

# 统计力学第六次作业

肖涵薄 31360164

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## 7.1

$\varepsilon = \frac{1}{2m} \left( \frac{2\pi\hbar}{L} \right)^2 \sum n^2 = aL^{-2} = aV^{-2/3}$ ,  $a$  为常数  $a = \frac{(2\pi\hbar)^2 \sum n_i^2}{2m}$ .  
则  $\frac{\partial \varepsilon}{\partial V} = -\frac{2}{3}\varepsilon/V$ . 用  $l$  取代  $\sum n_i^2$ .  $\frac{\partial \varepsilon}{\partial V} = -\frac{2}{3}\varepsilon/V$   
因此  $p = -\sum a_l \frac{\partial \varepsilon_l}{\partial V} = \frac{2}{3} \sum a_l \frac{\varepsilon}{V} = \frac{2}{3} \frac{U}{V}$

## 7.2

$\varepsilon = c \frac{2\pi\hbar}{L} \sqrt{\sum n^2} = aL^{-1} = aV^{-1/3}$ ,  $a$  为常数  $a = c2\pi\hbar \sqrt{\sum n^2}$ .  
则  $\frac{\partial \varepsilon}{\partial V} = -\frac{1}{3}\varepsilon/V$ . 用  $l$  取代  $\sqrt{\sum n^2}$ .  $\frac{\partial \varepsilon}{\partial V} = -\frac{1}{3}\varepsilon/V$   
因此  $p = -\sum a_l \frac{\partial \varepsilon_l}{\partial V} = \frac{1}{3} \sum a_l \frac{\varepsilon}{V} = \frac{1}{3} \frac{U}{V}$

## 7.4

$$S = Nk(\ln Z_1 - \beta \frac{\partial}{\partial \beta} \ln Z_1)$$

总的  $S$  为各状态的概率平均:

$$\begin{aligned} &= Nk \sum_s P_s (\ln Z_1 - \beta \frac{\partial}{\partial \beta} (-\beta \varepsilon_s - \ln P_s)) \\ &= \sum_s P_s Nk (\ln Z_1 + \beta \varepsilon_s) = \sum_s P_s Nk (-\beta \varepsilon_s - \ln P_s + \beta \varepsilon_s) = - \sum_s P_s \ln P_s \end{aligned}$$

当系统为非定域,

$$\begin{aligned} S &= Nk(\ln Z_1 - \beta \frac{\partial}{\partial \beta} \ln Z_1) - k \ln N! = - \sum_s P_s \ln P_s - k \ln N! \\ &= - \sum_s P_s \ln P_s - Nk(\ln N - 1) \end{aligned}$$

## 7.5

$A$  原子共有  $Nx$  个, 因此  $c$

$$\Omega = \binom{N}{Nx} = \frac{N!}{(Nx)!(N - Nx)!}$$

又由于定域系统中  $S = k \ln \Omega$ .

$$S = k \ln \frac{N!}{(Nx)!(N - Nx)!}$$

利用  $\ln N = N(\ln N - 1)$  可化简为:

$$S = -Nk[x \ln x + (1 - x) \ln(1 - x)]$$

## 7.6

(a)

总的状态数为

$$\Omega = \binom{N}{n}^2 = \left( \frac{N!}{n!(N - n)!} \right)^2$$

又由于定域系统中  $S = k \ln \Omega$ .

$$S = 2k \ln \frac{N!}{n!(N - n)!} = 2k(N \ln N - n \ln n - (N - n) \ln(N - n))$$

(b)

平衡态时  $F$  极小要求  $\frac{\partial F}{\partial n} \rightarrow 0 \implies u = T \frac{\partial S}{\partial n}$ .

$$\frac{\partial S}{\partial n} = 2k(-\ln n + \ln(N - n)) = 2k \ln(N/n - 1) \approx 2k \ln(N/n)$$

因此

$$\frac{u}{2kT} = \ln(N/n) \implies n = N e^{-\frac{u}{2kT}}$$

## 7.7

这种情况由于只有空位没有间隙原子出现, 因此  $S$  变为 7.5 的表达式, 即

$$S = k \ln \frac{N!}{n!(N-n)!} = k(N \ln N - n \ln n - (N-n) \ln(N-n))$$

同理, 平衡态时  $F$  极小要求  $\frac{\partial F}{\partial n} \rightarrow 0 \implies w = T \frac{\partial S}{\partial n}$ .

$$\frac{\partial S}{\partial n} = k(-\ln n + \ln(N-n)) = k \ln(N/n - 1) \approx k \ln(N/n)$$

