

# MAT292 Notes

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## 1 Existence and Uniqueness Theorem

1. We need  $f(t, y)$  continuous in the rectangle to get existence
2. We need  $f_y(t, y)$  continuous in the rectangle to get uniqueness
3. E & U Theorem is sufficient but not necessary. i.e. these conditions imply solution but not having these conditions doesn't mean there is no solution

## 2 Autonomous Equations and Population Dynamics

### 2.1 Logistic Growth

If uninhibited, we assume exp. growth however in the long run, population is limited to  $K$

Model:  $y' = rh(y)y$

We want  $h(y) \approx 1$  if  $y$  is small,  $h(y) < 1$  if  $y < k$ ,  $h(y) = 0$  if  $y = k$  and  $h(y) < 0$  if  $y > K$

This can thus be modelled as  $y' = r(1 - \frac{y}{k})y$ . This has two equilibria namely at  $y = 0$  and  $yk$ . The inflection points can be found by setting the derivative  $y''$  to 0.

## 3 Direction Fields and Orbits

### 3.1 Reducing non homogeneous systems to homogeneous systems

Lets take a solution  $x$  and write it as  $x = \phi + v$  where  $v$  is a constant. Then  $x' = Ax + b \rightarrow \phi' = A(\phi + v) + b$ . Since  $x_{eq} = A^{-1}b$ ,  $Av + b = 0$  by the equilibrium condition ( $\phi' - A\phi$ ) we have that  $\phi' = A\phi$ . So that  $x = \phi + x_{eq}$  where  $\phi$  is a solution of the homogeneous system.

Every solution of the non homogeneous problem can be written as a solution of the homogeneous problem plus the equilibrium.

## 4 Laplace Transform

- Remark: The laplace transform will allow us to reduce solving an ODE to solving an algebraic equation
- Solve algebraic equation and use the inverse laplace transform to get the solution to the ODE
- Definition: If  $f$  is defined on  $[0, \infty]$ , the Laplace Transform is defined as  $F(s) = \int_0^\infty e^{-st} f(t) dt$
- We write  $F = \mathcal{L}\{f\}$
- We use uppercase letters for Laplace transform e.g.  $G(s)$  is the LT of  $g(t)$
- Example: For  $f(t) = e^{at}$ , we get  $F(s) = \mathcal{L}\{f\}(s) = \int_0^\infty e^{-st} e^{at} dt = \lim_{b \rightarrow \infty} \int_0^b e^{(a-s)t} dt = \lim_{b \rightarrow \infty} \frac{1}{a-s} (e^{(a-s)b} - 1) = \frac{1}{s-a}$  if  $s > a$
- $\mathcal{L}\{1\} = \frac{1}{s}$
- Theorem:  $\mathcal{L}\{c_1 f_1 + c_2 f_2\} = c_1 \mathcal{L}\{f_1\} + c_2 \mathcal{L}\{f_2\}$
- To find  $\mathcal{L}\{\sin(at)\}$ , write  $\sin(at) = \frac{1}{2i}(e^{iat} - e^{-iat})$  and use the theorem above
- This will give  $\frac{1}{2i} \left( \frac{1}{s-ia} \right) - \frac{1}{2i} \left( \frac{1}{s+ia} \right) = \frac{a}{s^2+a^2}$  for  $s > 0$

