

Triangle Singularities in a Hilbert's House

NSTAR24, York

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Presentation Outline

- 1 Introduction
- 2 Analytic Aspects
- 3 Infinite Volume Unitarity Formalism
- 4 Conclusions

Introduction

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- Experimentally — peaks in invariant mass distributions.
- All observed peaks correspond to hadronic states? **No.**

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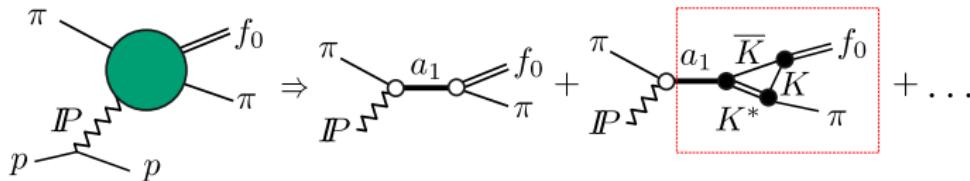
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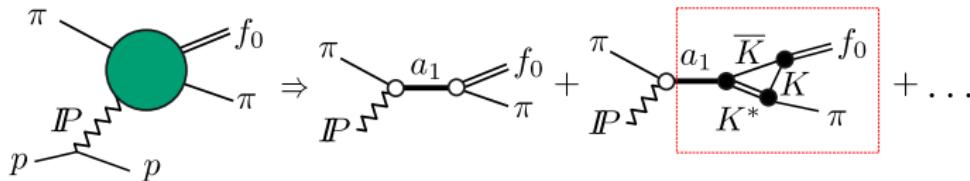
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¹Figure from: 10.1103/physrevlett.127.082501

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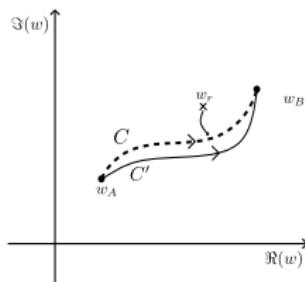
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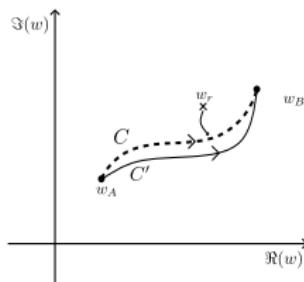


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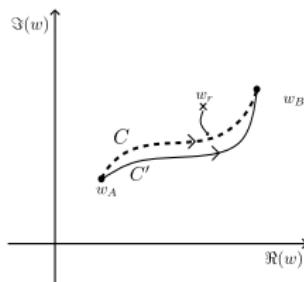
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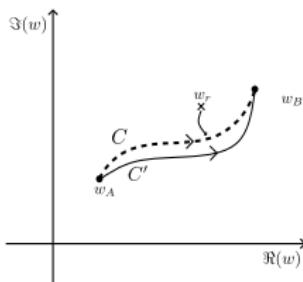
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- When is this not possible?
 - ① Endpoint singularities.
 - ② Pinch singularities.

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- For some real parameters α_i , the Landau equations read,

$$\alpha_i(q_i^2 - m_i^2) = 0 \implies \alpha_i = 0 \text{ or } q_i^2 = m_i^2,$$
$$\sum_{i \in \text{loop}} \alpha_i q_i^\mu(k_j) = 0.$$

