

# Nuclear Magnetic Resonance (NMR)

A Brief Introduction and Applications

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1. Principles of NMR
2. Apparatus and Experimental Procedures

# Setup

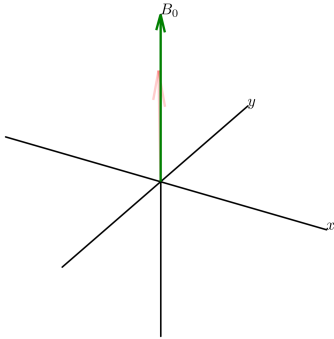


Figure: Typical  $\vec{B}_0$  and vector aligned with field

- Spin  $1/2$  nuclei
- Strong magnetic field, denoted  $\vec{B}_0$

## Energy and Hamiltonian

Magnetic moment of nucleus

$$\vec{\mu} = \gamma \vec{S} \quad (1)$$

Yields energy and Hamiltonian of

$$E = -\vec{\mu} \cdot \vec{B}_0 \quad (2)$$

$$H = -\gamma \vec{B}_0 \cdot \vec{S} \quad (3)$$

# Larmor Precession

Evolving an arbitrary spin vector

$$|\psi\rangle = \cos(\theta/2) |+\rangle + e^{i\phi} \sin(\theta/2) |-\rangle$$

yields, [1]

$$\langle S_x \rangle = \frac{\hbar}{2} \sin(\theta) \cos(\gamma B_0 t)$$

$$\langle S_y \rangle = \frac{\hbar}{2} \sin(\theta) \sin(\gamma B_0 t)$$

$$\langle S_z \rangle = \frac{\hbar}{2} \cos(\theta)$$

with the characteristic **Larmor frequency**

$$\omega_0 = \gamma B_0 \quad (4)$$

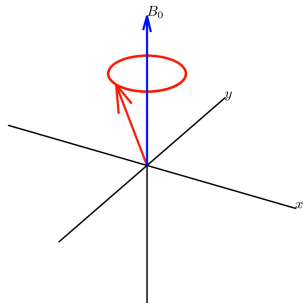


Figure: Larmor precession about  $\vec{B}_0$

# Net Magnetization and Relaxation

With many spins, any material can acquire a net magnetization.  
Along  $\hat{z}$ , statistical mechanics tells us that

$$\frac{N_+}{N_-} = \exp\left(\frac{\gamma\hbar B_0}{kT}\right) \quad (5)$$

with total **net magnetization**

$$M_z = \sum_i \gamma\hbar m_i = \frac{\hbar}{2}\gamma(N_+ - N_-) \quad (6)$$

Expect spins to eventually re-align with  $\hat{z}$  direction.

# Producing Transverse Magnetization

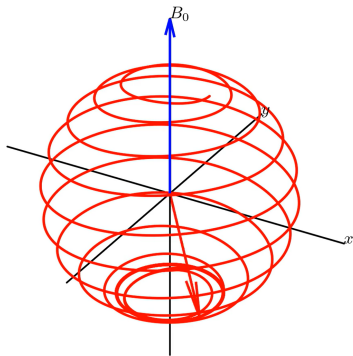
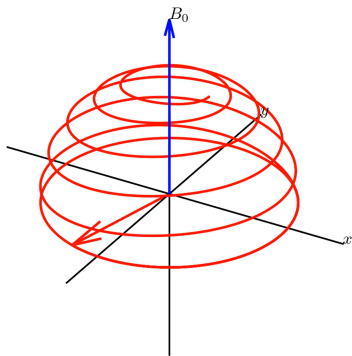
## Circularly Polarized Field

$$\vec{B}_1(t) = B_1 [\cos(\omega_0 t)\hat{x} + \sin(\omega_0 t)\hat{y}] \quad (7)$$

To “tip” magnetization into  $xy$  plane, apply a circularly polarized  $\vec{B}_1(t)$ . Classical torque result in a rotating frame yields,

$$\vec{B}^* = \left( \left[ \vec{B}_0 - \frac{\omega_0}{\gamma} \right] \hat{z}^* + B_1 \hat{x}^* \right) \quad (8)$$

## $\frac{\pi}{2}$ and $\pi$ pulses



(a)  $\frac{\pi}{2}$  pulse produces transverse magnetization (b)  $\pi$  pulse flips magnetization from  $+\hat{z}$  to  $-\hat{z}$

Figure:  $\langle \vec{S} \rangle$  over time during a  $\frac{\pi}{2}$  and  $\pi$  pulse

# Spin-Lattice Relaxation ( $T_1$ )

Time constant associated with decay of net magnetization back towards  $+\hat{z}$ .

$$\frac{d\vec{M}(t)}{dt} = \frac{\vec{M}(t) - \vec{M}_0}{T_1} \quad (9)$$

Referred to as the **Bloch equations**.



## Spin-Spin Relaxation ( $T_2$ )

Time constant associated with the decay of the relative “coherence” of spins in  $xy$  plane, with associated Bloch equation [2],

$$\frac{d\vec{M}_{x,y}}{dt} = -\frac{\vec{M}_{x,y}}{T_2} \quad (10)$$

# Free Induction Decay and $T_2^*$

- Observations performed with a pickup coil detecting flux in transverse direction
- Magnetization becomes increasingly decoherent; produces **free induction decay (FID)** signal
- Combined effects of spin-lattice, spin-spin, and field inhomogeneities [2]

$$\frac{1}{T_2^*} = \frac{1}{T_1} + \frac{1}{T_2} + \gamma \Delta B_0 \quad (11)$$

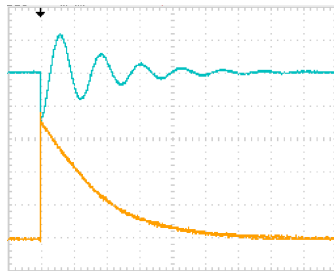





Figure: Basic FID signal of mineral oil [3]

# References

-  D. J. Griffiths and D. F. Schroeter, *Introduction to Quantum Mechanics*.  
Cambridge ; New York, NY: Cambridge University Press, 3 ed., 2018.
-  C. P. Slichter, *Principles of Magnetic Resonance*.  
USA: Springer Science, 3 ed., 1990.
-  *Pulsed Nuclear Magnetic Resonance*.