

数理方法 II 第三次作业

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(1)

$$x \frac{d^2 y}{dx^2} + \frac{dy}{dx} + (2 + \lambda/x)y = 0$$

$$\frac{d}{dx} \left(x \frac{dy}{dx} \right) - (-2)y + \frac{\lambda}{x}y = 0$$

即

$$k(x) = x, q(x) = -2, \rho(x) = \frac{1}{x}$$

(2)

原式化为 $y'' + \frac{a-bx}{x-x^2}y' - \frac{\lambda}{x-x^2}y = 0$

$$\exp\left(\int \frac{a-bx}{x-x^2} dx\right) = \exp\left(a \int \frac{1}{x} dx + (a-b) \int \frac{1}{1-x} dx\right) = \exp(a \ln x - (a-b) \ln(1-x)) = \frac{x^a}{(1-x)^{a-b}}$$

则最后可化为标准形式:

$$\frac{d}{dx}\left(\frac{x^a}{(1-x)^{a-b}}y'\right) + \lambda\left(\frac{x^{a-1}}{(1-x)^{a-b+1}}\right)y = 0$$

2

设 $y_m, y_n, n \neq m$ 是函数不同本征值的两个解.

$$\begin{cases} \frac{d}{dx}(py'_m) + (\lambda_m \rho - q)y_m = 0 \\ \frac{d}{dx}(py'_n) + (\lambda_n \rho - q)y_n = 0 \end{cases}$$

两式分别乘以 y_n, y_m , 相减, $y_n \frac{d}{dx}(py'_m) - y_m \frac{d}{dx}(py'_n) + (\lambda_m - \lambda_n)y_m y_n = 0$. 求区间 $[a, b]$ 积分,

$$\begin{aligned} & \int_a^b \left[y_n \frac{d}{dx}(py'_m) - y_m \frac{d}{dx}(py'_n) \right] dx + \int_a^b (py'_m \frac{d}{dx}y_n - py'_n \frac{d}{dx}y_m) dx + (\lambda_m - \lambda_n) \int_a^b \rho y_m y_n dx \\ &= \int_a^b \frac{d}{dx}(py_n y'_m - py_m y'_n) dx + (\lambda_m - \lambda_n) \int_a^b \rho y_m y_n dx \\ &= (py_n y'_m - py_m y'_n)|_{x=b} - (py_n y'_m - py_m y'_n)|_{x=a} + (\lambda_m - \lambda_n) \int_a^b \rho y_m y_n dx \end{aligned}$$

由于 $\begin{vmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{vmatrix}$, 有 $y_n(b)y'_m(b) = y_m(b)y'_n(b)$, $y_n(a)y'_m(a) = y_m(a)y'_n(a)$

$$= (\lambda_m - \lambda_n) \int_a^b \rho y_m y_n dx = 0$$

当 $\lambda_m \neq \lambda_n$, $\int_a^b \rho y_m y_n dx = 0$.

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(1)

根据定义, $\int \delta(\mathbf{r} - \mathbf{r}_0) d\mathbf{r}^3 = 1$, 在球坐标下即为

$$\begin{aligned} \int_0^\infty dr \int_0^{2\pi} r d\varphi \int_0^\pi r \sin \varphi \delta(\mathbf{r} - \mathbf{r}_0) d\theta &= 1 \\ \Rightarrow \int_0^\infty dr \int_0^{2\pi} d\cos \varphi \int_0^\pi r^2 \delta(\mathbf{r} - \mathbf{r}_0) d\theta &= 1 \end{aligned}$$

根据直角坐标下形式, 可知

$$\delta(r - r_0) \delta(\cos \theta - \cos \theta_0) \delta(\varphi - \varphi_0) = r^2 \delta(\mathbf{r} - \mathbf{r}_0)$$

移项即得

$$\delta(\mathbf{r} - \mathbf{r}_0) = \frac{1}{r^2} \delta(r - r_0) \delta(\cos \theta - \cos \theta_0) \delta(\varphi - \varphi_0)$$

(2)

$$\nabla^2 \frac{1}{|r - r_0|} = -\nabla \cdot \frac{1}{(r - r_0)^2}$$

由高斯定理可知

$$-\int \nabla \cdot \frac{1}{(r - r_0)^2} dV = -\int_{\Omega} \frac{1}{(r - r_0)^2} dS$$

取积分面为 $r - r_0 = a$ 的球壳, a 为任意常数.

$$-\int_{\Omega} \frac{1}{(r - r_0)^2} dS = -4\pi a^2 \frac{1}{a^2} = -4\pi$$

即

$$\int \nabla^2 \frac{1}{|r - r_0|} dr^3 = -4\pi$$

根据定义,

$$\nabla^2 \frac{1}{|r - r_0|} = -4\pi \delta(\mathbf{r} - \mathbf{r}_0)$$

4

(1) 先计算 $\mathcal{F}(e^{-a|t|})$

$$\begin{aligned}\mathcal{F}(e^{-a|t|}) &= \int_{-\infty}^0 e^{at-iwt} dt + \int_0^{\infty} e^{-at-iwt} dt \\ &= \int_0^{\infty} (e^{-(a+iw)t} + e^{-(a-iw)t}) dt = \frac{2a}{a^2 + w^2}\end{aligned}$$

