

Radiative Corrections

We are now (finally) ready to compute the first-order radiative correction to the magnetic moment of the electron

$$a_e^{(1)} = F_2^{(0)}$$

$$= \frac{1}{12m^2} t_c \left[\left(m^2 \gamma_r - \frac{P_r P}{2} - \frac{3}{2} m P_r \right) V_1^r \right. \\ \left. + \frac{m}{4} \left(\frac{P}{2} + m \right) [\gamma_r, \gamma_v] \left(\frac{P}{2} + m \right) \delta V_1^v \right]$$

where $V_1^r(P) = \Lambda_1^r(P, 0)$

$$\delta V_1^v(P) = \left. \frac{\partial \Lambda_1^r(P, q)}{\partial \Sigma_v} \right|_{q=0}$$

where $\Lambda_1^r(P, q) \equiv \Lambda_1^r(p', p)$ is the correction to the vertex function

$$\Lambda_1^r(p', p) = -e^2 \int \frac{d^4 k}{(2\pi)^4} \frac{-i^3 N^r(p', p, k)}{k^2 [(p' - k)^2 - m^2][(p - k)^2 - m^2]}$$

$$\text{with } N^r = \gamma_v [(\not{p} - \not{k}) + m] \gamma^r [(\not{p} - \not{k}) + m] \gamma^v$$

Let us now replace the kinematics

$$p' = \frac{1}{2}(P+q), \quad p = \frac{1}{2}(P-q)$$

S_1

$$\Lambda_1^{\mu}(P, q) = -ie^2 \int \frac{d^4 k}{(2\pi)^4} \frac{N^{\mu}(P, q, k)}{k^2 [((P+q)/2 - k)^2 - m^2] [((P-q)/2 - k)^2 - m^2]}$$

with

$$N^{\mu}(P, q, k) = \gamma_{\alpha} [(P+q)/2 - k + n] \gamma^{\mu} [(P-q)/2 - k + n] \gamma^{\alpha}$$

At $q=0$,

$$\Lambda_1^{\mu}(P, 0) = -ie^2 \int \frac{d^4 k}{(2\pi)^4} \frac{N^{\mu}(P, 0, k)}{k^2 [(P/2 - k)^2 - m^2]^2}$$

$$\text{with } N^{\mu}(P, 0, k) = \gamma_{\alpha} [P/2 - k + n] \gamma^{\mu} [P/2 - k + n] \gamma^{\alpha}$$

The derivative is given by

$$\begin{aligned} \left. \frac{\partial \Lambda_1^{\mu}}{\partial q_{\nu}} \right|_{q=0} &= -ie^2 \int \frac{d^4 k}{(2\pi)^4} \left\{ \frac{1}{k^2 [((P/2 - k)^2 - m^2)]^2} \frac{\partial N^{\mu}}{\partial q_{\nu}} \Big|_{q=0} \right. \\ &\quad + \left. \frac{N^{\mu}(P, 0, k)}{k^2} \times \frac{\partial}{\partial q_{\nu}} \left(\frac{1}{[((P+q)/2 - k)^2 - m^2] [((P-q)/2 - k)^2 - m^2]} \right) \right\} \end{aligned}$$

$$\begin{aligned} \text{Now, } \left. \frac{\partial N^{\mu}}{\partial q_{\nu}} \right|_{q=0} &= \gamma_{\alpha} \frac{\partial}{\partial q_{\nu}} [\gamma^{\mu} [(P+q)/2 - k + n]] \gamma^{\nu} [\gamma^{\mu} [(P-q)/2 - k + n]] \gamma^{\alpha} \\ &\quad + \gamma_{\alpha} [\gamma^{\mu} [(P/2 - k + n)] \gamma^{\nu} \frac{\partial}{\partial q_{\nu}} [\gamma^{\mu} [(P-q)/2 - k + n]]] \gamma^{\alpha} \end{aligned}$$

$$\text{with } \frac{\partial}{\partial q_{\nu}} [\gamma^{\mu} [(P+q)/2 - k + n]] = \frac{\gamma^{\mu}}{2} \frac{\partial}{\partial q_{\nu}} q_{\nu} = \frac{\gamma^{\mu}}{2}$$

$$\Rightarrow \left. \frac{\partial N^{\mu}}{\partial q_{\nu}} \right|_{q=0} = \gamma_{\alpha} \frac{\gamma^{\nu}}{2} \gamma^{\mu} [\gamma^{\nu} [\gamma^{\mu} [(P/2 - k + n)] \gamma^{\alpha} - \gamma_{\alpha} [\gamma^{\mu} [(P-q)/2 - k + n]] \gamma^{\nu} \frac{\gamma^{\alpha}}{2}]]$$