因此

$$\frac{w}{kT} = \ln(N/n) \implies n = N e^{-\frac{w}{kT}}$$

## 7.9

当气体没有整体运动时, 最概然分布为:

$$dN = e^{-\alpha - \frac{\beta}{2m}(p_x^2 + p_y^2 + p_z^2)} \frac{V dp_x dp_y dp_z}{h^3}$$

使坐标系做沿  $-z$  方向, 速度为  $v_0$  的运动, 设此时相对速度变化为  $v'_z = v_z + v_0 \iff p'_z = p_z + p_0$ , 分布为:

$$dN = e^{-\alpha - \frac{\beta}{2m}(p_x^2 + p_y^2 + p_z'^2)} \frac{V dp_x dp_y dp_z}{h^3}$$

用  $p_z$  代替  $p'_z$ ,

$$dN = e^{-\alpha - \frac{\beta}{2m}(p_x^2 + p_y^2 + (p_z - p_0)^2)} \frac{V dp_x dp_y dp_z}{h^3}$$

这等价于气体沿  $z$  方向运动.

## 7.12

根据速度分布律,

$$f(\vec{v}) = \left(\frac{m}{2\pi k_B T}\right)^{3/2} \exp\left(-\frac{mv^2/2}{k_B T}\right) = C \exp\left(-\frac{mv^2/2}{k_B T}\right)$$

则某一状态概率为

$$dW = C^2 \exp\left(-\frac{mv_1^2/2}{k_B T}\right) \exp\left(-\frac{mv_2^2/2}{k_B T}\right) d\vec{v}_1 d\vec{v}_2$$

修改积分变量  $(\vec{v}_1, \vec{v}_2) \rightarrow (\vec{v}_r, \vec{v}_e)$ . 其中  $v_e = \frac{1}{2}(v_1 + v_2)$ ,  $v_r = v_2 - v_1$  Jacobi 行列式为

$$\left| \frac{\partial(v_1, v_2)}{\partial(v_e, v_r)} \right| = \begin{vmatrix} 1 & 1 \\ -\frac{1}{2} & \frac{1}{2} \end{vmatrix} = 1$$

此时质量的含义发生变化, 用等效质量替换:  $m_e = 2m, m_\mu = m/2$ ,

$$dW = C_e C_\mu \exp\left(-\frac{m_e v_e^2/2}{k_B T}\right) \exp\left(-\frac{m_\mu v_r^2/2}{k_B T}\right) d\vec{v}_e d\vec{v}_r$$

对于相对速度  $v_r$ , 其分布为  $f(v_r) = C_\mu \exp\left(-\frac{m_\mu v_r^2/2}{k_B T}\right)$

$$\Rightarrow f(v_r) = \left(\frac{m_\mu}{2\pi k_B T}\right)^{3/2} \exp\left(-\frac{m_\mu v_r^2/2}{k_B T}\right)$$

由于对称性,  $|v_r|$  相同的各个  $v_r$  的分布是相同的, 因此  $d|v_r| = 4\pi|v_r|^2 dv_r$  速度分布为

$$f(|v_r|) d|v_r| = f(v_r) 4\pi|v_r|^2 dv_r = \left(\frac{m_\mu}{2\pi k_B T}\right)^{3/2} \exp\left(-\frac{m_\mu v_r^2/2}{k_B T}\right) 4\pi|v_r|^2 dv_r$$

$$\overline{|v|} = \int_0^\infty |v| f(|v|) d|v|$$

$$\begin{aligned} \overline{|v_r|} &= \int_0^\infty |v_r| f(|v_r|) d|v_r| = \int_0^\infty |v_r| \left(\frac{m_\mu}{2\pi k_B T}\right)^{3/2} \exp\left(-\frac{m_\mu v_r^2/2}{k_B T}\right) 4\pi|v_r|^2 dv_r \\ &= \int_0^\infty \frac{|v_r|}{\sqrt{2}} \frac{1}{2} \left(\frac{m}{2\pi k_B T}\right)^{3/2} \exp\left(-\frac{m(v_r/\sqrt{2})^2/2}{k_B T}\right) 4\pi 2 \times \left(\frac{|v_r|}{\sqrt{2}}\right)^2 \sqrt{2} d\frac{v_r}{\sqrt{2}} \end{aligned}$$

做变量替换  $\frac{v_r}{\sqrt{2}} = v$

$$= \sqrt{2} \left(\frac{m}{2\pi k_B T}\right)^{3/2} \exp\left(-\frac{mv^2/2}{k_B T}\right) 4\pi|v|^2 dv = \sqrt{2} f(|v|) d|v| = \sqrt{2} |v|$$

即  $\overline{|v_r|} = \sqrt{2} \overline{|v|}$

## 7.14

单位时间逃逸的分子数密度为:  $\Gamma = \pi n \left(\frac{m}{2\pi kT}\right)^{3/2} \exp\left(-\frac{mv^2}{2kT}\right) dv$ .