- Example: LT of  $f(t) = e^{2t}$  for  $0 \leq t < 1$  and  $f(t) = 4$  for  $1 \leq t$
- Divide the integral into two separate parts and evaluate it
- Exponential order: A function  $f(t)$  is of exponential order for  $M > 0$ ,  $K > 0$  and  $a \in \mathbb{R}$  if  $|f(t)| \leq Ke^{at}$  for  $t \geq M$  i.e.  $f$  eventually becomes between two exponential functions
- Theorem: Every bounded function is of exponential order
- A function  $f(t)$  is piecewise continuous on  $[a, b]$  iff there are finitely many "jump points" between  $a$  and  $b$   $a \leq t_0 < t_1 < \dots < t_{k-1} < t_k = b$  such that  $f$  is continuous on each of the intervals  $(t_i, t_{i+1})$  and  $f$  has finite limits at the jump points.
- Theorem: If for a function  $f(t)$ , we have that  $f$  is piecewise continuous on  $[0, A] \forall A \geq 0$  and  $f$  is of exponential order for  $M, k$  and  $a$ . Then  $\mathcal{L}\{f\}$  exists for all  $s > a$ .
- Theorem: If  $f(t)$  is of exponential order then we have:  $F(s) \rightarrow 0$  as  $s \rightarrow \infty$  where  $F(s)$  is the L.T. of  $f$
- Theorem: If  $f$  is continuous and  $f'$  is piecewise continuous on any interval  $[0, A]$  and  $f, f'$  are of exponential order for  $M, k, a$  then  $\mathcal{L}\{f'\}(s) = s\mathcal{L}\{f\}(s) - f(0)$  for  $s > a$ . Under the same conditions for  $n$  derivatives,  $\mathcal{L}\{f^{(n)}\}(s) = s^n \mathcal{L}\{f\}(s) - s^{n-1}f(0) - s^{n-2}f'(0) - \dots - sf^{(n-2)}(0) - f^{(n-1)}(0)$
- Proof:  $\mathcal{L}\{f'\}(s) = \int_0^\infty e^{-st} f'(t) dt = \lim_{b \rightarrow \infty} \left( \int_0^b e^{-st} f'(t) dt \right)$   
 $= \lim_{b \rightarrow \infty} \left( [e^{-st} f(t)]_0^b + \int_0^b f(t) s e^{-st} dt \right) = \lim_{b \rightarrow \infty} \left( e^{-bs} f(b) - f(0) + s \int_0^b f(t) e^{-st} dt \right)$   
 $= s\mathcal{L}\{f\}(s) - f(0)$  where  $s > a$  (by definition of exponential order)

## 5 Inverse Laplace Transform

- Theorem: If  $f(t), g(t)$  are piecewise continuous and of exponential order, then  $\mathcal{L}\{f\} = \mathcal{L}\{g\} \implies f(t) = g(t)$
- Technicality: Take  $f(t) = e^t$ ,  $g(t) = \begin{cases} e^t & t \neq 5 \\ 0 & t = 5 \end{cases}$ . Clearly  $\mathcal{L}\{f\} = \mathcal{L}\{g\}$  but  $f(t) \neq g(t) \forall t$
- Convention: We write  $f(t) = g(t)$  as long as they are the same whenever they are continuous
- Definition: If  $f$  is piecewise continuous and of exponential order and  $\mathcal{L}\{f\}(s) = F(s)$ , then we call  $f(t) = \mathcal{L}^{-1}\{F\}(t)$
- There is a complex analysis formula (Mellin Transform) to find  $\mathcal{L}^{-1}\{F\}$ . However this is rarely used in practice and we instead use tables

## 6 Solving ODEs with Laplace Transform

- Lets solve the IVP  $y'' + 2y' + 5y = e^{-t}$ ,  $y(0) = 1$ ,  $y'(0) = -3$ 
  - Use the Laplace transform:  $\mathcal{L}\{y'' + 2y' + 5y\} = \mathcal{L}\{e^{-t}\}$
  - $s^2Y(s) - sy(0) - y'(0) + 2(sY(s) - y(0)) + 5Y(s) = \frac{1}{s+1}$
  - $s^2Y(s) - s \cdot 1 - (-3) + 2(sY(s) - 1) + 5Y(s) = \frac{1}{s+1}$  using the initial conditions
  - Solving this gives  $Y(s) = \frac{s^2}{(s+1)(s^2+2s+5)}$
  - Simplifying and using partial fractions,  $Y(s) = \frac{1}{4} \frac{1}{s+1} + \frac{3}{4} \frac{s+1}{(s+1)^2+4} - \frac{2}{(s+1)^2+4}$
  - $y(t) = \mathcal{L}^{-1}\{Y(s)\} = \frac{1}{4} \mathcal{L}^{-1}\{\frac{1}{s+1}\} + \frac{3}{4} \mathcal{L}^{-1}\{\frac{s+1}{(s+1)^2+4}\} - 2 \mathcal{L}^{-1}\{\frac{1}{(s+1)^2+4}\}$
  - $= \frac{1}{4}e^{-t} + \frac{3}{4}e^{-t} \cos(2t) - e^{-t} \sin(2t)$
- Lets solve the IVP  $y'''' + 2y' + y = 0$ ,  $y(0) = 1$ ,  $y'(0) = -1$ ,  $y''(0) = 0$ ,  $y'''(0) = 2$ 
  - Applying the Laplace transform:  $s^4Y(s) - s^3y(0) - s^2y'(0) - sy''(0) - y'''(0) + 2(s^2Y(s) - sy(0) - y'(0)) + Y(s) = 0$
  - Using the initial conditions,  $s^4Y(s) - s^3 + s^2 - 2 + 2(s^2Y(s) - s + 1) + Y(s) = 0$
  - Solving gives  $Y(s) = \frac{s^3-s^2+2s}{(s^2+1)^2}$
  - We use a repeated partial fraction decomposition to write  $Y(s)$  as  $\frac{As+B}{s^2+1} + \frac{s+1}{(s^2+1)^2}$
  - We know  $y(t) = \mathcal{L}^{-1}\{\frac{s}{s^2+1} - \frac{1}{s^2+1} + \frac{s}{(s^2+1)^2} + \frac{1}{(s^2+1)^2}\}$
  - For the third term, we use  $\mathcal{L}\{t \cdot f(t)\} = \frac{-d}{ds}F(s)$  and then get  $y(t) = \cos t - \sin t + \frac{1}{2}t \sin t + \frac{1}{2} \sin t - \frac{1}{2}t \cos t$