则 $\mathcal{F}^{-1}(\frac{2a}{a^2+w^2}) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{2a}{a^2+w^2} e^{iwt} dw = e^{-a|t|}$, 两边求实部:

$$\int_{-\infty}^{\infty} \frac{1}{a^2 + w^2} \cos wt dw = \frac{\pi}{a} e^{-a|t|}$$

(2)

(i)

$$\begin{aligned}
\int \frac{1}{r} e^{ikr \cos \theta} dr &= \int_0^\infty \frac{1}{r} r^2 dr \int_1^{-1} e^{ikr \cos \theta} d \cos \theta \int_0^{2\pi} d\varphi \\
&= 2\pi \int_0^\infty \frac{1}{r} r^2 \frac{1}{ikr} (e^{ikr} - e^{-ikr}) dr = \frac{2\pi}{ik} \int_0^\infty (e^{ikr} - e^{-ikr}) dr \\
&= \frac{2\pi}{ik} (2i \int_0^\infty \sin kr dr) = \frac{4\pi}{k} \int_0^\infty \sin kr dr \\
\int_0^\infty \sin kr dr &= \lim_{\varepsilon \rightarrow 0^+} \operatorname{Im} \left(\int_0^\infty e^{\varepsilon r} e^{ikr} dr \right) \\
&= \lim_{\varepsilon \rightarrow 0^+} \operatorname{Im} \left(\int_0^\infty e^{(\varepsilon + ik)r} dr \right) \\
&= \lim_{\varepsilon \rightarrow 0^+} \operatorname{Im} \left(\frac{e^{(\varepsilon + ik)\infty} - 1}{\varepsilon + ik} \right) \\
&= \frac{1}{k}
\end{aligned}$$

则

$$\mathcal{F}\left(\frac{1}{r}\right) = \frac{1}{(2\pi)^{3/2}} \int \frac{1}{r} e^{ikr \cos \theta} dr = \frac{\sqrt{2}}{\sqrt{\pi}k} \int_0^\infty \sin kr dr = \frac{\sqrt{2}}{\sqrt{\pi}k^2}$$

(ii)

$$\begin{aligned}
\int \frac{\delta(r-a)}{r} e^{-ikr \cos \theta} dr &= \int_0^\infty \frac{\delta(r-a)}{r} r^2 dr \int_1^{-1} e^{-ikr \cos \theta} d \cos \theta \int_0^{2\pi} d\varphi \\
&= -2\pi \int_0^\infty \frac{\delta(r-a)}{ik} (e^{ikr} - e^{-ikr}) dr = -\frac{2\pi}{ik} (e^{ika} - e^{-ika}) = \frac{\pi}{k} \sin ka
\end{aligned}$$

则可得到

$$\mathcal{F}^{-1}\left(\frac{\sin ak}{k}\right) = \sqrt{\frac{\pi}{2}} \frac{\delta(r-a)}{r}$$

5

根据周期性有:

$$\mathcal{L}(f(t-a)) = \mathcal{L}(f(t)u(t-a))$$

又由于

$$\mathcal{L}(f(t-a)) = \int_0^\infty \frac{f(t-a)}{e^{ap}} e^{-pt+ap} dt = e^{-ap} F(p)$$

$$\mathcal{L}(f(t)u(t-a)) = F(p) - \int_0^a f(t)e^{-pt} dt$$

上面两式相等即可得到

$$(1 - e^{-ap})F(p) = \int_0^a f(t)e^{-pt} dt \implies F(p) = \frac{1}{1 - e^{-ap}} \int_0^a f(t)e^{-pt} dt$$

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(1)

$$\begin{aligned}
pU(x, p) - u(x, 0) &= a^2 \frac{\partial^2 U}{\partial x^2} + F(x, p) \\
\implies pU(x, p) - \varphi(x) &= a^2 \frac{\partial^2 U}{\partial x^2} + F(x, p)
\end{aligned}$$

有通解

$$\begin{aligned}
U(x, p) &= -\frac{a}{2\sqrt{p}} \int_{-\infty}^{\infty} \exp\left(-\frac{\sqrt{p}}{a}|x-x'|\right) \left(-\frac{F(x') + \varphi(x')}{a^2}\right) dx' \\
&= \frac{1}{2a} \int_{-\infty}^{\infty} \exp\left(-\frac{\sqrt{p}}{a}|x-x'|\right) \left(\frac{F(x') + \varphi(x')}{\sqrt{p}}\right) dx' \\
&= \frac{1}{2a} \int_{-\infty}^{\infty} \frac{1}{\sqrt{p}} \exp\left(-\frac{\sqrt{p}}{a}|x-x'|\right) (F(x') + \varphi(x')) dx'
\end{aligned}$$