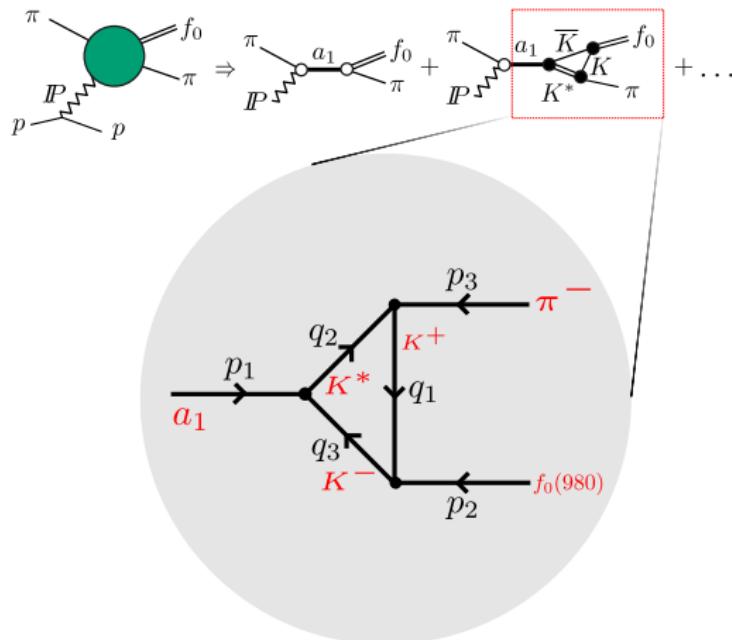
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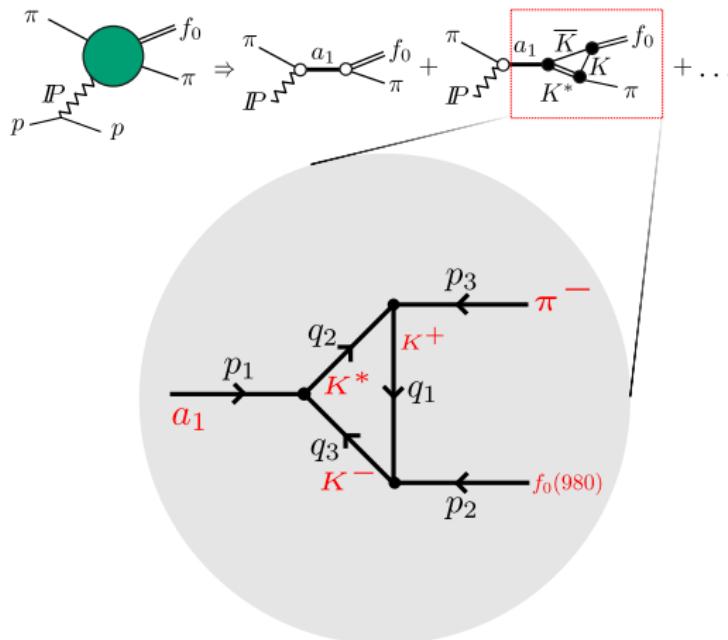
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- $\alpha_i = 0$ case is called the **subleading singularity** — equivalent to contracting the propagator.
- Interesting analogy: Feynman diagrams \rightarrow electrical circuits; propagators \rightarrow wires, with current q_i and resistance α_i . Then the Landau equations are identical to Kirchhoff's law!²

² $\Delta V = \sum IR$, for a triangle circuit with $\Delta V = 0$.

- The leading Landau singularity associated with the graph:



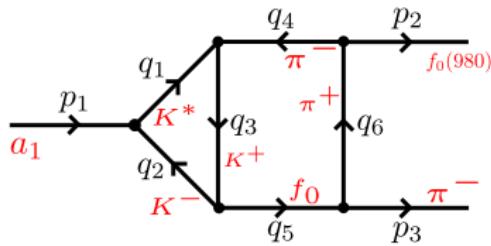
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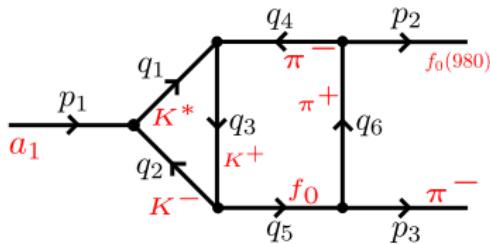
- Landau singularity for $p_3 = m_\pi$: $p_2/\text{GeV} \equiv \sqrt{\sigma}/\text{GeV} \in [0.986, 1.024]$ and $p_1/\text{GeV} \equiv \sqrt{s}/\text{GeV} \in [1.385, 1.436]$.

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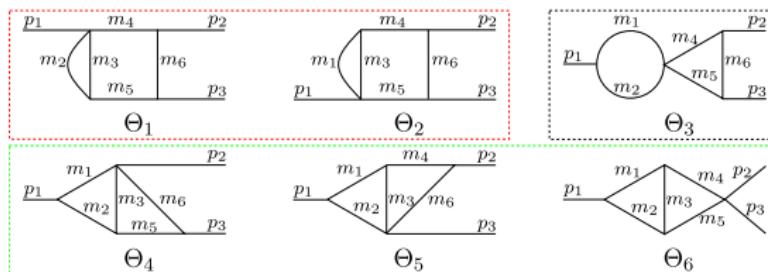


- Check for leading and subleading singularities. . .

