

Quasiconformal Mappings with Surface Domains

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Quasiconformal mappings from surfaces immersed in Euclidean space are discussed for the purposes of computing dilatation-optimal surface meshes with prescribed connectivity and Dirichlet boundary data. In particular, a quaternionic formulation of quasiconformality is proposed which facilitates a linear algorithm for computing least-squares quasiconformal maps from surfaces given as extrinsic mesh data. Additionally, an iterative procedure is given which computes optimal quasiconformal mappings with optional constraints on surface area and extrinsic geometry. Based on the established Quasiconformal Iteration method, the proposed algorithm produces high quality surface mappings which correctly capture boundary information while eliminating undesirable folds which appear during least-squares conformal mapping procedures.

CCS Concepts: • Computing methodologies → Shape modeling; • Mathematics of computing → Partial differential equations.

Additional Key Words and Phrases: quasiconformal mapping, Teichmüller mapping, mesh optimization, surface remeshing

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1 INTRODUCTION

Conformal maps between two-dimensional Riemann surfaces are widely recognized as useful tools in both theoretical and computational settings. In particular, enough interesting quantities are conformally invariant (e.g. the total Gaussian curvature and the Willmore energy of surfaces) that many difficult problems can be made considerably simpler by applying an appropriate conformal transformation. Moreover, at the present time there are several effective algorithms for computing conformal or nearly conformal mappings (see e.g. [Bobenko et al. 2015; Gu et al. 2004; Kharevych et al. 2006; Sawhney and Crane 2017; Springborn et al. 2008; Trefethen 2020] and references therein), the use cases of which cover everything from surface flattening to medical image registration.

Despite their numerous advantages, it is also well known that conformal maps are unsuitable for mapping problems with a pointwise boundary correspondence. In particular, there is very often no conformal map between two connected surfaces which maps boundaries to boundaries in a prescribed way, even in quite simple cases.

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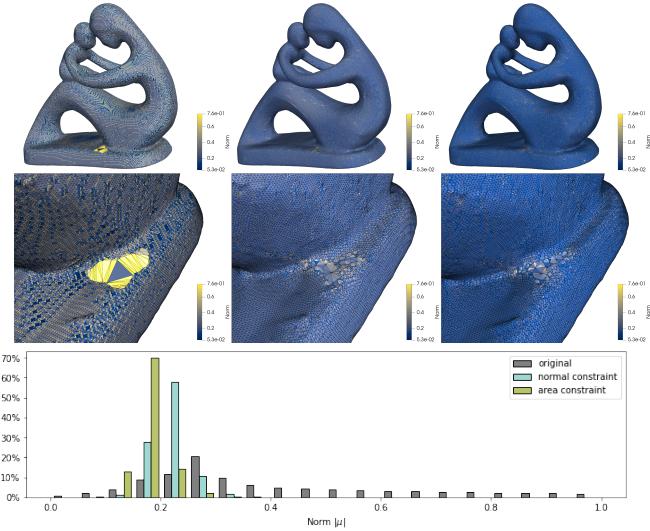


Fig. 1. Quasiconformal remeshings of a statue mesh with genus four constructed using the QC Iteration algorithm from Section 5. Left: original surface; Middle: remeshing with normal constraint (c.f. Section 4.1). Right: remeshing with area constraint (c.f. Section 4.2). Bottom: histogram of the norm of the Beltrami coefficient μ .

For example, it can be shown that there is no corner-preserving conformal mapping from a square onto a rectangle; no matter the ratio of lengths, there is an inevitable amount of shearing distortion that occurs during this process which is inherently non-conformal. This presents a significant challenge for computational applications which involve some form of mesh deformation, where a conformal (or close-to-conformal) mapping is desired which satisfies some given boundary data. While least-squares conformal mapping techniques can certainly be applied in this instance, they are known to produce undesirable folds if the target surface is non-convex (see e.g. Figure 2 or Figure 6).

On the other hand, if the condition of conformality is relaxed to that of quasiconformality instead (i.e. uniformly bounded conformality distortion), then such maps are relatively abundant. In fact, a mild growth condition on the distortion at the boundary guarantees the existence and uniqueness of a “best possible” quasiconformal map between two Riemann surfaces (see Theorem 2.2). Known as a Teichmüller mapping, this extremal mapping minimizes the maximal conformality distortion throughout the source domain. As these Teichmüller mappings can in many cases be realized as the unique minimizers of a particular energy functional (c.f. Section 2), they are also relatively computable using conventional numerical minimization techniques, making extremal quasiconformal mapping a valid

strategy in situations when an explicit boundary correspondence must be satisfied.

1.1 Related Work

Due to their advantageous properties, quasiconformal (QC) mappings have recently been investigated for a variety of computational applications. This usually involves working with some form of the Beltrami equation $f_{\bar{z}} = \mu f_z$ (see Section 2), which characterizes the quasiconformal mapping and can be discretized and solved on a manifold mesh or point cloud. Planar quasiconformal mappings are used in [Zeng and Gu 2011] for surface registration by pre-composing the computation of the Beltrami coefficient (BC) μ with a Ricci flow procedure to homogenize the domain. On the other hand, [Zeng et al. 2012] employ a discrete version of the Yamabe flow to compute planar Teichmüller mappings by evolving a given quasiconformal metric to a Teichmüller minimum, while [Nian and Chen 2016] compute similar mappings for isogeometric analysis using B-spline techniques. In a different direction, [Lipman 2012] formulates convex spaces of bounded distortion mappings which are computable and contain quasiconformal mappings as a subclass, while [Weber et al. 2012] make use of holomorphic quadratic differentials to formulate a minimization-based method for computing extremal Teichmüller maps between planar domains with boundary, or genus 0 surfaces with additional conformal mapping. Yet another approach is taken in [Lui et al. 2014] where an alternating minimization called “QC Iteration” for computing Teichmüller mappings of planar domains and genus 0 surfaces is developed. This QC iteration is connected to the theory of harmonic mappings in [Lui et al. 2015] and shown to converge under some assumptions. More recently, works such as [Choi 2021] have also used quasiconformal mappings as a way to obtain nearly-conformal parameterizations of planar domains.

The available techniques for computing quasiconformal maps have since been expanded to include other topological types as well as other forms of data. The work [Meng et al. 2016] develops an algorithm called TEMPO for computing Teichmüller mappings on point clouds. Moreover, quasiconformal maps of multiply-connected planar domains with prescribed distortion are computed in [Ho and Lui 2016], and [Lee et al. 2016] compute quasiconformal maps between 3D volumes for the purpose of surface registration. The variety of algorithms available for computing quasiconformal maps has also led to a number of interesting applications, such as a quasiconformal kernel for nearest neighbor calculations in machine learning algorithms [Peng et al. 2004] as well as a method for feature-preserving image resizing [Xu et al. 2018]. In addition, quasiconformal mappings have also been used to create origami-like surfaces with prescribed folds [Qiu et al. 2019], to compute quasiconformal rectilinear mappings for planar subdivision surfaces [Yang and Zeng 2020], and to study the morphometry of teeth [Choi et al. 2020].

1.2 Contributions

Despite the far-reaching interest into quasiconformal mappings, at present there are no methods which are adequate for computing quasiconformal mappings between surfaces of nonzero genus. Although the mentioned methods can conceivably be applied piece-wise after

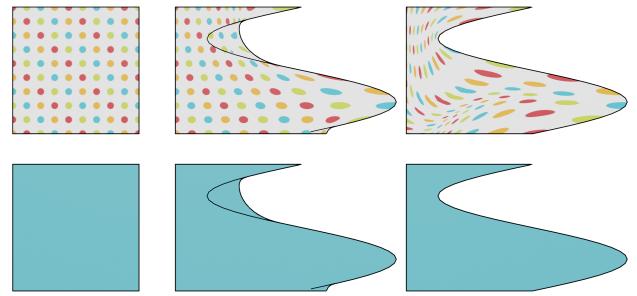


Fig. 2. Comparison between least-squares conformal mapping (LSCM, middle) and Teichmüller quasiconformal mapping (TQCM, right) on a planar domain with boundary. Note that the least-squares conformal mapping does not respect injectivity, despite producing less average texture distortion.

cutting each surface along the generators of its fundamental group, it is a nontrivial matter in practice to compose such a procedure with the quasiconformal mapping algorithms above. Additionally, there appear to be no algorithms in place for computing quasiconformal mappings directly from a non-planar surface of any genus. Instead, non-planar domains must first be mapped conformally to the plane, which is often nontrivial and requires additional computational resources.

The present work addresses these issues in the case that the source surface is embedded (or more generally, immersed) in \mathbb{R}^3 by providing a direct algorithm for computing quasiconformal mappings with prescribed Dirichlet boundary conditions, which can be applied to surfaces of arbitrary genus. Using a discretization based on finite elements, this procedure is essentially a single linear solve which computes a least-squares quasiconformal mapping with respect to some predefined BC μ given as a value-per-polygon. It is shown that the present algorithm for computing quasiconformal maps enables an extension of the Quasiconformal Iteration algorithm of Lui et al. [Lui et al. 2014] mentioned previously, which computes the extremal Teichmüller map $f : M \rightarrow P$ between two connected planar domains, or (with additional conformal mapping) between two connected, genus zero Riemann surfaces with boundary. The technology introduced here extends the applicability of the QC Iteration to maps from surfaces of arbitrary genus which are represented as (planar or non-planar) manifold meshes in \mathbb{R}^3 , giving a useful procedure for folding-free object deformation and surface remeshing. Specifically, the present contributions are:

- An intrinsic-to-extrinsic formulation of quasiconformal mappings $f : M \rightarrow f(M) \subset \mathbb{R}^3$ based on quaternionic theory which is convenient for computation.
- A self-contained, genus-agnostic algorithm for computing quasiconformal mappings whose domain is a manifold mesh in \mathbb{R}^3 .
- An extension of the QC Iteration algorithm for computing Teichmüller mappings to this setting, which computes optimal Teichmüller mappings from immersed surfaces which

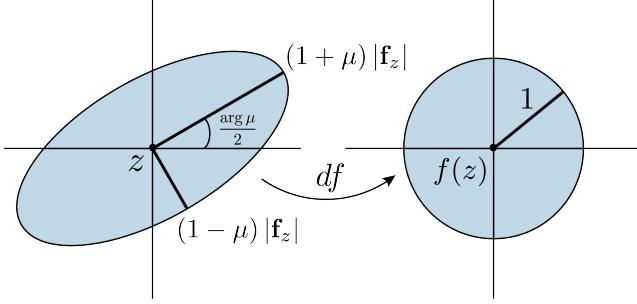


Fig. 3. The geometry of a quasiconformal mapping $f : M \rightarrow \mathbb{R}^3$.

are simply connected and dilatation-uniform quasiconformal mappings otherwise.

- Applications to surface remeshing and the construction of locally injective mappings satisfying prescribed Dirichlet boundary conditions.

The remainder of this work is structured as follows. The basics of quasiconformal maps and the QC Iteration are presented in Section 2. Section 3 gives the novel quaternionic formulation of quasiconformality which enables the main algorithm in Section 4. Section 5 then extends the QC Iteration to the present setting, and some applications of this are discussed in Section 6.

2 PRELIMINARIES

This Section recalls the standard presentation of extremal quasiconformal mappings which is necessary to describe the QC Iteration algorithm of [Lui et al. 2015, 2014]. More information regarding the theory of quasiconformal mappings can be found in [Gardiner and Lakic 2000; Hubbard 2006; Strebel 1984], and a good account of the connection between Teichmüller and harmonic mappings is given in [Daskalopoulos and Wentworth 2007].

2.1 Extremal Quasiconformal Mappings

Recall that a quasiconformal map $f : M \rightarrow P$ between Riemann surfaces is an orientation-preserving homeomorphism which has bounded conformality distortion with respect to a given complex structure J on TM . In standard notation, f is μ -quasiconformal provided it satisfies the *Beltrami equation*,

$$\bar{\partial}f = \partial f \circ \mu,$$

where $\partial, \bar{\partial}$ are the \mathbb{C} -linear resp. \mathbb{C} -antilinear parts of the natural derivative operator $df : TM \rightarrow TP$ and $\mu : TM \rightarrow TM$, $|\mu|_\infty < 1$, is the \mathbb{C} -antilinear *Beltrami coefficient* (also called Beltrami differential and abbreviated BC) of the mapping. In a local conformal coordinate $z : U \subset M \rightarrow \mathbb{C}$, this implies the expression

$$f_{\bar{z}} = \mu f_z,$$

where $f_z := \partial_z f$ (resp. $f_{\bar{z}} := \partial_{\bar{z}} f$) are the partial derivatives of the mapping f with respect to the conformal coordinate z and $\mu : U \rightarrow \mathbb{C}$ is the locally defined BC. Notice that f is conformal if and only if $\mu \equiv 0$, and this condition depends only on the conformal structure of M .

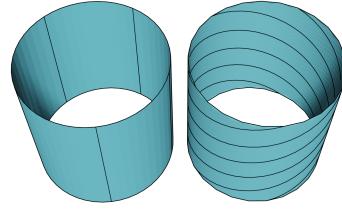


Fig. 4. Two non-homotopic families of maps from an interval to the cylinder (i.e. curves) with identical boundary data.