记 $m(x, x', p) = \frac{1}{p} e^{-\frac{|x-x'|}{a}p}$, 则 $\mathcal{L}^{-1}(m) = \eta(t - \frac{|x-x'|}{a})$. 那么

$$m(x, x', p) = \int_0^{\infty} \eta\left(t - \frac{|x-x'|}{a}\right) e^{-pt} dt$$

或写为

$$m(x, x', \sqrt{p}) = \int_0^{\infty} \eta\left(t - \frac{|x-x'|}{a}\right) e^{-\sqrt{p}t} dt = \int_0^{\infty} \eta\left(t - \frac{|x-x'|}{a}\right) \frac{e^{-\sqrt{p}t}}{e^{-pt}} e^{-pt} dt$$

因此 $\mathcal{L}^{-1}(m(x, x', \sqrt{p})) = \eta\left(t - \frac{|x-x'|}{a}\right) \frac{e^{-\sqrt{p}t}}{e^{-pt}}$

$$u(x, t) = \frac{1}{2a} \int_{-\infty}^{\infty} \eta\left(t - \frac{|x-x'|}{a}\right) \frac{e^{-\sqrt{p}t}}{e^{-pt}} (F(x') + \varphi(x')) dx'$$

$$u(x, t) = \int_{-\infty}^{\infty} \phi(\xi) K(x - \xi, t) d\xi + \int_0^t d\tau \int_{-\infty}^{\infty} f(\xi, \tau) K(x - \xi, t - \tau) d\xi$$

(2)

$$\frac{\partial U}{\partial t}(w, t) = -a^2 w^2 U(w, t) + F(w, t)$$

$$U|_{t=0} = \Phi(w)$$

则

$$U(w, t) = e^{-a^2 w^2 t} \Phi(w) + \int_0^t F(w, \tau) e^{-a^2 w^2 (t-\tau)} d\tau$$

做逆变换:

$$u(x, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \Phi(w) e^{-a^2 w^2 t} e^{iwx} dw + \mathcal{F}^{-1} \left\{ \int_0^t F(w, \tau) e^{-a^2 w^2 (t-\tau)} d\tau \right\}$$

$$u(x, t) = \int_{-\infty}^{\infty} \phi(\xi) K(x - \xi, t) d\xi + \int_0^t d\tau \int_{-\infty}^{\infty} f(\xi, \tau) K(x - \xi, t - \tau) d\xi$$

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先对 t 求拉普拉斯, $\mathcal{L}(u(x, t)) = U(x, p)$:

$$\begin{aligned} p^2 U(x, p) - pu(x, 0) - \frac{\partial u}{\partial t}(x, 0) &= c^2 \frac{\partial^2 U}{\partial x^2}(x, p) \\ \implies p^2 U(x, p) - p\varphi(x) - \psi(x) &= c^2 \frac{\partial^2 U}{\partial x^2}(x, p) \end{aligned}$$

再对 x 求傅里叶, $\mathcal{F}(U(x, p)) = U^*(w, p)$:

$$\begin{aligned} p^2 U^*(w, p) - p\Phi(w) - \Psi(w) &= -c^2 w^2 U^*(w, p) \\ \implies U^*(w, p) &= \frac{p\Phi(w) + \Psi(w)}{p^2 + c^2 w^2} \end{aligned}$$

拉普拉斯逆变换得到

$$\begin{aligned} U(w, t) &= \Phi(w) \cos(cwt) + \frac{\Psi(w)}{cw} \sin(cwt) \\ \mathcal{F}^{-1}[\Phi(w)e^{icwt}] &= \phi(x + ct) \\ \mathcal{F}^{-1}[\Phi(w)e^{-icwt}] &= \phi(x - ct) \\ \mathcal{F}^{-1}\left[\frac{1}{cw}\Psi(w)e^{icwt}\right] &= \frac{i}{c} \int_{-\infty}^{x+ct} \psi(x + ct) dx \\ \mathcal{F}^{-1}\left[\frac{1}{cw}\Psi(w)e^{-icwt}\right] &= \frac{i}{c} \int_{-\infty}^{x-ct} \psi(x - ct) dx \\ \implies u(x, t) &= \frac{1}{2} \left[\phi(x + ct) + \phi(x - ct) + \frac{1}{c} \int_{-\infty}^{x+ct} \psi(x + ct) dx - \frac{1}{c} \int_{-\infty}^{x-ct} \psi(x - ct) dx \right] \\ \implies u(x, t) &= \frac{1}{2} [\phi(x + ct) + \phi(x - ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(x + ct) dx \end{aligned}$$