

ElectroDynamics

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1 方程

真空麦克斯韦方程

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

$$\oint_S \mathbf{E} \, d\mathbf{s} = \frac{Q}{\epsilon_0}$$

高斯定律

$$\nabla \cdot \mathbf{B} = 0$$

$$\oint_S \mathbf{B} \, d\mathbf{s} = 0$$

高斯磁定律

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\oint_L \mathbf{E} \, d\mathbf{l} = -\frac{d\varphi_B}{dt}$$

法拉第电磁感应定律

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

$$\oint_L \mathbf{B} \, d\mathbf{l} = \mu_0 I + \mu_0 \epsilon_0 \frac{d\varphi_E}{dt}$$

安培定律

物质内麦克斯韦方程

$$\nabla \cdot \mathbf{D} = \rho_f$$

$$\oint_S \mathbf{D} \, d\mathbf{s} = Q_f$$

高斯定律

$$\nabla \cdot \mathbf{B} = 0$$

$$\oint_S \mathbf{B} \, d\mathbf{s} = 0$$

高斯磁定律

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\oint_L \mathbf{E} \, d\mathbf{l} = -\frac{d\varphi_B}{dt}$$

法拉第电磁感应定律

$$\nabla \times \mathbf{H} = \mathbf{J}_f + \frac{\partial \mathbf{D}}{\partial t}$$

$$\oint_L \mathbf{H} \, d\mathbf{l} = I_f + \frac{d\varphi_D}{dt}$$

安培定律

镜像法

距半径为 R_0 的球的球心距离为 a 处有一点电荷 q , 则镜像电荷 $-\frac{R_0}{a}q$ 距球心 $\frac{R_0^2}{a}$ 远, 且在靠近 q 方向.

球谐函数解拉普拉斯方程

$$\varphi(r, \theta) = \sum_{n=0}^{\infty} \left(a_n r^n + \frac{b_n}{r^{n+1}} \right) P_n(\cos \theta)$$

$$\begin{cases} P_0(\cos \theta) = 1 \\ P_1(\cos \theta) = \cos \theta \\ P_2(\cos \theta) = \frac{1}{2}(3 \cos^2 \theta - 1) \end{cases}$$

$$\frac{1}{|\mathbf{R} - \mathbf{a}'|} = \frac{1}{\sqrt{R^2 + a'^2 - 2Ra \cos \theta}} = \begin{cases} \frac{1}{a} \sum_{n=0}^{\infty} \left(\frac{R}{a}\right)^n P_n(\cos \theta), & (R < a) \\ \frac{1}{R} \sum_{n=0}^{\infty} \left(\frac{a}{R}\right)^n P_n(\cos \theta), & (R > a) \end{cases}$$

格林函数法

$$\nabla^2 G(x, x') = -\frac{1}{\varepsilon} \delta^3(x - x')$$

(1) 无界空间中

$$G(x, x') = \frac{1}{4\pi\varepsilon_0} \frac{1}{|r - r'|}$$

(2) 上半平面中

$$G(x, x') = \frac{1}{4\pi\varepsilon_0} \left(\frac{1}{|r - r'|} - \frac{1}{|r + r'|} \right)$$

(3) 球外空间 (R' 为电荷位置, α 为场点与电荷位置夹角, R_0 为球半径).

$$G(x, x') = \frac{1}{4\pi\varepsilon_0} \left(\frac{1}{R^2 + R'^2 - 2RR' \cos \alpha} - \frac{1}{\left(\frac{RR'}{R_0}\right)^2 + R_0^2 - 2RR' \cos \alpha} \right)$$

给定 $\rho(x')$, 第一类边值问题的解为 (G 交换了 x, x'):

$$\varphi(x) = \int_V G(x', x) \rho(x') dV' + \varepsilon_0 \oint_S \left(G(x', x) \frac{\partial \varphi}{\partial n'} - \varphi(x') \frac{\partial G(x', x)}{\partial n'} dS' \right)$$

若给定边界 φ , 则应使 G 在边界为 0, 若给定边界 $\frac{\partial \varphi}{\partial n}$, 则应使 $\frac{\partial G}{\partial n}$ 在边界为 0.

