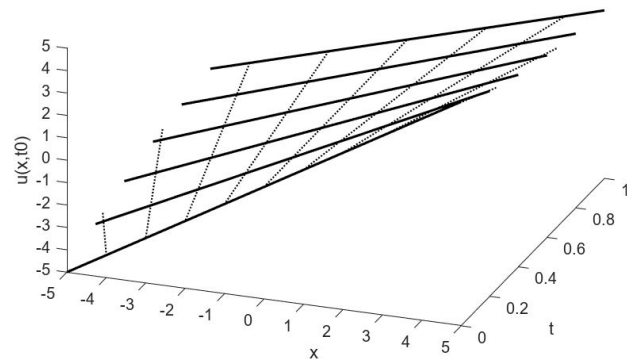
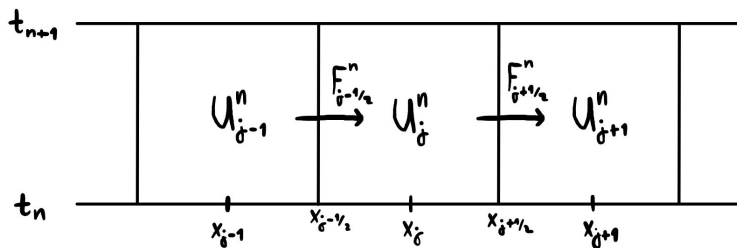
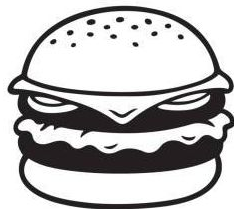


Solving the Inviscid Burgers' Equation Numerically

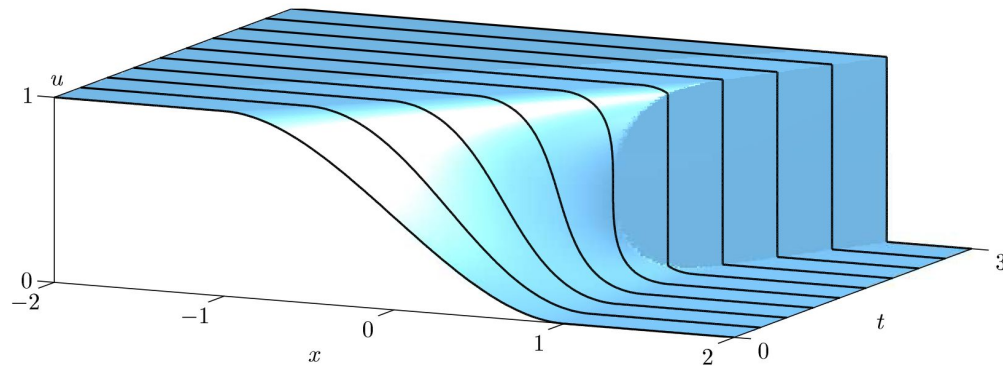
Andre Gormann and Ethan MacDonald



The Inviscid Burgers' Equation

- Modelling wave propagation
- Nonlinear Advection
- Shock formation (discontinuities)
- Applications include the modelling of fluids, sound waves, and traffic flow.

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = 0$$



The (Viscous) Burgers' Equation

- $\epsilon > 0$ is the diffusion coefficient
- The diffusion term causes the wave to dissipate as it moves
- The diffusion term disallows the formation of shocks
- Sharp but non-vertical gradients
- Inviscid Burgers' arises when we let $\epsilon \rightarrow 0$

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \epsilon \frac{\partial^2 u}{\partial x^2}$$

