# Assignment 3

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#### TASK GIVEN

- 1. Extend your unit cell by applying multiple of them from the previous assignment. Construct a bulk band structure  $\omega(k_B)$
- 2. Given that you construct a scattering matrix from the previous problem, terminate it from the endpoints. It is called a Ribbon problem.

#### SOLUTION 1

The easiest way I can think for solving such problem is to apply a recursive functions. Initially, I'll make a scattering matrix for (n-1) unit cell structure and combine it with another unit cell. Finally, I should get a nxn matrix. See Fig.

$$\begin{pmatrix} b_1 \\ b_2 \\ \vdots \\ b_{2n} \end{pmatrix} = M^{n-1} \begin{pmatrix} a_1 \\ a_2 \\ \vdots \\ a_{2n} \end{pmatrix}$$
 (0.1)

While for the another unit cell

$$\begin{pmatrix}
B_1 \\
B_2 \\
B_3 \\
B_4
\end{pmatrix} = S \begin{pmatrix}
A_1 \\
A_2 \\
A_3 \\
A_4
\end{pmatrix}$$
(0.2)

At the interface of these structures, we have  $B_1 = a_{2n}$  and  $A_1 = b_{2n}$ . By combining equations (0.1) and (0.2) we can now get (2n+2)x(2n+1) matrix.

$$\begin{pmatrix} b_1 \\ b_2 \\ \vdots \\ b_{2N+2} \end{pmatrix} = M^N \begin{pmatrix} a_1 \\ a_2 \\ \vdots \\ a_{2N+2} \end{pmatrix}$$

$$(0.3)$$

For simplicity, let us use notation  $M=M^n$ The Bloch matrix  $\tilde{B}$  can be found using relations

$$b_{2n+1} = a_{2n}e^{i\mathbf{k_b}\cdot\mathbf{a}}$$
$$b_{2n} = a_{2n+1}e^{-i\mathbf{k_b}\cdot\mathbf{a}}$$

where n varies from 1 to N. At the lower and upper boundaries

$$b_{2N+2} = a_1 e^{iN\mathbf{k_b} \cdot \mathbf{b}}$$
$$b_1 = a_{2N+2} e^{-iN\mathbf{k_b} \cdot \mathbf{b}}$$

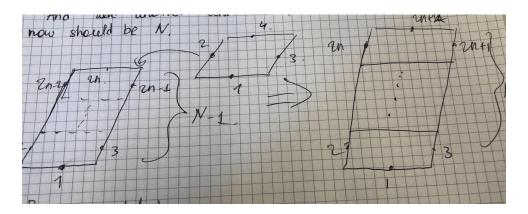


Figure 0.1: Our recursive strategy.

The bulk band structure, namely  $-\log |\det(M-\tilde{B})|$  represented in the figure below

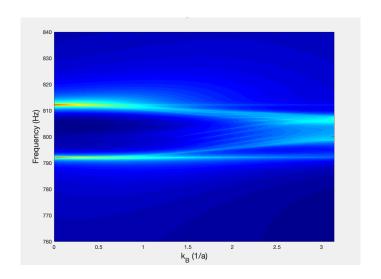


Figure 0.2: Bulk band structure for no bias velocity.

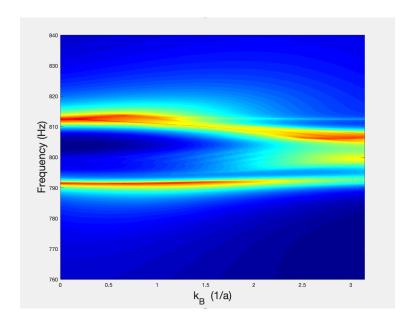


Figure 0.3: Bulk band structure for  $\frac{v}{D\gamma}$  = 0.2.

