

Landau Singularities: Scope and Applications to QCD

Ajay S. Sakthivasan ¹

¹HISKP & BCTP, University of Bonn

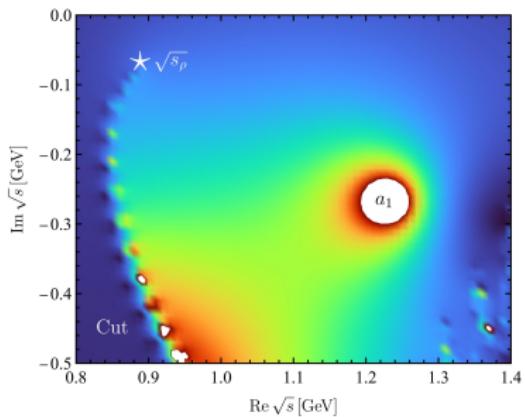
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Presentation Outline

- 1 Introduction
- 2 Analytic Aspects
- 3 Infinite Volume Unitary Formalism
- 4 Summary & Conclusions

Introduction

- Hadronic resonances correspond to poles of the S -matrix.

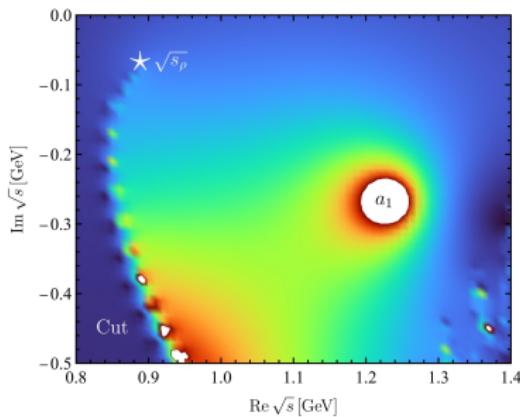


Pole position of the $a_1(1260)$.¹

¹Figure from: [10.1103/PhysRevD.105.054020](https://doi.org/10.1103/PhysRevD.105.054020)

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- Experimentally — peaks in invariant mass distributions.

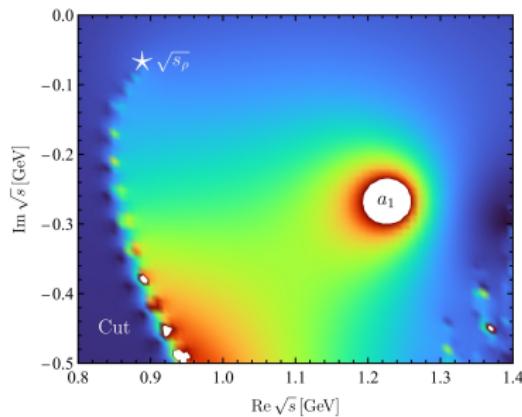


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Introduction

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- Experimentally — peaks in invariant mass distributions.
- All observed peaks correspond to hadronic states? **No.**



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- **Landau Singularities** may not correspond to poles of the S -matrix, but can *mimic* a resonance.
- Originally proposed by Landau; Recently used to explain a lot of observed peaks.

Singularities of Integrals

- Consider the complex-valued integral:

$$f(z) = \int_C dw g(z, w).$$

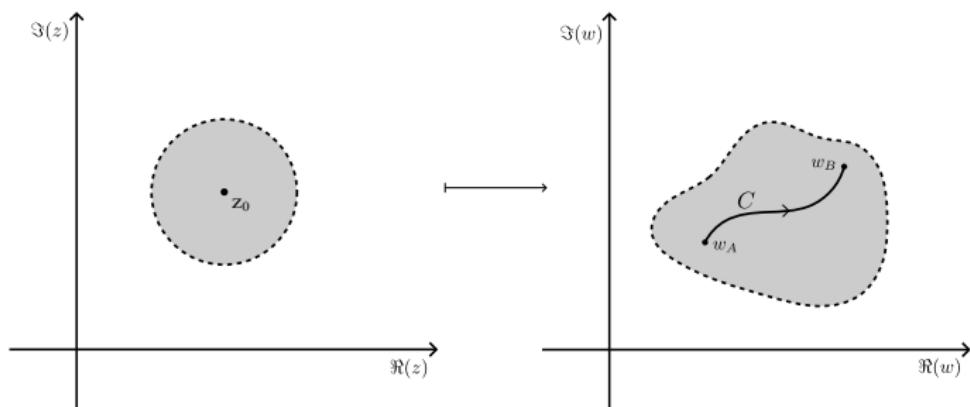
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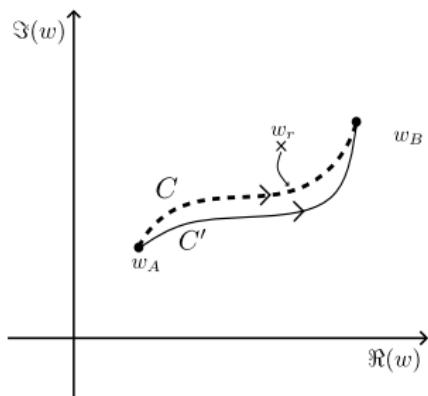
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- What can one say about the singularities of f , given the knowledge of the singularities of g ?

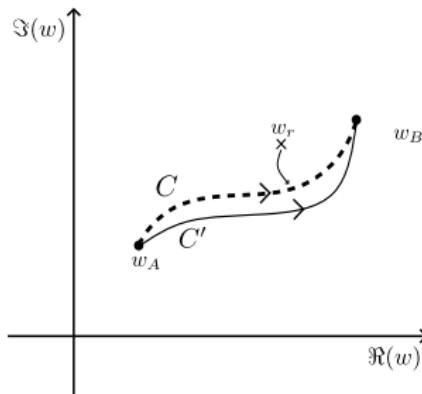
- Case: There are no singularities in the vicinity of the contour $\Rightarrow f$ is analytic.



- Case: Singularity approaches the contour \implies Cauchy tells us that f is still analytic.



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- Contour deformation makes evaluation possible.