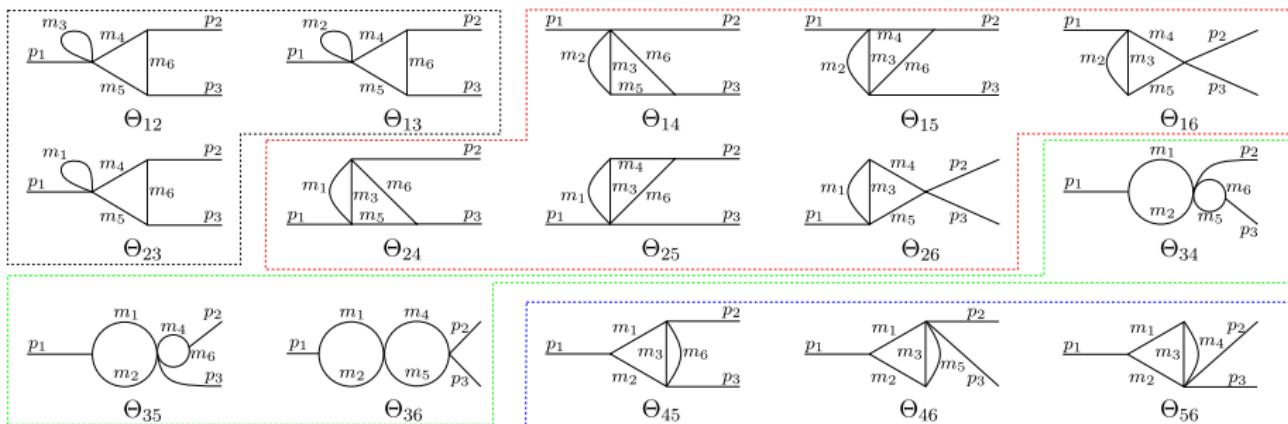
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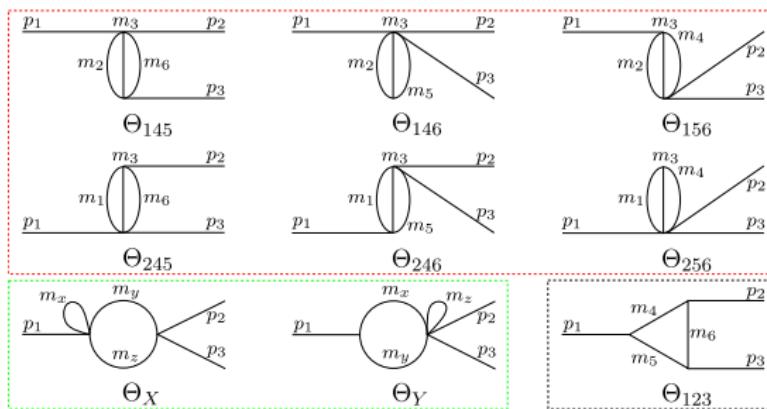
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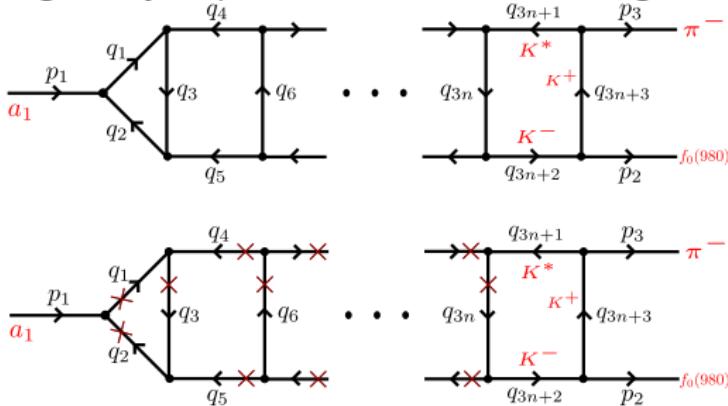


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- The triangle singularity is present and no other singularities.



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- One still needs to evaluate the corresponding Feynman integrals.

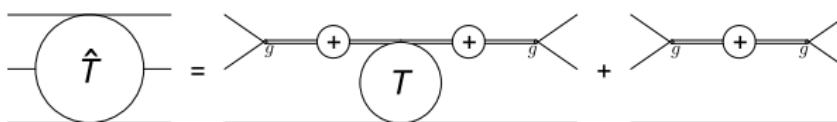
- Study the system using **Infinite Volume 3-Body Unitary Formalism**, due to Mai et al.³

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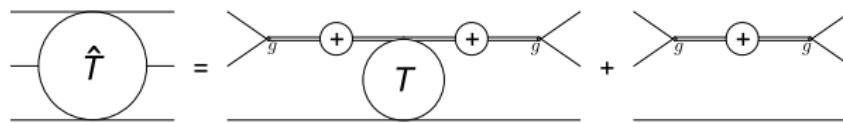
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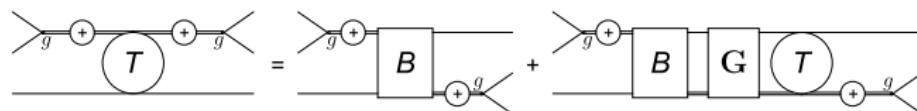
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- One considers the **Bethe-Salpeter ansatz**, $T = B + BGT$, to determine T .⁴



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- The equation to be solved,

$$T(p, q) = B(p, q) + \int \frac{d^3 l}{(2\pi)^3} \frac{1}{2E_l} B(p, l) \tau(\sigma(l)) T(l, q).$$

