

Optical Properties of Active and Passive Photonic Resonators



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Declaration

I, Ahmad Bilal (CIIT/FA15-BPH-019/ISB) hereby declare that this project neither as a whole nor as a part there of has been copied out from any source. It is further declared that I have developed this thesis and the accompanied report entirely on the basis of my personal efforts made under the sincere guidance of my supervisor. No portion of the work presented in this report has been submitted in support of any other degree of qualification of this or any other University or Institute of learning, if found I shall stand responsible.

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It is certified that Ahmad Bilal (Registration No. CIIT/FA15-BPH-019/ISB) has carried out all the work related to this thesis under my supervision at the Department of Physics, COMSATS University Islamabad and the work fulfills the requirement for award of BS degree.

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Dedication

This thesis is dedicated to my mother who brought me up all by herself, motivated me to always pursue my dreams, and made me the gentleman I am today.

Abstract

Since long, electronic integrated circuits has dominated our modern technology. Now with the dawn of photonics, which is basically using integrated circuits made up using optics, it is not far that our modern technology takes a new toll and slide into a new generation of digital devices. Photonics is the technology of generating and harnessing light and other forms of radiant energy whose quantum unit is a photon.

In this project, we have investigated the optical characteristics of coupled ring resonators, which may be either active (containing a gain element) or passive. In the case of passive coupled resonators, intrinsic and coupling losses affect the circulation, transmission, and dispersion of an input optical pulse. However, for active coupled resonators, gain excitation may be used to compensate for the losses. This alters the spectral and dispersive properties of the coupled-resonator system owing to the enhancement of the quality factor of the system. We demonstrated gain tunable all-optical analogs of Electromagnetically Induced Transparency (EIT) and Electromagnetically Induced Absorption (EIA). This allows us to precisely control the subluminal and superluminal group velocities of light pulses. Furthermore, we demonstrate the gain-mediated reversible transition between the sub and superluminal dispersion regimes.

The optimized coupled resonators enable continuous tuning of sub and superluminal group velocities. Furthermore, owing to gain manipulation, the transmission and light confinement duration may also be precisely tuned. These features make the coupled resonator system attractive for future applications in optical and quantum information science.

*Indeed, in the creation of the heavens,
and the earth and the alternation of
the night and the day, are signs for
those of understanding.*

The Nobel Quran [3:190]

Acknowledgement

In the name of Allah, who is the most beneficent and merciful. I would start off this extensive documentation with a quote from Carl Sagan, one of the greatest science educator ever, who created enough enthusiasm and curiosity in me to pursue my career in Physics. He said, "*Somewhere, something incredible is waiting to be known*". This is one of the reasons I chose to be a student of physics, it inspires me to search for the unknown clues that are hidden in the very fabric of reality. Physics gave mankind the power to dominate their world and use the best of nature for their benefit.

Since childhood, I had always been fascinated by computers and gadgets. Having the background of engineers in my family, I almost ended up joining computer engineering in High School. But the curiosity inside me had made me a stargazer. So I had questions about how do they get where they are, and what are they made of? These questions were those which made me switch my field to Physics which is a science of never-ending curiosity. In this process, a lot of people are included some directly and some indirectly, most of which is my family, because their never-ending support had made me chase my dreams.

So, to start off, I would personally like to thank my supervisor in this BS project, Dr. A. Naweed, whose outstanding supervision and guidance had helped me through thick and thin to complete this project and he also kept me motivated enough to continue my research in the field of photonics. I would like to thank my batch counselor Dr. A. H. Mujtaba, whose support and teachings made us all work harder and harder for the progression of science. Also, thanks to Ms. Zarqa for recommending me to Dr. Naweed. I would like to thank Dr. Seraj for motivating

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In the end, it is important to know that knowledge is a never-ending process, and Physics is such a beautiful field that every time I learn a new concept about the universe and its principles, it feels like I have been born again.

Ahmad B. Yousafzai
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Chapter 1

Introduction

Since the dawn of modern technology, the integrated circuits on which today our every electronic device operates, we have progressed a lot in developing faster and smaller computing devices. Decades have passed since electric circuits became integrated on microchips, also called ICs. This technology has no pause but the field of optical research which generated a great amount of research progress raised to a new form of technology on which we can operate our computing circuits called Photonics. Now is the time that we integrate photonic crystals and photonic structures on circuits and make use of them in communication, signal processing, biochemical sensing, slow and fast light structures [1], optical filters, optical buffers [12-13], wavelength-division-multiplex (WDM)[5-8] and on-chip optical interconnects [8]. Every phenomenon mentioned here is made possible by confining light in a very small volume. microresonators can be used to support the spectrum of optical modes with required polarization frequency and field patterns. These research phenomena will bring revolution to the digital technology, as we know today, with every hand-held device to corporate machines, all running on circuits made using photonic crystals and optical microresonators [9-11].

On a basic level, there are so far two settled components of light control and direction inside the volume of an optical microresonator. The rest is the ordinary system of Total Internal Reflection (TIR) and the presence of evanescent waves, where the directing medium must be optically denser, i.e., have a higher refractive index, than the en-

compassing one so as to accomplish light constrained. The second is the photonic bandgap (PBG) found in artificial optical media having a spatial periodicity in one, two, or three measurements, named photonic crystals (PC), which is a consequence of the phenomenon of Bragg reflection causing the arrangement of frequency bands where propagation of light is restricted by the destructive interference of field harmonics inside the crystal. Exceptionally bound optical modes can be accomplished in these bands when certain deformities are presented in the generally flawless intermittent crystals. With PC defect modes, the light can be found in a size similar to its wavelength (λ/n), where λ is the vacuum wavelength and n is the medium refractive index.[11]

These topics required a detailed study, which is what we are going to do in this Thesis. The scope of this thesis is not limited to the certain and most applicable type of optical resonator which is known as Whispering Gallery Resonators (WGR) [2], but we are going to extend this research on to different possible and quite promising arrangements and geometries of optical resonators known as Micro-Ring Resonators (MRR) [5]. In Microring resonators we mainly focus on the ring-shaped optical wave-guides introducing coupling and different modes in a single and composite system of resonators. This will allow us to collectively measure and observe the combined effects of such resonators by studying their optical properties. Broad numerical and exploratory investigations have been committed towards the investigation of at least two coupled cavities, and a few significant applications have been illustrated, including upgraded spontaneous emission inferable from mode-density enhancement at the photonic band edge [16], enhancement of cavity quantum electrodynamics effects [17], realization of quantum-optical Josephson interferometer [18], parametric oscillations in a triple microcavity system [12], dual wavelength lasing [19], and realization of photodetectors for high-performance wavelength demultiplexing [14]. Coupling effects have been observed in detail and have made possible to observe effects like Electromagnetically Induced Transparency [20] and Electromagnetically Induced Absorption in coupled resonator systems which are called Coupled Resonator Induced Transparency and Coupled Resonator Induced Absorption

[2].

This document is divided into different sections by compiling the work of 1 year long BS final year project. First, we will increase the understanding of the reader of what interferometers, resonators, optical resonators, and microring resonators are. Then, their underlying physics and relating phenomena that are followed by the regimes of these optical systems and what outcome could be achieved by using these optical systems and their applications in photonics. Then we will focus on the systems that we used in this research process and their basic physical explanations. After that, I will show you the results of what I have collected by modeling these systems in different conditions (parameters). This extensive documentation will be useful for anyone trying to get started in this field of research because it is written in such a manner that a newbie in the field of photonics can easily grasp the ideas and can learn from it.

1.1 Resonators

A resonator is a device that exhibits resonant behavior naturally (or artificially) on some resonant frequencies, that is, it oscillates at frequencies with higher amplitudes than others. These frequencies are called resonant frequencies. These oscillations can either be electromagnetic waves or mechanical waves as well. There are different uses of resonators, they can be used to filter some specific frequencies or can also be used to generate a specific frequency of the wave. A resonator in which the waves exist in hollow space is called a cavity resonator, which is used in electronics and radio signal processing, known as microwave cavities, to generate, transmit and receive electromagnetic signals. Acoustic cavity resonators, in which sound is produced by air vibrating in a cavity with one opening, are known as Helmholtz resonators.

1.1.1 Principle

The term resonator is most often used for a homogeneous object in which, vibrations travel as waves, at an approximately constant ve-

locity, bouncing back and forth between the sides of the resonator. The material of the resonator, through which the waves flow, can be viewed as being made of millions of coupled moving parts (such as atoms). Therefore, they can have millions of resonant frequencies, although only a few may be used in practical resonators. The oppositely moving waves interfere with each other, and their resonant frequencies reinforce each other to create a pattern of standing waves in the resonator. If the distance between the sides is d , the length of a round trip is $2d$. To cause resonance, the phase of a sinusoidal wave after a round trip must be equal to the initial phase so the waves self-reinforce. The condition for resonance in a resonator is that the round trip distance, $2d$, is equal to an integer number of wavelengths λ of the wave:

$$2d = N\lambda, \quad N \in \{1, 2, 3, \dots\}$$

If the velocity of a wave is c , the frequency is $f = c/\lambda$ so the resonant frequencies are:

$$f = \frac{Nc}{2d} \quad N \in \{1, 2, 3, \dots\}$$

So the resonant frequencies of resonators, called normal modes, are equally spaced multiples (harmonics) of the lowest frequency called the fundamental frequency. The above analysis assumes the medium inside the resonator is homogeneous, so the waves travel at a constant speed, and that the shape of the resonator is rectilinear. If the resonator is inhomogeneous or has a nonrectilinear shape, like a circular drumhead or a cylindrical microwave cavity, the resonant frequencies may not occur at equally spaced multiples of the fundamental frequency. They are then called overtones instead of harmonics. There may be several such series of resonant frequencies in a single resonator, corresponding to different modes of vibration. [1]

1.2 Optical Resonators

An optical resonator, also known as an optical cavity, is usually composed of two highly reflecting mirror held in front of each other par-

allely inside a vacuum so that the system exhibits resonant behavior which allows standing wave modes to exist with almost no loss. Thus optical resonator is a cavity with walls that are highly reflected for electromagnetic waves (i.e light).

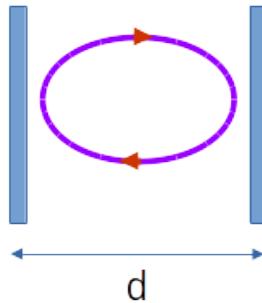


Figure 1.1: Illustration of a basic optical cavity.

1.3 Different Types of Optical Resonators

1.3.1 Fabry-Perot Resonator

A system of two mirrors held parallel to each other and both having high reflectivity's show a resonant behavior at some frequencies of the incident light. If both the mirrors have high reflectance, the incident light is still observed to pass through them without any decrease in the intensity and is detected, which occurs due to phenomenon's similar to quantum tunneling effects [14].

1.3.2 Gires-Tournois

It is basically a lossless Fabry-Perot resonator which has a 100% reflecting rear mirror, that means it reflects 100% at all frequencies. Still, some resonant frequencies stay between the mirrors for a longer period of time and thus descript resonant behavior and lead to ultra slow group velocities. This simple device is known for storing spectral power of light which is reflected from it while modifying its phase. That is why it is sometimes referred to as a "phase only" filter.

1.4 Micro Resonators

Microresonators are a special type of resonators made from a different type of materials which exhibit optical properties while being fabricated on a chip [8]. These kinds of resonators are actually useful in observing the effects of optical resonators on a device.

1.4.1 Different Geometries

There are many types of microresonators from which micro ring-resonators are very useful in making photonic devices and have a wide variety of application. Other kinds of resonators are also useful for different kind of applications and all have distinct optical properties based on their geometry. (See figure 1.2)

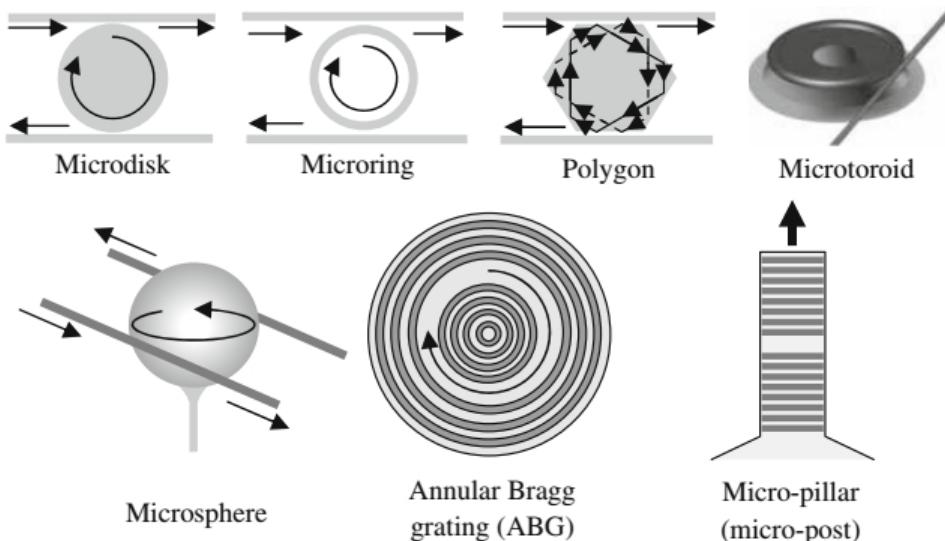


Figure 1.2: Different geometries of microresonators.[15]

1.5 Electromagnetically Induced Transparency and Absorption (EIT and EIA)

Electromagnetically Induced Transparency (EIT) is a coherent optical nonlinearity which makes a medium transparent to some narrow bandwidth of frequencies which were otherwise opaque to the incident radiation. This window leads to slow light at resonant frequencies in

an optical resonant system usually involving coupled system. This is observed due to the destructive quantum interference effects of the incident radiation in atomic levels [4].

Similarly, Electromagnetically Induced Absorption (EIA), is a similar phenomenon to EIT but in this nonlinearity, the medium becomes highly opaque to some bandwidth of frequencies at resonance. Thus blocking off completely the resonant frequency radiation and, causing a dip in the transmitted field. The quantum interference of light here is destructive and the atomic levels absorb the extra photons at such particular frequencies.[7]

1.6 Aims and Objectives

This thesis is a detailed study of the optical properties of such photonic resonators which are composed of passive and active material. We deeply study the changing behavior of active and passive resonators in different parameters. Active resonators are those resonators which are made from some gain medium and they also descript EIT and EIA like behavior in a similar and distinct fashion. We hope to achieve gain controlled variation between slow and fast group velocities of light and enhance the transmission of the system. Different scientific tools and utilities, such as Wolfram Mathematica and Python 3.5, are used to model these conditions and produce results.

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Chapter 2

Fundamental Characteristics of Optical Resonators

2.1 The Fabry-Perot Interferometer

Optical resonators were utilized as helpful gadgets as early as 1899 when Fabry and Perot depicted the utilization of a parallel-plate resonator as a multipass interferometer. Part of the incident light on this Fabry– Perot resonator is transmitted and another part is reflected, with power divisions that rely upon numerous factors. A simple illustration of the basic Fabry-Perot is shown in Figure 2.1, here r_1t_1 are the reflectivity constant and transmitivity constant of the mirror 1 respectively and r_2t_2 are the reflectivity and transmitivity constants of the mirror two respectively. Also, E_i is the incident Electromagnetic energy, E_t is the transmitted energy and E_r is the reflected energy. This is an asymmetric Fabry-Perot resonator:

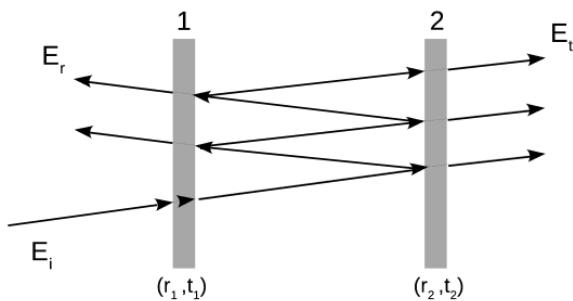


Figure 2.1: Illustrated energy diagram of a simple Fabry-Perot resonator

2.1.1 Theory of Fabry-Perot Interferometer

If the incident energy is in the form of white coherent light then at that point the transmission and reflection coefficients depend just on the mirror reflectivities. The total reflected power comprises of the power reflected from the principal mirror in addition to all the different reflections between the mirrors that add to the reflectivity in general. In summation, the equations are[1]:

$$\mathcal{R} = R_1 + T_1^2 R_2 \sum_{m=1}^{\infty} (R_1 R_2)^{m-1} = \frac{R_1 - 2R_1 R_2 + R_2}{1 - R_1 R_2} \xrightarrow[R_1=R_2 \equiv R]{\quad} \frac{2R}{1 + R} \quad (2.1)$$

Similarly, the transmitted energy in summation is:

$$\mathcal{T} = T_1 T_2 \sum_{m=1}^{\infty} (R_1 R_2)^{m-1} = \frac{T_1 T_2}{1 - R_1 R_2} \xrightarrow[R_1=R_2 \equiv R]{\quad} \frac{T^2}{1 - R^2} = \frac{1 - R}{1 + R} \quad (2.2)$$

Assuming, be that as it may, the incident light comprises of a transiently lucid (monochromatic) plane wave, at that point the reflected power will be relative to the square of the reasonable total of every reflected field. Since the fields convey phase information with amplitudes added, the division of reflected and transmitted light depends not just on the mirror reflectivities but in addition on the mirror separation and excitation wavelength. The rational total of fields is amplified when every one of the fields interferes constructively (in phase) and limited when they interfere destructively (out of phase). Phase gathers with propagation separation as $\phi(z) = \beta z$ and may likewise be gained upon communication with the mirrors. The sound forms of Eqs. 2.1 and 2.2 incorporate an aggregated stage factor for each round-trip that can be translated as a standardized detuning $\phi = T_R \omega$, where T_R is the cavity travel time, $T_R = n_{eff} L / c$ for the circumference, L and effective index n_{eff} . Presently, \tilde{r} speaks to the complex reflectivity:

$$\begin{aligned} \tilde{r} &= r_1 - t_1^2 r_2 \exp(im\phi) \sum_{m=1}^{\infty} (r_1 r_2 \exp(im\phi))^{m-1} \\ &= \frac{r_1 - r_2 \exp(i\phi)}{1 - r_1 r_2 \exp(i\phi)} \xrightarrow[r_1=r_2 \equiv r]{\quad} \frac{r(1 - \exp(+i\phi))}{1 - r^2 \exp(+i\phi)} \end{aligned} \quad (2.3)$$

and \tilde{t} represents the complex transmittivity:

$$\begin{aligned}\tilde{t} &= -t_1 t_2 \exp(im\phi/2) \sum_{m=1}^{\infty} (r_1 r_2 \exp(im\phi))^{m-1} \\ &= \frac{-t_1 t_2 \exp(im\phi/2)}{1 - r_1 r_2} \underset{r_1 = r_2 \equiv r}{=} \frac{-(1 - r^2) \exp(im\phi/2)}{1 - r^2} \quad (2.4)\end{aligned}$$

The square modulus of these perplexing amounts gives the reflection \mathcal{R} and transmission \mathcal{T} coefficients (shown in Fig. 2.2). Antiresonant wavelengths are more emphatically reflected than in the ambiguous case, while resonant wavelengths are transmitted 100% for adjusted reflectors ($r_1 = r_2$). For a fixed reflect dispersing, the transmission and reflection spectra in this manner show intermittent pinnales and valleys. Figure 2.2 presenting the transmission and reflection spectra for a lossless, adjusted Fabry–Perot resonator. The part of reflected and transmitted power for mixed up excitation is identical to the separate spectrally averaged reflection and transmission over time of the spectrum range.

