

Background: Asymptotic Freeness and Large Deviations

We start by recalling the basic asymptotic freeness result.

Proposition: Let (A_1^N, \dots, A_r^N) be an independent r -tuple of GUE $N \times N$ matrices. Then, the family A_1^N, \dots, A_r^N converge in distribution to r independent semicircular elements, $s_1, \dots, s_r \in B(\mathcal{F}(\mathbb{C}^r))$, in the sense that for all $m \geq 1$ and all $1 \leq i_1, \dots, i_m \leq r$, we have

$$\lim_{N \rightarrow \infty} E[\text{tr}(A_{i_1}^N \cdots A_{i_m}^N)] = \varphi(s_{i_1} \cdots s_{i_m}),$$

where φ is the vacuum state, $\varphi(T) = \langle T\Omega, \Omega \rangle$.

In fact, this collection is *almost surely* asymptotically free, in the following sense. Suppose we have two random matrices A^N and B^N defined on probability spaces (X_N, μ_N) . Define

$$X := \prod_{N \in \mathbb{N}} X_N$$

$$\mu := \prod_{N \in \mathbb{N}} \mu_N,$$

where the latter is the product measure on X . The matrices A^N and B^N are said to be almost surely asymptotically free if there exists a noncommutative probability space (A, φ) and $a, b \in A$, and for almost all $x = (x_N)_N \in X$, we have $A^N(x_N), B^N(x_N) \in (\mathbb{M}_N, \text{tr})$ converge in distribution to a, b .

Now, from here, we may ask a seemingly simple question: as N grows large, how likely are we to encounter other distributions? To make this sense more precise, we consider a random $N \times N$ self-adjoint matrix A , and let

$$\mu_A = \frac{1}{N} \sum_{i=1}^N \delta_{\lambda_i}$$

be its empirical spectral distribution. This is a random probability measure on \mathbb{R} , and as $N \rightarrow \infty$, the semicircle law gives that μ_A converges weakly to the semicircle distribution; this can be strengthened to almost sure convergence by an application of the argument for asymptotic freeness. The question then becomes, how quickly does the deviation between μ_A and any other probability distribution ν decrease as N increases? This is where the theory of large deviations starts to take shape.

Much of this exposition related to the classical notions of entropy will be centered around results discussed in [MS17, Ch. 7].

Large Deviations for Random Variables

We start with one of the classical examples of convergence of random variables to introduce large deviations. Consider a sequence of independent and identically distributed real-valued random variables $(X_i)_i$ with common distribution μ . Set

$$S_n = \frac{1}{n} \sum_{i=1}^n X_i.$$

$$m = E[X_1]$$

$$v = E[X_1^2] - m^2.$$

Then, we have that if $E[X_1^2] < \infty$, the central limit theorem says that $S_n \approx m + \frac{\sigma}{\sqrt{n}}N(0, 1)$.

If μ is the standard Gaussian distribution, then this gives that S_n is distributed as $N(0, 1/n)$; we then get that

$$P(S_n \in [x, x + dx]) \approx \sqrt{\frac{n}{2\pi}} e^{-nx^2} dx.$$

Asymptotically, this gives that the probability that S_n is near the value $x \in \mathbb{R}$ decays exponentially in n determined by a rate function $I(x) = x^2/2$.

We will now generalize this result. In particular, if we let μ be any distribution discussed above (rather than simply the normal distribution), then we will find a rate function $I(x)$ such that

$$e^{-nI(x)} \sim P(S_n > x)$$

whenever $x > m$, and whenever $x < m$

$$e^{-nI(x)} \sim P(S_n < x).$$

For a given distribution μ , we can compute the rate function by using a family of basic manipulations. If $x > m$, then for all $\lambda \geq 0$, we may use Markov's inequality to obtain