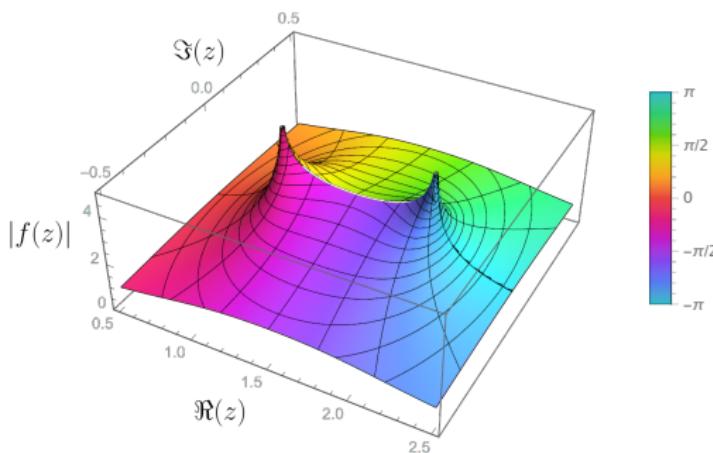
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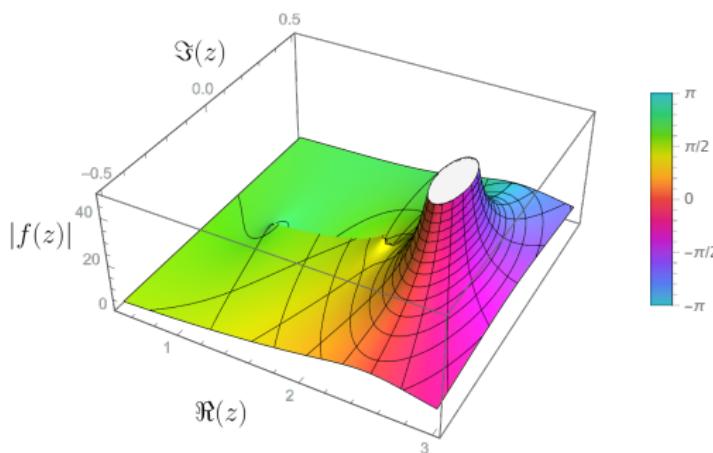


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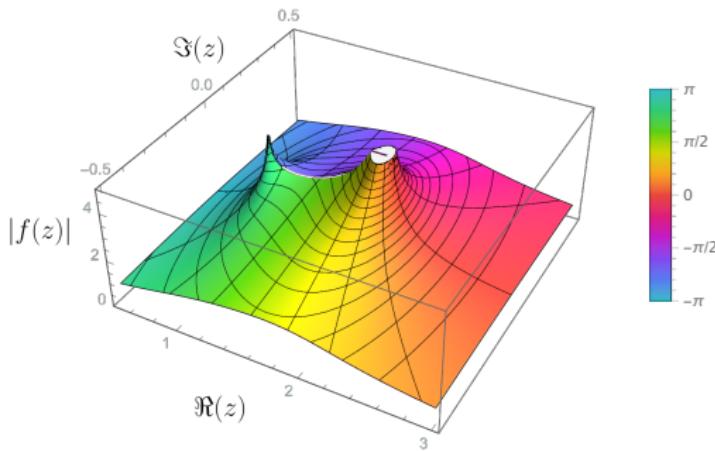
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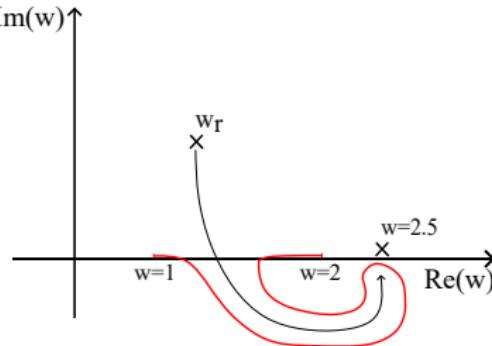
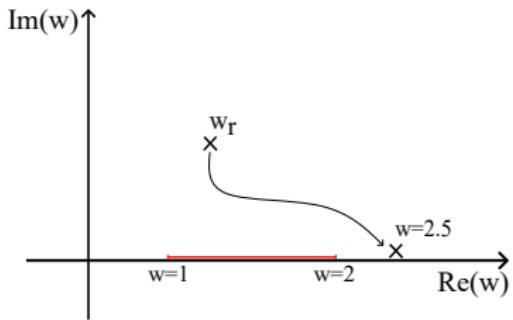


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- The situation was as follows:



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- We only motivate these generalisations.

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- For some real parameters α_i , the Landau equations read,

$$\begin{aligned}\alpha_i(q_i^2 - m_i^2) = 0 &\implies \alpha_i = 0 \text{ or } q_i^2 = m_i^2, \\ \sum_{i \in \text{loop}} \alpha_i q_i^\mu(k_j) &= 0.\end{aligned}$$

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- Third representation is interesting — one can use it to study the nature of the singularities. Also, reveals the existence of a different class of singularities.

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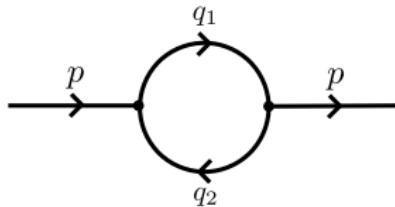
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 - ② For a graph with N internal propagators, $N - 1$ contractions results in a graph with trivial analyticity properties.
- Finally, as mentioned, we can also determine the nature of singularity, i.e., poles, algebraic/square root/logarithmic branch point.

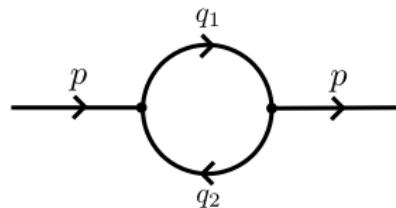
The Bubble Graph

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$$q_1^2 = m_1^2 \quad \& \quad q_2^2 = m_2^2,$$

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- The first singularity is exactly the **normal threshold**. Whereas, the second singularity is the **anomalous threshold**.
- One can determine the nature of the singularity as a **square root branch point**.

- A very important observation: If the integral is evaluated, one notices that the anomalous threshold **does not lie on the principal sheet**.

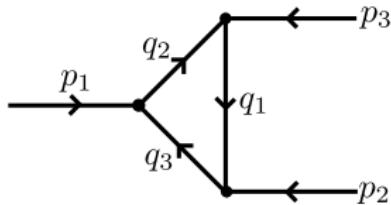
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- How can one identify the actually relevant singularities?
- Something that was briefly mentioned before: **we require the α 's to be positive**.