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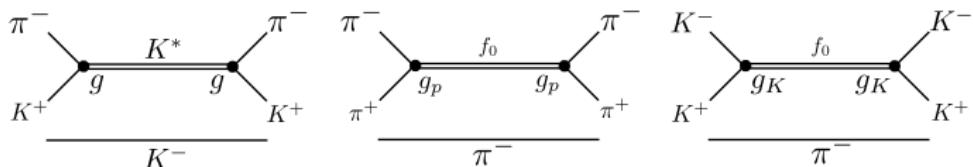
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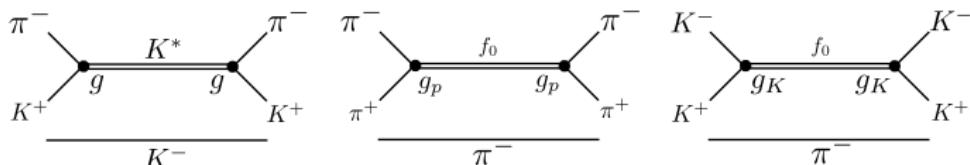
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- In this work, we consider scalar propagators, in relative s -wave.

- Need to carry out a coupled channel analysis.

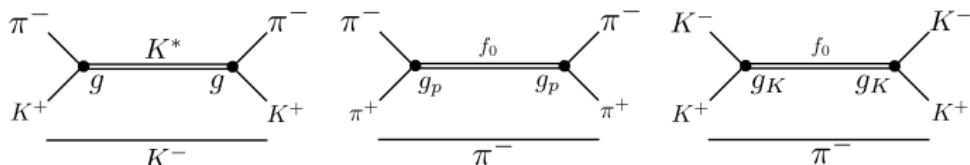


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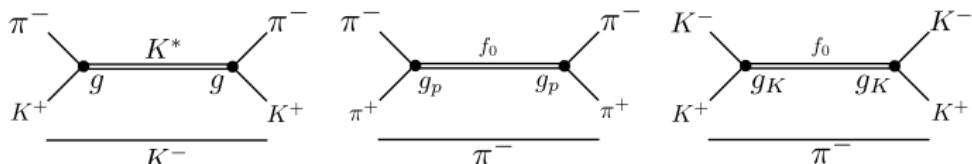
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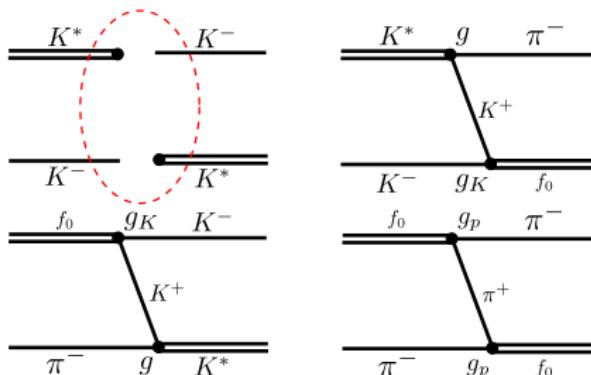
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- For f_0 we also make use of g_K/g_p ratio from BaBar collaboration.

- B is parametrised through the spectator masses.

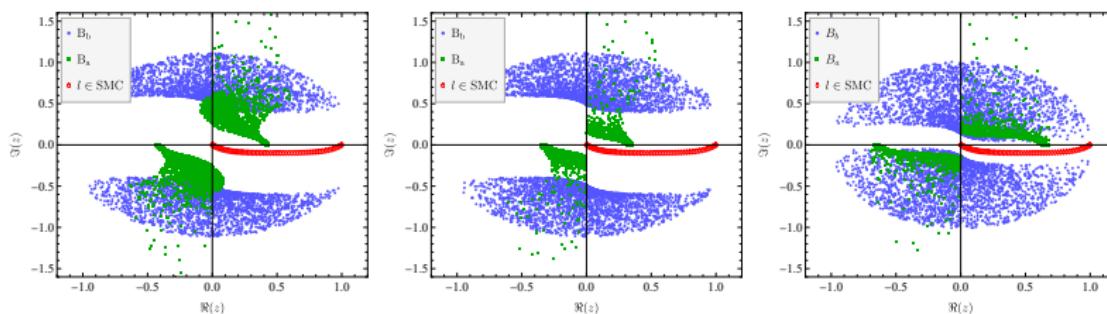
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- Diagrammatically,



