

## Report 6: Oblique Shock Wave

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# Contents

<b>1</b>	<b>Problem Statement</b>	<b>1</b>
<b>2</b>	<b>Mathematical Model</b>	<b>1</b>
<b>3</b>	<b>Assumptions</b>	<b>1</b>
<b>4</b>	<b>Analysis</b>	<b>1</b>
4.1	Working Procedure . . . . .	2
4.2	Results . . . . .	3
<b>5</b>	<b>Conclusion</b>	<b>4</b>
<b>A</b>	<b>Matlab Codes</b>	<b>5</b>
	<b>References</b>	<b>8</b>

# Nomenclature

N.S. Navier Stockes | w.r.t. with respect to

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# 1 Problem Statement

Produce charts that describe the change of supersonic flow properties when it turned away from itself.

## 2 Mathematical Model

The most general governing equation is N.S. equation. This is any dummy text just to show the capabilities of nomenclatures of L<sup>A</sup>T<sub>E</sub>X w.r.t. L<sup>A</sup>T<sub>E</sub>X.

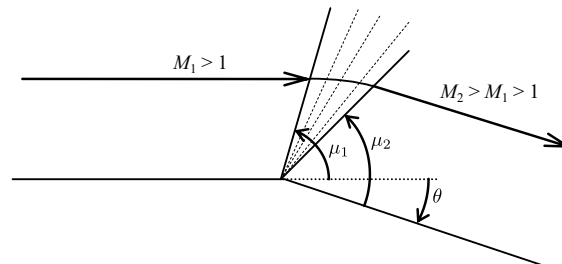
$$\mu = \sin^{-1} \left( \frac{1}{M} \right) \quad (1)$$

$$\left( \frac{a_0}{a} \right)^2 = 1 + \frac{\gamma - 1}{2} M^2 \quad (2)$$

## 3 Assumptions

- |  |  |  |
|--|--|--|
| <ul style="list-style-type: none"><li>1. Steady Flow</li><li>2. Quasi-dimensional flow; (area is variable with <math>x</math> only).</li><li>3. Body forces can be neglected; (weight of fluid).</li><li>4. Viscous stresses are absent.</li></ul> |  | <ul style="list-style-type: none"><li>5. Changes in potential energy are neglected.</li><li>6. Perfect gas.</li><li>7. Thermally perfect gas.</li><li>8. Adiabatic flow with no external work.</li></ul> |
|--|--|--|

## 4 Analysis



We can derive the formula that governs super flow expansion as:

$$\nu(M_1) = \nu(M_2) + \theta \quad (3)$$

$$\nu(M) = \sqrt{\frac{\gamma + 1}{\gamma - 1}} \tan^{-1} \left( \sqrt{\frac{\gamma - 1}{\gamma + 1}} (M^2 - 1) \right) \tan^{-1} \left( \sqrt{M^2 - 1} \right) \quad (4)$$

If you know  $M_1$  &  $\theta$  know you can use equation (3) to solve for  $M_2$  using Newton Raphson iteration scheme as described below:

$$M_2 = M_2 - \frac{f(M_2)}{f'(M_2)} \quad (5)$$

$$f(M_2) = \sqrt{\frac{\gamma+1}{\gamma-1}} \tan^{-1} \left( \sqrt{\frac{\gamma-1}{\gamma+1}} (M^2 - 1) \right) - \tan^{-1} \left( \sqrt{M^2 - 1} \right) - \theta - \nu(M_1) \quad (6)$$

$$f'(M_2) = \frac{M_2}{\sqrt{M_2^2 - 1} \left( 1 + \frac{\gamma-1}{\gamma+1} (M_2^2 - 1) \right)} - \frac{1}{M_2 \sqrt{M_2^2 - 1}} \quad (7)$$

After you get  $M_2$  you can get the pressure and temperature using the isentropic relations:

$$\frac{T_2}{T_1} = \frac{1 + \frac{\gamma-1}{2} M_1^2}{1 + \frac{\gamma-1}{2} M_2^2} \quad (8)$$

$$\frac{P_2}{P_1} = \left( \frac{T_2}{T_1} \right)^{\frac{\gamma}{\gamma-1}} \quad (9)$$

$$\frac{\rho_2}{\rho_1} = \left( \frac{T_2}{T_1} \right)^{\frac{1}{\gamma-1}} \quad (10)$$