Conservation Laws

$$u_t + f(u)_x = 0 \qquad f(u) = \frac{1}{2}u^2$$

$$\frac{d}{dt} \int_a^b u(x, t) dx = f(u(a, t)) - f(u(b, t)) \qquad (5)$$

Conservation Laws

If we integrate (5) from time t_n to t_{n+1} , we get

$$\int_{x_{j-1/2}}^{x_{j+1/2}} u(x, t_{n+1}) dx = \int_{x_{j-1/2}}^{x_{j+1/2}} u(x, t_n) dx - \left[\int_{t_n}^{t_{n+1}} f(u(x_{j+1/2}, t)) dt - \int_{t_n}^{t_{n+1}} f(u(x_{j-1/2}, t)) dt \right]. \quad (10)$$

Dividing by Δx and applying (9) yields

$$\bar{u}_j^{n+1} = \bar{u}_j^n - \frac{1}{\Delta x} \left[\int_{t_n}^{t_{n+1}} f(u(x_{j+1/2}, t)) dt - \int_{t_n}^{t_{n+1}} f(u(x_{j-1/2}, t)) dt \right]. \quad (11)$$

Conservation Laws

$$\mathcal{F}(U_j^n, U_{j+1}^n) \approx \frac{1}{\Delta t} \int_{t_n}^{t_{n+1}} f(u(x_{j+1/2}, t)) dt$$

$$\mathcal{F}(U_{j-1}^n, U_j^n) \approx \frac{1}{\Delta t} \int_{t_n}^{t_{n+1}} f(u(x_{j-1/2}, t)) dt.$$

$$U_j^{n+1} = U_j^n - \frac{\Delta t}{\Delta x} [\mathcal{F}(U_j^n, U_{j+1}^n) - \mathcal{F}(U_{j-1}^n, U_j^n)] .$$

First-Order Godunov Scheme (REA)

- **Reconstruct** a piecewise polynomial $p(x, t_n)$ from the approximate cell averages

$$\underline{U_j^n \approx \bar{u}_j^n} = \frac{1}{\Delta x} \int_{x_{j-1/2}}^{x_{j+1/2}} u(x, t_n) dx.$$

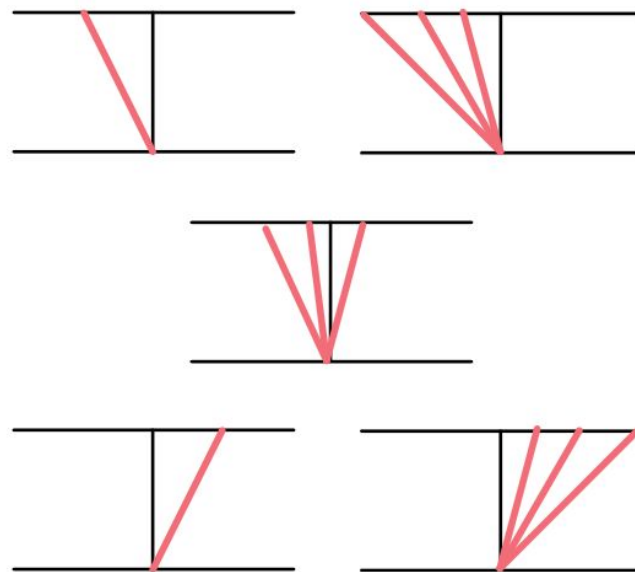
- **Evolve** the PDE from time t_n to t_{n+1} using $p(x, t_n)$ as initial profile
- **Average** the new $p(x, t_{n+1})$ to obtain new cell averages

$$U_j^{n+1} = \frac{1}{\Delta x} \int_{x_{j-1/2}}^{x_{j+1/2}} p(x, t_{n+1}) dx.$$

First-Order Godunov Scheme (REA)