The values of reflectivity coefficients r_1 , r_2 and the transmitivity coefficients t_1 , t_2 are mentioned in the figure. The plot is of the intensity of the Fabry- Perot resonator versus the round trip phase of the system. This displays a 100% transmission and 0% reflection on the resonant frequencies. Meaning all the incident light is detected on the other side of the resonator of these specific frequencies.

2.1.2 Effective Phase

Now let us look at the phase details of the transmission and the reflection spectra of the asymmetric Fabry-Perot resonator. The phase gives us a lot of details about the traveling light inside the resonator and gives other details about dispersion, group delay, and group index. Fig. 2.3 shows phases of both transmission and reflection of an asymmetric Fabry-Perot resonator.

The effective phase also called transmission phase of the system, which is the phase acquired upon transmission, gives us the information about the dispersion of the system. It is basically the argument

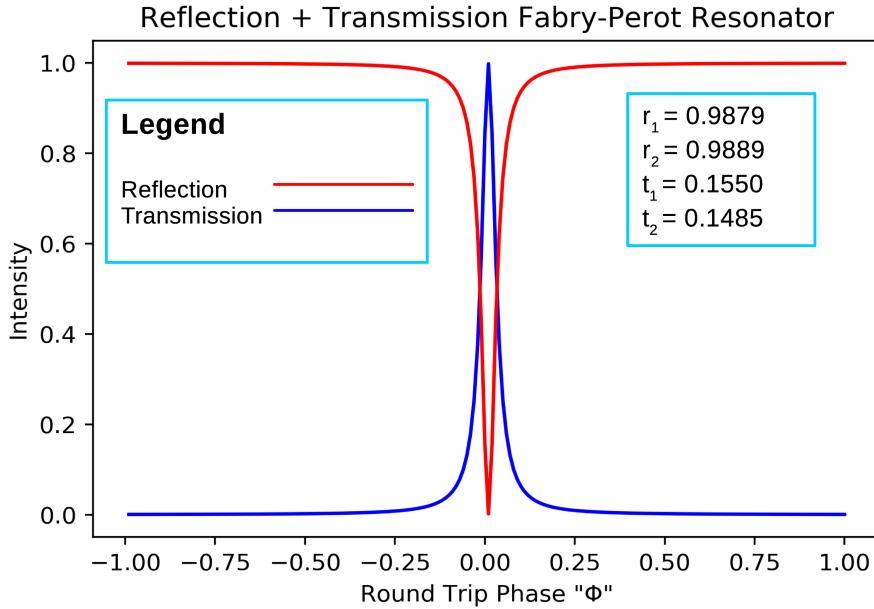


Figure 2.2: Transmitted and reflected field of an asymmetric Fabry-Perot resonator

of the complex transmittivity of the resonator (Eq. 2.4). Any complex number, $z = x + iy$, can be written as $z = |z|exp(i\phi)$. Where, ϕ is the argument of the complex number z given by $\phi = \tan^{-1}(y/x)$. Now we can write for our complex transmittivity as,

$$\frac{E_t}{E_i} = \left| \frac{E_t}{E_i} \right| exp(i\phi_{eff})$$

Where,

$$\phi_{eff} = \arctan\left(\frac{Im[E_t/E_i]}{Re[E_t/E_i]}\right) \quad (2.5)$$

2.1.3 Phasor Plots

Phaser plots are another useful way to study the behavior of light inside the optical cavity. The phasor plots are the complex plots between Real and Imaginary parts of the complex reflectivity and transmittivity (equation 2.3 and 2.4 respectively). Figure 2.4 shows the phasor plots of both transmittivity and reflectivity of an asymmetric Fabry-Perot resonator over the detuning period of 0 to 2π radians.

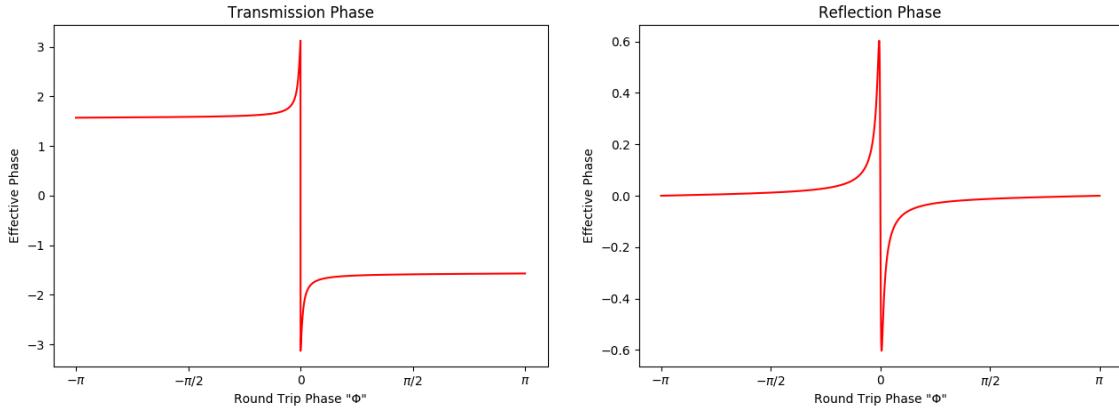


Figure 2.3: Transmission and Reflection phase vs normalized detuning of an asymmetric Fabry-Perot resonator critically coupled.

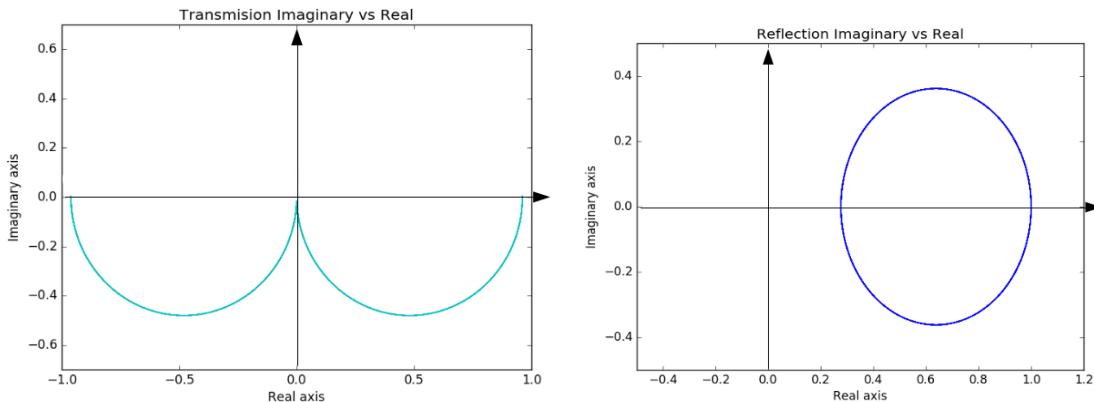


Figure 2.4: Phaser plots of complex Transmittivity and Reflectivity of an asymmetric Fabry-Perot resonator from 0 to 2π

2.1.4 Finesse, Q-factor

The resonance condition is fulfilled when the (compelling) circumference of the ring, or for the most part the round-trip length, is equivalent to a whole number numerous of the optical wavelength inside the medium. This means a progression of Lorentzian-molded transmission bends equally dispersed in recurrence by the FSR (Free Spectral Range), with the resonance linewidth portraying the capacity time of photons inside the cavity. The photon lifetime can be standardized to one optical cycle, known as the quality factor (Q), or the cavity round-trip time, known as the cavity Finesse (\mathcal{F}). The most extreme reachable Q-factor is characterized as Q_{int} , which is the in-

trinsic loss of the cavity such that, $\mathcal{Q}_{int} = 2\pi n/(\alpha_i \lambda_{res})$ [16] where α_i is the intrinsic loss coefficient and λ_{res} is the free-space resonant wavelength. At the point when the resonator is coupled to the outer world, the Q-factor further decreases because of the loss imported by the coupler (\mathcal{Q}_{ext}). This Q-factor, which defines mostly coupling and related external losses, can be given by $\mathcal{Q}_{ext} = (2n\pi L)/\lambda_{res} T$, where L is the circumference of the resonator (in case of a ring), λ_{res} is the resonant frequency, and T is the transmittivity coefficient. Thus the total quality factor \mathcal{Q}_{load} is comprised of these two parts: $\mathcal{Q}_{load}^{-1} = \mathcal{Q}_{int}^{-1} + \mathcal{Q}_{ext}^{-1}$.

$$\mathcal{F}ine se = \frac{2\pi}{\gamma}$$

$$\gamma = \frac{4}{\sqrt{F}}$$

$$F = \left(\frac{2r}{1-r}\right)^2$$

Thus,

$$\mathcal{F}ine se = \frac{\pi\sqrt{F}}{2} \quad (2.6)$$

Similarly,

$$\mathcal{Q}_{factor} = \frac{\lambda_{res}}{FWHM}$$

$$\mathcal{Q}_{factor} = \frac{nLf}{\lambda}$$

$$\mathcal{Q}_{factor} = mf \quad (2.7)$$

2.2 Ring Shaped Resonators

Optical interferometers such as Fabry-Perot or Gires-Tournois resonators are extremely useful in making devices that are compatible

in making spectroscopy tools, add-drop filters for specific optical frequencies, laser cavities as well as dispersion compensators. Due to their free space structural design, they are quite incompatible with planar integrated technology. Thus, we require another geometry of devices which have similar spectral properties. These can be fabricated easily and effectively on microchips and waveguiding geometries by molding of different waveguides into a ring shape.

Let us now discuss how ring resonators, whose principle is pretty much similar to the Fabry-Perot resonator and are more simple to make, operate. Basically, a ring resonator is a simple waveguide which is turned in the shape of a ring (see Fig. 2.5). This allows it to exhibit resonant behavior on very specific frequencies.

2.2.1 Evanescent Coupling

The optical systems discussed above experiences the passage of light across the ring from the optical waveguide through a process known as evanescent coupling. Evanescent coupling is a classical phenomenon whose quantum interpretation is exactly as photon tunneling. It is basically the power transfer of the wave which is directly dependant on the proximity of the bent waveguide and from the straight waveguide. Also, the length or area that has been exposed to the coupler also plays an important role in i.e how many parts of the waveguide/resonator overlaps; illustrated in Fig. 2.5.

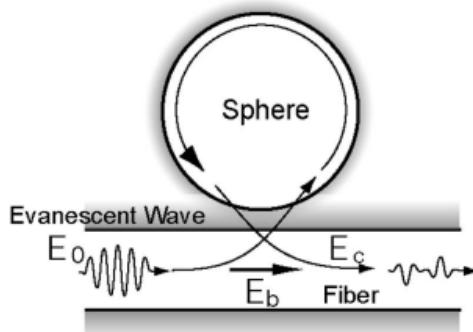


Figure 2.5: Schematic illustration of the microsphere-fiber-taper system.[2]

Coupling strength of an ideal coupler is dependant on the interaction length between the two optical modes which means that the power will transfer more efficiently when the two modes are matched on the basis of their respective phases. This allows us to observe distinct behaviors as well as multiple resonances in transmittance and reflectance.

The light is coupled inside the ring due to two sets of coupling approaches, one being lateral and other being vertical. Either one of these two approaches is necessary for the light from the optical waveguide to couple inside the bent waveguide (ring). The light stays, or resonates, inside the bent waveguide due to the effect of total internal reflection. Lateral coupling requires the same plane fabrication of waveguide and bent waveguide which are easier to fabricate. To ensure strong coupling, small separation is required between the waveguide and bent waveguide. Lateral configuration requires a single layer only but very accurate lithography and etching processes must be done to open up the gaps between straight and ring waveguides with high precision. In vertical coupling transport and bent waveguides are etched in various layers. Vertical coupling expels the index contrast restriction of lateral coupling, additionally, it is an empowering innovation for the ring radii less than $\approx 5\mu m$. Another bit of leeway of the vertical arrangement is that the ring and waveguide layers don't need to be a similar thickness, which improves the design opportunity. These resonators, when fabricated on the chip, display different effects of coupling and transfer of wave energy. These kinds of behaviors have been noticed in all kind of classical waves, such as sound waves, which was first observed inside a large cathedral's halls, thus it was named whispering gallery modes [4]. Also, these resonators can be made using different material but in the scope of this thesis, we used semiconductor silicon as the primary material.

2.2.2 Coupling Regimes

The transmission spectrum of the waveguide-ring coupled system now depends on the proximity and coupling effects. Let us discuss different coupling conditions for these resonators. Fig. 2.11 shows plots for an

all-pass resonator for under, over, and critically coupled system and also for a lossless all-pass system.

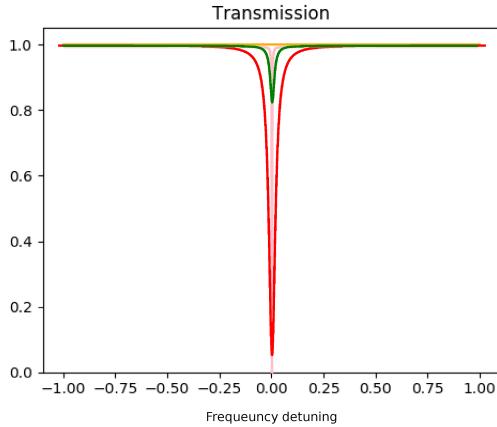


Figure 2.6: Different couplings shown in different colors.

Light when couples inside the ring will lose some part of its intensity. The amount of light that travels inside the ring will acquire a phase shift by completing one round trip and will now interfere destructively with the light which was not coupled inside the resonator. If the phase shifted light has amplitude weaker than the original or uncoupled light, the regime will be known as under-coupled. If the amplitude is higher than the original pulse, the regime would be over-coupled. Similarly, when the amplitude matches exactly with the outside light, the regime would be known as critically coupled. Also, if the light passes through the system without any attenuation or loss, we get 100% transmission.

2.3 Gain Incorporation in Resonators

Light, when travels through a medium, loses its intensity exponentially. This effect can be explained by Beer's Law for electromagnetic intensity. But some mediums, whose refractive index is such as they oppose the exponential decay of the light and rather increase the intensity in the propagation through the medium, are called natural gain medium. Also, there can be an artificial source to activate gain in a certain system. This is done by pumping energy or external light

source i.e. Lasers, to excite the atoms inside the cavity. This makes the stimulated emission releases of the photons increase exponentially and we see an increase in the incident intensity of the input light. We can use these gain mediums and build micro-resonators from them and observe different quantum optical phenomenons. First I will explain a bit about how gain works.

2.3.1 Beer's Law

The simple radiation law follows the beer's law in the absorption of any kind of radiation inside a medium. This tells us that the initial intensity of the light source depends on the variables of the medium it is passing through. For electromagnetic radiation, we can write this law as,

$$I(z) = I_o \exp(-\alpha z) \quad (2.8)$$

Here, I_o is the initial intensity of the radiation, α is the attenuation constant of the medium, z is the amount of distance traveled through the medium and $I(z)$ is the intensity of light after traveling the distance z .

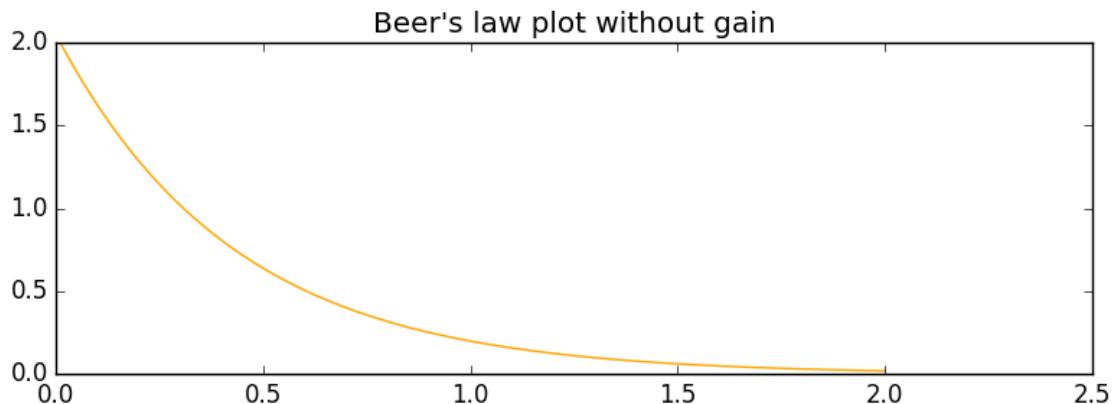


Figure 2.7: Beer's law plot with attenuation $0.01/\text{cm}$: y-axis shows the intensity of light and x-axis shows the distance traveled in meters.

