Solution (38.5): Copying the template equation, we have

$$\frac{\mathrm{d}v}{\mathrm{d}t} = -\frac{\mathrm{c}}{\mathrm{m}}v^2 + \mathrm{g},$$

where c is some constant. We see that the terminal velocity is

$$v_t = \sqrt{\frac{mg}{c}}$$
.

Separating variables, we have

$$\frac{dv}{-\frac{c}{m}v^2 + g} = dt$$

$$\frac{1}{g} \left( \frac{dv}{1 - \frac{c}{mg}v^2} \right) = dt$$

$$\frac{1}{g} \left( \frac{dv}{1 - (v/v_t)^2} \right) = dt.$$

Using the substitution  $\mathfrak{u}\coloneqq \nu/\nu_t,$  we have  $d\mathfrak{u}=\frac{1}{\nu_t}d\nu,$  meaning that

$$v_t \int \frac{1}{1 - u^2} du = \int g dt.$$

The integral of  $\frac{1}{1-u^2}$  is  $\frac{1}{2} \ln \left( \frac{1+u}{1-u} \right) = \operatorname{arctanh}(u)$ . Therefore, we have

$$\begin{aligned} \frac{v}{v_t} &= \tanh\left(\frac{g}{v_t}t\right) + K \\ v &= v_t \tanh\left(\frac{g}{v_t}t\right) + v_0 \\ &= \sqrt{\frac{mg}{c}} \tanh\left(\sqrt{\frac{c}{mq}}t\right) + v_0. \end{aligned}$$

## **Solution** (38.6):

(a) Using the chain rule and letting  $\frac{dm}{dt} = km^{2/3}$ , we have

$$\frac{dv}{dt} = km^{2/3} \frac{dv}{dm}$$
$$\frac{dv}{dm} + \frac{v}{m} = -\frac{b}{km}v + \frac{g}{km^{2/3}}.$$

With integrating factor  $m^{1+\frac{b}{k}}$ , we have

$$\begin{split} m^{1+\frac{b}{k}}\nu &= \frac{g}{k}\frac{m^{\frac{4}{3}+\frac{b}{k}}}{\frac{4}{3}+\frac{b}{k}} + C\\ \nu &= \frac{g}{k\Big(\frac{4}{3}+\frac{b}{k}\Big)}m^{\frac{1}{3}+\frac{b}{k}} + Cm^{-1-\frac{b}{k}}. \end{split}$$

We let  $v(m_0) = 0$ , so that

$$C = -\frac{g}{k(\frac{4}{3} + \frac{b}{k})} m_0^{\frac{4}{3} + \frac{b}{k}},$$

so

$$\nu = \frac{g}{\frac{4}{3}k + b} m^{\frac{1}{3}} \bigg( 1 - \bigg( \frac{m_0}{m} \bigg)^{\frac{4}{3} + \frac{b}{k}} \bigg).$$

Thus,

$$\begin{split} \frac{d\nu}{dt} &= g - \frac{1}{m} \frac{dm}{dt} \nu \\ &= g - \frac{1}{m} \Big( km^{2/3} \Big) \Bigg( \frac{g}{\frac{4}{3} k + b} m^{\frac{1}{3}} \bigg( 1 - \Big( \frac{m_0}{m} \Big)^{\frac{4}{3} + \frac{b}{k}} \bigg) \Bigg). \end{split}$$

(b) Using  $\frac{dm}{dt} = km^{2/3}v$ , and  $\frac{dv}{dt} = km^{2/3}v\frac{dv}{dm}$ , we obtain

$$\begin{split} m\frac{dv}{dt} + v\frac{dm}{dt} &= -bm^{2/3}v^2 + mg\\ v\,dv + \left(\frac{v^2}{m}\left(1 + \frac{b}{k}\right) - \frac{g}{km[2/3]}\right)dm &= 0. \end{split}$$