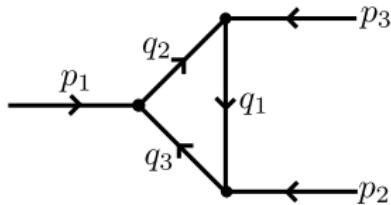
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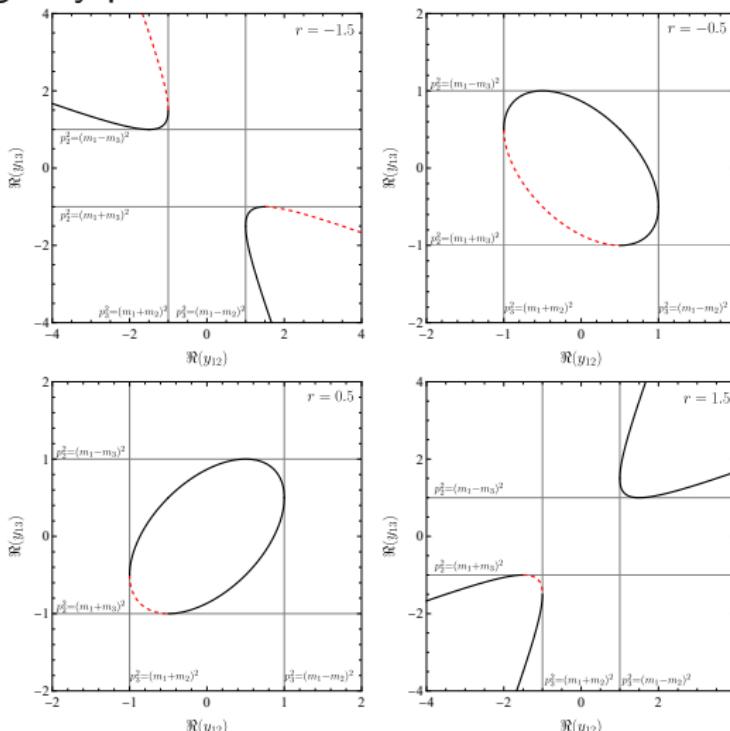
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- The above equation denotes a complex surface that can be understood for different cases.

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- We can attempt to gain a geometric understanding of the problem, with the imaginary parts set to 0:



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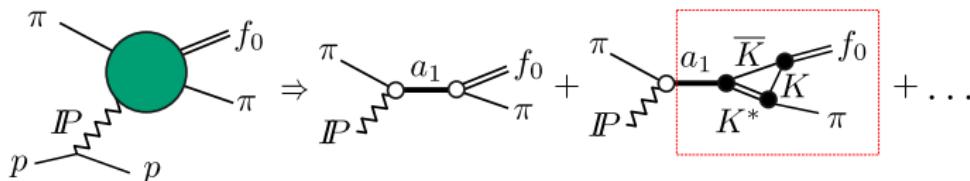
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 - ➊ To systematically study the effect of final-state rescattering.

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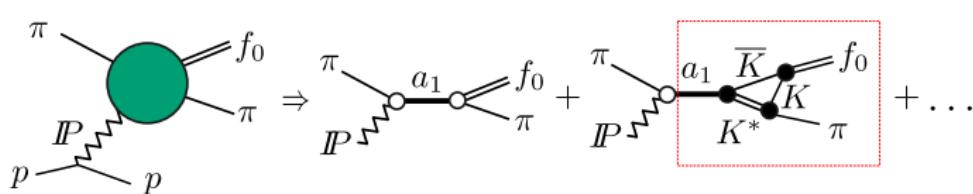
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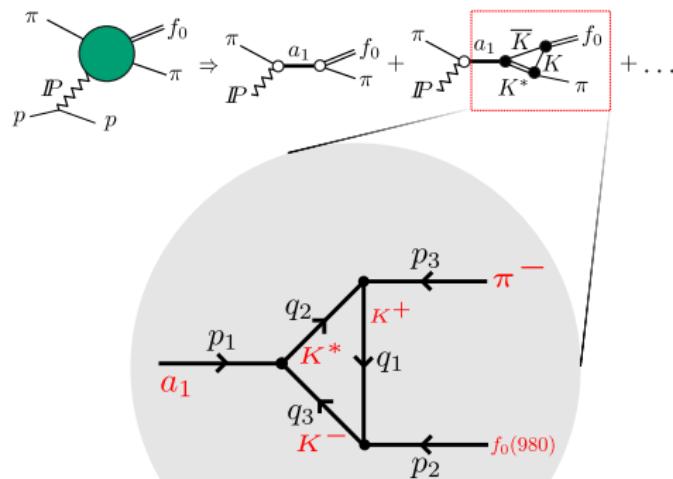
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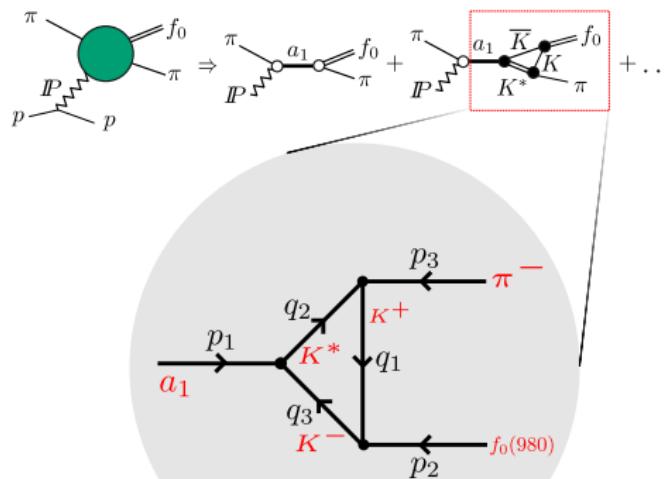
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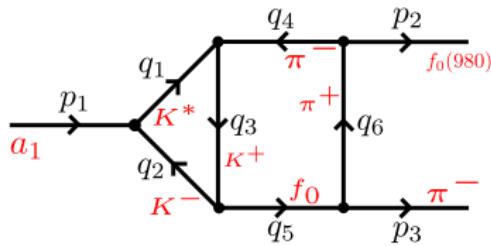
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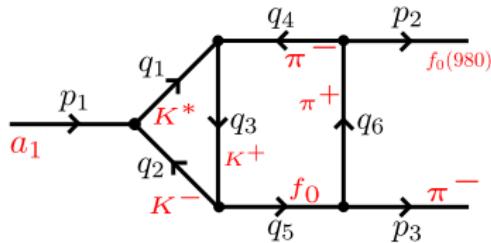
- Landau singularity for $p_3 = m_\pi$: $p_2/\text{GeV} \equiv \sqrt{\sigma}/\text{GeV} \in [0.986, 1.024]$ and $p_1/\text{GeV} \equiv \sqrt{s}/\text{GeV} \in [1.385, 1.436]$.