2.3.2 Beer's Law in Presence of Gain

In a gain medium, the intensity of the light will not decrease but it will gradually increase. This means that the attenuation α is negative

or we can introduce a new coefficient for such medium say g such that $-\alpha \rightarrow +g$ where g is some positive real number. This means that the intensity function now grows exponentially rather than decaying exponential.

$$I(z) = I_o \exp(+gz) \quad (2.9)$$

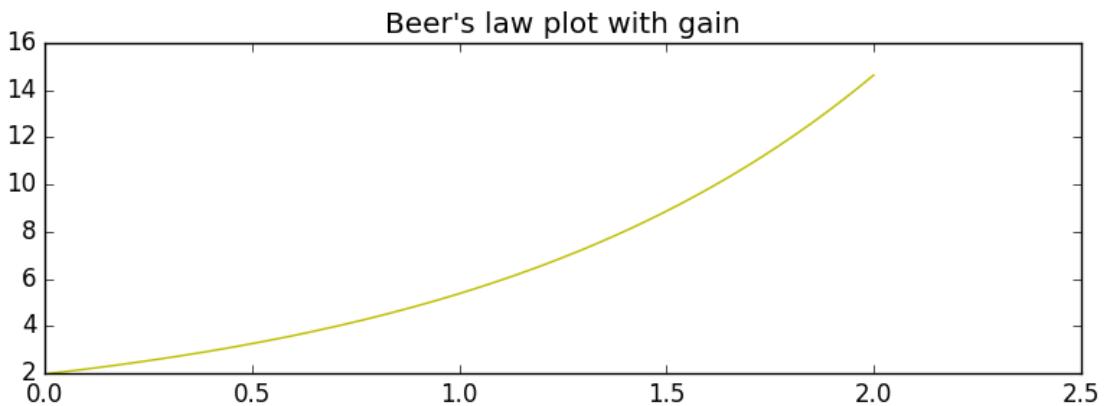


Figure 2.8: Beer's law plot with gain value 0.01/cm: y-axis shows the intensity of light and x-axis shows the distance traveled in meters.

2.3.3 Gain Medium

The active laser medium also called gain medium or lasing medium is the source of optical increase inside a laser. The gain is the result of stimulated emission of electronic or sub-atomic changes to a lower energy state from a higher energy state recently populated by a pump source. This gain in optical systems is usually used for amplification purposes and hence make optical amplifiers. Certain crystals, typically doped with rare-earth ions (e.g. neodymium, ytterbium, or erbium) or transition metal ions (titanium or chromium) can be used as a gain medium. Also, Semiconductors, e.g. gallium arsenide (GaAs), indium gallium arsenide (InGaAs), or gallium nitride (GaN) can also be used when doped [14]. Also, some material such as liquids in the form of dye solutions as used in dye lasers can also be used to make a gain element for active resonators [15].

2.4 Group Delay and Group Advance

The velocity of light is a universal constant $c \approx 299792458m/s$ in free space. However, the light velocity depends upon the medium when it propagates through it. The medium's dispersive properties are defined by the refractive index n of the material. The light pulse propagating from this material experiences delay in the arrival time as compared to the arrival time in a vacuum. To comprehend the results of scattering of the refractive index, we initially think about the proliferation of monochromatic light emission through a material. The phase velocity (v_p) depicts the speed at which the wavefronts travel through the material and is given by: $v_p = c/n$. Thus when the light pulse propagates through the medium, each component of the wave travels at different speed because n is a function of frequency ω . Here, group index is put into account due to the collective effects upon the wave packets of different frequencies, which is given by:

$$n_g = n + \omega \frac{dn}{d\omega}$$

Here, n_g is the group index of the material and ω is the frequency of the wave. Thus now we can define the group velocity of the propagating wave, which is the speed at which the envelope of the incoming pulses moves through the dispersive medium, given by:

$$v_g = \frac{c}{n_g}$$

Thus the reduction of the light experience in its velocity inside a dispersive medium. If the value of $n_g \approx 10^8$, then we experience ultra slow velocities of light which are measured up to $8m/s$ [17-19]. If group index becomes less than 1, such that $n_g < 1$, then we observe group velocities greater than the speed of light in vacuum, such that, $v_g > c$. This does not violate Einstein's Theory of Relativity [20] because there exist various proposals [5–8] for observing such effects by using anomalous dispersion near an absorption or gain lines [7-8]. The group index can also become negative, such that $n_g < 0$, resulting in negative group velocities which are sometimes also referred as 'Backward light' because it appears to propagate backward in the

medium. We can understand the negative group velocity by having a medium of length d , so it will take $d/v_g = n_g d/c$ amount of propagation time for a light pulse to cross it. And the vacuum transition time would be d/c , thus the time delay would be $\Delta T = d/v_g - d/c$ which can be written as, $(n_g - 1)L/c$. So for negative values of n_g , the delay time ΔT is negative which results in group advancement. This concept can be perceived as that if the light pulse is incident on such medium, the light pulse would appear on the other side sooner than if it had traveled through a vacuum. Figure 2.8 illustrates the propagation through a negative index medium [9].

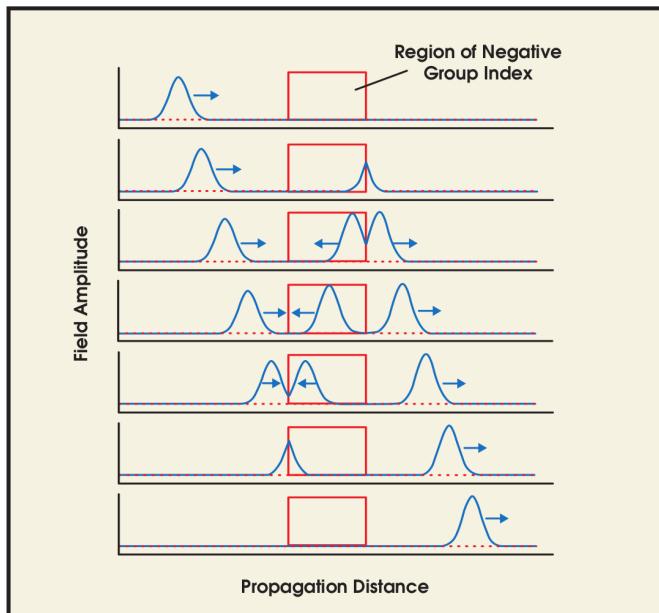


Figure 2.9: Time intervals of a wave propagating through a negative index material. Notice that the peak of entering pulse leaves before it enters the system.

The slow light propagation of light is enhanced by EIT which is a transparency window in the absorption line of the system. Slow light can be used in optical buffers where controllable slow light can dramatically increase the system's performance. Also, the spectral sensitivity of an interferometer[11] can be highly enhanced by introducing a slow-light medium, as well as slow light has a large number of uses in defense applications. Fast light, on the other hand, can be used to make ultra-sensitive gyroscopes [12] and gravitational wave detectors in the fields of astrophysics [13].

2.5 All-Pass Ring Resonator

A straightforward ring resonator is made by taking one yield of a conventional directional coupler and bolstering it once again into one input. Such a device displays periodic cavity resonance (reverberation) when light navigating the ring procures a phased move relating to a number numerous of 2π radians. The resonator is numerically defined from two parts: a coupling quality and an input way. In opposition to the limitless entirety inferences performed before for the Fabry–Perot and Gires–Tournois, in which we expected steady-state task and coordinating fields and derived basic spectral properties. Although both strategies are similarly substantial, the field-coordinating technique has the benefit of simplicity.

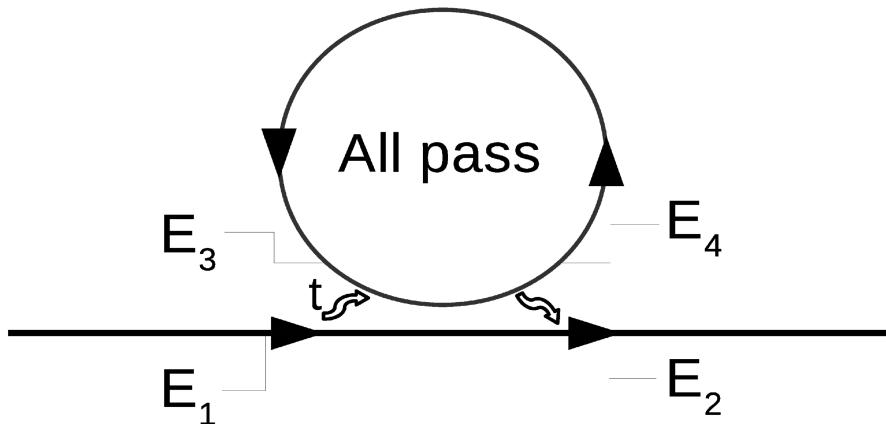


Figure 2.10: Illustrated fields of an all-pass resonator

Fig. 2.9 illustrates the basic geometry of the all-pass ring resonator with its respective energies. The incoming input power is E_1 , which couples inside the ring through the evanescent coupling. Then it travels as energy E_4 and makes a round trip phase. If this phase matches with the input signal, it makes constructed interference with the incoming light with energy E_3 and transfers its energy to the optical waveguide again with energy E_2 . The complex transmittivity and reflectivity are given by equation 2.10, as the reflection is not considered

in this geometry.

$$\frac{E_t}{E_i} = \frac{r - a_1 e^{i\phi_1}}{1 - r a_1 e^{i\phi_1}} \quad (2.10)$$

Here, r is the self-coupling coefficient, $a = e^{-\alpha L/2}$ is the intrinsic loss for circumference L and attenuation α , and $\phi = \omega T$ is the round trip phase of the ring or single pass phase shift.

2.5.1 Transmission For Passive and Active

Let us look at the transmission spectra of the passive and active all-pass ring resonator. Fig. 2.10 shows that the transmission peak is changed into a dip as was in the case of a symmetric Fabry-Perot resonator. On the right, it shows when the gain is introduced inside the system, the dip flips into a peak above the 1 mark on the graph telling us that we have more transmission on resonant frequencies than the actual input.

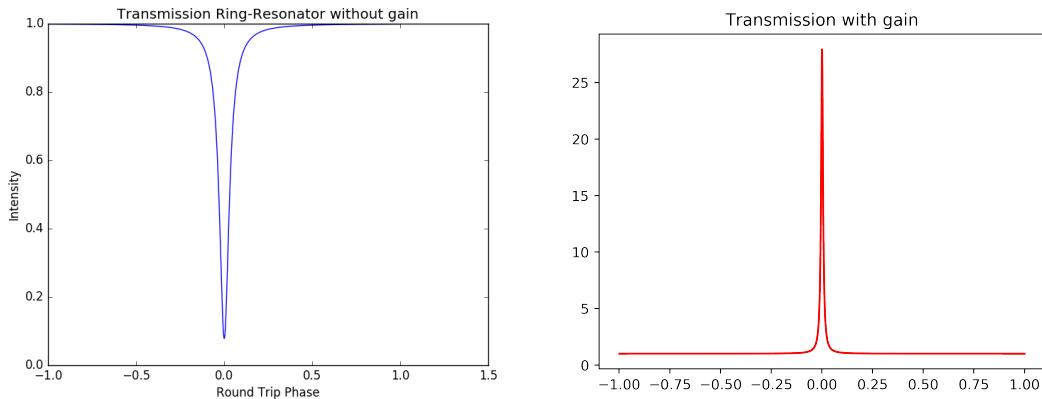


Figure 2.11: Transmission spectra of a passive All-pass ring resonator

2.5.2 Effective Phase

The phase of the All-pass ring resonator is shown in Figure 2.12. We can easily observe from this that with changing the values of the coupling r , the shape of the graph changes as that of a function of $\text{ArcTan}(\phi)$. The relation for phase is given by,

$$\Phi_{eff} = \pi + \phi + 2 \tan^{-1} \frac{r \sin \phi}{1 - r \cos \phi} \quad (2.11)$$

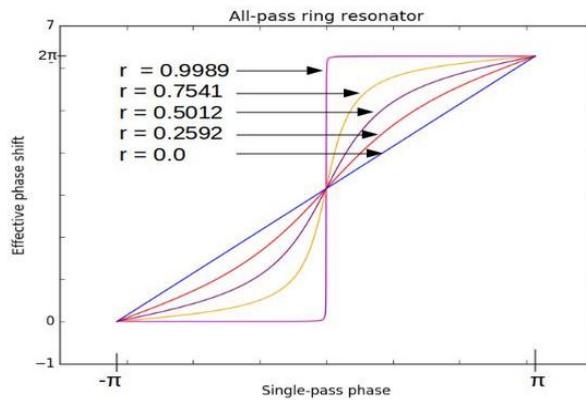


Figure 2.12: Phase diagram of an All-Pass ring resonator from 0 to π where r is the coupling parameter.

2.5.3 Phasor Plots

Now looking into some complex transmittivity of an All-pass ring resonator (Fig. 2.13). This plot is plotted over the complex plain of the detuning limits from 0 to 2π .

We observe that the transmission loop does not go to the negative real axis and touches exactly on the origin. The loop cuts the real axis twice on 0 and 1 and is symmetric above and below the axis.

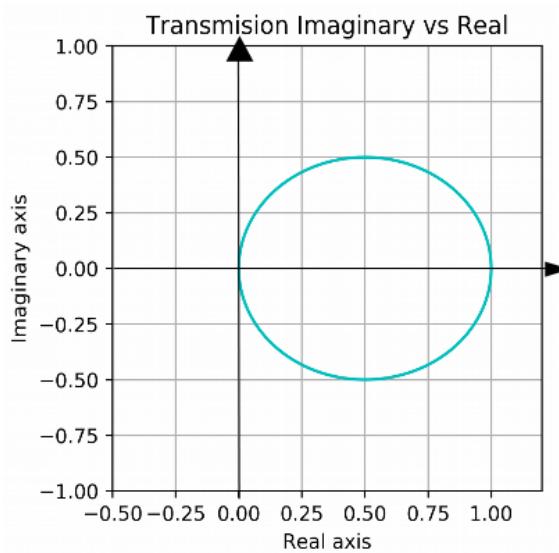


Figure 2.13: Phaser plot of complex transmittivity of an all-pass ring resonator

2.6 Add-Drop Ring Resonator

The immediate waveguide similarity of a free-space Fabry–Perot is acquired by including a second guide that side-couples to the resonator. Since this setup acts as a tight band abundance channel that can include or drop a recurrence band from an approaching sign, it is usually named as an add-drop filter. Fig. 2.14 shows the basic geometry of the add-drop ring resonator with its associated fields labeled accordingly. This resonator has an input, through and drop interfaces where t_1 is add and t_2 is drop coefficients. Input field is labeled as E_1 while the through field is labeled as E_2 . The drop field is on the left top corner labeled as E_5 . The ratio of these fields to the incident/input field defines the total transmittivity and total reflectivity of the filter. The transmission and reflection relations are given in equation 2.12 and 2.13.

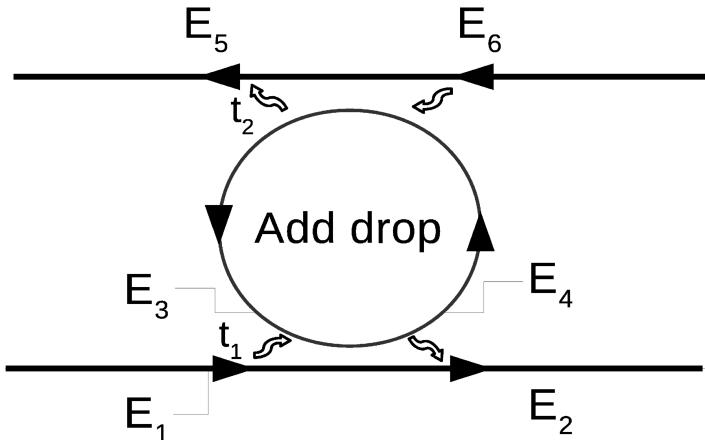


Figure 2.14: Illustrated fields of an add-drop resonator

$$\frac{E_t}{E_i} = \frac{r_1 - r_2 a_1 e^{i\phi_1}}{1 - r_1 r_2 a_1 e^{i\phi_1}} \quad (2.12)$$

$$\frac{E_r}{E_i} = \frac{-t_1 t_2 a_1 e^{i\phi_1/2}}{1 - r_1 r_2 a_1 e^{i\phi_1}} \quad (2.13)$$

2.6.1 Transmission and Reflection

Let us now look at some reflection and transmission spectra of a passive Add- drop filter. Fig. 2.15 shows that the transmission and reflection peaks are flipped as in case of an asymmetric Fabry-Perot resonator and the transmission phase is a direct function of the detuning. Reflection phase is also shown in Fig. 2.16.