- We evolve by solving Riemann problems

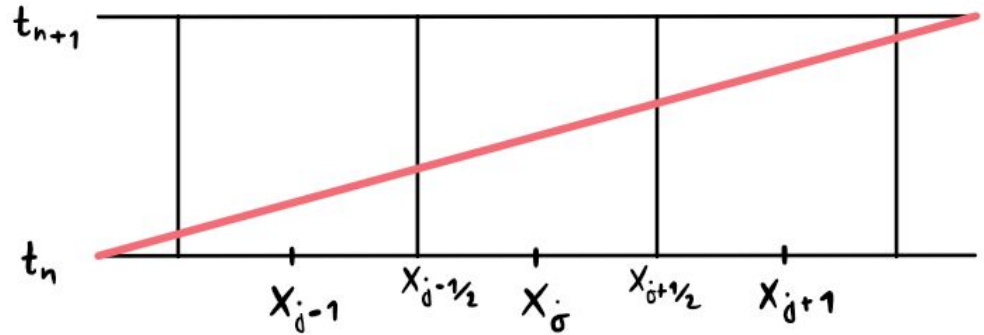
$$\mathcal{F}(u_L, u_R) = \begin{cases} \min_{u \in [u_L, u_R]} f(u) & u_L \leq u_R \\ \max_{u \in [u_R, u_L]} f(u) & u_L > u_R \end{cases}$$



Convergence

- CFL bounded by maximum wave-speed => variable time-step
- Bounds how far a characteristic can travel

$$\max |U_j^n| \frac{\Delta t_n}{\Delta x} \leq 1$$



Convergence

Theorem 12.3. Suppose $Q^{(\Delta t)}$ is generated by a numerical method in conservation form with a Lipschitz continuous numerical flux, consistent with some scalar conservation law. If the method is TV-stable, i.e., if $\text{TV}(Q^n)$ is uniformly bounded for all $n, \Delta t$ with $\Delta t < \Delta t_0, n \Delta t \leq T$, then the method is convergent, i.e., $\text{dist}(Q^{(\Delta t)}, \mathcal{W}) \rightarrow 0$ as $\Delta t \rightarrow 0$.

Riemann Problem

The Riemann problem is the IVP for Inviscid Burgers' with initial condition

$$u(x, 0) = \begin{cases} u_L & \text{if } x < 0 \\ u_R & \text{if } x \geq 0 \end{cases}$$

True Solutions:

Shockwave Case: $u_L > u_R$

$$u(x, t) = \begin{cases} u_L & \text{if } x < st \\ u_R & \text{if } x \geq st \end{cases}$$

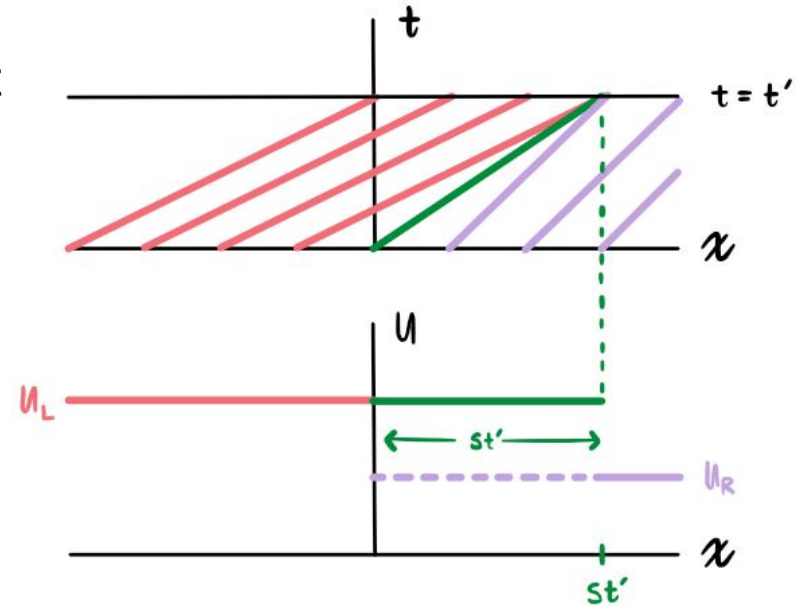
$$s = \frac{u_L + u_R}{2}$$

Rarefaction case: $u_L < u_R$

$$u(x, t) = \begin{cases} u_L & \text{if } x \leq u_L t \\ \frac{x}{t} & \text{if } u_L t < x < u_R t \\ u_R & \text{if } x \geq u_R t \end{cases}$$