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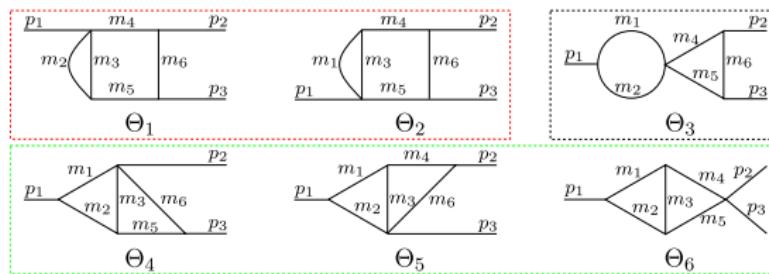


- Check for leading and subleading singularities. . .

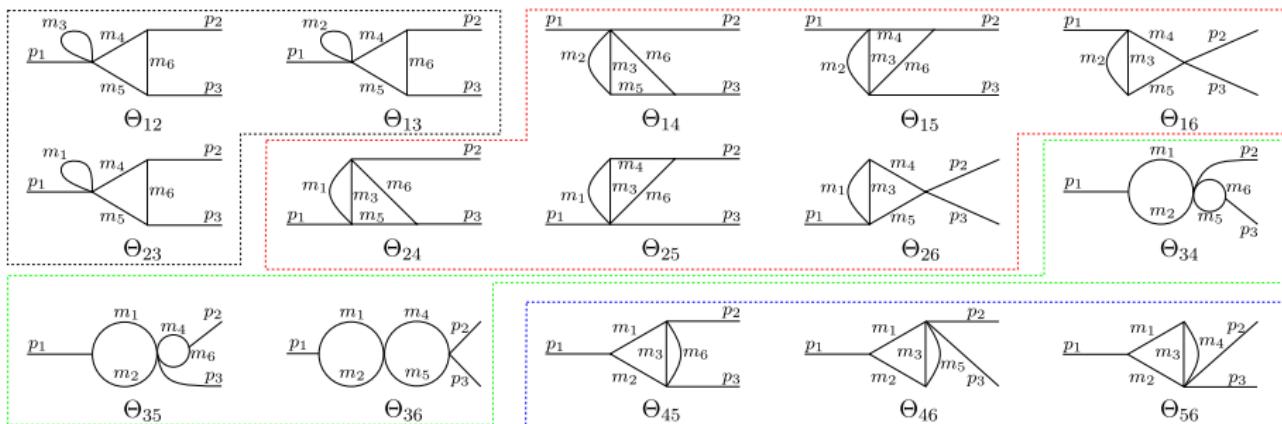
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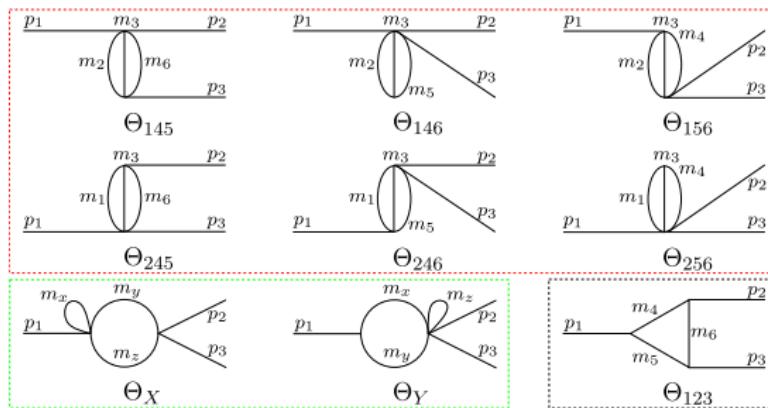
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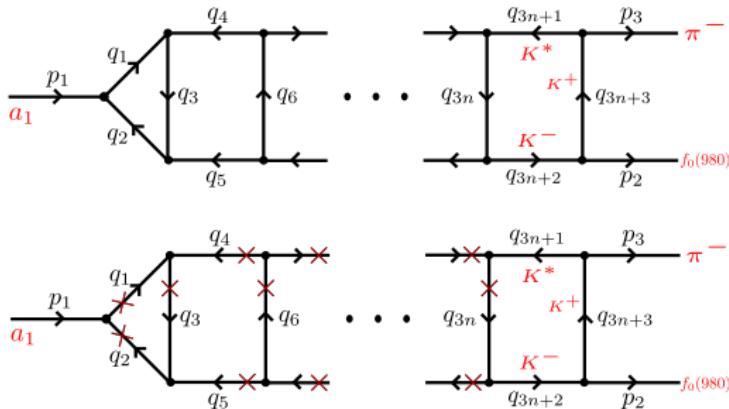


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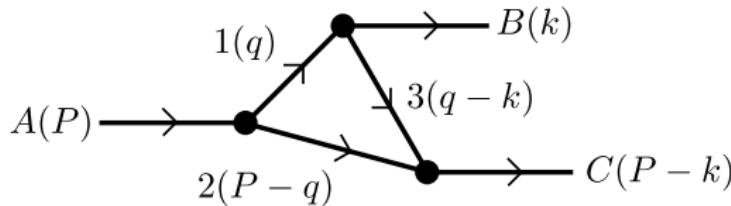