The derivative of the propagators is

$$\frac{\partial}{\partial \xi_\nu} \left(\frac{1}{[(P+q)_{1/2}-\mu]^2 - r^2} \frac{1}{[(P-q)_{1/2}-\mu]^2 - r^2} \right)_{\xi=0}$$

$$= \left[\frac{1}{[(P_{1/2}-\mu)^2 - r^2]} \frac{\partial}{\partial \xi_\nu} \left(\frac{1}{[(C(P+q))_{1/2}-\mu]^2 - r^2} \right)_{\xi=0} \right. \\ \left. + \frac{1}{[(C(P_{1/2}-\mu))^2 - r^2]} \frac{\partial}{\partial \xi_\nu} \left(\frac{1}{[(C(P-q))_{1/2}-\mu]^2 - r^2} \right)_{\xi=0} \right]$$

Now,

$$\frac{\partial}{\partial \xi_\nu} \left(\frac{1}{[(C(P\pm q))_{1/2}-\mu]^2 - r^2} \right)_{\xi=0} = \frac{1}{[(C(P_{1/2}-\mu))^2 - r^2]^2} \cdot \frac{\partial}{\partial \xi_\nu} ((P\pm q)_{1/2}-\mu)^2 \\ = \frac{1}{[(C(P_{1/2}-\mu))^2 - r^2]^2} \cdot 2(C(P_{1/2}-\mu))^{\alpha} \cdot \frac{1}{2} (\pm \delta_{\alpha}^{\nu}) \\ = \frac{\pm (C(P_{1/2}-\mu))^{\nu}}{[(C(P_{1/2}-\mu))^2 - r^2]^2}$$

↑ sign means that these terms
will sum to zero

$$\Rightarrow \frac{\partial}{\partial \xi_\nu} \left(\frac{1}{[(P+q)_{1/2}-\mu]^2 - r^2} \frac{1}{[(P-q)_{1/2}-\mu]^2 - r^2} \right)_{\xi=0} = 0 .$$

So, the derivative of the vertex function is

$$\frac{\partial \Lambda_i}{\partial q_\nu} \Big|_{q=0} = -ie^2 \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2 [(P_{12}-k)^2 - m^2]} \gamma^\mu \frac{\partial N^\mu}{\partial q_\nu} \Big|_{q=0}$$

with,

$$\frac{\partial N^\mu}{\partial q_\nu} \Big|_{q=0} = \gamma_\alpha \frac{\gamma^\nu}{2} \gamma^\mu [P_{12}-k+m] \gamma^\alpha - \gamma_\alpha [P_{12}-k+m] \gamma^\nu \frac{\gamma^\mu}{2} \gamma^\alpha$$

In the projection formula, we need to take traces.

For the first term, we need

$$T_1 = \text{tr} \left[\left(\underline{m^2 \gamma_\nu} - \underline{P_\nu P^\nu} - \frac{3}{2} \underline{m P_\nu} \right) N^\nu(P, 0, m) \right]$$

with

$$N^\nu(P, 0, m) = \gamma_\alpha [P_{12}-k+m] \gamma^\nu [P_{12}-k+m] \gamma^\alpha$$

so, the following components are

$$\gamma_\nu \gamma_\alpha \gamma^\alpha = 2 \gamma_\nu$$

$$\begin{aligned} & \cdot \text{tr} [\gamma_\nu \gamma_\alpha [P_{12}-k+m] \gamma^\nu [P_{12}-k+m] \gamma^\alpha] \\ &= -2 \text{tr} [\gamma_\nu [P_{12}-k+m] \gamma^\nu [P_{12}-k+m]] \\ &= -2 \text{tr} [\gamma_\nu (P_{12}-k) \gamma^\nu (P_{12}-k)] - 2m^2 \text{tr} [\gamma_\nu \gamma^\nu] \\ &= +4 \text{tr} [(P_{12}-k)(P_{12}-k)] - 8m^2 \text{tr} [\mathbb{1}] \\ &= 16(P_{12}-k)^2 - 32m^2 \\ &= 16 \left(\frac{P^2}{4} + k^2 - P \cdot k \right) - 32m^2 \\ &= 16(k^2 - P \cdot k - m^2) \end{aligned}$$