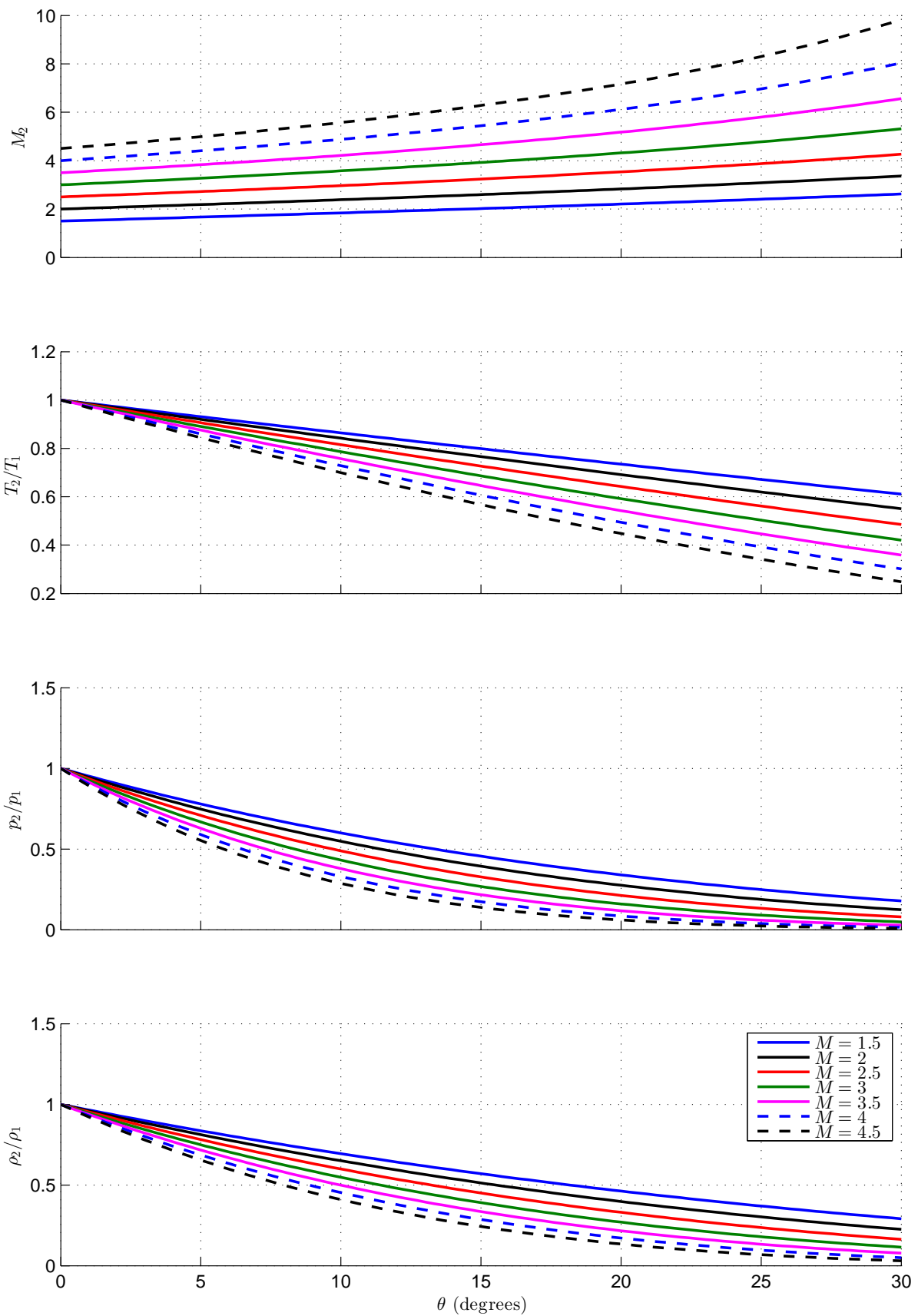
$$P_{01} = P_{02} = P_0 \quad (11)$$

function `M_2o.m` [code 2] illustrates this procedure. `Report6_main.m` [code 1] was used to plot the figures shown in section 4.2.

This is a line with some Arabic words. سطر انجليزي.  
Thus is some English words in an Arabic. الكلمات الانجليزية.  
line.

و بناءً عليه فقد تم اثبات المطلوب. و بناءً عليه فقد  
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## 4.2 Results



The following table is just for demonstration. It doesn't provide any useful information.

Angle( $\theta$ )	Temperature	Pressure	Density
0	123	111	444
20	500	222	444
40	640	222	444
90	200	222	444

## 5 Conclusion

We see that we can still obtain solutions for  $M_2$  for  $\theta > 90^\circ$ . But, however I think the solutions for  $\theta > 90^\circ$  aren't practical.

Pressure, temperature & density increases as the kinetic energy increase as in [1].



