

2017-03-01 ①

Suppose we have N particles of mass m . Suppose the total kinetic energy is $E = \frac{m}{2} \sum_{i=1}^N v_i^2 = \frac{M}{2} V^2$ where $v_i \in \mathbb{R}^3$.

Suppose the particles are in a small region of space so we don't care where they are. We only care about their velocity distribution. The total velocity state is given by

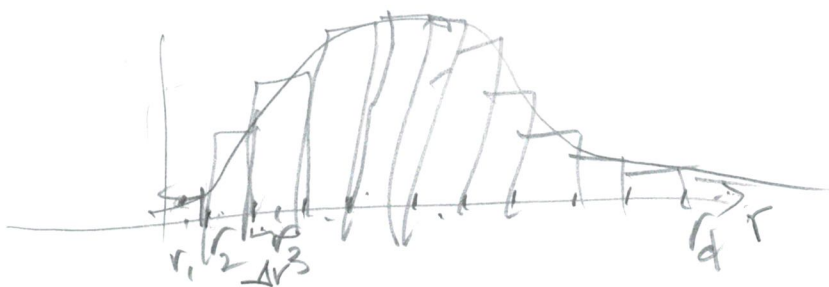
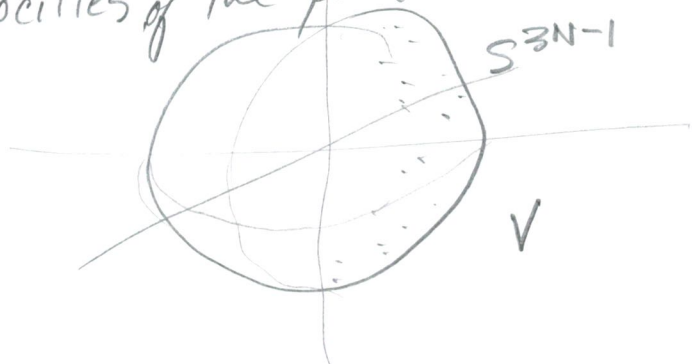
$$V \in \mathbb{R}^{3N}$$

In fact V lies on a sphere of radius $R = \sqrt{\frac{2E}{m}}$

$$\frac{mV^2}{2} = E \quad v^2 = \frac{2E}{m} = R^2$$

$$R^2 = \frac{2E}{m}$$

Consider the distribution of the individual velocities of the particles.



histogram

$$f(v)(r_j) = \# \text{ of particles that have } | \|v_i\| - r_j | \leq \Delta v$$

Try this for a 1-dimensional configuration space.

$$v \in \mathbb{R}^N$$

$$f: \mathbb{R}^N \rightarrow \left(\text{1-d} \rightarrow [0,1] \right)$$

a probability distribution after normalization by N .



S^{N-1}
microscopic
state

$\text{Prob}(\mathbb{R}^+)$
macroscopic
state.

Given the total energy E there is a maximum velocity that a single particle can have. Max. 1=3/1 (2)

$$\frac{m}{2} v_{\max}^2 = E \quad v_{\max} = \sqrt{\frac{2E}{m}} > 0$$

We can divide the single particle velocity space into d cells for some $d > 1$.

$$v_{\max}: \mathbb{R}^+$$

Use spherical coordinates. $\psi: \mathbb{R}^3 \cong S^2 \times \mathbb{R}^+$

$$\psi(x, y, z) = (r, \Omega)$$

$$r = \sqrt{x^2 + y^2 + z^2}$$

$$dx dy dz = r^2 dr \cdot d^2\Omega$$

$$\int_{S^2} d^2\Omega = 4\pi$$

$$\text{define } \Delta v = \frac{v_{\max}}{d}$$



\leftarrow regions in velocity space.

Given a $v \in \mathbb{R}^+$ define $\chi_X(v) = \begin{cases} 1 & \text{if } v \in X \\ 0 & \text{if } v \notin X \end{cases}$ $X \subseteq \mathbb{R}^+$
 $v \in \mathbb{R}^+$
 $v \leq v_{\max}$

The microstates = $\left\{ (v_1, v_2, \dots, v_N) \in (\mathbb{R}^3)^N \mid \sum_{i=1}^N \|v_i\|^2 = E \right\}$

Define the occupancy number $n_j: \text{ind} \rightarrow \mathbb{N}$

$$n_v(j) = \sum_{i=1}^N \chi_{R_j}(\|v_i\|)$$

Express $n_v(j)$ in terms of a density function ρ
assume ρ_v is smooth.

$$\rho_v: \mathbb{R}^+ \rightarrow \mathbb{R}^+$$

$$n_v(j) \cong \rho_v(j \Delta v) \cdot \Delta v$$

$$= \rho_v\left(\frac{j \cdot v_{\max}}{d}\right) \frac{v_{\max}}{d}$$

Improve notation.