Geometrically, the Beltrami equation implies that quasiconformal maps take small circles on the source space to small ellipses of bounded eccentricity on the target (see Figure 3). To see this, suppose $f : M \rightarrow f(M) \subset \mathbb{R}^3$ immerses the surface M in Euclidean space. Let g be a Riemannian metric on M with local expression $g = \sigma |dz|^2$ for some positive function $\sigma : U \rightarrow \mathbb{R}$, and denote the Euclidean metric on \mathbb{R}^3 by δ . Then, the Jacobian determinant describing the local area distortion of f has the expression

$$\text{Jac}(f) = |f_z|^2 - |f_{\bar{z}}|^2 = |f_z|^2 (1 - |\mu|^2) = \sqrt{\lambda_1 \lambda_2},$$

where λ_1, λ_2 are the eigenvalues of the pullback metric $f^*\delta$ on M relative to the flat metric $|dz|^2$. Moreover, since

$$f^*\delta = \langle df, df \rangle = |f_z|^2 |dz + \mu d\bar{z}|^2,$$

these eigenvalues are given by

$$\lambda_1 = |f_z|^2 (1 + |\mu|)^2, \quad \lambda_2 = |f_z|^2 (1 - |\mu|)^2.$$

These are the squared lengths of the major resp. minor axes of the ellipse in $T_z M$ which pushes forward under df to the unit circle in $T_{f(z)} f(M)$, quantifying the shearing distortion induced by quasiconformality. Their ratio defines the maximal dilatation of f ,

$$K(f) = \frac{1 + |\mu|_\infty}{1 - |\mu|_\infty},$$

which is $1 \leq K < \infty$ for orientation-preserving maps and $-\infty < K < -1$ for orientation-reversing maps. Note that $\text{Jac}(f) > 0$ when $f_z \neq 0$ and $|\mu| < 1$, reflecting the remarkable fact that quasiconformal mappings are locally injective. In the discrete setting, this implies that f cannot have fold-overs or wrap-ups (i.e. places where the mapping fails to be immersive, see Figures 2, 6, and 12), which is essential in applications such as medical device simulations where surfaces must remain embedded as they deform [cite?](#)

Although there are generally many quasiconformal mappings from one Riemann surface to another, there are relatively few which are distinguished as being extremal. These are the mappings f which minimize the maximal dilatation K , i.e.

$$K(f) \leq K(f'),$$

for any map $f' : M \rightarrow P$ such that $f' = f$ on ∂M . Such maps always exist for each homotopy class $[f] = \{f' : M \rightarrow P : f'|_{\partial M} = f|_{\partial M}\}$ relative to the boundary (see Figure 4 for an illustration of this concept), but need not be unique. On the other hand, in many circumstances there is a unique extremal map of special form which

is compatible with the given boundary data. This is called a Teichmüller map, and has the desirable property of uniform conformality distortion throughout the whole domain.

Definition 2.1. The quasiconformal mapping $f : M \rightarrow f(M)$ between Riemann surfaces is said to be Teichmüller provided there exists a constant $k \in \mathbb{R}$ and holomorphic quadratic differential $q \in T^*M \otimes T^*M$ such that the Beltrami differential $\mu : TM \rightarrow TM$ satisfies

$$\mu = k \frac{\bar{q}}{|q|}, \quad k = \|q\|_1 = \int_M |q|,$$

In this case, μ is said to be Teichmüller associated to f .

It is a fact that when a Teichmüller map exists, it is unique and extremal for its homotopy class. Moreover, existence is guaranteed in many practically-relevant situations thanks to classical results in complex analysis. In particular, define the boundary dilatation of the homotopy class $[f]$ as

$$H([f]) = \inf_{f' \in [f]} \left\{ \inf_{C \subset M} K(f'|_{M \setminus C}) \right\},$$

where C is any compact set strictly contained in M . The primary existence result for Teichmüller maps in terms of this criterion is due to Strebel [Strebel 1984].

THEOREM 2.2. [Gardiner and Lakic 2000, pg 127] Let M be a connected Riemann surface with potential boundary, and let $f : M \rightarrow f(M)$ be quasiconformal. Suppose the boundary dilatation $H([f]) < K(f)$. Then, $[f]$ contains an extremal Teichmüller mapping such that $\mu = k\bar{q}/|q|$ for a unique constant $k \in \mathbb{R}$ and integrable holomorphic quadratic differential $q \in T^*M \otimes TM$ which is unique up to multiplication by a positive constant.

From the above, it is clear that Teichmüller maps are remarkably well-behaved. In particular, any quasiconformal map $f : M \rightarrow P$ between Riemann surfaces satisfying “nice enough” boundary conditions contains a unique Teichmüller extremal map in its homotopy class. Conversely, it turns out that even when no extremal Teichmüller map exists, there is always a Teichmüller map with dilatation arbitrarily close to the extremal one (see [Strebel 1978, Theorem 8]). In some sense, this makes Teichmüller maps the best possible quasiconformal mappings between Riemann surfaces M and P , and this property has generated a large amount of interest in their computation. One such method for this is the QC Iteration algorithm of [Lui et al. 2015, 2014], which is a provably convergent (under some assumptions) and minimization-based strategy for computing Teichmüller mappings.

Remark 2.3. Note that the Teichmüller mappings from [Strebel 1978, Theorem 8] are allowed to be meromorphic with at most one simple pole, as opposed to purely holomorphic. This definition considerably widens the space of Teichmüller maps, at the cost of uniqueness in certain settings. Practically, QC Iteration may produce mappings with the same meromorphic structure (see e.g. Figure 8).

2.2 The Quasiconformal (QC) Iteration

The Quasiconformal (QC) Iteration algorithm introduced in [Lui et al. 2014] is an approach to computing extremal Teichmüller mappings between planar domains and between surfaces which are



Fig. 5. TQCM from a simply connected surface equipped with an ideal conformal structure (c.f. Section 6.2) to itself equipped with the usual conformal structure inherited from its embedding in \mathbb{R}^3 .

conformal to planar domains. Its general idea is to find a downward trajectory toward the optimal Teichmüller map in a given homotopy class by alternately minimizing the distortion induced by a particular BC μ and the deviation of μ from Teichmüller form as defined in Definition 2.1. Under somewhat restrictive assumptions mentioned in [Lui et al. 2015] (such as nonpositive Gaussian curvature of the target), this algorithm is guaranteed to converge to a unique Teichmüller mapping between Riemann surfaces of the same topology, which may be simply or multiply connected. On the other hand, good practical performance is seen even in cases when a unique Teichmüller extremum cannot be guaranteed, as the QC Iteration often converges to a mapping with near-constant dilatation regardless of these assumptions (see e.g. Figure 5). This Subsection recalls the details of this procedure, which will be necessary for its extension to general immersed surfaces in Section 5.

To describe the QC Iteration more precisely, it is useful to define an energy functional on the space of Beltrami differentials μ , denoted $\mathcal{B}(M, P)$. First, recall that any Riemann surface (M, J) with local conformal coordinate $z : U \subset M \rightarrow \mathbb{C}$ gives rise to a conformal class of Riemannian metrics $[g] = \{g_\sigma = \sigma |dz|^2 \mid \sigma : U \rightarrow \mathbb{R}^+\}$ compatible with J . Suppose (M, g) and (P, h) are Riemannian surfaces and the metrics g, h are compatible with the respective complex structures on M, P . Fix a mapping $f_0 : (M, g) \rightarrow (P, h)$ and suppose that P has nonpositive Gaussian curvature everywhere. Since Teichmüller mappings are harmonic with respect to a certain conformal class of metrics on M (c.f. Appendix A), it follows from [Eells and Sampson 1964, Section 11] that there is a unique Teichmüller mapping $f^{**} \in [f_0]$ with BC μ^{**} , depending only on the metric h and the conformal class $[g]$, which is extremal in the homotopy class $[f_0]$. On the other hand, each BC $\mu \in \mathcal{B}(M, P)$ induces its own conformal structure on M with canonical metric representative $g(\mu) = |dz + \mu d\bar{z}|^2$. Identifying $\mu_1 \sim \mu_2$ if there is a biholomorphism $(M, g(\mu_1)) \mapsto (M, g(\mu_2))$ homotopic to the identity map, it follows that every equivalence class $[\mu] \subset \mathcal{B}(M, P)$ of BCs also possesses a unique representative μ^* of Teichmüller form, which is associated to the unique Teichmüller mapping f^* depending only on the metric h and the conformal class of $g(\mu)$. This suggests a definition of the

Beltrami energy as

$$\mathcal{E}(\mu) := \mathcal{D}_{g(\mu)}(f^*) = \frac{1}{2} \int_M |\partial f^*|^2_h dS_{g(\mu)},$$

where $dS_{g(\mu)}$ and $\mathcal{D}_{g(\mu)}$ denote respectively the area element and Dirichlet energy with respect to the metric $g(\mu)$, $|\cdot|_h$ denotes the norm with respect to the metric h , and $f^* = f^*(\mu, h)$ is the unique minimizer of $\mathcal{D}_{g(\mu)}$ in the homotopy class $[f_0]$. Denoting $\mathcal{T}(M, P) \subset \mathcal{B}(M, P)$ the subset of BC's with constant norm, it can be shown (see [Lui et al. 2015, Theorem 3.2]) that the global minimizer of $\mathcal{E} : \mathcal{T}(M, P) \rightarrow \mathbb{R}$ is precisely the unique Teichmüller BC μ^{**} associated to $f^{**} : (M, g) \rightarrow (P, h)$ in the homotopy class $[f_0]$. Therefore, an optimal Teichmüller mapping between Riemannian surfaces in a given homotopy class can be computed by minimizing the energy \mathcal{E} . **check all this again**

With this framework in place, the QC Iteration procedure of [Lui et al. 2014] can be understood as an alternating descent algorithm which approximates the gradient flow of \mathcal{E} . Instead of working directly with this functional, the QC Iteration utilizes the “conformal part” of \mathcal{E} (c.f. Appendix A),

$$QC_\mu(f) = \int_M |\bar{\partial}f - \partial f \circ \mu|^2_h dS_g,$$

which we call the *quasiconformal distortion* of f with respect to μ . Clearly, $f : M \rightarrow P$ is quasiconformal with respect to μ if and only if $QC_\mu(f) = 0$, and it is shown in Appendix A that minimization of $\mathcal{D}_{g(\mu)}$ is equivalent to minimization of QC_μ when the norm $|\mu|$ is constant and the area of the target P is fixed. This leads to the iterative three-stage process in [Lui et al. 2014] loosely described as follows:

- (1) Minimize $QC_\mu(f)$ for $f : (M, g) \rightarrow (P, h)$.
- (2) Compute μ algebraically from f .
- (3) Locally post-process μ to bring it closer to Teichmüller form.

Steps 1 and 2 are relatively straightforward on planar domains M , which can be covered by a single local coordinate chart $U = M$ where $g = \sigma |dz|^2$. This is because QC_μ takes the simple expression,

$$\begin{aligned} 2QC_\mu(f) &= \int_U \frac{1}{\sigma} |\mathbf{f}_z - \mu \mathbf{f}_z|^2_h \sigma idz \wedge d\bar{z} \\ &= \int_U |\mathbf{f}_z - \mu \mathbf{f}_z|^2_h idz \wedge d\bar{z}, \end{aligned}$$

and the BC μ can be computed from any quasiconformal f as

$$\mu = \frac{\mathbf{f}_z}{\mathbf{f}_z},$$

where the vector quotient on the RHS is interpreted as the division of complex numbers. Step 3 on the other hand is an indirect way to move toward the global minimizer of \mathcal{E} , and is facilitated by the following fact.

LEMMA 2.4. *Any Teichmüller Beltrami coefficient $\mu : TM \rightarrow TM$ has harmonic norm, as well as harmonic argument when restricted to local charts $U \subset M$ where $q \neq 0$. Conversely, any pair (k, θ) where $k \in \mathbb{R}$ and $\theta : U \rightarrow \mathbb{R}$ is harmonic can be associated with a Teichmüller Beltrami coefficient μ on U .*

PROOF. Recall that in the local conformal coordinate $z : U \subset M \rightarrow \mathbb{C}$ the Teichmüller Beltrami coefficient satisfies

$$\mu = k \frac{\bar{q}}{|q|},$$

where $k \in \mathbb{R}$ and $q : U \rightarrow \mathbb{C}$ is a local holomorphic function. It follows immediately that the norm is constant, hence harmonic. Moreover, since q is holomorphic, so is $\log q = \log |q| + i \arg q$ (in an appropriate branch) and it follows that $\arg q$ is harmonic. Finally, we have the phase equality

$$e^{i \arg \mu} = \frac{\bar{q}}{|q|},$$

so that $\arg \mu = -\arg q$, which implies the first conclusion. Conversely, if $\theta : U \rightarrow \mathbb{R}$ is a harmonic function, then there is a conjugate harmonic function ϕ so that $\phi - i\theta$ is holomorphic. In this case, $q = e^{\phi - i\theta}$ is also holomorphic and

$$\mu = |\mu| e^{i\theta} = k \frac{\bar{q}}{|q|}$$

is Teichmüller. \square

By this result, the global minimizer μ^* of \mathcal{E} in any homotopy class $[f_0]$ should have constant norm and harmonic argument when restricted to any local coordinate system away from the zeros of the quadratic differential q . Therefore, Step 3 in the QC Iteration is to gently perturb the μ computed in Step 2 so as to bring its norm to constant while simultaneously smoothing its argument. This involves a straightforward projection of the norm onto its average value

$$|\mu| \leftarrow \frac{\int_M |\mu| dS_g}{\int_M dS_g},$$

followed by a single step of Jacobi iteration applied to $\arg \mu$. This usually (although not always) results in an iterate μ_{k+1} with lower energy \mathcal{E} which is used as the input to the next iteration, and these steps are repeated until convergence.

Remark 2.5. Notice that the argument of the BC μ is only well defined away from the zeros of the quadratic differential q , hence so is the smoothing procedure proposed in [Lui et al. 2014]. For this reason, the extended algorithm presented in Section 5 smooths on the phase $e^{i \arg \mu} = \mu / |\mu|$ which is well defined everywhere but not continuous, as can be seen in Figure 6. Since $\Delta e^{i \arg \mu} = e^{i \arg \mu} (\Delta \arg \mu - |\nabla \arg \mu|^2)$, these procedures are equivalent when M is simply connected.

While quasiconformal mappings remain important tools for computational tasks such as meshing and surface registration, their present utility is limited by the use of planar technology such as that which enables the QC Iteration. Indeed, to compute quasiconformal mappings between generic Riemann surfaces in this way requires pre- and post-processing with conformal flattening procedures, as well as attention to highly nontrivial issues of compatibility between solutions computed on different surface patches. The next Section discusses a primary contribution of this work, which is a formulation of quasiconformality using quaternionic surface theory which enables the direct computation of quasiconformal maps from immersed Riemann surfaces of any genus.

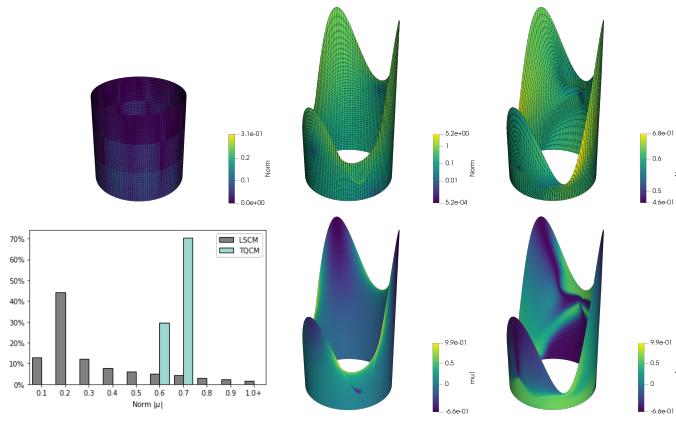


Fig. 6. Comparison of LSCM (mid) and TQCM (right) on a cylinder; top row colored by $|\mu|$ and bottom row colored by the first component μ^1 . Similar to the planar case, LSCM creates undesirable spillage and cannot guarantee injectivity. Conversely, TQCM remains injective despite a height distortion of 80% at the boundary; note the visible discontinuities in μ^1 as predicted by the theory.