# Appendices

## A Matlab Codes

Code 1: Report6\_main.m

```
1 clc
2 close all
3 clearvars
4
5 set(groot, 'DefaultAxesColorOrder', [0,0,1;0,0,0;1,0,0;0,0.5,0;1,0,1])
6 set(groot, 'DefaultAxesLineStyleOrder', '-|--|-.')
7 set(groot, 'DefaultLineLineWidth', 1.5);
8
9 gamma=1.4;
10 theta_d_vec=linspace(0,30);
11
12 M_1_vec=1.5:.5:4.5;
13
14 yAxesTitles_tex=["M_{2}", "T_{2}/T_{1}", "p_{2}/p_{1}", "\rho_{2}/\rho_{1}"]
15 ];
16 for ii=1:length(yAxesTitles_tex)
17     subplot(length(yAxesTitles_tex),1,ii);
18     hold on;
19     grid;
20     if ii~=length(yAxesTitles_tex), set(gca, 'XTickLabel', []);end
21     ylabel('$'+yAxesTitles_tex(ii)+'$', 'interpreter', 'latex');
22 end
23 xlabel('$\theta_{\text{d}}$ (degrees)', 'interpreter', 'latex'); %This applies to the
24     bottom subplot
25 M_1_legend_tex=strings(length(M_1_vec),1);
26 ii=1;
27 for M_1=M_1_vec
28     M_2=M_2o(M_1,theta_d_vec,gamma);
29     subplot(4,1,1);plot(theta_d_vec,M_2);
30
31     T_2_T_1=(1+(gamma-1)/2*M_1*M_1)./(1+(gamma-1)/2*M_2.^2);
32     subplot(4,1,2);plot(theta_d_vec,T_2_T_1);
33
34     p_2_p_1=T_2_T_1.^(gamma/(gamma-1));
35     subplot(4,1,3);plot(theta_d_vec,p_2_p_1);
36
37     rho_2_rho_1=T_2_T_1.^(1/(gamma-1));
```

```

38 subplot(4,1,4);plot(theta_d_vec,rho_2_rho_1);
39
40 M_1_legend_tex(ii)="$M="+M_1+'$';
41
42 ii=ii+1;
43 end
44
45 legend(M_1_legend_tex,'interpreter','latex');
46
47 set(groot,'DefaultAxesColorOrder','remove')
48 set(groot,'DefaultAxesLineStyleOrder','remove')
49 set(groot,'DefaultLineLineWidth','remove');
50
51 export_figure(gcf,'||',"figures",[],"pdf","emf")

```

#### Code 2: function M\_2o.m

```

1 function M_2_vec=M_2o(M_1,theta_d_vec, ...
2     gamma) %optional arguments
3 if nargin<3
4     gamma=1.4;
5 end
6
7 theta_vec=deg2rad(theta_d_vec);
8 n1=sqrt((gamma+1)/(gamma-1))*atan(sqrt((gamma-1)/(gamma+1)*(M_1^2-1)))-
9     atan(sqrt(M_1^2-1));
10 %Newton Raphson iteration
11 M_2_vec=1.1*ones(size(theta_d_vec));
12 for ii=1:length(theta_d_vec)
13     f=sqrt((gamma+1)/(gamma-1))*atan(sqrt((gamma-1)/(gamma+1)*(M_2_vec(ii)
14         ^2-1)))-atan(sqrt(M_2_vec(ii)^2-1))-theta_vec(ii)-n1;
15     fdash=1/(M_2_vec(ii)^2-1)^(1/2)*M_2_vec(ii)/(1+(gamma-1)/(gamma+1)*
16         (M_2_vec(ii)^2-1))-1/(M_2_vec(ii)^2-1)^(1/2)/M_2_vec(ii);
17     M_2_n=M_2_vec(ii)-f/fdash;
18     while abs(M_2_vec(ii)-M_2_n)>=100*eps %This is dangerous. Infinte loop
19         can occur!!
20         M_2_vec(ii)=M_2_n;
21         f=sqrt((gamma+1)/(gamma-1))*atan(sqrt((gamma-1)/(gamma+1)*(M_2_vec
22             (ii)^2-1)))-atan(sqrt(M_2_vec(ii)^2-1))-theta_vec(ii)-n1;
23         fdash=1/(M_2_vec(ii)^2-1)^(1/2)*M_2_vec(ii)/(1+(gamma-1)/(gamma+1)
24             *(M_2_vec(ii)^2-1))-1/(M_2_vec(ii)^2-1)^(1/2)/M_2_vec(ii);
25         M_2_n=M_2_vec(ii)-f./fdash;
26     end
27     M_2_vec(ii)=M_2_n;

```

24 | [end](#)

---

## References

- [1] J. D. Anderson, Modern Compressible Flow, McGraw-Hill, New York, 1990.
- [2] Report (1).
- [3] Report (3).