泊松方程

$$\nabla^2 \Phi = \nabla \cdot \mathbf{E} = -\frac{\rho}{\epsilon_0}$$

电多极矩

$$\varphi(\mathbf{x}) = \frac{1}{4\pi\epsilon_0} \left(\frac{q}{R} + \frac{\mathbf{p} \cdot \mathbf{R}}{R^3} + \frac{1}{6} \sum_{i,j} \mathcal{D}_{ij} \frac{\partial^2}{\partial x_i \partial x_j} \frac{1}{R} + \dots \right)$$

$$\mathbf{E} = -\frac{1}{4\pi\epsilon_0} \left(\frac{3(\mathbf{p} \cdot \mathbf{R}) \mathbf{R}}{R^5} - \frac{\mathbf{p}}{R^3} \right)$$

$$\mathbf{p} = \iiint_V \rho(x') x' d^3x'$$

$$\mathfrak{D} = \iiint_V 3\mathbf{x}'\mathbf{x}'\rho(\mathbf{x}')d^3x'$$

磁偶极矩

$$\varphi = \frac{\mathbf{m} \cdot \mathbf{R}}{4\pi R^3}$$

$$\mathbf{m} = \frac{1}{2} \iiint_V \mathbf{x}' \times \mathbf{J}(\mathbf{x}') d^3x'$$

保角变换 (z_1 为原来的点, a 为夹角出现的位置的横坐标, α 为边界夹角).

$$\frac{dz_1}{dz_2} = C_1 \prod_{i=1}^n (z_2 - a_i)^{\frac{\alpha_i}{\pi} - 1}$$

电荷

面电荷	电场	磁场
$\sigma_{polar} = P$	$\rho_p = -\nabla \cdot \mathbf{P}$	$\rho_M = -\mu_0 \nabla \cdot \mathbf{M}$
$\sigma_{free} = D$	$\rho_f = \nabla \cdot \mathbf{D}$	$\rho_f = 0$
$\sigma_{total} = \epsilon_0 E$	$\rho_{tot} = \epsilon_0 \nabla \cdot \mathbf{E}$	$\rho_{tot} = \rho_p = \mu_0 \nabla \cdot \mathbf{H}$
	$\epsilon_0 E = D - P$	$B = \mu_0 (H + M)$

边界条件

磁场	电场
$\varphi_1 = \varphi_2$	$\varphi_1 = \varphi_2$
$\mu_1 \frac{\partial \varphi_1}{\partial n} = \mu_2 \frac{\partial \varphi_2}{\partial n}$	$\epsilon_1 \frac{\partial \varphi_1}{\partial n} + \sigma_f = \epsilon_2 \frac{\partial \varphi_2}{\partial n}$
	$\sigma_{\text{电导}1} \frac{\partial \varphi_1}{\partial n} = \sigma_{\text{电导}2} \frac{\partial \varphi_2}{\partial n}$
$H_{1\parallel} + \alpha_f = H_{2\parallel}$	$E_{1\parallel} = E_{2\parallel}$
$B_{1\perp} = B_{2\perp}$	$D_{1\perp} + \sigma_f = D_{2\perp}$
$M_{1\perp} + \alpha_M = M_{2\perp}$	
$B_{2\perp} - B_{1\perp} = \mu_0 (\alpha_f + \alpha_M)$	$E_{2\perp} - E_{1\perp} = (\sigma_f + \sigma_p) / \epsilon_0$
$B_{\perp} = 0$ (超导球)	$D_{\perp} = \sigma_f$ (导体)