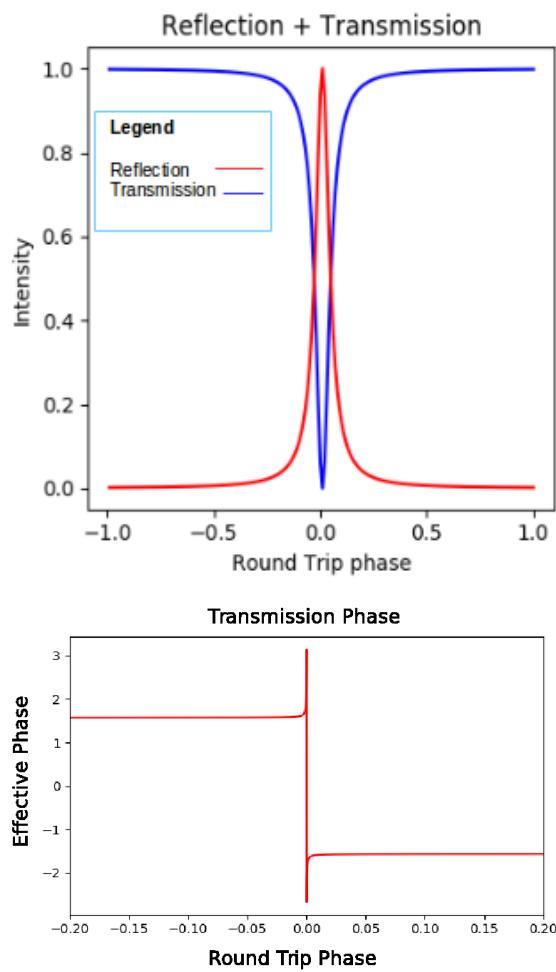


Figure 2.15: Reflection and Transmission spectra along with transmission phase

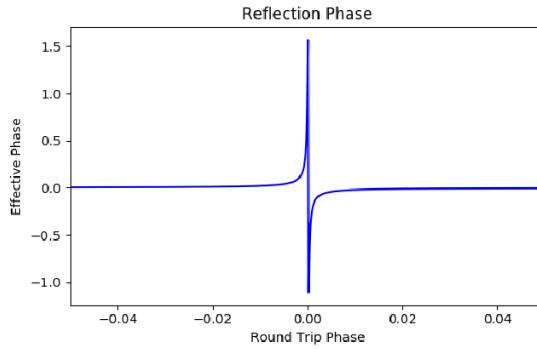


Figure 2.16: Reflection phase of the all-pass ring resonator.

2.6.2 Transmission and Reflection with gain

Now we introduce gain into the system and observe that the transmission dip also shifts into a peak which above the 1 mark (see Fig. 2.17). Meaning that it is greater than the initial intensity and the reflection peak is almost near zero meaning most of the incident light is being transmitted. We will study the transmission of some other different geometries of ring resonators with gain.

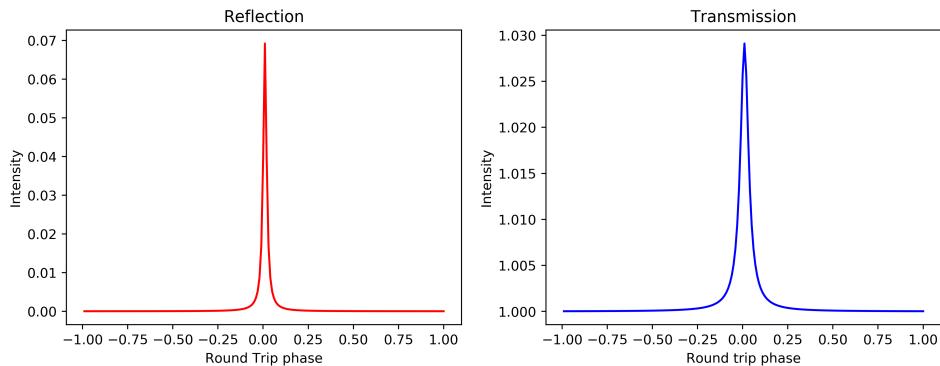


Figure 2.17: Gain introduced into an all-pass resonator: we see clear difference in the intensities.

2.6.3 Phasor Plots

Now let us see how complex plots of Add drop is different from the All-pass resonator. Fig. 2.18 shows that the loop goes towards the negative real axis as the phase is increased. This tells a lot about the distinct behavior.

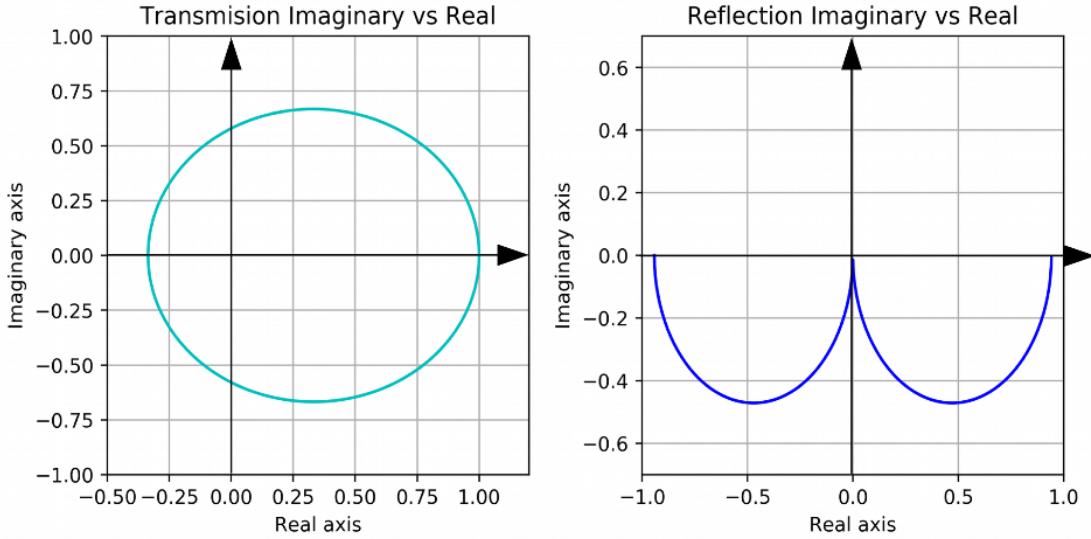


Figure 2.18: Phaser plots of complex Transmittivity and Reflectivity of an All-pass ring resonator from 0 to 2π

2.7 Two Coupled Ring Resonator

Now we turn another optical waveguide into a ring shape and install it on the top of the all-pass ring resonator such that now we have dual ring geometry and a waveguide coupler. This geometry does allow resonant behaviors and the spectra vary largely from an all-pass resonator. In this arrangement, the coupling between the two resonators (rings) also plays an important role in the spectra of the light that passes through the resonator. Fig. 2.19 displays the basic geometry of the couple ring system we are going to discuss along with their energies. The circumference of these resonators are same for both rings given by $b = 25\mu m$, and the refractive index for these resonators is $n = 3.45$ and the Quality Factors are given as $Q_1 = 1 \times 10^5$ and $Q_2 = 1 \times 10^6$

In this geometry, we will study the transmittance of the system as the symmetric case is its reflection in every case. The equation for complex transmittivity of this two resonator system is,

$$\frac{E_t}{E_i} = \frac{r_1 - r_{12}a_1 e^{i\phi_1}}{1 - r_1 r_{12}a_1 e^{i\phi_1}} \quad (2.14)$$

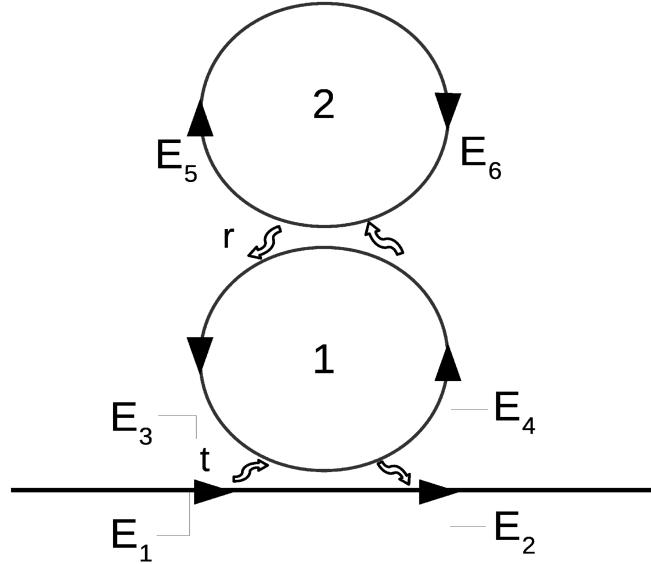


Figure 2.19: Illustrated fields and geometry of a coupled ring resonator

where r_{12} is the coupling parameter of the second resonator and is also a complex number given by,

$$r_{12} = \frac{r_2 - a_2 e^{i\phi_2}}{1 - r_2 a_2 e^{i\phi_2}}$$

2.7.1 Coupled Resonator Induced Transparency and Absorption

Coupled resonator systems, like the one above, have been observed to display the quantum interference effects like Electromagnetically Induced Transparency and Electromagnetically Induced Absorption (EIT and EIA) via their optical analog known as Coupled Resonator Induced Transparency and Absorption (CRIT and CRIA) [3,4]. These quantum interference effects are observed in atoms with at least 3 energy levels but the similar optical analogs of these effects can be seen in these classical resonator systems. This enables enhanced dispersion and transitions from superluminal and subluminal group velocities at room temperature and in micro level scale. These effects will be discussed in detail using ring geometry optical microresonators in later chapters.

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Chapter 3

Coupled Resonator Induced Transparency and Absorption

3.1 Electromagnetically Induced Transparency

Electromagnetically Induced Transparency (EIT) is a well-known phenomenon in atomic physics and its all-optical analog has generated tremendous interest in its physics. Basically, EIT is a transparency window in transmission and absorption spectrum. This narrow transparency window is the result of Fano-like interference among two transition pathways. There is another similar concept which is known as Autler-Townes Splitting (ATS) [1], which also shows a transparency window but it is the result of strong field-driven interactions in a two-level atomic system which causes the level to split [2].

EIT also enables us to hold control over the optical response of the medium. Basically, EIT is the result of having a strong connection between the light and the matter. Amplitudes of different pathways interfere due to quantum interference effects. These can be used in applications such as all-optical switching, slow light [8], optical sensing, light storage, and quantum information processing.

In photonics, EIT is said to be observed in plasmonic structures, photonic crystals, whispering gallery mode microcavities and coupled ring microresonators [5-8]. These devices can be summed up under one name, photonic devices and by seeing such effects we can say that we can get control of how information and energy travel through our device.

3.1.1 EIT in Atoms

A simple classical explanation for the EIT is as follows [2]. In the presence of light (electromagnetic field), the dipole inside the atom will start to oscillate due to the presence of the field. This dipole will oscillate with a certain amplitude and thus will re-emit radiation. If we introduce another field i.e light source which has the same amplitude but have a phase shift of π , then the combined effect of both the fields will cancel out. Thus then we can achieve a situation in which in spite of having an electromagnetic field, the dipole does not oscillate. Thus no radiation will be absorbed and the material becomes transparent. For EIT to occur, we must have a three-level atomic system, such as a lambda configuration, which is discussed below.

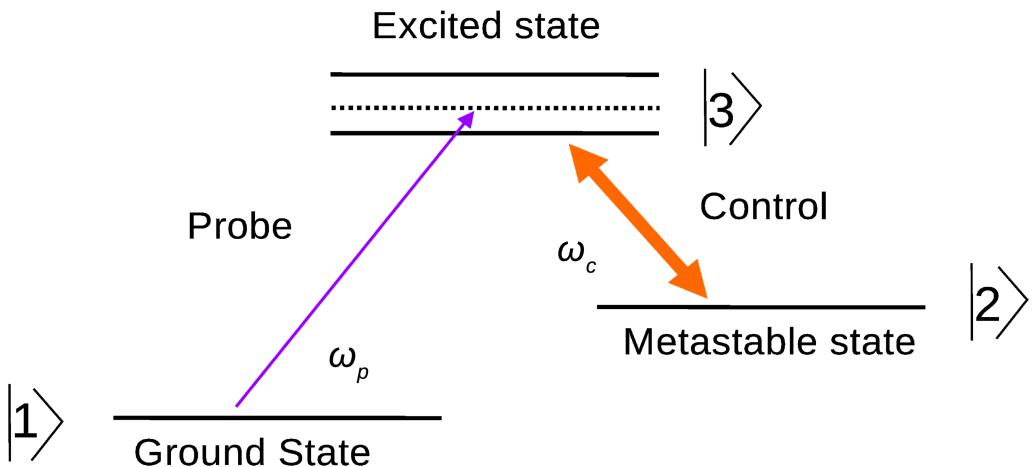


Figure 3.1: A three-level system where level 3 splits due to the much stronger field of control laser.

3.1.2 Three Level Atoms

A lambda-type three-level system is shown in Fig. 3.1. In this system, we will discuss what really happens quantum mechanically for the effect of EIT, without disrupting the essence of classical phenomena. The probability amplitudes of level $|3\rangle$ are driven by two terms in the system. One is the probability amplitude of the ground state $|1\rangle$ and the other is the oppositely phased probability amplitude of the state $|2\rangle$. These both driving forces are opposite in signs but

equal in magnitudes and have a frequency ω_p and are so balanced that probability amplitude of state $|3\rangle$ and the expected value of the amplitude of the sinusoidal motion at every frequency that has been applied is zero.

One may ask how that opposite phase for a transition from the coherent states $|1\rangle \rightarrow |2\rangle$ along with the applied field ω_c , makes absolute cancellation? Because we use the laser pulses that generate fast enough laser photons that the phase of transitions is maintained and is the correct phase for cancellation.

3.2 Coupled Resonator Induced Transparency (CRIT)

We can observe EIT-like properties in various types of coupled resonator systems. However, the scope of this thesis is limited to ring resonators systems only. The optical analog of EIT is referred to as Coupled Resonator Induced Transparency (CRIT). This kind of geometry (that we discussed in section 2.4) has been promising for a long time in the field of photonics. EIT can be observed in this system by mostly the explanation of classical wave travel.

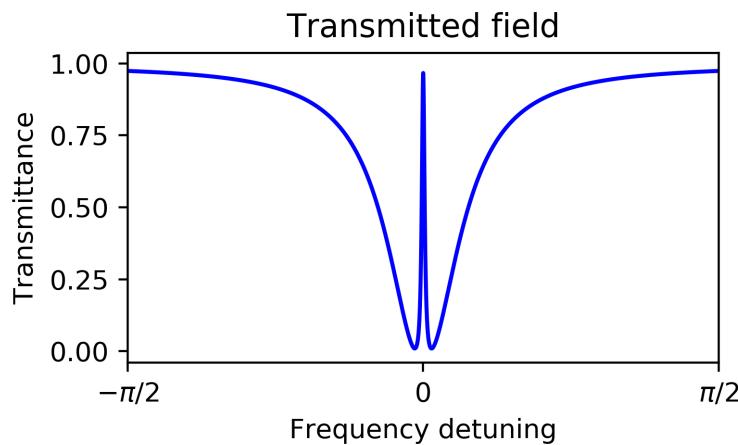


Figure 3.2: Electromagnetically Induced Transparency observed in a 2 ring resonator system.

The traveling photon is coupled inside the first ring through the evanescent wave and travels inside the ring and acquires a phase shift

equal to the round trip inside the optical cavity. When the light source and the phase shifted intracavity field matches so as that the constructive interference is amplified i.e their phases match perfectly, then at those frequencies there is a transparency window in the absorption spectrum i.e a narrow dip, or we see a sharp peak in the transmission spectrum [2].

Figure 3.2 displays the plot of transmitted intensity vs frequency detuning in a coupled resonator system (as shown in Fig. 2.16 in chapter 2). The parameters used here are couplings $r_1 = 0.9$ and $r_2 = 0.999$ and attenuations $a_1 = 0.88$ and $a_2 = 0.9999$ for first and second resonator respectively. Reproduced from the original work on "*Coupled resonator induced transparency*" [2] from 2004.

Now let us look at the phase response of such a coupled resonator system. Figure 3.3 shows the effective phase of the system in red and shows the coupling phase, which is the phase between the two coupled rings, in yellow.

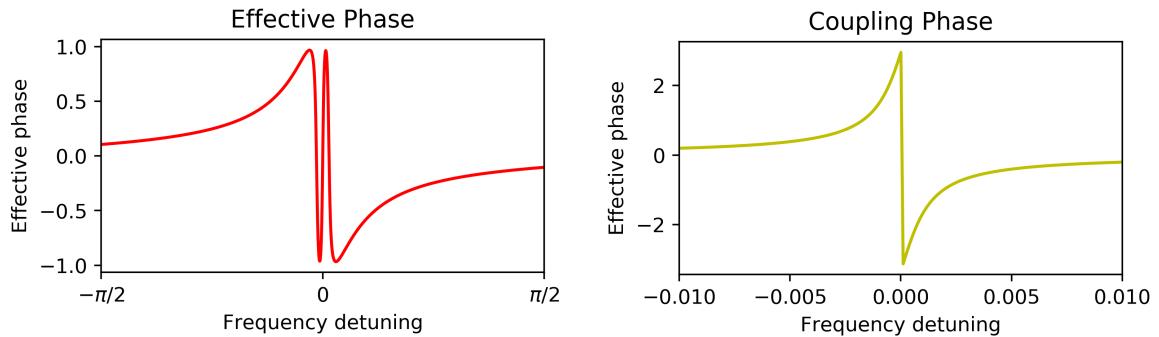


Figure 3.3: Effective phase of the system in red and coupling phase shown in yellow vs frequency detuning.

Figure 3.4 shows the derivative of the phase of the system which gives us great information about the group index and group velocity of the system.