## 7 Discontinuous Forcing Functions

- Finding  $\mathcal{L}\{t\} = \int_0^\infty e^{-st}t dt = \frac{1}{s^2}$  (IBP) or using  $\mathcal{L}\{t\} = \mathcal{L}\{t \cdot 1\}$  and  $\mathcal{L}\{t \cdot 1\} = -\frac{d}{ds}(\frac{1}{s}) = \frac{1}{s^2}$
- Consider  $y'' + 4y = g(t)$  with  $y(0) = 0$ ,  $y'(0) = 0$  and  $g(t) = \begin{cases} 0 & 0 < t < 5 \\ \frac{t-5}{5} & t = 5 \leq t < 10 \\ t & 10 \leq t \end{cases}$ 
  - Can split it up into 3 parts and find relevant boundary conditions in each case to be used in the next case
- Piecewise defined functions can simulate the activation of a signal
  - $u_c(t)$  is a step function that is 0 till  $x = c$  and then 1 afterwards
  - $u_{cd}(t)$  is an indicator function which is 1 between  $c$  and  $d$  and 0 everywhere else
  - Use a combination of step and indicator functions  $= u_{cd}(t) \frac{t-c}{d-c} + u_d(t)$  which is 0 less than  $c$ , increasing between  $c$  and  $d$  and 1 for values greater than  $d$

- We can find the Laplace transforms of these functions

$$\begin{aligned}
- \mathcal{L}\{u_c\} &= \int_0^\infty e^{-st} u_c(t) dt = \int_c^\infty e^{-st} \cdot 1 dt = \frac{e^{-cs}}{s} \\
- \mathcal{L}\{u_{cd}\} &= \mathcal{L}\{u_c - u_d\} = \frac{e^{-cs}}{s} - \frac{e^{-ds}}{s}
\end{aligned}$$

- The function  $g(t)$  from earlier can be written as  $u_{5,10}(t) \cdot \frac{t-5}{5} + u_{10}(t) \cdot 1$

$$\begin{aligned}
- \text{Rearrange this as } \frac{1}{5} [u_{5,10}(t)(t-5) + 5u_{10}] &= \frac{1}{5} [(u_5(t) - u_{10}(t))(t-5) + 5u_{10}(t)] = \\
&= \frac{1}{5} [u_5(t)(t-5) - u_{10}(t)(t-10)] \\
- h(t) &= u_5(t)(t-5) \text{ is the time shift of } f \text{ by } t=5
\end{aligned}$$

- Theorem (Laplace transform of time-shift): If  $F = \mathcal{L}\{f\}$  exists for  $s > a$  and  $c \geq 0$ , then  $\mathcal{L}\{u_c(t)f(t-c)\} = \int_0^\infty e^{-st} u_c(t)f(t-c) dt = \int_c^\infty e^{-st} f(t-c) dt = \int_0^\infty e^{-s(u+c)} f(u) du = e^{-sc} \int_0^\infty e^{-su} f(u) du = e^{-sc} \mathcal{L}\{f\}$