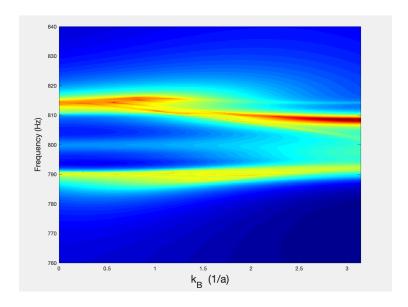


Figure 0.4: Bulk band structure for  $\frac{v}{D\gamma} = 0.5$ .

### SOLUTION 2

Here in Ribbon problem, let's say that we have a same unit cell and we just need to equate lower and upper bands. To be exact

$$\begin{pmatrix}
b_{1'} \\
b_{1} \\
\vdots \\
b_{2N} \\
b_{2'}
\end{pmatrix} = M \begin{pmatrix}
a_{1'} \\
a_{1} \\
\vdots \\
a_{2N} \\
a_{2'}
\end{pmatrix}$$
(0.4)

By using the relations  $b_{1'} = a_{1'}$  and  $b_{2'} = a_{2'}$ , we should transfer it to 2Nx2N matrix

$$\begin{pmatrix} b_1 \\ \vdots \\ b_{2N} \end{pmatrix} = T \begin{pmatrix} a_1 \\ \vdots \\ a_{2N} \end{pmatrix}$$
 (0.5)

After using "simple" calculations, we can actually find the coefficients of *T* matrix

$$T_{ij} = M_{ij} + \frac{M_{i2'}[M_{2'j}(1-M_{1'1'}) + M_{2'1'}M_{1'j}] + M_{i1'}[M_{1'j}(1-M_{2'2'}) + M_{1'2'}M_{2'j}]}{(1-M_{1'1'})(1-M_{2'2'}) + M_{1'2'}M_{2'1'}} \tag{0.6}$$

Below, bulk band structure for Ribbon problem are represented

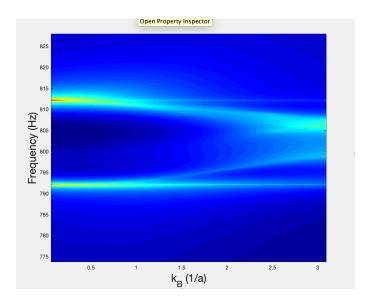


Figure 0.5: Bulk band structure for (Ribbon problem) for no bias velocity.

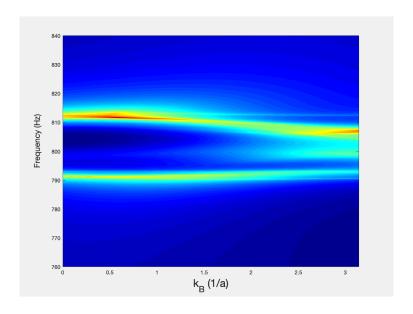


Figure 0.6: Bulk band structure for (Ribbon problem) for  $\frac{v}{D\gamma}=0.2$ .

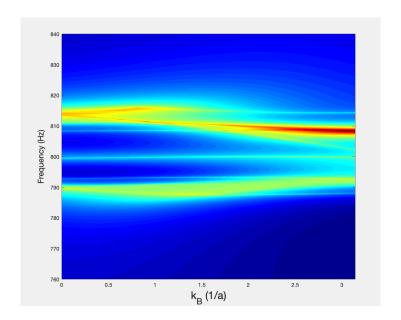


Figure 0.7: Bulk band structure for (Ribbon problem) for  $\frac{v}{D\gamma}=0.5$ .

## REFERENCES

- [1] Romain Fleury, MS. *Breaking Temporal Symmetries in Metamaterials and Metasurfaces*, August 2015.
- [2] Charles Kittel. Introduction to Solid State Physics, 1914.
- [3] Alexander B. Khanikaev, Romain Fleury, S. Hossein Mousavi and Andrea Alu. *Topologically robust sound propagation in an angular-momentum-biased graphene-like resonator lattice*, Nat. Comm. 2015