洛伦兹力:

$$\mathbf{F} = q\mathbf{E} + q\mathbf{v} \times \mathbf{B}$$

$$\mathbf{f} = \rho\mathbf{E} + \mathbf{J} \times \mathbf{B}$$

磁偶极子:

$$\vec{B} = \frac{\mu_0}{4\pi} \frac{3(\vec{m} \cdot \hat{R})\hat{R} - \vec{m}}{R^3}$$

$$\varphi = \frac{\vec{m} \cdot \vec{R}}{4\pi R^3}$$

电磁场:

$$\mathbf{S} = \mathbf{E} \times \mathbf{H}$$

$$w = \frac{1}{2} (\mathbf{E} \cdot \mathbf{D} + \mathbf{H} \cdot \mathbf{B}) = \frac{1}{2} (\rho\varphi + \mathbf{J}_f \cdot \mathbf{A})$$

电流:

$$\nabla \cdot \mathbf{J} = -\frac{\partial \rho}{\partial t}$$

$$\mathbf{J} = \sigma \mathbf{E}$$

毕奥——萨伐尔定律 $\mathbf{B} = \frac{\mu_0}{4\pi} \int \frac{I d\mathbf{l} \times \mathbf{e}_r}{r^2}$, 若 I 为直线, $B = \frac{\mu_0 I l}{4\pi r^2}$.

电磁波:

$$\left. \begin{aligned} \nabla \times \mathbf{B} - \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t} &= 0 \\ \nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} &= 0 \end{aligned} \right| \begin{aligned} \nabla^2 \mathbf{E} - \frac{1}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} &= \square \mathbf{E} = 0 \\ \nabla^2 \mathbf{B} - \frac{1}{c^2} \frac{\partial^2 \mathbf{B}}{\partial t^2} &= \square \mathbf{B} = 0 \end{aligned}$$

磁矢势:

库仑规范:

$$\nabla \cdot \mathbf{A} = 0$$

$$\nabla^2 \varphi = -\frac{\rho}{\epsilon_0}$$

$$\square \mathbf{A} = -\mu_0 \mathbf{J} + \frac{1}{c^2} \nabla \frac{\partial \varphi}{\partial t}$$

洛伦兹规范:

$$\nabla \cdot \mathbf{A} + \frac{1}{c^2} \frac{\partial \varphi}{\partial t} = 0$$

$$\square \varphi = -\frac{\rho}{\epsilon_0}$$

$$\square \mathbf{A} = -\mu_0 \mathbf{J}$$

2 数学

2.1 柱坐标系 (ρ, ϕ, z)

$$\nabla \varphi = \hat{e}_1 \frac{\partial \varphi}{\partial \rho} + \hat{e}_2 \frac{1}{\rho} \frac{\partial \varphi}{\partial \phi} + \hat{e}_3 \frac{\partial \varphi}{\partial z}$$

$$\begin{aligned}
\nabla \cdot \mathbf{A} &= \frac{1}{\rho} \frac{\partial(\rho A_1)}{\partial \rho} + \frac{1}{\rho} \frac{\partial A_2}{\partial \phi} + \frac{\partial A_3}{\partial z} \\
\nabla \times \mathbf{A} &= \hat{e}_1 \left(\frac{1}{\rho} \frac{\partial A_3}{\partial \phi} - \frac{\partial A_2}{\partial z} \right) + \hat{e}_2 \left(\frac{\partial A_1}{\partial z} - \frac{\partial A_3}{\partial \rho} \right) + \hat{e}_3 \frac{1}{\rho} \left(\frac{\partial(\rho A_2)}{\partial \rho} - \frac{\partial A_1}{\partial \phi} \right) \\
\nabla^2 \varphi &= \frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \varphi}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \varphi}{\partial \phi^2} + \frac{\partial^2 \varphi}{\partial z^2}
\end{aligned}$$