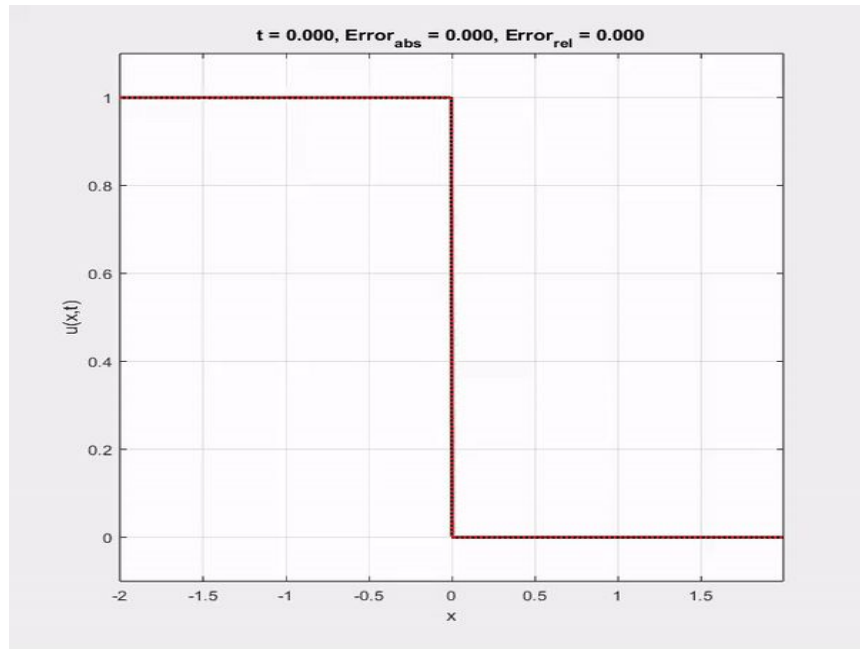
Riemann Shockwave Problem ($u_L > u_R$)

- u_L is moving faster than u_R
- u_L and u_R characteristic curves intersect
- Solution becomes discontinuous at the point in space where the curves intersect
- The discontinuity point moves through space as time progresses



Riemann Shockwave Problem ($u_L > u_R$)