This value is directly related to the group index of the system. From the graph, we can see that there are negative values for off resonances and positive values on resonances. This tells us that we

have superluminal light off resonance and subluminal on resonance.

$$\frac{1}{v_g} = \frac{n}{c} \frac{d\phi_{eff}}{d\omega} \quad (3.1)$$

Where, group index and group velocity are related by,

$$n_g = \frac{c}{v_g}$$

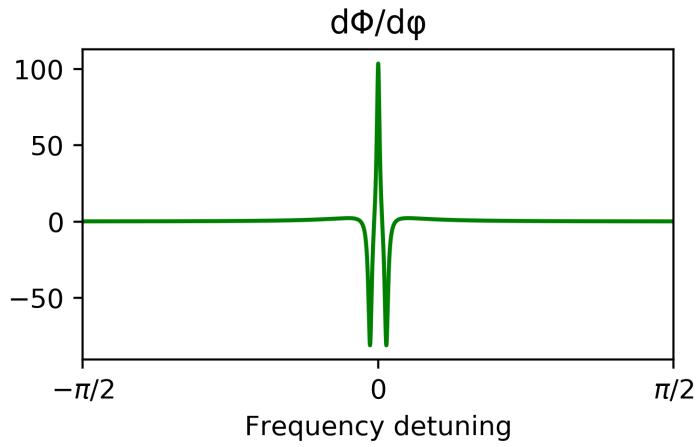


Figure 3.4: Derivative of the transmitted phase of the system vs frequency detuning.

3.3 Electromagnetically Induced Absorption

In contrast to EIT, in Electromagnetically Induced Absorption (EIA), constructive interference occurs between the transition probability amplitudes occurring across two transition pathways. This leads to a narrow dip in the center of the broader transmission dip in the spectrum of the system, known as Electromagnetically Induced Absorption (EIA) shown in Fig. 3.5. As a whole, EIA is not a very well understood phenomenon and its explanation lacks the true essence.

3.4 Coupled Resonator Induced Absorption

Coupled resonator system as discussed above also displays EIA-like effect. The optical analog of EIA is known as Coupled Resonator Induced Absorption. CRIA can give us both fast light and slow

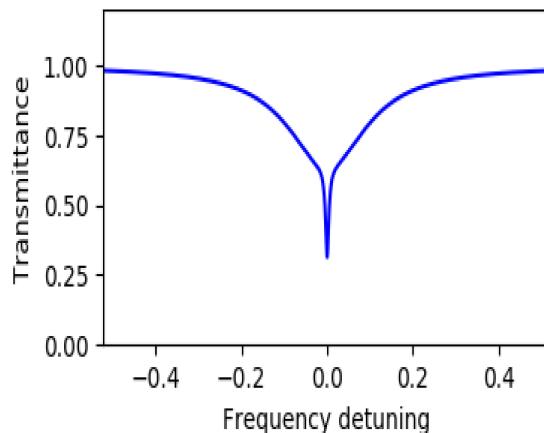


Figure 3.5: Electromagnetically Induced Absorption in a coupled two ring resonator system known as Coupled Resonator Induced Absorption.

light, most of the light leading to a different set of broad applications. The system discussed below is again a two ring resonator system as above and we will again follow the same procedure. Fig. 3.5 displays the CRIA in ring resonator system with coupling parameters as $r_1 = 0.8998$ and $r_2 = 0.9998$

3.4.1 CRIT in a Passive System

As before, now we are going to observe what changes does the system has when we introduce gain in it. For convenience, we will first see the passive results for a system and then introduce gain in either in second resonator or first resonator and then introducing gain in both of the resonators. This can be introduced by pumping some monochromatic light source or a laser, in either one of the rings which will drastically compensate the losses inside the resonator and will increase the overall output transmission of the system even above the incident light source. The system parameters are as follows, $r_1 = 0.9889$ and $r_2 = 0.9998$ with $Q_1 = 1 \times 10^5$ and $Q_2 = 1 \times 10^6$ for first and second resonator respectively (Fig. 3.6).

3.5 CRIT Results

We observe EIT in a coupled two resonator system, the transmission and effective phase of the system is shown in Fig. 3.6 displaying

normal dispersion meaning slow light in the system.

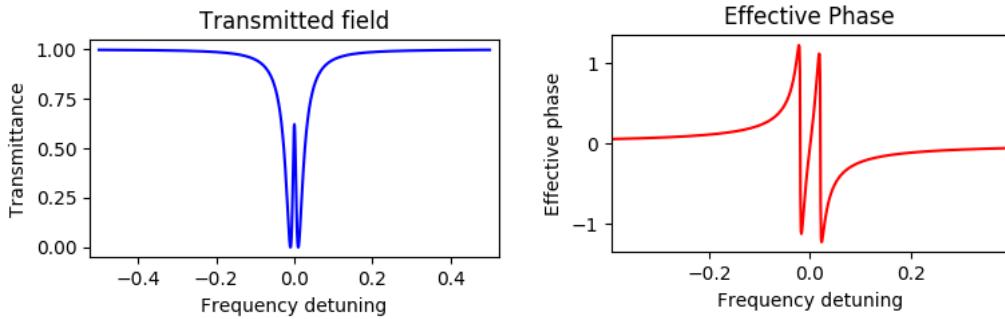


Figure 3.6: CRIT with its effective phase in a passive resonator system.

3.5.1 Introducing Gain Only In Second Resonator

Now we will activate gain in the second resonator, which has higher Q-factor (labeled 2 in Fig. 3.7). We increase the gain to compensate for the intrinsic losses, denoted by α_i , which is directly related to the Q-factor of the resonator. The transmission peak of the EIT starts to rise up gradually as g , the gain coefficient, is increased. The peak rises towards 1 mark up till $g \rightarrow \alpha_i$, where alpha is the attenuation constant for intrinsic loss.

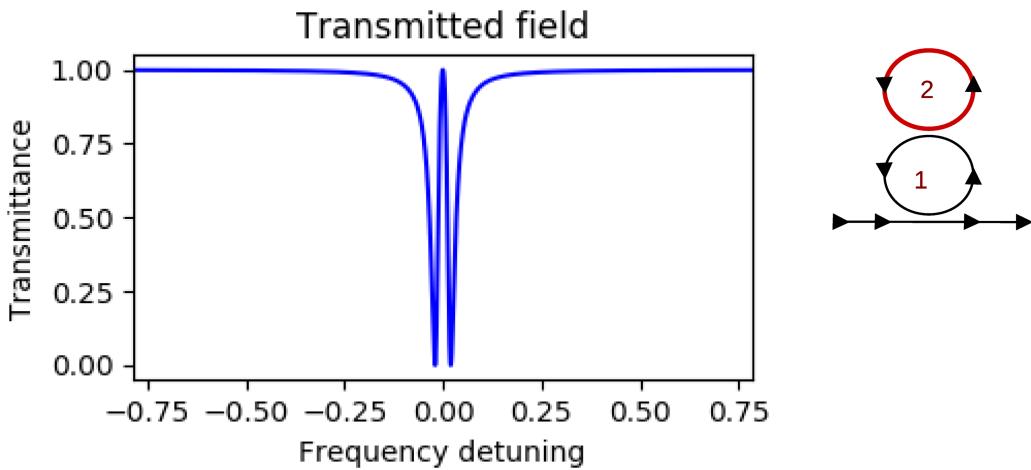


Figure 3.7: CRIT in an active coupled resonator system.

When gain coefficient equals the intrinsic lost coefficient, i.e $g = \alpha_i$, the peak of the transmission touches the unity on the graph meaning

now all of the incident light is verily being transmitted the other side i.e the medium has become completely transparent for resonant frequencies. This means that we have now compensated for all the losses inside the system which are intrinsic. When this happens, we see an abrupt change in the effective phase of the system. This system now gives us anomalous dispersion on-resonance meaning we get fast light in EIT (Fig. 3.8).

To get a brief idea of what is really happening, I have also calculated the group index of the system. In this case, we receive negative values for the group index n_g on resonance. Fig. 3.9 displays the group index for this particular case.

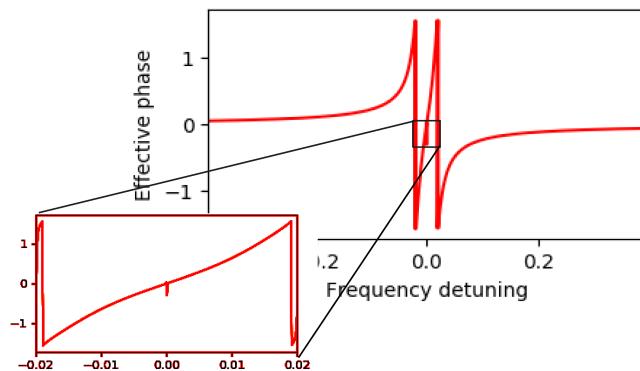


Figure 3.8: Effective phase of CRIT in an active coupled resonator system. Magnified view of phase near resonance frequency is shown.

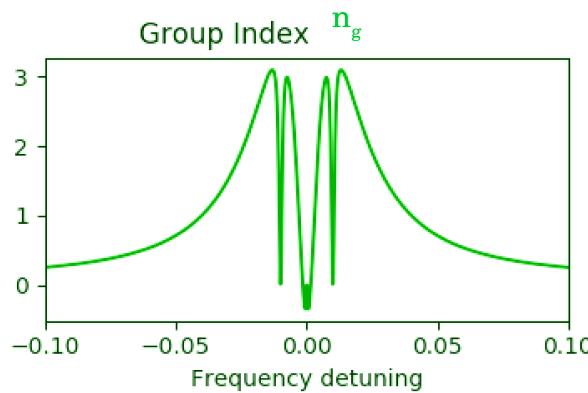


Figure 3.9: Group index of an active resonator system showing negative values on resonant frequencies.

3.5.2 Introducing Gain Only In First Resonator

Now we will activate gain in the first resonator (labeled 1), which has lower Q-factor in comparison to the second resonator. The spectrum of the EIT starts to rise up and displays a hanging EIT. (Fig. 3.10)

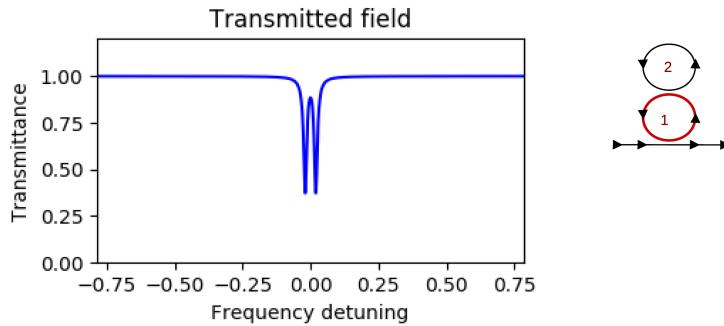


Figure 3.10: CRIT of the two resonator system with gain activated in resonator 1.

The effective phase of the system remains the same as a passive resonator, displaying normal dispersion (slow light). The index value on resonance calculated to be about ≈ 76.2 which shows that light has been slowed down only marginally for this case (Fig. 3.11).

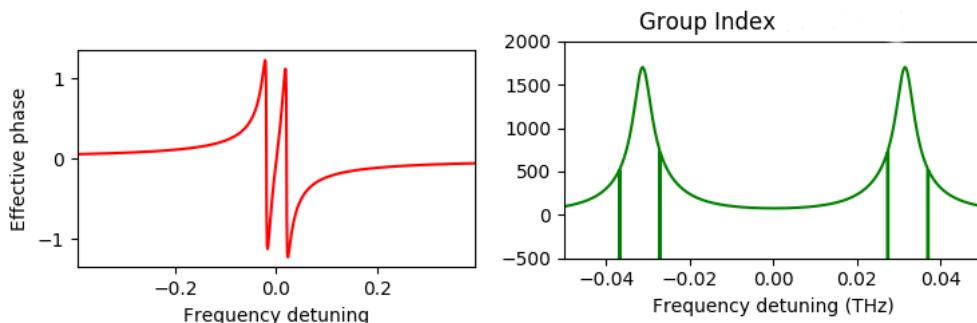


Figure 3.11: Effective phase shows normal dispersion (in red) and group index n_g shown in green.

3.5.3 Introducing Gain In Both Resonators

Now we will activate gain in both of the resonators simultaneously such that the ratio of the gain coefficients, g_1 and g_2 , are the same. We observe that the peak of the EIT, as well as the whole transmission, starts to rise up towards unity as $g \rightarrow \alpha_i$ (Fig. 3.12). The EIT transmission window also narrows down gradually.

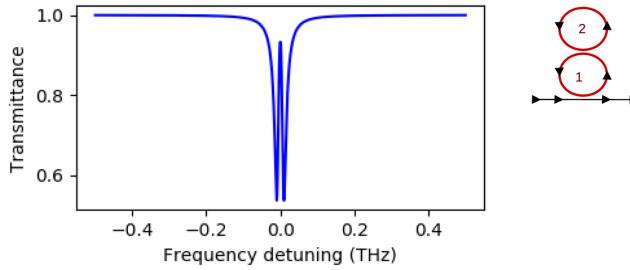


Figure 3.12: Transmission graph of two resonator system with gain activated in both.

The effective phase of the system shows rather distinct curves which are basically due to the artifacts in the system or computation errors. The on-resonance information tells us that we have normal dispersion and positive group index about ≈ 766.5 (Fig. 3.13).

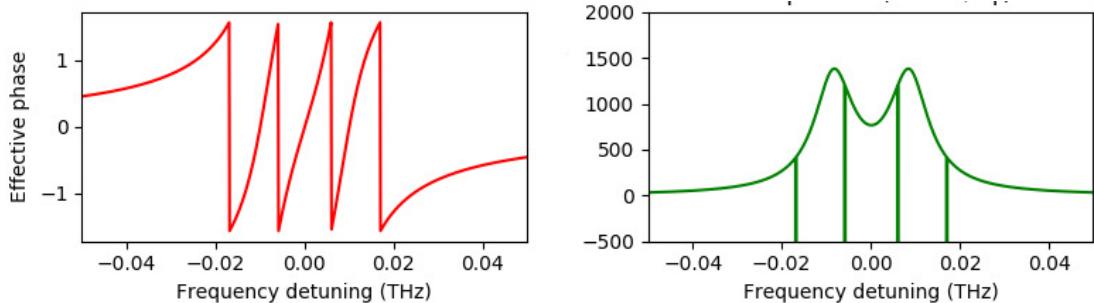


Figure 3.13: Phase and group index of a resonator system with gain in both resonators.

When the gain coefficient g becomes greater than the intrinsic loss coefficient α_i , then the whole transmission graph flips about the horizontal axis and we now see a distinct spectrum with a reversed EIT shown in Fig. 3.14.

This flipping of the transmission does not affect the dispersion of the medium. This means the effective phase and the group index of the system are similar in this case. However, we now have a transmission peak which is above unity. This shows that due to gain, intrinsic losses have been compensated and a transparent coupled resonator system is achieved.

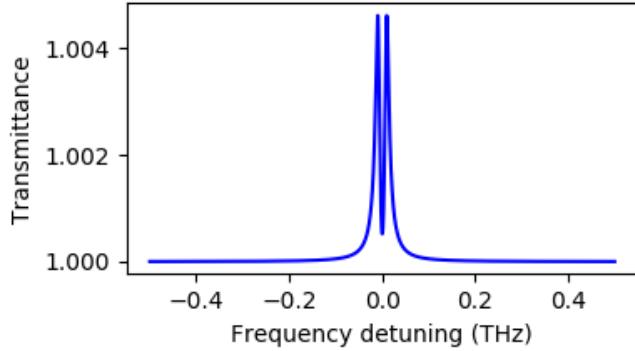


Figure 3.14: Flipping of the EIT spectrum when gain coefficient is larger than the attenuation coefficient.

3.6 CRIA Results

Now we will observe the behavior of the transmission spectrum of CRIA when we introduce gain in it. Similarly, first we are going to activate gain in second resonator, then into the first resonator and in last we activate and increase gain in both resonators simultaneously.

3.6.1 CRIA With Slow Light

We will first examine the response of the CRIA (Fig. 3.5) in a passive medium with having normal dispersion. shown in Fig. 3.15, with its group index.

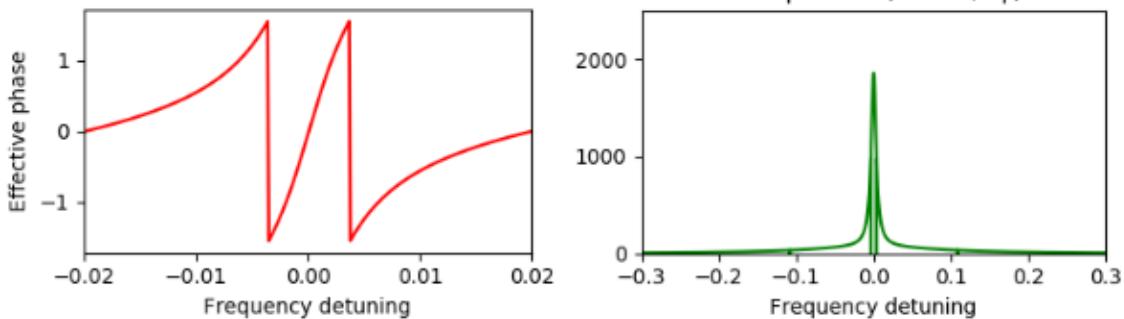


Figure 3.15: Phase and group index of CRIA.

3.6.2 Introducing Gain Only In Second Resonator

Now we activate gain in the second resonator such that $g \leq \alpha_i$. When the gain becomes closer to the value of α_i ($g \rightarrow \alpha_i$), the EIA spectrum

changes into EIT type transmission (Fig. 3.16).

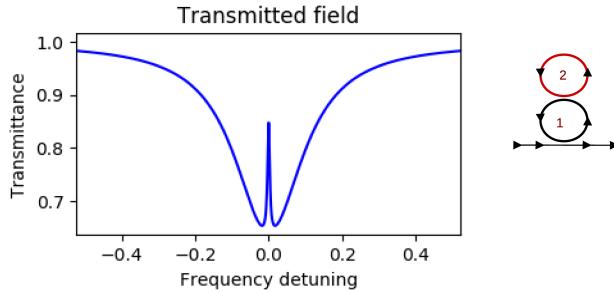


Figure 3.16: CRIA dip changes from Fig. 3.5 into an CRIT type transmission when gain is introduced in second resonator.