This gives  $\alpha = \nu$  and  $\beta = \frac{\nu^2}{m} \left( 1 + \frac{b}{k} \right) - \frac{g}{k m^{2/3}}$ . Solving for p(m), we get

$$p(m) = \frac{1}{\nu} \left( \frac{2\nu}{m} \left( 1 + \frac{b}{k} \right) \right)$$
$$= \frac{2}{m} \left( 1 + \frac{b}{k} \right).$$

Therefore, our integrating factor is

$$w(x) = m^{2 + \frac{2b}{k}}$$

This gives

$$\begin{split} &\frac{\partial\Phi}{\partial\nu}=\alpha\\ &\Phi=\frac{1}{2}m^{2+\frac{2b}{k}}\nu^2+c_1(m)\\ &\frac{\partial\Phi}{\partial m}=\beta\\ &\Phi=\frac{1}{2}m^{2+\frac{2b}{k}}\nu^2-\frac{g}{k\left(\frac{7}{3}+\frac{2b}{k}\right)}m^{\frac{7}{3}+\frac{2b}{k}}+c_2(\nu). \end{split}$$

Thus,  $c_2(v) = 0$ , and

$$\frac{1}{2}m^{2+\frac{2b}{k}}v^{2} - \frac{g}{k\left(\frac{7}{3} + \frac{2b}{k}\right)}m^{\frac{7}{3} + \frac{2b}{k}} = C.$$

Using  $v(m_0) = 0$ , we obtain the solution of

$$\frac{1}{2}m^{2+\frac{2b}{k}}\nu^2 = \frac{g}{k\left(\frac{7}{3} + \frac{2b}{k}\right)}m^{\frac{7}{3} + \frac{2b}{k}}\left(1 - \left(\frac{m_0}{m}\right)^{\frac{7}{3} + \frac{2b}{k}}\right).$$

Simplifying, this gives

$$v^{2} = \frac{2g}{k\left(\frac{7}{3} + \frac{2b}{k}\right)} m^{\frac{1}{3}} \left(1 - \left(\frac{m_{0}}{m}\right)^{\frac{7}{3} + \frac{2b}{k}}\right).$$

Therefore,

$$2\nu \frac{d\nu}{dm} = \frac{2g}{3k\left(\frac{7}{3} + \frac{2b}{k}\right)} m^{-2/3} \left(1 - \left(\frac{m_0}{m}\right)^{\frac{7}{3} + \frac{2b}{k}}\right) + \frac{2g}{km} \left(\frac{m_0}{m}\right)^{\frac{7}{3} + \frac{2b}{k}},$$

and

$$\begin{split} \frac{d\nu}{dt} &= \frac{k}{2} m^{2/3} \bigg( 2 \nu \frac{d\nu}{dm} \bigg) \\ &= \frac{g}{3 \bigg( \frac{7}{3} + \frac{2b}{k} \bigg)} \bigg( 1 - \bigg( \frac{m_0}{m} \bigg)^{\frac{7}{3} + \frac{2b}{k}} \bigg) + \frac{g}{m^{\frac{1}{3}}} \bigg( \frac{m_0}{m} \bigg)^{\frac{7}{3} + \frac{2b}{k}}. \end{split}$$

| **Solution** (38.7):

Solution (39.5): We take the derivative of

$$\frac{\mathrm{d} u_p}{\mathrm{d} x} = a_1(x) \frac{\mathrm{d} u_1}{\mathrm{d} x} + a_2(x) \frac{\mathrm{d} u_2}{\mathrm{d} x},$$

giving

$$\frac{d^2 u_p}{dx^2} = \alpha_1(x) \frac{d^2 u_1}{dx^2} + \frac{d\alpha_1}{dx} \frac{du_1}{dx} + \alpha_2(x) \frac{d^2 u}{dx^2} + \frac{d\alpha_2}{dx} \frac{du_2}{dx}.$$