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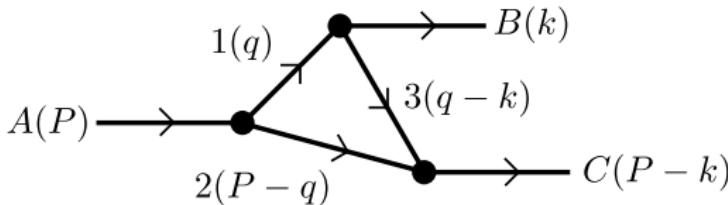
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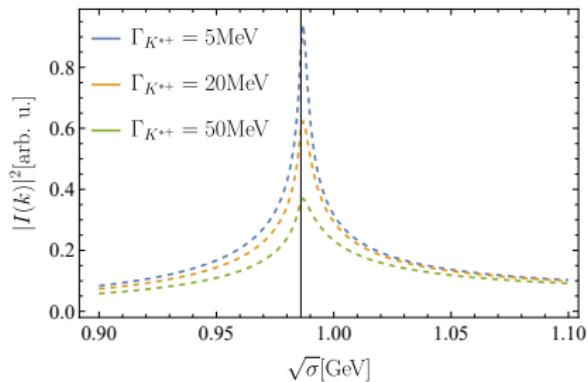


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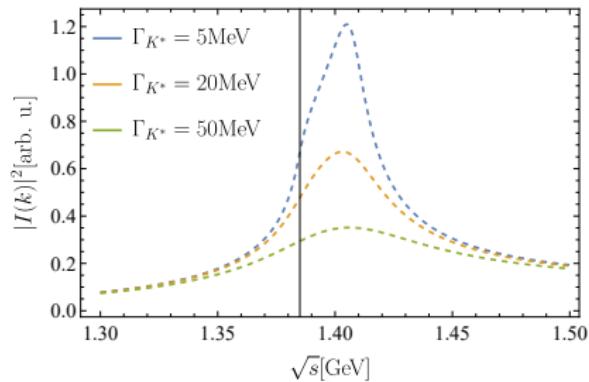
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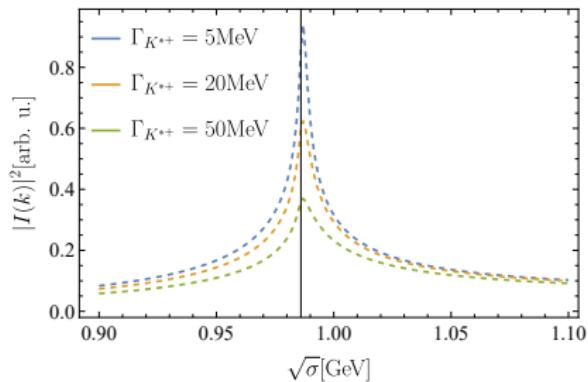
- Add an imaginary part to the masses to account for the instability (*width*) of the particles.



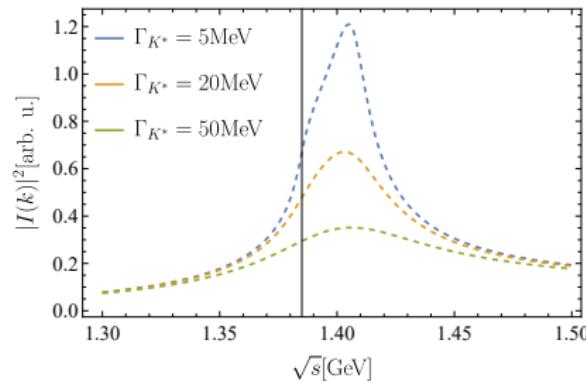
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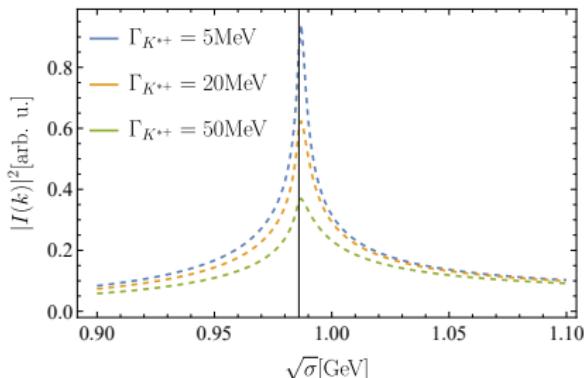


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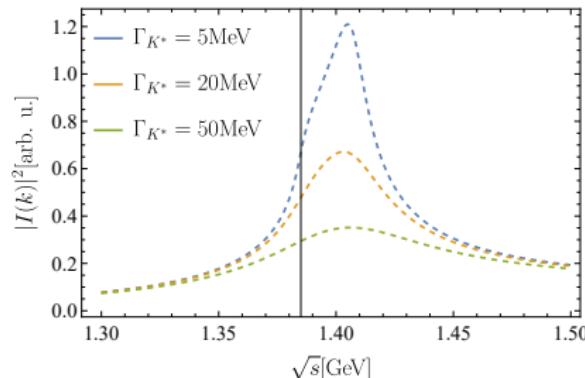


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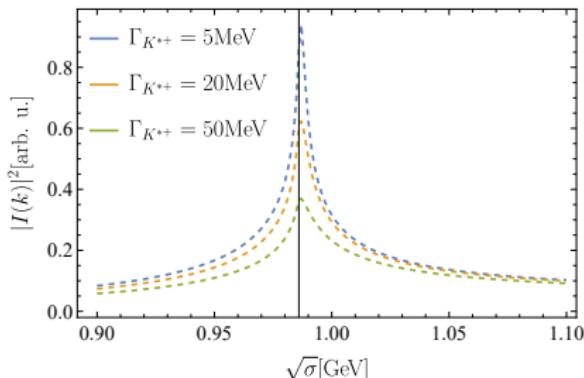


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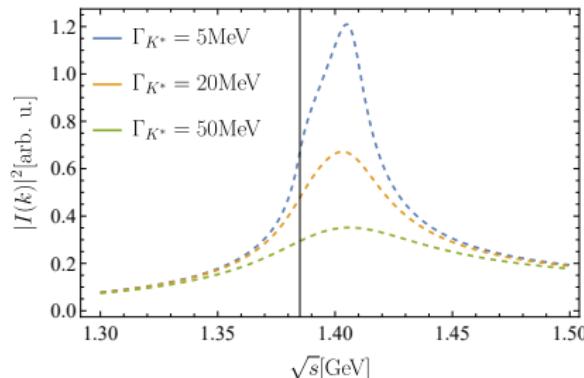


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- Accommodating another rescattering tedious; accommodating infinite rescattering is impossible this way.