The effective transmission phase changes from normal dispersion to anomalous dispersion. Corresponding group index displays a negative value of $n_g \approx -4505$ at resonance, implying a negative group delay and therefore, superluminal light (Fig. 3.17). Thus EIT which, in a passive coupled resonator system, always yields slow light, now enables superluminal (fast light) group velocities of resonant frequencies.

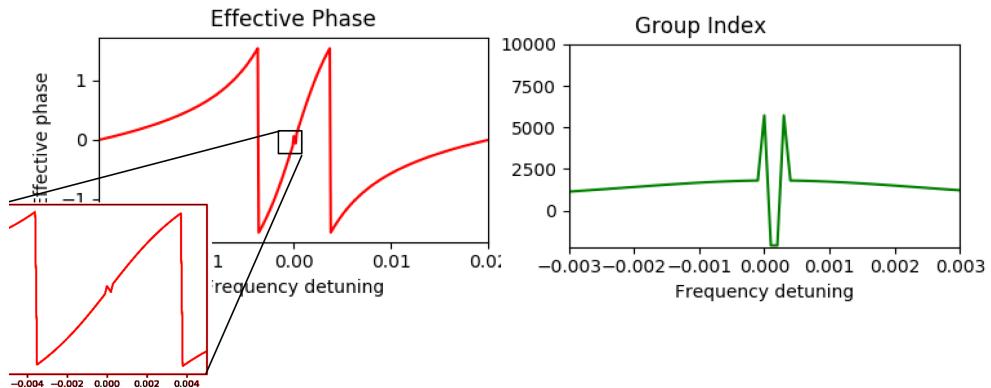


Figure 3.17: Phase of the system showing anomalous dispersion on resonance (zoomed) and group index showing negative values.

3.6.3 Introducing Gain Only In First Resonator

Now we activate gain in the first resonator (labeled 1 in Fig. 3.18) and observes that the EIA resonance narrows down and eventually becomes a sharp transmission dip. Fig. 3.18. Corresponding trans-

mission phase and group index reveal normal dispersion and slow light (Fig. 3.19).

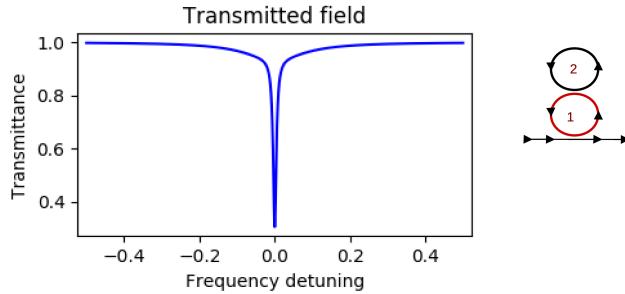


Figure 3.18: CRIA with gain activated in first resonator.

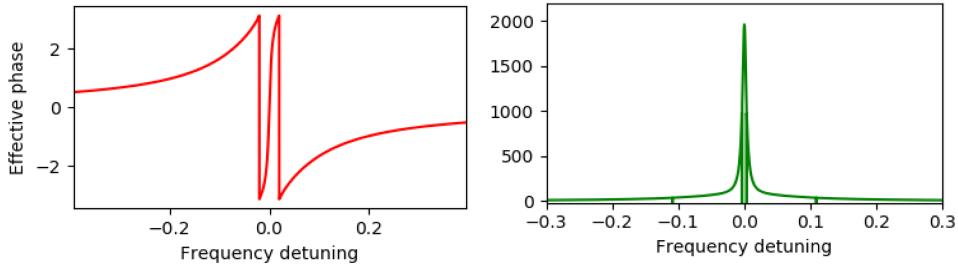


Figure 3.19: CRIA phase and group index.

3.6.4 Introducing Gain In Both Resonators

Now we consider the case where the gain is activated in both of the resonators simultaneously. No notable changes occurs in the transmission spectrum when g_1 and $g_2 < \alpha_i$. Fig. 3.20. The phase and group index is also shown in Fig. 3.21.

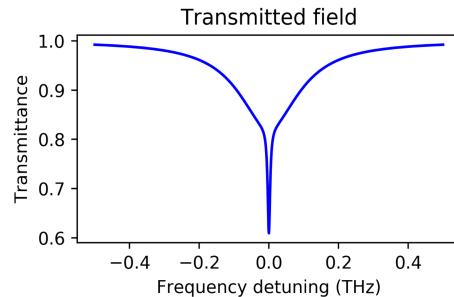


Figure 3.20: CRIA with gain in both resonators.

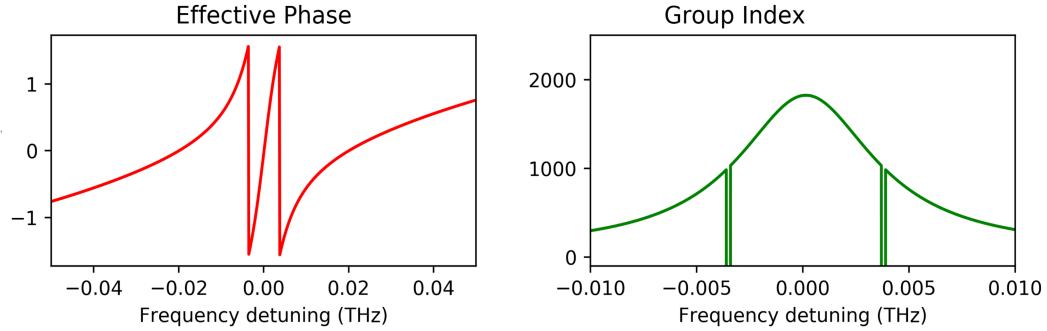


Figure 3.21: CRIA phase in red and group index in green.

However, if gain becomes closer to intrinsic loss, i.e $g \rightarrow \alpha_i$, the transmission spectrum values are very near to 1 now and we see anomalous dispersion in the effective phase of the system and negative group index of about ≈ -4550 (see Fig.3.22).

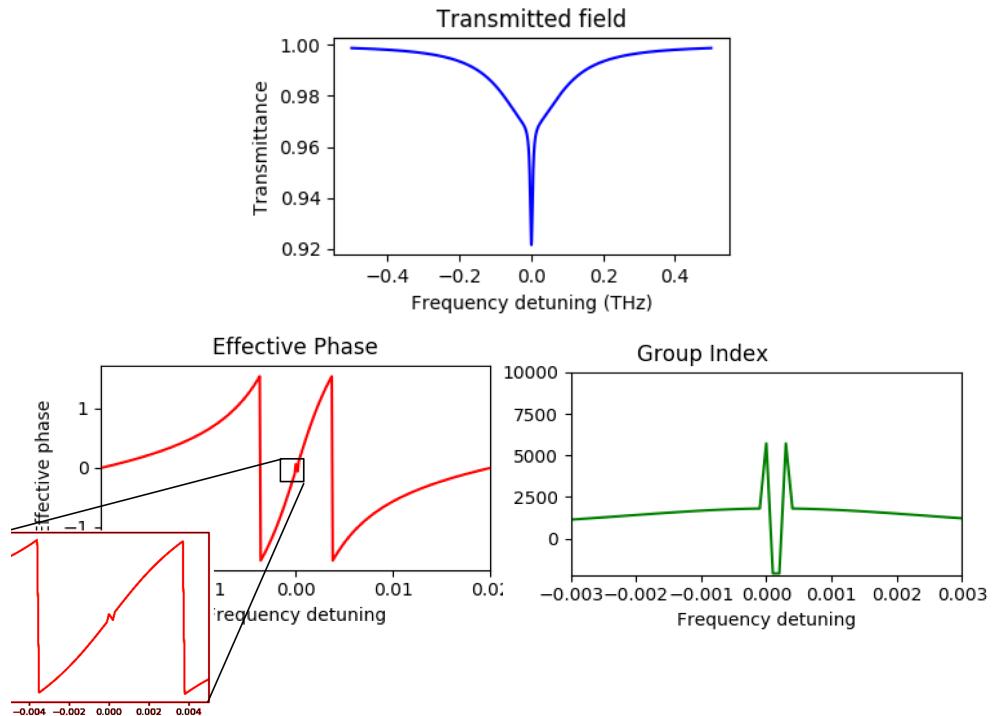


Figure 3.22: Transmission rises towards unity, dispersion becomes anomalous on resonance, and group index displays negative values.

With further increase in the gain, the spectrum flips over the horizontal axis and EIA turns into an EIT like transmission Fig. 3.23. The dispersion remains anomalous as the gain is further increased, i.e

$g > \alpha_i$. After the gain achieves a threshold value, again a transition occurs and anomalous dispersion is converted into normal dispersion (Fig. 3.24).

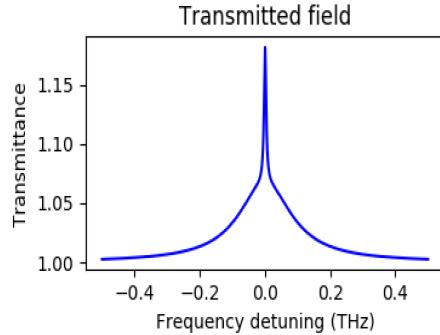


Figure 3.23: Transmission of the system

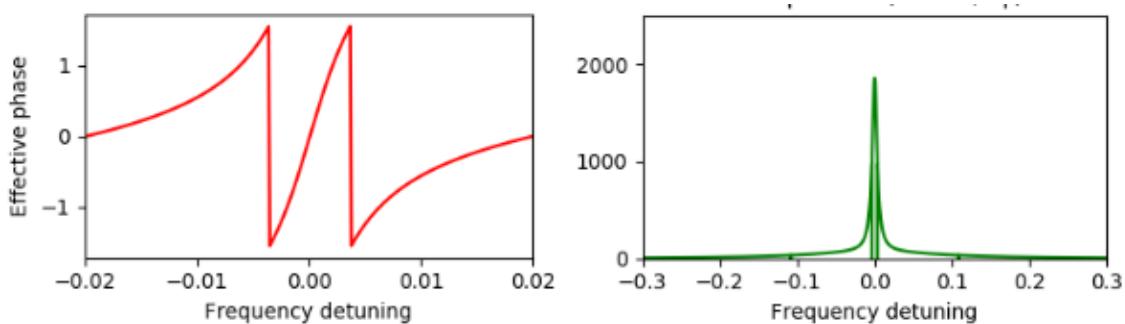


Figure 3.24: Phase and group index of the system.

It is worth noting that the EIT-like resonance shown in Fig. 3.23, usually gives slow light in passive resonators [8]. However, in the present case, fast light is obtained owing to the incorporation of a gain element into the system until a specific threshold gain value.

3.6.5 CRIA With Fast Light

Now we will consider a case featuring CRIA which leads to fast light in a passive coupled resonator system and studies the behavior when the gain is introduced in either one of the resonators. The coupling parameters are given as, $r_1 = 0.8998$ and $r_2 = 0.99998$ while the quality factors remains the same. (Fig. 3.25)

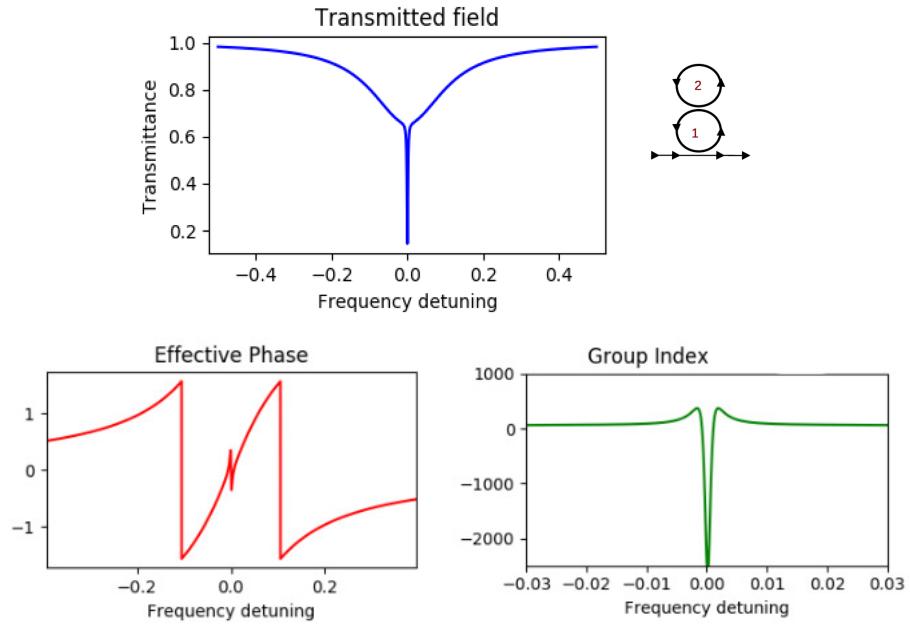


Figure 3.25: CRIA observed in a passive two resonator system with anomalous dispersion and negative group index.

3.6.6 Introducing Gain Only In Second Resonator

Now we activate gain in the second resonator, shown in the schematic (Fig. 3.26), and observes that the EIA resonance narrows down and becomes sharp. While the amplitude of the narrow CRIA feature decreases, resulting in higher transmission.

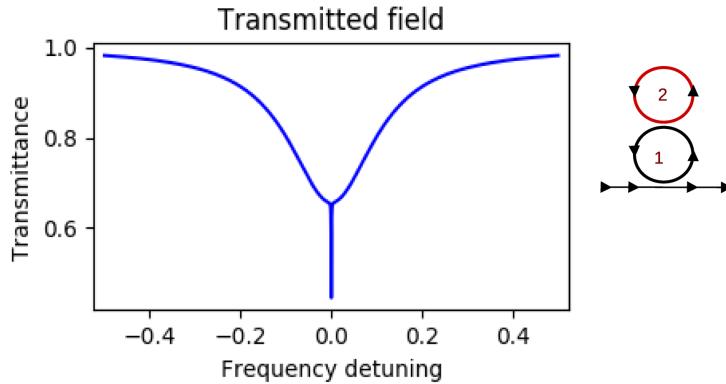


Figure 3.26: Transmission spectrum of CRIA with gain in second resonator.

The dispersion of the system changes as gain becomes close to intrinsic loss ($g \approx \alpha_i$) and normal dispersion and a positive group index are obtained (Fig. 3.27).

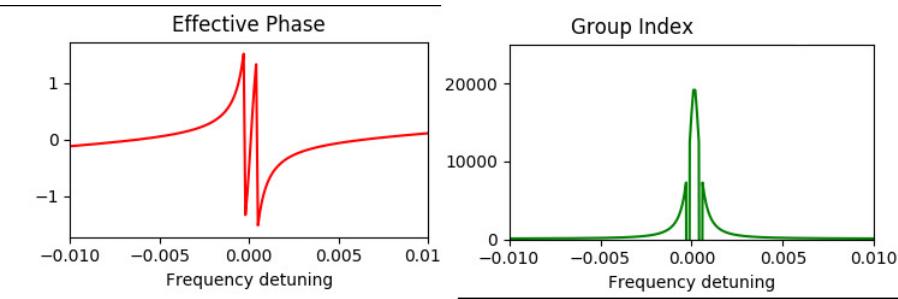


Figure 3.27: Transition from fast to slow light in CRIA.

However, when the gain is further increased and it exceeds the intrinsic losses, such that $g > \alpha_i$, the spectrum then shows a conversion from CRIA into CRIT (Fig. 3.28) with normal dispersion and positive group index.

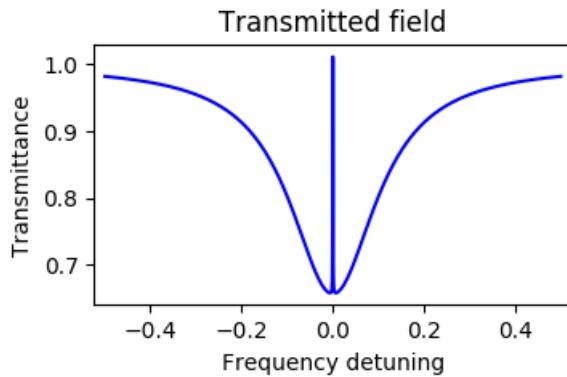


Figure 3.28: Transmission dip transforming into an transmission peak.

However, the dispersion of the system remains unaffected. It is worth noting that although the dispersion remains subluminal for higher transmission, the transmitted light becomes transparent due to CRIT resonance.

3.6.7 Introducing Gain Only In First Resonator

We again start with the passive CRIA of Fig 3.25. However, now we introduce gain in the first resonator only. We observe the CRIA resonance narrows down and becomes a sharp dip. Also, two off-resonances peaks appear in the spectrum. As $g \approx \alpha_i$, the resonance becomes further narrower and the system moves towards critical coupling. We still see fast light and negative group index from here but

most of the light is absorbed. The transmission spectrum in Fig. 3.29 is shown for $g > \alpha_i$.

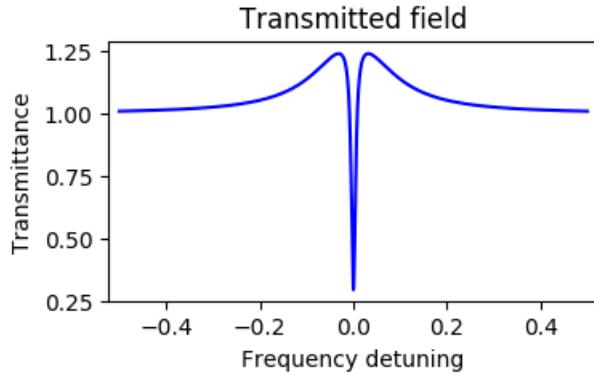


Figure 3.29: Transmission dip of CRIA with gain in first resonator.

