

PHYS417HW6

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This one was also a bit of a weird one.

1)

Given the publicly known key $P = 997$

a)

$N = 231$, so X is

$$\begin{aligned} X &= 2^{231} \bmod(997) \\ &= 72 \end{aligned}$$

b)

With $M = 337$, we have

$$\begin{aligned} Y &= 2^{373} \bmod(997) \\ &= 80 \end{aligned}$$

c)

I’ll be real, I calculated the previous mod functions using Desmos. Desmos is blowing up when I try to calculate K_1 or K_2 .

d)

To convert to binary, we will use 10 bits, and essentially flip 0’s to 1’s from left to right such that we don’t get any bigger sum than our target number. I unfortunately do not have the number, but it shouldn’t have been that bad to calculate.

2)

Given $\hat{\pi} = \hat{p} + e\hat{A}$ and $\hat{\tilde{\pi}} = \hat{p} - e\hat{A}$

And $\hat{a} = \frac{\ell_B}{\sqrt{2}\hbar} (\pi_x - i\pi_y)$ and $\hat{b} = \frac{\ell_B}{\sqrt{2}\hbar} (\tilde{\pi}_x - i\tilde{\pi}_y)$

$\vec{A} = \frac{B}{2} [-y\vec{e}_x + x\vec{e}_y]$

So;

$$\pi_x = p_x - \frac{eB}{2}y$$

$$\pi_y = p_y + \frac{eB}{2}x$$

$$\tilde{\pi}_x = p_x + \frac{eB}{2}y$$

$$\tilde{\pi}_y = p_y - \frac{eB}{2}x$$

a)

$$L_z = xp_y - yp_x$$

Rewriting the π operators in terms of the more familiar position and momentum (really just adding things together), we get

$$p_x = \frac{\pi_x + \tilde{\pi}_x}{2}$$

$$x = \frac{\pi_y - \tilde{\pi}_y}{eB}$$

$$p_y = \frac{\pi_y + \tilde{\pi}_y}{2}$$

$$y = \frac{-\pi_x + \tilde{\pi}_x}{eB}$$

Which means that

$$\begin{aligned}
L_z &= \left(\frac{\pi_y - \tilde{\pi}_y}{eB} \right) \left(\frac{\pi_y + \tilde{\pi}_y}{2} \right) - \left(\frac{-\pi_x + \tilde{\pi}_x}{eB} \right) \left(\frac{\pi_x + \tilde{\pi}_x}{2} \right) \\
&= \left(\frac{\pi_y^2 - \tilde{\pi}_y^2}{2eB} \right) - \left(\frac{-\pi_x^2 + \tilde{\pi}_x^2}{2eB} \right) \\
&= \frac{\pi_y^2 - \tilde{\pi}_y^2 + \pi_x^2 - \tilde{\pi}_x^2}{2eB} \\
&= \hbar (a^\dagger a - b^\dagger b)
\end{aligned}$$

When we throw this at a state $|n, m\rangle$, we will get

$$\hbar (a^\dagger a - b^\dagger b) |n, m\rangle = \hbar(n - m)|n, m\rangle$$

and so the expectation value is simply

$$\langle n, m | L_z | n, m \rangle = \hbar(n - m)$$

b)

We can do a similar sort of trick to calculate K_z .

If $K_z = (\vec{r} \times \vec{\pi})_z = x\pi_y - y\pi_x$, which has the same form as $L_z = (\vec{r} \times \vec{p})_z$. So then,

$$\begin{aligned}
K_z &= x\pi_y - y\pi_x \\
&= \pi_y \left(\frac{\pi_y - \tilde{\pi}_y}{eB} \right) - \pi_x \left(\frac{-\pi_x + \tilde{\pi}_x}{eB} \right) \\
&= \frac{\pi_y^2 - \pi_y \tilde{\pi}_y + \pi_x^2 - \pi_x \tilde{\pi}_x}{eB} \\
&= 2\hbar (a^\dagger a) + \frac{-\pi_x \tilde{\pi}_x - \pi_y \tilde{\pi}_y}{eB} \\
&= \hbar (2a^\dagger a - ab - a^\dagger b^\dagger)
\end{aligned}$$

If we throw this at a state, we get

$$\hbar (2a^\dagger a - ab - a^\dagger b^\dagger) |n, m\rangle = \hbar \left(2n|n, m\rangle + \sqrt{nm}|n-1, m-1\rangle + \sqrt{(n+1)(m+1)}|n+1, m+1\rangle \right)$$

The latter terms of which will be orthogonal with the state $|n, m\rangle$. The expectation value is thus

$$\langle n, m | K_z | n, m \rangle = 2\hbar n$$

c)

For coherent states α, β we can plug it into L_z and K_z and all that will really be different is the coefficients. Because of how the interaction between our operators and the coherent states are defined, we will get

$$\hbar (a^\dagger a - b^\dagger b) |\alpha\rangle |\beta\rangle = \hbar (|\alpha|^2 |\alpha\rangle - |\beta|^2 |\beta\rangle)$$

$$\langle \alpha, \beta | L_z | \alpha, \beta \rangle = \hbar (|\alpha|^2 - |\beta|^2)$$

and

$$\hbar (2a^\dagger a - ab - a^\dagger b^\dagger) |\alpha\rangle |\beta\rangle = \hbar (2|\alpha|^2 |\alpha\rangle |\beta\rangle - 2\alpha\beta |\alpha\rangle |\beta\rangle)$$

$$\langle \alpha, \beta | K_z | \alpha, \beta \rangle = 2\hbar (|\alpha|^2 - \alpha\beta)$$

I'll be real, not sure of the physical significance of this result.