Note that we must have

$$\frac{d^2u_p}{dx^2} + p(x)\frac{du_p}{dx} + q(x)u_p = r(x),$$

so we have

$$\begin{split} r(x) &= a_1(x) \frac{d^2 u_1}{dx^2} + \frac{d a_1}{dx} \frac{d u_1}{dx} + a_2(x) \frac{d^2 u}{dx^2} + \frac{d a_2}{dx} \frac{d u_2}{dx} \\ &+ p(x) \bigg( a_1(x) \frac{d u_1}{dx} + a_2(x) \frac{d u_2}{dx} \bigg) \\ &+ q(x) (a_1(x) u_1(x) + a_2(x) u_2(x)). \end{split}$$

Reordering and simplifying, we get

$$\begin{split} r(x) &= a_1(x) \left( \frac{d^2 u_1}{dx^2} + p(x) \frac{d u_1}{dx} + q(x) u_1(x) \right) + a_2(x) \left( \frac{d^2 u_2}{dx^2} + p(x) \frac{d u_2}{dx} + q(x) u_2(x) \right) + \frac{d a_1}{dx} \frac{d u_1}{dx} + \frac{d a_2}{dx} \frac{d u_2}{dx} \\ &= \frac{d a_1}{dx} \frac{d u_1}{dx} + \frac{d a_2}{dx} \frac{d u_2}{dx}. \end{split}$$

Pairing this expression with

$$\frac{\mathrm{d}a_1}{\mathrm{d}x}u_1(x) + \frac{\mathrm{d}a_2}{\mathrm{d}x}u_2(x) = 0,$$

we may solve for  $\frac{da_1}{dx}$  and  $\frac{da_2}{dx}$ , giving

$$\begin{split} \frac{da_1}{dx} &= -\frac{u_2(x)r(x)}{u_1(x)\frac{du_2}{dx} - u_2(x)\frac{du_1}{dx}} \\ \frac{da_2}{dx} &= \frac{u_1(x)r(x)}{u_1(x)\frac{du_2}{dx} - u_2(x)\frac{du_1}{dx}}. \end{split}$$

Therefore,

$$\begin{split} \alpha_1(x) &= -\int \frac{u_2(x)r(x)}{W(x)} \ dx \\ \alpha_2(x) &= \int \frac{u_1(x)r(x)}{W(x)} \ dx. \end{split}$$

**Solution** (39.7):

(a) We solve the homogeneous part to yield

$$u_1(x) = e^{-x}$$
  

$$u_2(x) = xe^{-x}$$

These give the Wronskian of

$$W(x) = e^{-x}(e^{-x} - xe^{-x}) + xe^{-2x}$$
  
=  $e^{-2x}$ .

We evaluate

$$a_{1}(x) = -\int e^{x}(xe^{-x})(e^{-x}) dx$$

$$= -\int xe^{-x} dx$$

$$= -(-xe^{-x} - e^{-x})$$

$$= xe^{-x} + e^{-x}$$

$$a_{2}(x) = \int e^{-x} dx$$

$$= -e^{-x}.$$

Thus, we have the general solution of

$$u(x) = c_1 e^{-x} + c_2 x e^{-x} + e^{-2x}$$
.

(b) Solving for the homogeneous solutions, we get

$$u_1(x) = e^x$$
  
$$u_2(x) = e^{-x},$$

with Wronskian

$$W(x) = -2$$
.

Setting up variation of parameters, we have

$$a_1(x) = -\int -\frac{1}{2} dx$$

$$= \frac{1}{2}$$

$$a_2(x) = -\frac{1}{2} \int e^{2x} dx$$

$$= -\frac{1}{4} e^{2x}.$$

Thus, we have the general solution of

$$u(x) = c_1 e^x + c_2 e^{-x} + \frac{1}{4} e^x.$$

(c) Solving for the homogeneous solution, we get

$$u_1(x) = \cos(x)$$
  
$$u_2(x) = \sin(x),$$

with Wronskian

$$W(x) = 1$$
.