$$\begin{aligned} P(S_n > x) &= P(nS_n > nx) \\ &= P\left(e^{\lambda(nS_n - nx)} \geq 1\right) \\ &\leq E\left[e^{\lambda(nS_n - nx)}\right] \\ &= e^{-\lambda nx} E\left[e^{\lambda(X_1 + \dots + X_n)}\right] \\ &= (e^{-\lambda x} E[e^{\lambda X}])^n, \end{aligned}$$

where X is identically distributed to each of the X_i , and we use the fact that the X_i are independent. We may then define

$$\Lambda(\lambda) = \ln E[e^{\lambda X}] \quad (*)$$

to be an extended real-valued function, but we only consider μ for which $\Lambda(\lambda)$ is finite for all λ in an open neighborhood of 0. The equation $(*)$ is known as the cumulant generating function for μ .

This gives the inequality

$$P(S_n > x) \leq e^{-n(\lambda x - \Lambda(\lambda))}.$$

Since \ln is a concave function, Jensen's inequality gives

$$\begin{aligned} \Lambda(\lambda) &\geq E[\ln(e^{\lambda X})] \\ &= E[\lambda X] \\ &= \lambda m. \end{aligned}$$

In particular, for any $\lambda < 0$ and $x > m$, we have $-n(\lambda x - \Lambda(\lambda)) \geq 0$, meaning this equation is valid for all λ . In particular, we have

$$P(S_n > x) \leq \inf_{\lambda \in \mathbb{R}} e^{-n\lambda x - \Lambda(\lambda)}.$$

Now, we observe that Λ is convex. This follows from Hölder's inequality

$$E\left[e^{(1-t)\lambda_1 x + t\lambda_2 x}\right] \leq E[e^{\lambda_1 x}]^{1-t} E[e^{\lambda_2 x}]^t$$

so that

$$\Lambda((1-t)\lambda_1 + t\lambda_2) \leq (1-t)\Lambda(\lambda_1) + t\Lambda(\lambda_2).$$

Defining the *Legendre transform* of Λ by

$$\Lambda^*(x) = \sup_{\lambda \in \mathbb{R}} (\lambda x - \Lambda(\lambda)),$$

we find that this is a convex function of x , as it is a supremum of a family of convex functions of x . Now, since $\Lambda(0) = 0$, it follows that $\Lambda^*(x) \geq 0$, and has $\Lambda^*(m) = 0$. In particular, this gives

$$P(S_n > x) \leq e^{-n\Lambda^*(x)}$$

whenever $x > m$.

It can also be shown that $e^{-n\Lambda^*(x)}$ is an asymptotic lower bound, in that

$$\liminf_{n \rightarrow \infty} \frac{1}{n} \ln P(x - \delta < S_n < x + \delta) \geq -\Lambda^*(x)$$

for all $x \in \mathbb{R}$ and all $\delta > 0$. The method for doing so is outlined in [MS17, Ch. 7, Section 2], and results in Cramér's theorem for real-valued random vectors.

Theorem (Cramér's Theorem): Let X_1, X_2, \dots be a sequence of independent and identically distributed random vectors in \mathbb{R}^d with common distribution μ . Define

$$\begin{aligned} \Lambda(\lambda) &= \ln E \left[e^{\langle \lambda, X_i \rangle} \right] \\ \Lambda^*(x) &= \sup_{\lambda \in \mathbb{R}^d} (\langle \lambda, x \rangle - \Lambda(\lambda)), \end{aligned}$$

and assume that $\Lambda(\lambda) < \infty$ for all λ . Set $S_n = \frac{1}{n}(X_1 + \dots + X_n)$. Then, the distribution μ_{S_n} satisfies has

- $x \mapsto \Lambda^*(x)$ is convex;
- $\{x \in \mathbb{R}^d \mid \Lambda^*(x) \leq \alpha\}$ is compact for all $\alpha \in \mathbb{R}$;
- for any closed $F \subseteq \mathbb{R}^d$,

$$\limsup_{n \rightarrow \infty} \frac{1}{n} \ln P(S_n \in F) \leq - \inf_{x \in F} \Lambda^*(x),$$

- and for any open $G \subseteq \mathbb{R}^d$,

$$\liminf_{n \rightarrow \infty} \frac{1}{n} \ln P(S_n \in G) \geq - \inf_{x \in G} \Lambda^*(x).$$

This theorem defines precisely the large deviation principle that the partial sums satisfy — namely, it is the Legendre transform of the cumulant-generating function.

