary components (red) and its decomposition into a genus-2 surface with four boundary components (purple) and a genus-2 surface with four boundary components (blue).

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# Large genus asymptotics

Uniformly in  $d_1, \dots, d_n$  as  $g \rightarrow \infty$ :

$$\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle_g = \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{\left(\frac{2}{3}\right)^{2g-2+n}} \left(1 + o(g^{-1})\right)$$

- Conjectured by [Delecroix–Goujard–Zograf–Zorich](#), 2019
- Proved by [Aggarwal](#), 2020  
(combinatorial/probabilistic analysis of Witten–Kontsevich topological recursion)
- Proved by [Guo–Yang](#), 2021  
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## Questions

- Universal strategy, adaptable to different problems?
- ‘Geometric’ meaning of the formula?
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# Large genus asymptotics: new perspective

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- Universal strategy: resurgent analysis of the determinantal formula
- Geometric meaning: Airy functions

$$y^2 - x = 0 \quad \xrightarrow{\text{quantisation}} \quad \left( \hbar^2 \frac{d^2}{dx^2} - x \right) \psi(x, \hbar) = 0$$

- Subleading corrections: algorithm + properties

Set  $(x)_k = x(x-1) \cdots (x-k+1)$ .

$$\begin{aligned} \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle_g = S \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{A^{2g-2+n}} & \left( 1 + \frac{A}{2g-3+n} \alpha_1 + \cdots \right. \\ & \left. + \frac{A^k}{(2g-3+n)_k} \alpha_k + O(g^{-k-1}) \right) \end{aligned}$$

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$$A = 2/3$$

$$\psi \sim \frac{1}{\sqrt{2x^{1/4}}} e^{\pm \frac{A}{\hbar} x^{-3/2}}$$



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Computable; polynomial in  $n$  and multiplicities of  $d_i$

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$$\alpha_1 = -\frac{17-15n+3n^2}{12} - \frac{(3-n)(n-p_0)}{2} - \frac{(n-p_0)2}{4}$$

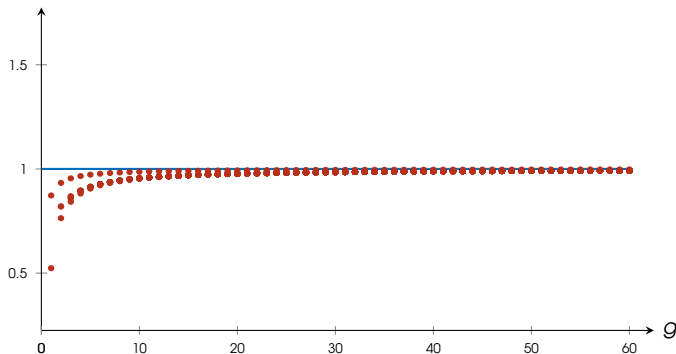
where  $p_0 = \#\{d_i = 0\}$

$$\frac{A^k}{(2g-3+n)_k} \alpha_k + O(g^{-k-1})$$

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$$\frac{\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle_g}{\frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{(2/3)^{2g-2+n}}} = 1 + O(g^{-1})$$

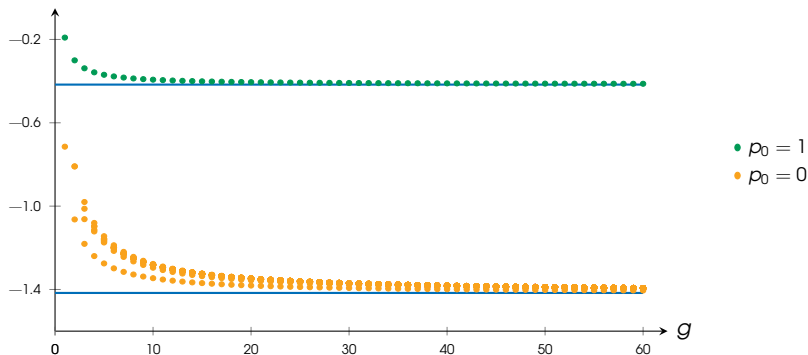
For  $n = 2$ :