$$\begin{aligned}
& \cdot \operatorname{tr}[\mathcal{P} \gamma_\alpha [\mathcal{P}_{12} - k + n] \mathcal{P} [\mathcal{P}_{12} - k + n] \gamma^\alpha] \\
&= -2 \operatorname{tr}[\mathcal{P} [\mathcal{P}_{12} - k + n] \mathcal{P} [\mathcal{P}_{12} - k + n]] \\
&= -2 \operatorname{tr}[\mathcal{P} (\mathcal{P}_{12} - k) \mathcal{P} (\mathcal{P}_{12} - k)] - 2n^2 \operatorname{tr}[\mathcal{P} \mathcal{P}] \\
&= -8 \left[2 \mathcal{P} \cdot \left(\frac{\mathcal{P}}{2} - k \right) \mathcal{P} \cdot \left(\frac{\mathcal{P}}{2} - k \right) - \mathcal{P}^2 \left(\frac{\mathcal{P}}{2} - k \right)^2 \right] - 8n^2 \mathcal{P}^2 \\
&= -8 \left[2 \left(\frac{\mathcal{P}^2}{2} - k \cdot \mathcal{P} \right)^2 - 4n^2 \left(\frac{\mathcal{P}^2}{4} - k \cdot \mathcal{P} + k^2 \right) \right] - 32n^4 \\
&= -16 \left[(2n^2 - k \cdot \mathcal{P})^2 - 2n^2 (k^2 - k \cdot \mathcal{P} + k^2) \right] - 32n^4 \\
&= -16 \left[4n^4 + (k \cdot \mathcal{P})^2 - 4n^2 k \cdot \mathcal{P} - 2n^4 + 2n^2 k \cdot \mathcal{P} - 2n^2 k^2 + 2n^4 \right] \\
&= -16 \left[4n^4 + (k \cdot \mathcal{P})^2 - 2n^2 k \cdot \mathcal{P} - 2n^2 k^2 \right]
\end{aligned}$$

$$\begin{aligned}
& \cdot \operatorname{tr}[\gamma_\alpha [\mathcal{P}_{12} - k + n] \mathcal{P} [\mathcal{P}_{12} - k + n] \gamma^\alpha] \\
&= 4 \operatorname{tr}[(\mathcal{P}_{12} - k + n) \mathcal{P} (\mathcal{P}_{12} - k + n)] \\
&= 4n \operatorname{tr}[(\mathcal{P}_{12} - k) \mathcal{P}] + 4n \operatorname{tr}[\mathcal{P} (\mathcal{P}_{12} - k)] \\
&= 32n (\mathcal{P}_{12}^2 - k \cdot \mathcal{P}) \\
&= 32n (2n^2 - k \cdot \mathcal{P})
\end{aligned}$$

So, the first trace is

$$\begin{aligned}
 T_1 &= m^2 \cdot 16(h^2 - P \cdot h - m^2) + 16(4m^4 - 2h^2 P \cdot h - 2h^2 m^2 + (P \cdot h)^2) \\
 &\quad - \frac{3}{2}m \cdot 32m(2m^2 - h \cdot P) \\
 &= -16m^2 h^2 + 16(P \cdot h)^2 - 48h^4 \\
 &= -16[3h^4 + m^2 h^2 - (h \cdot P)^2]
 \end{aligned}$$

The second trace needed is

$$T_2 = \text{tr}\left[\left(\frac{P}{2} + m\right)[\gamma_\mu, \gamma_\nu]\left(\frac{P}{2} + m\right)\frac{\partial N^\nu}{\partial g_\nu} \Big|_{g=0}\right]$$

with

$$\begin{aligned}
 \frac{\partial N^\nu}{\partial g_\nu} \Big|_{g=0} &= \frac{1}{2} \left\{ \gamma_\alpha \gamma^\nu \gamma^\mu (\not{P}_{1/2} - \not{k} + m) \gamma^\alpha \right. \\
 &\quad \left. - \gamma_\alpha (\not{P}_{1/2} - \not{k} + m) \gamma^\mu \gamma^\nu \gamma^\alpha \right\} \\
 &= \frac{1}{2} \left\{ -2 \left(\frac{P}{2} - k\right) \gamma^\mu \gamma^\nu + 4m g^{\mu\nu} \right. \\
 &\quad \left. + 2 \gamma^\nu \gamma^\mu \left(\frac{P}{2} - k\right) + 4m g^{\mu\nu} \right\}
 \end{aligned}$$