3/1

(3)

A microstate is a point $V \in \mathbb{R}^{3N}$ that has total kinetic energy $E > 0$. Let $N \in \mathbb{N}_1$ be the number of particles.

Let $P = \{1, 2, \dots, N\} = 1 \dots N$. P is the set of particle indices.

Regard \mathbb{R}^{3N} as $P \rightarrow \mathbb{R}^3$ = the set of microstates $= A$.

$$\mathbb{R}^{3N} \cong A = P \rightarrow \mathbb{R}^3 \text{ or } \text{Seq}(\mathbb{R}^3) \mid \#V = N.$$
$$N: \mathbb{N}_1 \text{ denote } A(i) = V_i$$

Let $m: \mathbb{R}^+$ be a mass.

Define the single particle kinetic energy

$$\begin{cases} T_i: \mathbb{R}^3 \rightarrow \mathbb{R} \\ T_i(x, y, z) = \frac{m}{2} (x^2 + y^2 + z^2) = \frac{m}{2} \|(x, y, z)\|^2 \end{cases}$$

The N -particle kinetic energy is T

$$T: A \rightarrow \mathbb{R}$$

$$T(V) = \sum_{i \in P} T_i(A(i)) = \sum_{i \in P} T_i(V_i)$$

Let $E: \mathbb{R}^+$ be a given total kinetic energy.

Define the energy sphere in A

$$S: \mathbb{R}^+ \rightarrow \mathbb{R}A \quad S_E = \{V: A \mid T(V) = E\}$$

Observe that $S_E \cong S^{3N-1}$ as a smooth manifold.

$$\dim S_E = 3N - 1.$$

Let $D: \mathbb{N}_1$ be the number of cells that we divide the single particle space into.

$$\text{Since } T(V) = \sum_i T_i(V_i) = E$$

we have $T_i(V_i) \leq E$ for all $i \in P$

Define $B_E \subseteq \mathbb{R}^3$ to be the E -energy ball in the single particle space

$$B_E = \{V: \mathbb{R}^3 \mid T_1(V) \leq E\}$$

$$A = \overline{\Phi} = \Phi^{\otimes N} = \Phi^{\otimes N}$$

$$\mathbb{R}^3 = \overline{\Phi}_1 = \overline{\Phi}$$

Given $V \in A$ let $E_i(V) = T_1(V_i)$

$$\text{Then } T(V) = \sum_i E_i(V)$$

To compute entropy, define the macro state as follows. Let D be some number of cells that we'll divide the single particle ball B_E into.

$$B_E = \bigcup_j B_E(D, j)$$

$$B_E(D, j) = \left\{ V: B_E \mid \left(\frac{j-1}{D} \right) \leq \frac{\|V\|}{v_{\max}} \leq \left(\frac{j}{D} \right) \right\}$$

$j \in 1 \dots D$



where $\frac{m}{2} v_{\max}^2 = E$ so $v_{\max} = \sqrt{\frac{2E}{m}}$

The velocity space volume of B_E is

$$\text{Volume}(B_E) = \frac{4\pi v_{\max}^3}{3} = \frac{4\pi \sqrt{\frac{2E}{m}}^3}{3}$$

$$= \frac{8\pi \sqrt{2} E^{3/2}}{3m^{3/2}}$$

The radius of $B_E(D, j)$ in velocity space is

$$\text{radius}(B_E(D, j)) \approx \frac{j}{D} v_{\max}$$

The volume of $B_E(D, j)$ is approximately

$$4\pi r^2 \underbrace{dr}_{\text{thickness}} = 4\pi \left(\frac{j}{D} v_{\max} \right) \frac{v_{\max}}{D}$$

$$= \frac{v_{\max}^3}{D}$$

The macrostate is the distribution of occupancies of the $B_E(D_{ij})$ for $j=1, \dots, D$ for $V \in S_E$ (5)

$$n_V(j) = \sum_{i=1}^N \chi_{B_E(D_{ij})}(V_i)$$

In order to approximate the situation with smooth functions, we need to take two limits.

- First suppose that instead of having a large number N of particles, suppose we have a smooth matter distribution. This actually loses a lot of information since we lose the identity of the particles. Can we recover that?
- Second, we need to treat space as continuous.

For the occupancy number, if d is very large we can approximate it using a smooth density function.

$$n_V(j) = \rho_V(r(j)) \cdot \Delta r$$

$$r(j) = \frac{j}{D} V_{\max} \quad \Delta r = \frac{V_{\max}}{D}$$

$$\text{So } n_V(j) = \rho_V\left(\frac{j V_{\max}}{D}\right) \frac{V_{\max}}{D}$$

9:30 — 12:00 pm PST } BayTech
 12:30 pm — 3:00 pm EST
 1:00 pm — 2:00 pm EST Assent &