ACM114 - Parallel Algorithms for Scientific Applications

Parallel implementation of the Boundary Element Method for the solution of the Laplace equation in unbounded domains

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Setup of the problem

We want to solve the Laplace equation in a general unbounded domain:

Find
$$u: \mathbb{R}^d \setminus \overline{\Omega} \to \mathbb{R}$$
, $u \in H^1_{\mathrm{loc}}(\mathbb{R}^d \setminus \overline{\Omega})$ $(d=2,3)$, such that:

:
$$(P) \left\{ \begin{array}{lll} -\Delta u &=& 0 & \text{in} & \mathbb{R}^d \setminus \overline{\Omega}, \\ \partial_{\boldsymbol{n}} u = q, & \text{or} & u = p & \text{on} & \Gamma, \\ + \text{ decay condition as } |\boldsymbol{x}| \to \infty \end{array} \right.$$



Applications

- Electrostatic and magnetostatic problems.
- Potential flow in fluid mechanics problems.
- Particle simulations (many body problem).
- The domain of definition of the solution is unbounded.
- Direct (naive) implementation of the Finite Differences or Finite Element method does not work in this case.
- We want to solve the problem for a general geometry.

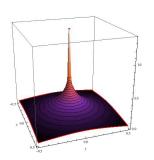
Boundary integral equation method

First we compute the Green's function of the problem:

$$G(\boldsymbol{x}, \boldsymbol{y}) = \begin{cases} -\frac{1}{2\pi} \log(|\boldsymbol{x} - \boldsymbol{y}|) & \text{if} \quad d = 2, \\ \\ \frac{1}{4\pi |\boldsymbol{x} - \boldsymbol{y}|} & \text{if} \quad d = 3, \end{cases}$$

solution of the following problem (in the sense of the distributions $\mathcal{D}'(\mathbb{R}^d)$):

$$-\Delta_{\boldsymbol{x}}G(\boldsymbol{x},\boldsymbol{y}) = \delta_{\boldsymbol{y}}(\boldsymbol{x}), \qquad \boldsymbol{x}, \boldsymbol{y} \in \mathbb{R}^d.$$



3 / 7

Integral representation & boundary integral equation

Using the second Green's identity we get:

$$\mathsf{BIR:} \quad u(\boldsymbol{x}) \quad = \quad \int_{\Gamma} \left\{ G(\boldsymbol{x}, \boldsymbol{y}) \boldsymbol{q}(\boldsymbol{y}) - \frac{\partial G}{\partial n_{\boldsymbol{y}}}(\boldsymbol{x}, \boldsymbol{y}) \boldsymbol{p}(\boldsymbol{y}) \right\} \, \mathrm{d}\sigma_{\boldsymbol{y}}, \quad \boldsymbol{x} \in \mathbb{R}^d \setminus \overline{\Omega},$$

$$\mathsf{BIE:} \quad \frac{p(\boldsymbol{x})}{2} \quad = \quad \int_{\Gamma} \left\{ G(\boldsymbol{x}, \boldsymbol{y}) q(\boldsymbol{y}) - \frac{\partial G}{\partial n_{\boldsymbol{y}}} (\boldsymbol{x}, \boldsymbol{y}) p(\boldsymbol{y}) \right\} \, \mathrm{d}\sigma_{\boldsymbol{y}}, \quad \boldsymbol{x} \in \Gamma.$$

Boundary element method

We discretize the boundary integral equation (BIE) by seeking a solution of the form

$$p(\boldsymbol{x}) pprox \sum_{j=1}^{N} p_j \mathbf{1}_{\Gamma_j}(\boldsymbol{x})$$

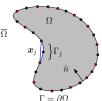
which leads to the following set of equations:

$$\frac{p_i}{2} = \sum_{i=1}^{N} \{g_{ij}q_j - f_{ij}p_j\}, \quad \forall i = 1, \dots, N$$

where

$$f(g_{ij} = \int_{\Gamma_i} G(\boldsymbol{x}_i, \boldsymbol{y}) \, \mathrm{d}\sigma_{\boldsymbol{y}}, \quad f_{ij} = \int_{\Gamma_i} \frac{\partial G}{\partial n_{\boldsymbol{y}}}(\boldsymbol{x}_i, \boldsymbol{y}) \, \mathrm{d}\sigma_{\boldsymbol{y}}.$$

 Γ_j : boundary elements x_j : boundary nodes



4 / 7

The approximate values of the solution or its normal derivative can be computed by solving the linear system

$$\left(\frac{\mathbf{1}}{2} + \mathbf{F}\right)\mathbf{p} = \mathbf{G}\,\mathbf{q}.$$

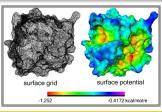
In other words, we have the linear system:

Dirichlet: $\mathbf{A} \mathbf{q} = \mathbf{b}$, $\mathbf{A} = \mathbf{G}$, $\mathbf{b} = (\mathbf{I}/2 + \mathbf{F}) \mathbf{p}$ Neumann: $\mathbf{A} \mathbf{p} = \mathbf{b}$, $\mathbf{A} = (\mathbf{I}/2 + \mathbf{F})$, $\mathbf{b} = \mathbf{G} \mathbf{q}$

Parallelization

Why parallelize?

- High-resolution meshes or large number of objects (e.g. galaxies or molecules) requires a large number of elements.
- ullet The computation of the matrix entries requires at least $\mathcal{O}(N^2)$ operations.
- Unlike FEM or FDM, the system matrix arising in BEM is nonsymetric and dense, so standard fast algorithms (e.g. cyclic reduction methods) do not work in this case.
- For large N the numerical solution of the resulting linear system may be infeasible expensive even for Krylov subspace methods, which require only $\mathcal{O}(N^2)$ operations to converge for a general well-conditioned matrix.
- \bullet The evaluation of the numerical solution at a single arbitrary point outside the boundary requires $\mathcal{O}(N^2)$ as well.



Parallel implementation

- $oldsymbol{0}$ The mesh data is made available to all p processors.
- ② Each processor computes and store $\delta w = N/p$ rows of ${\bf F}$ and ${\bf G}$. It allows to perform matrix-vector multiplications in parallel by computing the dot product between the matrix rows and the vector. This operation takes $\mathcal{O}(N^2/p)$ time units.
- The boundary condition vector is distributed to all the processors and then ${\bf b}$ is computed. Each processor computes δw entries of ${\bf b}$ and send them to the master process. This operation requires $\mathcal{O}(N^2/p)$ time units.
- **●** The linear system is solved using the GMRES method. It only requires matrix-vector products of \mathbf{A} . Each GMRES iteration requires $\mathcal{O}(N^2/p)$ time units. If \mathbf{A} is well-conditioned, the solution of the linear system takes $\mathcal{O}(N^2/p)$ time units.
- ① The integral representation is evaluated at N_p points. Each processor computes u at $\delta w_p = N_p/p$ different points. Subsequently, these values are sent to the master process where a vector stores u for visualization purposes.

Benchmark problem

- We consider the function $p(r,\theta) = \cos \theta/r$, which is analytical and satisfies $-\Delta p = 0$ for all $x = r(\cos \theta, \sin \theta) \in \mathbb{R}^2 \setminus \{(0,0)\}.$
- We solve the Laplace equation for the condition u=p on $\Gamma=\partial\Omega$. As result we recover $q=\partial u/\partial n$ on the boundary and the values of u at points $x\in\mathbb{R}^2\setminus\overline{\Omega}$ by using the integral representation formula.
- For N=60 the error obtained for the normal derivative is $\max_{x\in\Gamma}|q-\partial u/\partial n|=1$ 94F-03