3 QUASICONFORMALITY WITH QUATERNIONS

The present approach to computing quasiconformal mappings is facilitated by technology from quaternionic surface theory, the details of which are discussed in [Burstall et al. 2004; Kamberov et al. 2002]. Recall that the quaternions \mathbb{H} are the 4-dimensional division algebra over \mathbb{R} generated by the set of symbols $\{i, j, k\}$ satisfying the relations $i^2 = j^2 = k^2 = ijk = -1$. It turns out that the additional algebraic structure provided by \mathbb{H} is advantageous for studying the conformal geometry of immersed surfaces in a variety of settings. In particular, any immersion $f : M \rightarrow \mathbb{R}^3$ of the two-dimensional surface M can be regarded as taking values in the imaginary part of the quaternions $\text{Im } \mathbb{H} \cong \mathbb{R}^3$, and therefore multiplication in \mathbb{H} (conventionally acting on the right) can be applied. This leads to a coordinate-free theory of surfaces which is adaptable, concise, and well suited for computational applications.

To explain the basic idea behind this, let (M, J) be a connected Riemann surface with (possibly empty) boundary ∂M and complex structure $J : TM \rightarrow TM$ satisfying $J^2 = -\mathbb{1}_{TM}$. Then, M carries a natural conformal structure inherited from the scalar product on each tangent space, and an immersion $f : (M, J) \rightarrow \mathbb{R}^3$ is conformal provided it maps oriented orthogonal bases of TM to oriented orthogonal bases of $df(TM)$. By standard arguments (e.g. [Burstall et al. 2004, Lemma 2]), this implies the existence of a unit quaternion field $N : M \rightarrow S^2 \subset \text{Im } \mathbb{H}$ uniquely compatible with the orientation on M that stabilizes $df(TM)$, i.e. such that $df(TM) = \{v \in \text{Im } \mathbb{H} \mid NvN = v\}$. It follows that N is the Gauss map of the immersion f , and moreover that f is conformal when and only when there exists an $N : M \rightarrow S^2$ satisfying

$$*df = N df,$$

where $*df := df \circ J$ is the negative Hodge star operator on differential one-forms. This provides a useful link between the intrinsic

notion of conformality and its extrinsic expression in terms of the surface Gauss map.

The key to the results presented here is a description of quasiconformal mappings in terms of this quaternionic framework. First, note that the Gauss map $N : M \rightarrow S^2$ of any conformal immersion $f : M \rightarrow \text{Im } \mathbb{H}$ defines a canonical complex structure on TM (c.f. [Burstall et al. 2004, Section 2.1]) which conventionally acts on the left. Indeed, it is straightforward to check that the mapping $J : TM \rightarrow TM$ defined by

$$(a + Jb)v = adf(v) + bNdf(v),$$

for all $a + Jb \in \mathbb{C}$ and $v \in TM$ is a complex structure on TM compatible with N .

Remark 3.1. In fact, every conformal structure J on TM can be realized by the Gauss map of some conformal immersion $f : M \rightarrow \text{Im } \mathbb{H}$ (see [Kamberov et al. 2002, pg 8]).

Moreover, notice that any complex structure J induces a direct sum decomposition of the space $\text{Hom}(TM, \mathbb{R}^3)$ of \mathbb{R}^3 -valued functions on TM analogous to classical complex analysis. In particular, given $\alpha \in \text{Hom}(TM, \mathbb{R}^3)$ there are conformal resp. anticonformal parts of α with respect to f ,

$$\alpha^+ = \frac{1}{2}(\alpha - N * \alpha), \quad \alpha^- = \frac{1}{2}(\alpha + N * \alpha),$$

and it is easily checked that $*\alpha^+ = N\alpha^+$ while $*\alpha^- = -N\alpha^-$. This leads to a natural definition of quasiconformality in the quaternionic setting.

Definition 3.2. Let $M \subset \mathbb{R}^3$ be an immersed surface with complex structure $J : TM \rightarrow TM$ realized by its Gauss map $N : M \rightarrow S^2$. Then, a mapping $f : M \rightarrow \text{Im } \mathbb{H}$ is said to be quasiconformal provided there exists a measurable, \mathbb{C} -antilinear endomorphism $\mu : TM \rightarrow TM$ which satisfies $|\mu|_\infty < 1$ and

$$df^- = df^+ \circ \mu.$$

Definition 3.2 is the quaternionic analogue of the standard definition for abstract Riemann surfaces (M, J) given in Section 2. On the other hand, the algebraic structure of \mathbb{H} enables an equivalent criterion which is more useful for practical computation and does not require explicit function composition. First, note that the Beltrami differential can be interpreted as a function $\mu : TM \rightarrow (TM)^\perp$ taking values in the normal bundle to TM in \mathbb{H} , which is isomorphic to the complex plane at every point (simply choose N as the imaginary unit). To see this, consider a local section $v \in TM$ and write $\mu : TM \rightarrow TM$ as

$$\mu(v) = \mu^1(v)v + \mu^2(v)Jv,$$

for some “coordinate” functions $\mu^i : TM \rightarrow \mathbb{R}$. Consequently, it follows that

$$\begin{aligned} df^+ \circ \mu(v) &= df^+ \left(\mu^1(v)v + \mu^2(v)Jv \right) \\ &= \mu^1(v)df^+(v) + \mu^2(v) * df^+(v) \\ &= \left(\mu^1(v) + \mu^2(v)N \right) df^+(v) = \tilde{\mu}(v)df^+(v), \end{aligned}$$

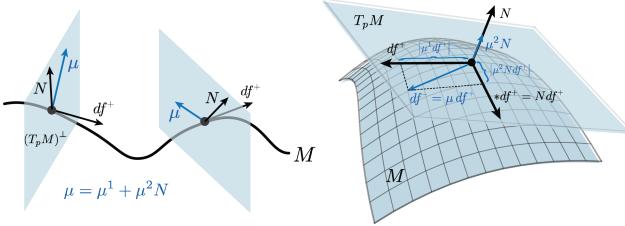


Fig. 7. Left: an illustration of μ as a normal-valued quaternion field. Right: an illustration of Definition 3.3.

where $\tilde{\mu} = \mu^1 + \mu^2 N$ is a $(TM)^\perp$ -valued function on TM . The reversibility of this argument establishes that normal-valued quaternionic functions $\tilde{\mu}$ satisfying $|\tilde{\mu}|_\infty < 1$ which transform appropriately (c.f. Lemma 3.4) correspond precisely to the Beltrami differentials of quasiconformal mappings, giving an alternative representation of these objects. This leads to the following equivalent to Definition 3.2 illustrated in Figure 7.

Definition 3.3. Let $M \subset \mathbb{R}^3$ be an immersed surface with complex structure $J : TM \rightarrow TM$ defined by its Gauss map $N : M \rightarrow S^2$. Then, a mapping $f : M \rightarrow \text{Im } \mathbb{H}$ is said to be quasiconformal provided there exists a measurable, normal-valued quaternion field $\mu : TM \rightarrow (TM)^\perp$ which satisfies $|\mu|_\infty < 1$ and

$$df^- = \mu df^+.$$

The equality in Definition 3.3 is suggestive of the coordinate-dependent expression $f_z = \mu f_{\bar{z}}$ seen before, but has the intuitive benefit of being expressed without reference to any underlying coordinate system. From this definition, it is straightforward to establish how $\mu : TM \rightarrow (TM)^\perp$ must transform under a change of basis for TM .

LEMMA 3.4. Let $a \in \mathbb{C}$ and suppose $\mu : TM \rightarrow (TM)^\perp$ is a Beltrami coefficient satisfying Definition 3.3. Then, for any local section $v \subset TM$ it follows that

$$\mu(av) = \frac{\bar{a}^2}{|a|^2} \mu(v).$$

PROOF. Due to the conformality (resp. anticonformity) of df^+ (resp. df^-), it follows that

$$\begin{aligned} a \mu(av) df^+(v) &= (\mu df^+)(av) = df^-(av) \\ &= \bar{a} df^-(v) = \bar{a} (\mu df^+)(v) = \bar{a} \mu(v) df^+(v). \end{aligned}$$

Therefore, $\mu(v) = (a/\bar{a}) \mu(av)$ and the conclusion follows. \square

In the literature on quasiconformal maps, a bounded measurable function μ that transforms in this way is said to be a “(-1,1)-form” on TM (c.f. [Hubbard 2006, Section 4.8]). In particular, given a local conformal coordinate z on a domain $U \subset M$, several authors write the suggestive expression $\mu = \hat{\mu} \frac{dz}{dz}$ for some local function $\hat{\mu} : U \rightarrow \mathbb{C}$ to illustrate the transformation rule seen in Lemma 3.4. More tensorially, any \mathbb{C} -antilinear $\mu : TM \rightarrow TM$ as in Definition 3.2 can be expressed locally as $\mu = \hat{\mu} d\bar{z} \otimes \frac{\partial}{\partial z}$, so that for any $a \in \mathbb{C}$,

$$\mu(av) = \hat{\mu} d\bar{z}(av) \otimes \frac{\partial}{\partial z} = \bar{a} \hat{\mu} d\bar{z}(v) \otimes \frac{\partial}{\partial z} = \bar{a} \mu(v).$$



Fig. 8. Another view of the Teichmüller extremal mapping from Figure 6 but with 95% deformation at the boundary. Note the visible structure of the quadratic differential governing the mapping, in particular the simple pole in front and simple zero on back which counteract to maintain zero Poincaré-Hopf index.

It is remarkable that the same transformation rule is readily established in the quaternionic setting without reference to any particular system of local coordinates.

Definition 3.3 provides a link between the intrinsic quasiconformal geometry of M and its extrinsic immersion in \mathbb{R}^3 which is useful for computing quasiconformal mappings between immersed surfaces. Before considering a minimization-based approach to accomplishing this and its consequences for QC Iteration, it is worth discussing the BC $\mu : TM \rightarrow (TM)^\perp$ in greater detail to gain some intuition about this new notion. In particular, let $(M, J) \subset \mathbb{R}^3$ be an immersed Riemann surface with complex structure J induced by its Gauss map $N : M \rightarrow S^2$ as before. Then, the differential of any other immersion $f : M \rightarrow \text{Im } \mathbb{H}$ splits as $df = df^+ + df^-$ with respect to J , so that the induced metric on M is given by

$$f^* \delta = |df|^2 = |df^+|^2 + |df^-|^2 + 2 \operatorname{Re} (df^+ \overline{df^-}),$$

where it was used that $\langle v, w \rangle_{\mathbb{H}} = (1/2)(v\bar{w} + w\bar{v}) = \operatorname{Re}(v\bar{w})$. The $(2,0)$ -part of this quantity is the classical *Hopf differential* of the mapping f ,

$$Q = df^+ \overline{df^-},$$

which quantifies deviation from conformality with respect to J . Interestingly, expanding the positive and negative parts of this expression yields the alternative representation

$$4Q = |df|^2 - |*df|^2 - 2 \langle df, *df \rangle N,$$

showing that the Hopf differential is normal-valued, i.e. locally expressible as $Q(v) = a(v) + b(v)N$ for some smooth functions $a, b : TM \rightarrow \mathbb{R}$. Moreover, it is clear from this expression that f is conformal or anticonformal if and only if $Q \equiv 0$. Conversely, when f is quasiconformal in the sense of Definition 3.3, simple algebraic manipulations show that $Q = \bar{\mu} |df^+|^2$ or alternatively

$$\mu = *Q,$$

demonstrating that the Beltrami differential $\mu : TM \rightarrow (TM)^\perp$ is simply the conjugate of the Hopf differential up to an application of the Hodge star. Indeed, the structure of this object is often visible in

the mapping (see e.g. Figure 8), where poles and zeroes necessarily appear according to the form of Q .

Remark 3.5. The relationship between μ and Q discussed here can also be considered a (rearranged) special case of the isomorphism $\tau \mapsto df^+ \tau$ between tangential-valued anticonformal one-forms and normal-valued quadratic differentials induced by the conformal immersion $df^+ : M \rightarrow \mathbb{R}^3$ (see [Kamberov et al. 2002][Section 1.8]).

4 MINIMIZING THE QUASICONFORMAL DISTORTION

It is now possible to discuss the present method for computing quasiconformal mappings, which is based on the quaternionic Definition 3.3. While the notion of quasiconformality in Definition 3.3 can be extended to cover the case of mappings $f : M \rightarrow P$ between fixed Riemann surfaces $M, P \subset \mathbb{R}^3$, the present discussion will focus on mappings $f : M \rightarrow f(M) \subset \mathbb{R}^3$ whose target is determined by the mapping itself, as this is the case most relevant for the remeshing and boundary deformation applications considered in Section 6 (see e.g. Figure 9). In this setting, the quasiconformal distortion has the expression,

$$QC_\mu(f) = \int_M |df^- - \mu df^+|^2 dS_g,$$

which (by Theorem A.1 in Appendix A) is related to the conformal part of the Dirichlet energy with respect to the metric $g(\mu)$ on M inherited from the BC μ . Since $d(\cdot)^-, d(\cdot)^+$, and the Hodge star are linear operators (for fixed N), this is a convex function of f (for fixed μ) whose global minimizer is a least-squares quasiconformal mapping approximately satisfying Definition 3.3. Therefore, quasiconformal mappings can be computed from BCs $\mu : TM \rightarrow (TM)^\perp$ by solving the minimization problem

$$\operatorname{argmin}_{f \in [f_0]} QC_\mu(f), \quad f|_{\partial M} = y,$$

where $y : \partial M \rightarrow \mathbb{R}^3$ is a prescribed function representing the desired boundary data and where $f_0 : M \rightarrow \mathbb{R}^3$ is a suitable mapping whose boundary is y . Expressing the unknown mapping f perturbatively in terms of a one-parameter family $f(x, t) = f(x) + t\varphi(x)$ for $t \in (-\varepsilon, \varepsilon)$ and some compactly supported variation $\varphi : M \rightarrow \mathbb{R}^3$ vanishing on the boundary ∂M , standard techniques from the calculus of variations yield a necessary and sufficient criterion for f to be a minimizer of QC_μ . In particular, using the fact that μ , the conformal structure, and the metric g are fixed, the derivative of the functional QC_μ at f in the direction φ is given by

$$\begin{aligned} \delta QC_\mu(f)\varphi &:= \frac{d}{dt} \Big|_{t=0} QC_\mu(f + t\varphi) \\ &= \int_M \langle df^- - \mu df^+, d\varphi^- - \mu d\varphi^+ \rangle dS_g, \end{aligned}$$

which can be written as a linear system of equations for the minimizer f once an appropriate basis of test functions is given. On the other hand, notice that the mapping f satisfying the above minimization does not constrain the surface area or extrinsic geometry of the image surface $f(M)$ in \mathbb{R}^3 , meaning the surface $f(M)$ may be quite different qualitatively than M (see e.g. Figure 9). While this is not necessarily a problem, it is often desirable (e.g. when remeshing) to have constraints on f which keep the target surface



Fig. 9. Two qualitatively different quasiconformal mappings (left, middle) from the torus with boundary (right) computed by minimizing QC_μ with Algorithm 1 with different boundary conditions. See Figure 15 in Section 6 for a quantitative view.

close to the original in some sense. Therefore, the remainder of this Section discusses two constraints which are useful for this purpose along with a finite element algorithm for computing least-squares quasiconformal mappings.