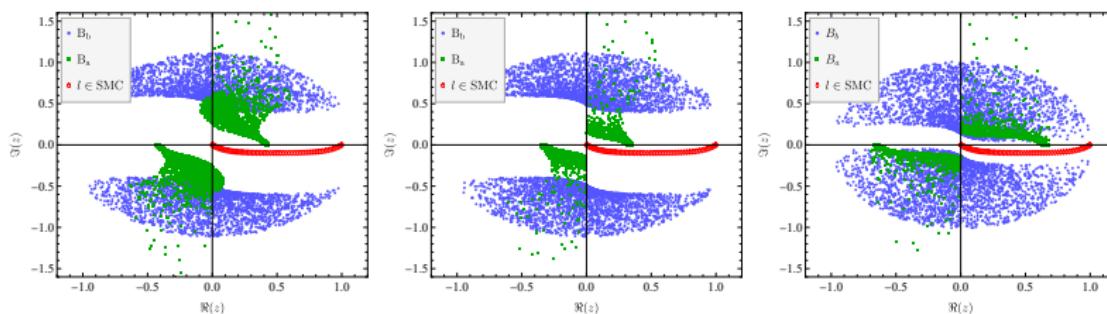
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- But, we need the amplitudes for $q \in \mathbb{R} \rightarrow$ can be done in different ways!

- The solution of the Bethe-Salpeter equation can be given in form of the **Born series**.

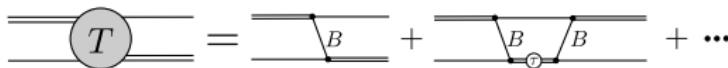
$$\text{Diagram } T = \text{Diagram } B + \text{Diagram } B \text{ (with loop)} + \dots$$


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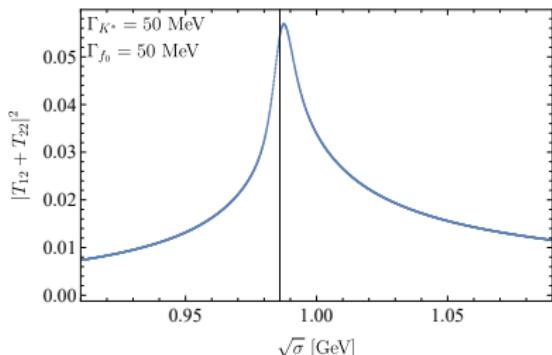


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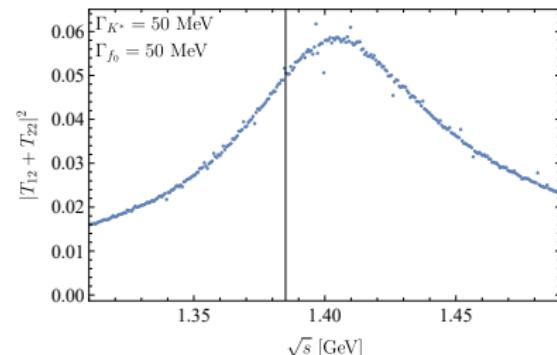
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Amplitude squared vs. $\sqrt{\sigma}$, for $\sqrt{s} = 1.42$ GeV.

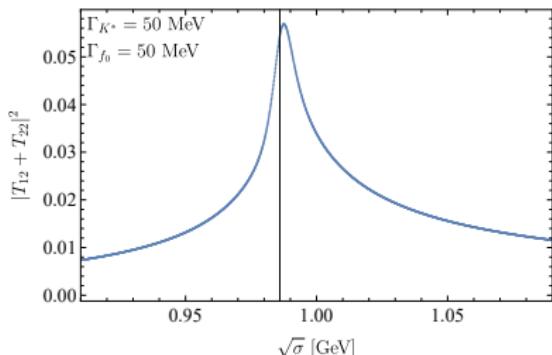


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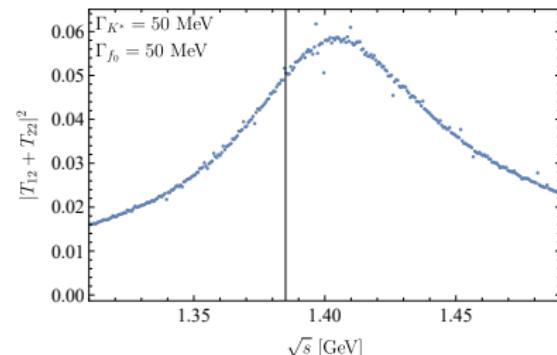
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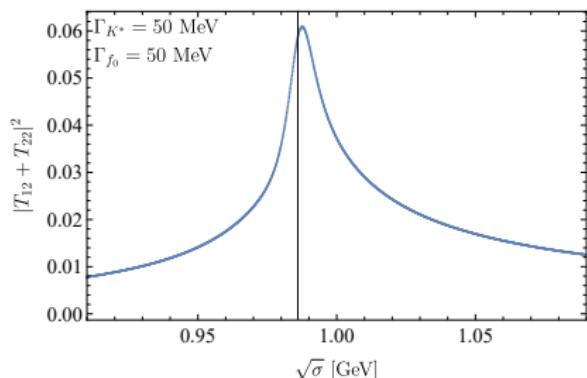
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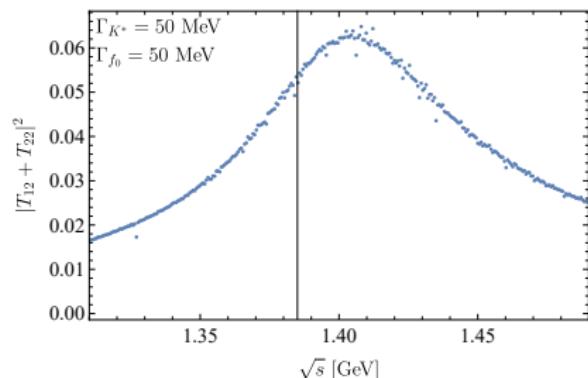
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- We see the (leading) Landau singularity in both the plots.