# Visualising the large genus asymptotics

$$\frac{2g-3+n}{2/3} \left( \frac{\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle_g}{\frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{(2/3)^{2g-2+n}}} - 1 \right) = \alpha_1(n, p_0) + O(g^{-1})$$

For  $n = 2$ :



# Borel meets Darboux

Borel's idea:

- Divergent power series:

$$\tilde{\varphi}(\hbar) = \sum_{m \geq 0} a_m \hbar^m$$

with  $|a_m| = O(R^{-m} m!)$ .

- The **Borel transform**

$$\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$$

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- Get a holomorphic function around the origin, take analytic continuation
- The large  $m$  **asymptotics** of  $a_m$  is totally controlled by the behaviour of  $\hat{\varphi}$  at its **singularities**

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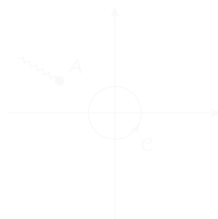
# Darboux's result: sketch of the proof

Take an abs. convergent power series:  $\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$

Suppose its analytic continuation has a single log singularity at  $s = A$ :

$$\hat{\varphi}(s) = (\text{holomorphic @ } A) \log(s - A) + \text{holomorphic @ } A$$

$$a_m = \frac{m!}{2\pi i} \oint_c \frac{\hat{\varphi}(s)}{s^{m+1}} ds$$



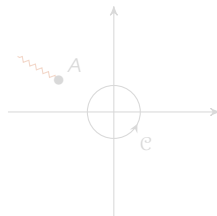
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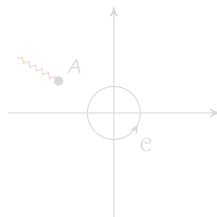
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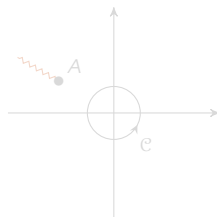
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Take an abs. convergent power series:  $\hat{\varphi}(s) = \sum_{m \geq 0} \frac{a_m}{m!} s^m$

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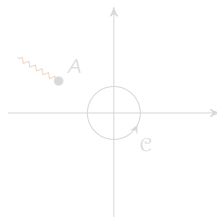
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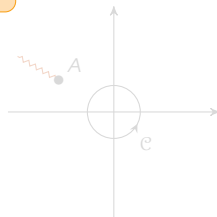
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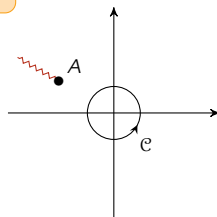
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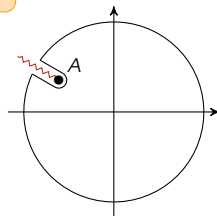
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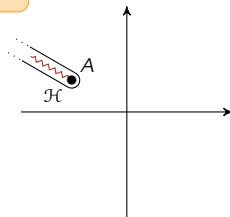
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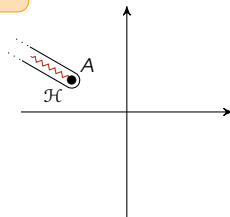
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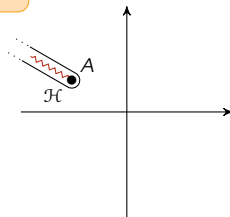
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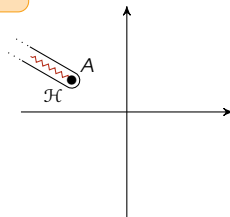
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# Borel meets Darboux: the algorithm

- Given:  $\tilde{\varphi}(\hbar) = \sum_m a_m \hbar^m$  divergent
- Borel transform:  $\hat{\varphi}(s) = \sum_m \frac{a_m}{m!} s^m$  abs. convergent
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- Large  $m$  **asymptotics**:

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# Properties of the resurgence method

- **Algorithmic.**

$$\tilde{\varphi} = \sum_m a_m \hbar^m \longrightarrow (S_A, \hat{\psi}_A)_{A \in \text{Sing}(\hat{\varphi})} \longrightarrow \text{asymptotic of } a_m$$

- **Exponential integrals.** The singularity structure of exponential integrals is well-understood:

$$\tilde{\varphi} = \text{Asym} \left( \int e^{-\frac{1}{\hbar} S(t)} dt \right) \longrightarrow (S_A, \hat{\psi}_A)_{A \in \text{Sing}(\hat{\varphi})}$$

- **Sums and products.** The singularity structure of sums and products of divergent series is well-understood:

$$\begin{aligned} \lambda_1 \tilde{\varphi}_1 + \lambda_2 \tilde{\varphi}_2 &\longrightarrow (S_A^+, \hat{\psi}_A^+)_{A \in \text{Sing}(\hat{\varphi}_1) \cup \text{Sing}(\hat{\varphi}_2)} \\ \tilde{\varphi}_1 \cdot \tilde{\varphi}_2 &\longrightarrow (S_A^\times, \hat{\psi}_A^\times)_{A \in \text{Sing}(\hat{\varphi}_1) \cup \text{Sing}(\hat{\varphi}_2)} \end{aligned}$$

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# Example: the Airy kernel

$$\begin{aligned}
 K(z, w; \hbar) &= \frac{\text{Ai}(z^2; \hbar) \text{Bi}'(w^2; \hbar) - \text{Ai}'(z^2; \hbar) \text{Bi}(w^2; \hbar)}{z^2 - w^2} = \sum_{m \geq 0} a_m \hbar^m \\
 &= \frac{1}{2\sqrt{zw(z-w)}} - \frac{1}{(zw)^{3/2}} \left( \frac{5}{96z^2} - \frac{7}{96zw} + \frac{5}{96w^2} \right) \hbar \\
 &\quad + \frac{1}{(zw)^{3/2}} \left( \frac{385}{9216z^5} - \frac{455}{9216z^4w} + \frac{385}{9216z^3w^2} - \frac{385}{9216z^2w^3} + \frac{455}{9216zw^4} - \frac{385}{9216w^5} \right) \hbar^2 + \dots
 \end{aligned}$$

$\widehat{\text{Ai}}(z^2; \hbar)$  and  $\widehat{\text{Ai}}'(z^2; \hbar)$  have

- a **log singularity** at  $\frac{4}{3}z^3$
- **Stokes** constant:  $S = 1$
- **minors**:  $\text{Bi}(z^2; \hbar)$  and  $\text{Bi}'(z^2; \hbar)$

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- **minors**:  $\text{Ai}(w^2; \hbar)$  and  $\text{Ai}'(w^2; \hbar)$

$$\Rightarrow \quad a_m = \frac{\Gamma(m)}{2\pi} \left( \frac{1}{(\frac{4}{3}z^3)^m} \frac{w-z}{2\sqrt{zw}(z^2-w^2)} + \frac{1}{(-\frac{4}{3}w^3)^m} \frac{z-w}{2\sqrt{zw}(z^2-w^2)} + \dots \right)$$

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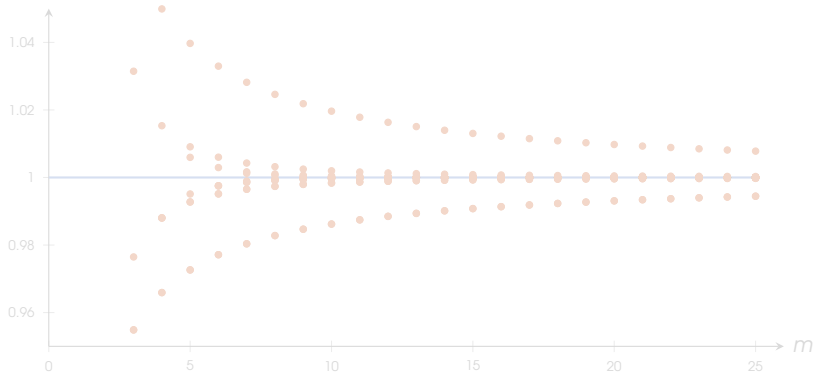
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$$\Rightarrow a_m = \frac{(-1)^m}{4\pi} \frac{\Gamma(m)}{(\frac{4}{3})^m} \left( \frac{1}{(zw)^{3/2}} h_{3m-1} \left( \frac{1}{z}, -\frac{1}{w} \right) + \dots \right)$$

