

The 12th HW of Electrodynamics

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1

Solve the equation:

$$\left(\nabla^2 - \frac{1}{\lambda^2}\right)\varphi = -\frac{q}{\epsilon_0}\delta(\vec{x})$$

solutions:

两边 fourier 变换后:

$$-(k^2 + \frac{1}{\lambda^2})\tilde{\varphi} = -\frac{q}{\epsilon_0}$$

因此

$$\tilde{\varphi} = \frac{\frac{q}{\epsilon_0}}{k^2 + \frac{1}{\lambda^2}}$$

求逆变换:

$$\varphi = \frac{q}{(2\pi)^3\epsilon_0} \iiint \frac{e^{i\mathbf{k}\cdot\mathbf{r} \cos\theta}}{k^2 + \frac{1}{\lambda^2}} d\mathbf{k}$$

化简:

$$\begin{aligned}
\varphi &= \frac{q}{(2\pi)^3 \varepsilon_0} \iiint \frac{e^{ikr \cos \theta}}{k^2 + \frac{1}{\lambda^2}} d\mathbf{k} \\
&= \frac{q}{(2\pi)^3 \varepsilon_0} \int_0^\pi d\theta \int_0^{2\pi} d\phi \int_0^\infty k^2 \sin \theta \frac{e^{ikr \cos \theta}}{k^2 + \frac{1}{\lambda^2}} dk \\
&= \frac{q}{(2\pi)^2 \varepsilon_0} \int_0^\pi d\theta \int_0^\infty k^2 \sin \theta \frac{e^{ikr \cos \theta}}{k^2 + \frac{1}{\lambda^2}} dk \\
&= \frac{q}{i(2\pi)^2 r \varepsilon_0} \int_0^\infty \frac{2k \sin kr}{k^2 + \frac{1}{\lambda^2}} dk \\
&= \frac{q}{i(2\pi)^2 r \varepsilon_0} \int_{-\infty}^\infty \frac{ke^{ikr}}{k^2 + \frac{1}{\lambda^2}} dk \\
&= \frac{q}{i(2\pi)^2 r \varepsilon_0} \int_{-\infty}^\infty \frac{ke^{ikr}}{k^2 + \frac{1}{\lambda^2}} dk \\
&= \frac{q}{i(2\pi)^2 r \varepsilon_0} 2\pi i \operatorname{Res} \left[\frac{ze^{izr}}{z^2 + \frac{1}{\lambda^2}}, z_0 \right] \\
&\quad \text{在 } z_0 = \frac{i}{\lambda} \text{ 取一阶极点} \\
&= \frac{q}{4\pi r \varepsilon_0} e^{-\frac{r}{\lambda}}
\end{aligned}$$

2

$$\begin{aligned}
\varphi(\vec{r}, t) &= \frac{1}{4\pi \varepsilon_0} \iiint d^3 \vec{r}' \frac{1}{|\vec{r} - \vec{r}'|} \rho \left(\vec{r}', t - \frac{1}{c} |\vec{r} - \vec{r}'| \right) \\
&= \frac{q}{4\pi \varepsilon_0} \iiint \frac{\delta(\mathbf{r}' - \mathbf{r}_e(t'))}{R} d^3 r'
\end{aligned}$$

做变量替换: $\mathbf{r}'' = \mathbf{r}' - \mathbf{r}_e(t')$. $J = \det(\nabla' \mathbf{r}'') = \det[\nabla' \mathbf{r}' - \nabla' \mathbf{r}_e(t')] = 1 - \frac{\mathbf{v} \cdot \mathbf{R}}{cR}$

$$\begin{aligned}
\varphi(\mathbf{r}, t) &= \frac{q}{4\pi \varepsilon_0} \iiint \frac{\delta(\mathbf{r}'')}{R} |J|^{-1} dV'' = \frac{q}{4\pi \varepsilon_0} \iiint \frac{\delta(\mathbf{r}'')}{R(1 - \frac{\mathbf{R} \cdot \mathbf{v}}{cR})} dV'' \\
&= \frac{q}{4\pi \varepsilon_0 [R(1 - \frac{\mathbf{v} \cdot \mathbf{R}}{cR})]_{r''=0}} = \frac{q}{4\pi \varepsilon_0 \left[|\vec{r} - \vec{r}'| \left(1 - \frac{v \cos \theta}{c} \right) \right]}
\end{aligned}$$

同理

$$\begin{aligned}
A(r, t) &= \frac{\mu_0}{4\pi} \iiint \frac{\mathbf{J}(\mathbf{r}', t - \frac{R}{c})}{R} dV' = \frac{\mu_0 q}{4\pi} \iiint \frac{v(t') \delta(\mathbf{r}' - \mathbf{r}_e(t'))}{R} dV' \\
&\Rightarrow \vec{A}(\vec{r}, t) = \frac{\vec{v}(t')}{c^2} \varphi(\vec{r}, t)
\end{aligned}$$