Infinite Volume Unitary Formalism

- Study the system using **Infinite Volume 3-Body Unitary Formalism**, due to Mai et al.³

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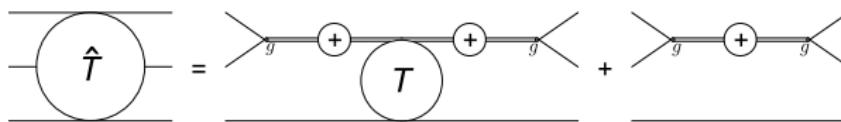
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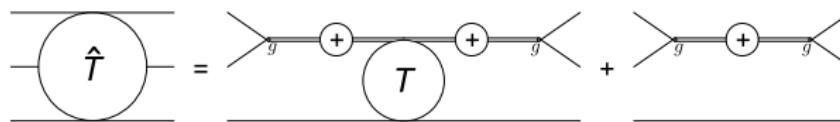


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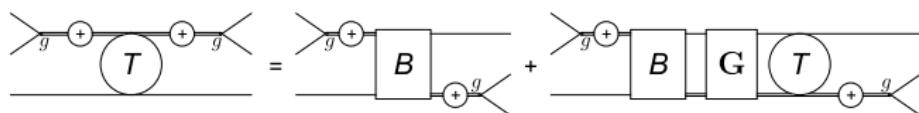


- It can be seen that the connected part contains the exchange dynamics, including the triangle diagram.

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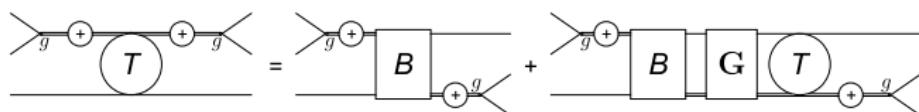
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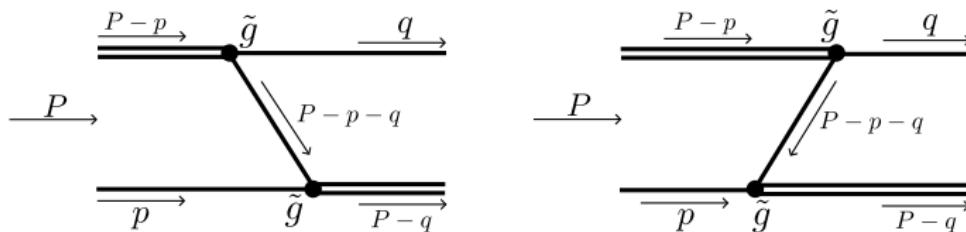
$$T(p, q) = B(p, q) + \int \frac{d^3 l}{(2\pi)^3} \frac{1}{2E_l} B(p, l) \tau(\sigma(l)) T(l, q).$$

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- Make use of **unitarity** (hence, the name) and dispersion relations to reconstruct B and τ .

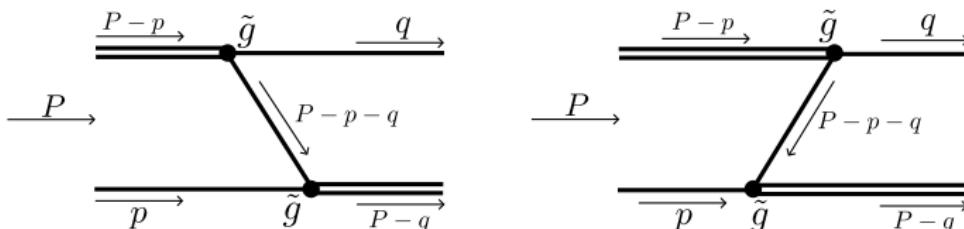
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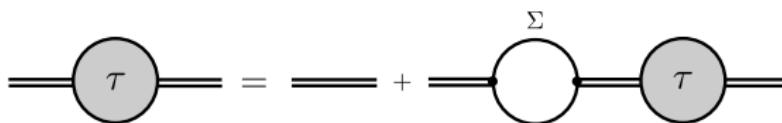


- The first diagram corresponds to a (time-)forward propagating exchange particle. A diagram of the second type can be added to the above equation to make it Lorentz covariant.

- Similarly, the explicit form of τ can also be determined.

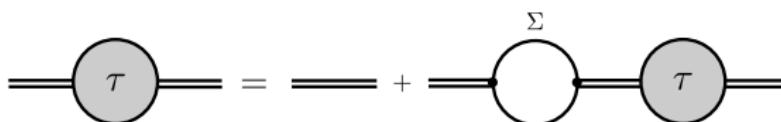
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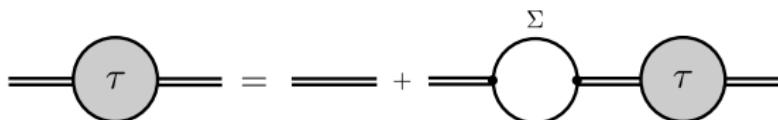
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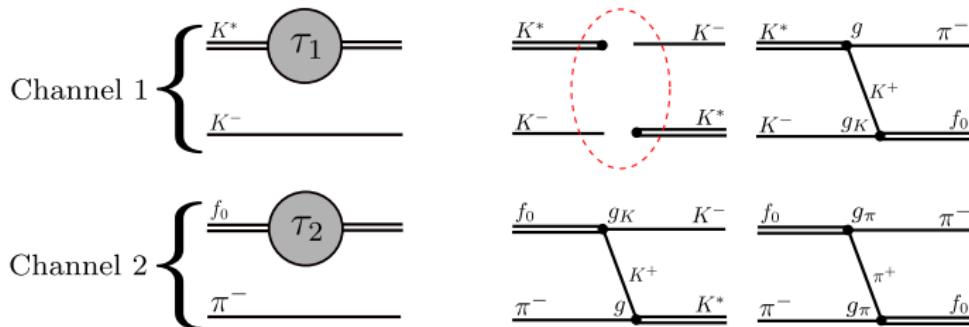
- A fine point: in principle, one also takes **3-body contact interactions** into account. This amounts to replacing B with $B + C$.
- In practice, C is suitably parametrised and experimental input is used.

Back to the $a_1(1420)$

- We work with a toy model: all particles and the source are scalar; isobars and spectators are always in relative s -wave; a little sloppy with isospin invariance. **Concentrate only on understanding the analytic structure within this framework.**

Back to the $a_1(1420)$

- We work with a toy model: all particles and the source are scalar; isobars and spectators are always in relative s -wave; a little sloppy with isospin invariance. **Concentrate only on understanding the analytic structure within this framework.**
- We have two possible isobars: K^* and f_0 *implies* need to carry out a coupled channel analysis.