Note,
 $[\gamma_\mu, \gamma_\nu] g^{\mu\nu} = [\gamma_\mu, \gamma^\nu] = 0$

$$\begin{aligned}
 \text{So, } T_2 &= -\text{tr}\left[\left(\frac{P}{2} + m\right) \underline{[\gamma_\mu, \gamma_\nu]} \left(\frac{P}{2} + m\right) \left(\frac{P}{2} - k\right) \gamma^\mu \gamma^\nu\right] \\
 &\quad + \text{tr}\left[\left(\frac{P}{2} + m\right) \underline{[\gamma_\mu, \gamma_\nu]} \left(\frac{P}{2} + m\right) \gamma^\nu \gamma^\mu \left(\frac{P}{2} - k\right)\right]
 \end{aligned}$$

$$\begin{aligned}
& \bullet \operatorname{tr} \left[\left(\frac{P}{2} + m \right) (Y_r Y_v - Y_v Y_r) \left(\frac{P}{2} + m \right) \left(\frac{P}{2} - k \right) Y^m Y^v \right] \\
& = m \operatorname{tr} \left[\frac{P}{2} (Y_r Y_v - Y_v Y_r) \left(\frac{P}{2} - k \right) Y^m Y^v \right] \\
& \quad + m \operatorname{tr} \left[(Y_r Y_v - Y_v Y_r) \frac{P}{2} \left(\frac{P}{2} - k \right) Y^m Y^v \right] \\
& = m \operatorname{tr} \left[\frac{P}{2} Y_r Y_v \left(\frac{P}{2} - k \right) Y^m Y^v \right] - m \operatorname{tr} \left[\frac{P}{2} Y_v Y_r \left(\frac{P}{2} - k \right) Y^m Y^v \right] \\
& \quad + m \operatorname{tr} \left[Y_r Y_v \frac{P}{2} \left(\frac{P}{2} - k \right) Y^m Y^v \right] - m \operatorname{tr} \left[Y_v Y_r \frac{P}{2} \left(\frac{P}{2} - k \right) Y^m Y^v \right] \\
& = 4m \operatorname{tr} \left[\frac{P}{2} \left(\frac{P}{2} - k \right) \right] + 2m \operatorname{tr} \left[\frac{P}{2} Y_v \left(\frac{P}{2} - k \right) Y^v \right] \\
& \quad - 8m \operatorname{tr} \left[\frac{P}{2} \left(\frac{P}{2} - k \right) \right] - 16m \operatorname{tr} \left[\frac{P}{2} \left(\frac{P}{2} - k \right) \right] \\
& = -20m \operatorname{tr} \left[\frac{P}{2} \left(\frac{P}{2} - k \right) \right] - 4m \operatorname{tr} \left[\frac{P}{2} \left(\frac{P}{2} - k \right) \right] \\
& = -24 \cdot 4m \frac{P}{2} \cdot \left(\frac{P}{2} - k \right) \\
& = -96m \left(\frac{P^2}{4} - k \cdot \frac{P}{2} \right) \\
& = -48m (2m^2 - k \cdot P)
\end{aligned}$$