Setting up variation of parameters, we then get

$$a_1(x) = -\int \sin(x)\cos(x) dx$$
$$= -\frac{1}{2}\cos(2x)$$
$$a_2(x) = \int \sin^2(x) dx$$
$$= \frac{1}{2}x + \frac{1}{2}\sin(2x).$$

Thus, we get the general solution of

$$u(x) = c_1 \cos(x) + c_2 \sin(x) + \frac{1}{2}(x + \sin(2x) - \cos(2x)).$$

Solution (39.8): We have the particular solutions of

$$u_1(t) = e^{\left(-\beta + \sqrt{\beta^2 - \omega_0^2}\right)t}$$
$$u_2(t) = e^{\left(-\beta - \sqrt{\beta^2 - \omega_0^2}\right)t}$$

Evaluating the Wronskian, we get

$$W(t) = -2\sqrt{\beta^2 - \omega_0^2} e^{-2\beta t},$$

so with variation of parameters, we have

$$\begin{split} \alpha_1(t) &= \frac{1}{2\sqrt{\beta^2 - \omega_0^2}} \int e^{\left(\beta - \sqrt{\beta^2 - \omega_0^2}\right)t} \delta \left(t - t'\right) \, dt \\ &= \frac{1}{2\sqrt{\beta^2 - \omega_0^2}} e^{\left(\beta - \sqrt{\beta^2 - \omega_0^2}\right)t'} \\ \alpha_2(t) &= -\frac{1}{2\sqrt{\beta^2 - \omega_0^2}} \int e^{\left(\beta + \sqrt{\beta^2 - \omega_0^2}\right)t} \delta \left(t - t'\right) \, dt \\ &= -\frac{1}{2\sqrt{\beta^2 - \omega_0^2}} e^{\left(\beta + \sqrt{\beta^2 - \omega_0^2}\right)t'}. \end{split}$$

Thus, we get the particular solution of

$$u_p(t) = \frac{1}{2\sqrt{\beta^2 - \omega_0^2}} \left( exp \left( \left(\beta - \sqrt{\beta^2 - \omega_0^2}\right) t' + \left(-\beta + \sqrt{\beta^2 - \omega_0^2}\right) t \right) - exp \left( \left(\beta + \sqrt{\beta^2 - \omega_0^2}\right) t' + \left(-\beta - \sqrt{\beta^2 - \omega_0^2}\right) t \right) \right).$$

## **Solution** (39.13):

(a) Setting up our differential equation of

$$\frac{d^2u}{dt^2} + 2\beta \frac{du}{dt} + \omega_0^2 t = F_0 \cos(\omega t),$$

we have homogeneous solutions of

$$\begin{split} u_1(t) &= e^{\left(-\beta + \sqrt{\beta^2 - \omega_0^2}\right)t} \\ u_2(t) &= e^{\left(-\beta - \sqrt{\beta^2 - \omega_0^2}\right)t} \,. \end{split}$$

Using the driving force  $F_0\cos(\omega t)$ , we use variation of parameters with the Wronskian of  $W(t)=e^{-2\beta\,t}$  to get

$$a_1(t) = -F_0 \int e^{\left(\beta - \sqrt{\beta^2 - \omega_0^2}\right)t} \cos(\omega t) dt$$

$$a_2(t) = F_0 \int e^{\left(\beta + \sqrt{\beta^2 - \omega_0^2}\right)t} \cos(\omega t) dt$$

and a particular solution of

$$u_{p}(t) = F_{0} \left( \frac{\omega_{0}^{2} \cos(\omega t) + 2\beta \omega \sin(\omega t) - \omega^{2} \cos(\omega t)}{\left(\omega_{0}^{2} - \omega^{2}\right)^{2} + (2\beta \omega)^{2}} \right).$$

## **Solution** (39.17):

(a) Using the power of the guess  $e^{\lambda t}$ , we find the solutions

$$u_1(t) = e^{-2t}$$
  
 $u_2(t) = e^{-t}$ .