4.1 Preserving extrinsic geometry

As mentioned, it is often useful to compute quasiconformal mappings which preserve the extrinsic geometry of the source surface $M \subset \mathbb{R}^3$. Particularly when remeshing, a quasiconformal map $f : (M, g_{\text{ref}}) \rightarrow (M, g)$ is desired which optimizes mesh element angles while preserving the original Gauss map N as well as possible. Certainly, preserving N exactly is not possible unless f is conformal, although reasonably good agreement can be enforced in any case by using an appropriate constraint alongside the minimization of QC_μ . One effective option for this was proposed in [Gruber and Aulisa 2020] for the purposes of preserving extrinsic features while minimizing the conformal distortion, which is equivalent to QC_μ when $\mu \equiv 0$. To describe this idea specifically, consider a point $x \in M$, an immersion $f : M \rightarrow \mathbb{R}^3$, and a curve $f(x + tv)$ for some $v \in T_x M$ and $|t| < \varepsilon$. This curve lies on the immersed surface $f(M)$, and letting N be normal to the image $f(M)$, Taylor expansion around $t = 0$ yields

$$\begin{aligned} f(x + tv) &= f(x) + t df(v) + \frac{t^2}{2} (\nabla df)(v, v) + O(t^3), \\ N\left(x + \frac{t}{2} v\right) &= N(x) + \frac{t}{2} dN(v) + O(t^2), \end{aligned}$$

where $df, \nabla df, dN$ are localized at x . Since $\langle N, df(v) \rangle = 0$ pointwise for all $v \in TM$, it follows by differentiation that

$$\langle dN(v), df(v) \rangle + \langle N, (\nabla df)(v, v) \rangle = 0,$$

and therefore the difference vector $f(x + tv) - f(x)$ satisfies

$$\left\langle f(x + tv) - f(x), N\left(x + \frac{t}{2} v\right) \right\rangle = 0 + O(t^3).$$

This means that the inner product of the difference vector $f(x + tv) - f(x)$ with the normal vector in the middle vanishes to second order, inspiring a constraint for the shape-controlled minimization of QC_μ . In particular, to keep the image of f close to the reference



Fig. 10. Optimal quasiconformal remeshings of a cat with genus 1 and 19 boundary components with respect to an equi-angular conformal structure (c.f. Section 6.2). The original surface (middle) is mapped through QC Iteration with preserved extrinsic geometry (right) and preserved area (left). See Figure 19 in Section 6 for a quantitative view.

immersion $\mathbb{1}_M$ it is reasonable to require that

$$\langle f(\mathbf{x}) - \mathbb{1}_M(\mathbf{x}), N_{mid}(\mathbf{x}) \rangle = 0,$$

where $N_{mid} = (1/2)(N_{old} + N_{new})$ approximates the normal halfway between the image of $\mathbb{1}_M$ and the image of f . This ensures that the difference vectors between the current points on M and $f(M)$ remain nearly tangential to the implicit surface in the middle, and can be implemented as a nonlinear constraint during the minimization of QC_μ using a Lagrange multiplier.

For the purposes of this work, it is enough to use $N = N_{old}$, giving a linearized version of this constraint and preserving the convexity of the problem. In particular, consider finding a pair of functions $v : M \rightarrow \mathbb{R}^3$ and $\rho : M \rightarrow \mathbb{R}$ which satisfy

$$\operatorname{argmin}_{v, \rho} \left(QC_\mu (\mathbb{1}_M + v) + \int_M \rho \langle v, N \rangle dS_g + \frac{\epsilon}{2} \int_M \rho^2 dS_g \right),$$

where $\epsilon > 0$ is a fixed penalty parameter and $v|_{\partial M} = 0$. Then, $f = \mathbb{1}_M + v$ minimizes QC_μ and satisfies the desired constraint. Formulated weakly, desired pair v, ρ should satisfy the system

$$\begin{aligned} 0 &= \delta QC_\mu (\mathbb{1}_M + v) \varphi + \int_M \rho \langle \varphi, N \rangle dS_g, \\ 0 &= \int_M \psi \langle v, N \rangle dS_g + \epsilon \int_M \psi \rho dS_g, \end{aligned}$$

for all suitable variations $\varphi : M \rightarrow \mathbb{R}^3$ and $\psi : M \rightarrow \mathbb{R}$. The mapping $f = \mathbb{1}_M + v$ which minimizes this modified problem is then quasiconformal with normal field N suitably close to that of the original surface M . Figure 10 and others in Section 6 show that this procedure is suitable for producing quasiconformal maps which preserve extrinsic features even around corners and delicate contours in the mesh such as facial expressions.

Remark 4.1. Here, the penalty parameter $\epsilon \approx 10^{-5}$ is used to ensure L^2 -regularity for the Lagrange multiplier ρ .

4.2 Preserving surface area

Recall that the minimizing the quasiconformal distortion QC_μ in the metric g is equivalent to minimizing the Dirichlet energy $\mathcal{D}_{g(\mu)}$ in the metric $g(\mu)$ when the area of the target is fixed (see Appendix A). Since this area is variable when the target surface is defined implicitly through the mapping f , it is useful to have a constraint which fixes this quantity. This provides a reliable way to produce constrained harmonic maps which often improves the numerical results, especially in the case of the QC Iteration seen in Section 3. The present constraint follows easily from the variation of the surface area functional \mathcal{A} , defined as

$$\mathcal{A}(f) = \int_M 1 dS_{f^*\delta} = \int_M \frac{N}{2} df \wedge \overline{df}.$$

Standard techniques give the derivative of \mathcal{A} at f in the direction φ ,

$$\delta \mathcal{A}(f)\varphi = \int_M \langle df, d\varphi \rangle dS_{f^*\delta},$$

giving a nonlinear constraint that can be implemented alongside the minimization of QC_μ . On the other hand, the applications in Section 6 will almost always consider $f = \mathbb{1}_M + v$ for some perturbation $v : M \rightarrow \mathbb{R}^3$ as before. To retain convexity of the minimization in this case, it is reasonable to consider a heuristic constraint which is linear and works well in practice. To that end, recall that the area element $dS_{f^*\delta} = |df|^2$ when f is conformal (see Appendix A), so that the area of M can be expressed as

$$A_0 := \mathcal{A}(\mathbb{1}_M) = \int_M |d\mathbb{1}_M|^2.$$

It follows that the difference in area between $f(M)$ and M can be approximated by quantities depending only on $g = \mathbb{1}_M^* \delta$,

$$\mathcal{A}(f) - A_0 = \int_M 1 dS_{f^*\delta} - \int_M |d\mathbb{1}_M|^2 \approx \int_M \langle d\mathbb{1}_M, dv \rangle dS_g.$$

With this, consider computing functions $f : M \rightarrow \mathbb{R}^3$ and $\lambda : M \rightarrow \mathbb{R}$ which satisfy

$$\operatorname{argmin}_{v, \lambda} \left(QC_\mu (\mathbb{1}_M + v) + \lambda \int_M \langle d\mathbb{1}_M, dv \rangle dS_g \right),$$

where $v|_{\partial M} = 0$ and $A_0 = \mathcal{A}(\mathbb{1}_M)$. This is a convex minimization problem whose global solution can again be computed from weak-form equations. In particular, the desired pair v, λ should satisfy the system

$$\begin{aligned} 0 &= \delta QC_\mu (\mathbb{1}_M + v) \varphi + \lambda \int_M \langle d\mathbb{1}_M, d\varphi \rangle dS_g, \\ 0 &= \psi \int_M \langle d\mathbb{1}_M, dv \rangle dS_g, \end{aligned}$$

for all suitable variations $\varphi : M \rightarrow \mathbb{R}^3$ and $\psi : M \rightarrow \mathbb{R}$. The unique mapping f satisfying this system is both a minimizer of QC_μ and a constrained (approximately) harmonic mapping in the metric $g(\mu)$. **WHY do we have to freeze one point when the surface is closed?????** Need to mention this.

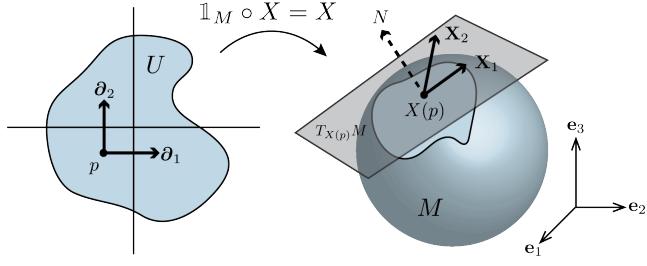


Fig. 11. The parametrization X coordinatizes a portion of the surface M . The vector fields ∂_i form a basis for $T_p U$, while their images $X_i = dX(\partial_i)$ form a basis for $T_{X(p)} M$.

4.3 Minimization Algorithm

The final goal of this Section is to describe in detail the present algorithm for computing minimizers of QC_μ in the framework of piecewise linear finite elements. Here it is assumed that the source surface $M_h \approx M$ is given as a structured or unstructured orientable manifold mesh of triangles or quadrilaterals which are not degenerate. This implies that the interior angles of each element must be bounded below, although this bound may be very close to zero in practice as in the case of Figures 1 and 19. Under this assumption, M_h can be expressed as the union

$$M_h = \bigcup_{T \in \mathcal{T}} T,$$

where each T is a polygonal element of the discrete surface. Additionally, it will be assumed that the vertices of M_h are embedded in \mathbb{R}^3 , so that $M_h \subset \mathbb{R}^3$ carries a submanifold structure via the usual inclusion mapping $1_{M_h} : M_h \rightarrow \mathbb{R}^3$. In practice, it is convenient to coordinatize M_h with local parametrizations $X_h : U_h \rightarrow M_h$ so that $1_{M_h} \circ X_h = X_h$ and each element T is the image of a reference (or parent) polygon $\hat{T} \subset \mathbb{R}^2$. This implies that functions on M_h may be discretized using a piecewise-linear polynomial basis (or tensor product of these) supported on each reference element (see e.g. [Dziuk and Elliott 2013, Section 4.3-4.4]). In particular, if the preimages through X_h of the n nodes of M_h are denoted by $\{\mathbf{v}_\alpha\}_{\alpha=1}^n$, the standard Lagrange nodal basis $\{\phi_\alpha\}_{\alpha=1}^n$ on U_h satisfies $\phi_\alpha(v_\beta) = \delta_{\alpha\beta}$ for all $1 \leq \alpha, \beta \leq n$. The space of piecewise-linear finite elements on U_h can then be expressed as

$$\text{Span}\{\phi_\alpha\} = \{\phi \in C^0(U_h) : \phi|_{\hat{T}} \in \mathbb{F}_1(\hat{T}), T \in \mathcal{T}\},$$

where \mathbb{F}_1 denotes the space of linear polynomials \mathbb{P}_1 when T is triangular and the tensor product space \mathbb{Q}_1 when T is quadrangular. Note that this defines an analogous (nonlinear) basis $\varphi = \phi \circ X_h$ for functions on the surface M_h .

Remark 4.2. To simplify the presentation, the remainder of the manuscript assumes the Einstein summation convention. Therefore, any tensor index appearing twice in an expression (once up and once down) is implicitly summed over its appropriate range. Similarly, the subscript h on discrete quantities will be omitted when the context is clear.

