Nuclear Magnetic Resonance (NMR)

A Brief Introduction and Applications

Ali Ahmed Neil Mandar Zain Kamal Department of Physcis & Astronomy, Rutgers University

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Overview

1. Principles of NMR

Setup

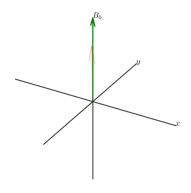


Figure: Typical \vec{B}_0 and vector aligned with field

- Spin 1/2 nuclei
- Strong magnetic field, denoted \vec{B}_0

Energy and Hamiltonian

Magnetic moment of nucleus

$$\vec{\mu} = \gamma \vec{S} \tag{1}$$

Yields energy and Hamiltonian of

$$E = -\vec{\mu} \cdot \vec{B_0} \tag{2}$$

$$H = -\gamma \vec{B_0} \cdot S \tag{3}$$

Larmor Precession

Evolving an arbitrary spin vector

$$\ket{\psi} = \cos\left(heta/2
ight)\ket{+} + e^{i\phi}\sin\left(heta/2
ight)\ket{-}$$

yields, [1]

$$\langle S_x \rangle = rac{\hbar}{2} \sin(\theta) \cos(\gamma B_0 t)$$

 $\langle S_y \rangle = rac{\hbar}{2} \sin(\theta) \sin(\gamma B_0 t)$
 $\langle S_z \rangle = rac{\hbar}{2} \cos(\theta)$

with the characteristic Larmor frequency

$$\omega_0 = \gamma B_0 \tag{4}$$

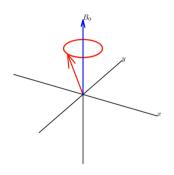


Figure: Larmor precession about \vec{B}_0

Net Magnetization and Relaxation

With many spins, any material can acquire a net magentization. Along \hat{z} , statistical mechanics tells us that

$$\frac{N_{+}}{N_{-}} = \exp\left(\frac{\gamma \hbar B_{0}}{kT}\right) \tag{5}$$

with total net magnetization

$$M_z = \sum_i \gamma \hbar m_i = \frac{\hbar}{2} \gamma (N_+ - N_-)$$
 (6)

Expect spins to eventually re-align with \hat{z} direction.

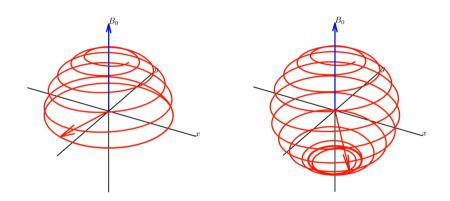
Producing Transverse Magnetization

Circularly Polarized Field

$$\vec{B}_1(t) = B_1 \left[\cos(\omega_0 t) \hat{x} + \sin(\omega_0 t) \hat{y} \right] \tag{7}$$

To "tip" magnetization into xy plane, apply a circularly polarized $\vec{B}_1(t)$. Classical torque result in a rotating frame yields,

$$\vec{B}^* = \left(\left[\vec{B}_0 - \frac{\omega_0}{\gamma} \right] \hat{z}^* + B_1 \hat{x}^* \right) \tag{8}$$



(a) $\frac{\pi}{2}$ pulse produces transverse magentization (b) π pulse flips magentization from $+\hat{z}$ to $-\hat{z}$

Figure: $\langle \vec{\mathcal{S}} \rangle$ over time during a $\frac{\pi}{2}$ and π pulse

Spin-Lattice Relaxation (T_1)

Time constant associated with decay of net magnetization back towards $+\hat{z}$.

$$\frac{d\vec{M}(t)}{dt} = \frac{\vec{M}(t) - \vec{M_0}}{T_1} \tag{9}$$

Referred to as the **Bloch equations**.

Spin-Spin Relaxation (T_2)

Time constant associated with the decay of the relative "coherence" of spins in xy plane, with associated Bloch equation,

$$\frac{d\vec{M}_{x,y}}{dt} = \frac{\vec{M}_{x,y}}{T_2} \tag{10}$$

Free Induction Decay and T_2^*

- Observations performed with a pickup coil detecting flux in transverse direction
- Magnetization becomes increasingly decoherent; produces free induction decay (FID) signal
- Combined effects of spin-lattice, spin-spin, and field inhomogeneities [2]

$$\frac{1}{T_2^*} = \frac{1}{T_1} + \frac{1}{T_2} + \gamma \Delta B_0 \tag{11}$$

References



D. J. Griffiths and D. F. Schroeter, *Introduction to Quantum Mechanics*. Cambridge; New York, NY: Cambridge University Press, 3 ed., 2018.



C. P. Slichter, *Principles of Magnetic Resonance*.

USA: Springer Science, 3 ed., 1990.