$$\begin{aligned}
& \cdot \operatorname{tr} \left[\left(\frac{P}{2} + h \right) \left(Y_u Y_v - Y_v Y_u \right) \left(\frac{P}{2} + h \right) Y^u Y^v \left(\frac{P}{2} - h \right) \right] \\
& = m \operatorname{tr} \left[\frac{P}{2} \left(Y_u Y_v - Y_v Y_u \right) Y^u Y^v \left(\frac{P}{2} - h \right) \right] \\
& \quad + m \operatorname{tr} \left[\left(Y_u Y_v - Y_v Y_u \right) \frac{P}{2} Y^u Y^v \left(\frac{P}{2} - h \right) \right] \\
& = m \operatorname{tr} \left[\frac{P}{2} Y_u Y_v Y^u Y^v \left(\frac{P}{2} - h \right) \right] - m \operatorname{tr} \left[\frac{P}{2} Y_v Y_u Y^u Y^v \left(\frac{P}{2} - h \right) \right] \\
& \quad + m \operatorname{tr} \left[Y_u Y_v \frac{P}{2} Y^u Y^v \left(\frac{P}{2} - h \right) \right] - m \operatorname{tr} \left[Y_v Y_u \frac{P}{2} Y^u Y^v \left(\frac{P}{2} - h \right) \right] \\
& = 16m \operatorname{tr} \left[\frac{P}{2} \left(\frac{P}{2} - h \right) \right] + 8m \operatorname{tr} \left[\frac{P}{2} \left(\frac{P}{2} - h \right) \right] \\
& \quad + 4m \operatorname{tr} \left[\frac{P}{2} \left(\frac{P}{2} - h \right) \right] - 4m \operatorname{tr} \left[\frac{P}{2} \left(\frac{P}{2} - h \right) \right] \\
& = 24 \cdot 4m \frac{P}{2} \cdot \left(\frac{P}{2} - h \right) \\
& = 96m \left(\frac{P^2}{4} - \frac{Ph}{2} \right) \\
& = 48m \left(2m^2 - h \cdot P \right)
\end{aligned}$$

$$\text{So, } \tau_2 = 96m(2m^2 - h \cdot P)$$

Therefore, we find for the anomaly

$$\begin{aligned} \tilde{a}_e^{(1)} &= \frac{1}{12m^2} \text{tr} \left[\left(m^2 \gamma_\mu - \not{P} - \frac{3}{2} m \not{P}_\mu \right) V_1^\mu \right. \\ &\quad \left. + \frac{m}{4} \left(\frac{\not{P}}{2} + m \right) [\gamma_\mu, \gamma_\nu] \left(\frac{\not{P}}{2} + m \right) \delta V_1^{\nu\mu} \right] \end{aligned}$$

$$= -i \frac{e^2}{12m^2} \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2 [(P/2-k)^2 - m^2]} \cdot \left[\tau_1 + \frac{m}{4} \tau_2 \right]$$

Now,

$$\begin{aligned} \tau_1 + \frac{m}{4} \tau_2 &= -16 [3m^4 + m^2 k^2 - (k \cdot P)^2] \\ &\quad + 24 m^2 (2k^2 - k \cdot P) \\ &= -16 m^2 k^2 + 16 (k \cdot P)^2 - 24 m^2 k \cdot P \end{aligned}$$

Notice that the integrand is finite as $k \rightarrow 0$, since the numerator is proportional to k . Naively, we see that the integral could be UV divergent, however we will see that this integral is finite.

So,

$$a_e^{(1)} = -i e^2 \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2 [(P/2-k)^2 - m^2]^2} \left(-\frac{4}{3} k^2 + \frac{4}{3} (k \cdot P)^2 - 2 k \cdot P \right)$$

We now introduce Feynman parameters

$$\begin{aligned}\frac{1}{A^2 B} &= 2 \int_0^1 dx_1 \int_0^1 dx_2 \delta(1-x_1-x_2) \frac{x_1}{[x_1 A + x_2 B]^3} \\ &= 2 \int_0^1 dx \frac{x}{[x A + (1-x) B]^3} \\ &= 2 \int_0^1 dx \frac{x}{[x(A-B)+B]^3}\end{aligned}$$

So, with

$$A = (P_{12} - k)^2 - m^2 = \frac{P^2}{4} + k^2 - k \cdot P - m^2$$

$$= k^2 - k \cdot P$$

and

$$B = k^2$$

$$\text{So, } \frac{1}{k^2 [(P_{12} - k)^2 - m^2]} = 2 \int_0^1 dx \frac{x}{[k^2 - k \cdot P x]^3}$$

So,

$$a_e^{(0)} = -2ie^2 \int_0^1 dx \times \int \frac{d^4 k}{(2\pi)^4} \frac{1}{(k^2 - k \cdot P_x)^3} \left(-\frac{4}{3} k^2 + \frac{4}{3m^2} (k \cdot P)^2 - 2k \cdot P \right)$$

To perform the momentum integral, let us perform this in d-dimensions to see the cancellation of the UV divergence.