Numerical Solution

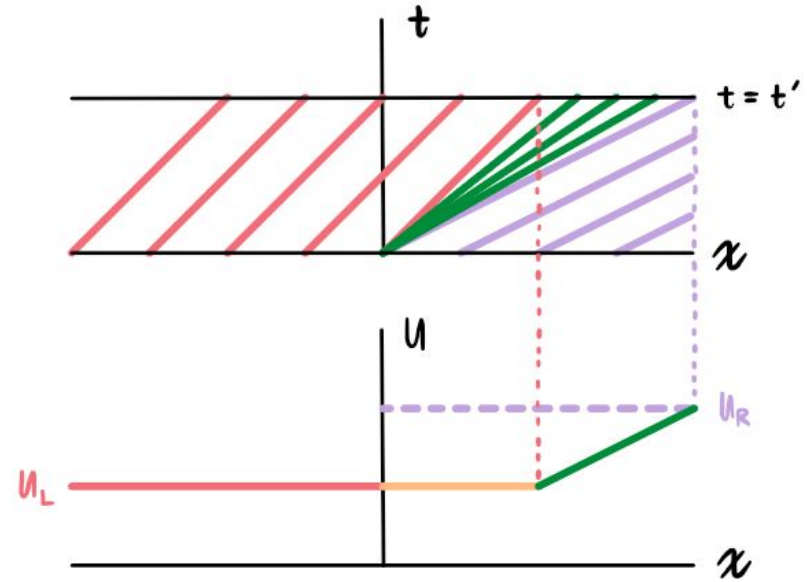


PINN Solution



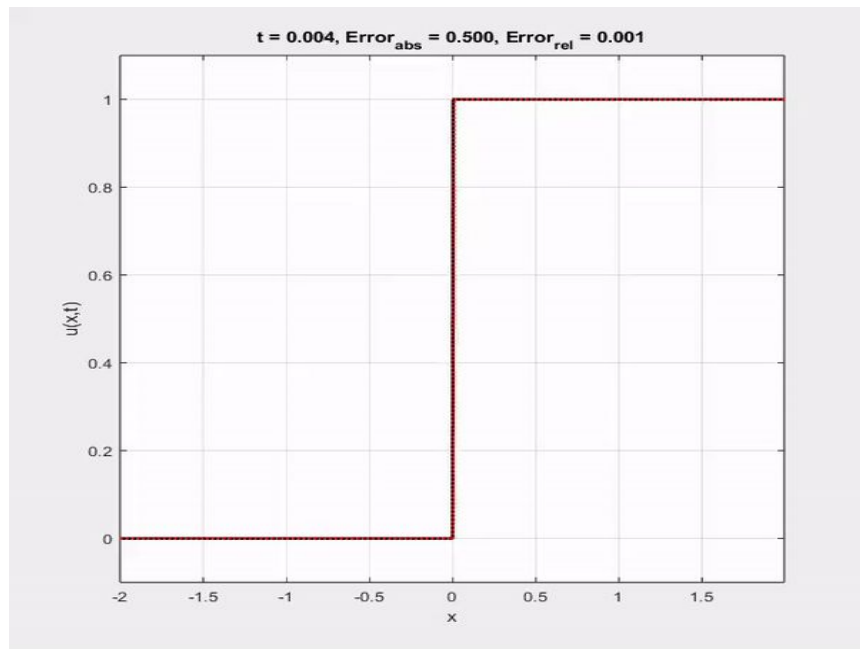
Riemann Rarefaction Wave Problem ($u_L < u_R$)

- u_L is moving slower than u_R
- u_L and u_R characteristic curves never intersect, so a shock never forms
- The characteristic curves become increasingly rarified as time progresses
- The solution is sloped between u_L and u_R
- This slope becomes longer and steeper as time progresses

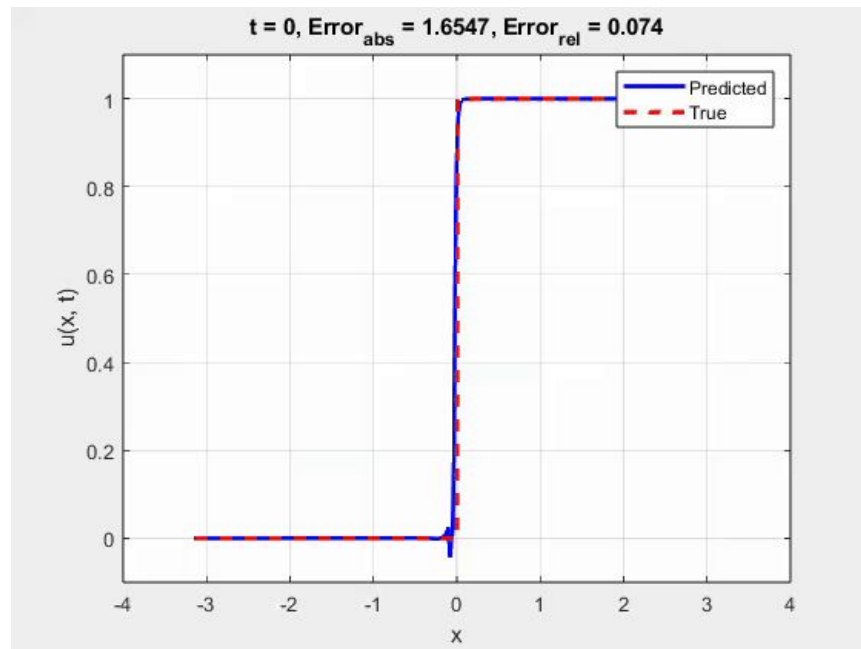


Riemann Rarefaction Wave Problem ($u_L < u_R$)