To describe the present discretization more precisely, suppose the smooth surface M is coordinatized by local parameterizations,

each of which looks like $X : U \rightarrow M$ for some domain $U \subset \mathbb{R}^2$ (see Figure 11). It follows that if $\{x^1, x^2\}$ are coordinates on U , then $\{\partial_1, \partial_2\}$ form the standard basis for TU where $\partial_i := \partial/\partial x^i$. Moreover, using the dual basis defined by $dx^i(\partial_j) = \delta_j^i$ for $1 \leq i, j \leq 2$ the differential of X has the expression $dX = X_i \otimes dx^i$ (Einstein summation assumed) where

$$X_i = X_i^j \partial_j, \quad X_i^j = \frac{\partial X^j}{\partial x^i}.$$

Clearly, this coordinate representation lies on the abstract surface M and need not exist in \mathbb{R}^3 . On the other hand, post-composition with the inclusion map $1_M : M \rightarrow \mathbb{R}^3$ yields $1_M \circ X = X$ enabling an alternative expression of the vectors X_i in terms of the standard basis $\{\mathbf{e}_I\}_{I=1}^3$ for \mathbb{R}^3 ,

$$X_i = X_i^J \mathbf{e}_J, \quad X_i^J = \frac{\partial X^J}{\partial x^i}.$$

This identification also provides local expressions on U of the Riemannian metric g and area element dS_g on $M \subset \mathbb{R}^3$, as well as the (outward-directed) unit normal field to M , given componentwise as

$$\begin{aligned} g_{ij} &= \langle \mathbf{X}_i, \mathbf{X}_j \rangle, \\ dS_g &= \sqrt{\det g} dx^1 \wedge dx^2, \\ N &= \frac{\mathbf{X}_1 \times \mathbf{X}_2}{|\mathbf{X}_1 \times \mathbf{X}_2|}. \end{aligned}$$

Letting $\{\omega^i\}$ be the dual basis to $\{\mathbf{X}_i\}$ and g^{ij} be the components of the metric inverse defined by $g^{ik} g_{kj} = \delta_j^i$, the differential of any function $f : M \rightarrow \mathbb{R}^3$ can be locally expressed on $TM|_{X(U)} \subset T\mathbb{R}^3$ as

$$df = F_i \otimes \omega^i = g^{ij} F_i \otimes \mathbf{X}_j.$$

Pulled back to TU through $F := f \circ X$, the vector fields $F_i = df(X_i) = dF(\partial_i)$ then have an expression in terms of the same standard basis $\{\mathbf{e}_J\}$,

$$F_i = F_i^J \mathbf{e}_J, \quad F_i^J = \frac{\partial (f \circ X)^J}{\partial x^i}.$$

These expressions enable a local description of the quasiconformal distortion QC_μ on each parameterization domain U . First, note that the discrete BC $\mu : TM_h \rightarrow (TM_h)^\perp$ can be conveniently stored as a single complex-valued function $\mu(\mathbf{X}_1)$, since Lemma 3.4 allows for the computation of any other $\mu(\mathbf{v})$ from this value by expressing $\mathbf{v} = a\mathbf{X}_1 + bJ\mathbf{X}_2$ for some $a, b \in \mathbb{R}$. In particular, it is convenient to define $\mu_F := \mu(\mathbf{X}_1) = \mu \circ dX(\partial_1)$ giving a representation of μ on each reference element. Since it may be assumed that X is conformal, it follows that $\mu(\mathbf{X}_2) = -\mu(\mathbf{X}_1) = -\mu_F$ (easily checked with Lemma 3.4), and the difference $df^- - \mu_F df^+$ applied to the basis vector fields $\{\mathbf{X}_i\}$ has the expression

$$\begin{aligned} (D_1 f^I) \mathbf{e}_I &:= df^-(\partial_1) - \mu \circ dX(\partial_1) df^+(\partial_1) \\ &= \left((1 - \mu_F) F_1^I + (1 + \mu_F) N F_2^I \right) \mathbf{e}_I, \\ (D_2 f^I) \mathbf{e}_I &:= df^-(\partial_2) - \mu \circ dX(\partial_2) df^+(\partial_2) \\ &= \left((1 + \mu_F) F_2^I - (1 - \mu_F) N F_1^I \right) \mathbf{e}_I, \end{aligned}$$

where the quantities μ_F, N, \mathbf{e}_I are all quaternionic. Using $\langle \mathbf{v}, \mathbf{w} \rangle = \text{Re}(\mathbf{v}\overline{\mathbf{w}})$ along with $\overline{\mathbf{v}\mathbf{w}} = \overline{\mathbf{w}}\overline{\mathbf{v}}$ and the fact that the \mathbf{e}_I are pure imaginary yields a local expression for the integrand of QC_μ ,

$$\begin{aligned} |df^- - \mu df^+|^2 &= g^{ij} \left((D_i f^K) \mathbf{e}_K, (D_j f^L) \mathbf{e}_L \right) \\ &= -g^{ij} \text{Re} \left(D_i f^K \mathbf{e}_{K \times L} \overline{D_j f^L} \right), \end{aligned}$$

where it was convenient to introduce the notation

$$\mathbf{e}_{K \times L} = \begin{cases} \mathbf{e}_K \times \mathbf{e}_L & K \neq L, \\ -1 & K = L. \end{cases}$$

Putting this together, the quasiconformal distortion of the local image $F_h(U_h)$ for any $U_h \subset M_h$ has the representation

$$\begin{aligned} QC_\mu(f) &= \int_{X(U)} |df^- - \mu df^+|^2 dS_g \\ &= \int_U -g^{ij} \text{Re} \left(D_i f^K \mathbf{e}_{K \times L} \overline{D_j f^L} \right) \sqrt{\det g} dx^1 \wedge dx^2, \end{aligned}$$

which is computed elementwise and summed to give the distortion of the mapping f on the entirety of M_h .

Generating linear systems for the minimization problems discussed in this section is now straightforward. Recall that $f : M_h \rightarrow \mathbb{R}^3$ can be represented component-wise in the nodal basis $\{\phi^\alpha\}$ for U_h as

$$F^K = (f \circ X)^K = F_\alpha^K \phi^\alpha,$$

where F_α^K denotes the value of the K^{th} component of F at (global) node α . It follows that the derivative of f in the direction X_i is then

$$df(X_i) = F_i^K \mathbf{e}_K = F_\alpha^K \phi_i^\alpha \mathbf{e}_K = \mathbf{F}_\alpha \phi_i^\alpha,$$

so that the variation of $QC_\mu(f)$ can be expressed as the linear system ($U = \cup U_h$)

$$\begin{aligned} \mathcal{J}_K^{\alpha\beta} F_\alpha^K &:= \delta QC_\mu(f) \varphi^\beta \\ &= F_\alpha^K \int_U -2g^{ij} \text{Re} \left(D_i \phi^\alpha \mathbf{e}_{K \times L} \overline{D_j \phi^L} \right) \sqrt{\det g} dx^1 \wedge dx^2, \end{aligned}$$

where \mathcal{J} is the Jacobian operator, and the vector test function $\varphi^\beta = \varphi^L \mathbf{e}_L$ is simply three independent copies of the nodal vector φ^β , since the test functions for each coordinate are the same. Expressing $f = \hat{f} + f_\alpha \varphi^\alpha$ for some known function $\hat{f} : M \rightarrow \mathbb{R}^3$ (such as $\hat{f} = \mathbb{1}_M$) then yields the linear system

$$\mathcal{J} \cdot S = -\mathcal{R}, \quad \mathcal{R}^\beta = \mathcal{J}^{\alpha\beta} \cdot \hat{\mathbf{F}}_\alpha = \mathcal{J}_K^{\alpha\beta} \hat{F}_\alpha^K,$$

which is readily solved for the nodal values F_α^K . The desired mapping $F_h = f_h \circ X_h$ is then reconstructed from these values by linear interpolation over the mesh. Note that the constraints from Sections 4.1 and 4.2 can be discretized analogously and included along with this procedure, leading to the least-squares quasiconformal mapping procedure described in Algorithm 1. The examples from this paper were implemented in the open source finite element library FEMuS [Aulisa et al. 2014], and the Boost library [Schäling 2011] was used to compute the relevant quaternion products.

Remark 4.3. Since Algorithm 1 is formulated for mappings with an implicit target, it is not immediately obvious if it also has advantages in the more traditional setting of mappings between fixed Riemann surfaces $M \mapsto P$ both given as e.g. manifold meshes in \mathbb{R}^3 . On the

other hand, since the geometry of P is known, constraints like those introduced here can certainly be implemented to keep the image of M close to P in an appropriate sense. This is essentially what is done in Figure 8 to keep the surface from “spilling out” of the original cylinder.

Algorithm 1 Computation of f given μ

Require: Surface $M \subset \mathbb{R}^3$, mapping $\hat{f} : M \rightarrow \mathbb{R}^3$. Beltrami coefficient μ .

- 1: Let $f = \hat{f} + f_\alpha \varphi^\alpha$, so that $F = \hat{F} + F_\alpha \phi^\alpha$.
- 2: Compute the Jacobian of QC_μ ,

$$\tilde{\mathcal{J}}_K^{\alpha\beta} = -2 \int_M g^{ij} \text{Re} \left(D_i \phi^\alpha \mathbf{e}_{K \times L} \overline{D_j \phi^L} \right),$$

where \int_M is shorthand for $\int_U \sqrt{\det g} dx^1 \wedge dx^2$.

- 3: **if** Normal constraint **then**
- 4: Initialize $\rho = \hat{\rho} + \rho_\alpha \psi^\alpha$.
- 5: Compute the block 2×2 Jacobian of the constraint

$$\mathcal{J} = \begin{pmatrix} \mathcal{J}_{11} & \mathcal{J}_{12} \\ \mathcal{J}_{21} & \mathcal{J}_{22} \end{pmatrix} = \begin{pmatrix} \tilde{\mathcal{J}} & \int_M \psi^\alpha \langle \phi^\beta, N \rangle \\ \int_M \psi^\beta \phi^\alpha N & \varepsilon \int_M \psi^\alpha \psi^\beta \end{pmatrix},$$

- 6: Form the residual vector
- 7: $\mathcal{R}^\beta = \mathcal{J}^{\alpha\beta} \cdot \hat{\mathbf{S}}_\alpha, \quad \hat{\mathbf{S}}_\alpha = (\hat{\mathbf{F}}_\alpha \quad \hat{\rho}_\alpha)^\top$.
- 8: Solve the linear system
- 9: $\mathcal{J} \cdot S = -\mathcal{R}, \quad S_\alpha = (\mathbf{F}_\alpha \quad \rho_\alpha)^\top$.
- 10: **else if** Area constraint **then**
- 11: Initialize $\lambda = \hat{\lambda}$.
- 12: Compute the block 2×2 Jacobian of the constraint

$$\mathcal{J} = \begin{pmatrix} \mathcal{J}_{11} & \mathcal{J}_{12} \\ \mathcal{J}_{21} & \mathcal{J}_{22} \end{pmatrix} = \begin{pmatrix} \hat{\lambda} \int_M g^{ij} \phi_i^\alpha \phi_j^\beta & \int_M g^{ij} \langle \hat{\mathbf{F}}_i, \phi_j^\beta \rangle \\ \int_M g^{ij} \phi_j^\alpha \hat{\mathbf{F}}_i & 0 \end{pmatrix}.$$

- 13: Augment $\mathcal{J}_{11} += \tilde{\mathcal{J}}$.
- 14: Form the residual vector
- 15: $\mathcal{R}^\beta = \mathcal{J}^{\alpha\beta} \cdot \hat{\mathbf{S}}_\alpha, \quad \hat{\mathbf{S}}_\alpha = (\hat{\mathbf{F}}_\alpha \quad \lambda)^\top$.
- 16: Solve the linear system
- 17: $\mathcal{J} \cdot S = -\mathcal{R}, \quad S_\alpha = (\mathbf{F}_\alpha \quad \lambda)^\top$.
- 18: **end if**
- 19: **return** $F = f \circ X$.

5 QUATERNIONIC QC ITERATION

Algorithm 1 provides a novel technical tool for computing quasiconformal mappings between immersed surfaces, and is applicable when the target is explicit or implicit (see Remark 4.3). As an interesting application of this, consider how the QC Iteration from Section 2.2 can be applied to generate Teichmüller quasiconformal maps between surfaces. First, note that all the theory regarding extremal

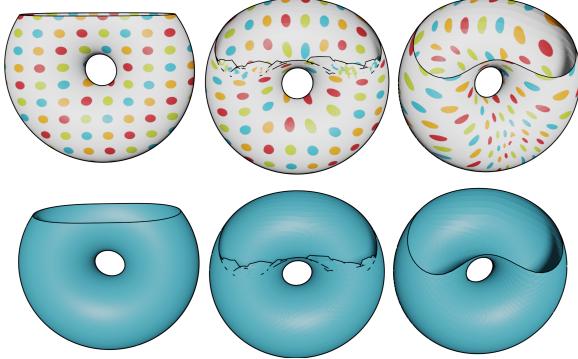


Fig. 12. Comparison of LSCM (mid) and TQCM (right) from a torus with boundary (left) to a similar surface with different boundary data. Here again the LSCM does not remain injective despite its much lower average distortion. Conversely, the TQCM produces a natural looking map with near-uniform texture distortion (99% of elements with $|\mu| \in [0.4, 0.5]$)

quasiconformal mappings discussed in Section 2.1 translates immediately to the quaternionic setting of maps $f : M \rightarrow f(M) \subset \mathbb{R}^3$. In particular, extremal Teichmüller mappings exist in every homotopy class of maps relative to ∂M , and these mappings are associated to the global minimizers of the energy $\mathcal{E}(\mu)$.

As previously mentioned, the QC iteration algorithm from [Lui et al. 2015, 2014] is an approach to computing Teichmüller mappings through the relationship of the quasiconformal distortion QC_μ to the energy \mathcal{E} . More precisely, the goal is to use minimizers of QC_μ (for fixed μ and fixed homotopy class $[f_0]$) to construct a sequence of BCs $\{\mu_k\}_{k=0}^{n_t}$ which converges to a function μ_{n_t} that is approximately Teichmüller, i.e. that approximately satisfies Definition 2.1. Accomplishing this requires a way to compute quasiconformal mappings f given μ , along with a procedure which encourages μ to have the correct local structure at each iteration k . It will now be shown that the quaternionic technology from Section 3 and the least-squares quasiconformal mapping Algorithm 1 from Section 4 enable the modified QC Iteration outlined in Algorithm 2, which directly computes an approximately Teichmüller quasiconformal mapping $f : M \rightarrow \mathbb{R}^3$ satisfying prescribed Dirichlet boundary data. The remainder of this Section discusses each step of this algorithm in detail.

Remark 5.1. It is important to mention that the quaternionic QC Iteration Algorithm 2 cannot approximate Teichmüller mappings to arbitrary accuracy when M is not simply connected. This is a consequence of the local phase-based smoothing referred to in Remark 2.5 and discussed in Section 5.3. In this case, the QC Iteration will optimize for a BC with approximately constant norm which may or may not be close to an actual Teichmüller mapping. Regardless, practical results illustrate that high-quality quasiconformal mappings between multiply connected surfaces are achieved with Algorithm 2 despite this concern, e.g. Figures 12, 9, 8 and others in Section 6.

Algorithm 2 Overview of the quaternionic QC Iteration

Require: Surface $M \subset \mathbb{R}^3$ and homotopy class $[f]$. Beltrami coefficient $\mu_0 = 0$. Stopping tolerance $\varepsilon > 0$ and maximum iteration number $n_t > 0$. [check this](#)

- 1: **while** $0 \leq k \leq n_t$ and $QC_{\mu_k}(f_k) > \varepsilon$ **do**
- 2: (1) Minimize $QC_{\mu_k}(f)$ for $f_k : M \rightarrow \mathbb{R}^3$ given μ_k .
- 3: (2) Compute μ_{k+1} algebraically given f_k .
- 4: (3) Post-process μ_{k+1} to bring it closer to Teichmüller form.
- 5: (4) Minimize $QC_\mu(f_k)$ for μ on the line between μ_k and μ_{k+1} , generating $\mu_{k+1} \leftarrow \mu$.
- 6: **end while**
- 7: **return** (f, μ)

5.1 Step 1: Computing Quasiconformal Mappings

Given a BC $\mu : TM \rightarrow (TM)^\perp$, the first task is to compute a corresponding map $f : M \rightarrow \mathbb{R}^3$ which minimizes the quasiconformal distortion QC_μ . The present method for accomplishing this was discussed at length in Section 4 and culminated in Algorithm 1, which is a linear procedure for computing a least-squares quasiconformal mapping when μ is stored as a value per-element aligned with the vector field $X_1 = dX(\partial_1)$. It is worth mentioning that the mapping computed in this process will generally not be harmonic in the metric $g(\mu)$ unless the target surface is made explicit or an area constraint such as that discussed in Section 4.2 is enforced. While this has consequences for the theory developed in [Lui et al. 2015], it does not seem to hinder the performance of the QC Iteration in practice.