- Consider the IVP  $y(0) = 0$ ,  $y'(0) = 0$  and  $y'' + 4y = u_1(t)$

$$\begin{aligned}
- \mathcal{L}\{y'' + 4y\} &= \mathcal{L}\{u_1(t)\} \rightarrow s^2 Y(s) - sy(0) + 4Y(s) = \frac{e^{-s \cdot 1}}{s} \\
- (s^2 + 4)Y(s) &= \frac{e^{-s}}{s} \rightarrow Y(s) = \frac{e^{-s}}{s(s^2 + 4)} \\
- Y(s) &= e^{-s} \left( \frac{1}{4s} + \frac{-\frac{1}{4}s}{s^2 + 4} \right) = \frac{1}{4} e^{-s} \left( \frac{1}{s} - \frac{s}{s^2 + 4} \right) \\
- \text{Let } H(s) &= \frac{1}{s} - \frac{s}{s^2 + 4} \text{ so that } h(t) = 1 - \cos(2t). \quad y(t) = \frac{1}{4} u_1(t) h(t-1) \text{ by the previous theorem}
\end{aligned}$$

- Consider  $y'' + 4y = g(t)$  with  $y(0) = 0$ ,  $y'(0) = 0$  and  $g(t) = \begin{cases} 0 & 0 < t < 5 \\ \frac{t-5}{5} & t = 5 \leq t < 10 \\ t & 10 \leq t \end{cases}$  as

before

$$\begin{aligned}
- g(t) &\text{ is equivalent to } \frac{1}{5} (u_5(t)(t-5) - u_{10}(t)(t-10)) \\
- \mathcal{L}\{y'' + 4y\} &= \mathcal{L}\{g(t)\} \rightarrow (s^2 + 4)Y(s) = \frac{1}{5} [e^{-5s} \cdot \frac{1}{s^2} - e^{-10s} \cdot \frac{1}{s^2}] \\
- Y(s) &= \frac{1}{5} \left[ e^{-5s} \frac{1}{s^2(s^2 + 4)} - e^{-10s} \frac{1}{s^2(s^2 + 4)} \right] \text{ so that } y(t) = \frac{1}{5} [u_5(t)h(t-5) - u_{10}(t)h(t-10)] \\
&\text{ where } h(t) = \mathcal{L}^{-1}\left\{\frac{1}{s^2(s^2 + 4)}\right\} = \frac{1}{4}t - \frac{1}{8}\sin(2t)
\end{aligned}$$

## 8 Impulse Functions

- $\delta_\epsilon(t) = \frac{1}{\epsilon} \cdot [u_0(t) - u_\epsilon(t)]$
- Note:  $\int_0^\infty \delta_\epsilon(t) dt = 1$
- Equivalently  $\delta_\epsilon(t) = \begin{cases} \frac{1}{\epsilon} & 0 \leq t < \epsilon \\ 0 & \text{else} \end{cases}$
- Let  $\delta = \lim_{\epsilon \rightarrow 0} \delta_\epsilon$  such that  $\delta(t) = 0$  when  $t \neq 0$

- Then  $\delta(t) = \begin{cases} \infty & t = 0 \\ 0 & t \neq 0 \end{cases}$
- $\int_0^\infty \delta(t) dt = \lim_{\epsilon \rightarrow 0} 1 = 1$
- For any function  $f(t)$  continuous on  $a \leq 0 < b$ , we have  $\int_a^b f(t)\delta(t) dt = \lim_{\epsilon \rightarrow 0} \int_a^b f(t)\delta_\epsilon(t) dt = \lim_{\epsilon \rightarrow 0} \int_0^\epsilon f(t) \cdot \frac{1}{\epsilon} dt = \lim_{\epsilon \rightarrow 0} \frac{1}{\epsilon} \int_0^\epsilon f(t) dt = \lim_{\epsilon \rightarrow 0} \frac{1}{\epsilon} (\epsilon - 0)f(t^*)$  for some point  $t^*$  s.t.  $0 \leq t^* \leq \epsilon$  (MVT for integrals)
- $= \lim_{\epsilon \rightarrow 0} f(t^*) = f(0)$  since  $t^* \rightarrow 0$  as  $\epsilon \rightarrow 0$
- $\delta$  is a generalized function (aka distribution). It is also called the Dirac delta function
- We can also shift the impulse function i.e.  $\delta(t - t_0) = 0$  if  $t \neq t_0$
- If  $f(t)$  is continuous on  $a \leq t_0 < b$ ,  $\int_a^b f(t)\delta(t - t_0) dt = f(t_0)$
- Theorem  $\mathcal{L}\{\delta(t - t_0)\} = \int_0^\infty e^{-st}\delta(t - t_0) dt = f(t_0) = e^{-st_0}$
- Consider an undamped oscillator  $y'' + y = I_0\delta(t)$   $y(0) = 0$ ,  $y'(0) = 0$ 
  - Laplace:  $s^2Y(s) + Y(s) = I_0e^{-s \cdot 0} \implies Y(s) = I_0 \frac{1}{s^2+1}$
  - Inverse LT:  $y(t) = I_0 \sin(t)$
  - However this gives  $y'(0) = I_0$  but since we only consider  $t \geq 0$ ,  $y(t)$  is actually  $u_0(t)I_0 \sin(t)$
  - $\lim_{t \rightarrow 0^-} y'(t) = 0$