- At $1 + \infty$ -loop level,

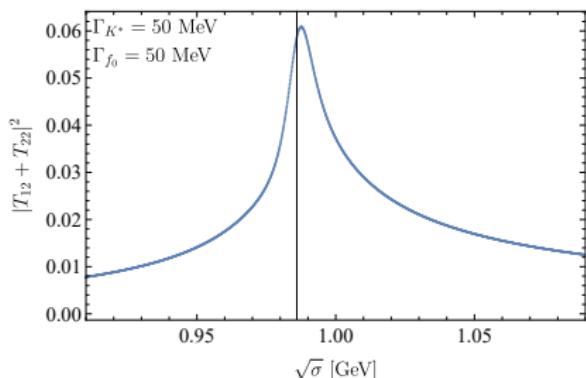


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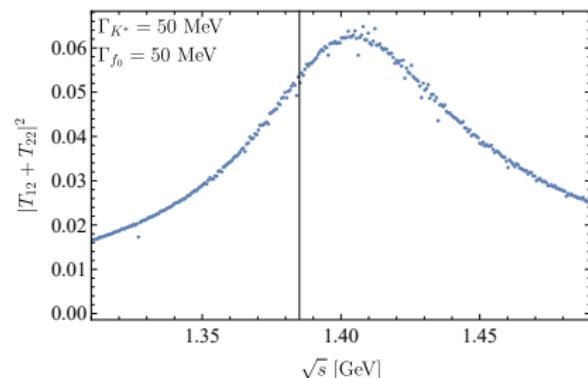


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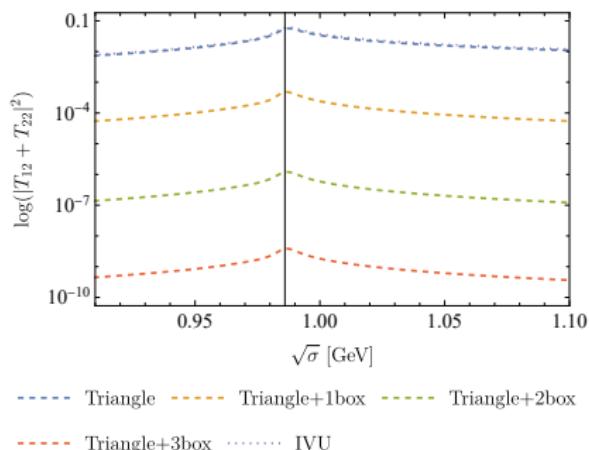


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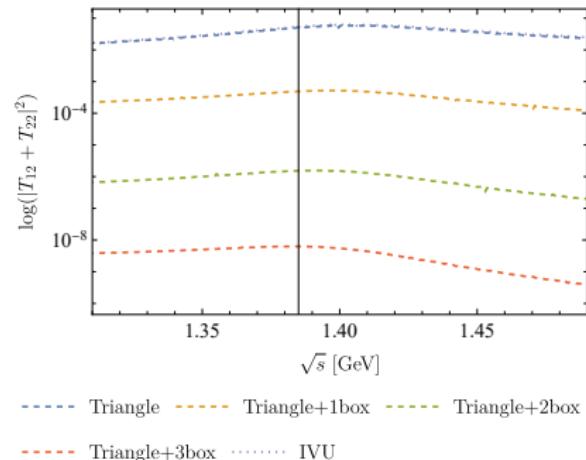
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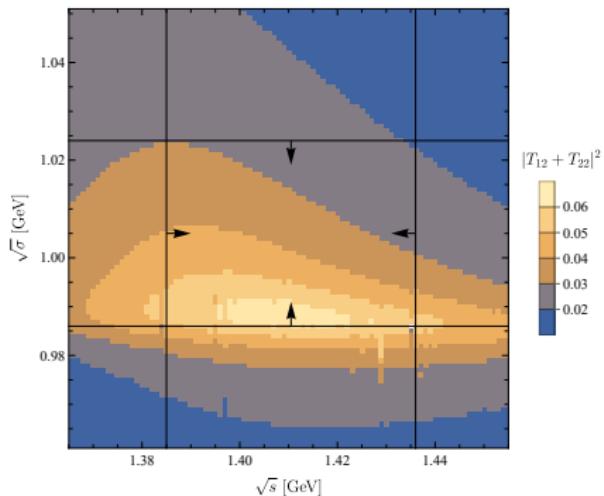


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- Characteristic of triangle singularities: very sensitive to invariant masses!