Large Deviations for the Empirical Distribution

Now, our next step is to develop an analogous large deviation principle for the empirical distribution of the random variables X_1, X_2, \dots . This will give us the idea of classical entropy.

We start by considering the case of (independent and identically distributed) random variables $X_i: \Omega \rightarrow A$ taking values in a finite set $\{a_1, \dots, a_d\}$, and define $p_k := P(X_i = a_k)$. We expect that, as $n \rightarrow \infty$, the empirical distribution of the X_i should converge to the probability measure (p_1, \dots, p_d) on A .

Toward this end, let $Y_i: \Omega \rightarrow \mathbb{R}^d$ be defined by

$$Y_i := (\chi_{\{a_1\}}(X_i), \dots, \chi_{\{a_d\}}(X_i)).$$

We observe that p_k is the probability that Y_i will have 1 in position k and 0 elsewhere, and that $\frac{1}{n}(Y_1 + \dots + Y_n)$ gives the relative frequency of a_1, \dots, a_d — i.e., this has the same information as the empirical distribution of X_1, \dots, X_n .

Any probability measure on A is a d -tuple of positive real numbers satisfying $q_1 + \dots + q_d = 1$. By Cramér's theorem and our discussion above, we have

$$P\left(\frac{1}{n}(\delta_{X_1} + \dots + \delta_{X_n}) \approx (q_1, \dots, q_d)\right) = P\left(\frac{1}{n}(Y_1 + \dots + Y_n) \approx (q_1, \dots, q_d)\right) \\ \sim e^{-n\Lambda^*(q_1, \dots, q_d)}.$$

Applying our definitions for Λ and Λ^* , we have

$$\Lambda(\lambda_1, \dots, \lambda_d) = \ln(p_1 e^{\lambda_1} + \dots + p_d e^{\lambda_d}) \\ \Lambda^*(q_1, \dots, q_d) = \sup_{(\lambda_1, \dots, \lambda_d)} (\lambda_1 q_1 + \dots + \lambda_d q_d - \ln(p_1 e^{\lambda_1} + \dots + p_d e^{\lambda_d})).$$

To compute the supremum over all tuples, we find that the partial derivatives with respect to each λ_i are given by

$$q_i - \frac{1}{p_1 e^{\lambda_1} + \dots + p_d e^{\lambda_d}} p_i e^{\lambda_i},$$

so by concavity, we get that the maximum value occurs when

$$\lambda_i = \ln\left(\frac{q_i}{p_i}\right) + \Lambda(\lambda_1, \dots, \lambda_d).$$

We thus get

$$\Lambda^*(q_1, \dots, q_d) = \sum_{i=1}^d q_i \ln\left(\frac{q_i}{p_i}\right).$$

The quantity on the right is the relative Shannon entropy $H((q_1, \dots, q_d)|(p_1, \dots, p_d))$, which is strictly positive except when $q_1 = p_1, \dots, q_d = p_d$.

In particular, we have that (p_1, \dots, p_d) admits a large deviation principle with rate function given by the relative Shannon entropy. That this holds for any distribution is known as Sanov's theorem.

Theorem (Sanov's Theorem): Let X_1, X_2, \dots be a sequence of independent and identically distributed real random variables with common distribution μ , and let

$$\nu_n := \frac{1}{n} \sum_{i=1}^n \delta_{X_i}$$

be the empirical distribution. Then, the family $\{\nu_n\}_{n \geq 1}$ satisfies a large deviation principle given by the rate function

$$I(\nu) = \begin{cases} \int p \ln(p) d\mu & d\nu = p d\mu \\ +\infty & \text{else.} \end{cases}$$

One-Dimensional Free Entropy: A Heuristic Approach

The next logical step after defining a large deviation principle for a sequence of random variables is to define such a quantity for single free random variables.