Numerical Solution



PINN Solution

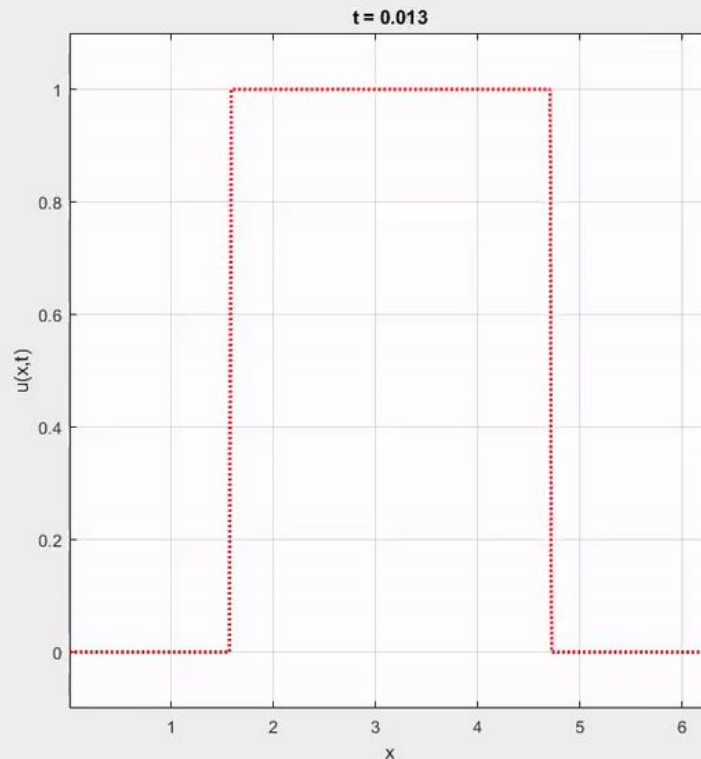


Square Wave Problem

$$\text{IVP: } \begin{cases} u_t + uu_x = 0, & x \in [0, 2\pi], t > 0 \\ u(x, 0) = u_0(x) \\ u(0, t) = u(2\pi, t), & t > 0 \end{cases}$$