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- Outlook:
 - ① Realistic model — spin with pseudovector source.

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- The higher order diagrams did not have significant contributions \Rightarrow explains why no need to take final state interactions into account.
- Outlook:
 - ① Realistic model — spin with pseudovector source.
 - ② $a_1(1420)$ on the lattice? Implement finite volume unitarity.⁵

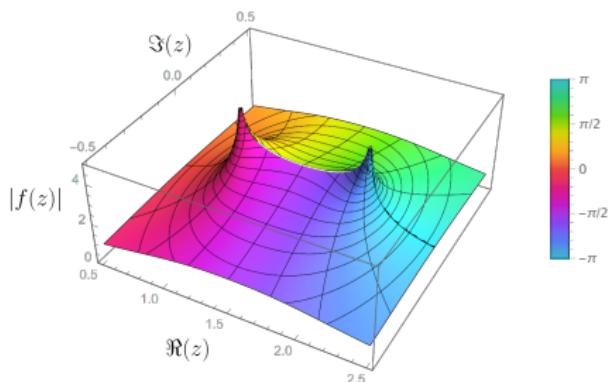
⁵10.1140/epja/i2017-12440-1

Thanks for listening!

Landau singularities

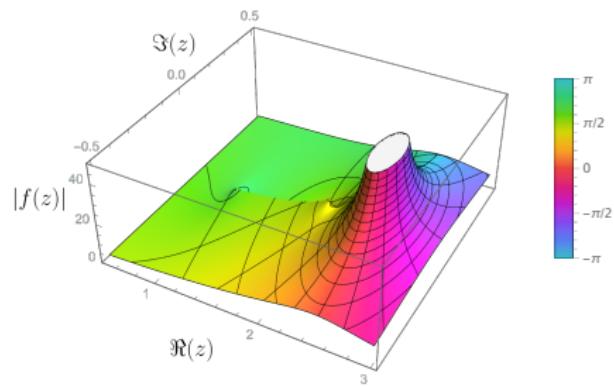
Endpoint singularities: singularities in g , hitting the contour C . E.g.,

$$f(z) = \int_1^2 dw \frac{1}{w-z}. \quad (1)$$

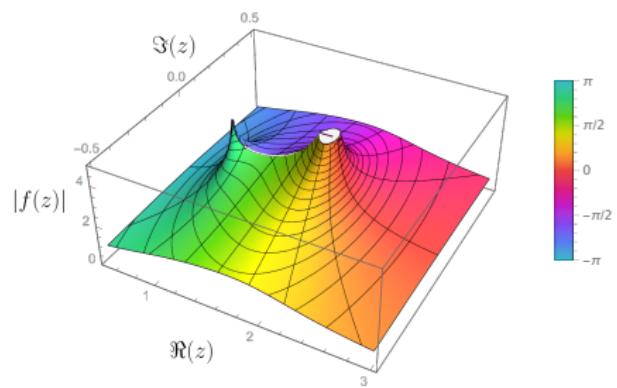


Pinch singularities: Contour C , gets trapped between two singularities in g . E.g.,

$$f(z) = \int_1^2 dw \frac{1}{(w-z)(w-5/2)}. \quad (2)$$



Pinch singularity avoided when the singularities approach from the same side of the contour.



The Other Method

- **Complex Contour on Riemann Surfaces (Cahill & Sloan)⁶:**
Instead of a continued fraction, use the knowledge of the **location of the branch cuts** to analytically continue the amplitude for $q \in \mathbb{R}$.⁷

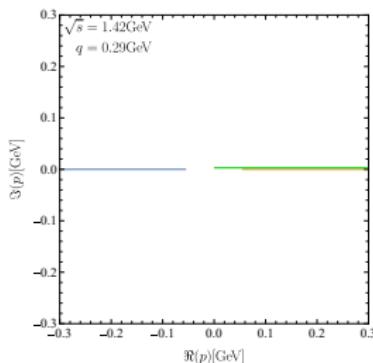
⁶10.1016/0375-9474(71)90156-4

⁷For a comprehensive review: *The Quantum Mechanical Three-Body Problem*, Schmid & Ziegelmann.