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- These are fixed through PDG values of physical mass and width.
- For f_0 we also make use of g_K/g_p ratio from BaBar collaboration.

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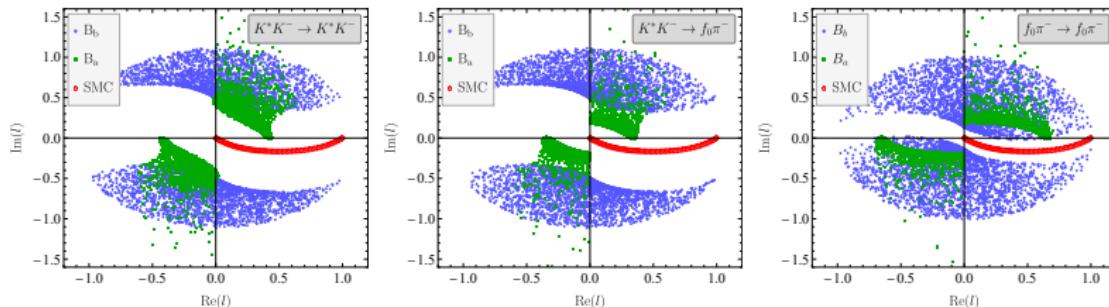
- This defines the **loop expansion** for the invariant mass distributions.
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- To note, the discrete version of the integral can be cast into a matrix equation, and then inverted to obtain the connected amplitude, given that the matrix is non-singular, i.e.,

$$T = B(\mathbb{1} - BG)^{-1}$$

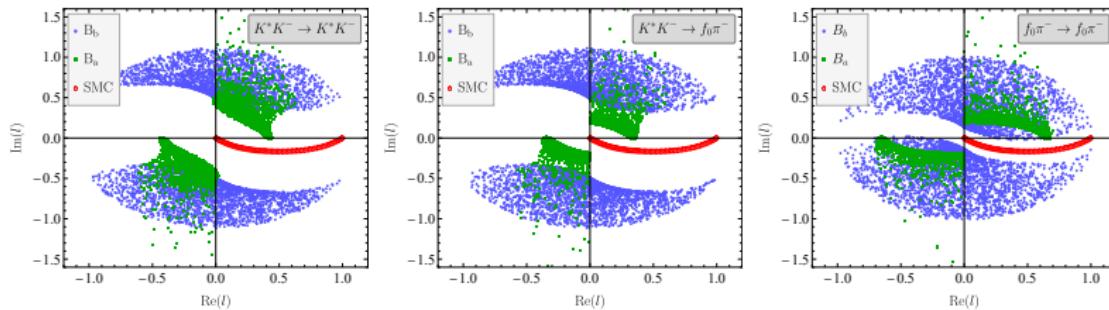
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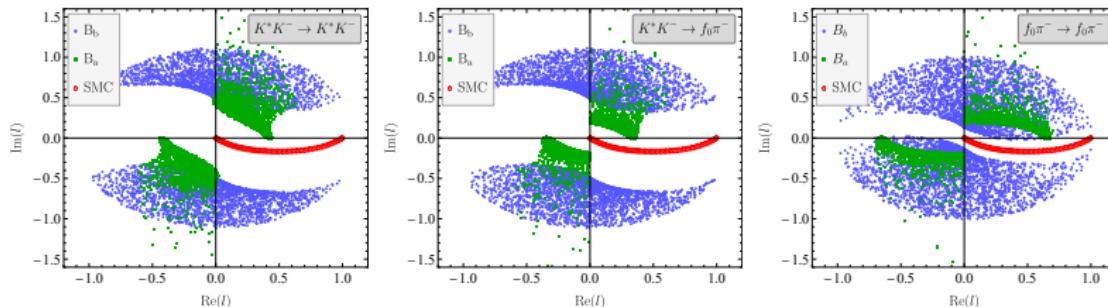


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- The integration contour for the self-energy integral should be consistent with the primary integration contour. The choice of contours is a matter of taste.
- But, we need the amplitudes for $q \in \mathbb{R} \rightarrow$ can be done in different ways!

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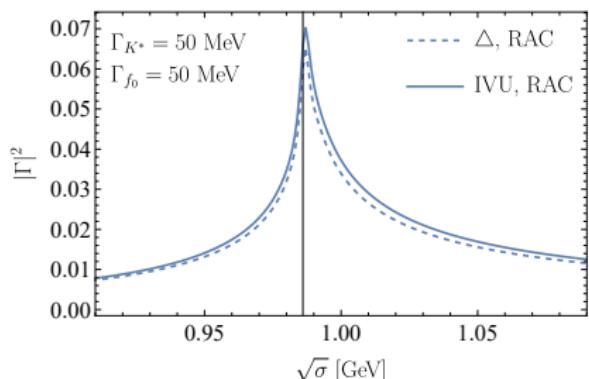
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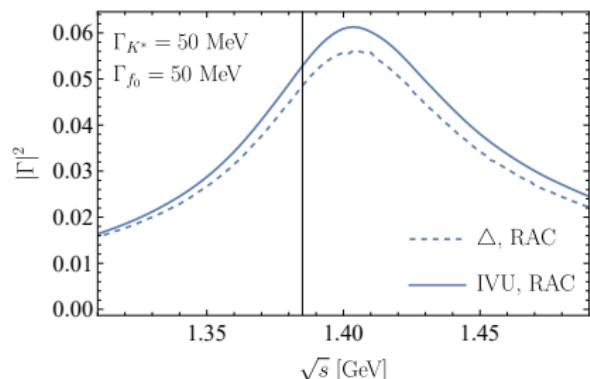
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- Why not Padé? Padé was unable to reproduce sharper features like cusps.