5.2 Step 2: Computing the Beltrami Coefficient

Once a particular quasiconformal mapping $f : M \rightarrow \mathbb{R}^3$ has been determined, it is necessary to compute its Beltrami coefficient μ in the complex structure determined by the Gauss map N . This reveals another substantial benefit of the quaternionic approach from Section 3: since $\mathbb{R}^3 \cong \text{Im } \mathbb{H}$, computing the desired function $\mu : TM \rightarrow (TM)^\perp$ as in Definition 3.3 becomes a simple algebraic exercise captured by the follow Lemma.

LEMMA 5.2. *Let $f : M \rightarrow \text{Im } \mathbb{H}$ be a quasiconformal mapping with respect to the complex structure induced by $N : M \rightarrow S^2$. Then, the Beltrami coefficient $\mu : TM \rightarrow (TM)^\perp$, $\mu = \mu^1 + \mu^2 N$ can be expressed component-wise as*

$$\mu^1 = \text{Re} \left(\frac{df^- \overline{df^+}}{|df^+|^2} \right), \quad \mu^2 = \text{Re} \left(\frac{df^- N \overline{df^+}}{|df^+|^2} \right),$$

PROOF. Since $df^- = \mu df^+$, the definition of the complex structure on TM induced by N yields

$$\frac{df^- \overline{df^+}}{|df^+|^2} = \mu^1 + \mu^2 N.$$

The desired representation now follows from the fact that $\bar{vw} = \bar{w}\bar{v}$ for quaternions $v, w \in \mathbb{H}$. \square

This result provides a quick and straightforward way to extract the BC μ of the least-squares minimizer f from Step 1 and is given in pseudocode as Algorithm 3. On the other hand, recall that the

components of μ are not invariant under changes of basis for TM , which has consequences for the present implementation since $\mu_F = \mu(\mathbf{X}_1)$ is stored as a complex function per-element. In particular, the local calculations in Step 3 will require averaging μ_F between mesh elements, making it necessary to compare values of μ which are discretized using different basis vectors \mathbf{X}_1 . This requires careful consideration of Lemma 3.4, which is the analogue of coordinate-dependence for quaternionic BCs.

Algorithm 3 Computation of $\mu_F = \mu(\mathbf{X}_1)$ from $f : M \rightarrow \mathbb{R}^3$.

Require: Quasiconformal mapping $f : M \rightarrow \mathbb{R}^3$. Local parameterizations $X : U \rightarrow M$ such that $F = f \circ X$.

- 1: Compute $\mathbf{v}^+ = dF^+(\partial_1)$ and $\mathbf{v}^- = dF^-(\partial_1)$.
- 2: Compute the BC

$$\mu_F = \frac{\mathbf{v}^+ \bar{\mathbf{v}}^+}{|\mathbf{v}^+|^2}.$$

- 3: **return** $\mu_F^1 = \operatorname{Re}(\mu_F)$ and $\mu_F^2 = -\operatorname{Re}(\mu_F N)$.

5.3 Step 3: Post-processing the Beltrami Coefficient

Once the BC μ has been computed from the most recent mapping f , the final step in the original QC iteration is to perturb it in a way which lowers its Beltrami energy $\mathcal{E}(\mu)$. More precisely, a combination of projection on the norm $|\mu|$ along with local Laplace smoothing on the phase $e^{i\arg\mu}$ is performed in order to produce a BC which is closer to locally optimal as defined by Lemma 2.4. To accomplish this in the quaternionic setting, it remains to discuss the details of the smoothing operations described in Section 2.2. First, recall that the BC μ corresponding to the f from Step 1 is stored as a per-element value $\mu_F = \mu(\mathbf{X}_1)$ which is aligned with the image of the basis vector ∂_1 coming from the bottom edge of each reference element. From this data, it is straightforward to extract the polar decomposition $\mu_F = |\mu_F|(\mu_F/|\mu_F|)$, giving the per-element norm and phase of the BC which can be manipulated separately to decrease the energy \mathcal{E} . In particular, since Teichmüller BCs $\mu = k\bar{q}/|q|$ have constant norm and harmonic argument away from the zeros of the quadratic differential q by Lemma 2.4, it is desirable to perturb these quantities in a way which encourages this structure. In the case of the norm $|\mu_F|$, this is not difficult, and the present method comes immediately from [Lui et al. 2014]. If $n_T := |T|$ is the number of mesh elements, the average of the $|\mu_F|$ is computed across the mesh,

$$|\mu|_{\text{avg}} = \frac{1}{n_T} \sum_{F=1}^{n_T} |\mu_F|,$$

which then replaces the value of $|\mu_F|$ on each individual element of the mesh (c.f. Algorithm 4). Seen abstractly, this can be considered a projection operation $\mathcal{B} \mapsto \mathcal{T}$ (c.f. Section 2.2) to the space of constant-norm BCs.

The local smoothing operation performed on the phase $e^{i\arg\mu_F} = \mu_F/|\mu_F|$ is more delicate due to the basis-dependence of μ and is illustrated in Figure 13. While the main idea is still basic Jacobi iteration as in [Lui et al. 2014], here it is necessary to compare the BCs μ_F which are discretized over local vectors \mathbf{X}_1 which are

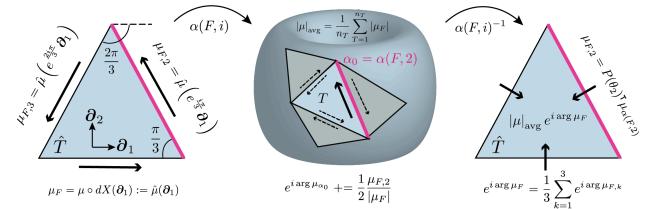


Fig. 13. Step 3 of the QC iteration in the discrete setting applied to an element with equilateral angles. Note the local averaging of the norm and phase as outlined in Algorithm 4.

globally distinct. To describe this more precisely, consider a reference element on the parametrization domain U_h with basis vector ∂_1 aligned with its bottom edge. Then, an adjacent edge making an angle θ with the bottom edge carries a natural direction vector given by $e^{i\theta}\partial_1 = \cos\theta\partial_1 + i\sin\theta\partial_1$. Now, unless the image of this adjacent edge lies on the boundary ∂M_h , there is some other element of M_h which shares it, and as the mesh is generally unstructured the preimage of the corresponding edge belonging to this element may not make an angle of θ with respect to its bottom edge vector ∂_1 . Therefore, it is not sufficient to simply average the representations $\mu = \mu(\mathbf{X}_1)$ directly, as there is no guarantee that the tangent spaces of neighboring elements on M_h are properly aligned. On the other hand, it follows from Lemma 3.4 that the BC discretized on any tangent vector is computable from $\mu(\mathbf{X}_1)$, i.e. if $a = e^{i\theta}$ then $\mu(a\mathbf{X}_1) = \bar{a}^2\mu(\mathbf{X}_1)$ (recall that X is conformal). Therefore, a local Jacobi iteration as in [Lui et al. 2014] is still feasible provided the discrete representations of μ are aligned beforehand. Moreover, it is easy to see that reflection $a\mathbf{X}_1 \mapsto -a\mathbf{X}_1$ induces no change in μ , so that the values of μ on shared edges will be component-wise comparable.

In view of this discussion, the present smoothing on the phase of μ involves first localizing $\mu_F/|\mu_F| = e^{i\arg\mu_F}$ to the edges of each mesh element, averaging its values there, and finally re-projecting these averaged values back into the elements (c.f. Algorithm 4). More precisely, consider an element \hat{T} and the positive angle $\theta_i = \angle(\partial_1, \mathbf{E}_i)$ which each edge vector \mathbf{E}_i of \hat{T} makes with edge vector ∂_1 . Then, the projection matrix from \hat{T} to edge i follows directly from Lemma 3.4 and is built as

$$P(\theta_i) = \begin{pmatrix} \cos^2\theta_i - \sin^2\theta_i & -2\cos\theta_i\sin\theta_i \\ 2\cos\theta_i\sin\theta_i & \cos^2\theta_i - \sin^2\theta_i \end{pmatrix},$$

so that $\mu_{F,i} := P(\theta_i)\mu_F$ gives the appropriate value of μ_F on the i^{th} edge of \hat{T} . Doing this over all elements $1 \leq \hat{T} \leq n_T$ gives the local per-edge phases $\mu_{F,i}/|\mu_F|$ (note that $\det P(\theta_i) = 1$) that can be averaged according to how many elements share each edge. In particular, since reversing the orientation of an edge does not change the value of the phase, it suffices to add the contributions of each local edge $\mu_{F,i}/|\mu_F|$ to the global edge phase $e^{i\arg\mu_\alpha}$, $1 \leq \alpha \leq n_E$ and then divide by the number of elements sharing edge α —either 1 or 2, for boundary or internal edges, respectively. This action yields a smoothed phase in each edge, which can be averaged back from the edges to the elements \hat{T} with the matrices $P(\theta_i)^T$. Finally, once the norm and the phase of μ_F have been independently averaged,

the new BC on the each reference element \hat{T} is reconstructed as

$$\mu_F \leftarrow |\mu|_{\text{avg}} \mathcal{S} \left(e^{i \arg \mu_F} \right),$$

where \mathcal{S} denotes the local smoothing operation just discussed. Pseudocode for the entire procedure outlined this Subsection is given in Algorithm 4.

Remark 5.3. Notice that only one step of Jacobi iteration is performed during the post-processing of μ_F . Empirically, more smoothing than this generally produces a QC Iteration terminating in a “worse” BC with less uniform distortion and a higher maximum value. Theoretically, this is related to the fact that the perturbation $\mathcal{E}(\mu + \varepsilon v)$ is not guaranteed to be smaller than $\mathcal{E}(\mu)$ when $\operatorname{Re} \int_M q(f) v dS_g(\mu) < 0$ where $q(f)$ is the Hopf differential (see [Lui et al. 2015, Lemma 4.1]).

Algorithm 4 Post-processing of μ_F .

Require: Triangulation (or quadrangulation) \mathcal{T} of the surface M with cardinality n_T , per-element values μ_F , local-to-global edge mapping $\alpha(F, i)$.

- 1: Initialize $|\mu|_{\text{avg}} = 0$
- 2: **for** element $1 \leq \hat{T} \leq n_T$ **do**
- 3: Update $|\mu|_{\text{avg}} += \frac{|\mu_F|}{n_T}$
- 4: Normalize $\mu_F \leftarrow \frac{\mu_F}{|\mu|_{\text{avg}}}$
- 5: **end for**
- 6: **for** edge $1 \leq \alpha \leq n_E$ **do**
- 7: Initialize $\mu_\alpha = 0$ and $l_\alpha = 0$
- 8: **end for**
- 9: **for** element $1 \leq \hat{T} \leq n_T$ **do**
- 10: **for** edge $1 \leq i \leq n_e$ **do**
- 11: Compute $\theta_i = \angle(\partial_1, E_i)$
- 12: Update $\mu_{\alpha(F,i)} += P(\theta_i) \mu_F$ and $l_{\alpha(F,i)} += 1$
- 13: **end for**
- 14: **end for**
- 15: **for** edge $1 \leq \alpha \leq n_E$ **do**
- 16: Average $\mu_\alpha \leftarrow \frac{\mu_\alpha}{l_\alpha}$
- 17: **end for**
- 18: **for** element $1 \leq \hat{T} \leq n_T$ **do**
- 19: Initialize $\mu_F = 0$
- 20: **for** edge $1 \leq i \leq n_e$ **do**
- 21: Compute $\theta_i = \angle(\partial_1, E_i)$
- 22: Update $\mu_F += \frac{1}{n_e} P(\theta_i)^\top \mu_{\alpha(F,i)}$
- 23: **end for**
- 24: Compute $\arg \mu_F = \operatorname{atan2}(\mu_F^2, \mu_F^1)$
- 25: **end for**
- 26: **return** $\mu_F \leftarrow |\mu|_{\text{avg}} e^{i \arg \mu_F}$

5.4 Step 4: Ensuring the Distortion is Non-increasing

The final step in the quaternionic QC Iteration presented here is to ensure that the distortion QC_μ does not increase during the perturbative Step 3. Recall from Remark 2.5 that the local smoothing on the phase of μ_F is really a surrogate for similar smoothing on

the argument, and these procedures are inequivalent for multiply connected M . Therefore, it is not guaranteed that Step 3 actually decreases the energy \mathcal{E} in every case, and it is useful to include a line search which forces the iteration to converge when no more progress can be made. This is done by searching for the energetically optimal Beltrami coefficient on the straight line between the iterates μ_k and μ_{k+1} when the previously computed mapping f_k is fixed. In particular, consider the linear combination

$$\mu(t) = t\mu_{k+1} + (1-t)\mu_k,$$

of the BC μ_k from the beginning of iteration stage k and the smoothed BC μ_{k+1} from Step 3. Then, using $\dot{\mu} = \mu_{k+1} - \mu_k$ to denote differentiation with respect to the parameter t , the derivative of $QC(f_k)$ is given by

$$\frac{d}{dt} QC_\mu(f_k) = - \int_M \langle \dot{\mu} df_k^+, df_k^- - \mu df_k^+ \rangle dS_g.$$

Setting the above equal to zero and solving for t then yields the parameter value

$$t_0 = \frac{\int_M \langle \dot{\mu} df_k^+, df_k^- - \mu_k df_k^+ \rangle dS_g}{\int_M |\dot{\mu} df_k^+|^2 dS_g},$$

which corresponds to the optimal BC along this linear interpolation. Setting $\mu_{k+1} \leftarrow \mu(t_0)$ then gives an updated BC which is used as input to the next iteration.