2.2 球坐标系 (r, θ, φ)

$$\begin{aligned}
\nabla \varphi &= \hat{e}_1 \frac{\partial \varphi}{\partial r} + \hat{e}_2 \frac{1}{r} \frac{\partial \varphi}{\partial \theta} + \hat{e}_3 \frac{1}{r \sin \theta} \frac{\partial \varphi}{\partial \phi} \\
\nabla \cdot \mathbf{A} &= \frac{1}{r^2} \frac{\partial(r^2 A_1)}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta A_2) + \frac{1}{r \sin \theta} \frac{\partial A_3}{\partial \phi} \\
\nabla \times \mathbf{A} &= \hat{e}_1 \frac{1}{r \sin \theta} \left[\frac{\partial}{\partial \theta} (\sin \theta A_3) - \frac{\partial A_2}{\partial \phi} \right] + \hat{e}_2 \left[\frac{1}{r \sin \theta} \frac{\partial A_1}{\partial \phi} - \frac{1}{r} \frac{\partial}{\partial r} (r A_3) \right] + \hat{e}_3 \frac{1}{r} \left[\frac{\partial}{\partial r} (r A_2) - \frac{\partial A_1}{\partial \theta} \right] \\
\nabla^2 \varphi &= \frac{1}{r^2 \sin \theta} \left[\sin \theta \frac{\partial}{\partial r} \left(r^2 \frac{\partial \varphi}{\partial r} \right) + \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \varphi}{\partial \theta} \right) + \frac{1}{\sin \theta} \frac{\partial^2 \varphi}{\partial \phi^2} \right]
\end{aligned}$$

2.3 矢量变换

$$\begin{aligned}
\nabla \mathbf{r} &= -\nabla' \mathbf{r} = \mathbf{e}_r \\
\nabla \frac{1}{\mathbf{r}} &= -\nabla' \frac{1}{\mathbf{r}} = -\frac{1}{r^2} \mathbf{e}_r \\
\nabla \times \frac{1}{\mathbf{r}^2} &= \nabla \cdot \frac{1}{\mathbf{r}} = 0 \\
\nabla \cdot \varphi \mathbf{A} &= \varphi \nabla \cdot \mathbf{A} + \mathbf{A} \cdot \nabla \varphi \\
\nabla \times \varphi \mathbf{A} &= \varphi \nabla \times \mathbf{A} + \nabla \varphi \times \mathbf{A} \\
\nabla \cdot (\mathbf{A} \cdot \mathbf{B}) &= (\mathbf{A} \cdot \nabla) \mathbf{B} + (\mathbf{B} \cdot \nabla) \mathbf{A} + \mathbf{A} \times (\nabla \times \mathbf{B}) + \mathbf{B} \times (\nabla \times \mathbf{A}) \\
\nabla \cdot (\mathbf{F} \times \mathbf{G}) &= (\nabla \times \mathbf{F}) \cdot \mathbf{G} - \mathbf{F} \cdot (\nabla \times \mathbf{G}) \\
\mathbf{A} \times (\mathbf{B} \times \mathbf{C}) &= (\mathbf{A} \cdot \mathbf{C}) \mathbf{B} - (\mathbf{A} \cdot \mathbf{B}) \mathbf{C} \\
(\mathbf{A} \times \mathbf{B}) \times (\mathbf{C} \times \mathbf{D}) &= [\mathbf{A} \cdot (\mathbf{B} \times \mathbf{D})] \mathbf{C} - [\mathbf{A} \cdot (\mathbf{B} \times \mathbf{C})] \mathbf{D} \\
(\mathbf{A} \times \mathbf{B}) \cdot (\mathbf{C} \times \mathbf{D}) &= \begin{vmatrix} \mathbf{A} \cdot \mathbf{C} & \mathbf{A} \cdot \mathbf{D} \\ \mathbf{B} \cdot \mathbf{C} & \mathbf{B} \cdot \mathbf{D} \end{vmatrix} \\
\mathbf{A} \times (\mathbf{B} \times \mathbf{C}) + \mathbf{B} \times (\mathbf{C} \times \mathbf{A}) + \mathbf{C} \times (\mathbf{A} \times \mathbf{B}) &= 0
\end{aligned}$$