## 9 Convolution Integrals and Their Applications

- If  $f$  and  $g$  are piecewise continuous on  $[0, \infty)$ , we define their convolution as  $(f * g)(t) = \int_0^t f(t - \tau)g(\tau) d\tau$
- Example:  $f(t) = t$  and  $g(t) = e^{-2t}$ 
  - $(f * g)(t) = \int_0^t (t - \tau)e^{-2\tau} d\tau = [(t - \tau)(-\frac{1}{2})e^{-2\tau}]_{\tau=0}^{\tau=t} + \int_0^t (-\frac{1}{2})e^{-2\tau} d\tau = \frac{t}{2} + \frac{1}{4}e^{-2t} - \frac{1}{4}$
- Rules for Convolutions
  - $f * g = g * f$
  - $f * (g + h) = f * g + f * h$
  - $f * 0 = 0$
  - $(cf) * g = c(f * g)$
  - Note:  $f * 1 = f$
- Proof that the convolution is commutative:

$$- f * g = \int_0^t f(t - \tau)g(\tau) d\tau = - \int_t^0 f(u)g(t - u) du = \int_0^t g(t - u)f(u) du = g * f$$

- Convolution Theorem: If  $F(s) = \mathcal{L}\{f(t)\}$  and  $G(s) = \mathcal{L}\{g(t)\}$  both exist for  $s > a \geq 0$  then  $F(s)G(s) = \mathcal{L}\{(f * g)(t)\}$  and equivalently  $\mathcal{L}^{-1}\{F(s)G(s)\} = (f * g)(t)$

## 10 Introduction to Partial Differential Equations

- The Heat Equation
  - Consider a metal rod of length  $L$ . Let  $u(x, t)$  be the temperature in the cross section at location  $x$  and time  $t$
  - $u_t(x, t)$  represents the change in temperature over time for a fixed slice
  - $u_x(x, t)$  represents the change over the rod for a fixed time
  - $u_{xx}(x, t)$  represents how  $u_x(t)$  changes over the rod for a fixed time
  - The Heat Equation is given by  $u_t = \alpha^2 u_{xx}$  where  $\alpha^2$  is the thermal diffusivity
- Example: Solving the Heat Equation
  - We can solve this by the separation of variables
  - Consider a metal rod with length 50cm, insulated on the ends and with an initial temp of 20 C throughout with the ends maintained at 0C
  - The homogeneous problem is  $u_t = \alpha^2 u_{xx}$  with  $u(0, T) = 0$ ,  $u(L, t) = 0$  for  $t > 0$  and  $u(x, 0) = 20$
  - Assume that we can write  $u(x, t) = X(x)T(t)$ . This is a strong assumption since many functions cannot be written in this form e.g.  $u(x, t) = \sin(xt)$
  - Using this assumption, we can rewrite the homogeneous problem as  $X(x)T'(t) = \alpha^2 X''(x)T(t)$ ,  $X(0)T(t) = 0$ ,  $X(L)T(t) = 0$  and  $X(x)T(0) = 20$
  - This leads to  $X(0) = 0$  and  $X(L) = 0$  since  $T(t)$  cannot always be 0
  - We rearrange the equation above as  $\frac{T'(t)}{\alpha^2 T(t)} = \frac{X''(x)}{X(x)} = -\lambda$
  - Since the two sides are equal, this expression cannot change over  $x$  or  $t$  so we let the expression be equal to  $-\lambda$  where  $\lambda$  is a constant
  - We can find a solution for  $X''(x)$  since  $\frac{X''(x)}{X(x)} = -\lambda$  which has a solution  $X(x) = c_1 \cos(\sqrt{\lambda}x) + c_2 \sin(\sqrt{\lambda}x)$  and using  $X(0) = 0$  and  $X(L) = 0$ , we get  $c_1 = 0$  and  $0 = c_2 \sin(\sqrt{\lambda} \cdot L)$  so that  $\sqrt{\lambda} \cdot L = n \cdot \pi$ . Equivalently  $\lambda = \frac{n^2 \cdot \pi^2}{L^2}$
  - We can similarly find a solution for  $T(t)$  where the ODE is  $\frac{T'(t)}{\alpha^2 T(t)} = -\lambda$
  - $T'(t) = -\frac{n^2 \pi^2}{L^2} \alpha^2 T(t) \implies T(t) = e^{-\frac{n^2 \pi^2}{L^2} \alpha^2 t}$
  - Combining into solutions of the heat equation:  $u(x, t) = X(x)T(t) = \sin(\frac{n\pi}{L}x)e^{-\frac{n^2 \pi^2}{L^2} \alpha^2 t}$  (there can be any constant factor at the front)
  - Fundamental solutions to the Heat Equation:  $u_n(x, t) = \sin(\frac{n\pi x}{L})e^{-\frac{n^2 \pi^2}{L^2} \alpha^2 t}$