Remark 5.4. It is worth mentioning again that Step 4 is not present in the original formulation of [Lui et al. 2014], and is not strictly necessary for the QC iteration to be applied. On the other hand, faster convergence and better overall performance is observed when it is included.

6 EXAMPLES AND APPLICATIONS

Now that the QC Iteration has been extended to the quaternionic setting of maps from immersed surfaces $M \subset \mathbb{R}^3$, some useful applications are discussed. While there are many varied uses for computational quasiconformal maps, the two discussed at present relate to object deformation and surface remeshing.

6.1 Object deformation

Certain tasks in animation and graphics require a natural looking map between objects which satisfies some prescribed boundary data. It will now be shown that Algorithm 1 and the quaternionic QC Iteration Algorithm 2 can be useful for this purpose by producing evenly distorted mappings which nicely interpolate a given boundary. Some qualitative examples of this have already been seen in Figures 8, 9, and 12, but it is enlightening to discuss quantitative aspects of Algorithms 1 and 2 as well. Figure 14 shows how a genus zero surface with 6 boundary components is deformed by an optimal Teichmüller map computed with QC Iteration to satisfy some given boundary conditions. Here the structure of the mapping is nicely visible: the norm $|\mu|$ is nearly constant and μ is discontinuous at the zeroes and poles of its associated quadratic differential. Similarly, Figure 15 shows two optimal mappings satisfying different prescribed boundary data computed with Algorithm 2 from an open (genus 1) torus. Again, it can be seen that the BC μ is discontinuous and relatively uniform across the target surfaces, as expected.

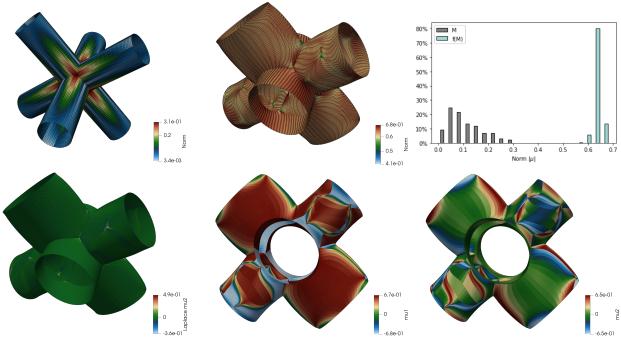


Fig. 14. An optimal quasiconformal mapping of a multiply connected domain computed using the (unconstrained) QC Iteration Algorithm 2 for prescribed boundary data. Top row: source and target surfaces (middle, right) colored by $|\mu|$, histogram. Bottom row: Laplacian $\Delta\mu^1$, components of the BC μ . Notice that μ is discontinuous and harmonic away from its zeros as expected.

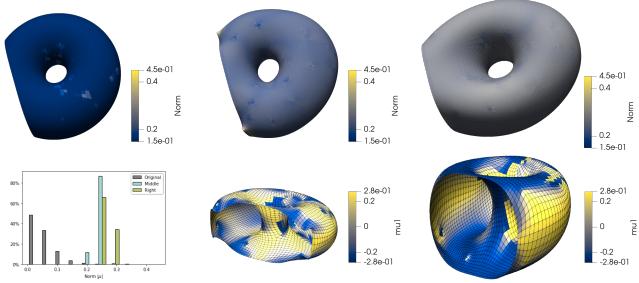


Fig. 15. A quantitative view of the tori from Figure 9 computed with the (unconstrained) QC Iteration Algorithm 2. Top row: source surface (left) and target surfaces computed for different given boundary data, colored by $|\mu|$. Bottom row: histogram, μ^1 -values for each surface above.

Finally, Figure 16 provides a quantitative view of the mapping in Figure 2, showing that the quaternionic QC Iteration is also applicable to planar domains using e.g. $N = (0 \ 0 \ 1)^T$ and produces the expected results. As desired, the optimal Teichmüller mappings computed with Algorithm 2 respect injectivity by construction, leading to a more visually pleasing correspondence.

6.2 Surface remeshing

Most methods for solving PDEs on surface data are only as accurate as the mesh which discretizes the surface M . Here we show that the quaternionic QC Iteration can be useful for generating a nicer discretization of a given surface $M \subset \mathbb{R}^3$ by optimizing the angles between mesh elements. To that end, suppose (M, g) is given where $g = \mathbb{1}_M^* \delta$ is the metric inherited from the embedding of the vertex positions. The idea behind remeshing with Algorithm 1 is to find an optimal quasiconformal mapping $f : (M, g_{\text{ref}}) \rightarrow (M, g)$ in the homotopy class of $\mathbb{1}_M$ which preserves the extrinsic geometry of M but changes its conformal class in a beneficial way, effectively making the mesh elements more regular. This can be done by carrying out the QC Iteration from Section 5 using the constraint formulated

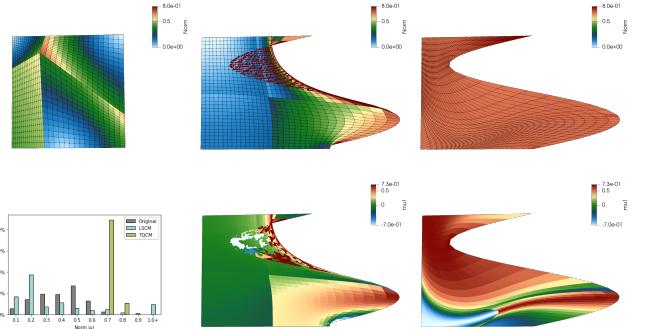


Fig. 16. A quantitative view of the LSCM and TQCM mappings in Figure 2 computed using a target conformal structure of regular squares. Top row: source surface (left), LSCM (middle), and TQCM (right) all colored by $|\mu|$. Bottom row: histogram, plots of μ^1 for the mappings above.

in Section 4.1, which will search for an optimal Teichmüller map from the conformal class of a desired reference metric g_{ref} to the immersed surface which respects the normal N . Some examples of this are displayed in Figures 1, 17, 18, and 19, which show how the QC Iteration produces surfaces with near-uniform distortion despite the use of very irregular initial data.

To elaborate on this procedure, note that the remeshing described above is dependent on the conformal class of the metric g_{ref} , i.e. on the set of reference angles provided by the user. This reflects the fact that QC_μ will measure the μ -quasiconformal distortion of a mapping $f \circ X_h : U_h \rightarrow \mathbb{R}^3$ with respect to a particular discretization of U_h . In particular, the deviation from μ -quasiconformality will be reflected in the way mesh elements are distorted as they move from U_h into \mathbb{R}^3 under this mapping, so it is important to specify a good configuration of angles on U_h to which the target surface $f(M)$ should be quasiconformal. There is an obvious choice for this which comes from pulling back the initial triangulation/quadrangulation of M_h onto the reference domains U_h , so that $g_{\text{ref}} \in [g]$, but this is undesirable for two reasons. First, it is often the case that the angles coming from the initial configuration are nearly degenerate (see e.g. Figure 19 and Figure 1), so this choice is not useful if the goal is to improve the mesh on the surface. Second and most importantly, the identity mapping is always conformal (hence quasiconformal with distortion zero), so in fact this choice will not remesh the surface at all.

Another possibility for the reference discretization is to require the interior angles to be as close as possible to $\pi/3$ in triangular elements and $\pi/2$ for quadrangular elements. This encourages the target to be quasiconformally flat when the connectivity of M is perfectly regular, and provides a reasonable general purpose choice for remeshing applications (see Figure 17). However, this choice can also lead to a sub-optimal mapping for highly irregular meshes, as vertices of the surface M_h may lie on a variable number of elements. In this case, it can be more beneficial to use the “optimized” Algorithm 5 to determine shape of each reference element. This procedure chooses interior angles in such a way that the largest angle of each element dominates the choice of conformal structure. To describe this in more detail, consider that any vertex v_i away

Algorithm 5 Generation of reference angles

Require: Reference discretization \mathcal{T} of the closed surface M .

- 1: **for** $T \in \mathcal{T}$ **do**
- 2: **if** T is triangle **then**
- 3: $N_T = 3, \Theta = \pi$
- 4: **else if** T is quadrilateral **then**
- 5: $N_T = 4, \Theta = 2\pi$
- 6: **end if**
- 7: **for** vertex $1 \leq i \leq N_T$ **do**
- 8: Compute $m_i = \#$ of adjacent elements
- 9: $\alpha_i \leftarrow 2\pi/m_i$
- 10: **end for**
- 11: Determine maximum vertex angle α_i .
- 12: **if** $\alpha_i > \alpha_j$ for all $j \neq i$ **then**
- 13: $\beta_i = \alpha_i$
- 14: **for** vertices $j \neq i$ **do**
- 15: $\beta_j = \alpha_j (\Theta - \alpha_i) / (\sum_{k \neq i} \alpha_k)$
- 16: **end for**
- 17: **else**
- 18: **for** vertices $1 \leq j \leq N_T$ **do**
- 19: $\beta_j = \alpha_j \Theta / (\sum_{k=1}^{N_T} \alpha_k)$.
- 20: **end for**
- 21: **end if**
- 22: **end for**
- 23: **return** β

from the boundary is surrounded by a fixed number of elements m_i , and for each element containing v_i it is desirable for the corresponding interior angle α_i to be as close as possible to $2\pi/m_i$. Moreover, the sum of the interior angles α_i in each element should equal π or 2π for triangular or quadrilateral elements, respectively. Therefore, if a leading node v_i exists in an element such that $\alpha_i > \alpha_j$ for all $j \neq i$, Algorithm 5 fixes α_i and the remaining angle sum $\pi - \alpha_i$ (or $2\pi - \alpha_i$ for quadrilateral elements) is redistributed to the other nodes v_j proportionally to the value of α_j . If a leading node does not exist in a particular element, then the angle sum π (or 2π for quadrilateral elements) is instead redistributed to all nodes v_j proportionally to the value of α_j . This provides a heuristic choice of conformal structure $[g_{\text{ref}}]$ which leads to qualitatively different behavior than the equi-angular choice described before. An illustration of this can be seen in Figure 17, where changing the conformal structure corresponding to g_{ref} yields a noticeable change in the results. Unsurprisingly, a more dilatation-uniform result is obtained by using the conformal structure coming from Algorithm 5 which is adapted to the source surface.

Remark 6.1. It is worth mentioning that the use of a user-defined reference metric g_{ref} makes the result of Algorithm 1 (and hence Algorithm 2) independent of the conformal structure on the source surface $M \subset \mathbb{R}^3$ induced by its embedding. This is highly useful in many cases: as computational meshes often contain elements which are nearly degenerate (e.g. Figures 1 and 19), it is desirable to have algorithms which do not depend strongly on mesh quality.

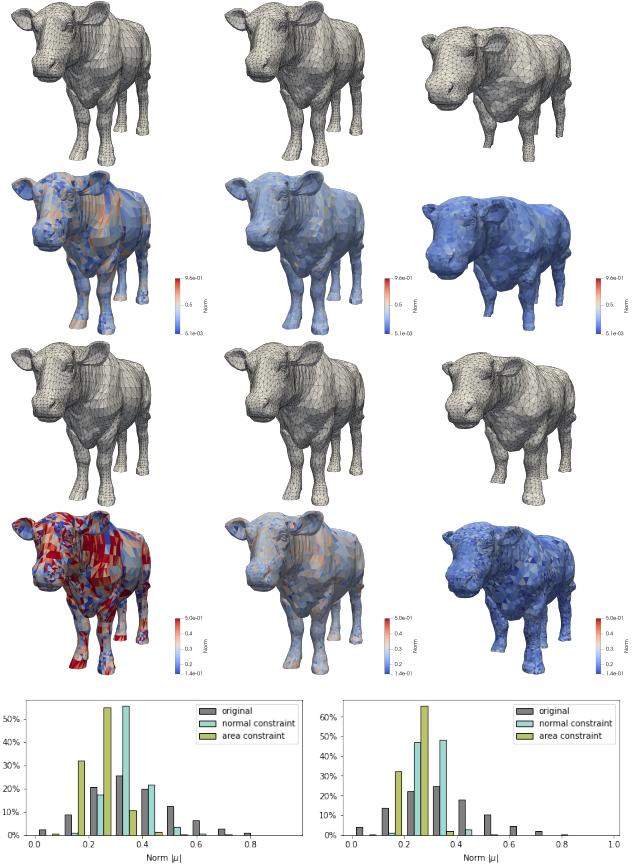


Fig. 17. Top two rows: equi-angular conformal structure; original surface (left), TQCM with normal constraint (middle), TQCM with area constraint (right). Next two rows: same as above but with conformal structure found by Algorithm 5. Bottom row: histogram corresponding to top two rows (left), histogram corresponding to next two rows (right).

This Section is finished with a few more remeshing examples. Figure 18 illustrates how the constraints on area and extrinsic geometry from Section 4 can be used along with the QC Iteration to produce optimal quasiconformal maps which respect sharp features such as corners in the original mesh. Notice that there is a trade-off observed between conformal distortion and geometry preservation; the rounder area-constrained mapping has a norm $|\mu|$ which is significantly left-shifted compared to the sharper normal-constrained mapping. Perhaps this is also related to the fact that minimizing QC_μ is equivalent to minimizing $\mathcal{D}_{g(\mu)}$ when the area is fixed. A similar comparison is carried out in Figure 19, which represents a more challenging example for Algorithm 2. Again, the original mesh is greatly improved (although not quite dilatation-uniform), and the constraints have a noticeable effect on the optimal mapping. Interestingly, here the area-constrained mapping is both more distorted and more uniform than its normal-constrained equivalent. The final example presented here is Figure 20, which shows that optimal quasiconformal mappings may or may not be preferred over

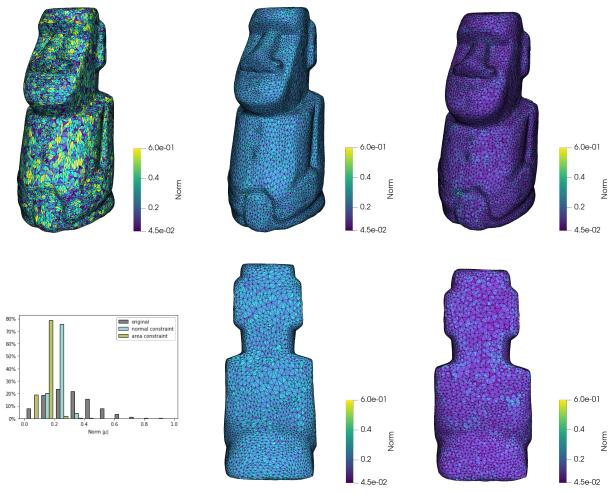


Fig. 18. QC remeshings of a moai statue (left) colored by $|\mu|$, constructed by QC Iteration with different constraints and with respect to the conformal structure given by Algorithm 5. Note that the normal constraint from Section 4.1 (middle) preserves extrinsic features with only minimal rounding of sharp features, while the area constraint from Section 4.2 (right) produces lower conformal distortion at the cost of extrinsic shape.

least-squares conformal procedures when high-quality LSCMs can be produced. In this case, a double torus with constrained area is remeshed using both Algorithm 2 and Algorithm 1 with $\mu = 0$ (i.e. least-squares conformal mapping). Notice that the LSCM produces a very low-distortion mapping at the cost of uniformity, while the result of QC Iteration is highly uniform but more distorted. Therefore, the decision of which technique is most advantageous for any given application is ultimately the choice of the practitioner.