Note the following integrals (No proof)

$$I(n) = \int d^d k \frac{1}{(k^2 + 2\vec{P} \cdot k - M^2)^n} = i\pi^{d/2} \frac{(-1)^n}{\Gamma(n)} \frac{\Gamma(n-d/2)}{(\vec{P}^2 + M^2)^{n-d/2}}$$

$$I_\mu(n) = \int d^d k \frac{k_\mu}{(k^2 + 2\vec{P} \cdot k - M^2)^n} = -\hat{P}_\mu I(n)$$

$$I_{\mu\nu}(n) = \int d^d k \frac{k_\mu k_\nu}{(k^2 + 2\vec{P} \cdot k - M^2)^n} = \left[\hat{P}_\mu \hat{P}_\nu - \frac{1}{2} g_{\mu\nu} \left(\frac{\vec{P}^2 + M^2}{n-1-d/2} \right) \right] I(n)$$

For us, $\hat{P} = -\frac{P_x}{2}$, $M^2 = 0$, $n = 3$, and let $d = 4 - 2\varepsilon$

$$\text{so, } I(3) = -\frac{i\pi^2}{2\hat{P}^2} = -\frac{2i\pi^2}{P^2 x^2} = -\frac{i\pi^2}{2m^2 x^2}$$

↑
take $\varepsilon \rightarrow 0$
at end

and

$$I_\mu(3) = +\frac{x}{2} \hat{P}_\mu I(3)$$

$$\begin{aligned} I_{\mu\nu}(3) &= \left[\frac{x^2}{4} \hat{P}_\mu \hat{P}_\nu - \frac{1}{2} g_{\mu\nu} \frac{(P^2 x^2 / 4)}{2 - d/2} \right] I(3) \\ &= \frac{x^2}{4} \left[\hat{P}_\mu \hat{P}_\nu - \frac{2m^2}{\varepsilon} g_{\mu\nu} \right] I(3) \end{aligned}$$

Therefore,

$$\begin{aligned} \int d^4 k \frac{1}{(k^2 - k \cdot P_x)^3} &\left(-\frac{4}{3} k^2 + \frac{4}{3m^2} (k \cdot P)^2 - 2k \cdot P \right) \\ &= -\frac{4}{3} g^{\mu\nu} I_{\mu\nu}(3) + \frac{4}{3m^2} P^\mu P^\nu I_\mu(3) - 2P^\mu I_\mu(3) \end{aligned}$$

$$\begin{aligned}
& \int d^4k \frac{1}{(k^2 - k \cdot P_x)^3} \left(-\frac{4}{3} k^2 + \frac{4}{3m^2} (k \cdot P)^2 - 2k \cdot P \right) \\
&= -\frac{4}{3} g^{\mu\nu} I_{\mu\nu}(3) + \frac{4}{3m^2} P^\mu P^\nu I_{\mu\nu}(1) - 2P^\mu I_\mu(3) \\
&= -\frac{4}{3} g^{\mu\nu} \frac{x^2}{4} \left[P_\mu P_\nu - \frac{2m^2}{\epsilon} g_{\mu\nu} \right] I(3) \\
&\quad + \frac{4}{3m^2} P^\mu P^\nu \frac{x^2}{4} \left[P_\mu P_\nu - \frac{2m^2}{\epsilon} g_{\mu\nu} \right] I(3) \\
&\quad - 2P^\mu \frac{x}{2} P_\mu I(3) \\
&= \left\{ -\frac{4}{3} \frac{x^2}{4} \left[P^2 - \frac{2m^2 \cdot 4}{\epsilon} \right] + \frac{4}{3m^2} \frac{x^2}{4} \left[P^2 P^2 - \frac{2m^2}{\epsilon} P^2 \right] - x P^2 \right\} I(3) \\
&= \left\{ \frac{x^2}{3} \left[-P^2 + \frac{2m^2 \cdot 4}{\epsilon} \right] + \frac{P^2 P^2}{m^2} - \frac{2m^2}{\epsilon} \frac{P^2}{m^2} \right\} I(3) - x P^2 I(3) \\
&= \left\{ \frac{x^2}{3} \left[-4m^2 + 16m^2 \right] - x \cdot 4m^2 \right\} I(3) \quad \text{Naive UV divergence cancels.} \\
&= 4m^2 x (x-1) I(3) \\
&= -2 \frac{(x-1)}{x} i\pi^2
\end{aligned}$$