- At $1 + \infty$ -loop level,

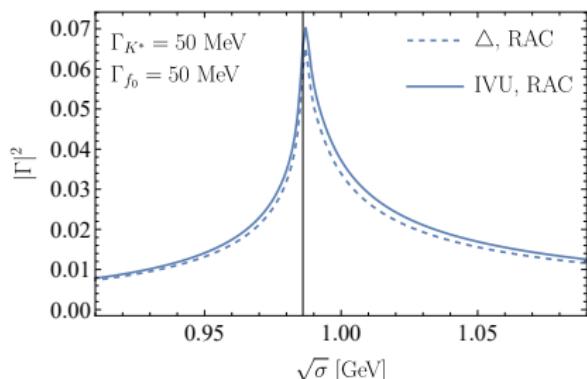


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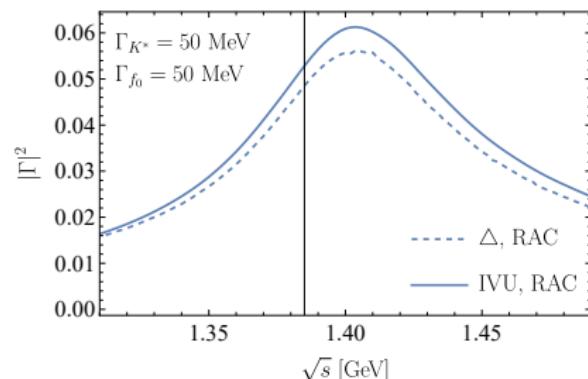


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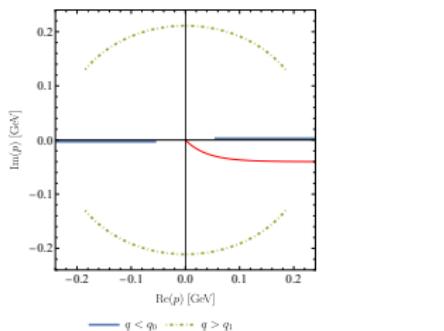
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- We see the Landau singularity in both the plots.

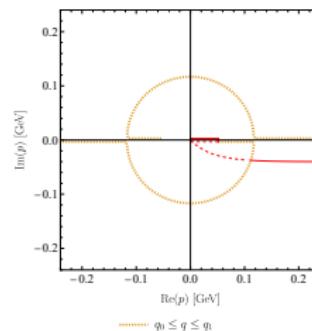
- Step2.2: **Cahill & Sloan Method**:⁶ We make use Cauchy's integral theorem instead.

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- To expound, instead of a continued fraction, use the knowledge of the **location of the branch cuts** to analytically continue the amplitude for $q \in \mathbb{R}$.⁷



Branch cuts of B for q outside the triangle singularity region.

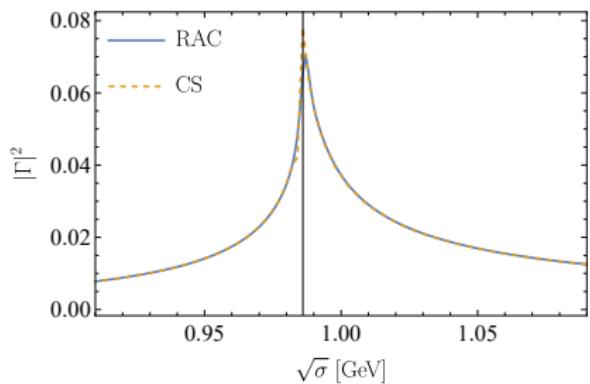


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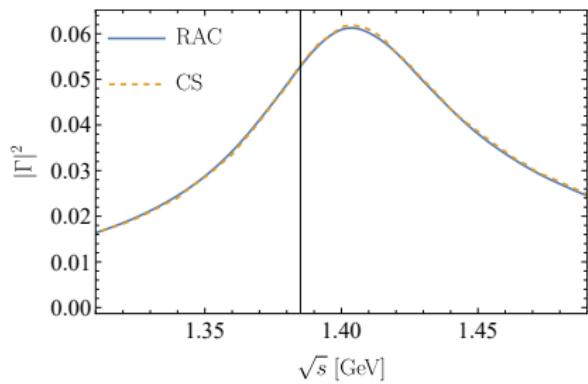
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⁷For a comprehensive review: *The Quantum Mechanical Three-Body Problem*, Schmid & Ziegelmann.

- Matches with the other method.



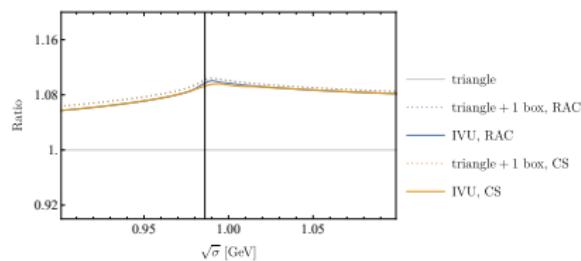
Amplitude squared vs. $\sqrt{\sigma}$, for
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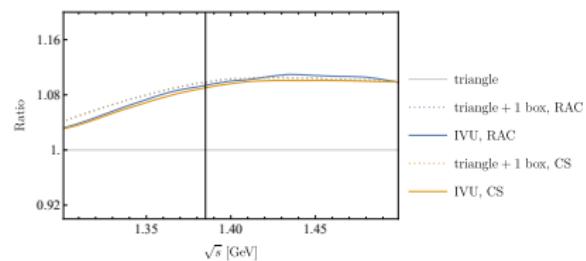
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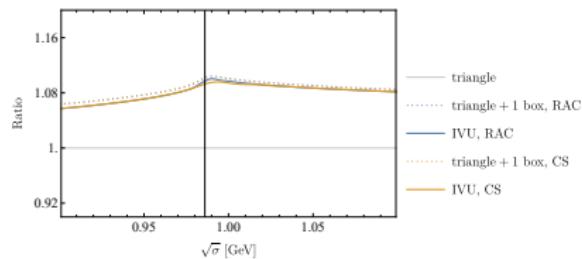


Ratios of amplitudes to triangle amplitudes vs. $\sqrt{\sigma}$, for $\sqrt{s} = 1.42$ GeV.

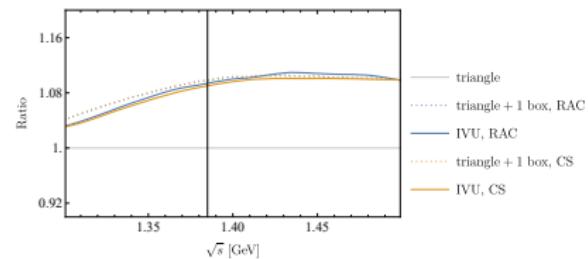


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- The amplitudes at triangle +1 box are about 1.1x the triangle diagram amplitudes.

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Thanks for listening!