- Theorem (Superposition Principle): If  $a(x, t)$  and  $b(x, t)$  solve the heat equation,  $c_1 a(x, t) + c_2 b(x, t)$  solves the heat equation  $\forall c_1, c_2 \in \mathbb{R}$
- General solution to the heat equation is  $u(x, t) = \sum_{n=1}^{\infty} c_n \sin(\frac{n\pi x}{L}) e^{-\frac{n^2 \pi^2 \alpha^2}{L^2} t}$
- Using the initial conditions, we have  $20 = u(x, 0) = \sum_{n=1}^{\infty} c_n \sin(\frac{n\pi x}{L})$
- This will work for  $c_n = \frac{2}{L} \int_0^L 20 \cdot \sin(\frac{n\pi x}{L})$  (Fourier Series)

## 11 Fourier Series

- We define an inner product on  $PC[a, b]$  (piecewise continuous from  $a$  to  $b$ ) by  $\langle f, g \rangle = \int_a^b f(x)g(x) dx$
- We say  $f$  and  $g$  are orthogonal if  $\langle f, g \rangle = 0$
- Consider  $f(x) = \cos(x)$  and  $g(x) = \sin(x)$  in  $PC[-\pi, \pi]$ .  $\langle \cos x, \sin x \rangle = \int_{-\pi}^{\pi} \cos x \sin x dx = 0 \implies \cos x, \sin x$  are orthogonal on  $PC[-\pi, \pi]$
- Consider  $f(x) = \cos(x)$  and  $g(x) = x^2$  in  $PC[-\pi, \pi]$   $\langle \cos x, x^2 \rangle = \int_{-\pi}^{\pi} \cos x x^2 dx = -4\pi$
- Theorem: The following set of functions is an orthogonal family in  $PC[-L, L]$ :  $\{\frac{1}{2}, \sin(\frac{m\pi x}{L}), \cos(\frac{m\pi x}{L}) : m = 1, 2, 3, \dots\}$
- Theorem (Fourier's Theorem): Suppose  $f$  is periodic with period  $2L$  and both  $f, f'$  belonging to  $PC[-L, L]$ . Then  $f$  has a Fourier Series:  

$$f(x) = \frac{a_0}{2} + \sum_{m=1}^{\infty} (a_m \cos \frac{m\pi x}{L} + b_m \sin \frac{m\pi x}{L}), \text{ where } a_0 = \frac{2}{L} \langle f(x), \frac{1}{2} \rangle = \frac{1}{L} \int_{-L}^L f(x) dx,$$

$$a_m = \frac{1}{L} \langle f(x), \cos \frac{m\pi x}{L} \rangle = \frac{1}{L} \int_{-L}^L f(x) \cos \frac{m\pi x}{L} dx \text{ and } b_m = \frac{1}{L} \langle f(x), \sin \frac{m\pi x}{L} \rangle = \frac{1}{L} \int_{-L}^L f(x) \sin \frac{m\pi x}{L} dx$$
- The Fourier series converges to  $f(x)$  at all points where  $f$  is continuous and to the midway point between left and right limit at all points where  $f$  is discontinuous