7 CONCLUSIONS AND FUTURE WORK

A direct representation of quasiconformal mappings whose domain is an immersed surface in \mathbb{R}^3 has been introduced, and a least-squares algorithm for computing these objects has been discussed. When the Beltrami coefficient μ is known, this procedure requires only the solution of a linear system and reconstructs the quasiconformal mapping which best represents μ in the homotopy class of the initial data. This has enabled a quaternionic version of the QC Iteration Algorithm from [Lui et al. 2015, 2014], which computes optimal Teichmüller mappings in the case of simply connected surfaces and dilatation-uniform quasiconformal mappings otherwise. The present QC Iteration is applicable to both planar and non-planar manifold meshes and is amenable to optional constraints on surface area and extrinsic geometry. Illustrative examples have demonstrated the utility of these algorithms for object deformation and surface remeshing applications.

Despite the advances discussed here, there remain several challenges and avenues for future work in this area. For one, there is still no satisfactory algorithm for computing Teichmüller extremal mappings between Riemann surfaces of arbitrary genus. Since it has been proven (see [Lui et al. 2015]) that these mappings can

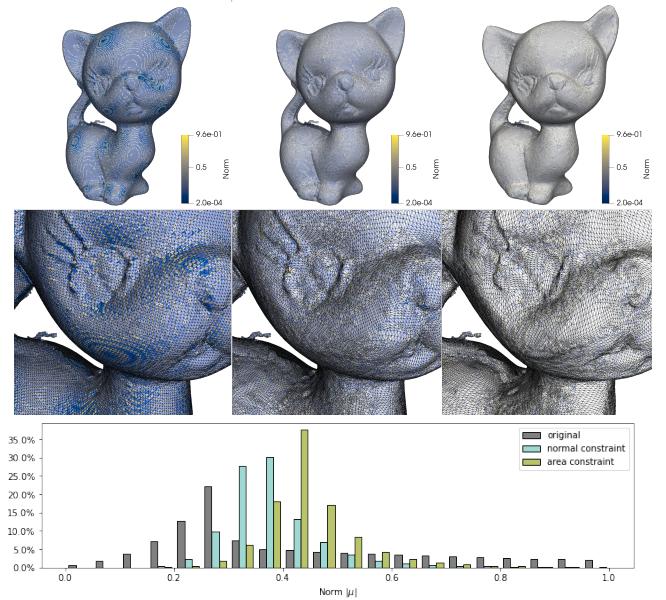


Fig. 19. QC remeshings of a kitten mesh (colored by $|\mu|$) with genus one and 19 boundary components constructed using the QC Iteration algorithm from Section 5 with respect to an equi-angular conformal structure. Left: original surface; Middle: remeshing with normal constraint (c.f. Section 4.1). Right: remeshing with area constraint (c.f. Section 4.2). Bottom: histogram of $|\mu|$.

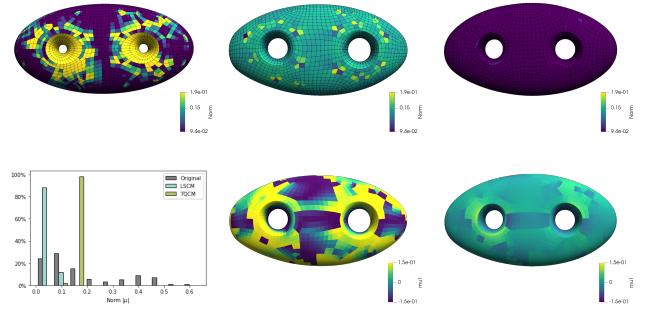


Fig. 20. A comparison of optimal quasiconformal remeshing procedures when area is constrained. Top row: least-squares conformal mapping (top right) and QC Iteration (top middle) on a double torus (top left), colored by $|\mu|$ and with respect to the conformal structure given by Algorithm 5. Bottom row: histogram, plots of μ^1 for the surfaces above. Notice that the LSCM produces less distortion at the cost of uniformity in the mapping.

be computed in many cases by an iterative procedure such as the QC Iteration, it remains important to find a rigorous alternative to the heuristic Step 3 in this procedure which is guaranteed to always decrease the Beltrami energy. Moreover, it would be even better to have a direct, single-stage minimization algorithm which computes the optimal Teichmüller mapping given an initial surface, boundary data, and homotopy class. Finally, it should be mentioned

that every treatment of computational quasiconformal mapping to date relies on solving a discretized problem which agrees with the smooth theory only in the limit of mesh refinement. It would be highly interesting to see progress toward a fully discrete theory for quasiconformality similar to what is currently being developed for conformality.

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A HARMONIC VS TEICHMÜLLER

This Appendix gives a brief description of the connection between Teichmüller and harmonic maps which may be useful for understanding some results in the body. First, let $f : (M, g) \rightarrow \mathbb{R}^3$ be a conformal immersion with normal N . Then, the area form is the trace of the metric tensor $\text{tr } (f^* \delta) = |df|^2$, since

$$\begin{aligned} \frac{N}{2} df \wedge \overline{df} &= -\frac{N}{2} (df * df - * df df) \\ &= -\frac{N}{2} (df N df - N df df) = |df|^2. \end{aligned}$$

On the other hand, the Dirichlet energy with respect to the metric $g = f^* \delta$ also has the simplified expression,

$$\begin{aligned} \mathcal{D}_g(f) &= \frac{1}{2} \int_M df \wedge * df = \frac{1}{2} \int_M |df|^2 + |* df|^2 \\ &= \int_M |df|^2 = \mathcal{A}(f), \end{aligned}$$

so that $\mathcal{D}_g = \mathcal{A}$ when f is conformal, and minimizing the Dirichlet energy is equivalent to minimizing the surface area. More generally, suppose $f : M \rightarrow \mathbb{R}^3$ is a mapping which is not conformal, but there is a conformal structure on M specified by some reference immersion (say z) satisfying $* dz = N dz$. In this case, $df = df^+ + df^-$ relative to this structure and the area form becomes

$$\begin{aligned} \frac{N}{2} df \wedge \overline{df} &= -\frac{N}{2} (df^+ \wedge df^+ + df^- \wedge df^- + df^+ \wedge df^- + df^- \wedge df^+) \\ &= |df^+|^2 - |df^-|^2, \end{aligned}$$

where it was used that

$$df^+ \wedge df^- = -\overline{df^- \wedge df^+} = df^+ * df^- - * df^+ df^- = 0.$$

In this case, the Dirichlet energy is

$$\mathcal{D}_g(f) = \frac{1}{2} \int_M df \wedge * df = \int_M |df^+|^2 + |df^-|^2,$$

which follows because

$$df^- \wedge *df^+ = df^+ \wedge *df^- = 0.$$

This gives a more general decomposition

$$\mathcal{D}_g(f) = \int_M |df^+|^2 - |df^-|^2 + 2 \int_M |df^-|^2 = \mathcal{A}(f) + 2C\mathcal{D}(f),$$

where \mathcal{A} is signed area and $C\mathcal{D}$ is conformal distortion. Note that \mathcal{D}_g depends only on the metric δ of \mathbb{R}^3 and the conformal class of g , and that f is conformal if and only if $C\mathcal{D}(f) = 0$, in which case this expression reduces to the one above. Moreover, when the target of f is fixed, the area term does not vary and so the minimizers of $C\mathcal{D}$ and \mathcal{D}_g are identical. This implies that any conformal map is also harmonic, which has proven to be quite useful for computing minimal surfaces and least-squares conformal mappings (see e.g. [Gruber and Aulisa 2020; Lévy et al. 2002; Pinkall and Polthier 1993]).

It turns out that a similar decomposition is quite useful for understanding Teichmüller quasiconformal mappings. Write $g = \sigma |dz|^2$ for some local conformal coordinate $z : U \subset M \rightarrow \mathbb{C}$ and recall that the quasiconformal distortion of the mapping $f : M \rightarrow \mathbb{R}^3$ with Beltrami coefficient $\mu : TM \rightarrow (TM)^\perp$ (c.f. Definition 3.3) is

$$\begin{aligned} 4QC_\mu(h) &:= 2 \int_M |dh^- - \mu dh^+|^2 dS_g \\ &= \int_U \frac{1}{\sigma} |\mathbf{h}_{\bar{z}} - \mu \mathbf{h}_z|^2 \sigma idz \wedge d\bar{z} = \int_U |\mathbf{h}_{\bar{z}} - \mu \mathbf{h}_z|^2 idz \wedge d\bar{z}, \end{aligned}$$

where we have abused notation by equating the global object QC_μ with its representation on U . Since QC_μ measures deviation from quasiconformality, it is natural to wonder if it appears as the conformal part of some Dirichlet energy functional. To examine this, recall that any mapping $f : (M, g) \rightarrow f(M) \subset (\mathbb{R}^3, \delta)$ which is quasiconformal with respect to μ induces a new metric on M , expressed locally as $g(\mu) := f^* \delta = \rho |d\zeta|^2$, where $\rho = |\mathbf{f}_z|^2$ and $d\zeta = dz + \mu d\bar{z}$. This is simply the image metric expressed on the source surface M , which is not conformal to g unless $\mu \equiv 0$. There is then the usual notion of Dirichlet energy with respect to $g(\mu)$, expressed locally as

$$\begin{aligned} 2\mathcal{D}_{g(\mu)}(f) &= \int_M |df|^2 dS_{g(\mu)} = \int_U \frac{1}{\rho} \left(|\mathbf{f}_z|^2 + |\mathbf{f}_{\bar{z}}|^2 \right) \rho idz \wedge d\bar{z} \\ &= \int_U \left(|\mathbf{f}_z|^2 + |\mathbf{f}_{\bar{z}}|^2 \right) idz \wedge d\bar{z}, \end{aligned}$$

which depends on the metric δ and the conformal class of $g(\mu)$. The next result (essentially due to [Lui et al. 2015, Lemma 3.3]) establishes a decomposition similar to that discussed earlier, which connects the critical points of $\mathcal{D}_{g(\mu)}$ to those of QC_μ in the case that μ has constant norm.

THEOREM A.1. *Let $\mu : TM \rightarrow TM$ be a Beltrami differential with constant norm $|\mu| < 1$ corresponding to a quasiconformal map $f : M \rightarrow f(M) \subset \mathbb{R}^3$. Then, the Dirichlet energy $\mathcal{D}_{g(\mu)}(h)$ of any mapping $h : M \rightarrow \mathbb{R}^3$ decomposes as*

$$\mathcal{D}_{g(\mu)}(h) = \frac{2}{1 - |\mu|^2} QC_\mu(h) + \mathcal{A}(h).$$

In particular, when the area of the image is fixed, quasiconformal maps with $BC\mu$ are also harmonic with respect to the metric $g(\mu)$.

PROOF. Consider local coordinates z, ζ as above so that $g = \sigma |dz|^2$, $g(\mu) = \rho |d\zeta|^2 = |\mathbf{f}_z|^2 |dz + \mu d\bar{z}|^2$. A straightforward computation establishes the Jacobian determinants

$$\begin{aligned} id\zeta \wedge d\bar{\zeta} &= \left(|\zeta_z|^2 - |\zeta_{\bar{z}}|^2 \right) idz \wedge d\bar{z} = \left(1 - |\mu|^2 \right) idz \wedge d\bar{z}, \\ dh \wedge d\bar{h} &= \left(|\mathbf{h}_z|^2 - |\mathbf{h}_{\bar{z}}|^2 \right) id\zeta \wedge d\bar{\zeta}, \end{aligned}$$

and using the representations of $d\zeta, d\bar{\zeta}$ in terms of $dz, d\bar{z}$ leads to the partial derivatives

$$\mathbf{h}_\zeta = \frac{1}{1 - |\mu|^2} (\mathbf{h}_z - \bar{\mu} \mathbf{h}_{\bar{z}}), \quad \mathbf{h}_{\bar{\zeta}} = \frac{1}{1 - |\mu|^2} (\mathbf{h}_{\bar{z}} - \mu \mathbf{h}_z).$$

Therefore, when μ has constant norm, the energy $\mathcal{D}_{g(\mu)}$ may be expressed as

$$\begin{aligned} 2\mathcal{D}_{g(\mu)}(h) &= \int_U 2 \left| \mathbf{h}_{\bar{\zeta}} \right|^2 id\zeta \wedge d\bar{\zeta} + \int_U \left(\left| \mathbf{h}_z \right|^2 - \left| \mathbf{h}_{\bar{z}} \right|^2 \right) id\zeta \wedge d\bar{\zeta} \\ &= \frac{2}{1 - |\mu|^2} \int_U |\mathbf{h}_{\bar{z}} - \mu \mathbf{h}_z|^2 idz \wedge d\bar{z} + \int_U dh \wedge d\bar{h} \\ &= \frac{4}{1 - |\mu|^2} QC_\mu(h) + 2\mathcal{A}(h), \end{aligned}$$

where $\mathcal{A}(h)$ is the signed area of the image surface, counted with multiplicity. Therefore, when this area is fixed $\mathcal{D}_{g(\mu)}$ and QC_μ have the same set of critical points. \square

Since Teichmüller mappings are unique when they exist and have Beltrami coefficients with constant norm, Theorem A.1 implies that any homotopy class of maps from M into a fixed target which contains a Teichmüller extremum also contains a unique harmonic map, and that these maps actually coincide. Conversely, in this case it follows that minimizing QC_μ for a constant-norm μ is equivalent to minimizing the Dirichlet energy $\mathcal{D}_{g(\mu)}$. These facts are what underlie the QC Iteration algorithm, and may explain why the mappings with fixed area computed through this process generally have more optimal distortion profiles (see e.g. Figures 1, 18, and 19).