3.6.8 Introducing Gain In Both Resonator

If we introduce gain in both of the resonators of an initially passive system (Fig. 3.25), simultaneously, we observe the narrowing of CRIA dip and rise of the transmission towards unity. However, the transmission dip changes into a peak when the gain is greater than intrinsic losses ($g > \alpha_i$) owing to the transparency of the system (Fig. 3.30). The CRIT resonance of this transparent system leads to normal dispersion and positive group index of $\approx 2 \times 10^4$ (Fig.3.31).

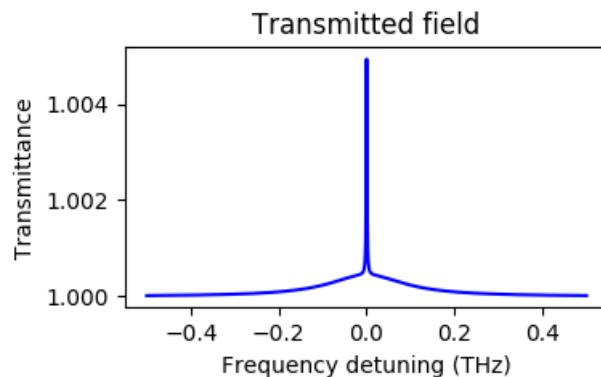


Figure 3.30: Transmission dip transforming into a transmission peak.

These observations lead to the conclusion that when we introduce gain into the system, with it being activated in either one of the resonators or both, we observe drastic changes in the transmission and

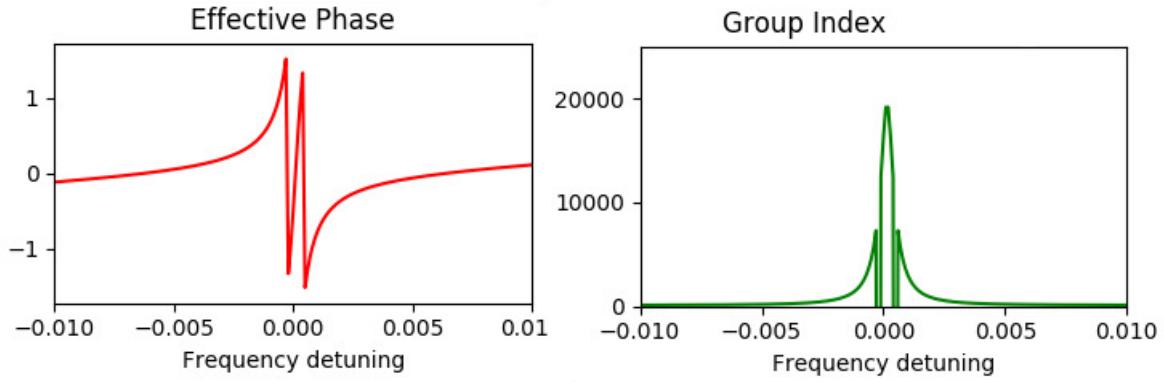


Figure 3.31: Respective phase and group index of the system.

phase spectrums thus affecting the group delay and dispersion of the system. This allowed achieving gain-assisted tunability of the dispersion of the system, meaning the transitions between fast and slow light can be controlled by simply introducing gain into a passive coupled resonator system.

3.7 Proposed Applications

3.7.1 Gain-controlled Optical Pulse Storage and On-demand Retrieval

Coupling parameter between the resonators r_2 , which is the internal coupling of the two rings, can induce very distinct changes into the system.

Now we will discuss a case in which the coupling between the first and the second resonator is very weak and the resonance of the second resonator is almost zero. Meaning we have a single resonance transmission at critically coupled i.e all of the light is absorbed. Fig. 3.32 displays the transmission spectrum along with its phase displaying anomalous dispersion.

Now we will excite gain into the second resonator such that the gain coefficient g_2 , is equal to the intrinsic loss coefficient α_i . We observe a transmission EIT peak of 100 percent intensity meaning now we have the light pulse storage achieved via gain tuning. This process happens at nanoseconds scale meaning we can switch from zero to maximum intensity in nanoscale durations. Fig. 3.33 displays the gain excited

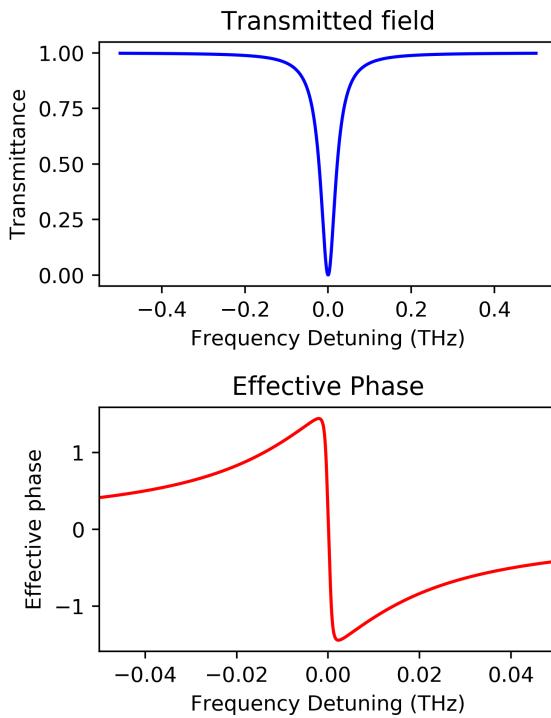


Figure 3.32: Single resonance displayed in a coupled resonator system with its transmission phase.

transmission with normal dispersion meaning we have slow light which was the main ambition.

Thus the light pulse is stored inside the coupled resonator system and to retrieve the stored pulse back, we just simply need to excite gain into the second resonator of the system. The true shape of the input signal is not disturbed and the narrow peak of EIT shows us that very narrow bandwidth of frequencies can be stored.

3.7.2 Gain Controlled Slow and Fast Light Tuning

We have observed from these results that the transitions from super and subluminal velocities of light have been in the control of the optical gain of the system. We have observed superluminal regimes of transmission reverting into subluminal transmission by the excitation of gain inside either one or both of the coupled two resonator system. Also, the subluminal transmissions such as EIT which always displays slow light in passive systems provided us with fast light and anomalous dispersion in the presence of gain in either of the resonators. This tun-

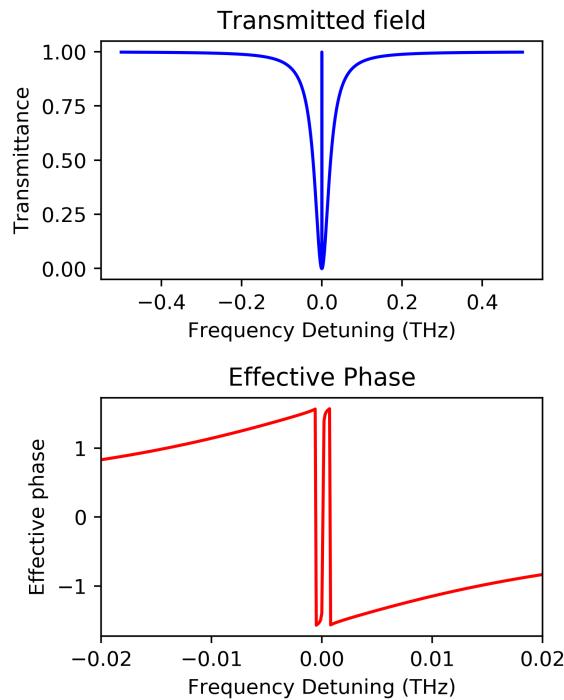


Figure 3.33: EIT peak emerged when gain is excited in second resonator in a coupled resonator system with its transmission phase.

ing of gain made possible the transitions between fast and slow light in both ways, meaning we have achieved control over how and when do we want fast and slow light for our device.

3.8 Discussion

We exhibit increased tunability of optical analogs of Electromagnetically Induced Transparency (EIT) and Electromagnetically Induced Absorption (EIA). This enabled us to unequivocally control superluminal and subluminal group velocities of light in a coupled ring resonator framework, with a direct increase excitation and empowered reversible advances between them. Besides, we watched sub and superluminal light including all-optical EIT reverberation, which in every single past examination, in view of inactive coupled ring resonators, has been seen to yield just subluminal light.

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Chapter 4

Cascaded Resonances in Three Coupled Resonators

In this chapter, we will now extend our study on composite resonator systems. Such as, we will increase the number of resonators in our system. These structures, that will be discussed here, were not optimized to achieve enhanced dispersion. Rather the purpose of these analyses is to demonstrate the versatility and distinct optical characteristics of cascaded resonances which are obtained by cascading three resonators (Fig. 4.1). For the scope of this thesis, the system will consist of ring-shaped resonators and thus we will study properties of such ring resonators and their mutual coupling effects.

4.1 Triple Ring Resonator System

Now we introduce a new resonator geometry. Basically, we are to simply add another ring above the coupled two resonator system which was discussed in greater detail Chapter 3. Now we have a three-resonator system with each having their own distinct resonant frequencies. These resonators show distinct properties due to the introduction of an additional ring with a higher Quality-factor, coupled to the two resonator system. This allows us to observe multiple resonances and observe phenomenons like CRIT and CRIA with another perspective. This also enables us to simultaneously measure these effects in a single system and thus obtain versatile transmission and dispersion.

The illustration in Fig. 4.1 shows us that the three rings are mutually coupled to each other and only the first ring is coupled to the

optical waveguide. The incident energy from the input is labeled as E_1 which couples to the first resonator due to evanescent coupling. The energy is then transferred to E_3 . This energy travels the first ring and is again coupled into the as E_5 and then again into the third resonator with energy E_7 . Then it loops back into the waveguide as E_8 , E_6 and E_4 respectively to couple back to the waveguide, each acquiring a distinct phase shift and outputs the signal with energy E_2 . Here, E_3 and E_4 describes the circulating fields of the first resonator. Evanescent coupling between the first and the second resonator results in the circulating fields of the resonator two labeled E_5 and E_6 . A similar mechanism leads to the excitation of the resonant mode of the third cavity, with circulating fields described by E_7 and E_8 .

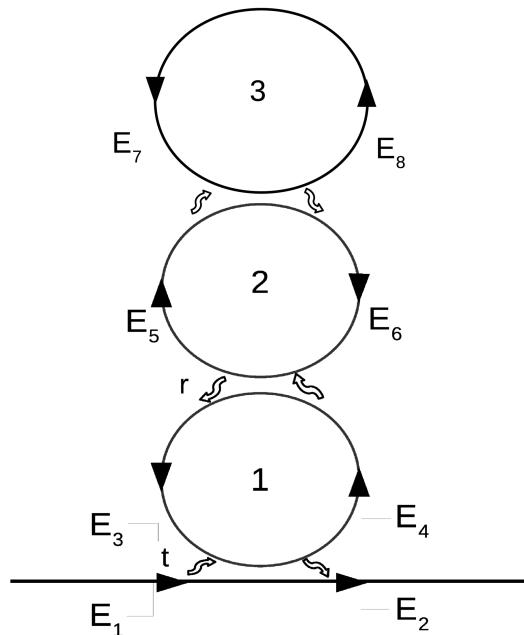


Figure 4.1: Basic illustration of three ring resonator geometry along with its respective fields. Here, first resonator is labeled as 1, second is labeled as 2, and third is labeled as 3.

4.1.1 Transmission and Phase relations

The complex transmittivity and its respective phase can be determined by equations 4.1 and 4.2 respectively.

$$\frac{E_t}{E_i} = \frac{r_1 - a_1 r_{12} e^{i\phi_1}}{1 - r_1 r_{12} a_1 e^{i\phi_1}} \quad (4.1)$$

where,

$$r_{12} = \frac{r_2 - a_2 r_{23} e^{i\phi_2}}{1 - r_2 r_{23} a_2 e^{i\phi_2}}$$

and,

$$r_{23} = \frac{r_3 - a_3 e^{i\phi_3}}{1 - r_3 a_3 e^{i\phi_3}}$$

Similarly, the effective phase of the complex transmittivity is given by:

$$\phi_{eff} = \arctan\left[\frac{r_1|r_{12}|a_1 \sin(\phi_1 + \phi_{12})}{1 - |r_{12}|a_1 \cos(\phi_1 + \phi_{12})}\right] - \arctan\left[\frac{|r_{12}|a_1 \sin(\phi_1 + \phi_{12})}{r_1 - |r_{12}|a_1 \cos(\phi_1 + \phi_{12})}\right] \quad (4.2)$$

Now we use these equations to observe characteristics of triple cavity resonances.

4.2 Passive Three Resonator Results

We can obtain very interesting results from a passive three resonator systems some of which are discussed in this section. Fig. 4.2 displays CRIT inside a CRIA resonance hence negative group index and fast light are obtained on resonant frequencies while large values of subluminal group indexes are clearly apparent at the off-resonance transmission. When resonator two is allowed to couple with the third resonator, due to its effects there is an on-resonance peak inside the CRIA dip. These cascaded resonances can help play an important role in the tunability of fast and slow light and/or in large and small absorption of resonant frequencies. This may lead to new applications

in communication technology and related fields. These kinds of effects which were observed in atomic systems had shortcomings due to a large amount of absorption and had low-temperature maintenance problems. Now, these effects can be observed in photonic resonators which are operatable at room temperatures.

4.2.1 CRIT inside CRIA

In Fig. 4.2, we see that the transmission spectrum looks like a CRIA dip, meaning it will display the properties of CRIA. The system parameters are given as, $r_1 = 0.898945$, $r_2 = 0.999958$, and $r_3 = 0.999999$. The Q -factors are 1×10^5 , 1×10^6 and 1×10^7 for first second and third resonator respectively. When we zoom into the graph (as shown on the right), we observe that there is another peak rising within the CRIA, showing the feature due to the third resonance of the third resonator. Thus now we have CRIT inside a CRIA transmission. This will allow us to have transmission intensity larger while maintaining the characteristics of CRIA dispersion. Now we can have more light in CRIA mediate transmission. This means we can have fast light dispersion on on-resonant frequencies, which can be enhanced by adjusting mutual coupling of the resonators.

The effective phase of the system is shown in red where we see a normal curve stretching from positive to negative horizontal and vertical axes. But when we zoom into the middle of the curve, we see two features instead of one which characterize the presence of three cavities. The zoomed version is shown on the right in Fig. 4.2 and it displays a negative slope on resonance meaning the light we are receiving in the transmission is fast light. Thus we conclude that we have superluminal light along with large transmission owing to cascaded CRIT and CRIA.

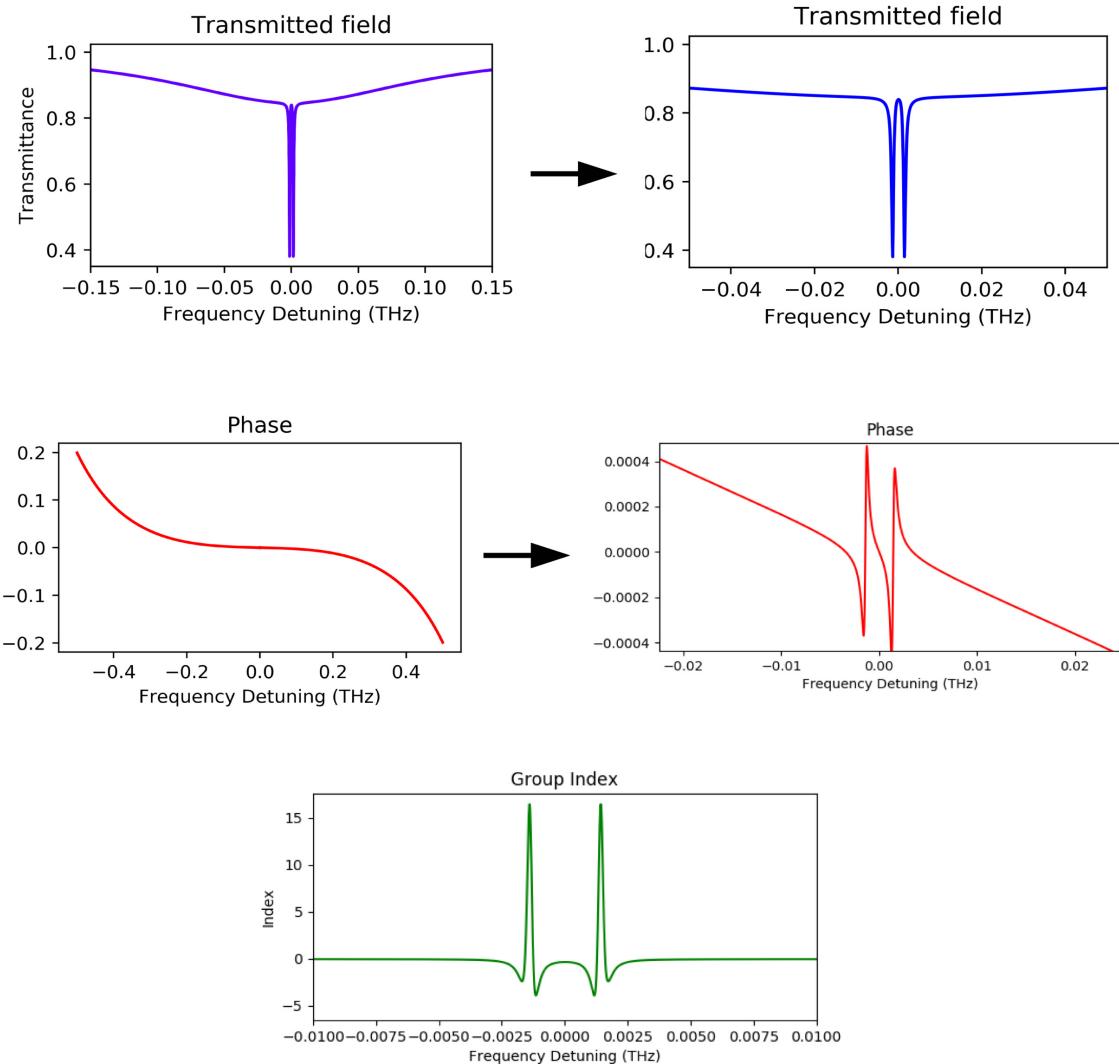


Figure 4.2: CRIT observed in an CRIA transmission in three resonator system with its phase in red and group index in green.