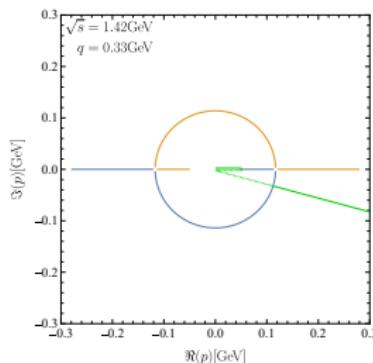
The Other Method

- **Complex Contour on Riemann Surfaces (Cahill & Sloan)⁶:**

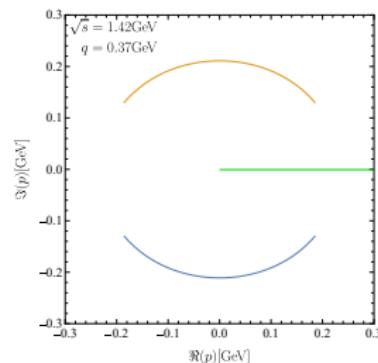
Instead of a continued fraction, use the knowledge of the **location of the branch cuts** to analytically continue the amplitude for $q \in \mathbb{R}$.⁷



Branch cuts of B for small q .



Branch cuts of B for intermediate q .

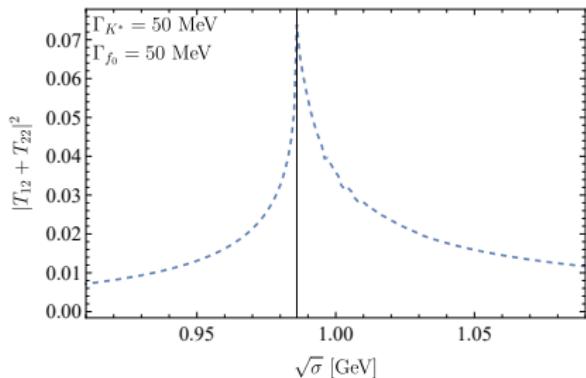


Branch cuts of B for large q .

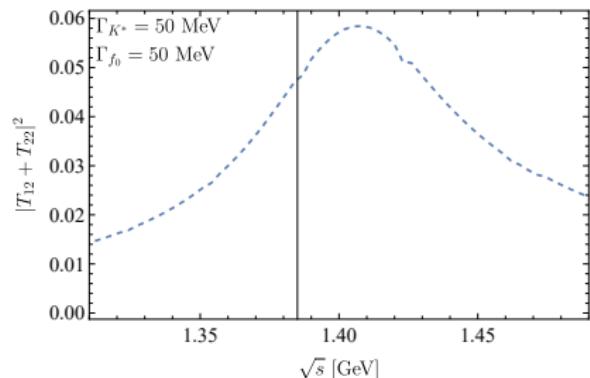
⁶10.1016/0375-9474(71)90156-4

⁷For a comprehensive review: *The Quantum Mechanical Three-Body Problem*, Schmid & Ziegelmann.

- At $1 + \infty$ -loop level,

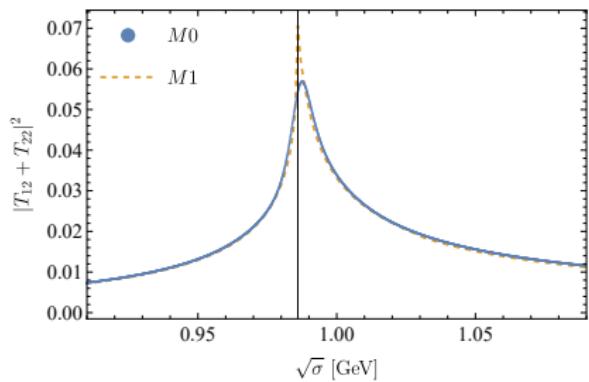


Amplitude squared vs. $\sqrt{\sigma}$, for
 $\sqrt{s} = 1.42$ GeV.

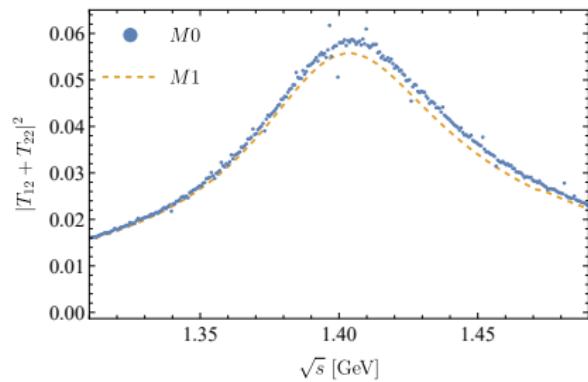


Amplitude squared vs. \sqrt{s} , for
 $\sqrt{\sigma} = 0.99$ GeV.

- Matches with the other method.



Amplitude squared vs. $\sqrt{\sigma}$, for
 $\sqrt{s} = 1.42$ GeV.



Amplitude squared vs. \sqrt{s} , for
 $\sqrt{\sigma} = 0.99$ GeV.