Therefore, the simplified expression for $\alpha_e^{(1)}$ is

$$\begin{aligned}
\alpha_e^{(1)} &= -2 \frac{i e^2}{(2\pi)^4} \int_0^1 dx \times \left\{ \frac{2}{x} (1-x) i\pi^2 \right\} \\
&= + \frac{(2\pi)^2}{(2\pi)^4} e^2 \int_0^1 dx (1-x) = \frac{e^2}{4\pi^2} \left[x - \frac{x^2}{2} \right]_0^1 = \frac{1}{2} \left(\frac{e^2}{4\pi^2} \right)
\end{aligned}$$

So,

$$\alpha_e^{(1)} = \frac{1}{2} \left(\frac{e^2}{4\pi^2} \right)$$

Recall the fine-structure constant , $\alpha = \frac{e^2}{4\pi^2}$

$$\Rightarrow \boxed{\alpha_e^{(1)} = \frac{\alpha}{2\pi}}$$

Schwinger's triumph!

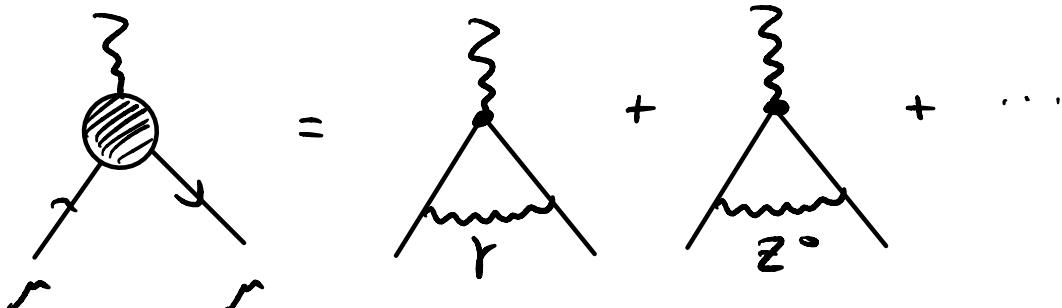
here, $\alpha_e^{(1)} = 1.16 \dots \times 10^{-3}$

The measured value is $(\alpha_e)^{\text{exp}} = 1.15965218073(28) \times 10^{-3}$
while the current theoretical prediction ($O(\alpha^{10})$!) is
 $(\alpha_e)^{\text{SM}} = 1.15965218188(78) \times 10^{-3}$

A remarkable prediction!

For muons, the leading order is identical , $\alpha_\mu^{(1)} = \frac{\alpha}{2\pi}$,
but since $\frac{m_\mu}{m_e} \sim 200$, the value of $\frac{g-2}{2}$ is susceptible

to searches for new physics.



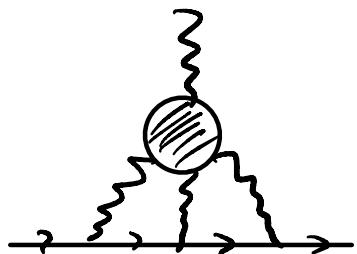
$$\text{For the } Z^0\text{-boson, } \alpha_\mu \sim \left(\frac{m_\mu}{m_Z}\right)^2 \cdot \frac{g_e^2 e}{16\pi} \cdot \log\left(\frac{m_Z^2}{m_\mu^2}\right)$$

↓ EW coupling.

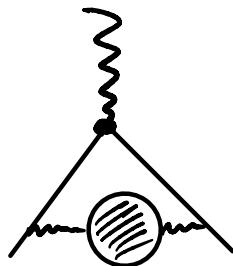
There is an intense search for new physics by looking at the muon anomaly. The current (as of Aug. 2023) experimental measurement is

$$(\alpha_\mu)^{\text{exp}} = 116592.059(22) \times 10^{-11}$$

The SM theory is a little bit muddled. There are discrepancies between lattice methods and data-driven analyses. There is an apparent 4.2σ discrepancy with experiment. \mathcal{BD} , two hadronic processes are the leading sources



and



Hadronic light-to-light

Hadronic vacuum polarization