(b) We find the Wronskian

$$W(t) = -3e^{-3t}$$

from which we are able to find

$$\begin{aligned} a_1(t) &= \frac{1}{3} \int e^{2t} \cos(t) dt \\ &= \frac{1}{15} e^{2t} (\sin(t) + 2\cos(t)) \\ a_2(t) &= -\frac{1}{3} \int e^t \cos(t) dt \\ &= -\frac{1}{6} e^t (\sin(t) + \cos(t)). \end{aligned}$$

Thus, the particular solution is

$$u_p(t) = \frac{1}{15}(\sin(t) + 2\cos(t)) - \frac{1}{6}(\sin(t) + \cos(t)).$$

(c) We find the full solution such that

$$c_1 + c_2 = \frac{31}{30}$$
$$-2c_1 - c_2 = \frac{1}{10}.$$

Therefore, we have

$$c_1 = -\frac{34}{30}$$
$$c_2 = \frac{13}{6}.$$

Our solution is

$$-\frac{34}{30}e^{-2t} + \frac{13}{6}e^{-t} - \frac{1}{30}\cos(t) - \frac{1}{10}\sin(t).$$

**Solution** (39.18): We start with the ansatz  $x^{\alpha}$ . Plugging this into our homogeneous equation, we get

$$x^{\alpha} \left( \alpha^2 - \alpha - 2 \right) = 0.$$

Therefore, we get that  $\alpha = 2, -1$ , giving the homogeneous solutions of

$$u_1(x) = x^2$$

$$u_2(x) = \frac{1}{x}.$$

We calculate the Wronskian to be W(x) = -3, so we use variation of parameters to obtain

$$a_1(x) = \frac{1}{3} \int \left(\frac{1}{x}\right) x \, dt$$
$$= \frac{1}{3} x$$
$$a_2(x) = -\frac{1}{3} \int x^3 \, dx$$

$$=-\frac{1}{12}x^4.$$

Therefore,

$$u(x) = a_1 x^2 + a_2 x^{-1} + \frac{1}{4} x^3.$$

Plugging in the initial conditions, we have

$$0 = a_1 + a_2 + \frac{1}{4}$$
$$0 = 2a_1 - a_2 + \frac{3}{4}.$$

This resolves to

$$a_1 = -\frac{1}{3}$$

$$a_2 = \frac{1}{12},$$

so we have the solution

$$u(x) = -\frac{1}{3}x^2 + \frac{1}{12}x^{-1} + \frac{1}{4}x^3.$$

**Solution** (39.21):

**Solution** (39.22 (b)):

Solution (39.28): We begin with the assumption that we have a power series of the form

$$u(x) = x^{\alpha} \sum_{k=0}^{\infty} c_k x^k.$$

Differentiating, we get

$$\begin{split} \frac{d^2u}{dx^2} &= \sum_{k=2}^{\infty} c_k(\alpha+k)(\alpha+k-1)x^{\alpha+k=2} \\ xu &= \sum_{k=1}^{\infty} c_{k-1}x^{\alpha+k}. \end{split}$$

Plugging this into Airy's equation, we are able to extract

$$c_2(\alpha+1)(\alpha+2) + \sum_{k=1}^{\infty} (c_{k+2}(\alpha+k+2)(\alpha+k+1) - c_{k-1})x^{\alpha+k} = 0.$$

Thus, we are left with the indicial equation of

$$c_2(\alpha+1)(\alpha+2)=0$$

and recurrence relation of

$$c_{k+2} = \frac{c_{k-1}}{(\alpha+k+2)(\alpha+k+1)}.$$

Since  $c_2$  is not the first term of the series, we are allowed to assume that  $c_2=0$  and  $\alpha=0$ . This gives chains  $c_0\to c_3\to \cdots$  and  $c_1\to c_4\to \cdots$  given by the recurrence relation. Therefore, we find the expressions

$$c_{3n} = c_0 \left( \prod_{j=1}^{n} (3j)(3j-1) \right)^{-1}$$

$$c_{3n+1} = c_1 \left( \prod_{j=1}^n (3j+1)(3j) \right)^{-1},$$

whose corresponding series are linearly independent.