In [Voi93], Voiculescu used a heuristic method based on random matrix theory to establish the value $\Sigma(x)$ for some random variable x in a C^* -probability space (\mathcal{A}, φ) distributed according to a compactly supported measure ν on \mathbb{R} (i.e., its spectral measure with respect to the state φ).

It can be established (as in [MS17, Exercise 1.8]) that if X is a $N \times N$ GUE matrix, then the density inside the real vector space $\mathbb{M}_N(\mathbb{C})_{\text{s.a.}}$ is given by

$$dP_N(X) = K e^{-N^2/2 \text{tr}(X^2)} dm,$$

where m is the Lebesgue measure on \mathbb{R}^{N^2} and tr denotes the normalized trace. It can be shown, as in [AGZ10, Theorem 2.5.2], that the joint distribution of the eigenvalues $\lambda_1(X) \leq \dots \leq \lambda_N(X)$ is absolutely continuous with respect to the Lebesgue measure and has density given by

$$dQ_N(\lambda_1, \dots, \lambda_N) = \frac{N^{N^2/2}}{(2\pi)^{N/2} \prod_{j=1}^N j!} e^{-(N \sum_{i=1}^N \lambda_i^2)/2} \prod_{i < j} (\lambda_i - \lambda_j)^2 \prod_{i=1}^N d\lambda_i$$

Heuristically, letting μ_A denote the empirical spectral distribution on $\mathbb{M}_N(\mathbb{C})_{\text{s.a.}}$ given by

$$\mu_A = \frac{1}{N} (\delta_{\lambda_1(A)} + \dots + \delta_{\lambda_N(A)}),$$

we may consider a large deviation principle with $P_N(\mu_A \approx \nu) \sim e^{-N^2 I(\nu)}$, and write it out as

$$\begin{aligned} P_N(\mu_A \approx \nu) &= Q_N \left(\frac{1}{N} (\delta_{\lambda_1} + \dots + \delta_{\lambda_N}) \right) \\ &= \frac{N^{N^2/2}}{(2\pi)^{N/2} \prod_{j=1}^N j!} \int_E e^{-(N \sum_{i=1}^N \lambda_i^2)/2} \prod_{i < j} (\lambda_i - \lambda_j)^2 \prod_{i=1}^N d\lambda_i, \end{aligned}$$

where E is the given set in the Q_N . Whenever $\frac{1}{N} (\delta_{\lambda_1(A)} + \dots + \delta_{\lambda_N(A)}) \approx \nu$, we have that

$$-\frac{N}{2} \sum_{i=1}^N \lambda_i^2 = -\frac{N^2}{2} \left(\frac{1}{N} \sum_{i=1}^N \lambda_i^2 \right)$$

is a Riemann sum for the integral of t^2 with respect to ν . Furthermore, we have

$$\prod_{i < j} (\lambda_i - \lambda_j)^2 = e^{\sum_{i \neq j} \ln |\lambda_i - \lambda_j|}$$

The sum inside the argument of the exponential is a Riemann sum for the quantity $N^2 \iint \ln |s - t| d\nu(s) d\nu(t)$. All in all, we get the following large deviation principle, which was proven rigorously in [BAG97] for the general case of Gaussian random matrices with Dyson index β .

Theorem: Let

$$I(\nu) = - \iint \ln |s - t| d\nu(s) d\nu(t) + \frac{1}{2} \int t^2 d\nu(t) - \frac{3}{4}.$$

Then, the following hold.