4.2.2 CRIA inside CRIT

After these interesting results, let us now move towards another useful transmission spectrum of our triple resonator system (see Fig. 4.3). The system parameters are given as, $r_1 = 0.998845$, $r_2 = 0.999988$, and $r_3 = 0.999999$. The Q – factors are 1×10^5 , 1×10^6 and 1×10^7 for first second and third resonator respectively. This spectrum is also achieved by the same arrangement of the resonators and now we have changed the couplings once more to obtain an interesting result.

In Fig. 4.3, we see a CRIT spectrum. When we again zoom into

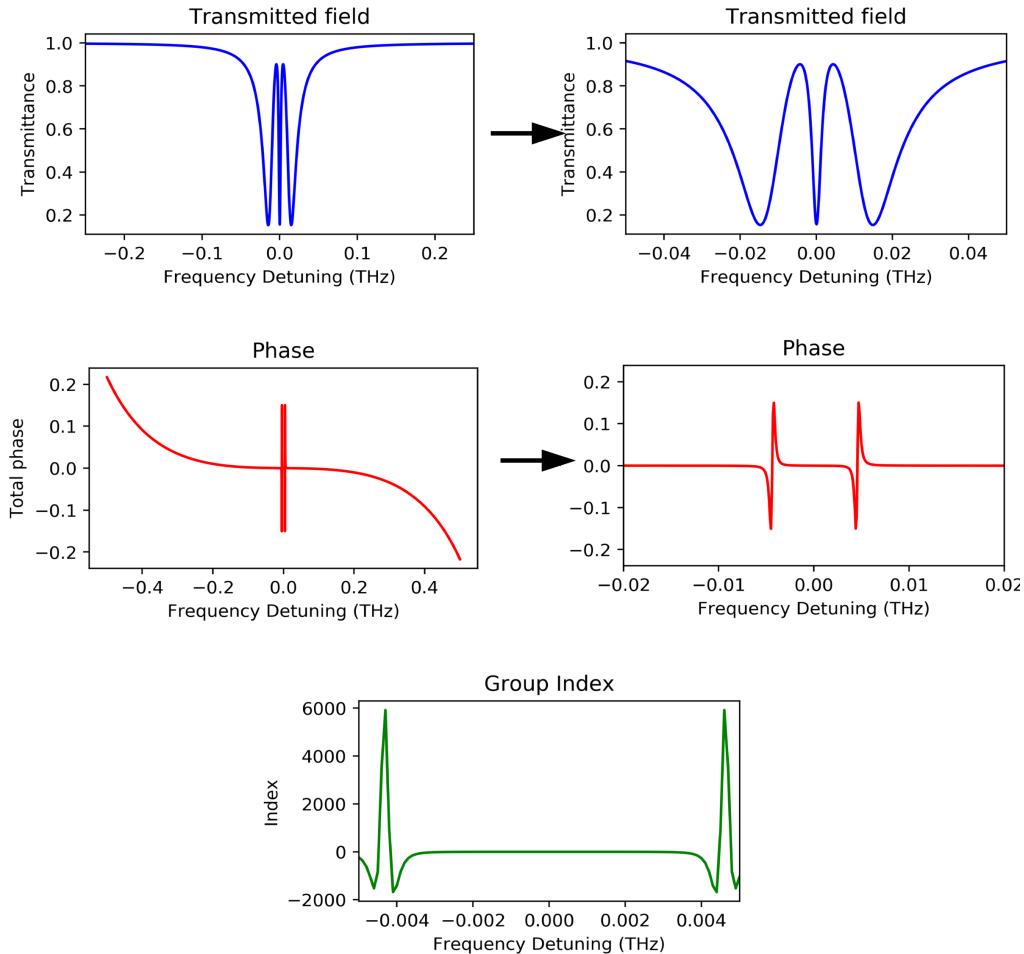


Figure 4.3: CRIA observed in an CRIT transmission in three resonator system with its phase in red and group index in green.

the transmission, we see a narrow dip on resonance. This narrow dip shows us that we actually have a CRIA resonance at the center of a CRIT resonance.

We see that the graph dips to zero almost, meaning all of the light is absorbed. This transmission dip tells us that we can filter out exactly this narrow spectrum of resonant frequencies. The phase of the system is shown in red and we see features of superluminous on a similar phase curve this time of quite high value. The slope of this graph is almost zero on resonance.

The group index, shown in green, also displays group index very close to zero. Although its value is ≈ -0.37 meaning superluminal velocities of resonant frequencies. However, almost all of the reso-

nant light circulates in the coupled resonator system and eventually dissipates. Thus this value is practically useless.

4.2.3 Cascaded CRIA and CRIT in Absorption

In Fig. 4.4, we will observe that the transmission can be a lot influenced if we were to change the coupling effects between the resonators. The system parameters are given as, $r_1 = 0.999945$, $r_2 = 0.999999$, and $r_3 = 0.999999$. The *Q – factors* are 1×10^5 , 1×10^6 and 1×10^7 for first second and third resonator respectively. This tells us a lot about how our signal is transmitted and how much use can we achieve from the single system by changing a few of its properties. In Fig. 4.3, we change the coupling parameters in all of the resonators and have achieved a rather interesting transmission spectrum.

We can clearly see what looks like an absorption spectrum of a single resonator, has a narrow, actually CRIT peak within the broader CRIA transmission dip (zoomed on the right). This peak also has a dip on resonance due to the coupling of the third resonator. Now we observed cascaded results of CRIT within an absorption and CRIA inside that CRIT.

The effective phase of the system, at first, also seems like that of a single resonator system but zooming into the graph tells another story. On the right, we can see three distinct features caused by coupling between the three resonators and have a negative slope on resonance. This negative slope refers that superluminal light is achieved the resonance.

The group index plot, shown in green, also displays a negative group index on resonance and positive index peaks off resonances. The negative group delay also predicts superluminal velocities of the resonant frequencies.

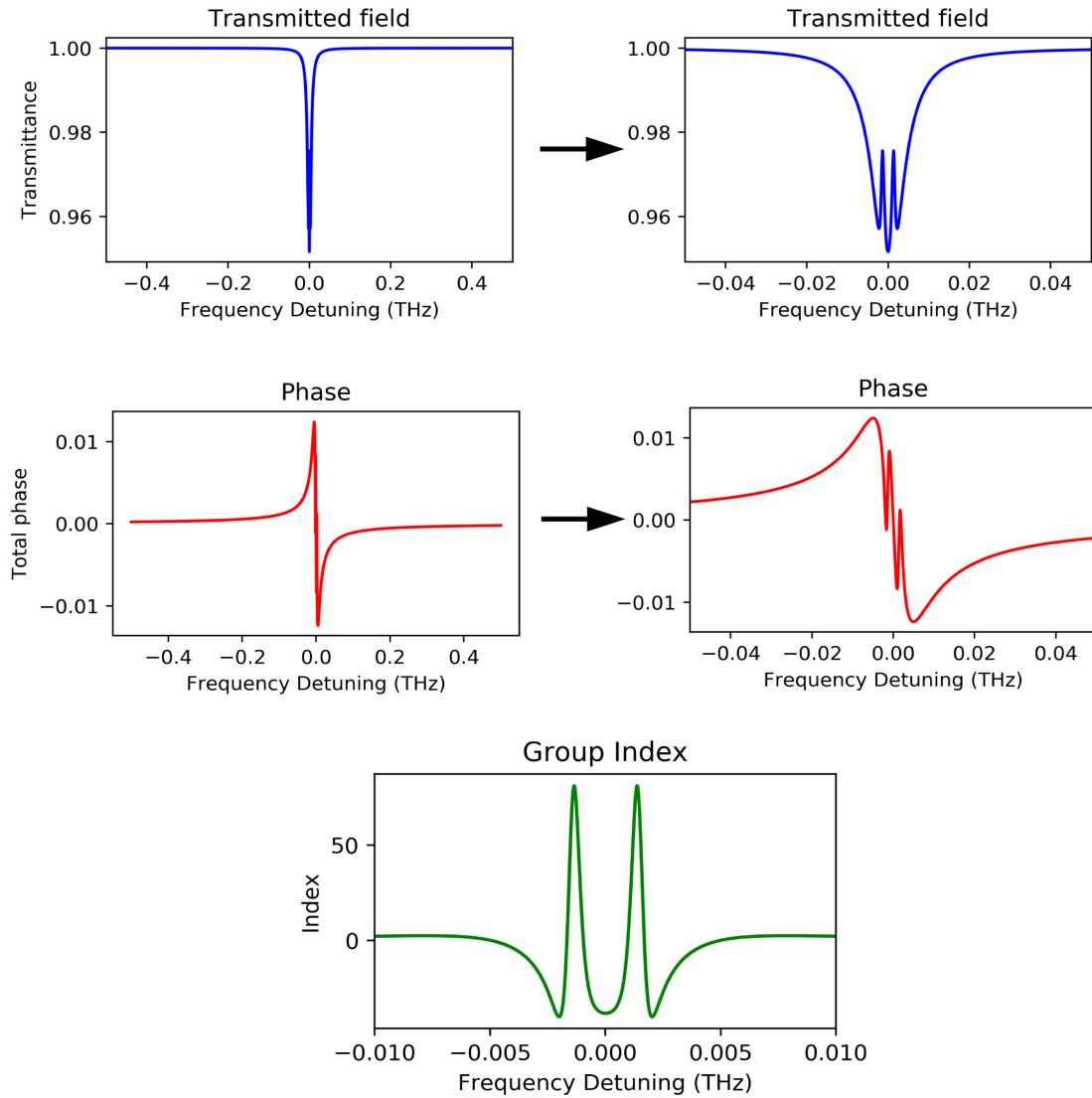


Figure 4.4: Cascaded resonance effects in the three-resonator system with its phase in red and group index in the green. Magnified views of resonances and phase are also given for clear illustration.

4.2.4 Double CRIA

We now demonstrate a double CRIA resonance (see Fig. 4.5). This means that two narrow dips are realized within a broader dip. This is in contrast to a typical CRIA where only one narrow dip is present inside a broader dip. This double EIA or CRIA is obtained by slight detuning the system and changing the coupling effects.

In Fig. 4.5 we can clearly see that two narrow peaks, which are

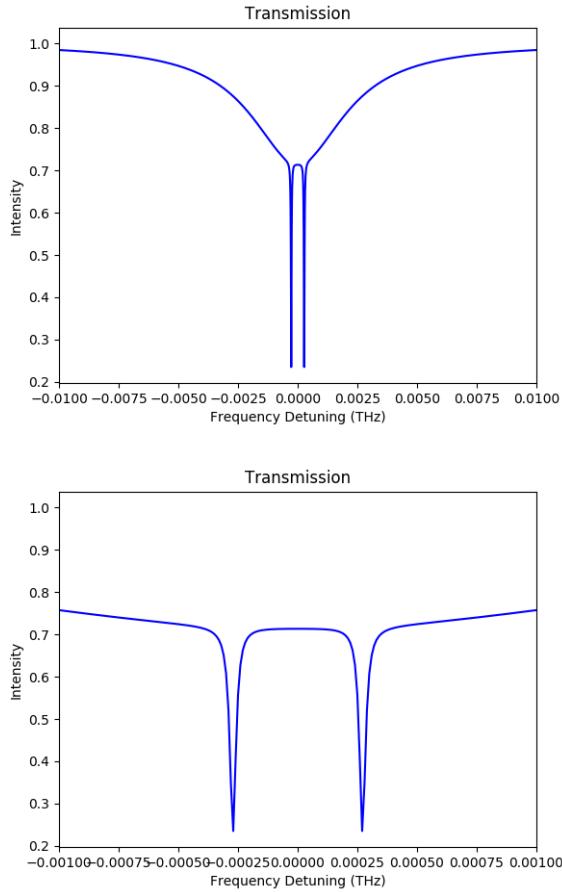


Figure 4.5: Double transmission dips observed inside an EIA like transmission off resonant to the spectrum.

caused by the two resonances of the coupled resonators, 2 and 3 respectively. And the broader dip above them is caused by the resonance of the first resonator.

This allowed us to observe an effect which resembles the CRIA in two resonator system but this time now we have two narrow dips off resonances. This tells us that we have CRIA like properties and transmission have two narrow absorption lines but on the off-resonance. The resonant frequencies will experience very little absorption and will be mostly transmitted.

The phase and group velocities of this case are not discussed as the purpose of this study were not to discuss the enhanced dispersion. The dispersive properties of these resonances will be similar to the ones discussed above. The true meaning is to show enhanced transmittance.

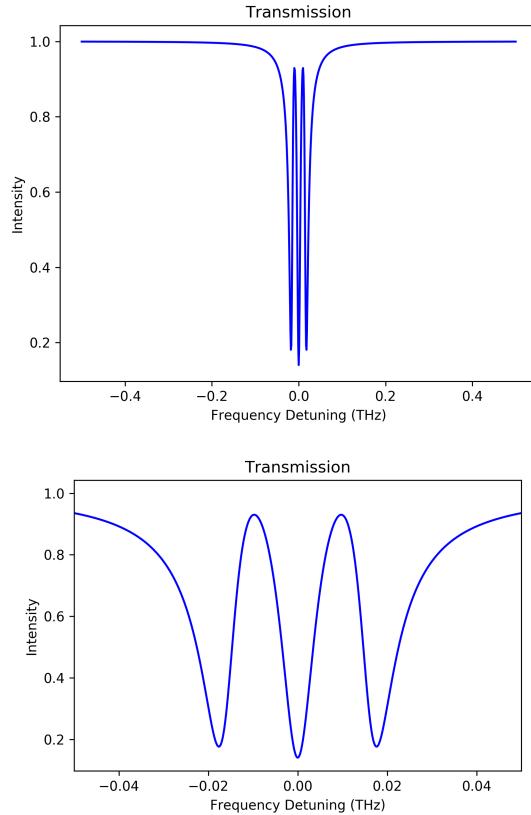


Figure 4.6: Double transmission peaks observed inside an EIT like transmission off resonant to the spectrum.

4.2.5 Double CRIT

We now demonstrate a double CRIT resonance in Fig. 4.6. This means that two narrow peaks are realized within another narrow dip on-resonance. This is in contrast to a typical CRIT where only one narrow peak is present. This double EIT or CRIT is obtained by changing the coupling effects.

From these results, we concluded that increasing the number of resonators in the system and mutually coupling them with each other, enables us to demonstrate the versatility of the cascaded resonances. These resonances can be then utilized and achieved in practical applications to make use of them in important tasks. Optical tunability and signal filtration of specific frequencies can be achieved by using similar systems and having similar effects on them.

4.3 Discussion

We have examined the spectral and dispersive properties of a course of three mutually coupled ring resonators. We have discovered various one of a kind cascaded resonances which show particular spectral and dispersive conduct. In view of our request, we propose new uses of coupled resonators for future quantum and optical information and communication technologies. Interacting triple cavities were previously investigated using one and two-dimensional photonic crystals for parametric oscillations [5] and group delay control [6], respectively. Furthermore, double anti-crossing behavior was recently demonstrated using triple microtoroid cavities [7]. Tunable slow and fast light in triple one-dimensional photonic crystal microcavities has also been demonstrated owing to the tuning of the resonant wavelength.

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Chapter 5

Conclusion

We have demonstrated gain tunable optical analogs of Electromagnetically Induced Transparency (EIT) and Electromagnetically Induced Absorption (EIA). This allows us to precisely control sub and superluminal group velocities of light pulses that propagate through a coupled ring resonator system, which incorporates a linear gain element.

Furthermore, gain incorporating coupled resonators also enable reversible transitions between the sub and superluminal dispersion regimes. Additionally, we observed sub and superluminal light featuring all-optical EIT resonance, which in all previous studies based on passive coupled resonators, has been observed to yield only subluminal light. We also studied the spectral and dispersive properties of a cascade of three mutually coupled ring resonators. We have found a number of unique cascaded resonances, which display distinct spectral and dispersive behavior. The optimized coupled resonators demonstrate continuously tunable sub and superluminal group velocities owing to excitation of linear optical gain. These features of coupled resonators also allow enabling control of transmission and optical pulse storage times. Furthermore, cascades of ring resonators may be used to further enhance these characteristics.

Based on our study, we proposed and discussed new applications of coupled resonators. The findings reported in this thesis seem to be significant for the future optical and quantum communication systems.

Appendix A

Abrevations

EIT: Electromagnetically Induced Transparency

EIA: Electromagnetically Induced Absorption

CRIT: Coupled Resonator Induced Transparency

CRIA: Coupled Resonator Induced Absorption

FSR: Free Spectral Range

MRR: Micro Ring Resonator

ICs: Integrated Circuits

FWHM: Full width at half maximum

WDM: Wavelength division multiplex

TIR: Total Internal Reflection

PBG: Photonic Band Gap

PC: Photonic Crystals

WGR: Whispering Gallery Resonator

APPENDIX A. ABREVATIONS

FPI: Fabry- Perot Interferometer

GTI: Gires- Tournois Interferometer

ABG: Anullar Bragg Grating

ATS: Autler- Townes Splitting

CR: Cascaded Resonances

Appendix B

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APPENDIX B. BIBLIOGRAPHY

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