分子平均速率为

$$\bar{v} = \frac{\int_0^\infty \Gamma v dv}{\int_0^\infty \Gamma dv} = \frac{\int_0^\infty v^4 \exp\left(-\frac{mv^2}{2kT}\right) dv}{\int_0^\infty v^3 \exp\left(-\frac{mv^2}{2kT}\right) dv} = \sqrt{\frac{9\pi kT}{8m}}$$

方均根速率为

$$\sqrt{\overline{v^2}} = \sqrt{\frac{\int_0^\infty \Gamma v^2 dv}{\int_0^\infty \Gamma dv}} = \sqrt{\frac{\int_0^\infty v^5 \exp\left(-\frac{mv^2}{2kT}\right) dv}{\int_0^\infty v^3 \exp\left(-\frac{mv^2}{2kT}\right) dv}} = \sqrt{\frac{4kT}{m}}$$

平均能量为:

$$\bar{E} = \frac{1}{2} m \overline{v^2} = 2kT$$

## 7.18

简谐振动能级为:

$$\varepsilon = (n + \frac{1}{2})\hbar\omega$$

简并度为 1, 则

$$\begin{aligned} Z_1 &= \sum \exp(-\beta(n + \frac{1}{2})\hbar\omega) = \frac{e^{-\frac{1}{2}\beta\hbar\omega}}{1 - e^{-\beta\hbar\omega}} \\ \ln Z_1 &= -\frac{1}{2}\beta\hbar\omega - \ln(1 - e^{-\beta\hbar\omega}) \\ S &= Nk \left( \ln Z_1 - \beta \frac{\partial}{\partial \beta} \ln Z_1 \right) \\ &= Nk \left( -\frac{1}{2}\beta\hbar\omega - \ln(1 - e^{-\beta\hbar\omega}) - \beta \left( -\frac{1}{2}\hbar\omega - \frac{1}{1 - e^{-\beta\hbar\omega}} e^{-\beta\hbar\omega} (-\hbar\omega) \right) \right) \\ &= Nk \left( -\ln(1 - e^{-\beta\hbar\omega}) + \frac{\beta\hbar\omega}{e^{\beta\hbar\omega} - 1} \right) \end{aligned}$$

代入  $\theta = T\beta\hbar\omega$ ,

$$= Nk \left( -\ln(1 - e^{-\theta/T}) + \frac{\theta/T}{e^{\theta/T} - 1} \right)$$

## 7.21

$$\begin{aligned} Z_1 &= e^{-\beta\varepsilon_1} + e^{-\beta\varepsilon_2}, N = e^{-\alpha-\beta\varepsilon_1} + e^{-\alpha-\beta\varepsilon_2} \\ U &= -N \frac{\partial}{\partial \beta} \ln Z_1 = -N \frac{\partial}{\partial \beta} \ln(e^{-\beta\varepsilon_1} + e^{-\beta\varepsilon_2}) \\ &= -N \left( \frac{1}{e^{-\beta\varepsilon_1} + e^{-\beta\varepsilon_2}} (-\varepsilon_1 e^{-\beta\varepsilon_1} - \varepsilon_2 e^{-\beta\varepsilon_2}) \right) \\ &= N \left( \frac{1}{e^{-\beta\varepsilon_1} + e^{-\beta\varepsilon_2}} (\varepsilon_1 e^{-\beta\varepsilon_1} + \varepsilon_2 e^{-\beta\varepsilon_2}) \right) \\ &= N\varepsilon_1 + \frac{N(\varepsilon_2 - \varepsilon_1)}{1 + e^{\beta(\varepsilon_2 - \varepsilon_1)}} \\ S &= Nk \left( \ln Z_1 - \beta \frac{\partial}{\partial \beta} \ln Z_1 \right) \\ &= Nk \left( -\beta\varepsilon_1 + \ln(1 + e^{-\beta(\varepsilon_2 - \varepsilon_1)}) + \beta\varepsilon_1 + \frac{\beta(\varepsilon_2 - \varepsilon_1)}{1 + e^{\beta(\varepsilon_2 - \varepsilon_1)}} \right) \\ &= Nk \left( \ln(1 + e^{-\beta(\varepsilon_2 - \varepsilon_1)}) + \frac{\beta(\varepsilon_2 - \varepsilon_1)}{1 + e^{\beta(\varepsilon_2 - \varepsilon_1)}} \right) \end{aligned}$$