Wavelets in the Hamiltonian

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Abstract

Here a standard Hamiltonian is simplified and made diagonal in the momentum domain by using a Wavelet frame to approximate the potential.

1 Introduction

Throughout this document, $\hbar = 1$, as it should be. (q, p) are generalized position and momentum respectively.

We take Hamiltonians of the form,

$$\hat{H} = \frac{\hat{P}^2}{2m} + \hat{V}$$

where \hat{V} is diagonal in q-space. Thus we can write $\langle q|\hat{V}|q\rangle=V(q)\langle q|q\rangle$.

2 Diagonalization

We wish to find the eigenvalues of the time-independent Schrodinger equations,

$$\hat{H}|E\rangle = E|E\rangle$$

so that we can take the time dependent eigenvector as $e^{-iEt}|E\rangle$. Expanding the equation in p space and letting $|E\rangle = |p\rangle$,

$$\begin{split} \frac{1}{2m}\langle p|\hat{P}^2|E\rangle + \langle p|\hat{V}|E\rangle &= E\langle p|E\rangle \\ \frac{p^2}{2m} + \langle p|\hat{V}|p\rangle &= E \end{split}$$

The \hat{V} is not naturally diagonal in p-space, but we can use discrete Wavelet frames to solve this. Take a "snug" wavelet frame [1],

3 Wavelet Frames

Let a "snug" wavelet frame be defined by,

$$\sum_{\alpha} |\psi_{\alpha}\rangle\langle\psi_{\alpha}| \approx A\hat{\mathbf{1}}$$

for $|\psi_{\alpha}\rangle = \hat{U}(\alpha)|\psi_{0}\rangle$ where \hat{U} is the unitary representation of a group and $||\psi_{0}||^{2} = 1$.

We can use these frames to further reduce the Hamiltonian. Looking at the potential part,

$$\langle p|\hat{V}|p\rangle \approx \frac{1}{A^2} \sum_{\alpha,\beta} \langle p|\psi_{\alpha}\rangle \langle \psi_{\alpha}|\hat{V}|\psi_{\beta}\rangle \langle \psi_{\beta}|p\rangle$$
$$= \frac{1}{A^2} \sum_{\alpha,\beta} \psi_{\alpha}(p) \overline{\psi_{\beta}(p)} V_{\alpha\beta}$$

where $V_{\alpha\beta} = \langle \psi_{\alpha} | \hat{V} | \psi_{\beta} \rangle$ follows the reasonably obvious notation. We can write this as a matrix vector product,

$$\langle p|\hat{V}|p\rangle = \vec{\psi}(p)^T \mathbf{V} \vec{\psi}(p)$$

4 Finite Frames

We can get excellent approximations of the potential with just finitely many frame vectors. This has to do with the redundancy and space-momentum localization of the wavelet vectors [1].

This is done case by case and heuristically if the $V_{\alpha\beta}$ is computed numerically. Closed form solutions could just simply look at the way the closed form falls off and pick termination point based on that.

5 Simple Example

Take the step rectangle function as the potential,

$$V(x) = V_0 \mathbf{1}_{|x| \in (-x_0, x_0)}$$

We use the Mexican Hat Wavelet,

$$\langle p|\psi_0\rangle = \frac{2}{\sqrt{3}\sqrt[4]{\pi}}\langle p|\hat{P}^2|g\rangle$$

$$= \frac{2}{\sqrt{3}\sqrt[4]{\pi}}p^2e^{-p^2/2}$$

$$\langle q|\psi_0\rangle = -\frac{2}{\sqrt{3}\sqrt[4]{\pi}}\partial_q^2e^{-q^2/2}$$

$$= \frac{2}{\sqrt{3}\sqrt[4]{\pi}}(1-q^2)e^{-q^2/2}$$

which is unit length.

We use the affine group, with representation,

$$\begin{split} \langle q|\hat{U}(a,b)|f\rangle &= f(\frac{q-b}{a})/\sqrt{|a|}\\ \langle p|\hat{U}(a,b)|f\rangle &= \tilde{f}(ap)e^{-ibp}\sqrt{|a|} \end{split}$$

References

 $[1]\,$ Ingrid Daubechies et al. Ten lectures on wavelets, volume 61. SIAM, 1992.