- (i) The function $I: \mathcal{M} \rightarrow [0, \infty]$ is a well-defined convex function with compact level sets on the space of probability measures on \mathbb{R} that has minimum value 0 at Wigner's semicircle law.
- (ii) The empirical spectral distribution satisfies a large deviation principle with respect to Q_N with rate function I . That is,

$$\liminf_{N \rightarrow \infty} \frac{1}{N^2} \ln Q_N \left(\frac{1}{N} (\delta_{\lambda_1} + \dots + \delta_{\lambda_N}) \in G \right) \geq - \inf_{\nu \in G} I(\nu)$$

for any open $G \subseteq \mathcal{M}$, and

$$\limsup_{N \rightarrow \infty} \frac{1}{N^2} \ln Q_N \left(\frac{1}{N} (\delta_{\lambda_1} + \dots + \delta_{\lambda_N}) \in F \right) \leq - \inf_{\nu \in F} I(\nu)$$

for any closed $F \subseteq \mathcal{M}$.

Microstates Free Entropy

Voiculescu's original definition of free entropy in the case with more than one free random variable can be viewed as a generalization of the microstates approach to entropy for a classical discrete random variable, discussed in the survey [Voi01]. For the sections following, I will assume a certain level of proficiency with the theory of von Neumann algebras, as this is the setting where Voiculescu developed the theories around free entropy in [Voi94].

Consider a discrete random variable with output values in $\{1, \dots, n\}$ assigned with probabilities p_1, \dots, p_n . The microstates of this discrete random variable for a fixed N are then the set

$$\{f \mid f: \{1, \dots, N\} \rightarrow \{1, \dots, n\}\},$$

and for a fixed ε , the microstates that approximate this distribution are those f with

$$\left| \frac{|f^{-1}(\{j\})|}{N} - p_j \right| < \varepsilon,$$

where $|\cdot|$ denotes the (necessarily finite) cardinality. We denote the collection of such f by $\Gamma(p_1, \dots, p_n; \varepsilon, N)$. One may find the Shannon entropy by evaluating the limit

$$\lim_{\varepsilon \rightarrow 0} \lim_{N \rightarrow \infty} \frac{1}{N} \ln |\Gamma(p_1, \dots, p_n; \varepsilon, N)|$$

Understanding the Microstates Free Entropy

Instead of a classical random variable, we will let (M, τ) be a tracial von Neumann algebra,¹ and let (x_1, \dots, x_n) be a tuple of self-adjoint free random variables in M .

For a fixed R , we let the set of microstates $\Gamma_R(X_1, \dots, X_n; m, k, \varepsilon)$ admit three degrees of approximation: m denotes the level of approximation of mixed moments, k denotes the size of the (self-adjoint) approximation matrices that have operator norm at most R , and ε denotes the closeness of the approximation. Put into symbols, we select all n -tuples $(A_1, \dots, A_n) \in (\mathbb{M}_k(\mathbb{C})_{\text{s.a.}})^n$ with each $\|A_j\| < R$ satisfying

$$|\tau(x_{i_1} \cdots x_{i_p}) - \text{tr}_k(A_{i_1} \cdots A_{i_p})| < \varepsilon$$

for all $1 \leq p \leq m$ and all multi-indices $\mathbf{i}: \{1, \dots, p\} \rightarrow \{1, \dots, n\}$. Here, tr_k is the normalized trace.

As in the case of entropy for classical random variables, free entropy emerges from a certain large deviation principle. Specifically, recall that the distribution of a family of noncommutative random variables is given by the collection of mixed moments with respect to the trace,

$$\Delta(x_1, \dots, x_n) = \{\tau(x_{i_1} \cdots x_{i_p}) \mid p \in \mathbb{N}, \mathbf{i}: \{1, \dots, p\} \rightarrow \{1, \dots, n\}\}.$$

If we have an n -tuple of independent GUE matrices $(A_1, \dots, A_n) \in (\mathbb{M}_k(\mathbb{C})_{\text{s.a.}})^n$, we know that as $k \rightarrow \infty$, that there is almost sure convergence in distribution to a family (s_1, \dots, s_n) of free semicirculars in M . The large deviations will then be given by

$$P_N(\Delta(A_1, \dots, A_n) \approx \Delta(x_1, \dots, x_n)) \sim e^{-k^2 I(x_1, \dots, x_n)},$$

where I is the free entropy.