# Example visualised

Write  $a_m = \frac{(-1)^m}{(zw)^{3/2}} \sum_{k+\ell=3m-1} a_{k,\ell} \frac{1}{z^k (-w)^\ell}.$

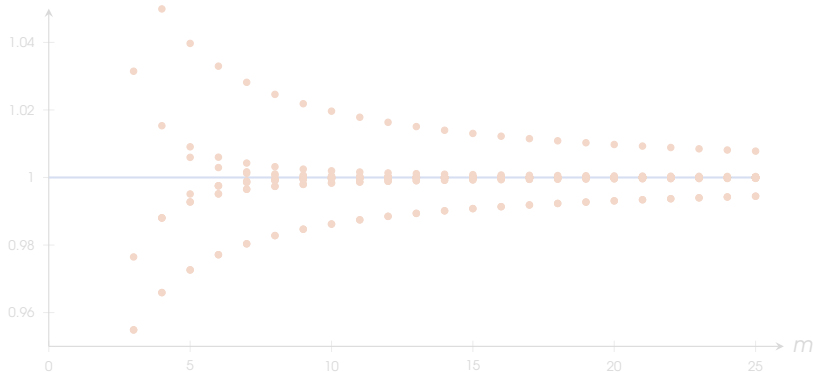
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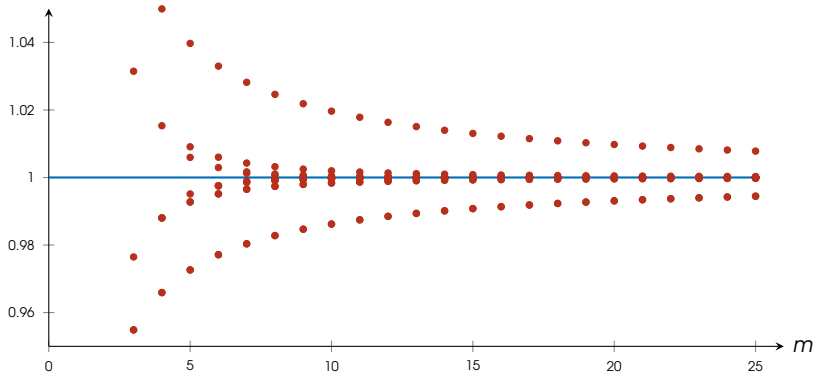
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# Strategy towards large genus asymptotics

## Goal

Compute the large genus asymptotics of  $\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle_g$

- 1 Take the  $n$ -pnt fnct

$$W_n(z_1, \dots, z_n; \hbar) = \sum_{g \geq 0} \hbar^{2g-2+n} W_{g,n}(z_1, \dots, z_n)$$

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# Determinantal formula

Define the disconnected  $n$ -pnt fnct and recall the Airy kernel

$$W_n^\bullet(z_1, \dots, z_n; \hbar) = \sum_{P \in \text{Part}(n)} W_{\ell(P)}(z_P; \hbar),$$

$$K(z, w; \hbar) = \frac{\text{Ai}(z^2; \hbar) \text{Bi}'(w^2; \hbar) - \text{Ai}'(z^2; \hbar) \text{Bi}(w^2; \hbar)}{z^2 - w^2}.$$

Determinantal formula (Bergère–Eynard, Bertola–Dubrovin–Yang):