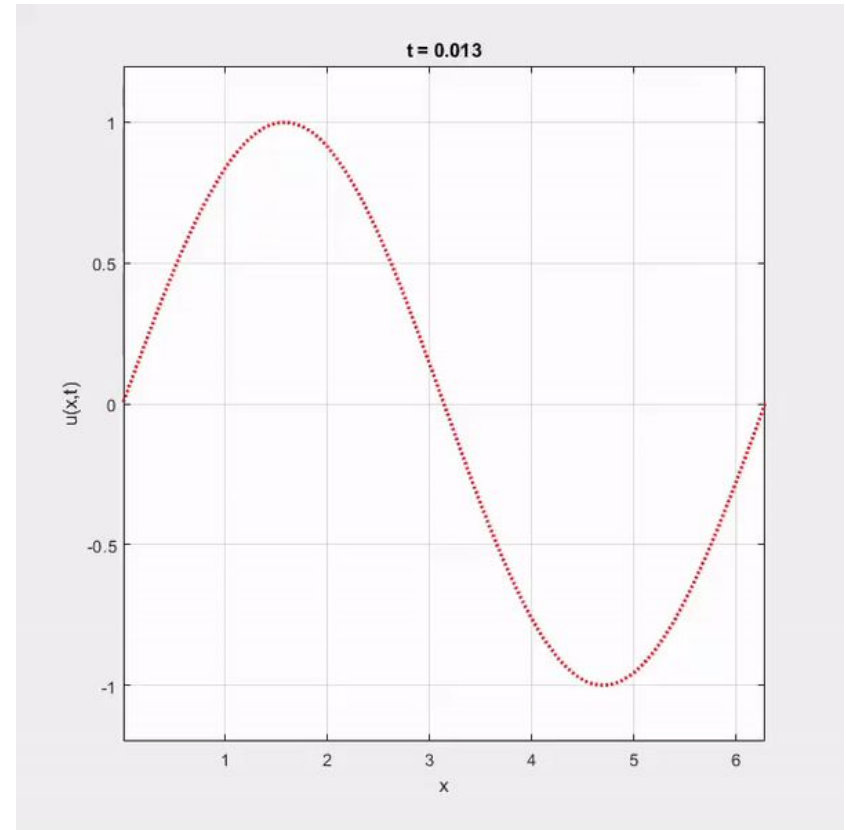
Where

$$u_0(x) = \begin{cases} 1 & x \in [\pi/2, 3\pi/2] \\ 0 & \text{otherwise} \end{cases}$$



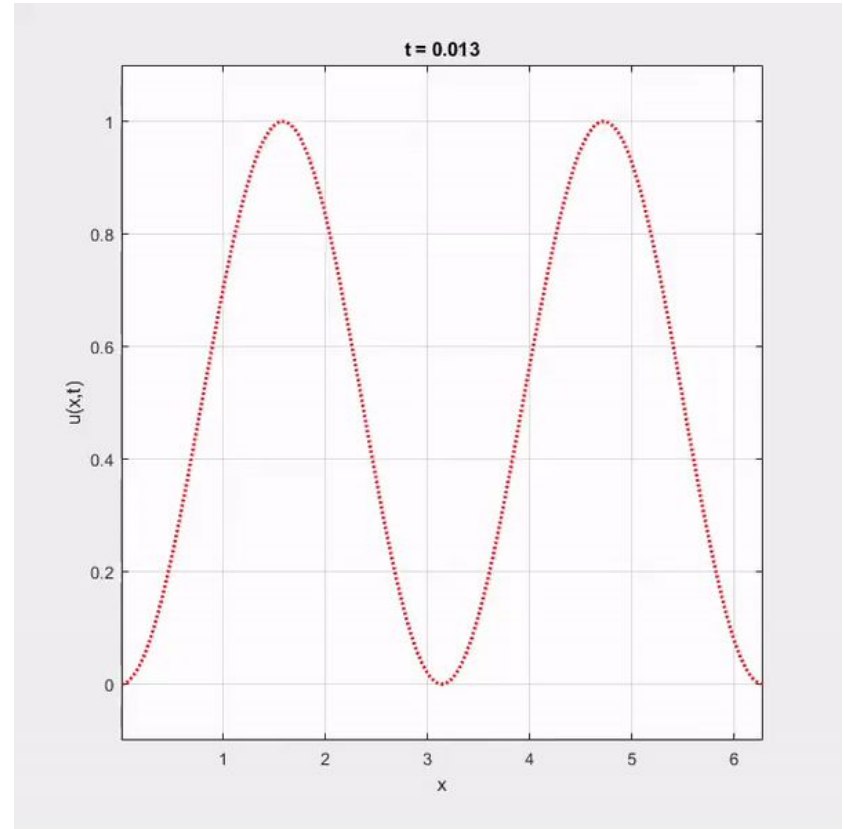
Sine Wave Problem

$$\text{IVP: } \begin{cases} u_t + uu_x = 0, & x \in [0, 2\pi], t > 0 \\ u(x, 0) = \sin(x) \\ u(0, t) = u(2\pi, t), & t > 0 \end{cases}$$



Sine-Squared Wave Problem

IVP:
$$\begin{cases} u_t + uu_x = 0, & x \in [0, 2\pi], t > 0 \\ u(x, 0) = \sin^2(x) \\ u(0, t) = u(2\pi, t), & t > 0 \end{cases}$$



Q & A