To compute this value, we let λ denote the Lebesgue measure on $(\mathbb{M}_k(\mathbb{C})_{\text{s.a.}})^n \cong \mathbb{R}^{nk^2}$, and define

$$\chi(x_1, \dots, x_n) = \sup_{R > 0} \inf_{m \in \mathbb{N}} \inf_{\varepsilon > 0} \limsup_{k \rightarrow \infty} \left(\frac{1}{k^2} \ln \lambda(\Gamma_R(x_1, \dots, x_n; m, k, \varepsilon)) + \frac{n}{2} \ln k \right).$$

It turns out that the value of R has minor influence, and we only need choose a fixed R greater than the norm of each x_i .

¹A von Neumann algebra equipped with a faithful (injective on positive elements), normal (w^* -continuous), tracial ($\tau(xy) = \tau(yx)$) state ($\tau(1) = 1$) $\tau: M \rightarrow \mathbb{C}$.

Proposition ([Voi94, Proposition 2.2]): Let $C^2 = \tau(x_1^2 + \cdots + x_n^2)$. Then,

$$\chi(x_1, \dots, x_n) \leq 2^{-1} n \ln(2\pi e n^{-1} C^2).$$

Proof. We will instead prove a slightly different inequality. Define

$$\chi_R(x_1, \dots, x_n; m, k, \varepsilon) = \ln(\lambda(\Gamma_R(x_1, \dots, x_n; m, k, \varepsilon))),$$

and we will show that

$$\chi_R(x_1, \dots, x_n; m, k, \varepsilon) \leq 2^{-1} n k^2 (\ln(2\pi e n^{-1} (C^2 + n\varepsilon)) - \ln k),$$

assuming that $m \geq 2$. Applying the operations $\limsup_{k \rightarrow \infty}$, $\inf_{\varepsilon > 0}$, $\inf_{m \in \mathbb{N}}$, and $\sup_{R > 0}$ in succession will give us the free entropy.

For this, we use the p -dimensional Shannon entropy inequality

$$-\int f \ln(f) d\lambda_p \leq 2^{-1} p \ln(2\pi e p^{-1} a^2),$$

where f is some probability density function on \mathbb{R}^p and

$$a^2 = \int (x_1^2 + \cdots + x_p^2) f d\lambda_p,$$

where λ_p denotes the Lebesgue measure on \mathbb{R}^p .

We apply this to the special case of the Lebesgue measure on $(\mathbf{M}_k(\mathbb{C})_{\text{s.a.}})^n$, which is induced by the Hilbert–Schmidt metric

$$\langle (A_1, \dots, A_n), (B_1, \dots, B_n) \rangle = \sum_{i=1}^n \text{tr}(A_i B_i),$$

and take the indicator function

$$f(A_1, \dots, A_n) = \begin{cases} 0 & (A_1, \dots, A_n) \notin \Gamma_R(x_1, \dots, x_n; m, k, \varepsilon) \\ (\lambda(\Gamma_R(x_1, \dots, x_n; m, k, \varepsilon)))^{-1} & (A_1, \dots, A_n) \in \Gamma_R(x_1, \dots, x_n; m, k, \varepsilon), \end{cases}$$

giving

$$\chi_R(x_1, \dots, x_n; m, k, \varepsilon) \leq 2^{-1} n k^2 (\ln(2\pi e n^{-1} k^{-2} a^2) - \ln k).$$

Finally, by the definition of the microstate space, we have

$$\left| \frac{1}{k} \int \text{Tr} \left(\sum_{j=1}^n A_j^2 \right) f d\lambda - \tau \left(\sum_{j=1}^n x_j^2 \right) \right| < n\varepsilon,$$

meaning that $a^2 \leq C^2 + n\varepsilon$. □

Applications: Structural Properties of Free Group Factors

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