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# Singularity structure of $W_n$

Singularity strct  
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- $2n \log$  **singularities** of  $\widehat{W}_n$ , located at

$$+ \frac{4}{3} z_i^3 \quad \text{and} \quad - \frac{4}{3} z_i^3, \quad i = 1, \dots, n$$

- **Stokes constants**:  $S = 1$

- **Minors**:

Ⓐ at  $+\frac{4}{3} z_i^3$ : replace each  $(A_i, A_i')$  with  $(B_i, B_i')$

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Uniformly in  $d_1, \dots, d_n$  as  $g \rightarrow \infty$ :

$$\langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle_g = S \frac{2^n}{4\pi} \frac{\Gamma(2g-2+n)}{A^{2g-2+n}} \left( 1 + \frac{A}{2g-3+n} \alpha_1 + \cdots \right. \\ \left. + \frac{A^k}{(2g-3+n)_k} \alpha_k + O(g^{-k-1}) \right)$$

where:

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- $A = 2/3$   
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## Bessel

Norbury's int. nmbrs (BGW  $\tau$ -fnct (Chidambaram–Garcia–Failde–AG)):

$$\begin{aligned} \langle\langle \tau_{d_1} \cdots \tau_{d_n} \rangle\rangle_g^\Theta &= \int_{\overline{\mathcal{M}}_{g,n}} \Theta_{g,n} \prod_{i=1}^n \psi_i^{d_i} (2d_i + 1)!! \\ &= S \frac{2^n}{4\pi} \frac{\Gamma(2g - 2 + n)}{A^{2g-2+n}} \left( 1 + \frac{A}{2g - 3 + n} \alpha_1 + \cdots \right. \\ &\quad \left. + \frac{A^k}{(2g - 3 + n)^k} \alpha_k + O(g^{-k-1}) \right) \end{aligned}$$

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# $r$ -Airy

Witten  $r$ -spin int. nmbrs ( $r$ -KdV  $\tau$ -fnct (Faber–Shadrin–Zvonkine)):

$$\begin{aligned}
 \langle\langle \tau_{a_1, d_1} \cdots \tau_{a_n, d_n} \rangle\rangle_g^{r\text{-spin}} &= \int_{\overline{\mathcal{M}}_{g,n}} c_w(a_1, \dots, a_n) \prod_{i=1}^n \psi_i^{d_i} (rd_i + a_i)!_{(r)} \\
 &= \frac{2^n}{2\pi} \frac{\Gamma(2g-2+n)}{r^{g-1-|d|}} \left[ \frac{S_{r,1}}{|A_{r,1}|^{2g-2+n}} \left( \alpha_0^{(r,1)} + \frac{|A_{r,1}|}{2g-3+n} \alpha_1^{(r,1)} + \cdots \right) \right. \\
 &\quad + \cdots \\
 &\quad + \frac{S_{r, \lfloor \frac{r-1}{2} \rfloor}}{|A_{r, \lfloor \frac{r-1}{2} \rfloor}|^{2g-2+n}} \left( \alpha_0^{(r, \lfloor \frac{r-1}{2} \rfloor)} + \frac{|A_{r, \lfloor \frac{r-1}{2} \rfloor}|^K}{2g-3+n} \alpha_1^{(r, \lfloor \frac{r-1}{2} \rfloor)} + \cdots \right) \\
 &\quad \left. + \frac{\delta_r^{\text{even}}}{2} \frac{S_{r, \frac{r}{2}}}{|A_{r, \frac{r}{2}}|^{2g-2+n}} \left( \alpha_0^{(r, \frac{r}{2})} + \frac{|A_{r, \frac{r}{2}}|^K}{2g-3+n} \alpha_1^{(r, \frac{r}{2})} + \cdots \right) \right]
 \end{aligned}$$

where  $S_{r,i}$ ,  $A_{r,i}$ ,  $\alpha_k^{(r,i)}$  are obtained the  $r$ -Airy ODE.

Thank you for the attention!